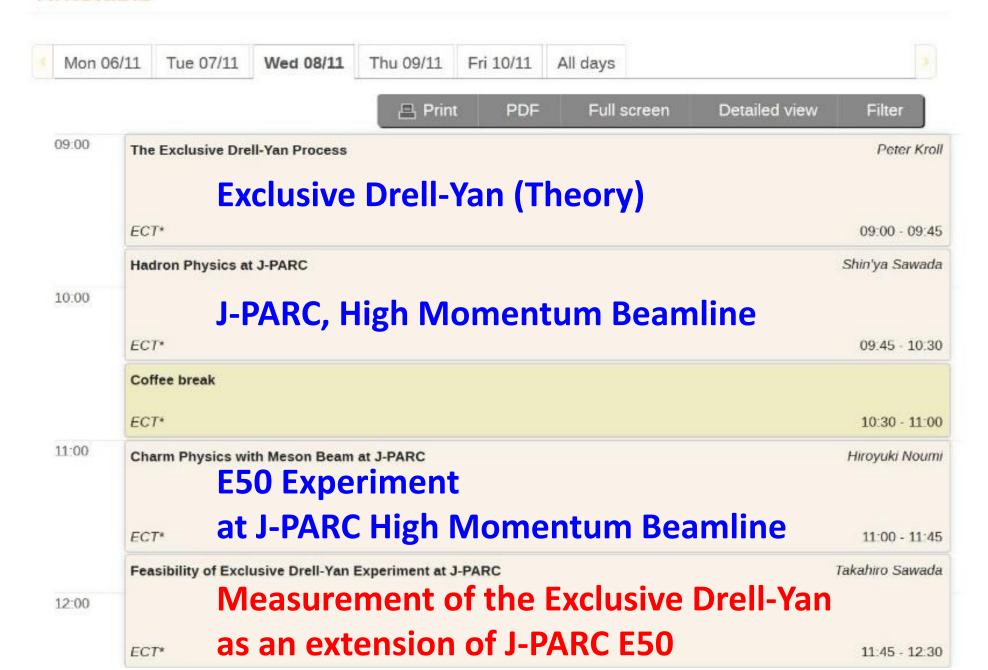
Feasibility of Exclusive Drell-Yan Experiment at J-PARC

Takahiro Sawada University of Michigan

Based on the paper of
T. Sawada, W. C. Chang, Kumano, J. C. Peng, S. Sawad, K. Tanaka,
Phys. Rev. D93 (2016) 114034, arXiv:1605.00364

Timetable



Outline

- Extraction of PDFs and GPDs via Space-like and Time-like processes
- Exclusive pion-induced Drell-Yan process ($\pi^- p \rightarrow \mu^+ \mu^- n$)
- Related Issues:

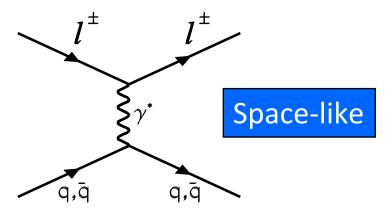
Nucleon GPD, Pion DA, Transition DA, Transition GPD, Pion-pole, Soft nonfactorizable mechanism

- High momentum beamline at J-PARC and E50 experiment
- Feasibility study of exclusive pion-induced Drell-Yan experiment at J-PARC E50
- Summary

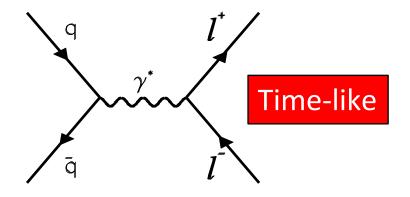
Extraction of **PDFs**

Complementarity between **Space-like** and **Time-like** processes

Deep Inelastic Scattering (DIS)



Drell-Yan (DY) process



Main Processes in Global PDF Analysis

Eur. Phys. J. C (2009) 63: 189–285

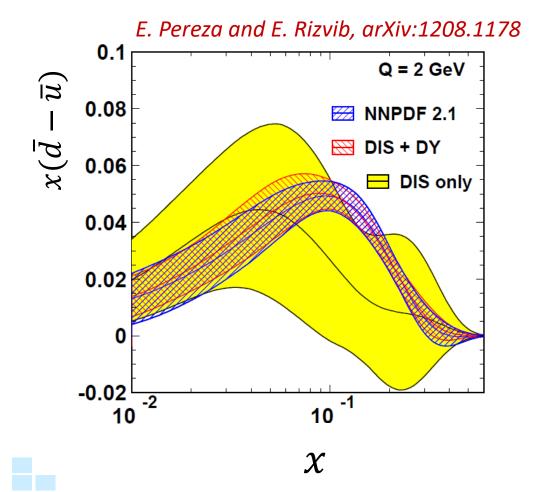
Process	Subprocess	Partons	x range	
$\ell^{\pm}\{p,n\} \to \ell^{\pm}X$	$\gamma^* q \to q$	q, \bar{q}, g	$x \gtrsim 0.01$	DIC
$\ell^{\pm}n/p \to \ell^{\pm}X$	$\gamma^* d/u \to d/u$	d/u	$x \gtrsim 0.01$	DIS
$pp \rightarrow \mu^+\mu^- X$	$u\bar{u}, dd \rightarrow \gamma^*$	$ar{q}$	$0.015 \lesssim x \lesssim 0.35$	
$pn/pp \rightarrow \mu^+\mu^- X$	$(u\bar{d})/(u\bar{u}) \to \gamma^*$	\bar{d}/\bar{u}	$0.015 \lesssim x \lesssim 0.35$	DY
$\nu(\bar{\nu})N \to \mu^-(\mu^+)X$	$W^*q \rightarrow q'$	$q,ar{q}$	$0.01 \lesssim x \lesssim 0.5$	
$\nu N \to \mu^- \mu^+ X$	$W^*s \to c$	S	$0.01 \lesssim x \lesssim 0.2$	
$\bar{\nu}N \to \mu^+\mu^- X$	$W^*\bar{s} \to \bar{c}$	\bar{s}	$0.01 \lesssim x \lesssim 0.2$	חוכ
$e^{\pm}p \rightarrow e^{\pm}X$	$\gamma^* q \to q$	$g,q,ar{q}$	$0.0001 \lesssim x \lesssim 0.1$	DIS
$e^+p \rightarrow \bar{\nu} X$	$W^+\{d,s\}\to\{u,c\}$	d, s	$x \gtrsim 0.01$	
$e^{\pm}p \rightarrow e^{\pm}c\bar{c}X$	$\gamma^*c \to c, \gamma^*g \to c\bar{c}$	c, g	$0.0001 \lesssim x \lesssim 0.01$	
$e^{\pm}p \to \text{jet} + X$	$\gamma^* g \to q \bar{q}$	g	$0.01 \lesssim x \lesssim 0.1$	
$p\bar{p} \to \text{jet} + X$	$gg, qg, qq \rightarrow 2j$	g,q	$0.01 \lesssim x \lesssim 0.5$	
$p\bar{p} \to (W^{\pm} \to \ell^{\pm} \nu) X$	$ud \to W, \bar{u}\bar{d} \to W$	u,d,\bar{u},\bar{d}	$x \gtrsim 0.05$	
$p\bar{p} \to (Z \to \ell^+\ell^-)X$	$uu, dd \rightarrow Z$	d	$x \gtrsim 0.05$	

Both DIS and Drell-Yan process are powerful tools to probe the quark and anti-quark structure in hadrons (Universality, Factorization)

Extraction of **PDFs**

Complementarity between **Space-like** and **Time-like** processes

The asymmetry of the light sea " $\bar{d}(x) - \bar{u}(x)$ " and its one standard deviation uncertainty



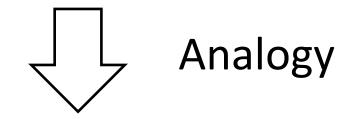
Example:

Constraint of $x(\bar{d} - \bar{u})$ in Global Analysis

- When only DIS data are included in the fit
 - → Yellow contour
- When DY data are included in addition
 - → Red hashed contour
- Reference NNPDF2.1 fit
 - → Blue hashed contour

Use of both DIS and DY provides the strong constrain of $x(\bar{d} - \bar{u})$

The measurements of both Space-like process (DIS) and Time-like process (DY) are helpful to probe the PDFs.



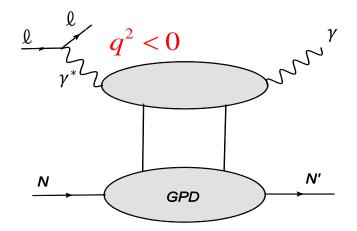
The measurements of both Space-like process and Time-like process are helpful to probe the GPDs?

