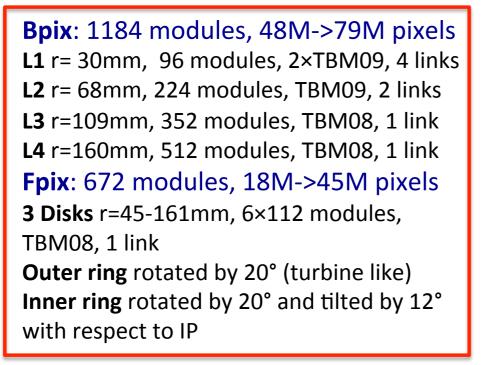
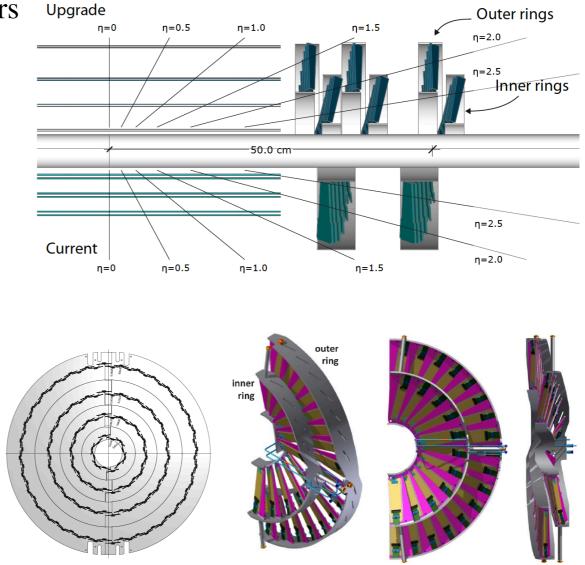
# Sensor Simulations in CMS

M. Swartz Johns Hopkins University

## **Phase 1 Pixel Detector**

- Upgrade Pixel detector: 4 barrel layers (instead of 3) and 3 forward disks on each side (instead of 2)
- Installation: during extended winter shutdown in 2016/17

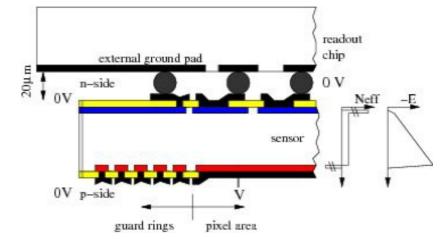


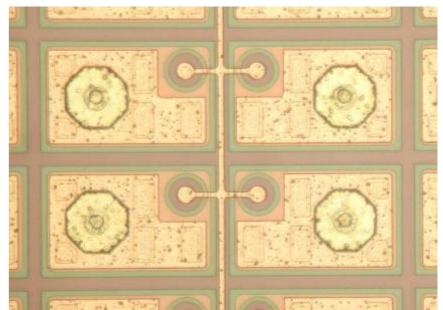


- New 4 layer barrel [BPix] and 3 disk forward [FPix] detectors
  - \* new digital ROCs for BPix L2-4/FPix and BPix L1
  - \* mixed phase CO2 cooling
  - \* DC DC powering

#### Sensors

- BPix Sensors [CiS]
  - \* same modified p-spray design as for Phase 0
  - \* 100x150 um cell size
  - \* bias grid with punch though resistors
  - \* DOFZ <111> substrate
  - \* resistivity 5±0.5 kΩcm
  - \* polished to 285 µm thickness
- FPix Sensors [Sintef]
  - \* modified Phase 0 p-stop design
    - better HV performance
  - \* 100x150 um cell size
  - resistivity ~8 kΩcm
  - \* 290 µm thickness





# **Readout Chips**

- PSI46DIGI: BPix L2-4, FPix
  - 8-bit digital charge info [analog in Phase 0]
  - \* Readout speed 160Mbit/s [40 MHz]
  - Time stamp buffer size 24 [12]
  - \* Data buffer size 80 [24]
  - Six metal layers [5]
  - In time threshold <2000e [3500e]</li>
  - \* Data loss 1.6%@150MHz/cm<sup>2</sup> [5-6%]
- PROC600: BPix L1
  - \* Dynamic Cluster Column Drain [2x2 pixels]
  - \* Transfer speed increased 20->40 MHz
  - \* Deadtime free data buffer management
  - \* Data loss 2.5%@585MHz/cm<sup>2</sup>

