

Recent COMPASS results on Transverse Spin Asymmetries in SIDIS

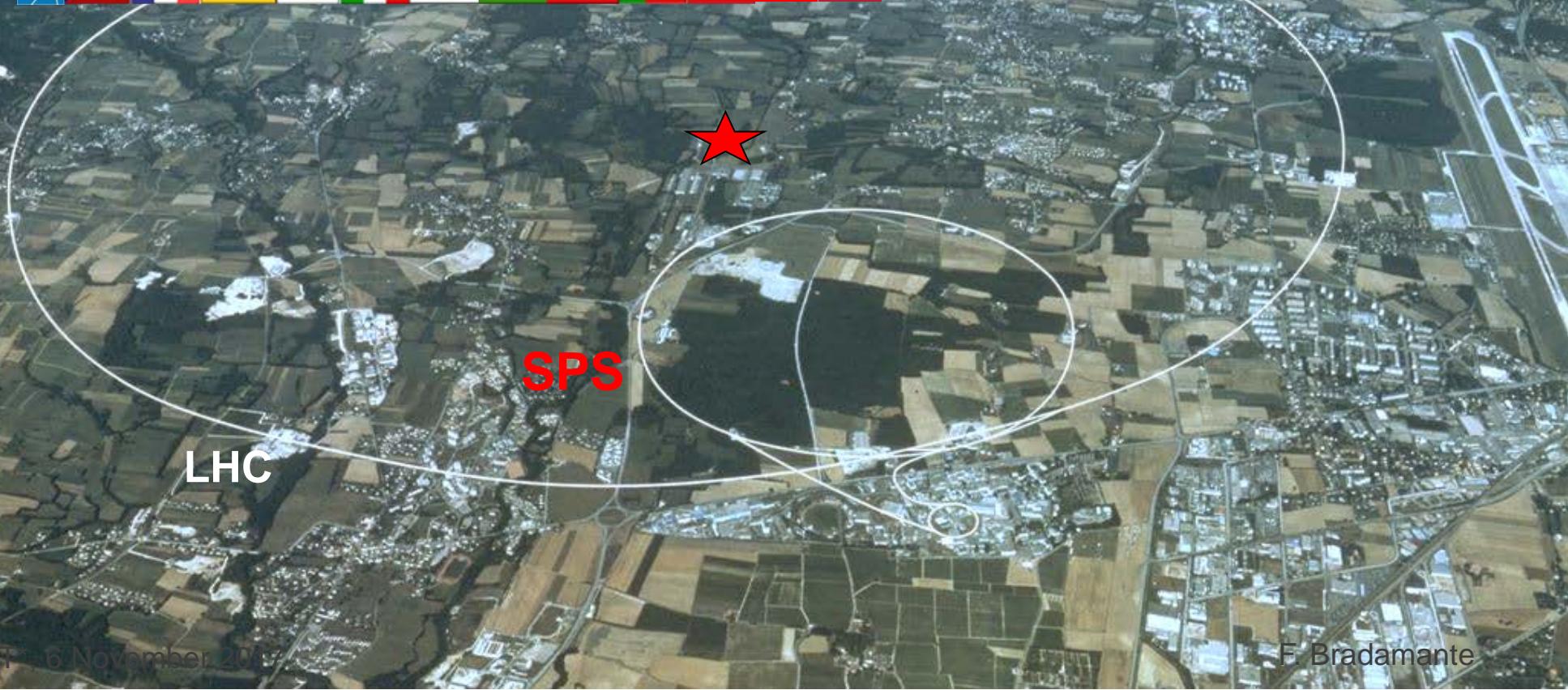
Franco Bradamante

INFN Trieste, Italy

on behalf of the COMPASS Collaboration



fixed target experiment at the CERN SPS





*CO*mmun
*M*uon and
*P*roton
*A*pparatus for
*S*tructure and
*S*pectroscopy

fixed target experiment at the CERN SPS



physics programme:

hadron spectroscopy (p, π, K)

- light mesons, glue-balls, exotic mesons
- polarisability of pion and kaon

nucleon structure (μ)

- longitudinal spin structure
- transverse momentum and transverse spin structure



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this talk

COMPASS spectrometer



designed to

- **use high energy beams**
- **have large angular acceptance**
- **cover a broad kinematical range**

COMPASS spectrometer



designed to

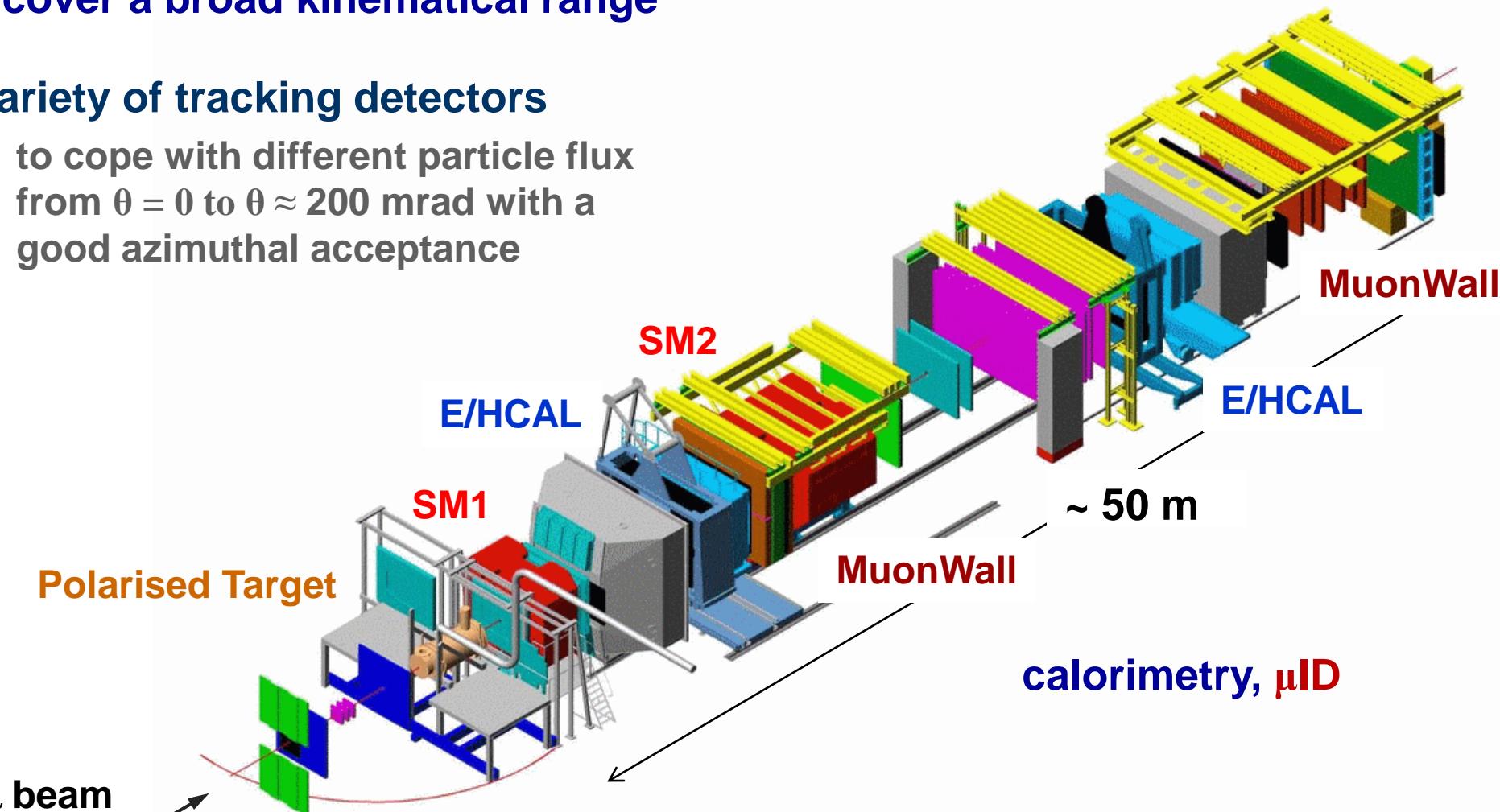
- use high energy beams
- have large angular acceptance
- cover a broad kinematical range

two stages spectrometer

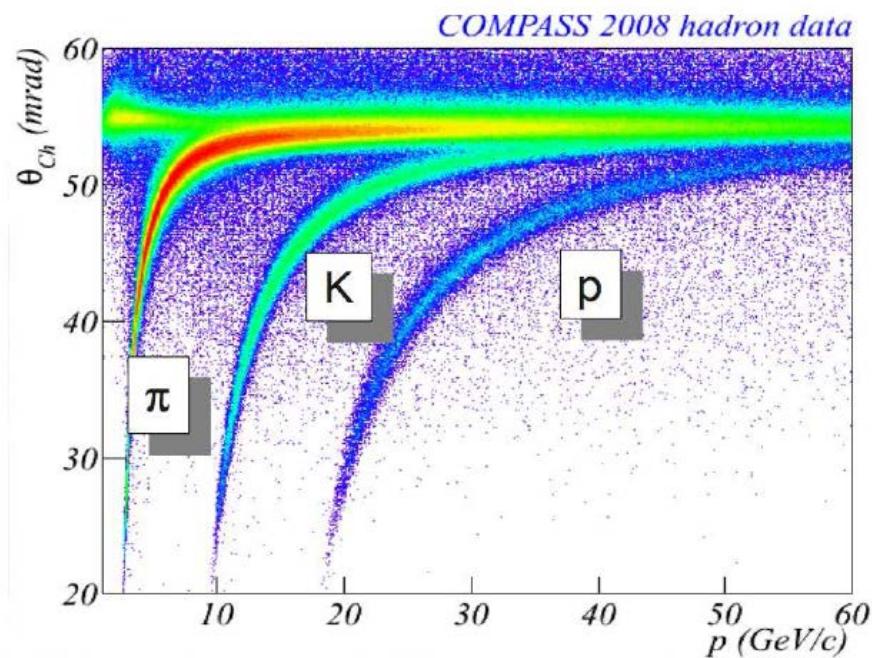
- Large Angle Spectrometer (**SM1**)
- Small Angle Spectrometer (**SM2**)

variety of tracking detectors

to cope with different particle flux
from $\theta = 0$ to $\theta \approx 200$ mrad with a
good azimuthal acceptance

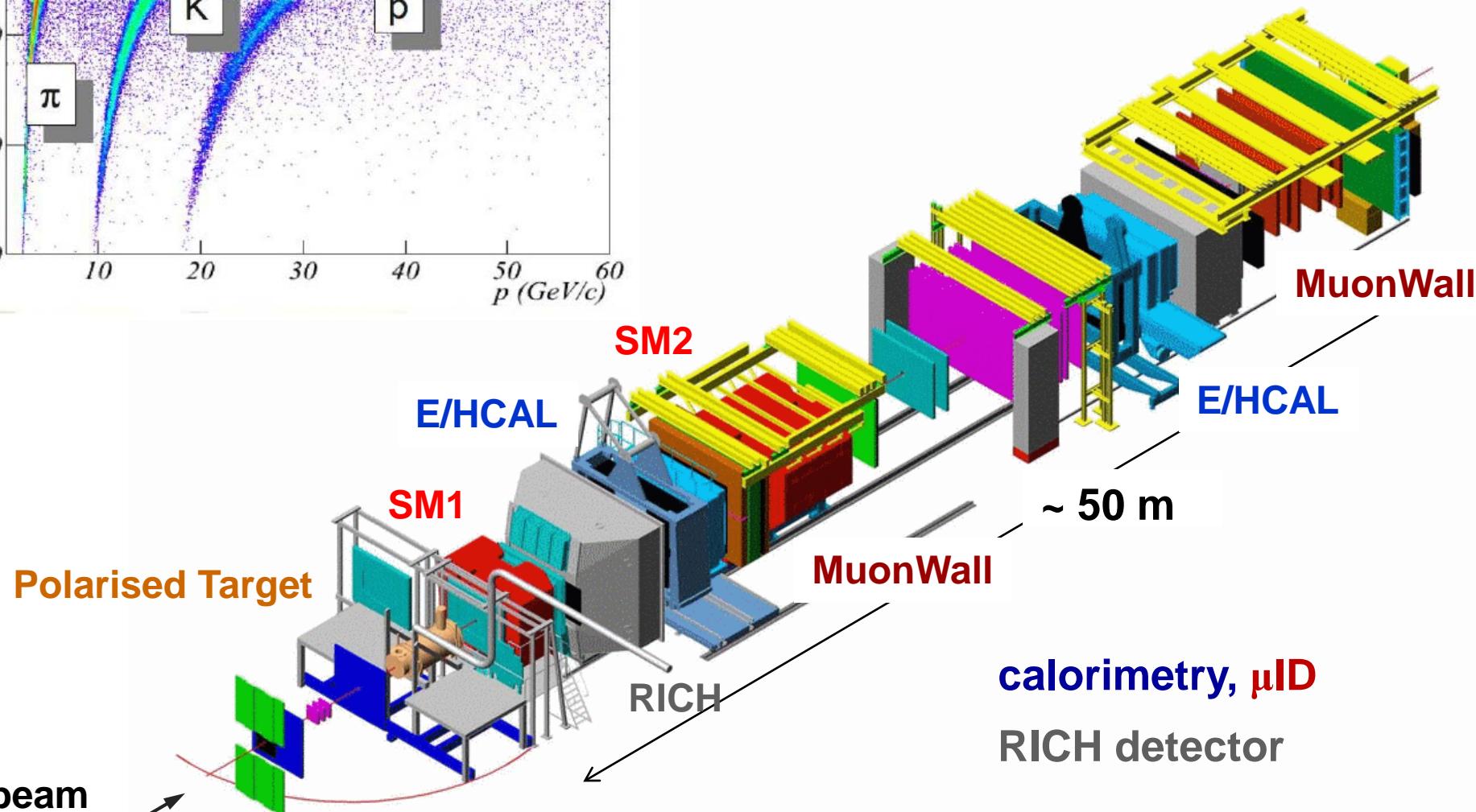


COMPASS spectrometer



two stages spectrometer

- Large Angle Spectrometer (**SM1**)
- Small Angle Spectrometer (**SM2**)

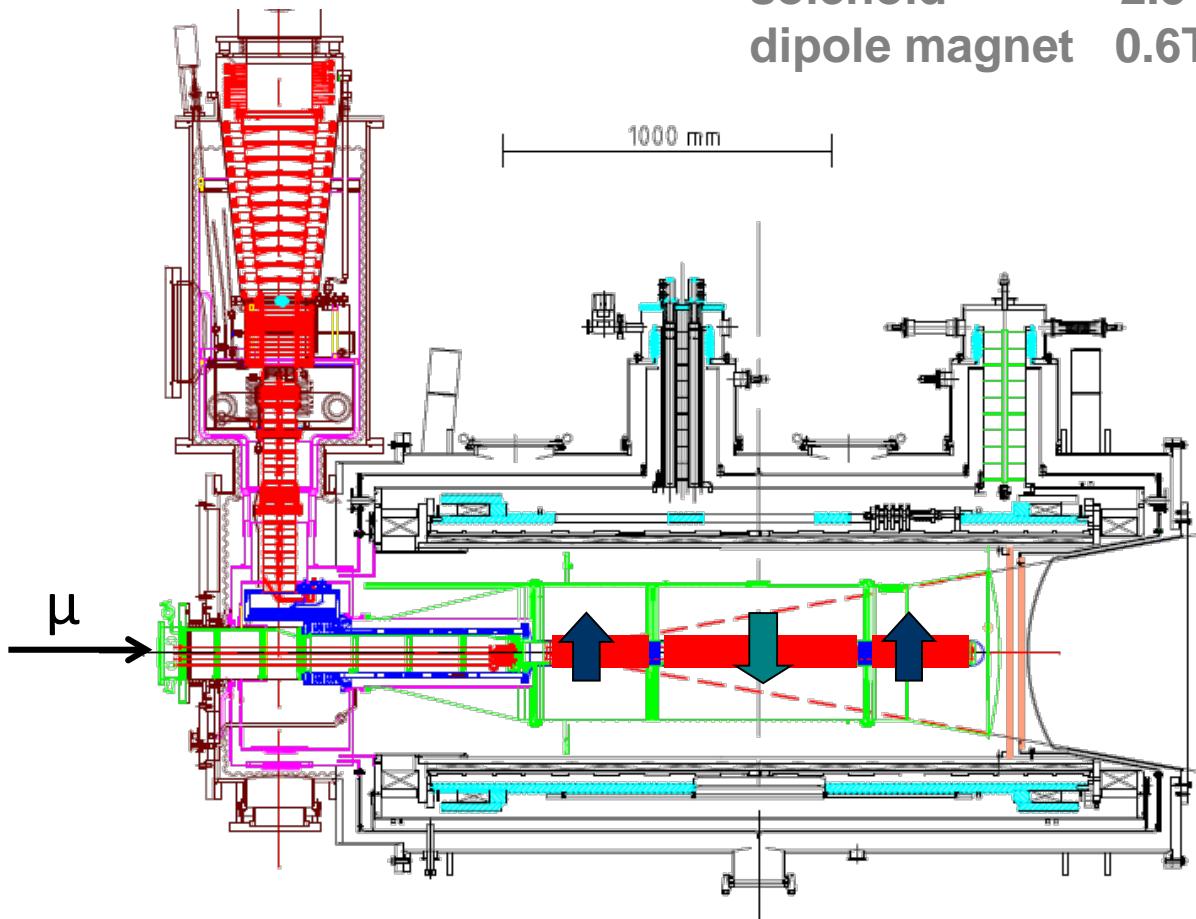


the polarized target system (>2005)



${}^3\text{He} - {}^4\text{He}$ dilution refrigerator ($T \sim 50\text{mK}$)

solenoid 2.5T
dipole magnet 0.6T



acceptance $> \pm 180$ mrad

3 target cells
30, 60, and 30 cm long

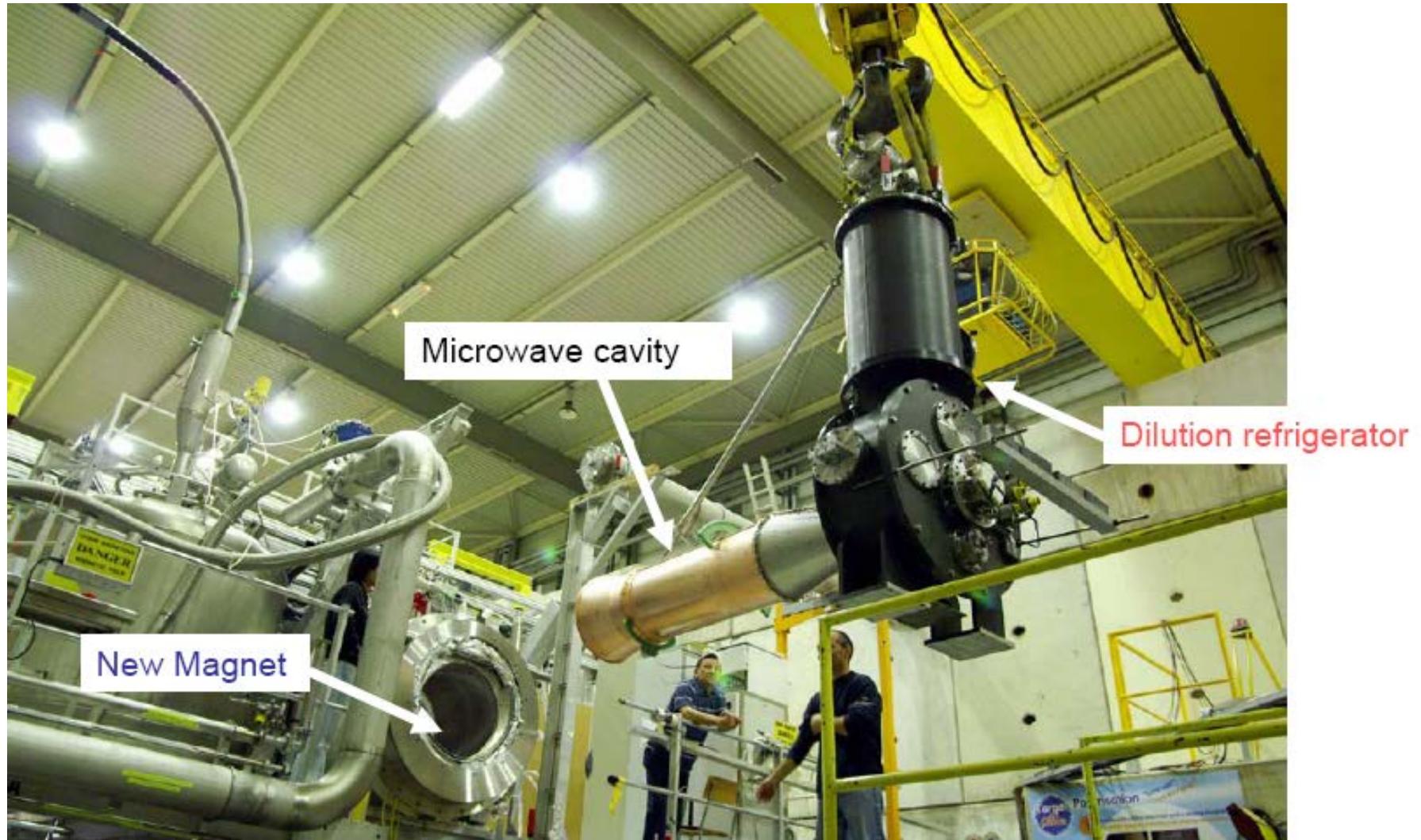
opposite polarisation

d (${}^6\text{LiD}$)	p (NH_3)
50%	90%
40%	16%

polarization
dilution factor

***no evidence for relevant
nuclear effects (160 GeV)***

the polarized target system



COMPASS data taking



COMPASS data taking



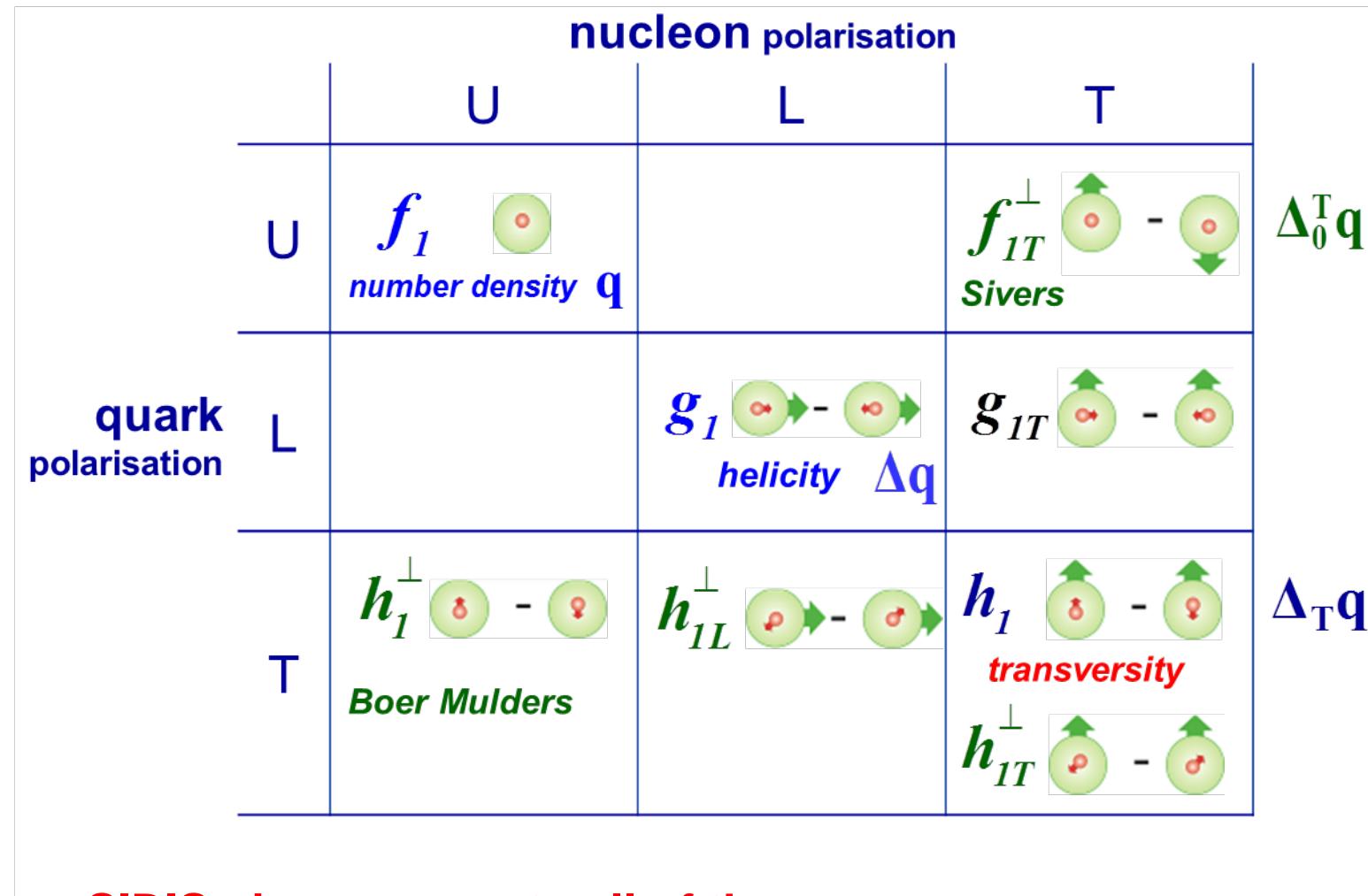
2002	nucleon structure with	160 GeV μ	L&T	polarised deuteron target
2003	nucleon structure with	160 GeV μ	L&T	polarised deuteron target
2004	nucleon structure with	160 GeV μ	L&T	polarised deuteron target
2005	<i>CERN accelerators shut down</i>			
2006	nucleon structure with	160 GeV μ	L	polarised deuteron target
2007	nucleon structure with	160 GeV μ	L&T	polarised proton target
2008	<i>hadron spectroscopy</i>			
2009	<i>hadron spectroscopy</i>			
2010	nucleon structure with	160 GeV μ	T	polarised proton target
2011	nucleon structure with	190 GeV μ	L	polarised proton target
2012	Primakoff & DVCS / SIDIS test			
2013	<i>CERN accelerators shut down</i>			
2014	Test beam Drell-Yan process with π beam and T polarised proton target			
2015	Drell-Yan process with π beam and T polarised proton target			
2016	DVCS / SIDIS with μ beam and unpolarised proton target			
2017	DVCS / SIDIS with μ beam and unpolarised proton target			
2018	Drell-Yan process with π beam and T polarised proton target			