Extraction of **GPDs**

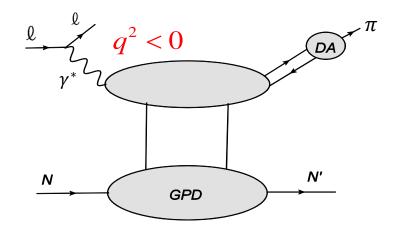
by Lepton Beam (Space-like Processes)

Muller et al., PRD 86 031502(R) (2012)

Deeply Virtual Compton Scattering (DVCS)



Deeply Virtual Meson Production (DVMP)



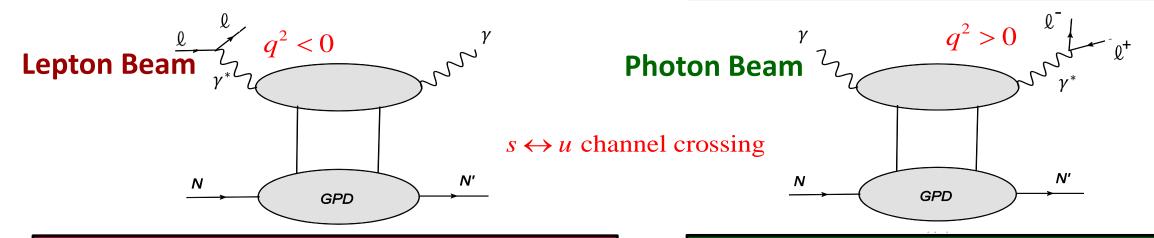
Extraction of **GPDs**

Space-like vs. **Time-like** Processes

Muller et al., PRD 86 031502(R) (2012)

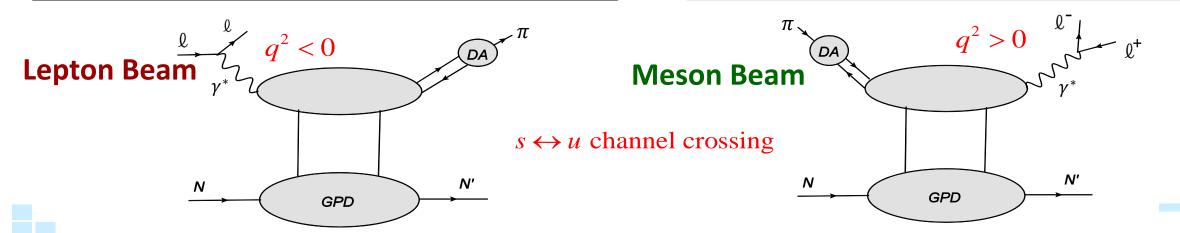
Deeply Virtual Compton Scattering (DVCS)

Time-like Compton Scattering (TCS)



Deeply Virtual Meson Production (DVMP)

Exclusive meson-induced DY



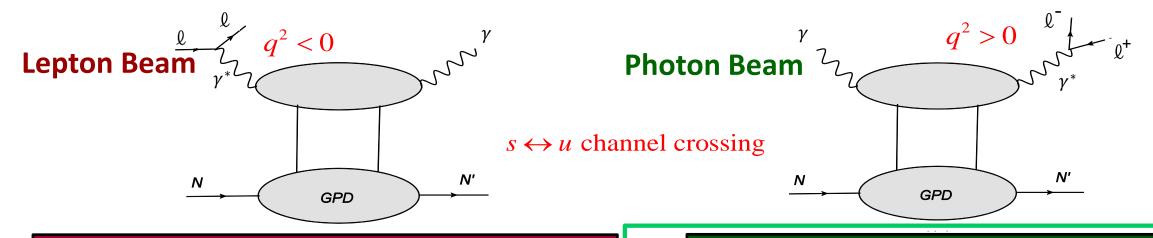
Extraction of **GPDs**

Space-like vs. **Time-like** Processes

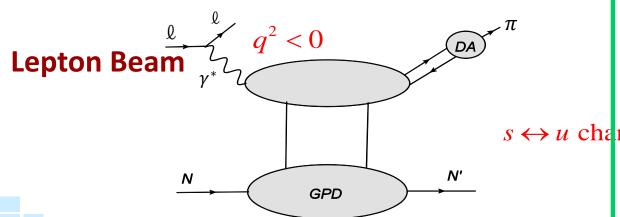
Muller et al., PRD 86 031502(R) (2012)

Deeply Virtual Compton Scattering (DVCS)

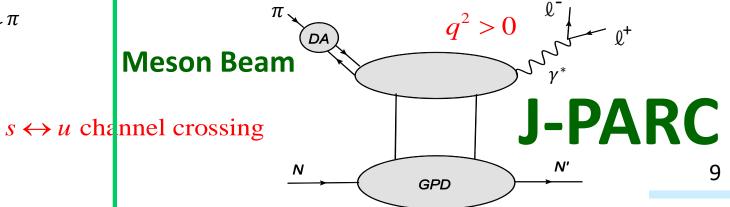
Time-like Compton Scattering (TCS)



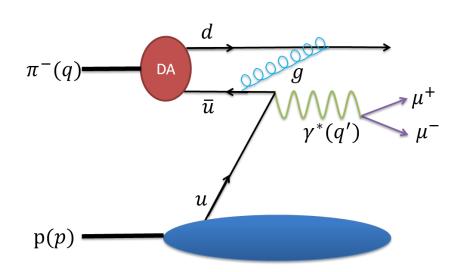
Deeply Virtual Meson Production (DVMP)



Exclusive meson-induced DY



Semi-exclusive pion-induced DY at large x_{π}



When Longitudinal momentum fraction $x_{\pi} \rightarrow 1$

- Virtual photon polarization:

Transverse → Longitudinal

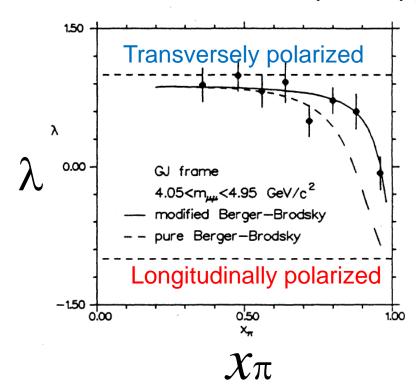
- Dimuon angular distribution:

$$(1 + \cos^2\theta) \rightarrow \sin^2\theta$$

Dominance of higher-twist contributions in the forward production

Drell-Yan decay angular distributions

E615, PRD 39, 92 (1989)

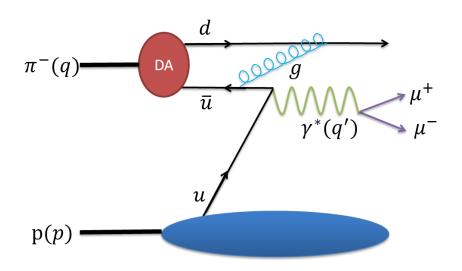


Semi-exclusive pion-induced DY at large x_{π}



Exclusive pion-induced DY $\pi N \rightarrow \gamma^* N'$







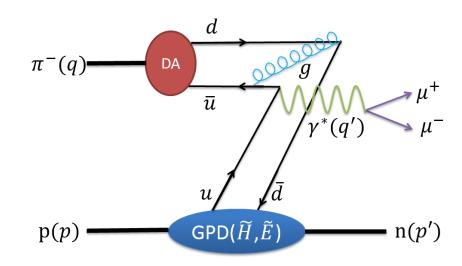
- Virtual photon polarization:

Transverse → Longitudinal

- Dimuon angular distribution:

$$(1 + \cos^2\theta) \rightarrow \sin^2\theta$$

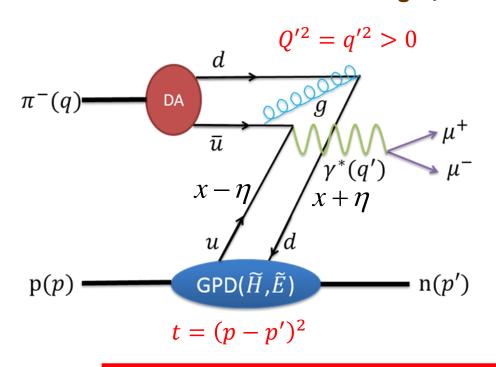
Dominance of higher-twist contributions in the forward production



- Spectator quark originating from the pion may be absorbed by the remnant of the target
- The target matrix element is given by a GPD with skewness

Exclusive Drell-Yan process: $\pi^- p \rightarrow \mu^+ \mu^- n$

E.R. Berger, M. Diehl, B. Pire, PLB 523 (2001) 265



$$\tau = \frac{Q'^2}{2pq} \approx \frac{Q'^2}{s - M_N^2} \qquad \xi = \frac{(p - p')^+}{(p + p')^+} = \frac{\tau}{2 - \tau} \quad \text{(skewness)}$$

$$2pq \quad s - M_N^2 \qquad (p + p')^+ \quad 2 - \tau$$

$$\chi^*(q') \qquad \mu^+ \qquad \tilde{x} = -\frac{(q + q')^2}{2(p + p') \cdot (q + q')} \approx -\frac{Q'^2}{2s - Q'^2} = -\xi \quad \text{(scaling variable)}$$

$$\tilde{\mathcal{H}}^{du}(\tilde{x}, \xi, t) = \frac{8}{2}\alpha_s \int_0^1 dz \, \frac{\phi_\pi(z)}{z} \, dz \, \frac{\phi_\pi(z)}{z} \, dz$$
Twist-two Pion DA

$$\tilde{\mathcal{H}}^{du}(\tilde{x},\xi,t) = \frac{8}{3}\alpha_s \int_{-1}^{1} dz \frac{\phi_{\pi}(z)}{1-z^2} \qquad \text{Twist-two Pion DA}$$

$$\times \int_{-1}^{1} dx \left(\frac{e_d}{\tilde{x}-x-i\epsilon} - \frac{e_u}{\tilde{x}+x-i\epsilon} \right) \qquad \text{Electric charges}$$

$$\times (\tilde{H}^d(x,\xi,t) - \tilde{H}^u(x,\xi,t)), \longleftarrow \qquad \text{GPD } \tilde{H}$$
for $p{\to}n$ transition