## **Radiation Exposure during 2017**

The absolute charged particle fluences were detemined from Fluka and independently from counting clusters at a large enough radius 10 cm to minimize track angle effects. The relative fluences come from measuring the total cluster charge/volume in different subdetectors

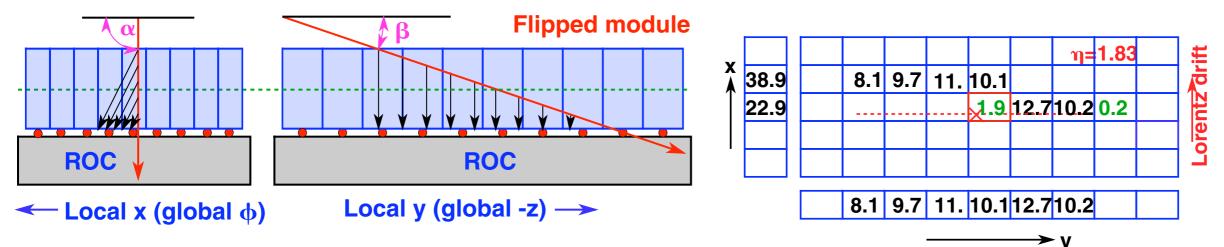
- BPix L1 saw over 5x10<sup>14</sup> cm<sup>-2</sup>
  - \* multiply by 0.6 to get  $\Phi_{neq}$ ?
  - \* effects of neutrals?
- BPix L2 and FPix R1 are quite similar
  \* same fluence in each FPix disk
- BPix L3 and FPix R2 are also quite similar
  - \* same fluence in each FPix disk

The appropriate hardness factor will actually be quite important in determining the longevity of the detector ...

Layer	Charged fluence
<b>BP LI</b>	10x10 <sup>12</sup> cm <sup>-2</sup> /fb
BP L2	2.6x10 <sup>12</sup> cm <sup>-2</sup> /fb
BP L3	1.4x10 <sup>12</sup> cm <sup>-2</sup> /fb
BP L4	0.8x10 <sup>12</sup> cm <sup>-2</sup> /fb
FP R I	3.1x10 <sup>12</sup> cm <sup>-2</sup> /fb
FP R2	1.4x10 <sup>12</sup> cm <sup>-2</sup> /fb

#### **Hit Reconstruction**

Tracks deposit distinct patterns of charge on the pixel sensors



- Hit position estimation is based on 1D projections of the 2D cluster
  - \* factorizes due to field configurations and cell periodicity
  - projected shapes depend upon the projected angles a and b
  - \* reconstruction algorithms use angle information iteratively
- Two techniques used in track reconstruction
  - \* "Generic" technique is h-like, uses end pixel charges of projection
    - faster, less precise algorithm used for all but last tracking pass
    - needs external Lorentz drift calibration [from detailed simulation]
  - \* "Template" technique fits projections to simulated profiles
    - slower, more precise algorithm used for final fitting pass
    - needs full cluster shape calibration
    - generates probabilities that test the consistency of the shapes

### **Pixelav Detailed Simulation**

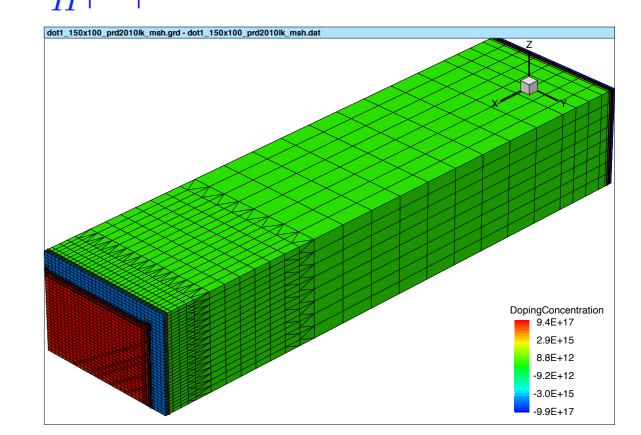
Created to interpret beam tests of irradiated sensors, now used to perform Lorentz calibrations and generate template profile shapes:

- charge deposition model based on Bichsel π-Si x-sections
- delta ray range: Continuous Slowing Down Approx + Nist Estar dedx
- plural scattering and magnetic curvature of delta ray tracks
- carrier transport from Runge-Kutta integration of saturated drift

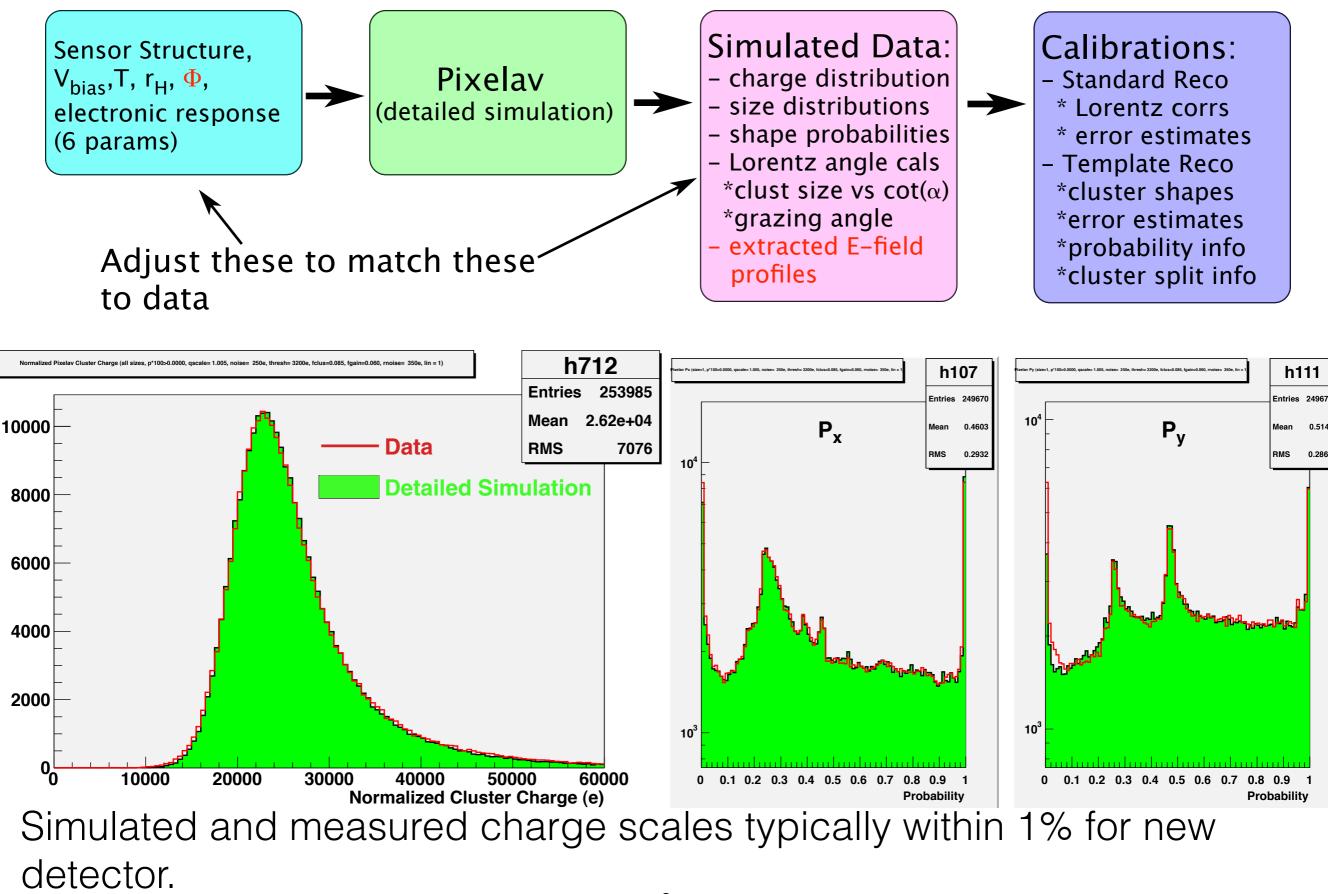
$$\frac{d\vec{x}}{dt} = \vec{v} = \frac{\mu \left[ q\vec{E} + \mu r_H \vec{E} \times \vec{B} + q\mu^2 r_H^2 (\vec{E} \cdot \vec{B}) \vec{B} \right]}{1 + \mu^2 r_H^2 |\vec{B}|^2}$$