MUON beam PROGRAM:

**TRANSVERSITY and
Transverse Momentum Dependent PDFs**

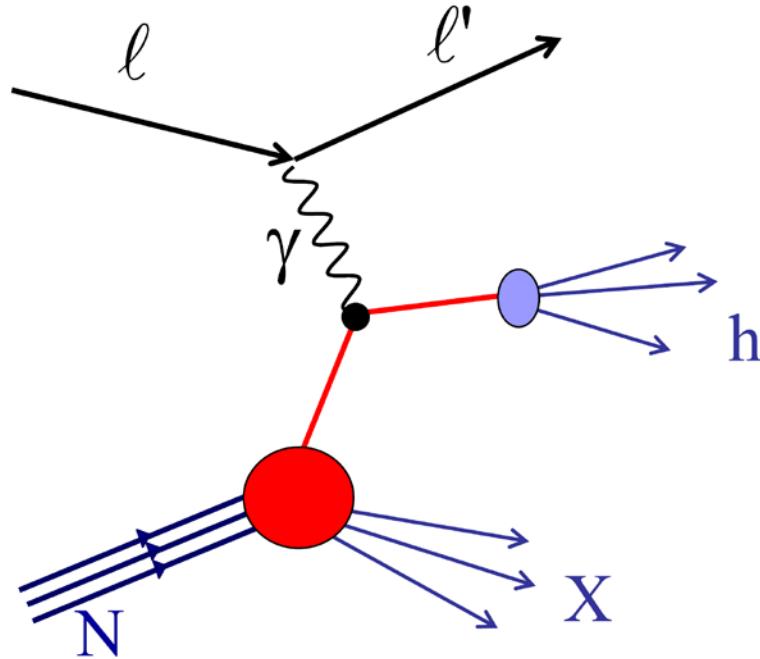
the structure of the nucleon

taking into account the quark intrinsic transverse momentum k_T ,
at leading order other 6 TMD PDFs are needed for a full description
of the nucleon structure



Semi-Inclusive Deep Inelastic Scattering

hard interaction of a lepton with a nucleon via virtual photon exchange

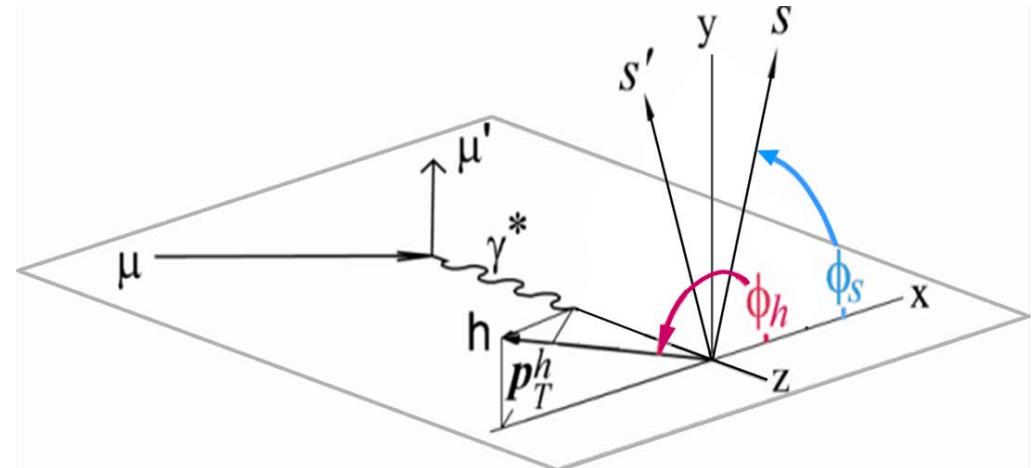


$$\sigma^{lN \rightarrow lhX} \propto \sum_q f(x) \otimes \sigma^{lq \rightarrow lq} \otimes D_q^h(z)$$

$$x = \frac{Q^2}{2P \cdot q} \quad y = \frac{P \cdot q}{P \cdot \ell} =_{LAB} \frac{E - E'}{E}$$

$$Q^2 = -q^2 \quad W^2 = (P + q)^2$$

$$z = \frac{P \cdot P_h}{P \cdot q} =_{LAB} \frac{E_h}{E - E'}$$



Semi-Inclusive Deep Inelastic Scattering

$$\begin{aligned}
& \frac{d\sigma}{dx dy d\psi dz d\phi_h dP_{h\perp}^2} = \\
& \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos \phi_h F_{UU}^{\cos \phi_h} \right. \\
& \quad \left. + \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin \phi_h F_{LU}^{\sin \phi_h} \right. \\
& \quad \left. + S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin \phi_h F_{UL}^{\sin \phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right] + S_{\parallel} \lambda_e \left[\sqrt{1-\varepsilon^2} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos \phi_h F_{LL}^{\cos \phi_h} \right] \right. \\
& \quad \left. + |\mathbf{S}_{\perp}| \left[\sin(\phi_h - \phi_S) \left(F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) \right. \right. \\
& \quad \left. \left. + \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} + \varepsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)} \right. \right. \\
& \quad \left. \left. + \sqrt{2\varepsilon(1+\varepsilon)} \sin \phi_S F_{UT}^{\sin \phi_S} + \sqrt{2\varepsilon(1+\varepsilon)} \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)} \right] \right. \\
& \quad \left. + |\mathbf{S}_{\perp}| \lambda_e \left[\sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} + \sqrt{2\varepsilon(1-\varepsilon)} \cos \phi_S F_{LT}^{\cos \phi_S} \right. \right. \\
& \quad \left. \left. + \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \right] \right\},
\end{aligned}$$

unpol target

→ pol target

↑ pol target

18 structure functions

Semi-Inclusive Deep Inelastic Scattering

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& + \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin \phi_h F_{LU}^{\sin \phi_h} \\
& + S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin \phi_h F_{UL}^{\sin \phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right] + S_{\parallel} \lambda_e \left[\sqrt{1-\varepsilon^2} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos \phi_h F_{LL}^{\cos \phi_h} \right] \\
& + |S_{\perp}| \left[\sin(\phi_h - \phi_S) \left(F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) \right. \\
& + \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} + \varepsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)} \\
& \left. + \sqrt{2\varepsilon(1+\varepsilon)} \sin \phi_S F_{UT}^{\sin \phi_S} + \sqrt{2\varepsilon(1+\varepsilon)} \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)} \right] \\
& + |S_{\perp}| \lambda_e \left[\sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} + \sqrt{2\varepsilon(1-\varepsilon)} \cos \phi_S F_{LT}^{\cos \phi_S} \right. \\
& \left. + \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \right] \left. \right\},
\end{aligned}$$

14 independent azimuthal modulations

Semi-Inclusive Deep Inelastic Scattering

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& + S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin \phi_h F_{UL}^{\sin \phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right] + S_{\parallel} \lambda_e \left[\sqrt{1-\varepsilon^2} \right. \\
& + |S_{\perp}| \left[\sin(\phi_h - \phi_S) \left(F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) \right. \\
& + \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} + \varepsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)} \\
& + \sqrt{2\varepsilon(1+\varepsilon)} \sin \phi_S F_{UT}^{\sin \phi_S} + \sqrt{2\varepsilon(1+\varepsilon)} \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)} \left. \right] \\
& + |S_{\perp}| \lambda_e \left[\sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} + \sqrt{2\varepsilon(1-\varepsilon)} \cos \phi_S F_{LT}^{\cos \phi_S} \right. \\
& \left. \left. + \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \right] \right\},
\end{aligned}$$

14 independent azimuthal modulations

amplitudes of the modulations
→ TMD PDFs

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& + |S_{\perp}| \left[\sin(\phi_h - \phi_S) \left(F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon \right. \right. \\
& + \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} + \varepsilon \cos \phi_h F_{UT}^{\cos \phi_h} \\
& + \sqrt{2\varepsilon(1+\varepsilon)} \sin \phi_S F_{UT}^{\sin \phi_S} + \sqrt{2\varepsilon(1-\varepsilon)} \cos \phi_S F_{UT}^{\cos \phi_S} \\
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& \left. \left. + \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \right] \right\},
\end{aligned}$$

14 independent azimuthal modulations

amplitudes of the modulations
→ TMD PDFs

SIDIS

- allows to disentangle the effects related to the different TMD PDFs and to access all of them
- by identifying the final state hadrons and using different targets allows for flavour separation
→ very powerful tool

all the amplitudes (AA)
have been measured
in COMPASS

TMDs in unpolarised SIDIS

$$\frac{d\sigma}{dx dy d\psi dz d\phi_h dP_{h\perp}^2} = \frac{\alpha^2}{xy Q^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ \begin{array}{l} k_T \\ k_T \\ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos \phi_h F_{UU}^{\cos \phi_h} \\ + \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin \phi_h F_{LU}^{\sin \phi_h} \\ h_I^\perp H_I^\perp \\ + \dots \end{array} \right\}$$

unpolarised SIDIS

Relevance for TMDs:

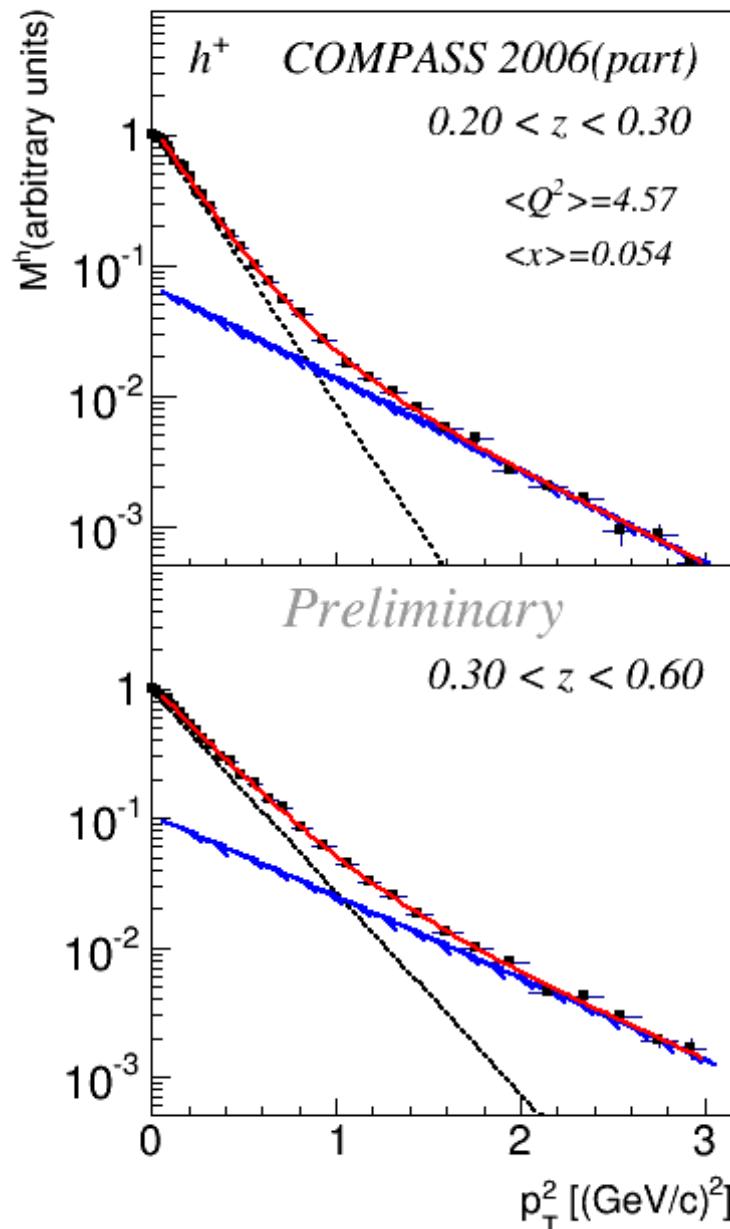
- the cross-section dependence on p_{Th} comes from:
 - intrinsic k_T of the quarks
 - p_\perp generated in the quark fragmentation
$$\langle p_{Th}^2 \rangle = \langle p_\perp^2 \rangle + z^2 \langle k_T^2 \rangle$$
- the azimuthal modulations in the unpolarized cross-sections comes from:
 - intrinsic k_T of the quarks
 - Boer-Mulders PDF

combined analysis should allow to disentangle the different effects

COMPASS

- has produced results on 6LiD ($\sim d$) from 2004/6 data
- will measure SIDIS on LH_2 in parallel with DVCS

unpolarised SIDIS – p_{Th} distributions



deuteron

Fit distributions with

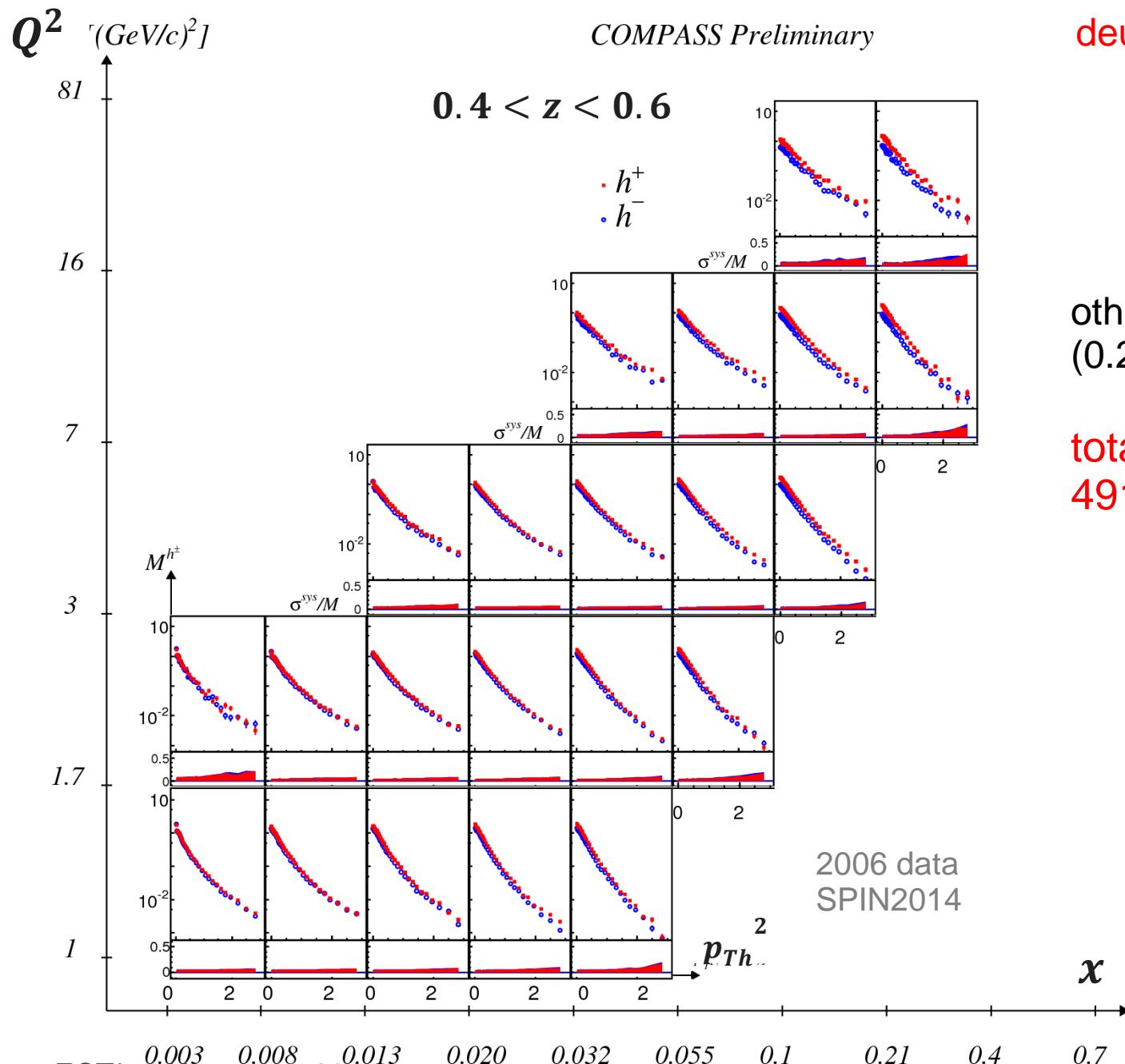
- 1 exponential for $p_{Th}^2 \in [0.05, 0.68]$
- 2 exponentials for $p_{Th}^2 \in [0.05, 3]$



