

# EDM correlations

## in SUSY

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## Motivation

EDMs of fundamental particles have not been observed,

BUT...

- ① remarkable sensitivity to NEW PHYSICS

$$d_e \approx \left( \frac{300 \text{ GeV}}{M_{\text{SUSY}}} \right)^2 \sin \Phi_{CP} \times 10^{-25} \text{ e.cm}$$

$\Rightarrow M_{\text{SUSY}}$  up to 100 TeV

for  $d_e \sim 10^{-30} \text{ e.cm}$

- ② experimental progress

$$\begin{aligned} d_e &\rightarrow 10^{-30} \text{ e.cm} \\ d_n &\rightarrow 10^{-28} \text{ e.cm} \end{aligned} \quad \left. \begin{array}{l} \\ \end{array} \right\} \begin{array}{l} \text{in a few} \\ \text{years} \end{array}$$

- ③ complementary to collider data

$$\text{E.g. } \Phi_{CP} \sim 10^{-5}$$

- ④ probe fundamental sources of CP, possibly baryogenesis, ...

# Relativistic EDMs

$$H_{\text{non-rel.}} = -d \vec{S} \cdot \vec{E}$$



$$\mathcal{L}_{\text{rel.}} = -\frac{i}{2} d \bar{\Psi} (F\sigma) \gamma_5 \Psi$$

$\left\{ \begin{array}{l} F_{\mu\nu} = \text{photon field strength} \\ \Psi = \text{fermion} \end{array} \right.$

For composite objects ( $n$ , atoms, ...)  
also relevant

$$\mathcal{L}_{\text{eff}} = \frac{g^2}{32\pi^2} \Theta G_{\mu\nu} \tilde{G}^{\mu\nu} + \frac{w}{3} f^{abc} G_a \tilde{G}_b G_c$$

$$- \frac{i}{2} \tilde{d} \bar{\Psi} g (G\sigma) \gamma_5 \Psi$$

$$+ \sum_{ij} C_{ij} (\bar{\Psi}_i \Psi_i) (\bar{\Psi}_j i \gamma_5 \Psi_j)$$

$\left\{ \begin{array}{l} G_{\mu\nu} = \text{gluon field strength} \\ \Theta = \text{QCD } \Theta\text{-term} \\ w = \text{Weinberg operator} \\ \tilde{d} = \text{colour EDM} \end{array} \right.$

## Neutron EDM

NDA:

Arnowitt, ... '90

$$d_n \sim d_d - 0.25 d_u + e(0.4 \tilde{d}_d - 0.1 \tilde{d}_u) + 0.3 \text{ GeV} \cdot eW$$

All Wilson coefficients at the EW scale.

Sum rules:

Pospelov, Ritz '01

$$d_n \sim d_d - 0.25 d_u + e(\tilde{d}_d + 0.5 \tilde{d}_u) + 0.1 \text{ GeV} \cdot eW$$

Other models

⋮

strangeness  $\rightarrow$  Ellis, Flores '96  
Misano, Shimizu '04

## Thallium / electron EDM

Liu, Kelly '92  
Bouchiat '75

$$d_{Tl} = -585 d_e - C_s \cdot 43 \text{ e.GeV}$$

$$\rightarrow \mathcal{Z} \sim C_s \bar{e} \gamma^5 e \bar{N} N$$

$$\text{When } C_s \rightarrow 0, \quad d_{Tl} \leftrightarrow d_e \quad \Rightarrow |d_e| < 1.5 \cdot 10^{-27} \text{ e cm}$$

# Model - dependence

$d_n$	Model	Ref.
2.7	MIT bag model	Baluni
3.6	Current algebra	Crewther, ...
3.3	Effective chiral approach	Pich, ...
6.7	HB ch PT	Borasoy
3.0	Chiral bag model	Musakhanov, ...
1.4	Cloudy bag model	Morgan, ...
1.2	Chiral quark-meson model	McGovern, ...
2.4	QCD sum rules	Pospelov, ...
1.4	perturbative chiral model	Kuckei, ...