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- includes diffusion, trapping, and charge induction on implants
- Electronic Simulation: noise, linearity, thresholds, mis-calibration



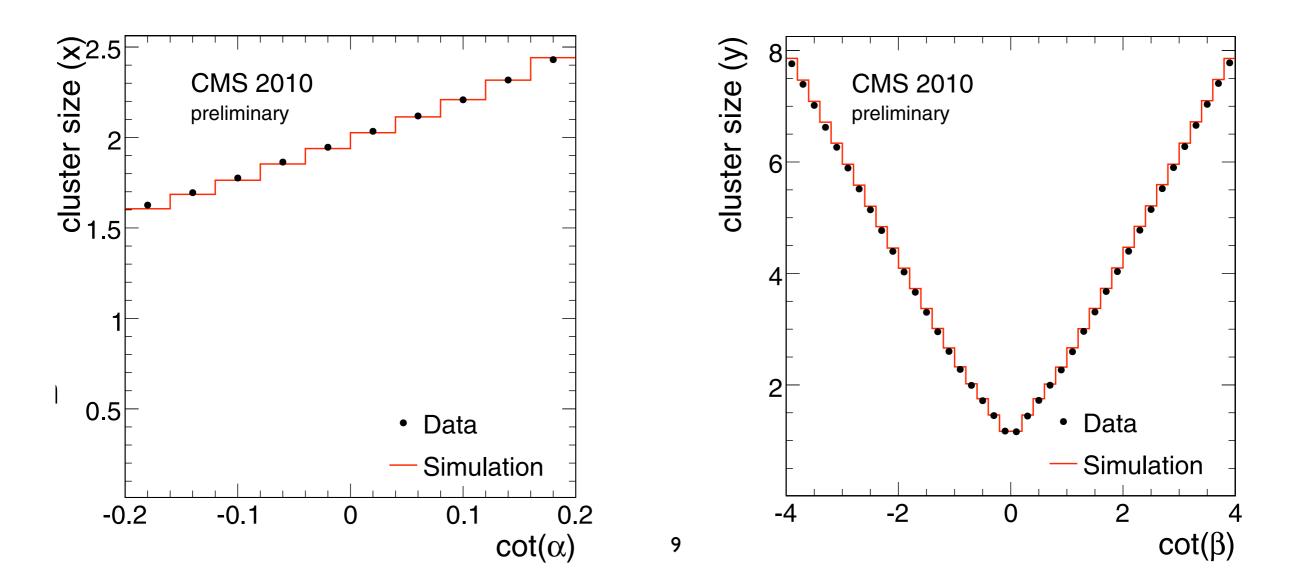
## **Calibration of Reconstruction Algs**



#### **Calibration of "In-Time" Thresholds**

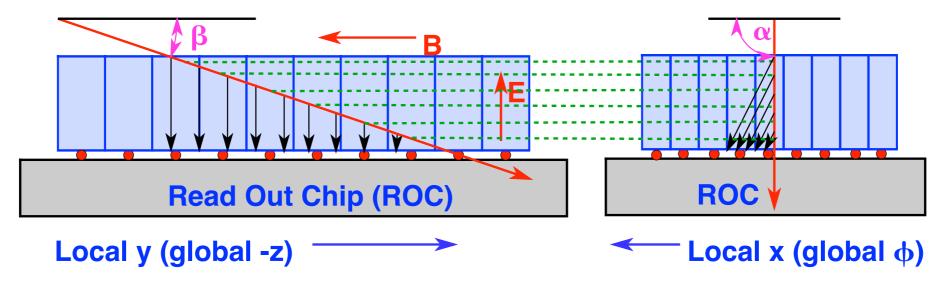
The average x/y cluster sizes for each bin in  $\cot(\alpha)/\cot(\beta)$  depend upon the effective threshold.