needed to describe the shape
of p_{Th}^2 the COMPASS data

Transversity 2014

unpolarised SIDIS – p_{Th} distributions



deuteron

other 3 z bins
 $(0.2-0.3, 0.3-0.4, 0.6-0.8)$

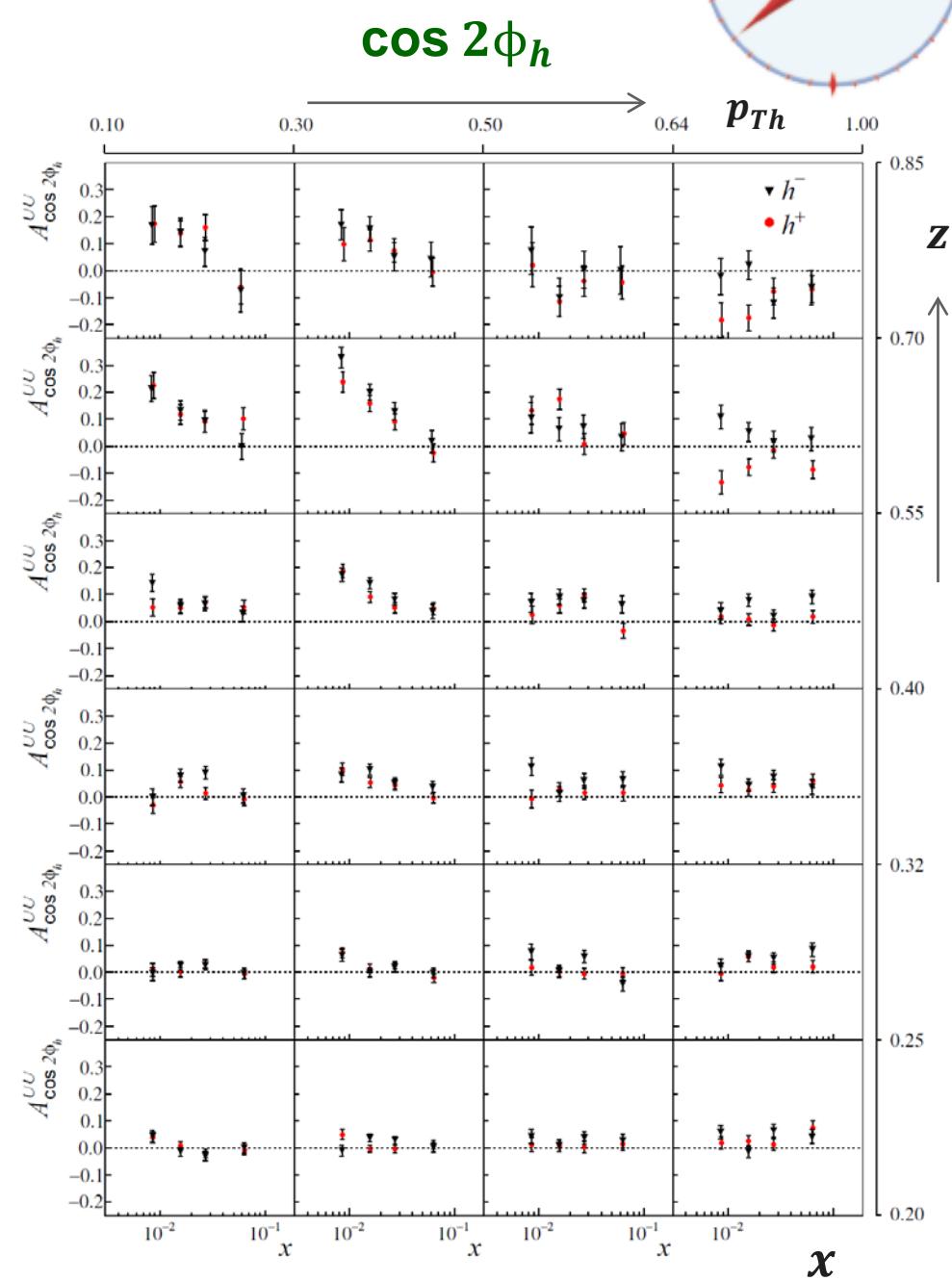
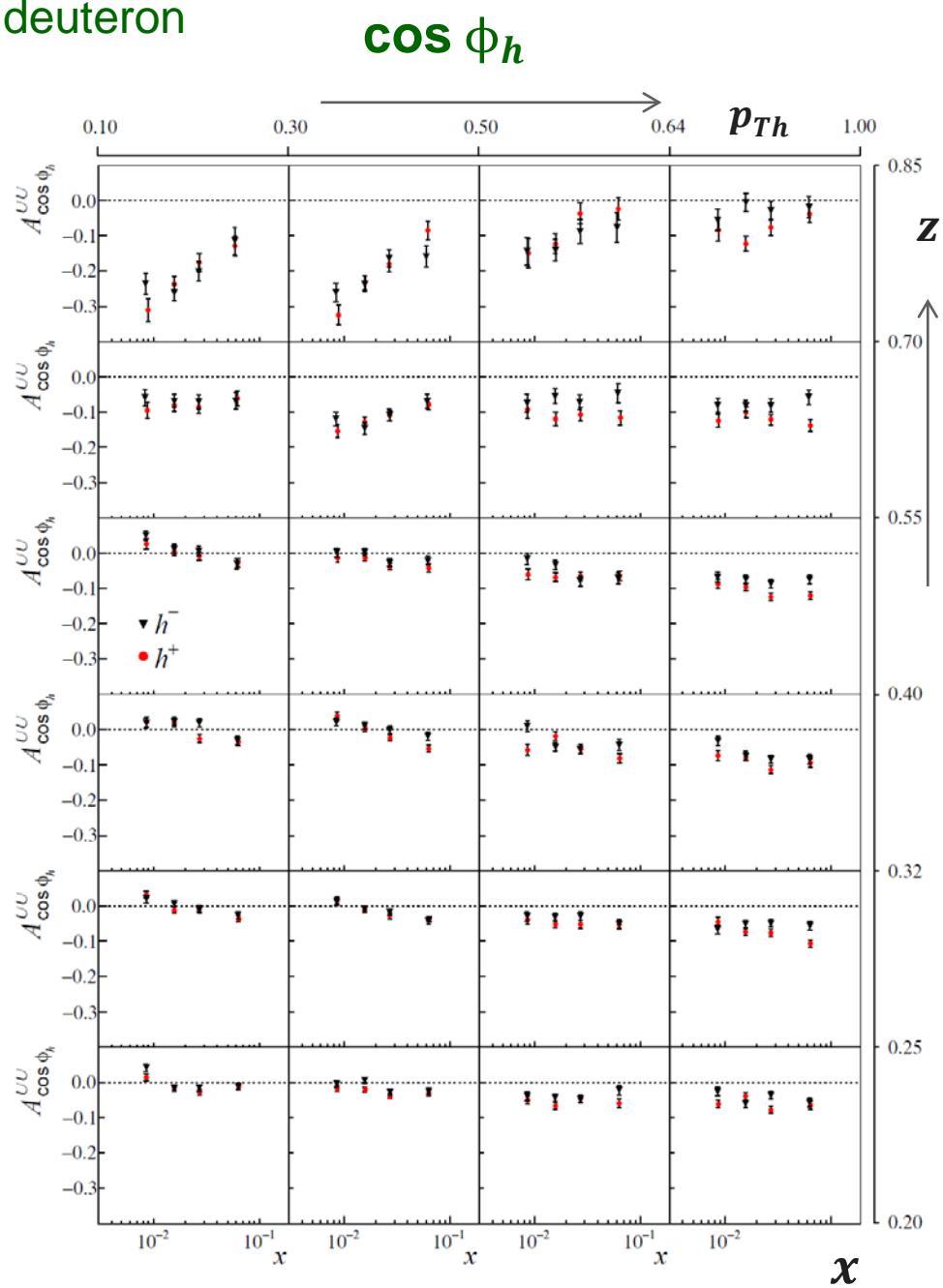
total:
 4918 data points

paper ready,
 to be sent to PRD

unpolarised SIDIS – azimuthal asymmetries



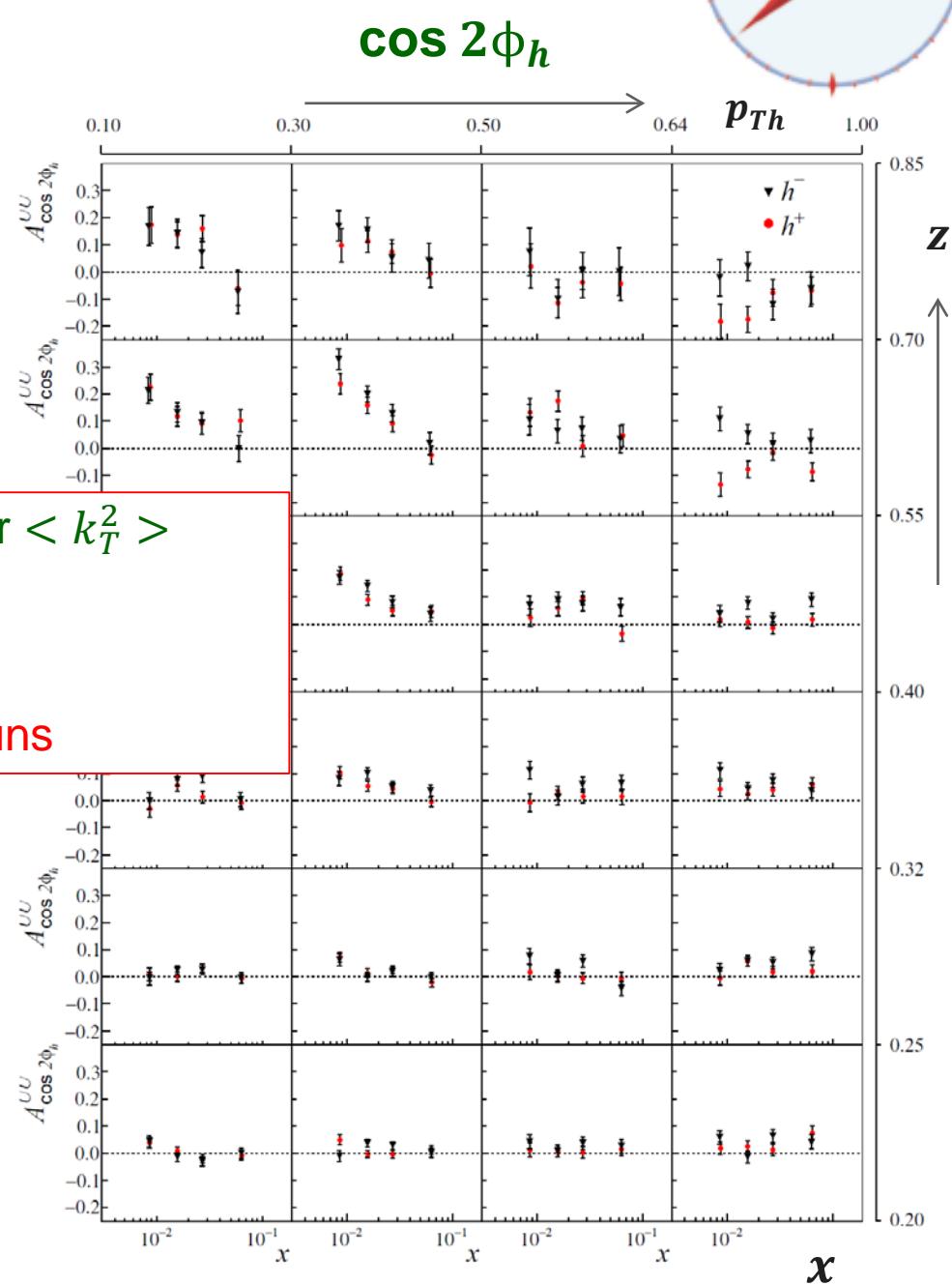
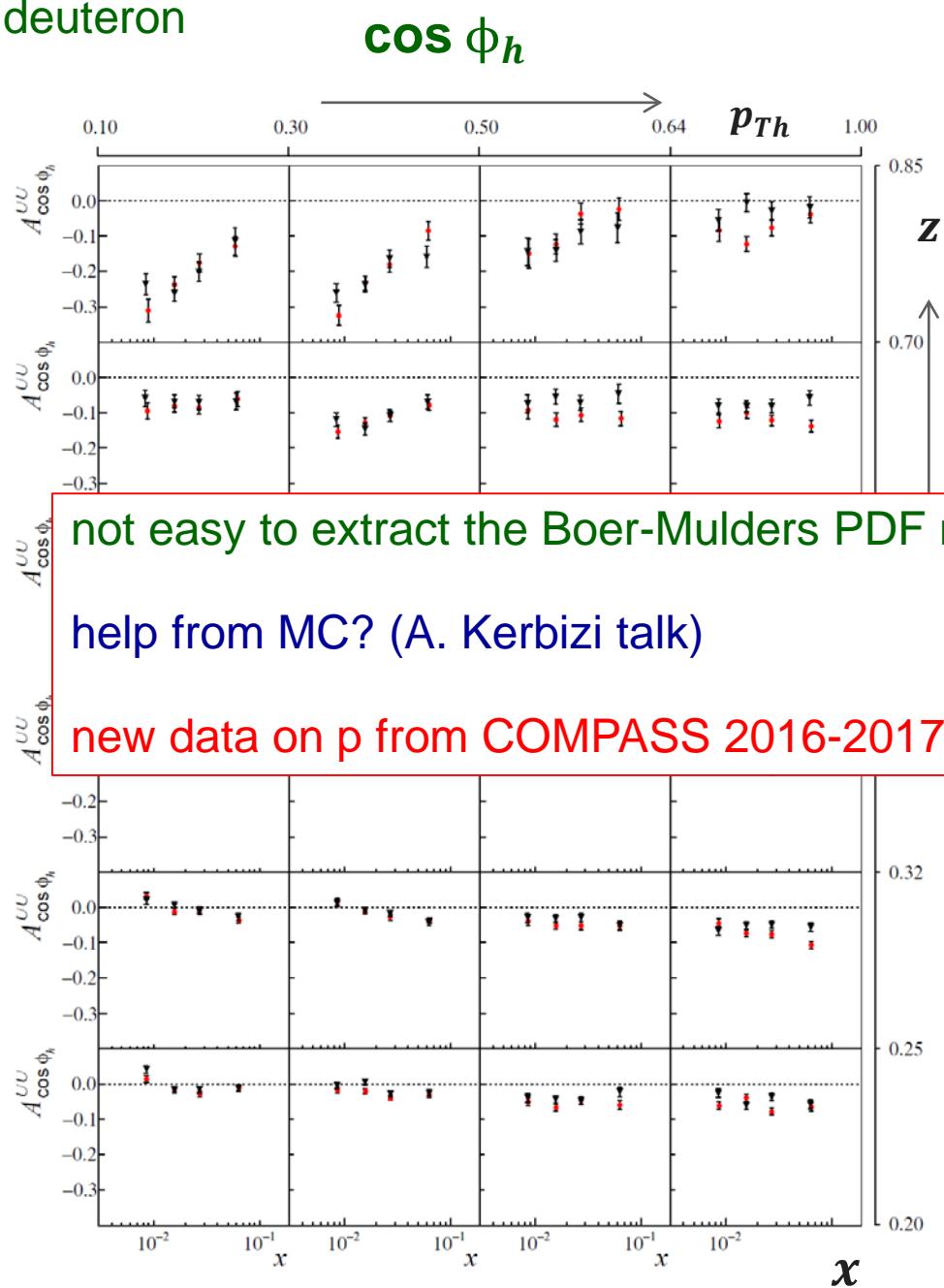
deuteron



unpolarised SIDIS – azimuthal asymmetries



deuteron



not easy to extract the Boer-Mulders PDF nor $\langle k_T^2 \rangle$

help from MC? (A. Kerbizi talk)

new data on p from COMPASS 2016-2017 runs

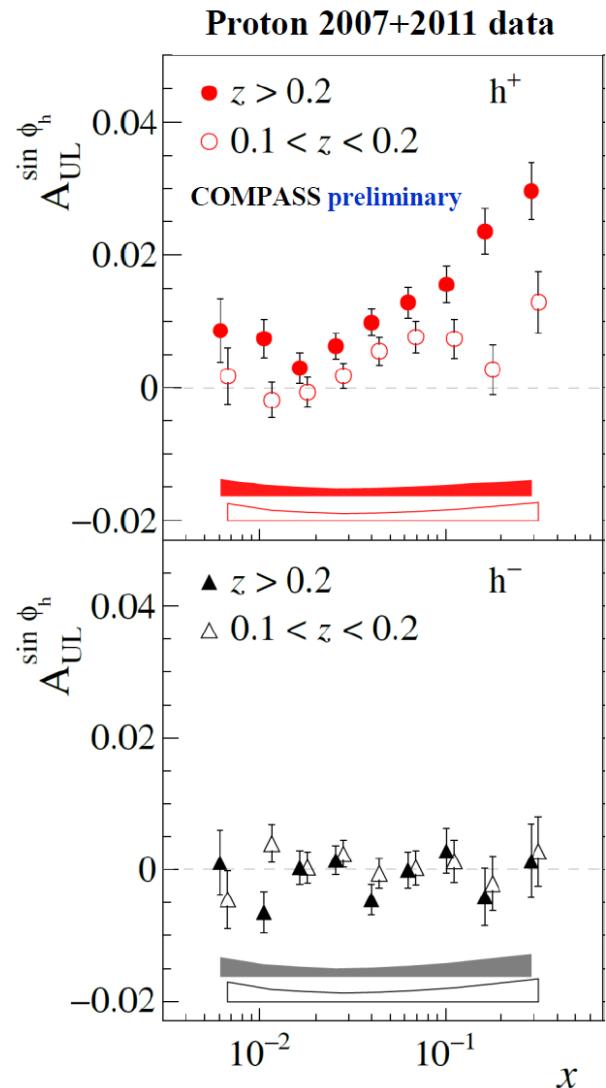
TMDs in SIDIS off longitudinally polarised N

$$\frac{d\sigma}{dx dy d\psi dz d\phi_h dP_{h\perp}^2} = \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ \dots \right.$$
$$+ S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin \phi_h F_{UL}^{\sin \phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right]$$
$$+ S_{\parallel} \lambda_e \left[\sqrt{1-\varepsilon^2} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos \phi_h F_{LL}^{\cos \phi_h} \right]$$
$$\left. + \dots \right\}$$

SIDIS off longitudinally polarised p



$$A_{UL}^{\sin \phi_h} = F_{UL}^{\sin \phi_h} / F_{UU}$$



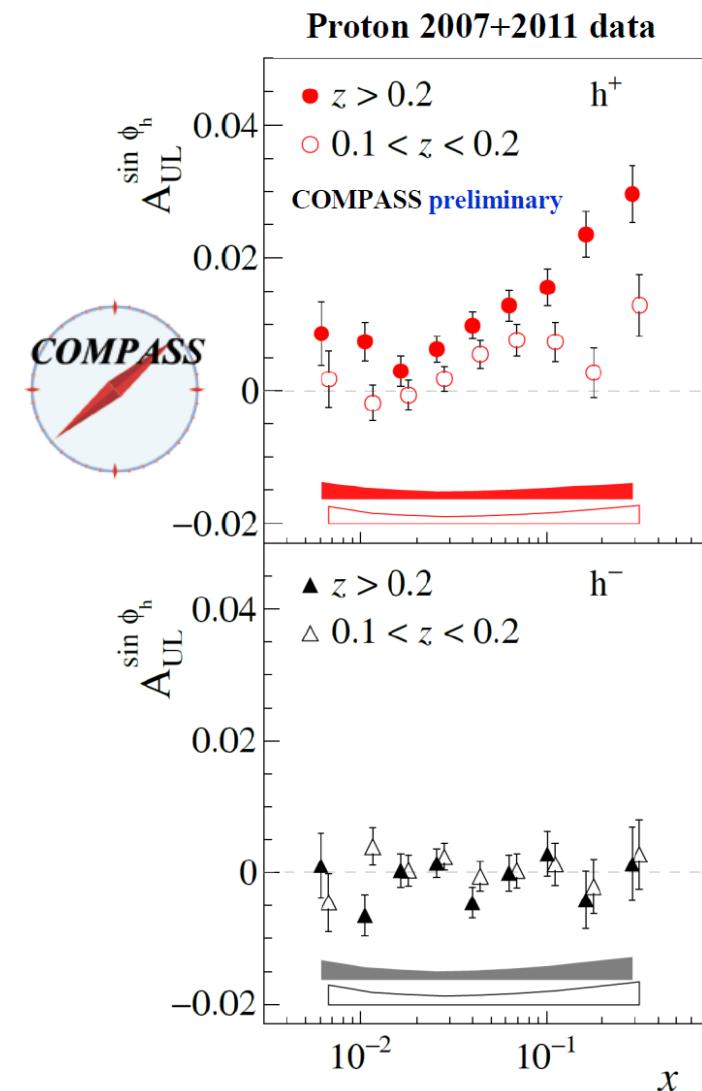
SIDIS off longitudinally polarised p

$$A_{UL}^{\sin \phi_h} = F_{UL}^{\sin \phi_h} / F_{UU}$$

$$F_{UL}^{\sin \phi_h} = \frac{2M}{Q} \mathcal{C} \left\{ -\frac{\hat{\boldsymbol{h}} \cdot \boldsymbol{p}_T}{M_h} \left(x h_L^q H_{1q}^{\perp h} + \frac{M_h}{M} g_{1L}^q \frac{\tilde{G}_q^{\perp h}}{z} \right) \right.$$

$$\left. + \frac{\hat{\boldsymbol{h}} \cdot \boldsymbol{k}_T}{M} \left(x f_L^{\perp q} D_{1q}^h - \frac{M_h}{M} h_{1L}^{\perp q} \frac{\tilde{H}_q^h}{z} \right) \right\}$$

Q-suppressed,
different “twist” contributions



SIDIS off longitudinally polarised p

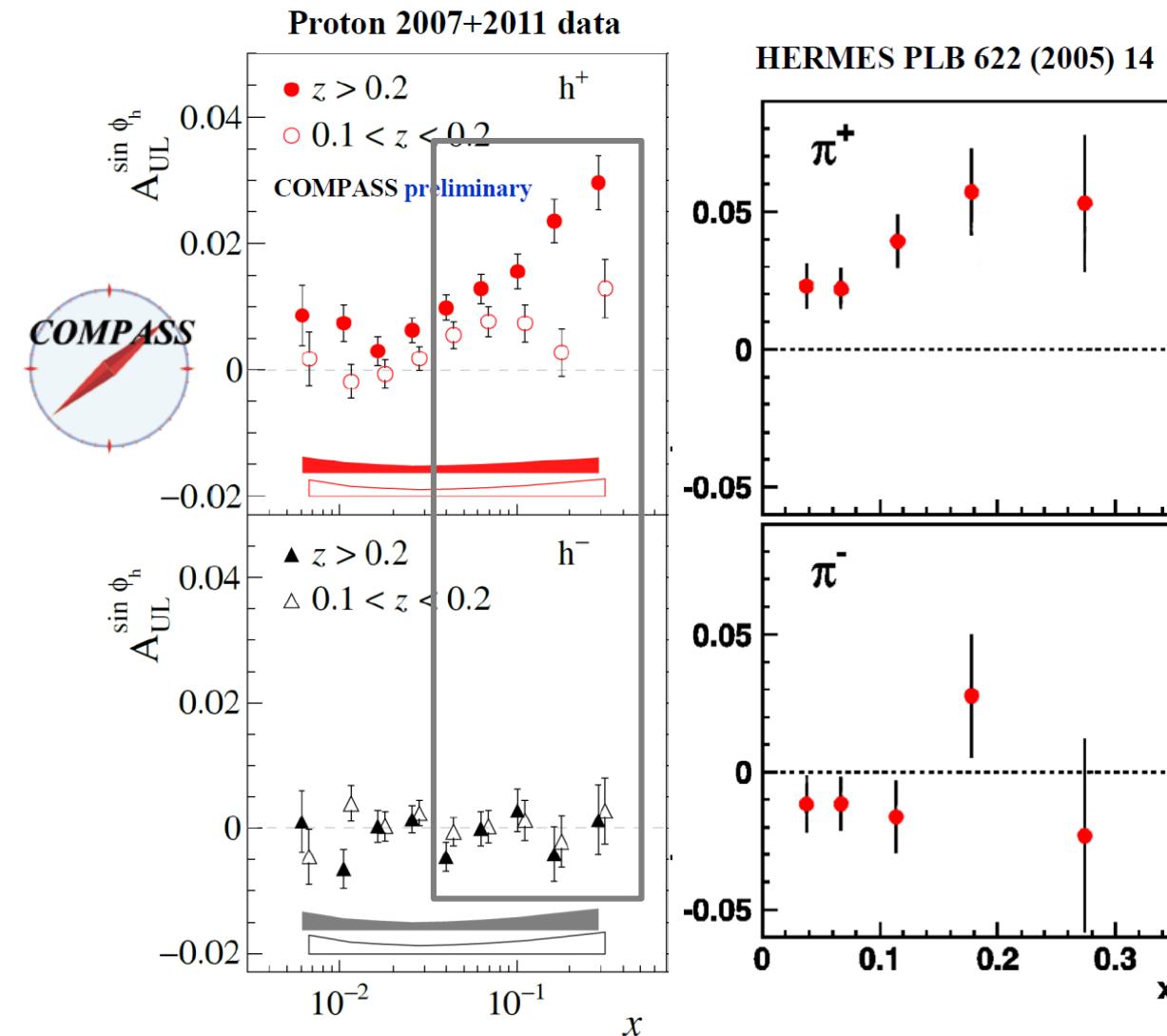
$$A_{UL}^{\sin \phi_h} = F_{UL}^{\sin \phi_h} / F_{UU}$$

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$$\left. + \frac{\hat{h} \cdot \mathbf{k}_T}{M} \left(x f_L^{\perp q} D_{1q}^h - \frac{M_h}{M} h_{1L}^{\perp q} \frac{\tilde{H}_q^h}{z} \right) \right\}$$

Q-suppressed,
different “twist” contributions

$$\sqrt{2 \epsilon(1 + \epsilon)}$$



SIDIS off transversely polarised N

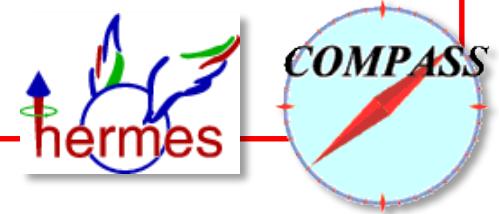
$$\begin{aligned}
\frac{d\sigma}{dx dy d\psi dz d\phi_h dP_{h\perp}^2} = & \dots \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \Bigg\{ \\
& + |\mathbf{S}_\perp| \left[\begin{aligned}
& f_{IT}^\perp D_L \\
& \sin(\phi_h - \phi_S) \left(F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) \\
& h_{IT}^\perp H_L^\perp \\
& + \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} + \varepsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)} \\
& h_L H_L^\perp \\
& + \sqrt{2\varepsilon(1+\varepsilon)} \left(\sin \phi_S F_{UT}^{\sin \phi_S} + \sqrt{2\varepsilon(1+\varepsilon)} \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)} \right) \\
& g_{IT} D_L \\
& + |\mathbf{S}_\perp| \lambda_e \left[\begin{aligned}
& \sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} + \sqrt{-\varepsilon(1-\varepsilon)} \cos \phi_S F_{LT}^{\cos \phi_S} \\
& + \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \end{aligned} \right] \Bigg\},
\end{aligned}$$