( $d_n$  in units of  $|q| \times 10^{-16}$  e·cm)

# Experimental results:

all negative ...

$$d_{Tl} < 9 \times 10^{-25} \text{ e.cm} \quad (90\% \text{ CL})$$

$$d_{Hg} < 2 \times 10^{-28} \text{ e.cm} \quad (95\% \text{ CL})$$

$$d_n < 6 \times 10^{-26} \text{ e.cm} \quad (90\% \text{ CL})$$

Note :

SM predictions are  
very small,

$$d_n \sim 10^{-32} \text{ e.cm}$$

...

# EDMs in SUSY

(7)

Induced by new CP phases:

$$\begin{aligned}
 \Delta Z = & \left( \mu \bar{\psi}_{H_1} \psi_{H_2} + \text{h.c.} \right) \\
 & + \frac{1}{2} \left( m_3 \bar{\lambda}_3 \lambda_3 + m_2 \bar{\lambda}_2 \lambda_2 + m_1 \bar{\lambda}_1 \lambda_1 \right) \\
 & + \left( A_{ij}^d H_i \tilde{q}_{Lj} \tilde{q}_{Rj}^* + \dots \right)
 \end{aligned}$$

$\dots$  = complex

2 phases eliminated by  $U(1)_R$ ,  $U(1)_{PQ}$

Physical phases:  $\text{Arg}(m_i^* A)$ ,  $\text{Arg}(B^* A)$ , ...

Typical EDM contribution:

Ellis, Ferrara, Nanopoulos '82

$\mathcal{G}_{\text{susy}} \sim 10^{-2} \iff f_L - \frac{i}{\lambda} f_R$ 
  
 (CP problem)

The CP problem appears already in the most minimalistic models, e.g.

mSUGRA :

$\mu, A = \text{complex!}$

This is due to holomorphicity of SUSY.

(Unlike the FCNC problem)



EDMs are special observables, probe even minimal scenarios

## Suppression of EDMs

① Small CP-phases ,  $y \leq 10^{-2}$

(e.g. phase alignment  $g_M \approx g_A \approx -g_\mu + \text{corrections}$ )

② Heavy SUSY spectrum ,  $m \sim \text{a few TeV}$

(motivated by a heavy Higgs  $\rightarrow$  heavy stop)

