Coherence effects in a QCD parton cascade

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Motivation: jet substructure



- Ideal techniques for heavy ion collisions.
- More direct access to the underlying dynamics:
 - QGP properties.
 - Energy loss.
 - Coherence.

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Color coherence in vacuum

• Is radiation independent?: $q\bar{q}$ antenna as a laboratory.



$$dN = \frac{d\omega}{\omega} \frac{d\Omega}{2\pi} \frac{\alpha_s C_F}{2\pi} \Big[R_q + R_{\bar{q}} - 2\mathcal{J} \Big]$$

- The spectrum is suppressed at large angles due to the presence of destructive inteferences (coherence).
- Angular ordering.

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Color coherence in a medium

• How does the medium change this picture?



• A parton can change color through interaction with the medium, breaking the correlation between emitted gluons.

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Particle propagation in matter



$$W(ec{x}) = \mathcal{P} \; exp \Big[ig \int dx_+ A_-(x_+, ec{x}) \Big]$$

- The effect of the medium is to induce color rotation at each scattering center.
- The quark (a high energy quark) loses a negligible amount of energy and propagates in straight lines (*eikonal* propagation).

In-medium antenna radiation

• To study the degree of coherence we a take a very soft gluon $\omega \rightarrow 0$ (out-out radiation).



Image: A matrix and a matrix

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The decoherence parameter

• The interaction of the $q\bar{q}$ pair with the medium is described by the survival probability S.

$$egin{aligned} \mathcal{S} &\equiv rac{1}{N_c^2-1} \Big\langle W(ec{x}_\perp) W^\dagger(ec{y}_\perp) \Big
angle \ &\mathcal{S} &\equiv 1-\Delta_{med}(t) \ \end{aligned}$$
 $\Delta_{med} &\equiv 1-exp \Big[-rac{1}{4} \hat{q} L(ec{x}_\perp-ec{x}_\perp)^2 \Big] \end{aligned}$

• This factor determines a characteristic time-scale for decoherence of the $q\bar{q}$ pair.

The resulting spectrum

$$dN = \frac{d\omega}{\omega} \frac{d\Omega}{2\pi} \frac{\alpha_s C_F}{2\pi} \Big[R_q + R_{\bar{q}} - (1 - \Delta_{med}) \ 2\mathcal{J} \Big]$$

$$\Delta_{med}
ightarrow 0: dN \sim R_q + R_{ar q} - 2\mathcal{J}$$

Dilute medium : coherence (angular ordering)

$$\Delta_{med}
ightarrow 1$$
 : dN $\sim R_q + R_{ar q}$
Opaque medium : decoherence (two independent emitters)

[The radiation pattern of a QCD antenna in a diluite/dense medium, Yacine Mehtar-Tani, Carlos A. Salgado and Konrad Tywoniuk]

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Main limitations

- We have to deal with more realistic settings:
 - Non-eikonal antenna.
 - Multiple emissions.
 - Finite formation time.

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Multiple emissions

- The antenna provides a simple and intuitive picture.
- Does it hold for more than two emitters?



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Direct terms



 $|\mathcal{M}_1|^2 \propto C_F^2$ $|\mathcal{M}_2|^2 \propto C_F^2$ $|\mathcal{M}_3|^2 \propto N_c C_F^2$

• The direct terms are proportional to a color factor, i.e., no medium effects appear.

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Interference terms





Large N_c limit

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 $\mathcal{M}_1 \otimes \mathcal{M}_3^* \propto \mathcal{S}(t,L)$

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Interference terms



Large N_c limit

$$\mathcal{M}_2\otimes \mathcal{M}_3^*\propto \mathcal{S}(0,t)\ \mathcal{S}(t,L)$$

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Multiple emissions results

- We have considered the case of three emitters.
- The interference terms are proportional to the survival probabilities S in the (0, t) and (t, L) regions: the general result of the antenna is valid for each of the smaller antennas.
- If coherence is not preserved after the in-medium splitting, the antenna won't radiate coherently in the following emission.
- These computations can be generalized to the problem of *n* emitters.

Main limitations

- We have to deal with more realistic settings:
 - Non-eikonal antenna.
 - Multiple emissions.
 - Finite formation time.

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Finite formation time antenna



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Finite formation time antenna

- Region I:
 - q and \bar{q} phases: $exp\left\{i\frac{p_{1\perp}^2}{2E_i}(t_2-t_1)\right\}$.
 - Average of the Wilson lines: $exp\left\{-\frac{1}{12}\hat{q}n_{12}^2(t_2-t_1)^3\right\}$.
 - Competing process between t_f and t_d :
 - $t_f \ll t_d$: vacuum propagation.
 - $t_f >> t_d$: medium effects.

Finite formation time antenna

- Region II:
 - All phases cancel out.
 - Average of a trace of four Wilson lines:

$$Q(t_L, t_2) = \frac{1}{N_c} \Big\langle Tr \Big[W_1(t_L, t_2) W_2^{\dagger}(t_L, t_2) W_{\bar{1}}(t_L, t_2) W_{\bar{1}}^{\dagger}(t_L, t_2) \Big] \Big\rangle$$

$$Q(t_L, t_2) = e^{-\frac{1}{4}\hat{q}(n_1^2 + n_2^2)(t_2 - t_1)^2(t_L - t_2)} \Big[1 - \hat{q}(n_1 \cdot n_2)(t_2 - t_1)^2 \int_{t_2}^{t_L} dt_3 \ e^{-\frac{1}{6}n_{12}^2(t_3 - t_2)^3} \Big]$$

• Competing process between p_T and Q_s .

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Finite formation time antenna

- We have studied a singlet antenna with **short formation time** considering separately two different regions:
 - **Region I** contains information about **local scales** (t_f vs. t_d).
 - **Region II** compares global scales $(p_T \text{ vs. } Q_s)$.
- Both regions are well connected.
- We are studying the details of the relation between these regions to obtain general results about finite formation time setups.

A hard quark propagating through a medium



$$\propto \frac{1}{N_c \ (N_c^2 - 1)} \left\langle W^{ai}(\vec{0}) \ W^{ai}(\vec{r_3}) \right\rangle_{(t,t')} \left\langle f^{ijc} \ f^{\alpha bz} \ W^{i\alpha}(\vec{r_3}) \ W^{jb}(\vec{0}) \ W^{\dagger zc}(\vec{r_3}) \right\rangle_{(t',L)}$$

Summary

- HIC are key to go further in the study of the QCD matter formed in the collisions.
- Color coherence is essential to understand the jet constituents' energy loss (are they independent or not?).
- In spite of the singlet antenna limitations (eikonal propagation, zero formation time, only one splitting...), it is a very convenient *laboratory*.
- The general result of the singlet antenna is valid for the subsequent antennas in the multiple emissions case.
- Finite formation time setups showed us some interesting preliminary results about the evolution of these systems.
- These computations go a step forward to obtain a complete description of a QCD cascade.

Thanks for your attention

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