Semi-Inclusive Deep Inelastic Scattering

MAJOR RESULT:

in the past 10 years 2 of these new PDF's have been measured and shown to be different from zero

by COMPASS and HERMES



the transversity PDF

amplitude of the sine modulation in $\phi_h + \phi_s - \pi$
Collins asymmetry $\sim h_1 \otimes H_1^\perp$

the Sivers PDF

amplitude of the sine modulation in $\phi_h - \phi_s$
Sivers asymmetry $\sim f_{IT}^\perp \otimes D_1$

A STEP TOWARDS
THE 3-D STRUCTURE OF THE NUCLEON

Collins asymmetry

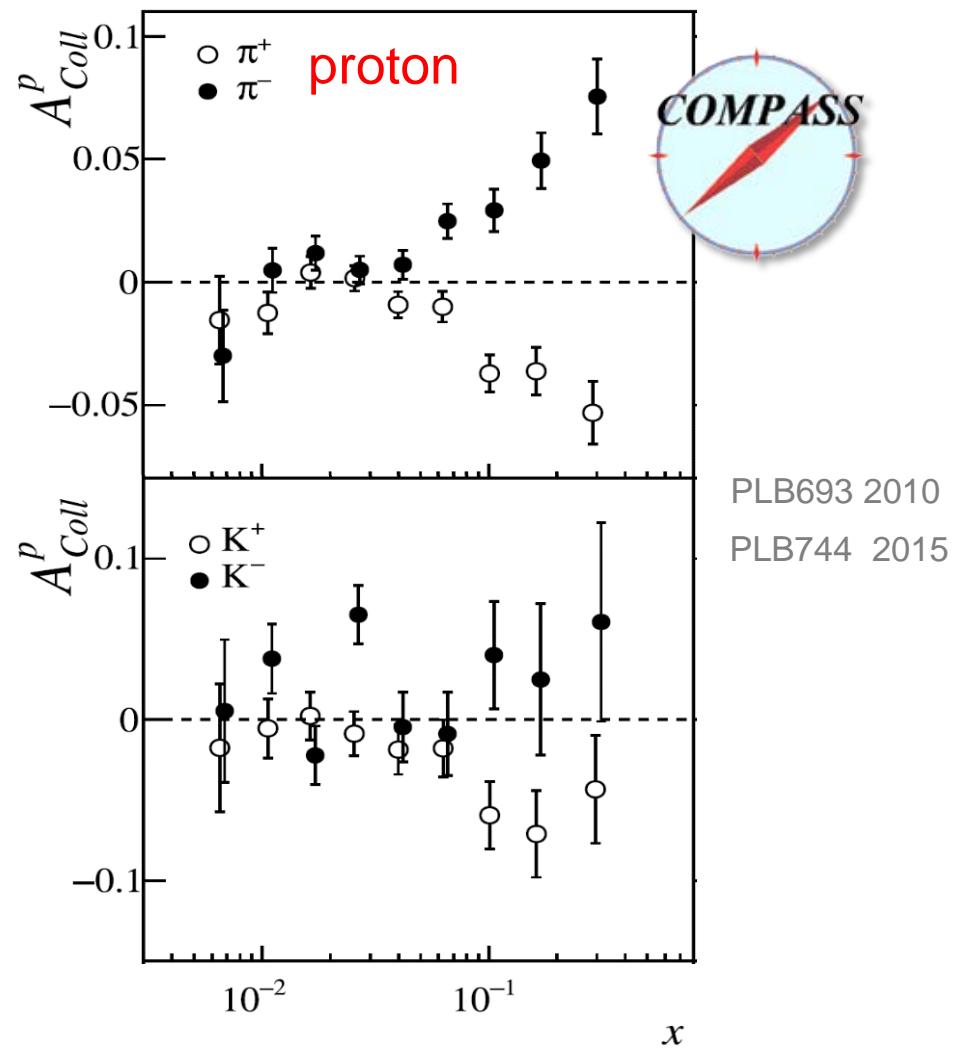
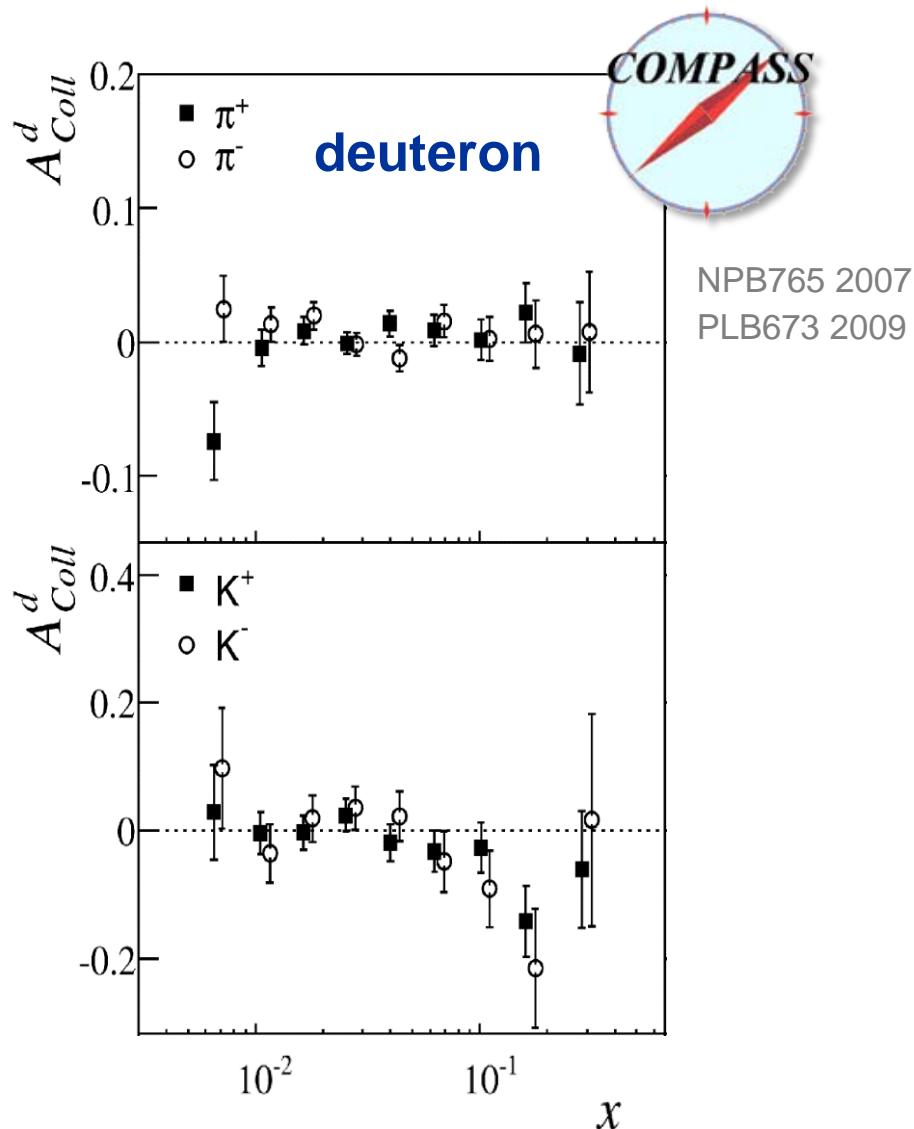
Collins asymmetry

$$\sim h_l \otimes H_l^\perp$$

2004: first evidence for non-zero Collins asymmetry on p from HERMES



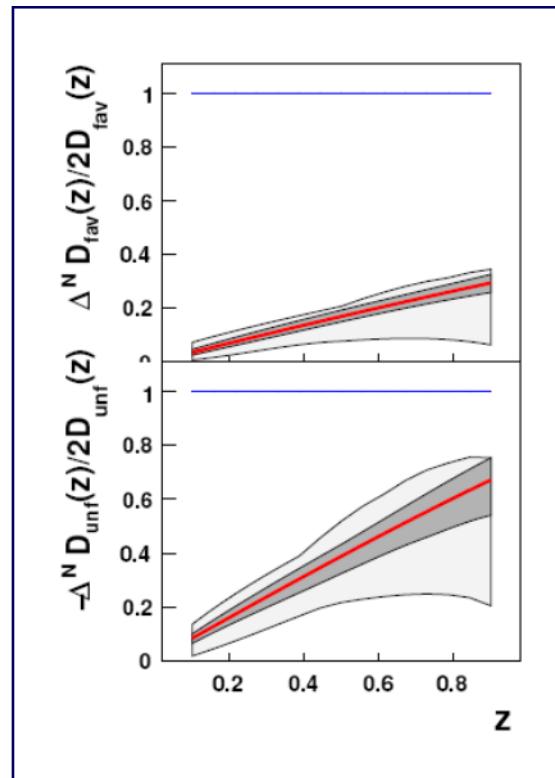
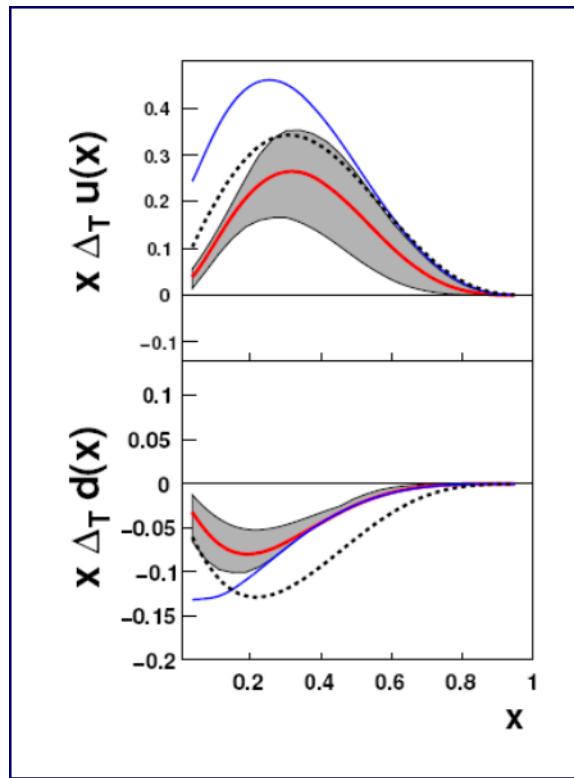
final COMPASS results



Transversity from SIDIS

M. Anselmino et al., Nucl. Phys. Proc. Suppl. 2009

fit to HERMES p, COMPASS d, Belle e+e- data



and many others ...

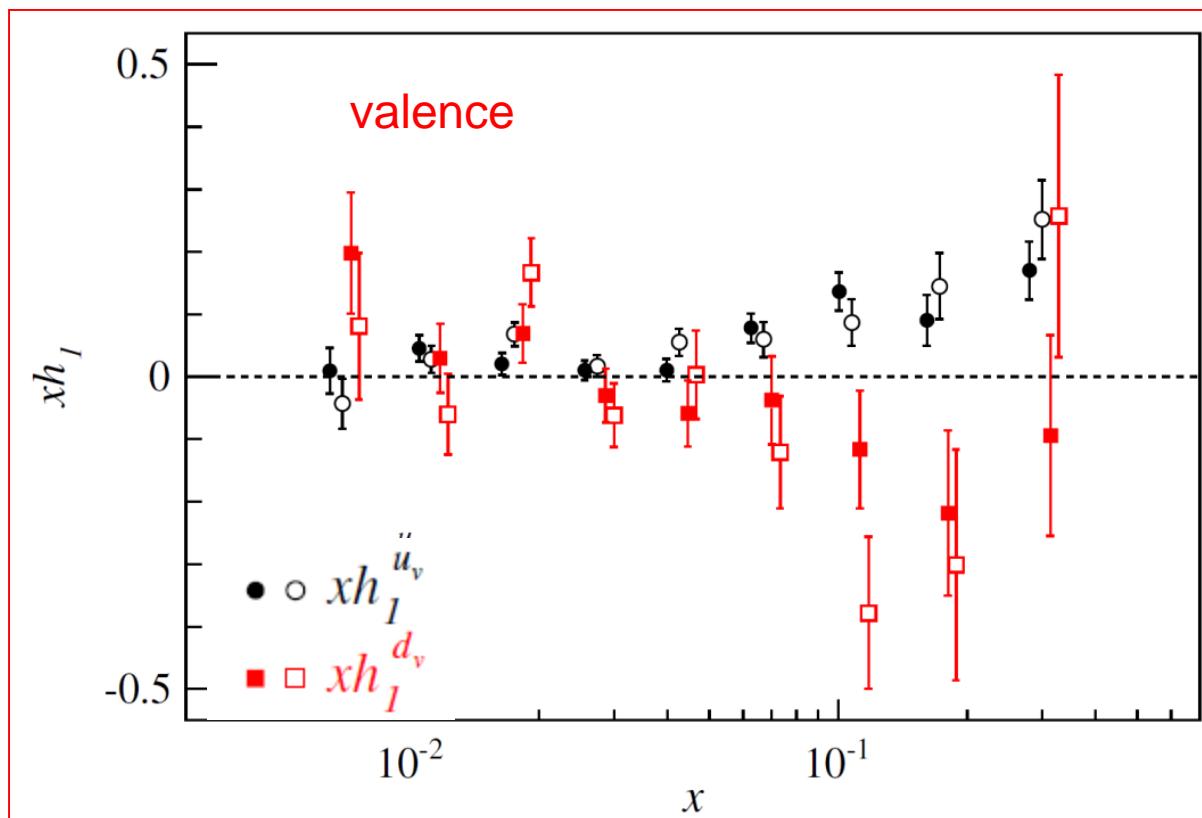
Transversity from SIDIS

Collins and di-hadron asymmetries

point by point extraction

one can use directly the COMPASS p and d asymmetries,
and the Belle data to evaluate the analysing power
(with some “reasonable” assumptions)

advantage: no MC nor parametrisation is needed



open points: dihadron
closed points: Collins

large uncertainties
on the d distribution
due to the poor
deuteron/neutron
data

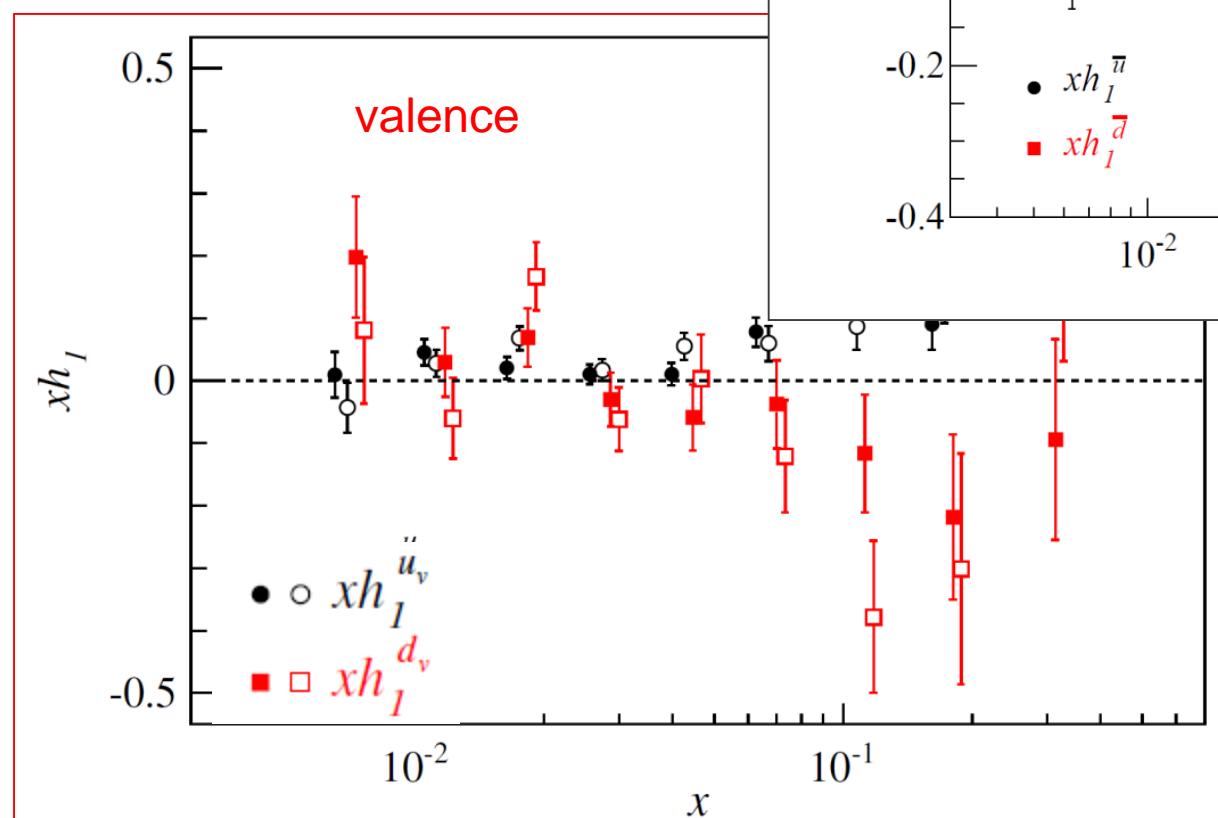
A. Martin F. B. V. Barone
PRD91 2015

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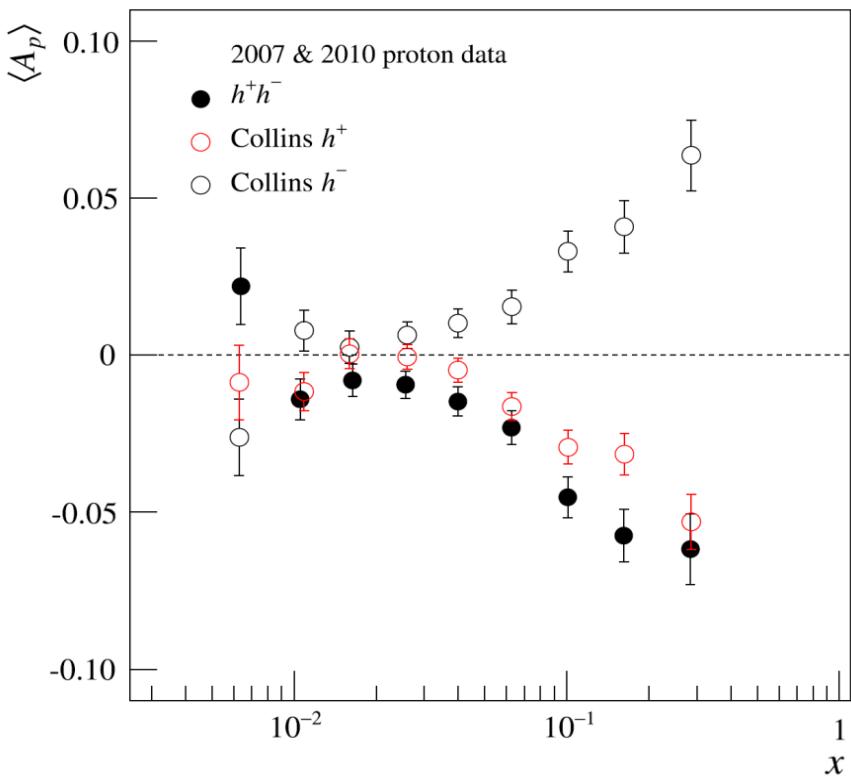
A. Martin F. B. V. Barone
PRD91 2015

large uncertainties
on the d distribution
due to the poor
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data

Transversity from SIDIS

not shown here

- di-hadron asymmetries [PLB 713 (2012) 10, PLB 736 (2014) 124]
- interplay among transversity induced asymmetries
[PLB 753 (2016) 406]