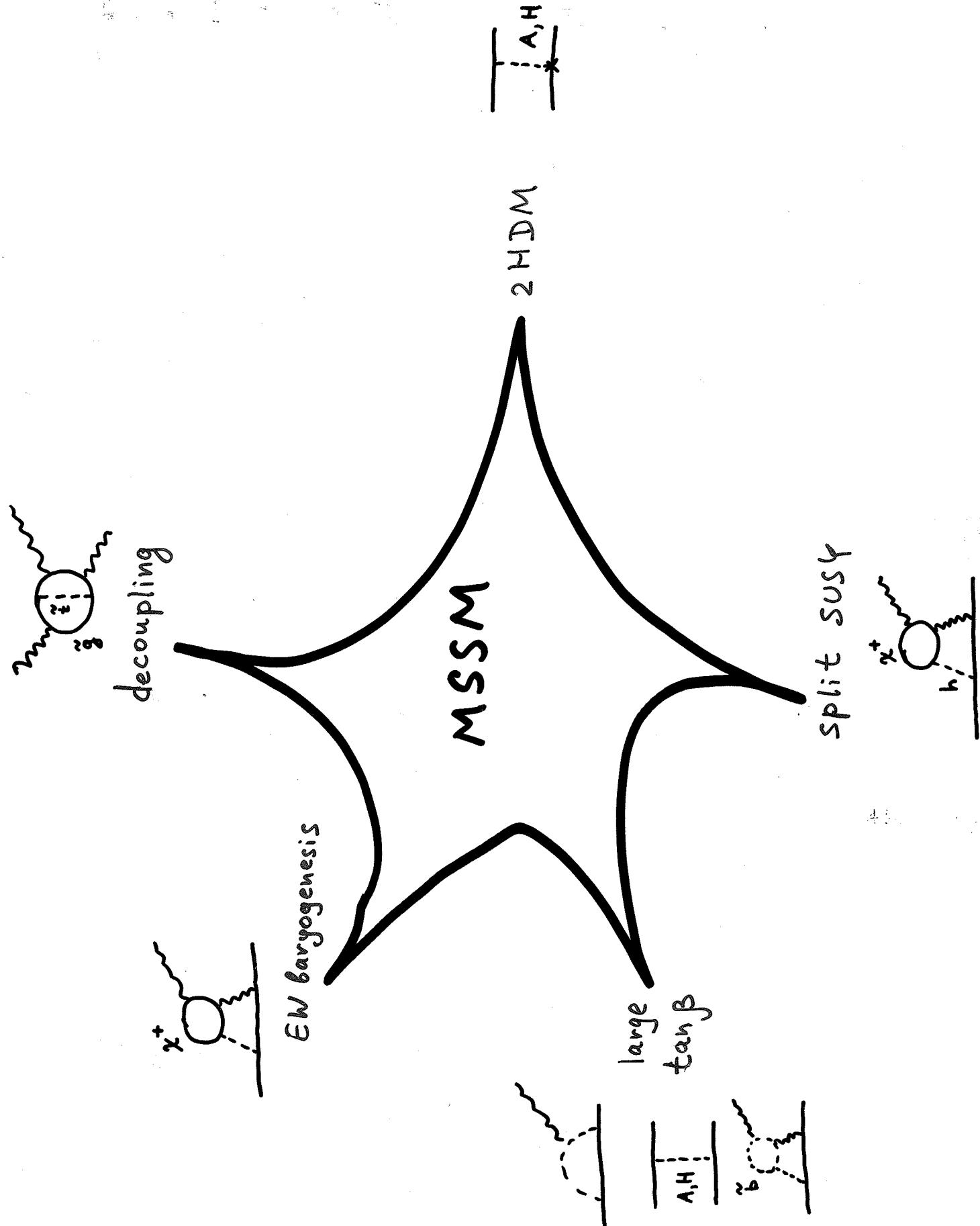
③ Decoupling ,  $\tilde{m}_{1,2} \gtrsim 10 \text{ TeV}$

(but the stop mass  $\lesssim \text{TeV}$ )

③' Decoupling'  $\simeq$  split SUSY

(only fermions are light ,  $\simeq \text{TeV}$ )

The CP problem is difficult , but possible to avoid.



## EDM correlations

(11)

SUSY predicts certain correlations between

$$\underline{d_n, d_e, (d_{Hg}, d_\mu, d_{D,\dots})}$$



indirect signature of SUSY

Typically expect:

$$d_n \sim 10 d_e \quad \left( \frac{m_q}{m_e} \sim 10 \right)$$

even though

$$d_e \sim \frac{1}{x^+, x^0}$$

$$d_n \sim \frac{1}{\tilde{g}, x^+, x^0} + \dots$$

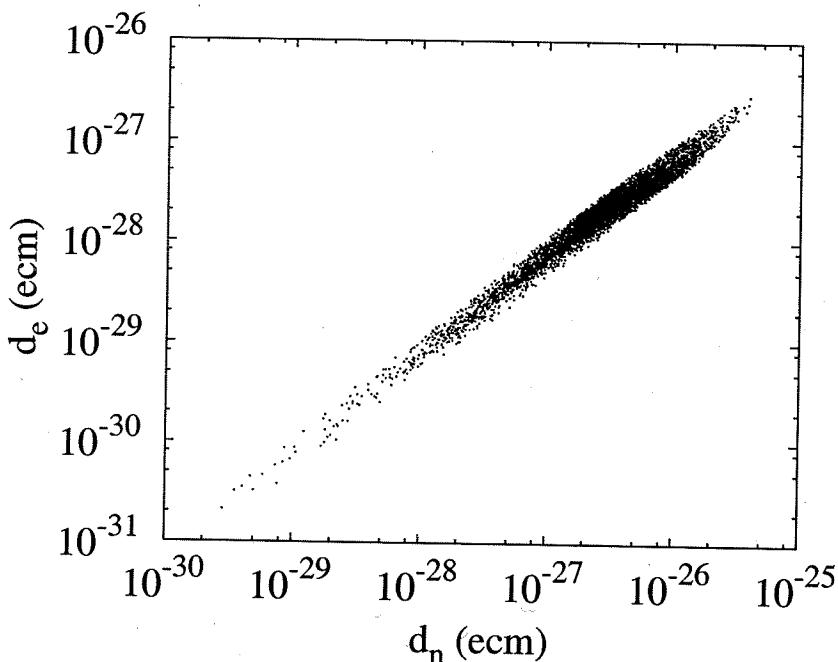
# $d_e - d_n$ correlation

## in m SUGRA

$m_0, m_{1/2}, |A| \in [200 \text{ GeV}, 1 \text{ TeV}]$

$\tan\mu \in [-\pi/500, \pi/500]$

$\tan\beta = 5$



$$d_e \sim 10^1 d_n$$

↓  
indirect  
evidence  
for  
SUSY!

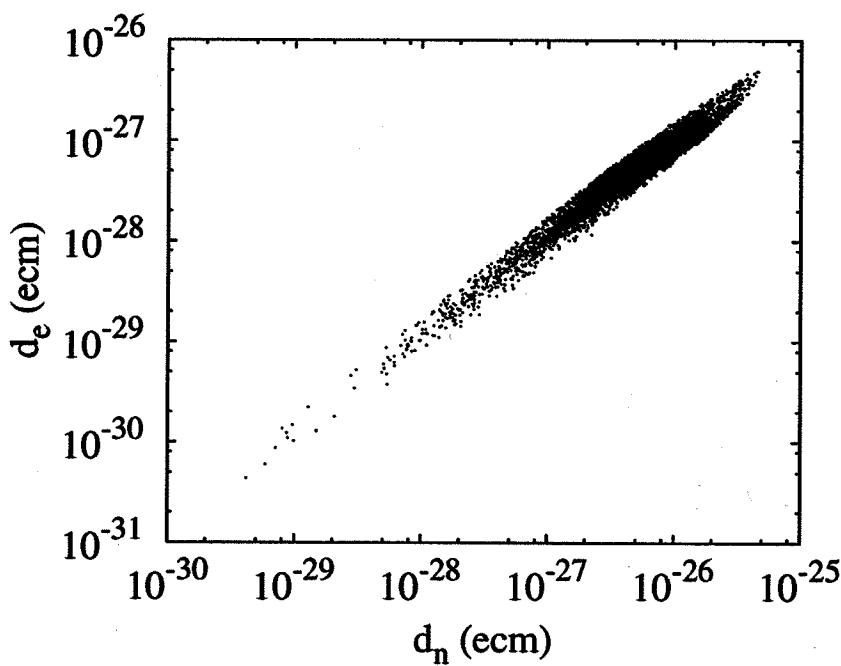
Note: persists for  $\tan\beta \sim 35$ ; also for  $m_{\tilde{g}} \neq m_{\tilde{e}}$  (GUT)

# Heavy m SUGRA

$$m_0, m_{1/2}, |A| \in [2 \text{ TeV}, 10 \text{ TeV}]$$

$$\Phi_A, \Phi_\mu \in [-\pi, \pi]$$

A very similar picture:



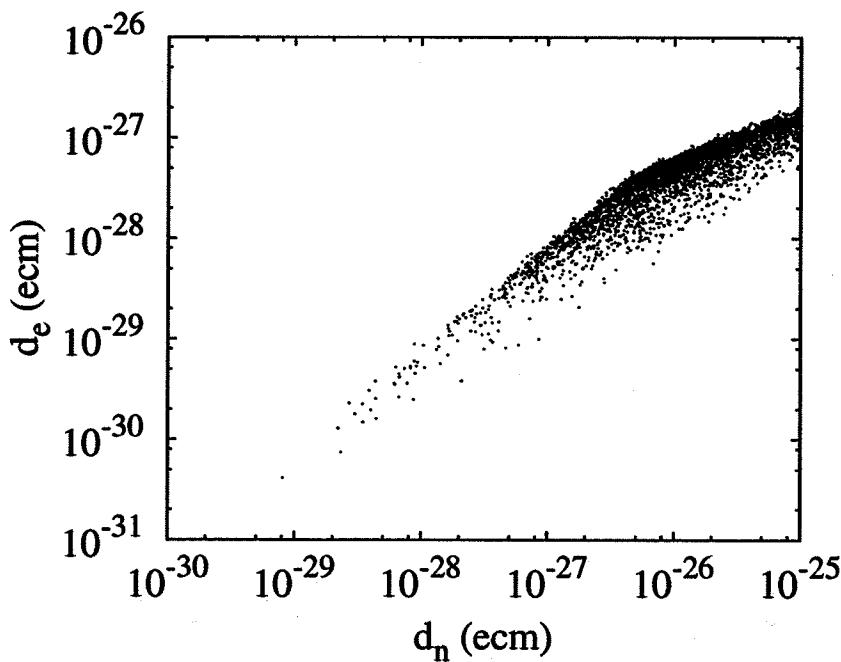
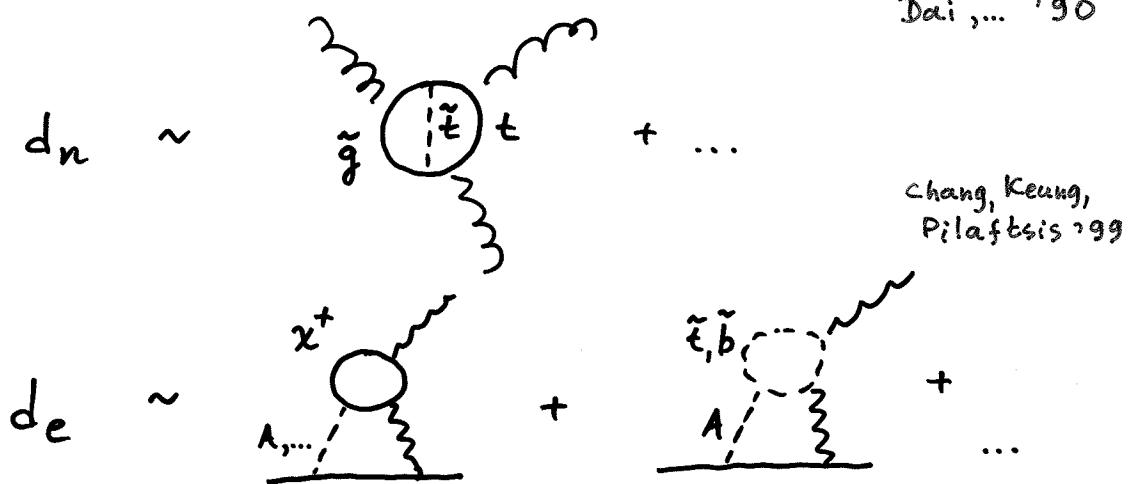
Most points  
are  
in the  
observable  
range!

## Decoupling

$$m_{\tilde{q}_3}, m_{1/2}, |A| \in [200 \text{ GeV}, 1 \text{ TeV}]$$

$$\vartheta_A, \vartheta_\mu \in [-\pi, \pi]$$

$$m_{\tilde{q}_{1,2}} \rightarrow \infty$$



$$d_n \sim 10 \div 100 d_e$$

Should be  
observed  
soon!