Leading-twist cross section:

$$\begin{split} \frac{d\sigma_L}{dt dQ'^2}\bigg|_{\tau} &= \frac{4\pi\alpha_{\rm em}^2}{27}\frac{\tau^2}{Q'^8}f_{\pi}^2\left[(1-\xi^2)|\tilde{\mathcal{H}}^{du}(\tilde{x},\xi,t)|^2 \right. \\ &\left. -2\xi^2{\rm Re}\left(\tilde{\mathcal{H}}^{du}(\tilde{x},\xi,t)^*\tilde{\mathcal{E}}^{du}(\tilde{x},\xi,t)\right) - \xi^2\frac{t}{4m_N^2}|\tilde{\mathcal{E}}^{du}(\tilde{x},\xi,t)|^2\right], \end{split}$$

Two GPDs inputs

BMP2001

E. R. Berger, M. Diehl and B. Pire, Eur. Phys. J. C 23 (2002) 675

E. R. Berger, M. Diehl and B. Pire, Phys. Lett. B 523 (2001) 265

$$\tilde{H}^{d,u}(x,\xi,t) = \tilde{H}^{d,u}(x,\xi,0) [g_A(t)/g_A(0)]$$

Here, $\tilde{H}^q(x,\xi,0)$ constructed from an ansatz based on double distributions as an integral of $\tilde{H}^q(x,0,0)=\Delta q(x)$ combined with a certain profile function generating the skewness ξ dependence

$$\phi_\pi(z) o (3/4)(1-z^2)$$
 asymptotic form

GK2013

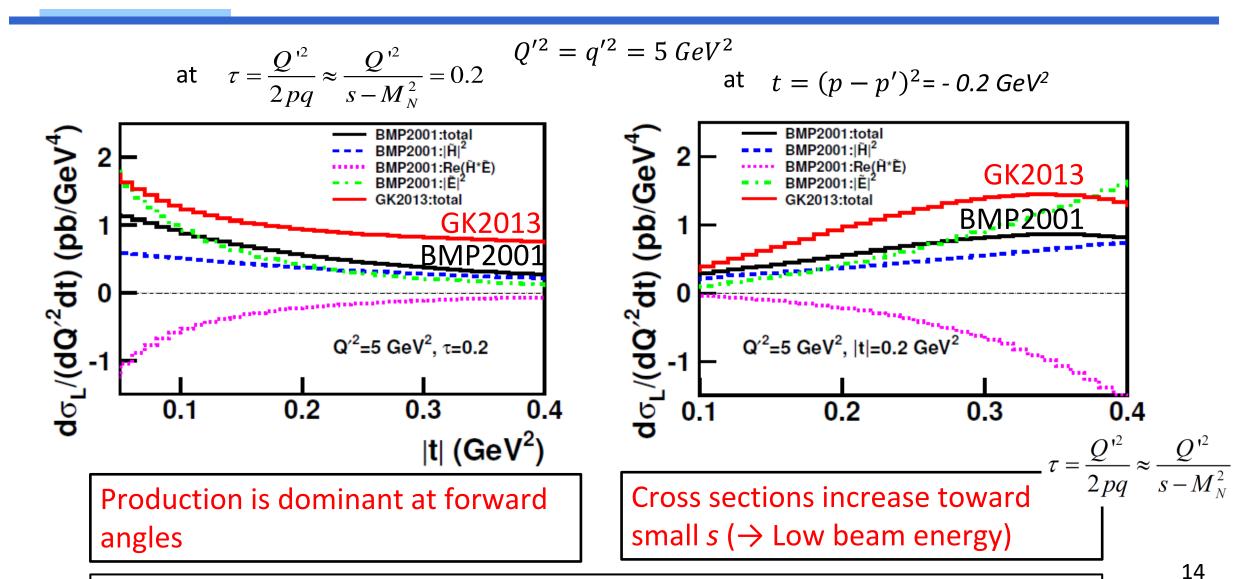
P. Kroll, H. Moutarde and F. Sabatie, Eur. Phys. J. C 73 (2013) 1, 2278

The parameters are determined from the HERMES data on the cross sections and target asymmetries for pi+ electroproduction

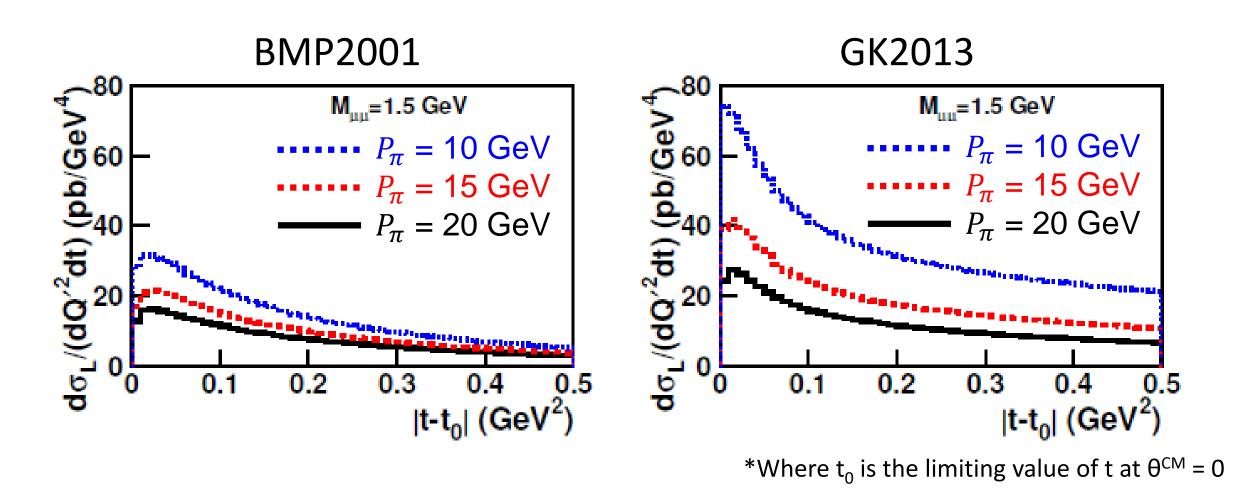
$$\phi_{\pi}(z) = (3/4)(1-z^2)[1+a_2C_2^{(3/2)}(z)]$$

with a_2 ($\mu = 2$ GeV) = 0.22

Differential cross sections of exclusive Drell-Yan

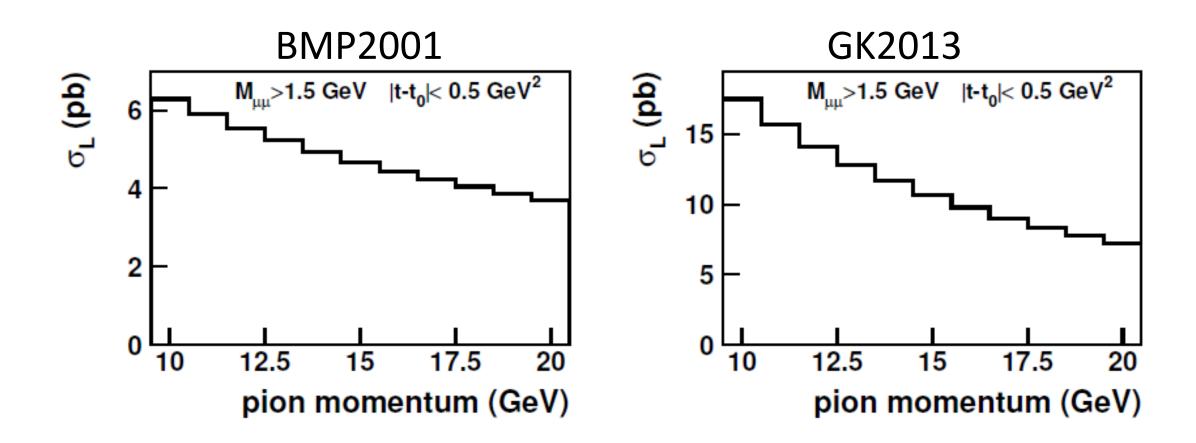


Differential cross sections of exclusive Drell-Yan



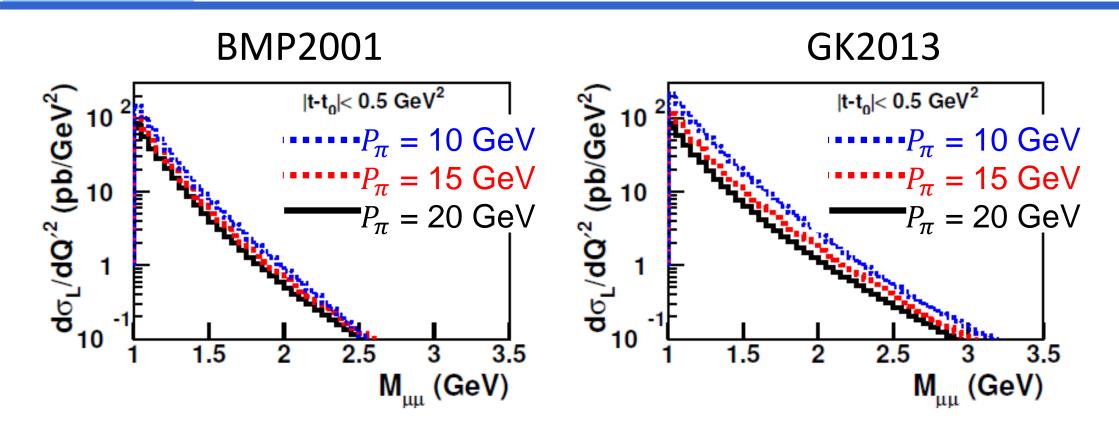
Pion beam with lower momentum → Larger cross section