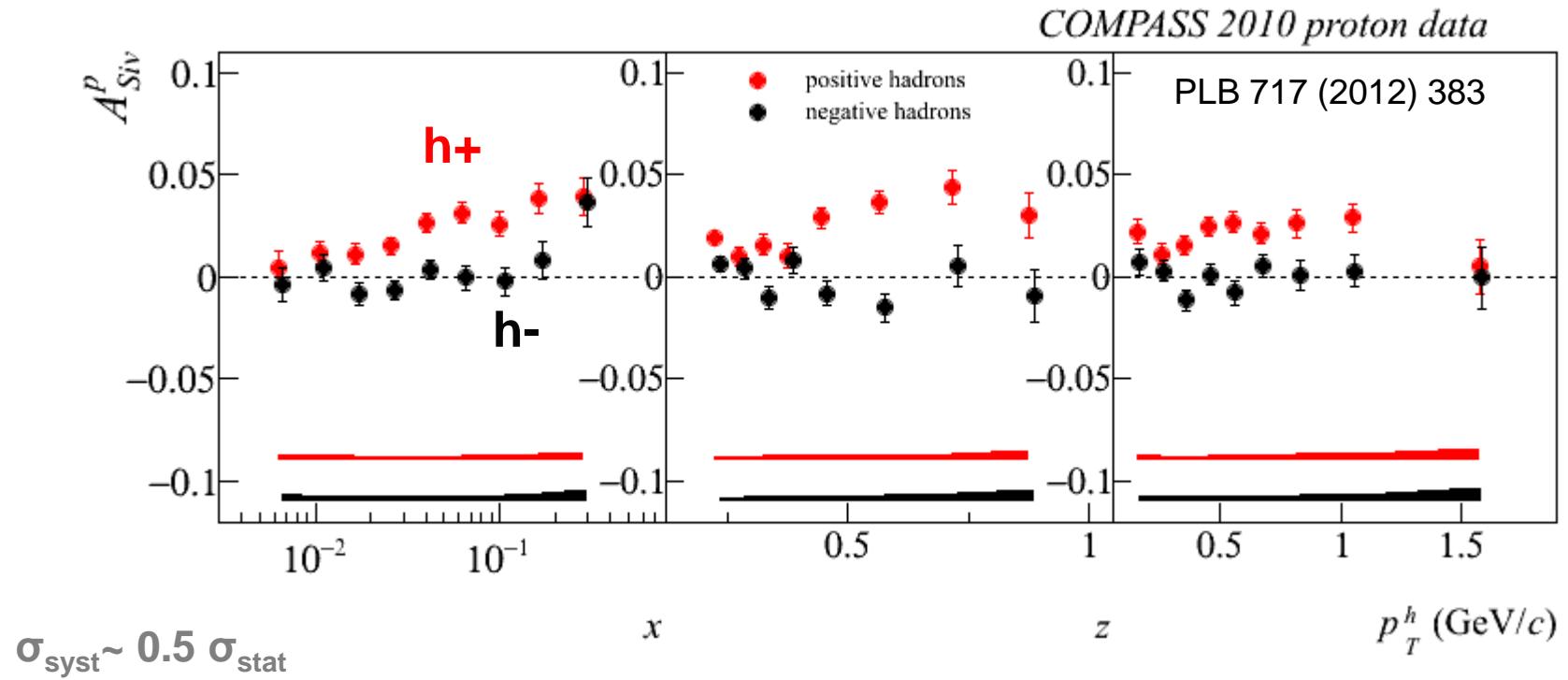


Sivers asymmetries

Sivers asymmetries on proton



charged hadrons
2010 data

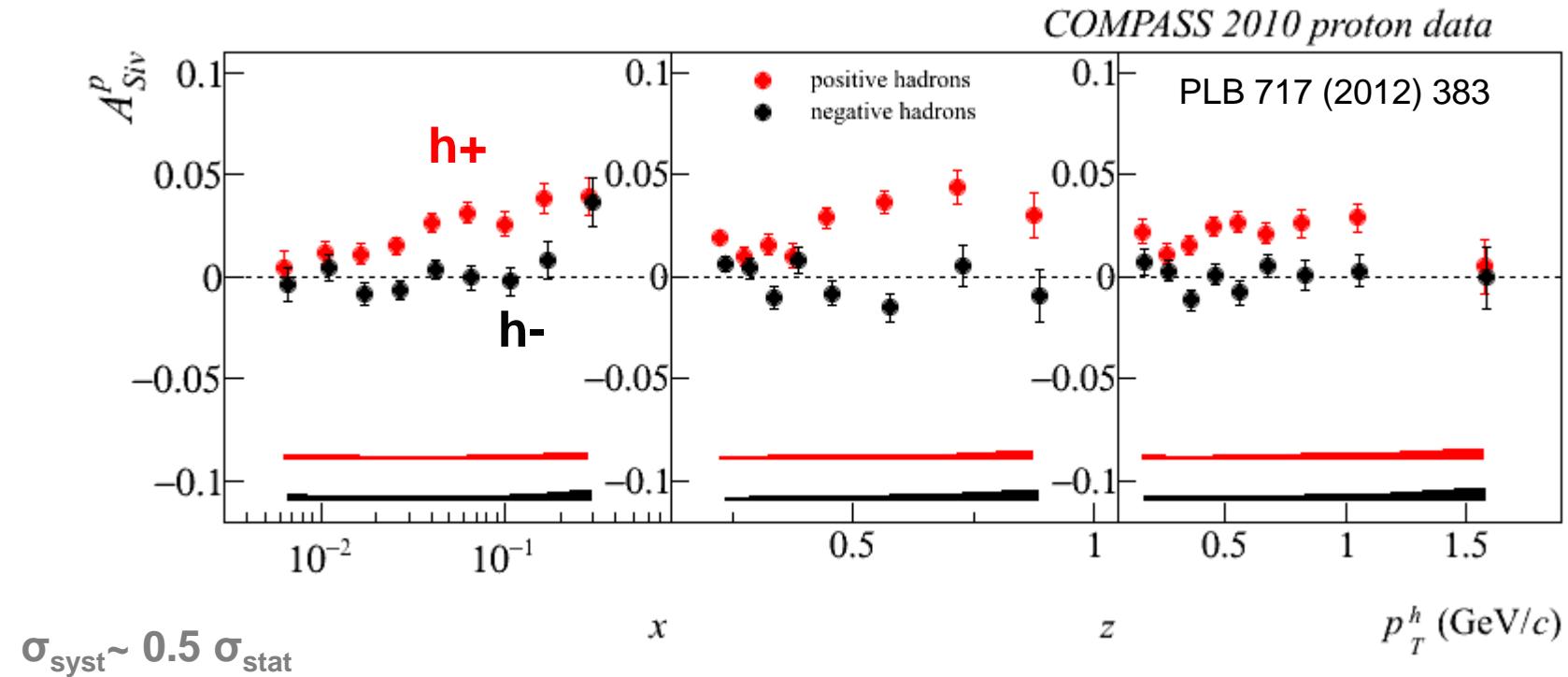


clear evidence for a positive signal for h^+ , which extends to small x

Sivers asymmetries on proton



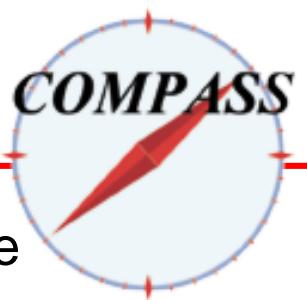
charged hadrons
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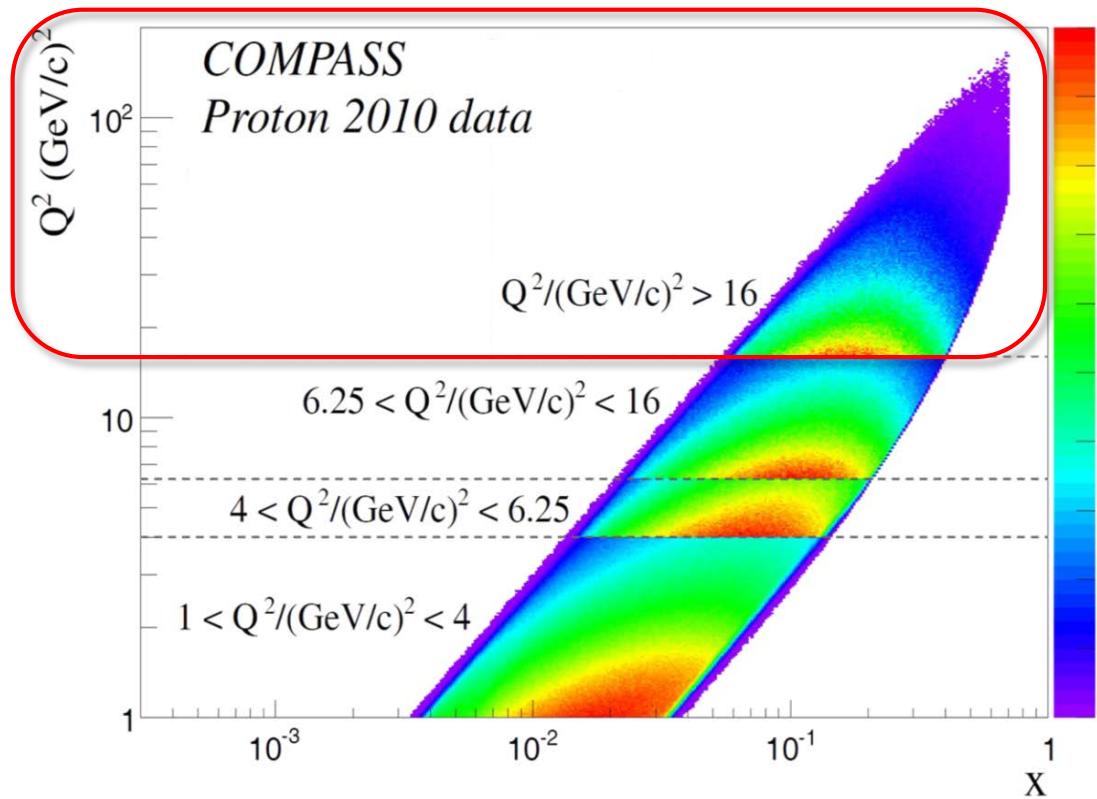
clear evidence for a positive signal for h^+ , which extends to small x

d data compatible with zero but
with large statistical uncertainties

Sivers asymmetries on proton

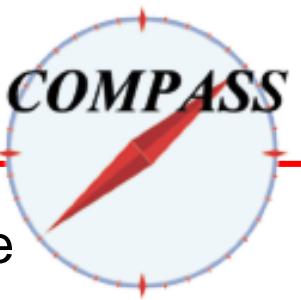


COMPASS has measured the SIDIS TSA in the four Q^2 ranges of the Drell-Yan measurement

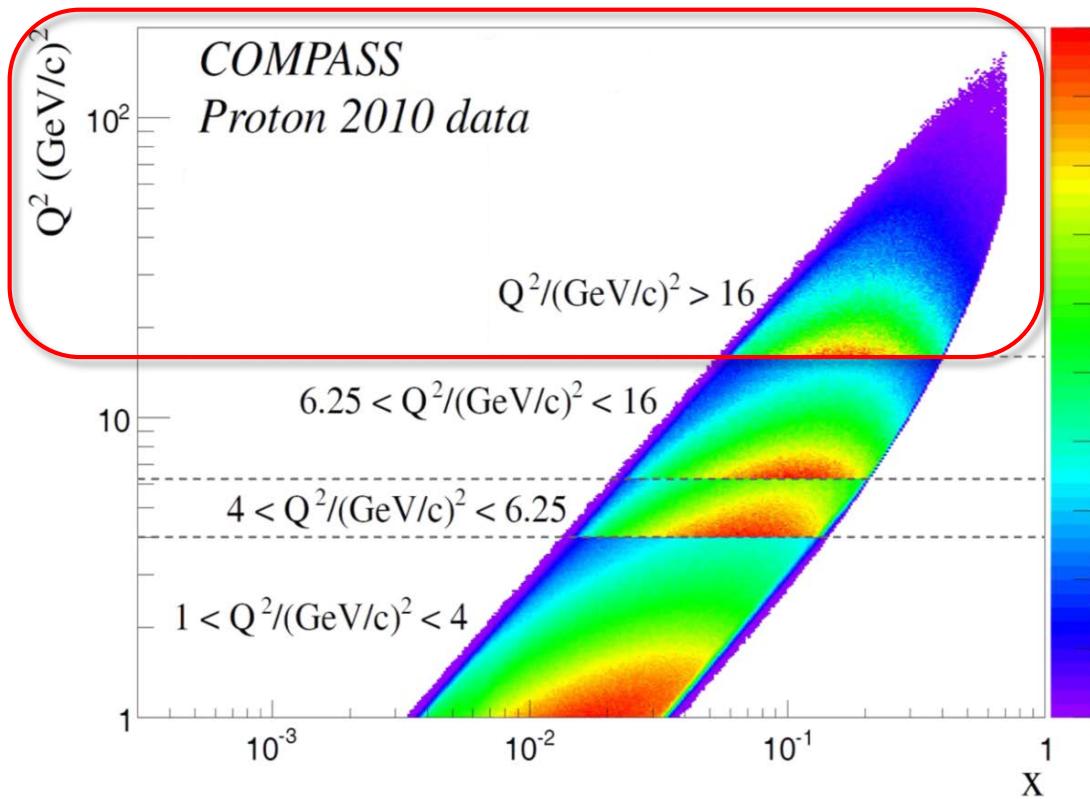


“golden” region for DY: $Q^2 > 16 \text{ GeV}^2$

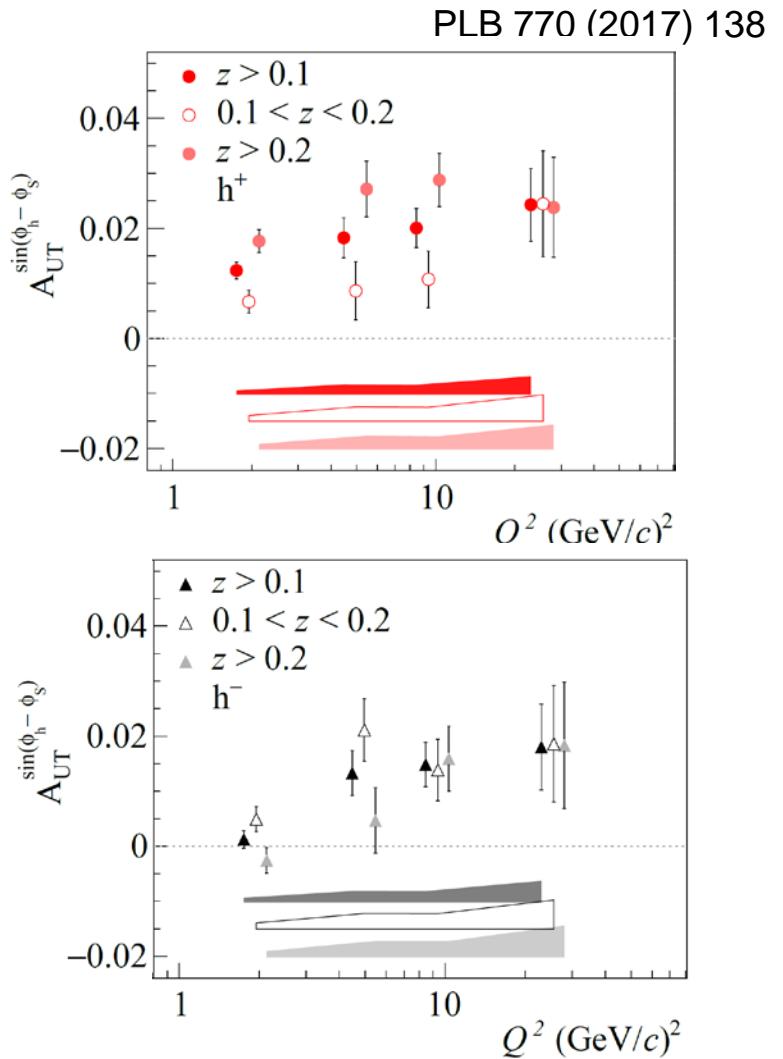
Sivers asymmetries on proton



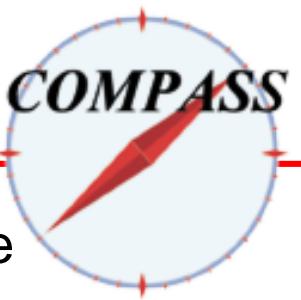
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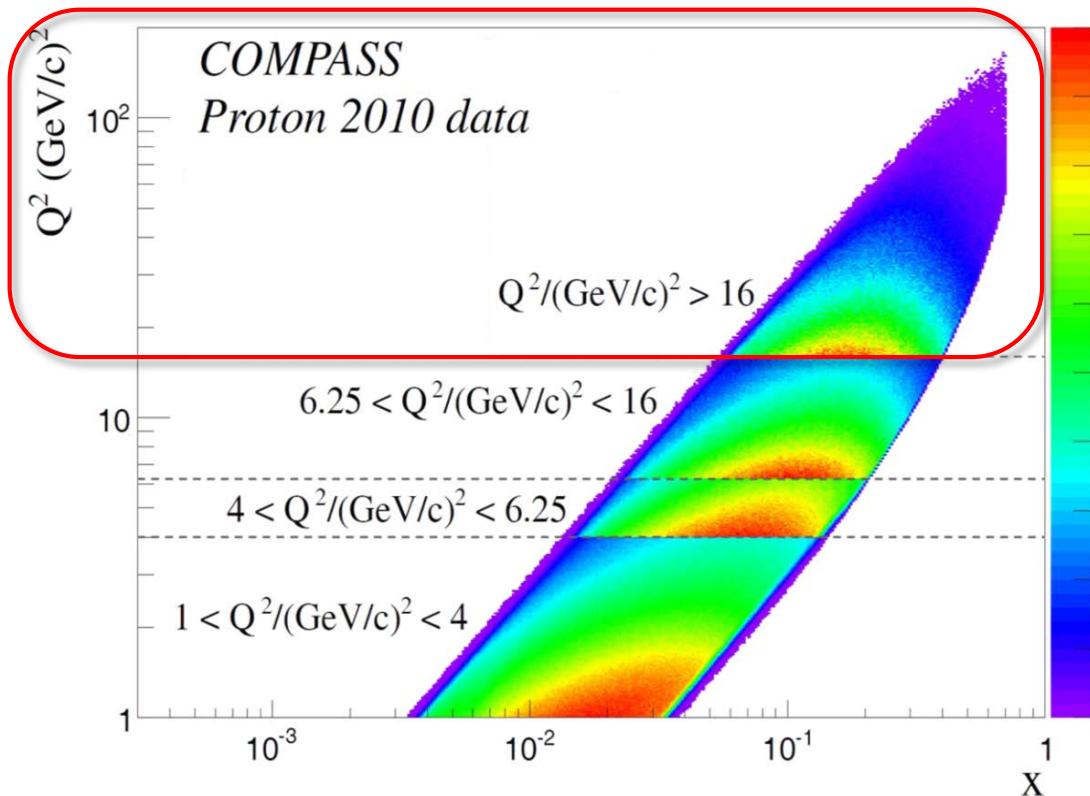
“golden” region for DY: $Q^2 > 16$ GeV 2



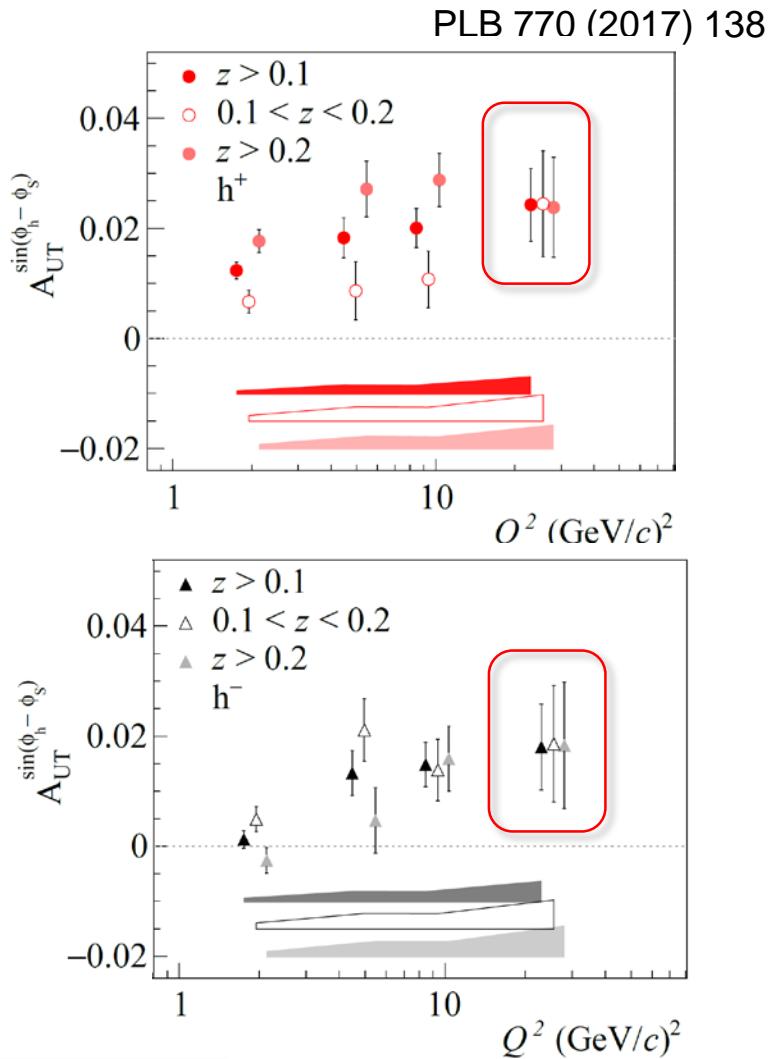
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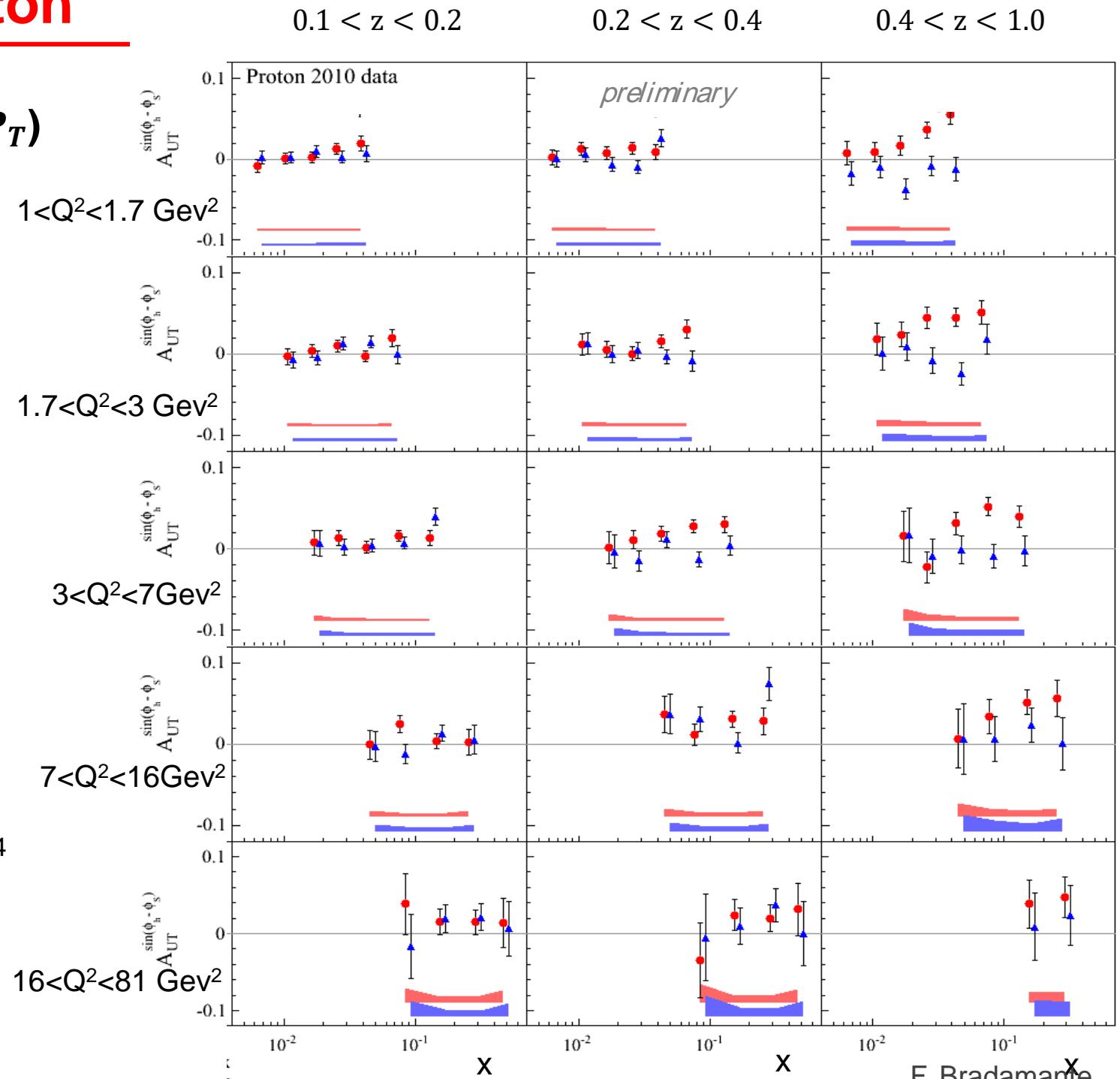
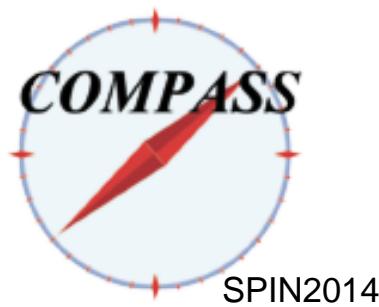
clearly positive
test of change of sign feasible

TSA on proton

multiD ($x, Q^2; z, P_T$) analysis

an example:
Sivers
asymmetry

$P_T > 0.1 \text{ GeV}/c$



Sivers asymmetries



not shown here: gluon Sivers

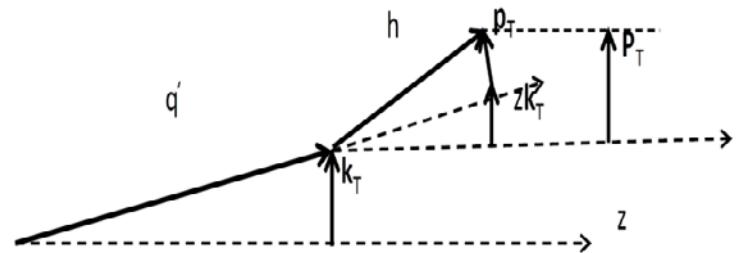
- J/Ψ asymmetry [J. Matousek, DSPIN-15]
- high p_T hadron pair asymmetry [K. Kurek, DSPIN-15, PLB 772 (2017) 854]