- Simulate the same track angles, momenta as reconstructed in the data
  - charge per unit  $\cot(\alpha/\beta)$  is same for simulated/measured samples
- Adjust threshold to achieve best agreement
  - x-size vs  $\cot(\alpha)$  is also sensitive to the Lorentz angle (meas separately)
  - thresholds vary from ~1600e [L3/4] to ~2500e [L2] to 3500-5000e



#### **Old E-field Calibration Technique**

Use signal trapping as a function of pixel column [depth] to probe Efield shape across the substrate of irradiated sensors



- Need to run at a series of low bias voltages
- 2003-2005 beam tests used a ROC with no zero suppression

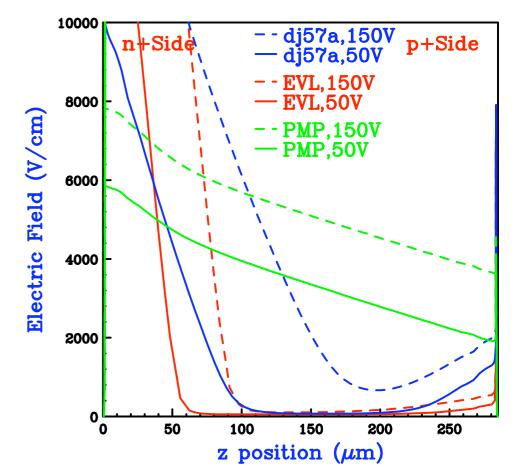
could see very small [even wrong sign] signals

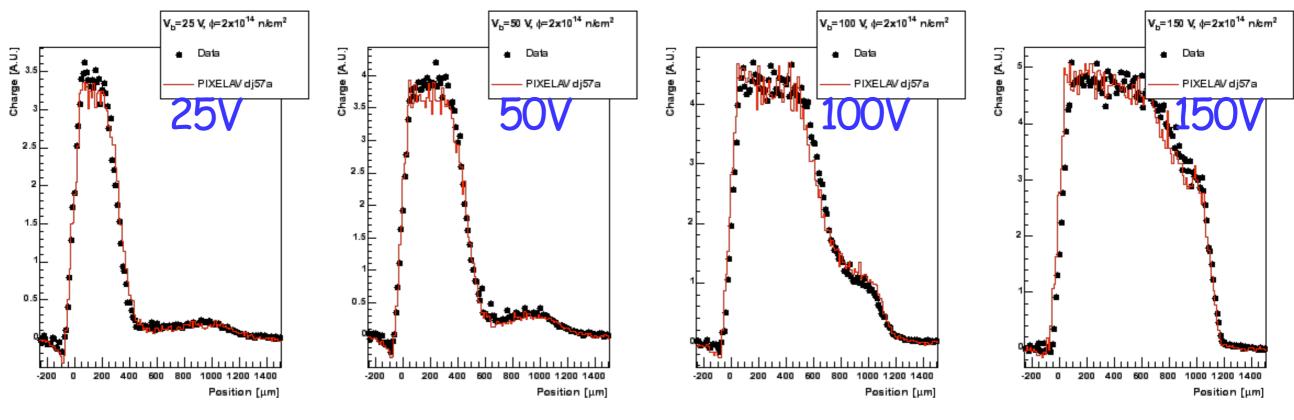
- Model results using TCAD with SRH statistics and 2 midgap defects [Eremin, Verbitskaya, Li]
  - \* TCAD defects model the E-field shape
  - \* Pixelav signal trapping independently adjusted

#### From RD50 Workhop 6 [2005]

Best fit to 2.0x1014 neq/cm2: labelled dj57a

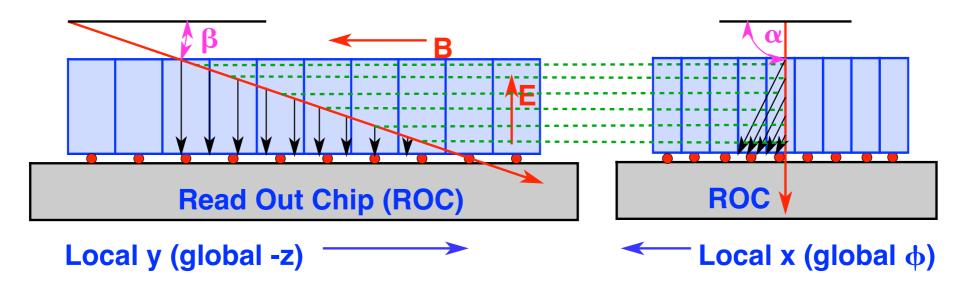
- ∗ N<sub>A</sub>/N<sub>D</sub>=0.68
- \*  $\Gamma_{e/h} = 0.8$  \* Ljubljana trapping rates
- \*  $\sigma_{Ah}/\sigma_{Ae}=0.25, \sigma_{Dh}/\sigma_{De}=1.00,$
- \* E-field still doubly-peaked