Total LO cross sections of exclusive Drell-Yan



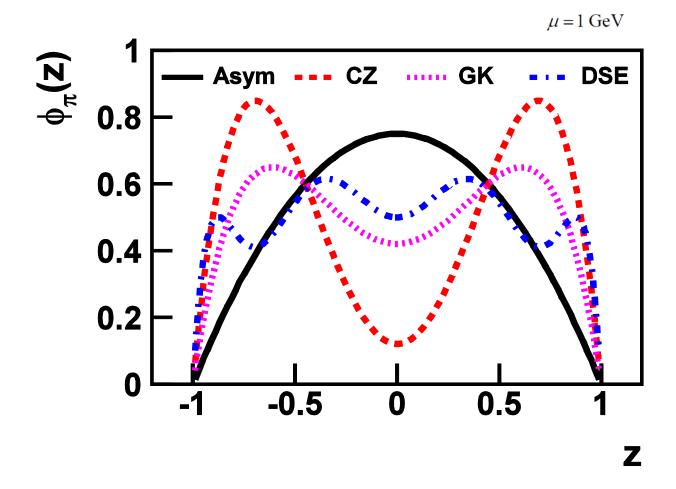
J-PARC (P
$$\pi$$
 = 10-20 GeV) σ = 5 ~ 15 pb CERN COMPASS (P π = 190 GeV) σ = 0.65 pb

Differential cross sections of exclusive Drell-Yan



Dimuons with $1.5 < M_{\mu^+\mu^-} < 2.9$ GeV (below J/ ψ mass) will be used for the exclusive DY analysis. The contribution from exclusive DY could be well separated from others (Accidental BG, Inclusive DY, etc) by using the missing mass technic.

Sensitivity to Pion DAs



G. P. Lepage et al., Phys. Lett. B 87, 359 (1979)
A. V. Efremov et al., Phys. Lett. B 94, 245 (1980)
V. L. Chernyak et al., Phys. Rept. 112 (1984) 173
P. Kroll et al., Eur. Phys. J. C 73 (2013) 1, 2278
I. C. Cloet et al., Phys. Rev. Lett. 111 (2013) 092001

TABLE I: Modeling of pion DAs

$\phi_{\pi}(z,\mu)$	Asymptotic [41]	CZ [42]	GK [51]	DSE [52]
a_2	0	2/3	0.22	0.20
a_4	0	0	0	0.093
a_6	0	0	0	0.055
$\mu^2 ({\rm GeV^2})$	1	0.25	4	4

Asym: Asymptotic form

CZ: by Chernyak and Zhitnitsky

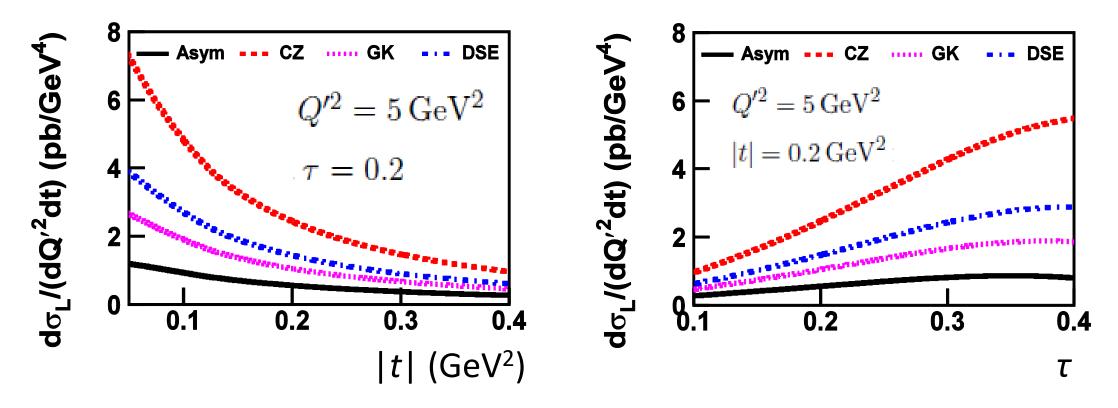
GK: by Goloskokov-Kroll

DSE: using the Dyson-Schwinger

equation framework

Sensitivity to Pion DAs

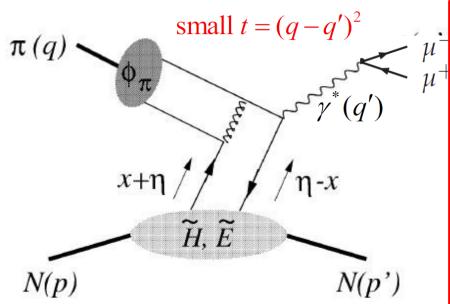
Fixed the GPDs to "BMP2001", and replaced the pion DAs



Differential cross section of exclusive DY is strongly sensitive to the input of pion DAs

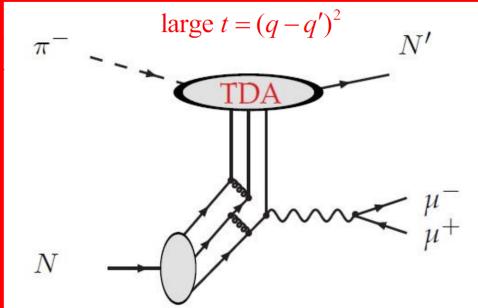
Exclusive DY at large momentum transfer | t |

Bernard Pire, IWHS2011



 ϕ_{π} : pion distritribution amplitude (DA)

- •DA characterizes the minimal valence Fock state of hadrons.
- •DA of pion are also explored by pionphoton transition form factor in Belle and Barbar Exps.



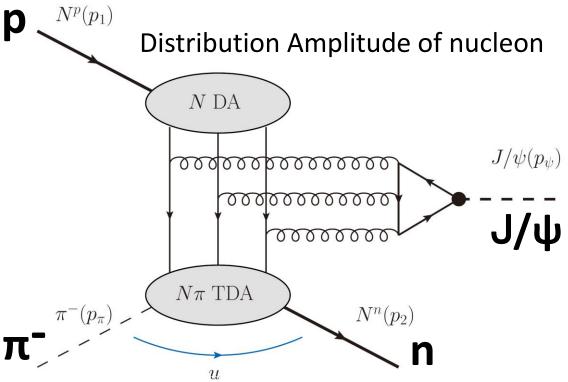
TDA : π -N transistion distritribution amplitude

- •TDA characterizes the next-to-minimal valence Fock state of hadrons.
- •TDA of pion-nucleon is related to the pion cloud of nucleons.

52

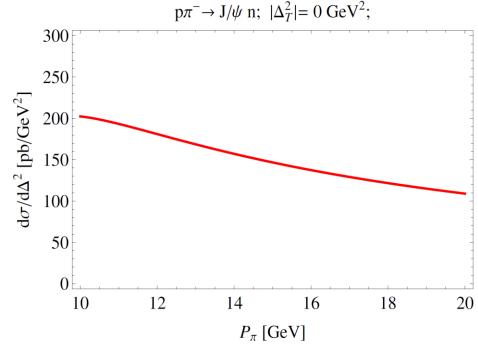
Pion-induced exclusive backward J/ψ production

B. Pire et.al Phys. Rev. D 95, 034021 (2017)



Transition Distribution Amplitude for pion → nucleon

Differential cross section at J-PARC energy



Momentum of Pion Beam (GeV)

might be observed in J-PARC?