**new:
weighted
Sivers asymmetries
in SIDIS**

weighted Sivers asymmetries - why

“standard” Sivers asymmetry

$$A_{Siv}(x, z) = \frac{\sum_q e_q^2 x f_{1T}^{\perp q}(x) \otimes D_{1q}(z)}{\sum_q e_q^2 x f_1^q(x) D_{1q}(z)}$$



to evaluate the convolution, models are needed
e.g. Gaussian model

with several assumptions, one gets

$$A_{Siv,G}(x, z) \cong \frac{\pi M}{2\langle P_T \rangle} \frac{\sum_q e_q^2 x f_{1T}^{\perp(1)q}(x) z D_{1q}(z)}{\sum_q e_q^2 x f_1^q(x) D_{1q}^h(z)}$$

$$f_{1T}^{\perp(1)q}(x) = \int d^2 \vec{k}_T \frac{k_T^2}{2M^2} f_{1T}^{\perp q}(x, k_T^2)$$

weighted Sivers asymmetries

by weighting the spin dependent part of the cross-section with P_T
one can solve the convolution

there are two slightly different possibilities:

$$w = P_T/zM$$

$$A_{Siv}^w(x, z) = 2 \frac{\sum_q e_q^2 x f_{1T}^{\perp(1)q}(x) D_{1q}(z)}{\sum_q e_q^2 x f_1^q(x) D_{1q}(z)}$$

easier interpretation

$$w' = P_T/M$$

$$A_{Siv}^{w'}(x, z) = 2 \frac{\sum_q e_q^2 x f_{1T}^{\perp(1)q}(x) z D_{1q}(z)}{\sum_q e_q^2 x f_1^q(x) D_{1q}(z)}$$

easier comparison with the “gaussian ansatz”

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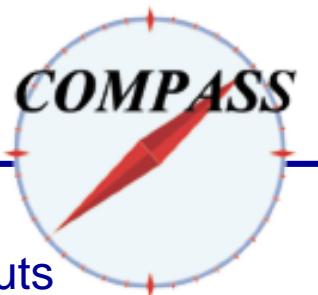
$$A_{Siv}^{w'}(x, z) = 2 \frac{\sum_q e_q^2 x f_{1T}^{\perp(1)q}(x) z D_{1q}(z)}{\sum_q e_q^2 x f_1^q(x) D_{1q}(z)}$$

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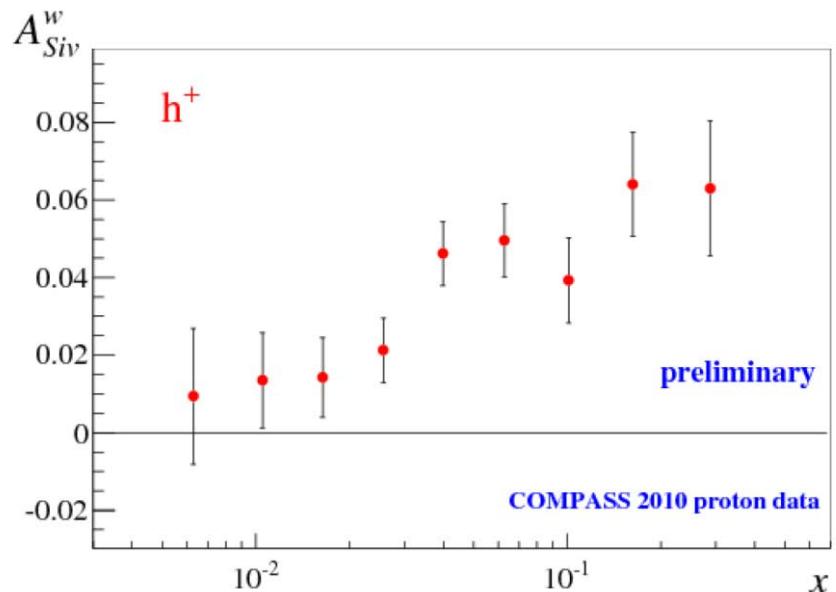
COMPASS has measured both, in bins of x and z

P_T/zM weighted Sivers asymmetry

$$A_{Siv}^w(x) = 2 \frac{\sum_q e_q^2 x f_{1T}^{\perp(1)q}(x) \int D_{1q}(z) dz}{\sum_q e_q^2 x f_1^q(x) \int D_{1q}(z) dz} \quad w = P_T/zM$$

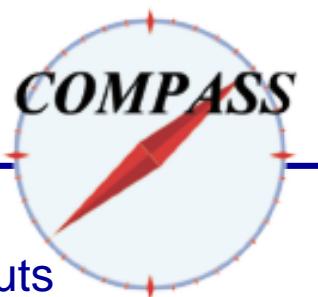


standard cuts
 $z > 0.2$
arXiv:1702.00621



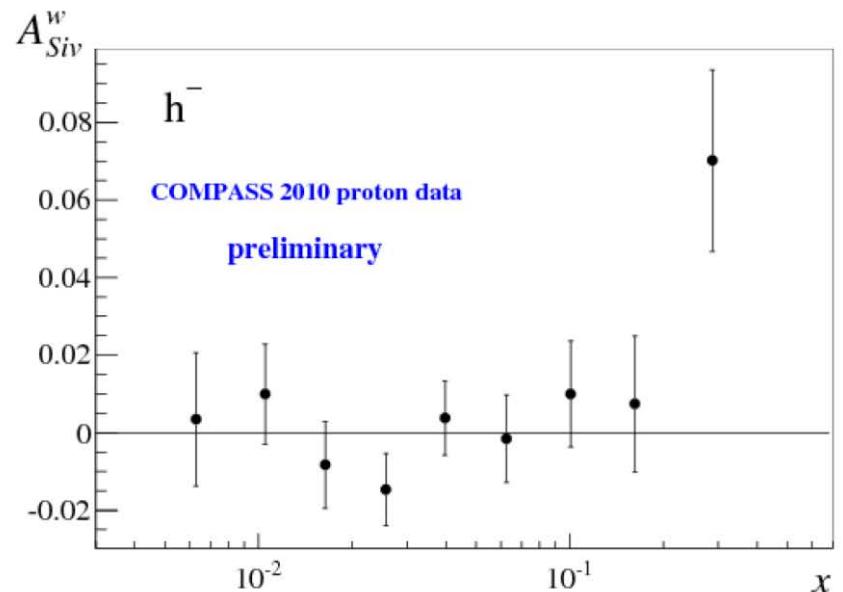
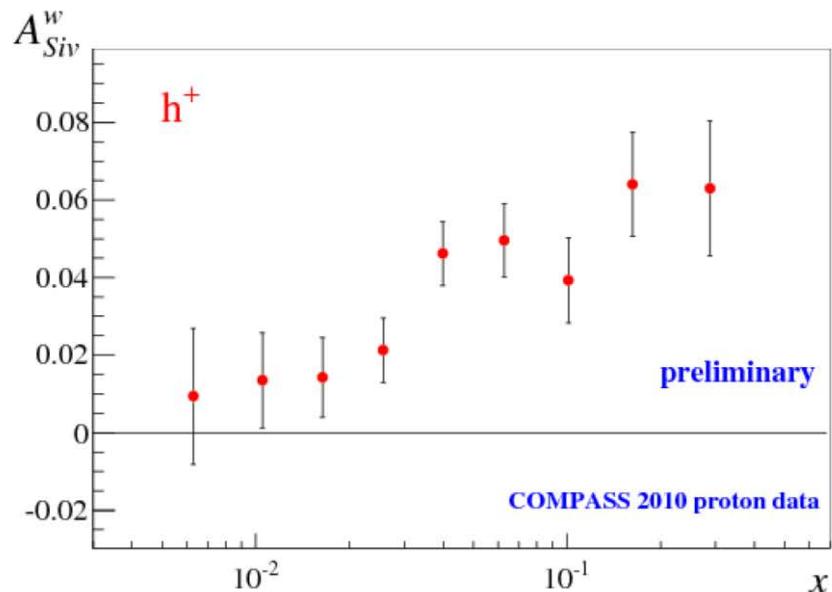
$$\sim 2 \frac{f_{1T}^{\perp(1)u}(x)}{f_1^u(x)}$$

P_T/zM weighted Sivers asymmetry



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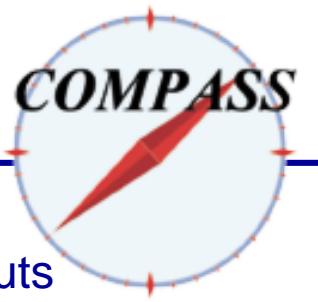
$$\sim 2 \frac{f_{1T}^{\perp(1)u}(x)}{f_1^u(x)}$$

both $f_{1T}^{\perp(1)u}$ and $f_{1T}^{\perp(1)d}$ contribute

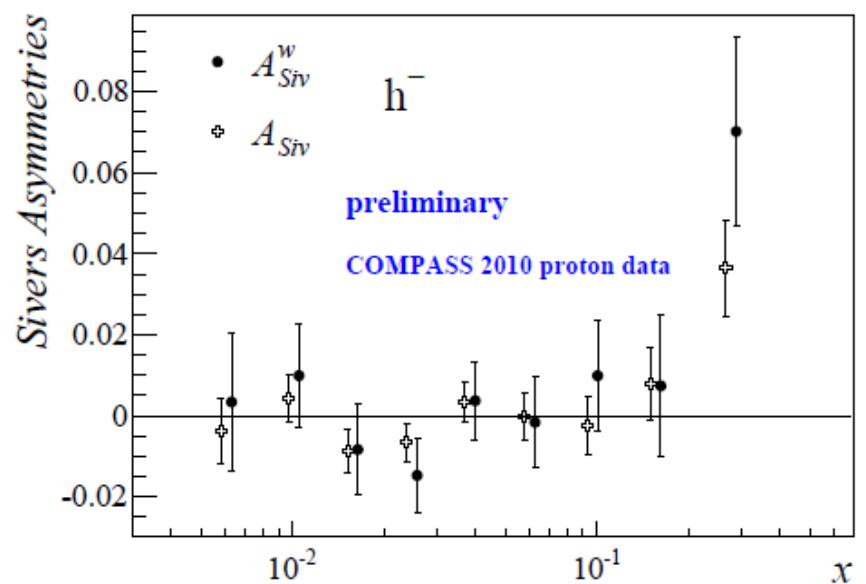
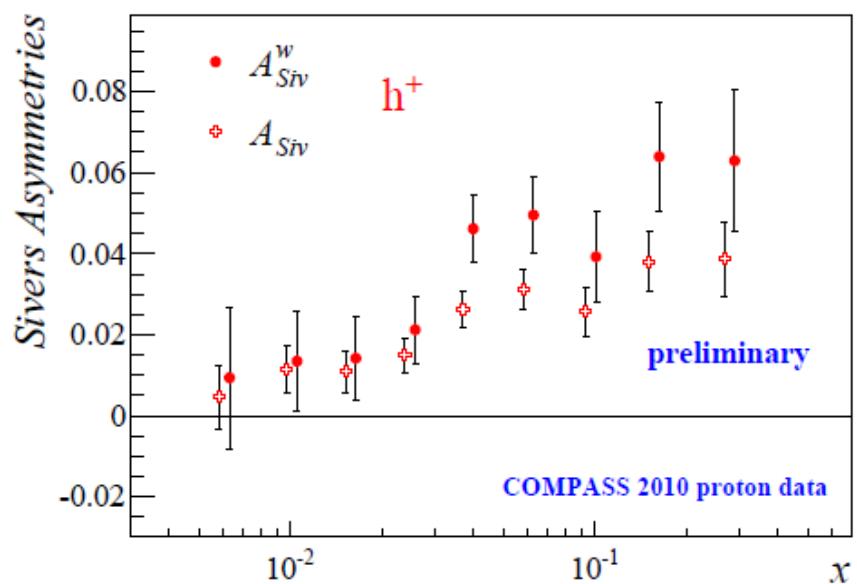
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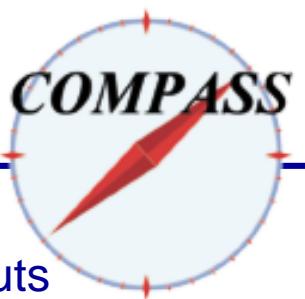


comparison with the published “standard” asymmetries



open points: standard asymmetries, PLB 717 (2012) 383

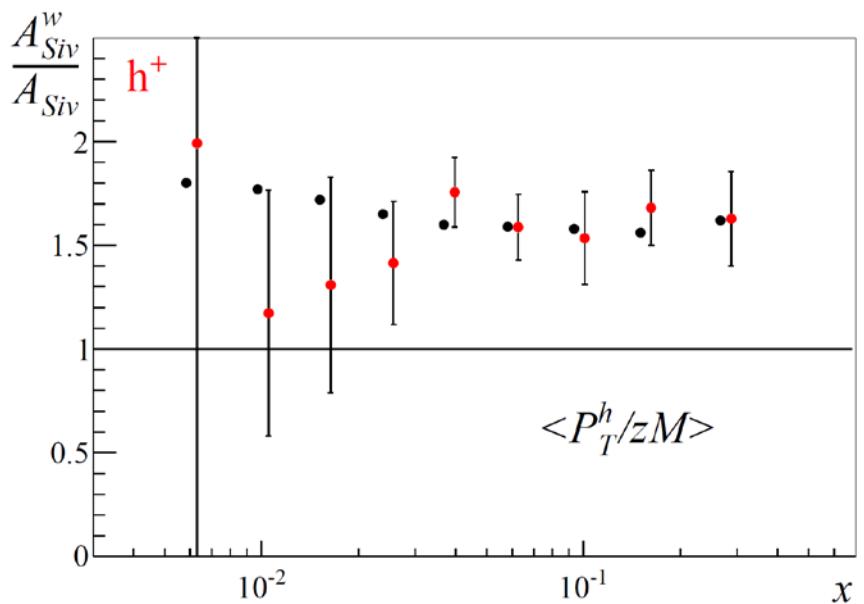
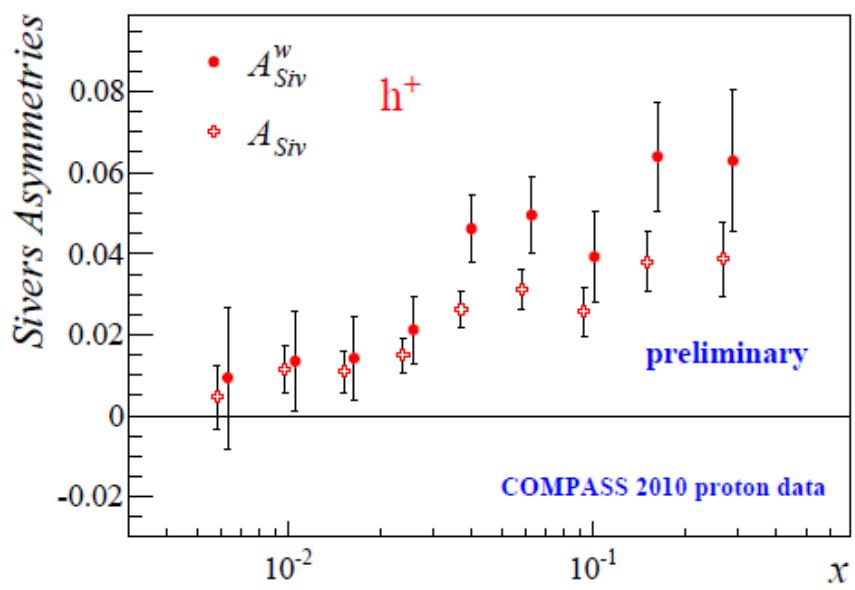
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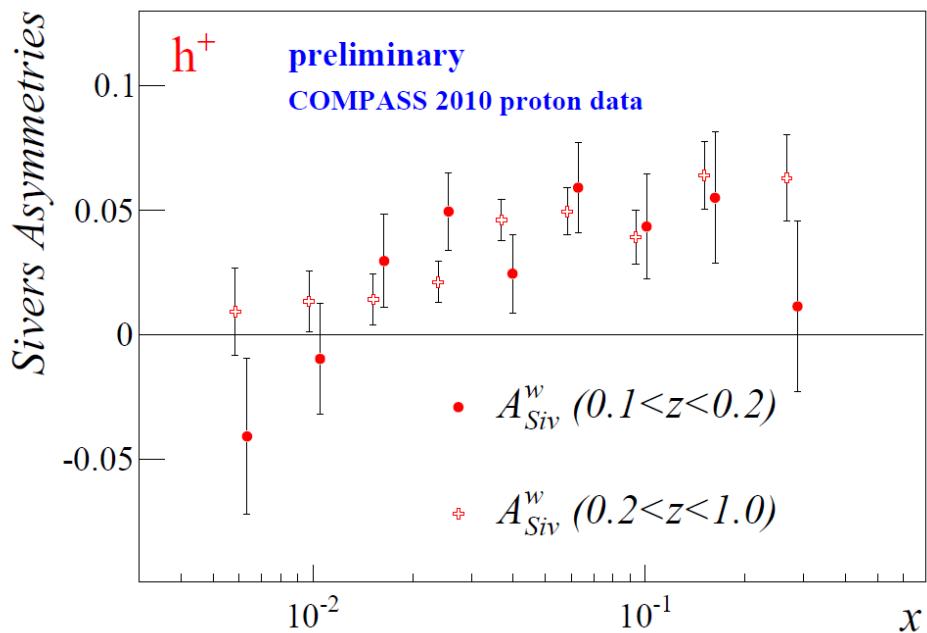
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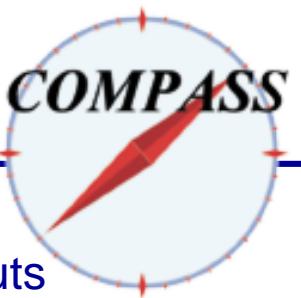
different z ranges



$$\sim 2 \frac{f_{1T}^{\perp(1)u}(x)}{f_1^u(x)}$$

ok

P_T/zM weighted Sivers asymmetry

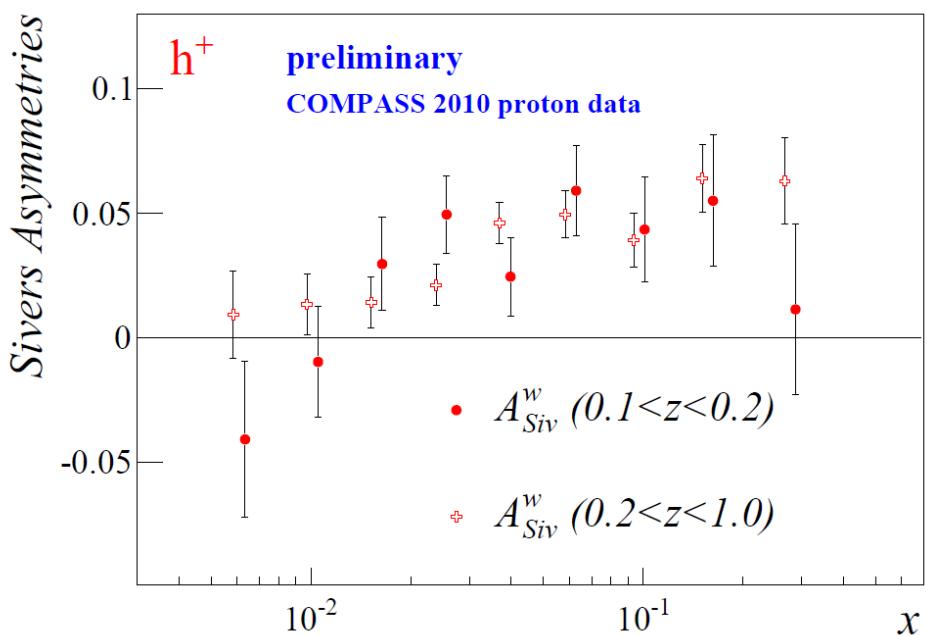


standard cuts

arXiv:1702.00621

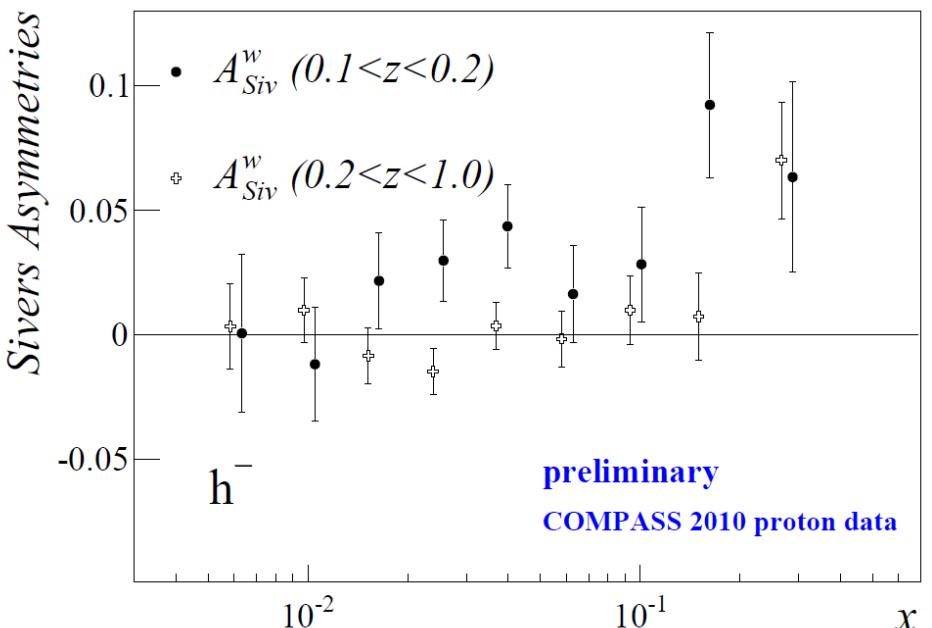
$$A_{Siv}^w(x) = 2 \frac{\sum_q e_q^2 x f_{1T}^{\perp(1)q}(x) \int D_{1q}(z) dz}{\sum_q e_q^2 x f_1^q(x) \int D_{1q}(z) dz} \quad w = P_T/zM$$

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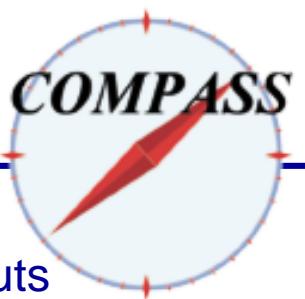
ok



both $f_{1T}^{\perp(1)u}$ and $f_{1T}^{\perp(1)d}$ contribute
the relative contribution depends on z