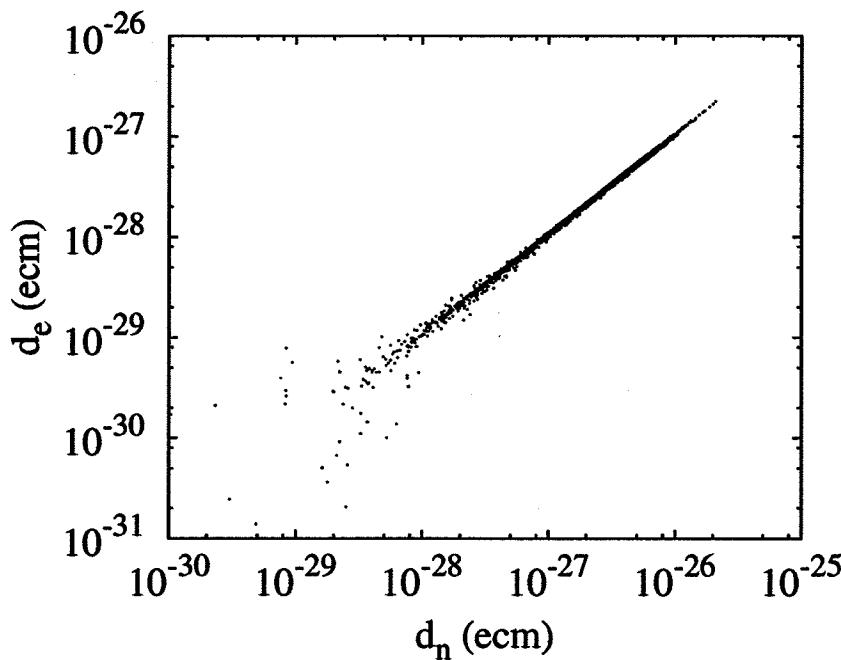
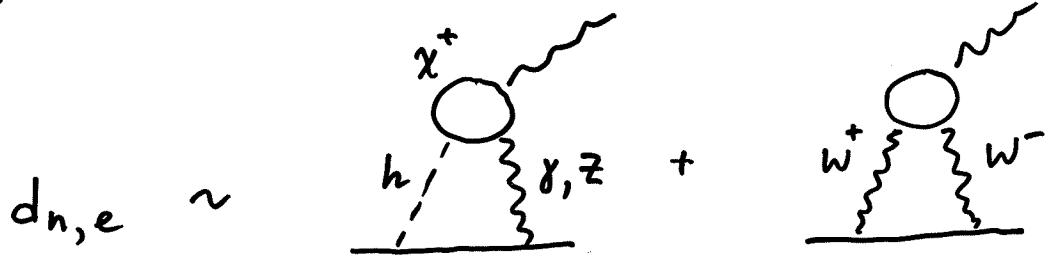
# Split SUSY

$$M_{1,2,3} \in [200 \text{ GeV}, 1 \text{ TeV}]$$

$$|\mu| \in [200 \text{ GeV}, 1 \text{ TeV}]$$

$$m_h \in [100 \text{ GeV}, 300 \text{ GeV}]$$

$$m_{\text{scalars}} \rightarrow \infty$$



$$d_n \sim 10 d_e$$

Observable!