Exclusive DY at small momentum transfer | t |

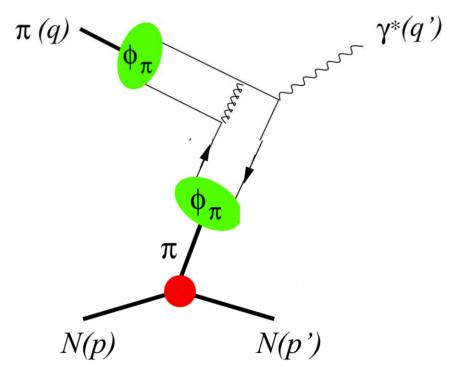
Leading-Twist Diagram

N(p)

$\pi(q)$ ϕ_{π} $x+\eta$ $\eta-x$ H, \widetilde{E}

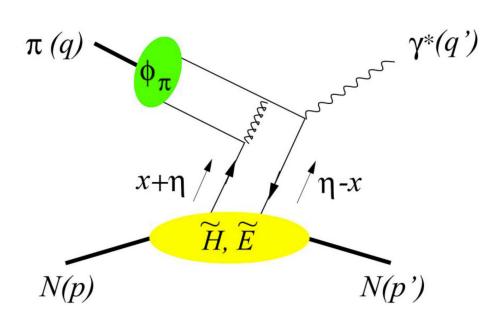
N(p')

Pion-pole Dominance for \tilde{E}

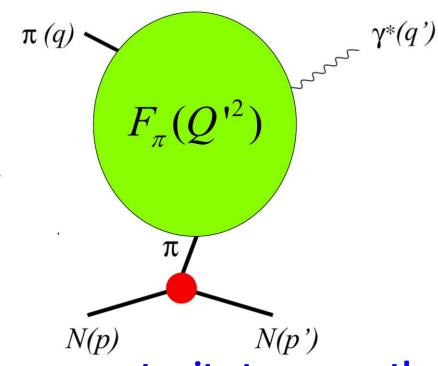


Exclusive DY at small momentum transfer |t|

Leading-Twist Diagram



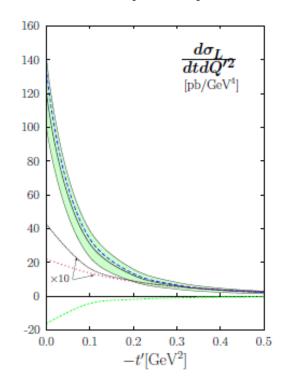
Time-like Pion FF

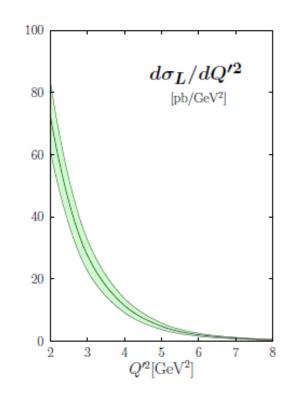


Unique opportunity to access the pion timelike FF, compared with the other measurements in e^+e^-

Differential cross sections with an updated time-like pion FF

S.V. Goloskokov, P. Kroll, PLB 748 (2015) 323





 $Q'^2=4\,{\rm GeV^2}$ and $s=20\,{\rm GeV^2}$ solid lines with error bands: full result pion pole, $|\langle\widetilde{H}^{(3)}\rangle|^2$, interference, short dashed: leading-twist contribution

time-like pion FF: $Q'^2|F_\pi(Q'^2)| = 0.88 \pm 0.04 \, \mathrm{GeV^2}$ (CLEO, BaBar, $J/\Psi \to \pi^+\pi^-$) phase $(\exp{[i\delta(Q'^2)]})$ from disp. rel. Belicka et al(11) for $Q'^2 < 8.9 \, \mathrm{GeV^2}$ $\delta = 1.014\pi + 0.195 (Q'^2/\mathrm{GeV^2} - 2) - 0.029 (Q'^2/\mathrm{GeV^2} - 2)^2$ for $Q'^2 \geq 8.9 \, \mathrm{GeV^2}$: $\delta = \pi$, the LO pQCD result

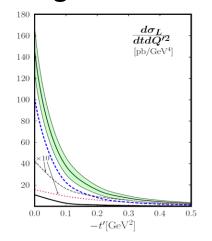
Beyond the Leading Twist

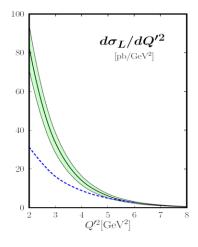
S.V. Goloskokov, P. Kroll, PLB 748 (2015) 323

$$\frac{d\sigma}{dt dQ'^2 d\cos\theta d\phi} = \frac{3}{8\pi} \left\{ \sin^2\theta \frac{d\sigma_L}{dt dQ'^2} + \frac{1 + \cos^2\theta}{2} \frac{d\sigma_T}{dt dQ'^2} + \frac{\sin(2\theta)\cos\phi}{\sqrt{2}} \frac{d\sigma_{LT}}{dt dQ'^2} + \sin^2\theta\cos(2\phi) \frac{d\sigma_{TT}}{dt dQ'^2} \right\}$$

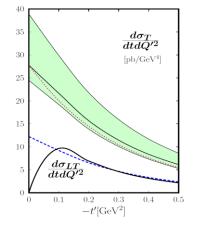
The measurement of angular distributions provides the each terms

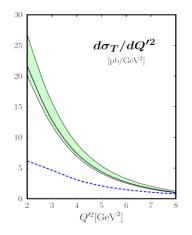
Longitudinal cross sections



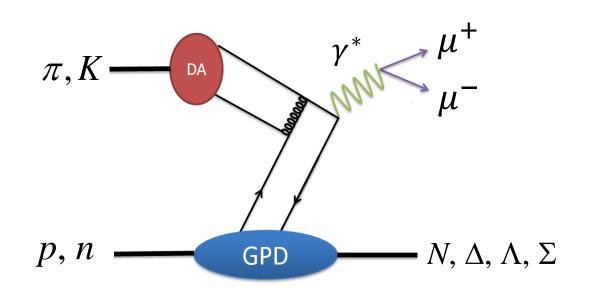


Transverse cross sections





Accessing transition GPD via the exclusive dimuon measurement



•
$$\pi^- p \rightarrow \gamma^* n$$

•
$$\pi^- p \rightarrow \gamma^* \Delta^0$$

•
$$\pi^- n \rightarrow \gamma^* \Delta^-$$

•
$$\pi^+ n \rightarrow \gamma^* p$$

•
$$\pi^+ p \rightarrow \gamma^* \Delta^{++}$$

•
$$\pi^+ n \rightarrow \gamma^* \Delta^+$$

Talked by S. Kumano

•
$$K^-p \rightarrow \gamma^* \Lambda$$

•
$$K^-p \rightarrow \gamma^* \Lambda(1405)$$

•
$$K^- p \rightarrow \gamma^* \Lambda(1520)$$

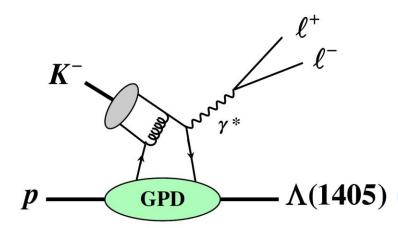
•
$$K^- n \rightarrow \gamma^* \Sigma^-$$

•
$$K^+ n \rightarrow \gamma^* \Theta^+$$

The exotic hadrons cannot exist as stable fixed targets.

→ The GPDs of the exotic hadrons could not be observed directly.

However, the transition GPDs such as for proton \rightarrow an exotic hadron can be investigated. For example, the transition GPDs of $p \rightarrow \Lambda(1405)$ should reflect the exotic nature of $\Lambda(1405)$.



Soft Nonfactorizable Mechanism

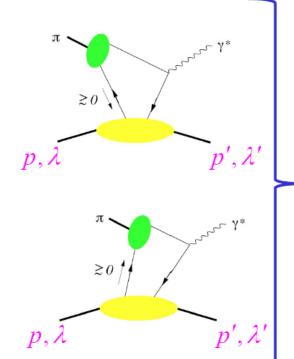
LCSR for nonfactorizable amplitude

SNM > factorization

interplay of soft/hard QCD mechanism

K. Tanaka, arXiv:1703.02190

"nonfactorizable" mechanism



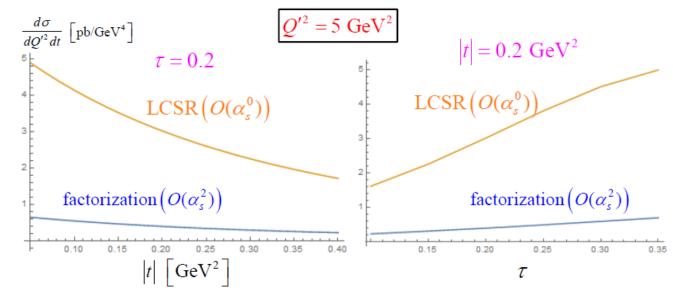
"Light-cone QCD SR (LCSR)"

$$= g_{\nu}^{-} \int_{\eta}^{x_{0}} dx \, e^{-\frac{x-\eta \, Q'^{2}}{x+\eta \, M_{B}^{2}}} \tilde{C}_{H}(x,\eta,Q'^{2})$$

$$\times \left[e_{u} \tilde{H}^{du}(x,\eta,t) - e_{d} \tilde{H}^{du}(-x,\eta,t) \right]$$