$0.1 < z < 0.2 \quad h^+ \sim h^-$

P_T/zM weighted Sivers asymmetry

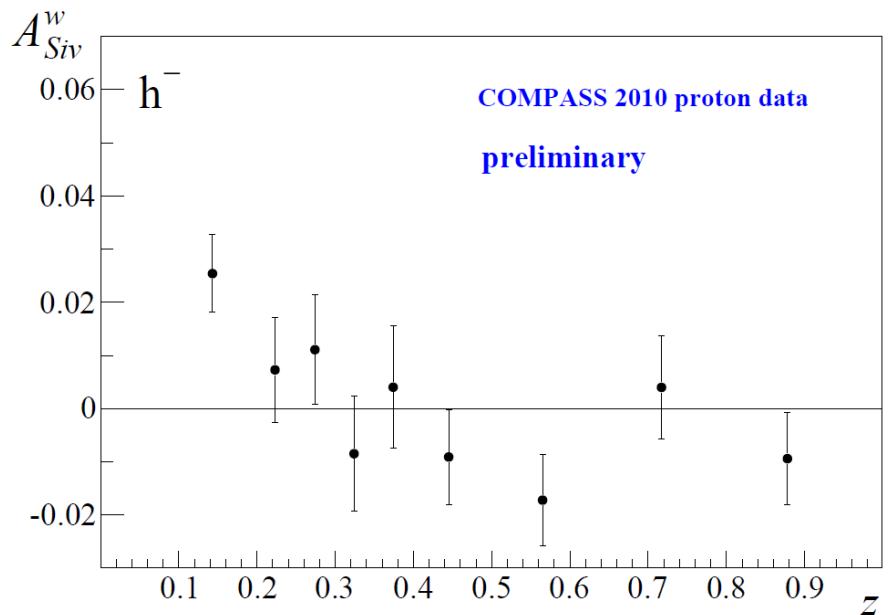
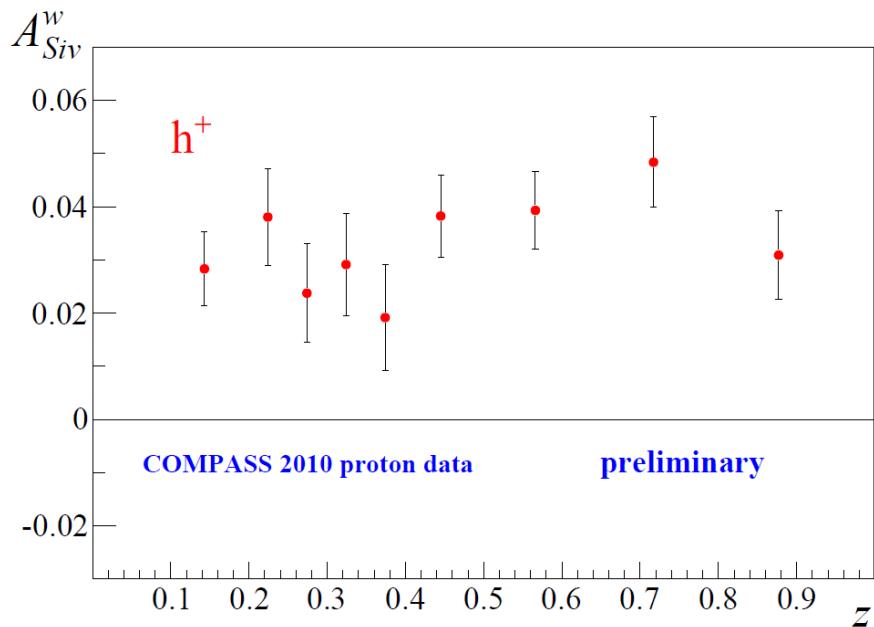


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$$A_{Siv}^w(x) = 2 \frac{\sum_q e_q^2 x f_{1T}^{\perp(1)q}(x) \int D_{1q}(z) dz}{\sum_q e_q^2 x f_1^q(x) \int D_{1q}(z) dz} \quad w = P_T/zM$$

z dependence

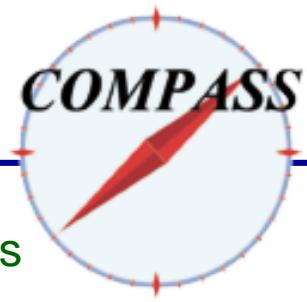


$$\sim 2 \frac{\int C(x) f_{1T}^{\perp(1)q}(x) dx}{\int C(x) f_1^q(x) dx}$$

ok

both $f_{1T}^{\perp(1)u}$ and $f_{1T}^{\perp(1)d}$ contribute
the relative contribution depends on z

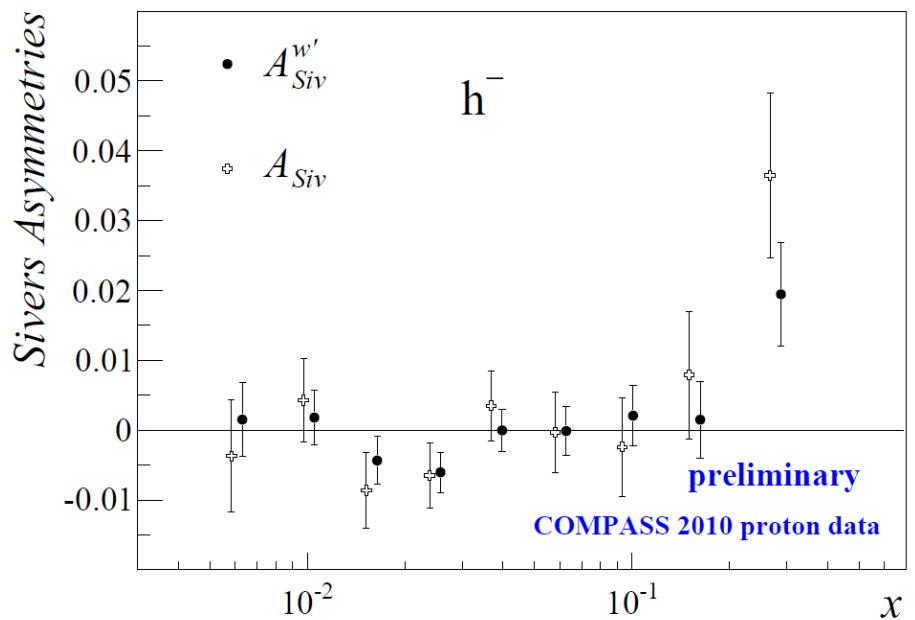
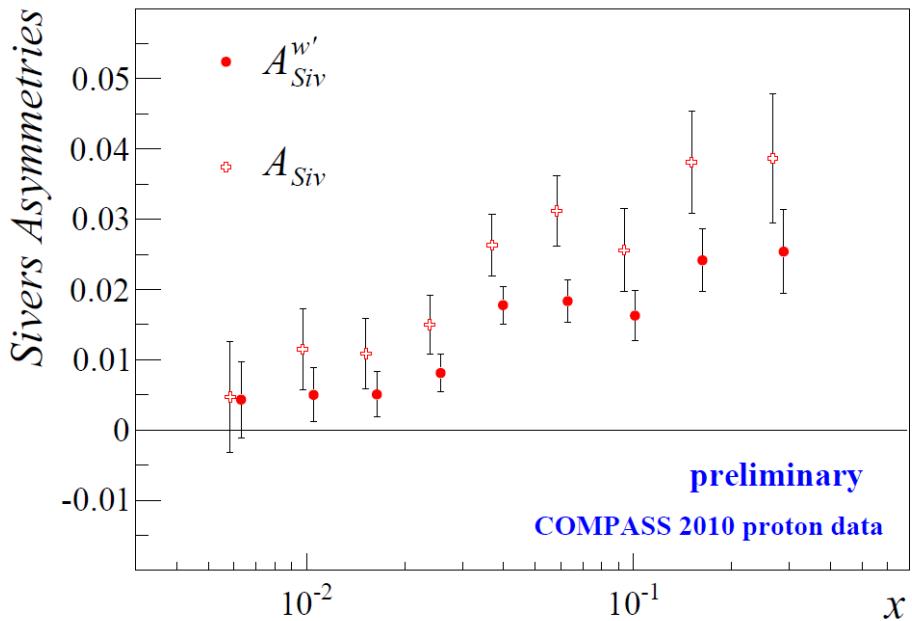
P_T/M weighted Sivers asymmetry



$$A_{Siv}^{w'}(x) = 2 \frac{\sum_q e_q^2 x f_{1T}^{\perp(1)q}(x) \int z D_{1q}(z) dz}{\sum_q e_q^2 x f_1^q(x) \int D_{1q}(z) dz} \quad w' = P_T/M$$

standard cuts
 $z > 0.2$

comparison with the published “standard” asymmetries’



P_T/M weighted Sivers asymmetry

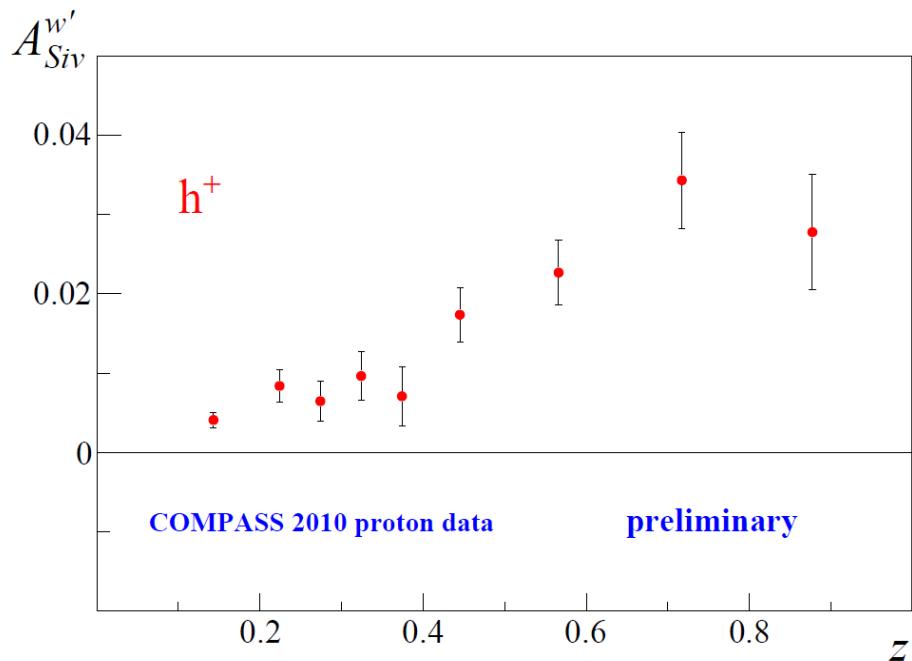


$$A_{Siv}^{w'}(z) = 2 \frac{\sum_q e_q^2 \int C(x) f_{1T}^{\perp(1)q}(x) dx \ z D_{1q}(z)}{\sum_q e_q^2 \int C(x) f_1^q(x) dx \ D_{1q}(z)}$$

standard cuts

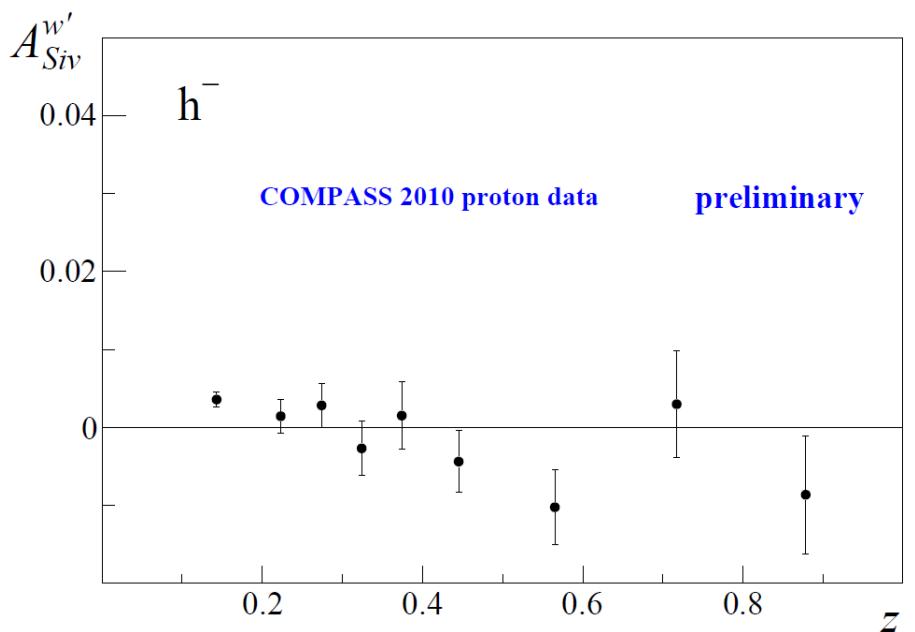
$$w' = P_T/M$$

z dependence



COMPASS 2010 proton data

preliminary



COMPASS 2010 proton data

preliminary

$$\sim 2z \frac{\int C(x) f_{1T}^{\perp(1)q}(x) dx}{\int C(x) f_1^q(x) dx}$$

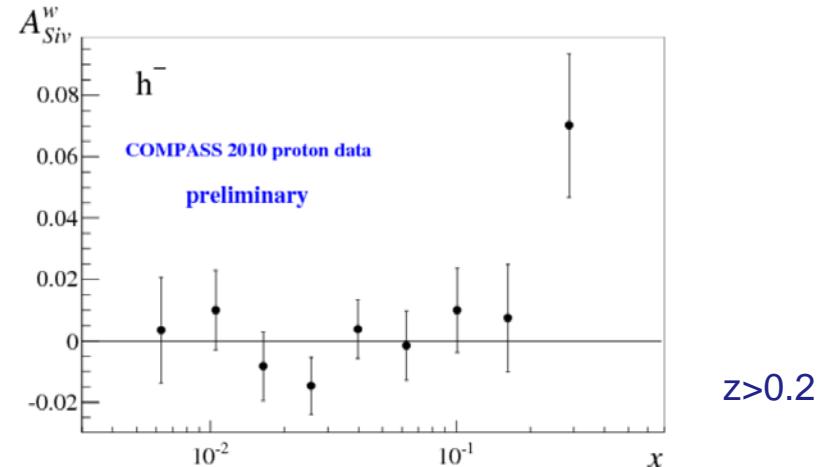
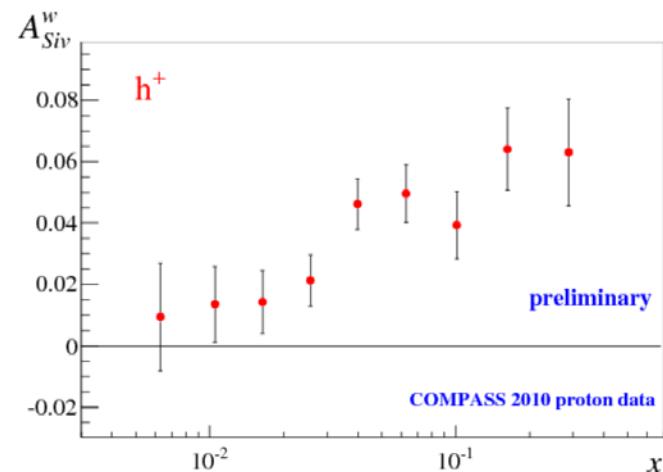
ok

1st moment of the Sivers functions

from P_T/zM weighted Sivers asymmetry



$$A_{Siv}^w(x) = 2 \frac{\sum_q e_q^2 x f_{1T}^{\perp(1)q}(x) \int D_{1q}(z) dz}{\sum_q e_q^2 x f_1^q(x) \int D_{1q}(z) dz}$$



used to extract the u- and d-quark Sivers functions

(J. Matousek talk)

assumptions:

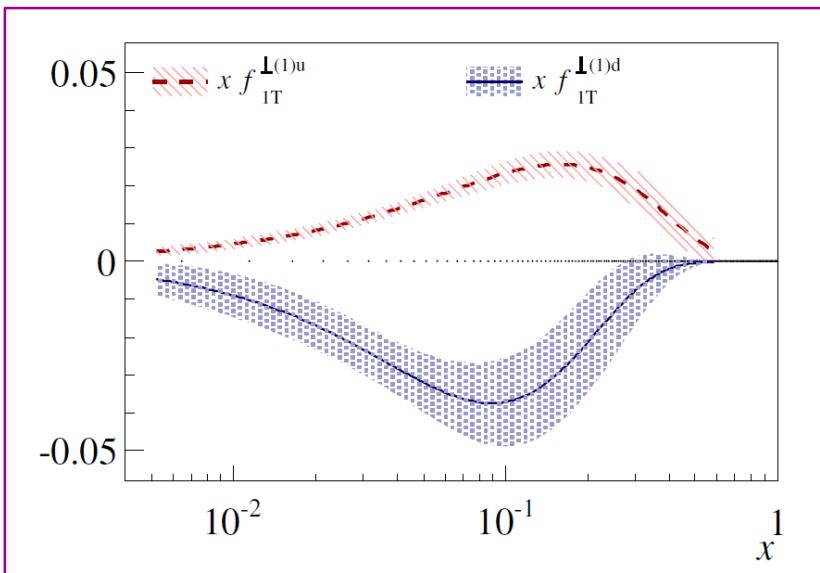
- the Sivers functions of all the sea quarks are zero

$$A_{Siv}^w(x) = 2 \frac{x f_{1T}^{\perp(1)u_v} \int D_{1u} dz + x f_{1T}^{\perp(1)d_v} \int D_{1d} dz}{\sum_q e_q^2 x f_1^q(x) \int D_{1q}(z) dz}$$

- $f_{1T}^{\perp(1)q} = a_q x^{b_q} (1-x)^{c_q}$ with $q = u_v, d_v$

1st moment of the Sivers functions

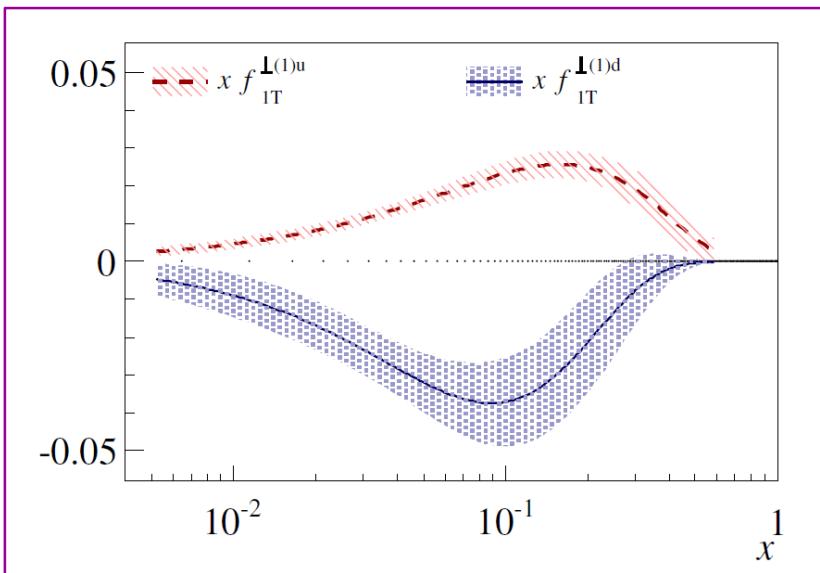
from P_T/zM weighted Sivers asymmetry



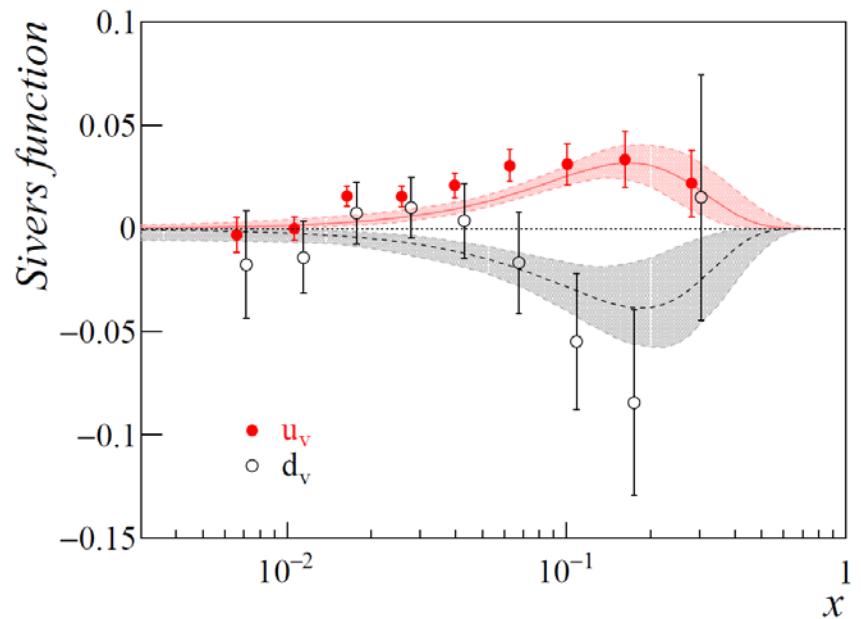
bands: 1σ statistical errors only

1st moment of the Sivers functions

from P_T/zM weighted Sivers asymmetry



bands: 1σ statistical errors only



point by point extraction
COMPASS standard
Sivers asymmetries
p and d, pion and K
M B B 2017

curves: fit of COMPASS
and HERMES data,
Anselmino et al 2012

P_T weighted Sivers asymmetry



to summarise:

- the results for the asymmetries for p look very interesting
in particular
at first order no deviation from naïve expectation based on factorisation
- a first simple attempt to extract the 1st moment of the Siver functions gives quite reasonable results
- new d data needed also in this case to complete the exploratory COMPASS program

**measurement of the
transversity transmitted
 Λ polarisation**

NEW!