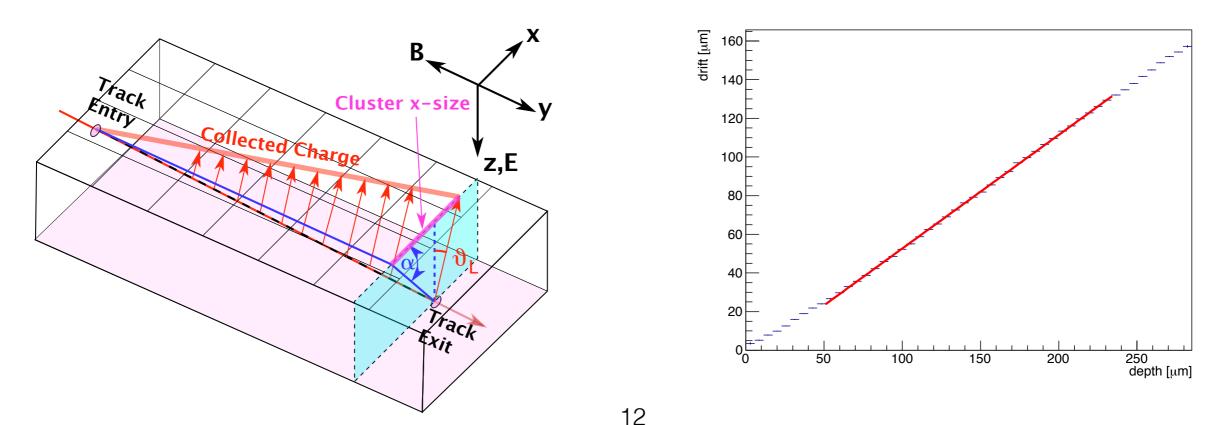


#### **Lorentz Angle Calibration**

Drift vs depth [grazing angle technique] was developed by UniZ colleagues to calibrate the Lorentz angle

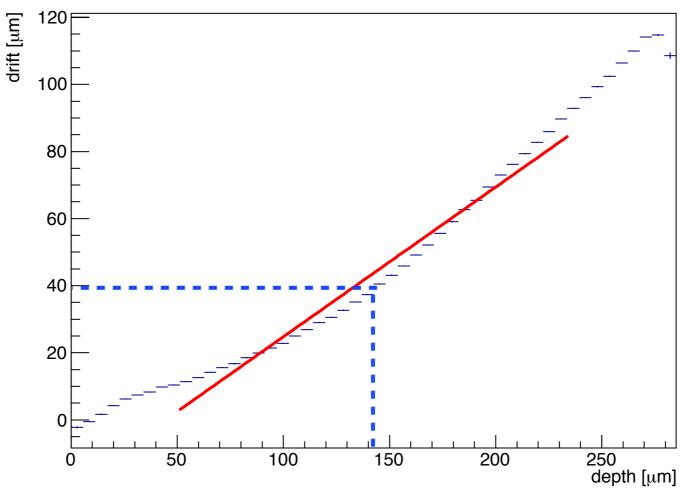


Accumulate the charge centroid [drift] vs depth for a sample of highly inclined tracks. The angle is the average Lorentz angle