$$\times \overline{u}(p'\lambda') \gamma^{+} \gamma_{5} u(p\lambda) + \cdots$$

$$\tilde{H}^{du}(x,\eta,t) = \tilde{H}^{u}(x,\eta,t) - \tilde{H}^{d}(x,\eta,t)$$



$$\frac{d\sigma}{dQ'^{2}dt}(\pi^{-}p \to \gamma^{*}n)$$

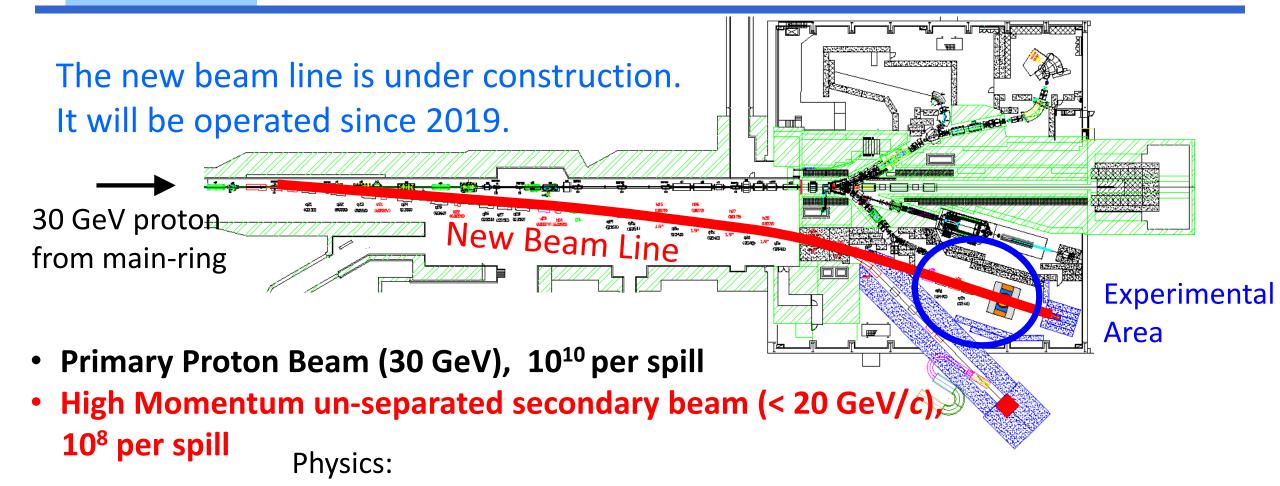
$$= \frac{4\pi\alpha_{\text{em}}^{2}}{27} \frac{\tau^{2}}{Q'^{8}} f_{\pi}^{2} \left[(1-\eta^{2}) |\mathcal{J}^{du}|^{2} - 2\eta^{2} \operatorname{Re}(\mathcal{J}^{du^{*}} \mathcal{E}^{du}) - \eta^{2} \frac{t}{4M^{2}} |\mathcal{E}^{du}|^{2} \right]$$

Experiment at J-PARC



High Momentum Beam Line at J-PARC

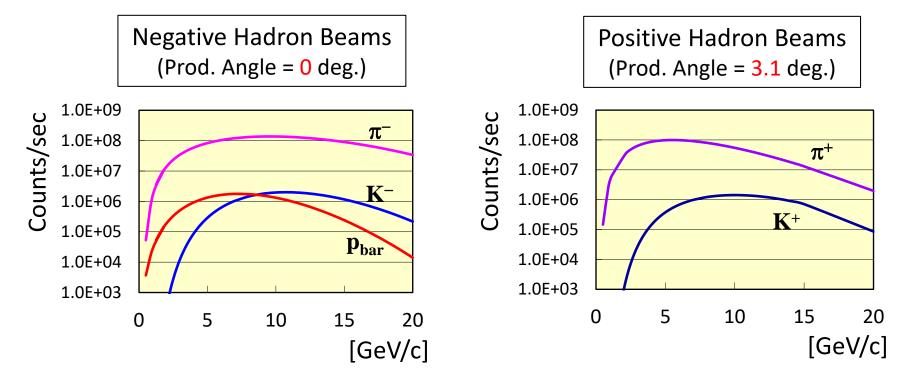
Talked by Shin'ya Sawada



Vector meson modification in the nuclear matter Charmed Baryon spectroscopy Nucleon Structure

Unseparated secondary beams

- High-intensity secondary Pion beam
- High-resolution beam: $\Delta p/p \sim 0.1\%$



³¹

Uniqueness of pion-induced Drell-Yan physics studied at High Momentum Beam Line in J-PARC

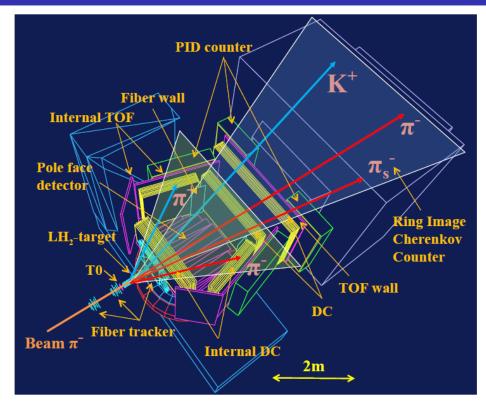
1970s -1980s	Beam Particle	Beam Momentum(Ge	eV) Target	
CERN-NA3 CERN-WA11 CERN-WA39 Fermilab-E326 Fermilab-E537 Fermilab-E615 2010s CERN-COMPASS	π^{\pm} π^{-} π^{-} π^{-} π^{-} π^{-}	150/200/280 140/194/286 150/175 39.5 225 125 252	H, Pt D, W Be W W W W	Inclusive Drell-Yan $\pi^{\pm}N \to \mu^{+}\mu^{-}X$
J-PARC	π^- 1	10 ~ 20	Н	Exclusive Drell-Yan $\pi^-p o \mu^+\mu^-n$ + Inclusive Drell-Yan $\pi^-p o \mu^+\mu^-X$

J-PARC E50 Experiment

for Charmed Baryon Spectroscopy

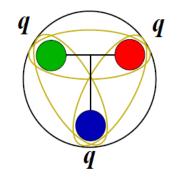
Talked by Hiroyuki Noumi

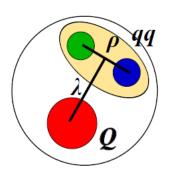
Spectrometer (Charmed baryon prod. and decay)



Stage-1 approved by J-PARC PAC-18, August 12, 2014.

Quark correlation





- color-spin interaction between quarks $\propto 1/m_i m_j$
- 3 light diquark pairs ⇒ difficult to distinguish
- Heavy $Q \Rightarrow$ separate to Q and q q

We will investigate the diquark correlation by measurement of charmed baryon's properties

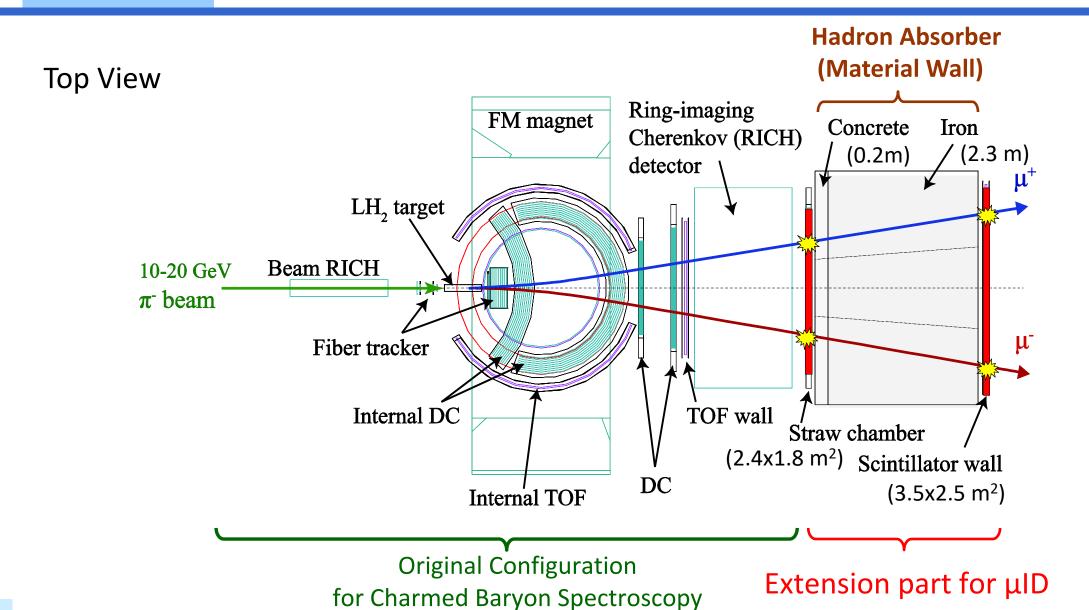
- Level structure
- Production rate

15/24

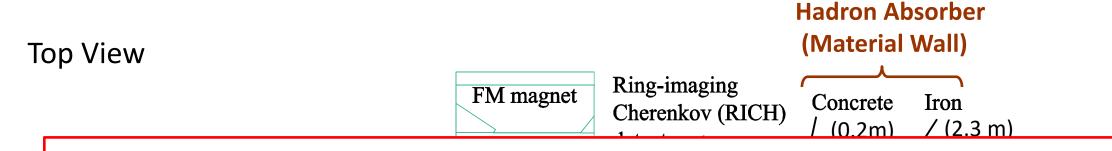
Decay branching ratio

4/24

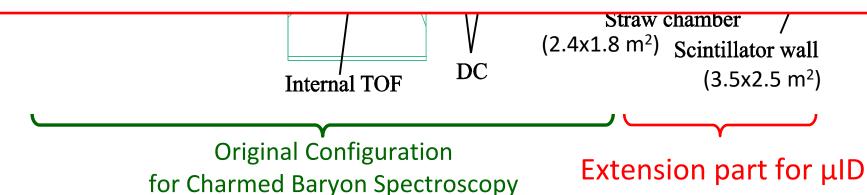
Extension of J-PARC E50 experiment for Drell-Yan measurement