# $d_n$ - $d_e$ correlations

Abel, OL '05

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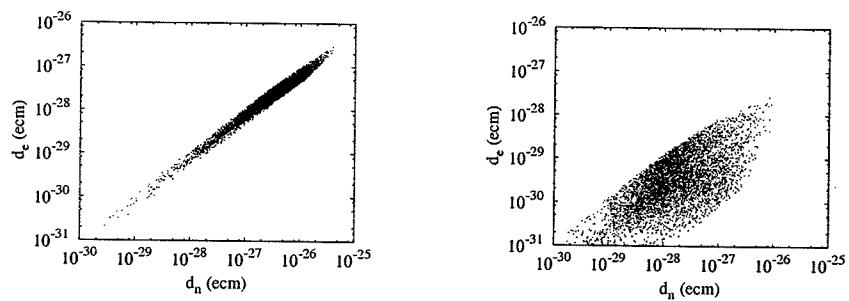


Figure 6:  $d_e$  vs  $d_n$  in mSUGRA with small phases,  $\tan \beta = 5$ . Left:  $\phi_\mu \neq 0$ , right:  $\phi_A \neq 0$ .

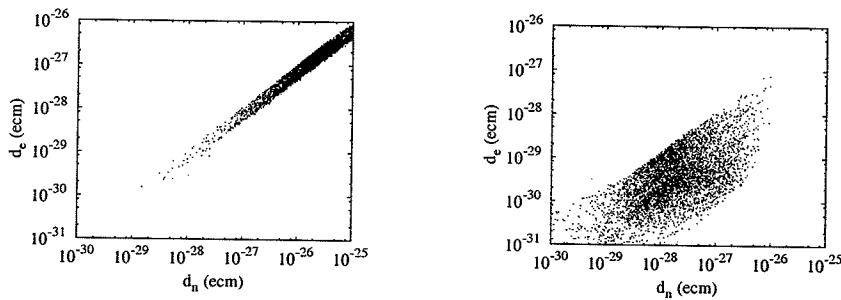


Figure 7: As in Fig. 6, but for  $\tan \beta = 35$ .

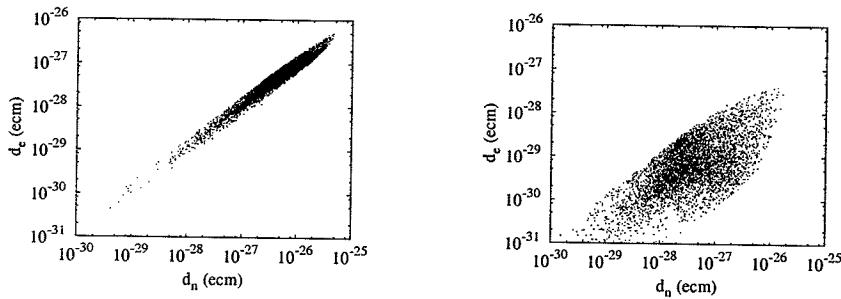


Figure 8:  $d_e$  vs  $d_n$  in mSUGRA with a heavy spectrum,  $\tan \beta = 5$ . Left:  $\phi_\mu \neq 0$ , right:  $\phi_A \neq 0$ .

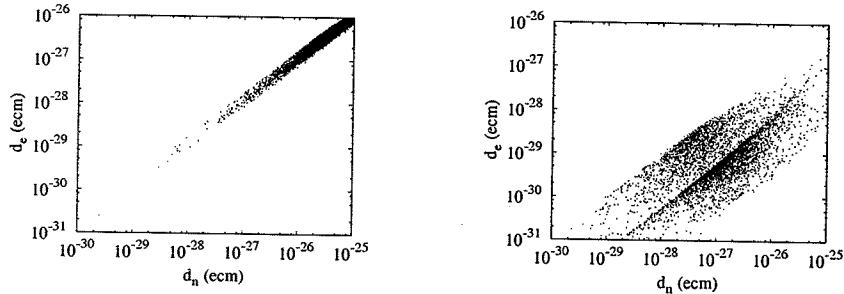


Figure 9: As in Fig. 8, but for  $\tan \beta = 35$ .