## **Lorentz Angle Calibration**

When the detector is irradiated, this technique is not useful



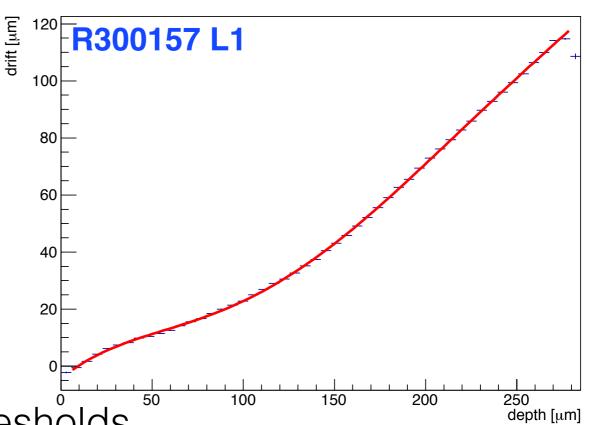
- the slope becomes steeper, but the actual offset from Lorentz drift [needed for eta-like reco] becomes smaller
  - \* Lorentz drift correction is the offset at the detector midplane
  - \* steeper slope would imply a larger drift correction
- we need a better method to calibrate the sensor simulation

# **E-Field Measurement and Template/LA Calibration**

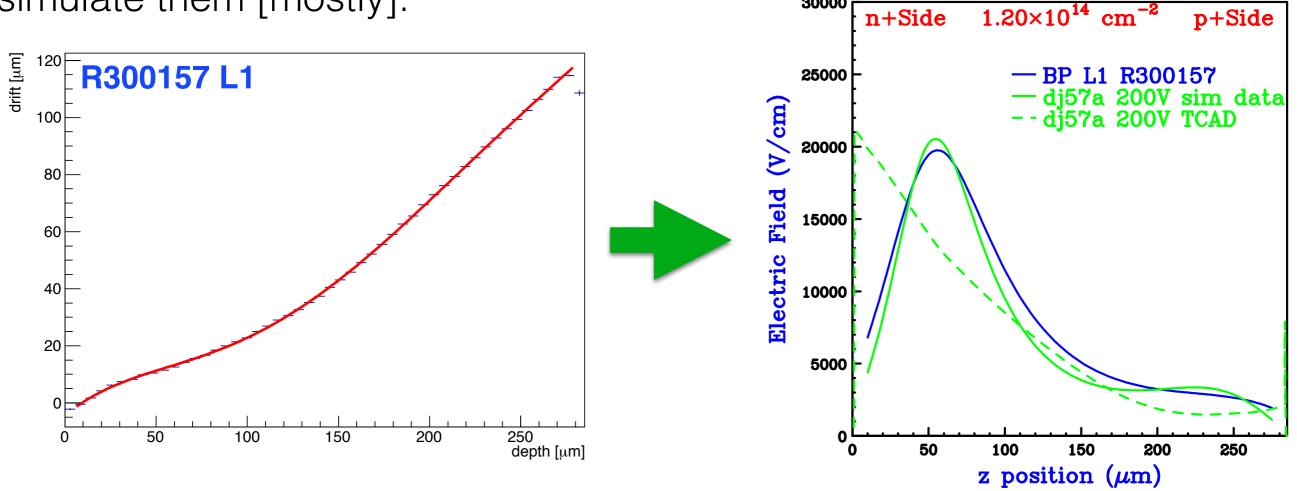
Take our drift (x) vs Depth (D) data, fit to a polynomial [5th order] and then calculate a local slope [Lorentz Angle] vs D. We then convert it to an E vs D curve from the expression  $E^{120}$ 

$$\tan \theta_L = \frac{dx}{dD} = r_H \mu(E) B_y$$
$$\rightarrow E = \mu^{-1} \left( \frac{1}{r_H B_y} \frac{dx}{dD} \right)$$

- depends upon the slope dx/dD
  \* insensitive to alignment effects
- insensitive to the knowledge of thresholds
- insensitive to trapping [displacement is measured at fixed depth]!
- can be done at operating voltage: no need for bias scans
- extracts information that is sort of comparable to the simulated E-field
  \* still need to simulate the extracted fields in this procedure
- Q vs D distributions can then be used to independently adjust the trapping rates for e/h



The extracted electric field profile is distorted by focusing near the n+ implant and other systematic effects. The good news is that we can simulate them [mostly]:



- Run 300157 was taken after 11.8 fb<sup>-1</sup>:  $\Phi_Q = 1.2 \times 10^{14} \text{ cm}^{-2}$ 
  - \* the neutron equivalent flux [0.6 hardness]  $\Phi_{eq} = 0.72 \times 10^{14} \text{ cm}^{-2}$
  - \* the electric field is well described by our old model dj57a?
    - it was from a sensor that had been exposed to  $\Phi_{eq} = 2 \times 10^{14} \text{ cm}^{-2}$

# Differences

## **Beam Test**

- Sensors are irradiated in a short time [hours/days]
  - defect-defect interactions
    occur at rates ~ density<sup>2</sup>
- Sensors are "standard annealed"
  - in 2003, sensors warmed to 30C for 30 hours?
    - they were always kept cold after that

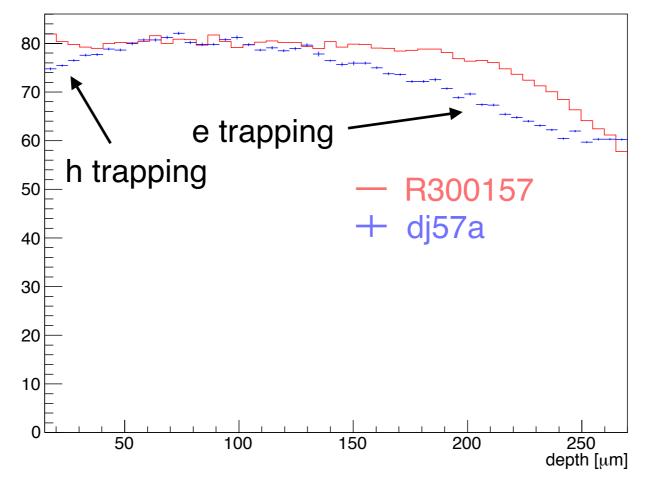
# Operation

- Sensors are irradiated in a long time [months]
  - defects have more time to interact with impurities
- Sensors are not annealed at all
  - in CMS, everything has been kept at -10C [ROCs generating heat] or -20C [readout off]
    - there will be some annealing when maintenance is done if not before

Could these differences affect the evolution of the sensor E-fields and trapping rates?

#### **Trapping Measurement**

Compare the measured depth profile with the simulated profile



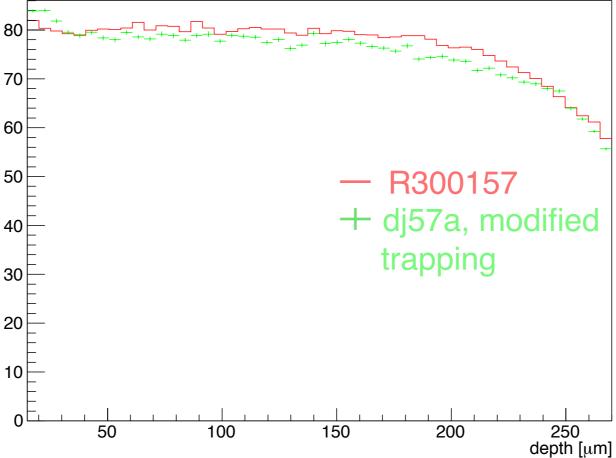
The trapping rates for e and h are both too large!

How much trapping do we expect for  $\Phi_Q = 1.2 \times 10^{14} \text{ cm}^{-2}$ ?

In our test beam models, the trapping rates should scale as  $0.8\Phi_{eq} = 0.48\Phi_Q = 0.6 \times 10^{14} \text{ cm}^{-2}$ ?

# **Trapping Measurement**

Simulate the dj59a E-field with trapping rates corresponding to 0.6x10<sup>14</sup> cm<sup>-2</sup>



- The electric field is evolving faster [differently] than expectations from the beam test models
- Trapping rates appear to be evolving according to the fluence calculation with a hardness factor of 0.6
- The slower evolution of the trapping rates has important consequences for the longevity of the detector

#### Summary

- Sensor modeling is a key element in the calibration of CMS' pixel hit reconstruction
- Lorentz drift vs depth provides information that is used in tuning the sensor models
  - \* it is performed with full bias voltage collision data
- The space charge effects have onset more quickly than might have been expected from beam test data
  - \* the effects differ from beam test expectations
  - could be due to different radiation profiles or different annealing history
- The signal trapping effects seem to be behaving according to expectations from beam test data