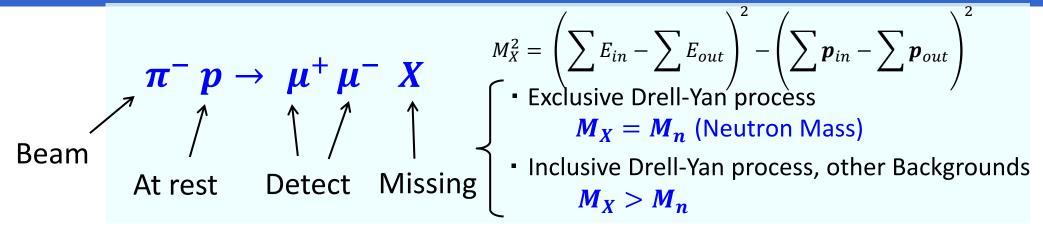
Extension of J-PARC E50 experiment for Drell-Yan measurement



DY trigger rate is expected to be very limited. So then, the DY measurement and the charmed-baryon spectroscopy could be carried out together in the E50 experiment

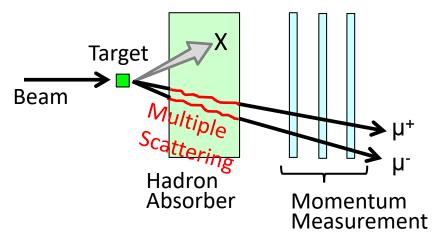


Identifying the Exclusive Drell-Yan Process with Missing Mass Technique



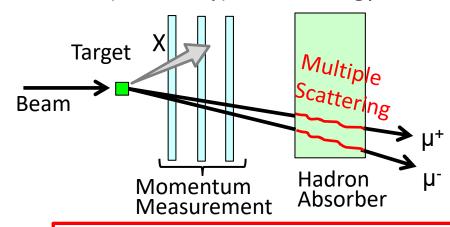
Typical Drell-Yan experiments

with high energy beam



Proposed Drell-Yan experiment at J-PARC

with (Relatively) lower energy beam



Open-aperture spectrometers

 \rightarrow Good resolution for M_{χ}

Yield Estimation

15 GeV/c 20 GeV/c 10 GeV/c Beam Energy Beam Intensity High Low **Total Cross Section** of Exclusive DY High Low High of Inclusive DY Low Acceptance High Low for exclusive DY π⁻ beam π^+ beam (prod. angle 0 deg) (prod. angle 3.1 deg) Beam Intensity High ← **Total Cross Section of DY** High ← Low

Simulation

Assumptions:

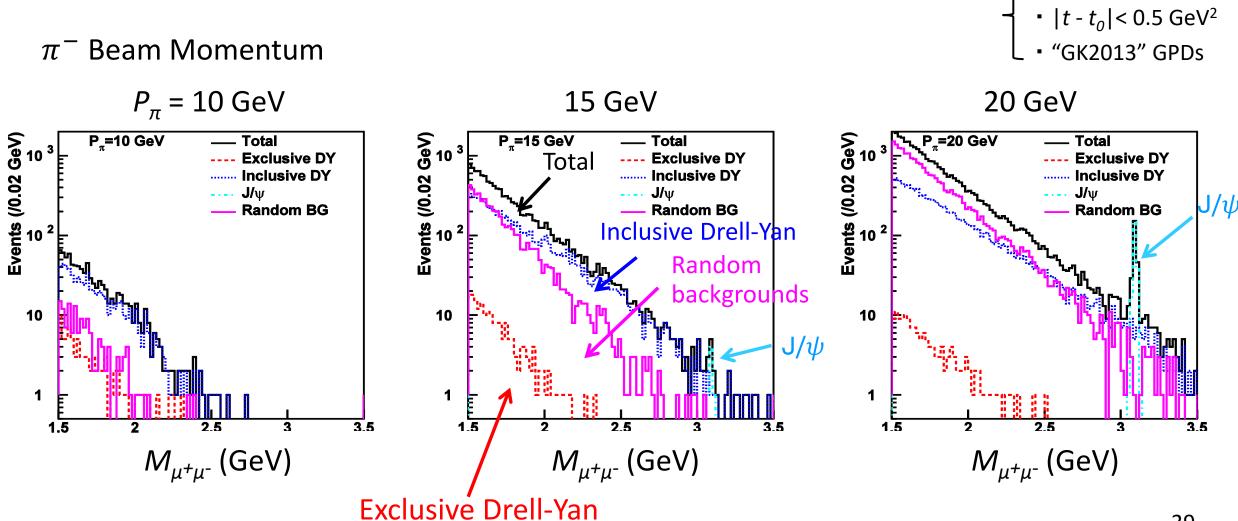
- Target: 57cm LH₂ $(n_{TGT} = 4 \text{ g/cm}^2)$
- Beam momentum resolution($\Delta p/p$) = 0.1 %
- $1.83/1.58/1.00 *10^7 \pi^-$ /spill for 10/15/20 GeV beam
- Data Taking: 50 days (→ 100 days or more)
- E50 spectrometers + μ ID system

Expected cross sections for the exclusive/inclusive Drell-Yan processes

	Exclusive Drell-Yan		Inclusive Drell-Yan
	$\left(\begin{array}{l} M_{\mu^{+}\mu^{-}} > 1.5 \text{ GeV}, \\ t - t_{0} < 0.5 \text{ GeV}^{2} \end{array} \right)$		$(M_{\mu^+\mu^-} > 1.5 \text{ GeV})$
	"BMP2001"	"GK2013"	
P_{π} = 10 GeV	6.29 pb	17.53 pb	2.11 nb
P_{π} = 15 GeV	4.67 pb	10.65 pb	2.71 nb
P_{π} = 20 GeV	3.70 pb	7.25 pb	3.08 nb

• Total hadronic interaction cross sections of π^-p is about 20-30 mb while the production of ${\rm J}/\psi$ is about 1-3 nb

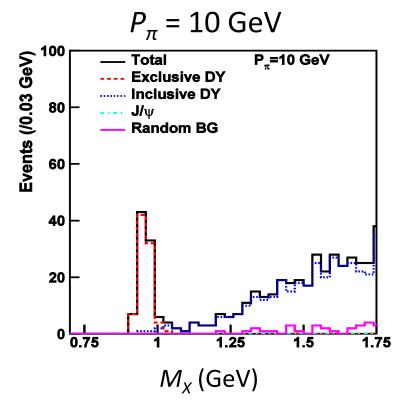
Simulated invariant mass $M_{\mu^+\mu^-}$ spectra of the $\mu^+\mu^-$ events