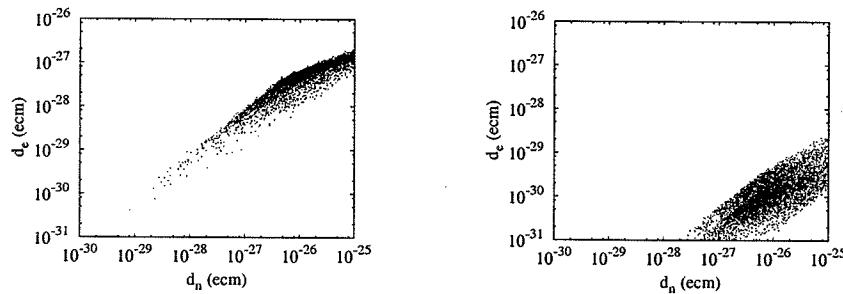
$d_n$ - $d_e$  correlations

Figure 10:  $d_e$  vs  $d_n$  in the decoupling scenario,  $\tan \beta = 5$ . Left:  $\phi_\mu \neq 0$ , right:  $\phi_A \neq 0$ .

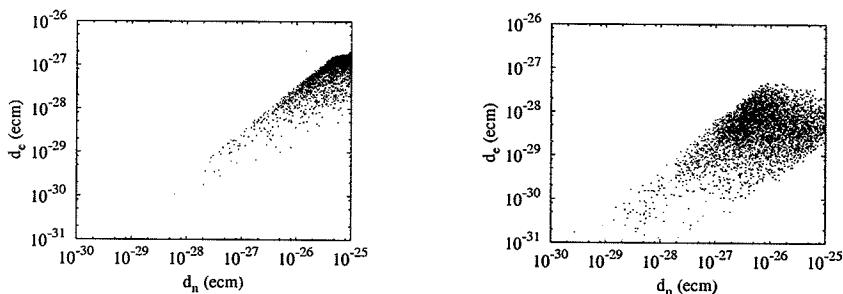


Figure 11: As in Fig.10, but for  $\tan \beta = 35$ .

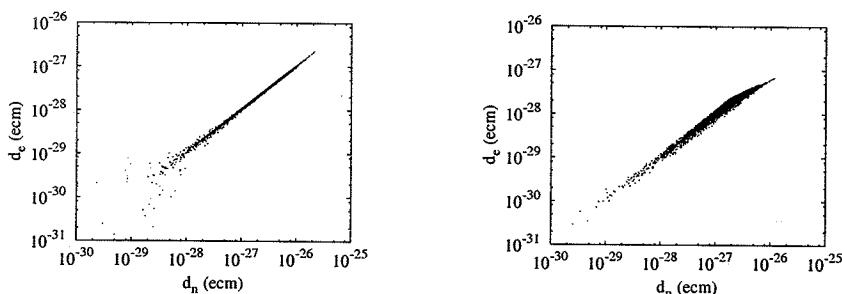


Figure 12:  $d_e$  vs  $d_n$  in split SUSY. Left:  $\tan \beta = 5$ , right:  $\tan \beta = 35$ .

## Message:

typically

$$\underline{d_n \sim 10 \div 100 d_e}$$

(rather insensitive to  $\tan\beta$ ,  
universality assumptions,  
... )

Thus,

$$\begin{array}{ccc} d_e \gtrsim d_n & \Rightarrow & \text{SUSY} \\ d_e \ll d_n & & \text{disfavored} \end{array}$$

Also,

complementary to collider data:

$$\left. \begin{array}{l} M_{\text{SUSY}} \approx \text{TeV}, \quad g_{\text{CP}} \sim 10^{-5} \\ M_{\text{SUSY}} \gtrsim 5 \text{ TeV}, \quad g_{\text{CP}} \sim 1 \end{array} \right\} \begin{array}{l} \text{out} \\ \text{of} \\ \text{colliders' reach} \end{array}$$

## Conclusion:

- even in the minimal SUSY scenarios EDMs are observable within 5-... years
- $d_n - d_e, \dots$  correlations  $\rightarrow$   
 $\rightarrow$  indirect evidence for SUSY
- complementary to collider data