Λ polarisation

the Λ polarisation can be measured from the angular distribution of the proton produced in the decay $\Lambda \rightarrow p \pi^-$

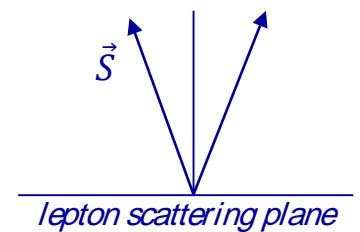
in the Λ c.m.s., the proton angular distribution is $\frac{dN}{d \cos\theta} \propto 1 + \alpha P_\Lambda \cos\theta$

where $\alpha = 0.642 \pm 0.013$ and

θ is the angle between the proton direction and the direction of the Λ polarisation

in SIDIS off transversely polarised nucleons, the Λ polarisation measured using the “reflected” direction of the nucleon spin can be written as

$$P_\Lambda = \frac{\sum_q e_q^2 h_1^q H_1^{\Lambda/q}}{\sum_q e_q^2 f_1^q D_1^{\Lambda/q}}$$



“transversity transmitted Λ polarisation”

the “transverse polarisation” is measured using as axis the normal to the lepton scattering plane

Λ polarisation

$$P_\Lambda = \frac{\sum_q e_q^2 h_1^q H_1^{\Lambda/q}}{\sum_q e_q^2 f_1^q D_1^{\Lambda/q}}$$

today transversity is somewhat known, while $H_1^{\Lambda/q}$ is completely unknown

with different assumptions,

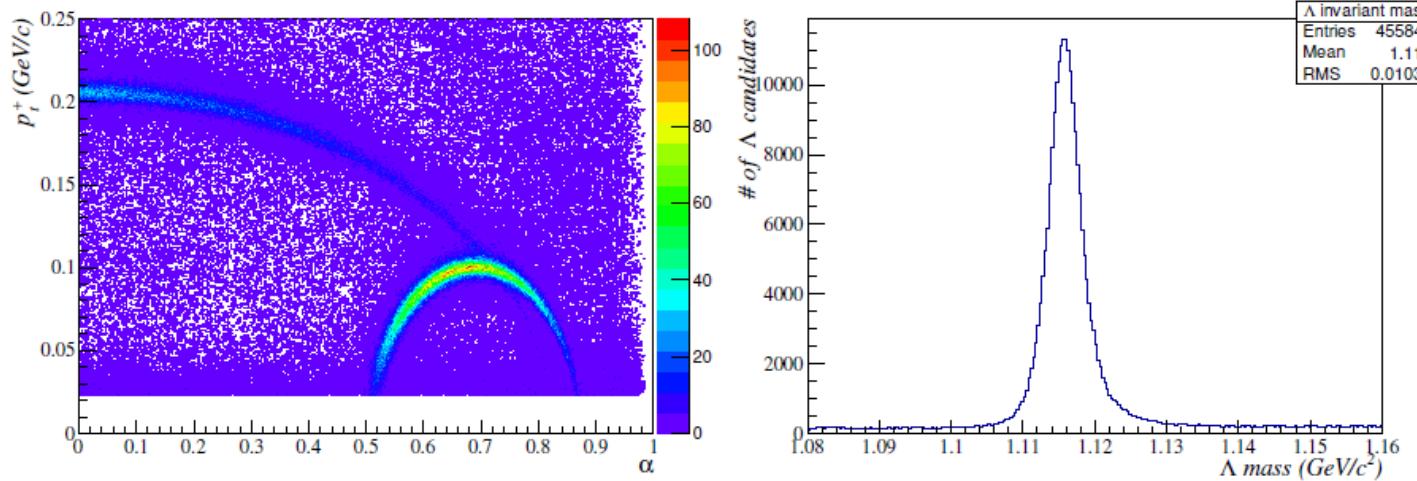
this measurement can give information either on h_1^s or on $H_1^{\Lambda/q}/D_1^{\Lambda/q}$

COMPASS has measured the Λ and $\bar{\Lambda}$ polarisation
from the complete proton data set (2007 and 2010)

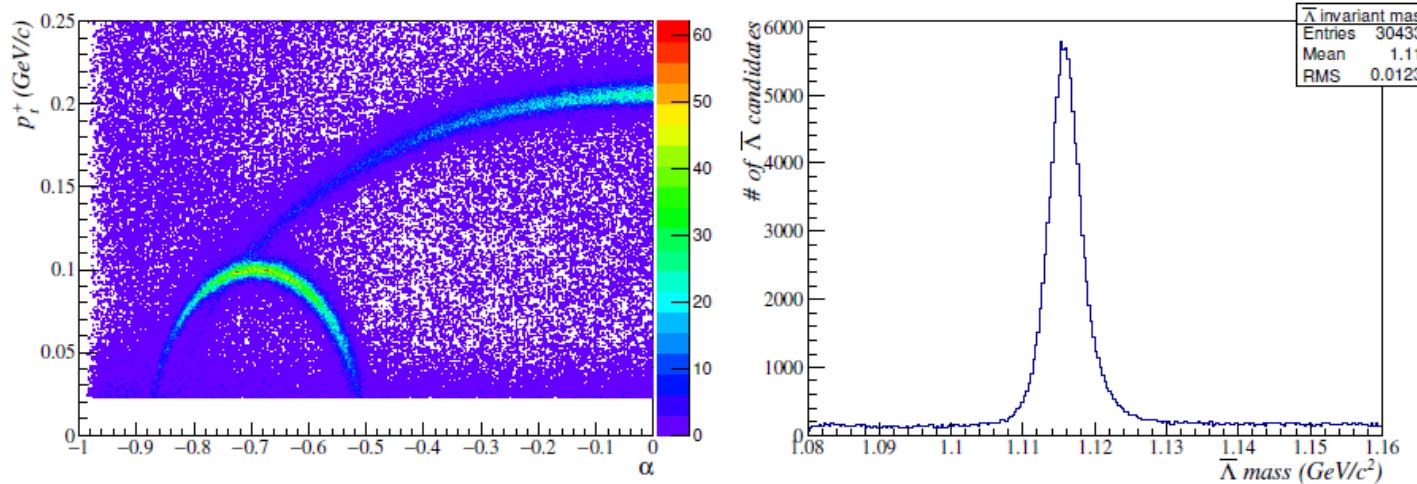
$\Lambda/\bar{\Lambda}$ polarisation



$Q^2 > 1 \text{ (GeV/c)}^2$



$\sim 300k$



$\sim 150k$

$\Lambda/\bar{\Lambda}$ polarisation



the polarisation has been measured in x , z and P_T bins
from the complete p data set (2007 and 2010)

for all Λ and $\bar{\Lambda}$ candidates and

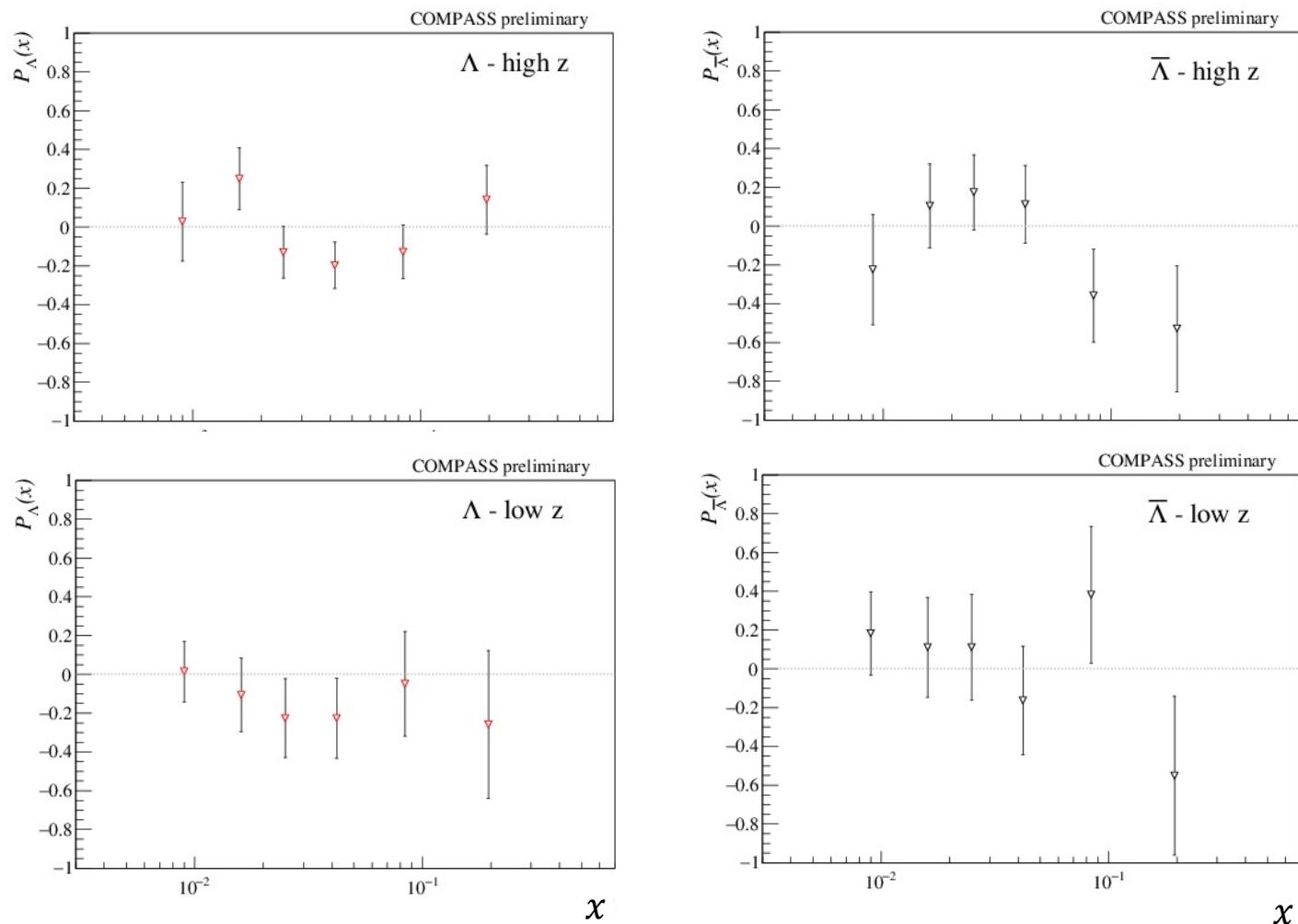
for

- $z > 0.2$ and $x_F > 0$ high z region
- $z > 0.2$ or $x_F > 0$ low z region

- $x > 0.032$ high x region
- $x < 0.032$ low x region

- $P_T > 0.5$ GeV/c high P_T region
- $P_T < 0.5$ GeV/c low P_T region

$\Lambda/\bar{\Lambda}$ polarisation



statistically limited
still the only existing measurement
interpretation work ongoing in COMPASS

COMPASS
Common Muon and Proton Apparatus for
Structure and Spectroscopy

Long-Term plans

A large-scale photograph of a particle accelerator facility, likely the COMPASS experiment at CERN. The image shows complex yellow steel trusses, various pipes, and equipment. In the foreground, there's a large cylindrical detector component. The background features several tall, light-colored vertical panels. The number "6" is visible on one of the panels.

1. COMPASS QCD facility
2. Beyond 2020 Workshop (March 2016)

3. Long term plans

- RF separated beam
- Spectroscopy
- Drell-Yan
- Exclusive measurements with muon and hadron beams

4. Shorter term plans

- SIDIS
- Drell-Yan
- Astrophysics

5. Summary

→ LoI



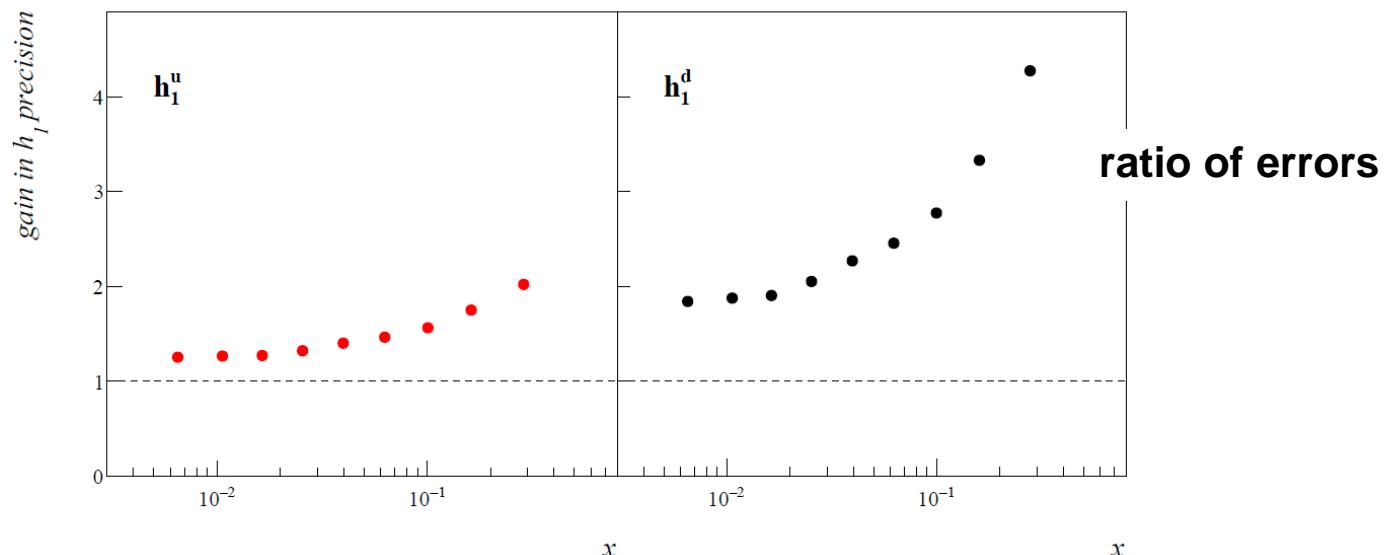
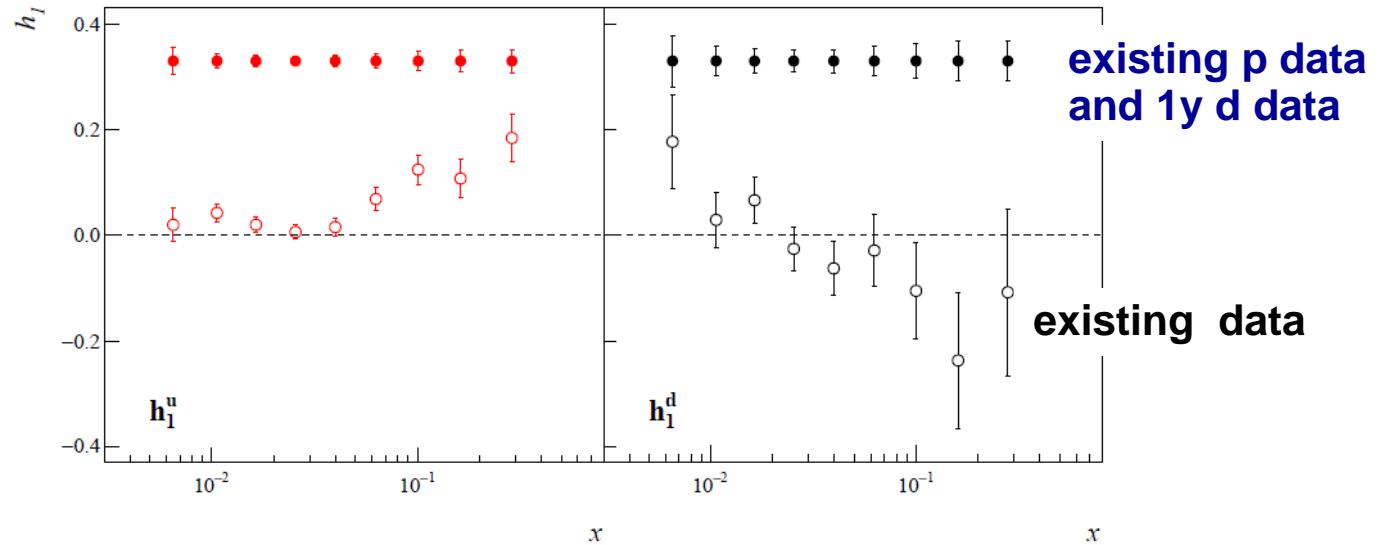
Oleg Denisov INFN(Torino)/CERN for the COMPASS Coll.

latest developments: a proposal for

1. a one year measurement of SIDIS on transversely polarised deuteron



u and d quark
Transversity PDFs



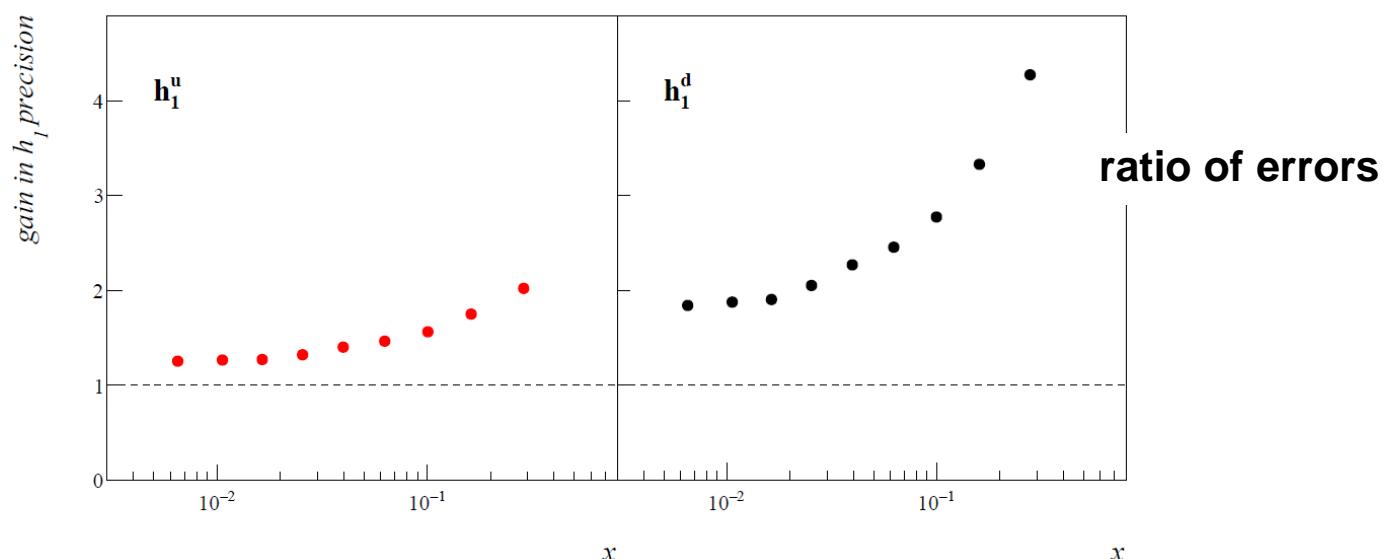
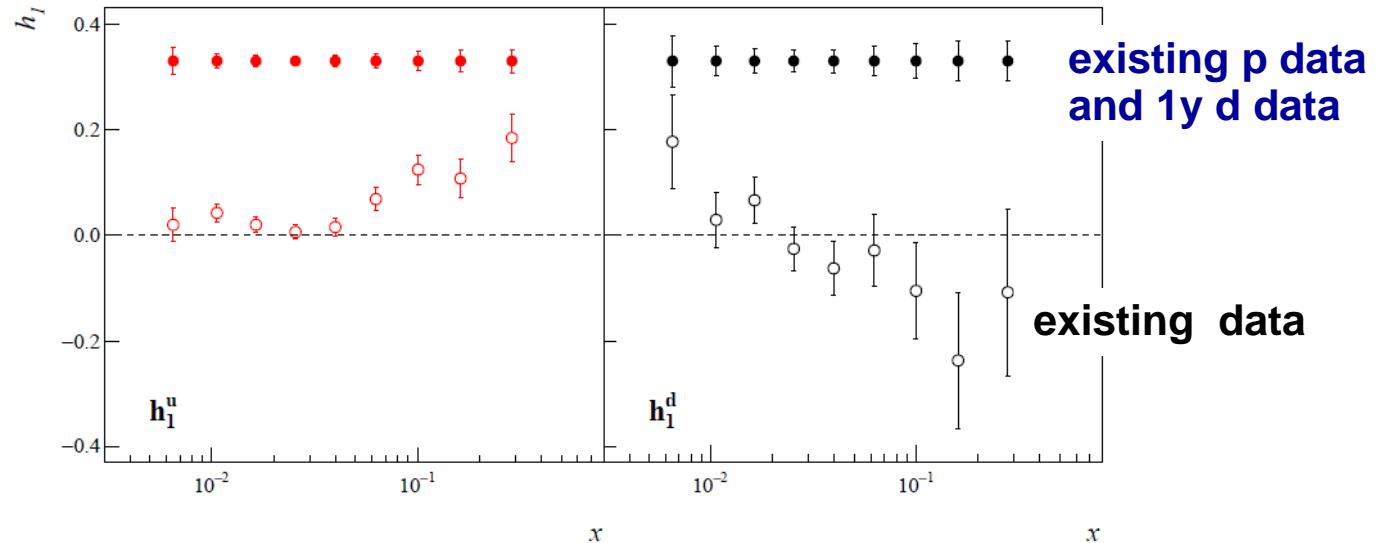
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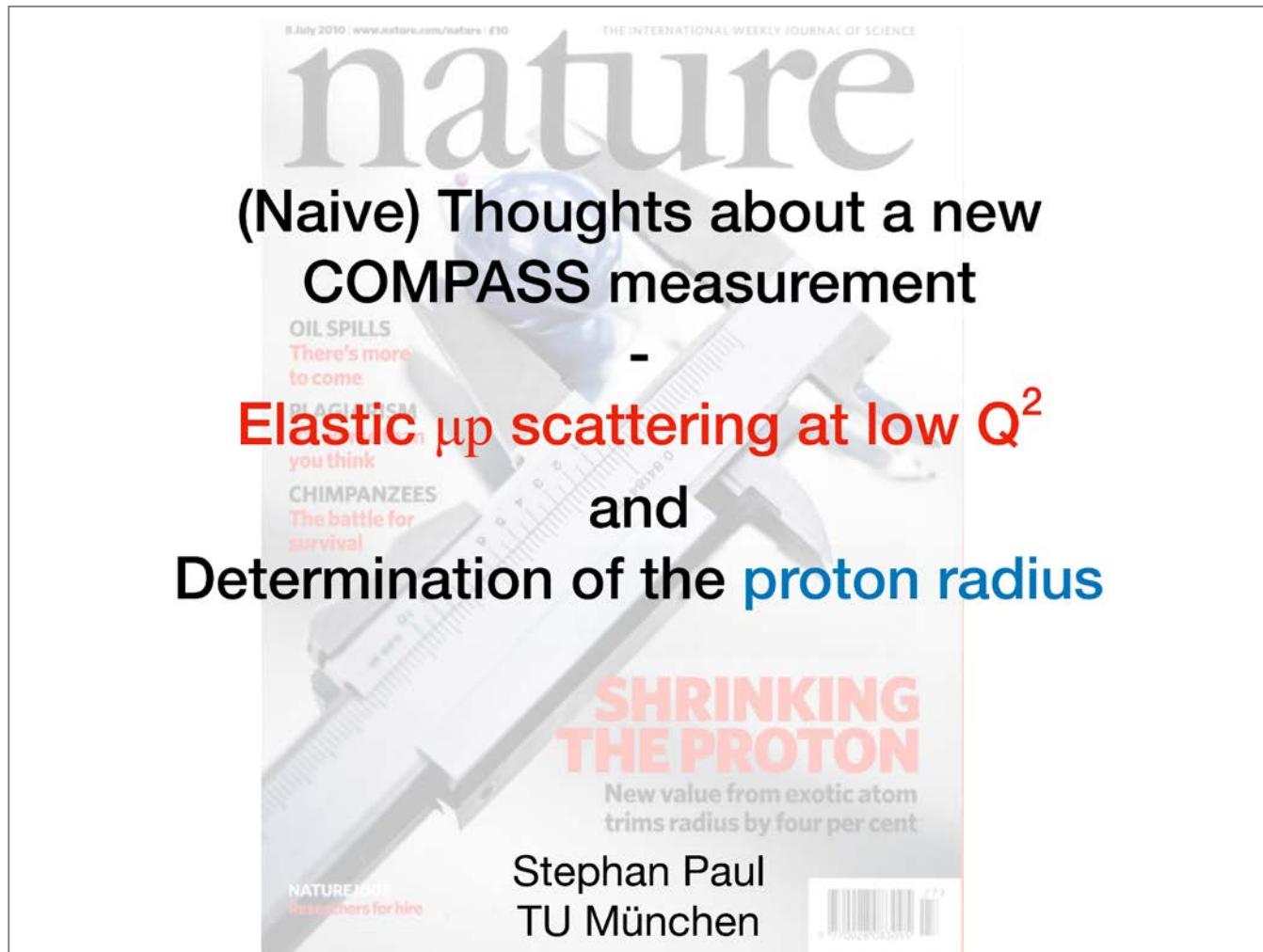


u and d quark
Transversity PDFs

needed to complete
the COMPASS
transverse
spin program



latest developments: a proposal for 2. a measurement of μp elastic scattering



in parallel with an LoI for the longer term projects