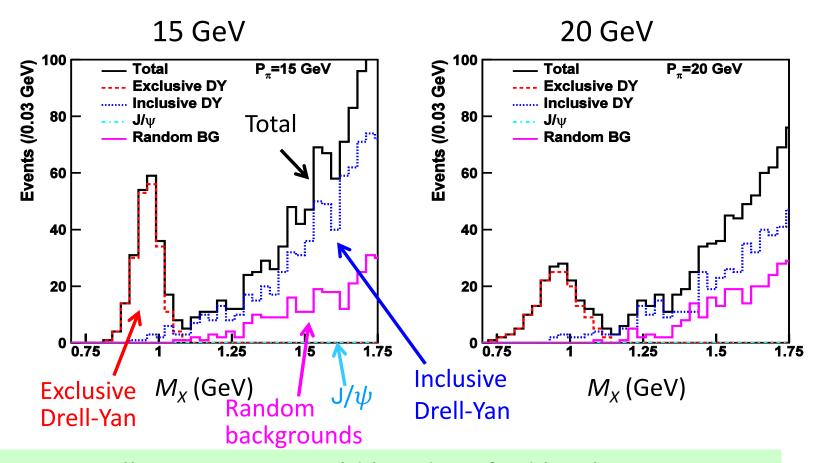


Data Taking: 50 days

Simulated missing-mass M_X spectra of the $\mu^+ \mu^-$ events

π^- Beam Momentum



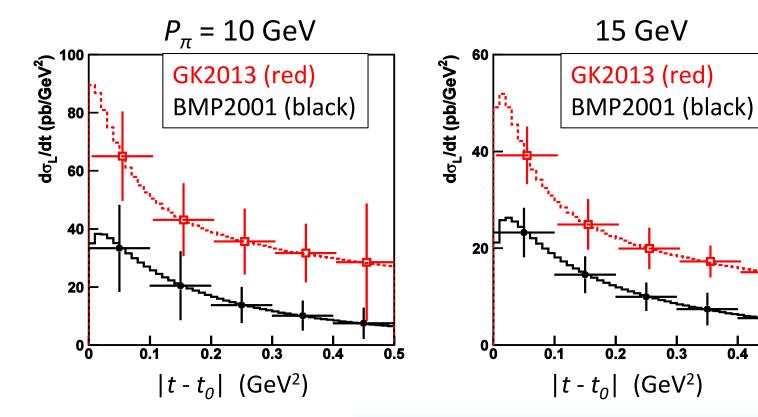


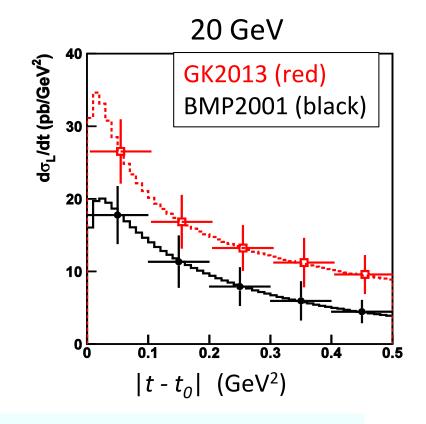
- Data Taking: 50 days
- $1.5 < M_{\mu^+\mu^-} < 2.9 \text{ GeV}$
- $|t t_0| < 0.5 \text{ GeV}^2$
- "GK2013" GPDs

The exclusive Drell-Yan events could be identified by the signature peak at the nucleon mass in the missing-mass spectrum for all three pion beam momenta.

The expected statistical errors of the exclusive Drell-Yan measurement for two GPDs inputs

π^- Beam Momentum

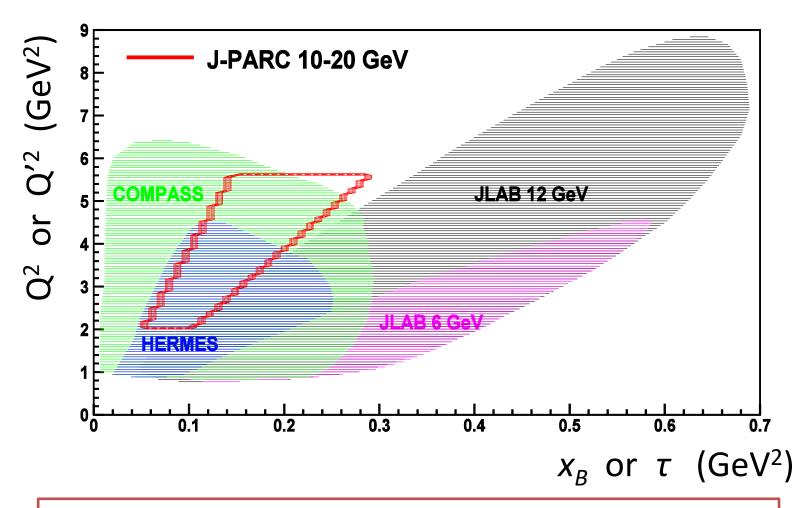




- Data Taking: 50 days
- $1.5 < M_{\mu^+\mu^-} < 2.9 \text{ GeV}$
- $|t t_0| < 0.5 \text{ GeV}^2$

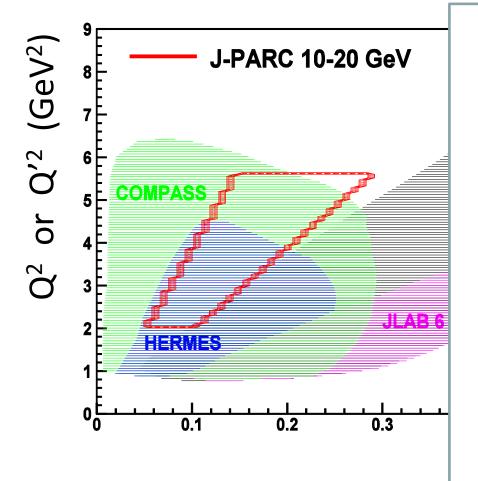
The statistics accuracy is adequate for discriminating between the predictions from two current GPD modelings.

Kinematic regions of GPDs explored by space-like and time-like processes



- JLAB, HERMES, COMPASS: Spece-like approach
- J-PARC: Time-like approach

Kinematic regions of GPDs explored by space-like and time-like processes

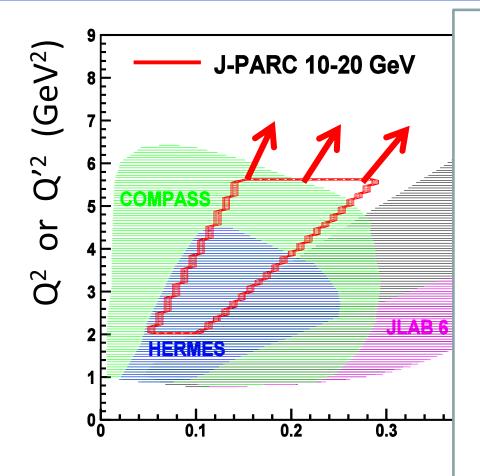


Impacts of GPD measurement at J-PARC:

- Test of universality of GPD in space-like and time-like processes
- Test of factorization of exclusive Drell-Yan process

- JLAB, HERMES, COMPASS: 3
- J-PARC: Time-like approach

Kinematic regions of GPDs explored by space-like and time-like processes



- JLAB, HERMES, COMPASS:
- J-PARC: Time-like approach

Impacts of GPD measurement at J-PARC:

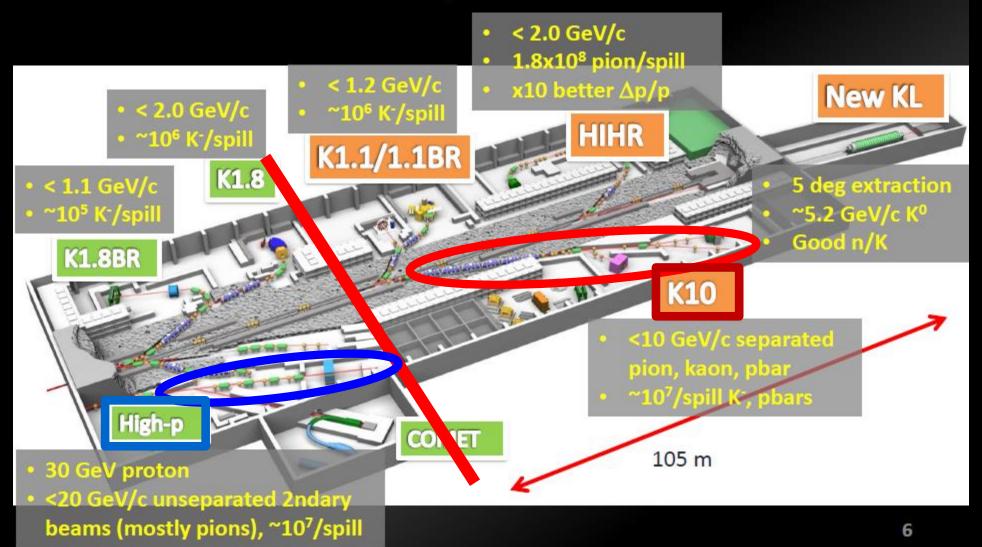
- Test of universality of GPD in space-like and time-like processes
- Test of factorization of exclusive Drell-Yan process

Further possibilities:

- Information of GPD at large-Q'² region
- Test of QCD-evolution properties of GPD

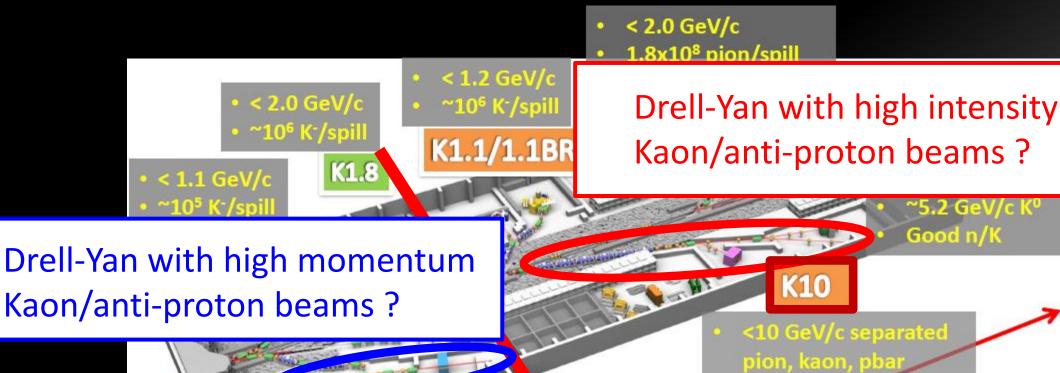
Talked by Shin'ya Sawada

Overview: Extension of the Hadron Experimental Facility



Talked by Shin'ya Sawada

Overview: Extension of the Hadron Experimental Facility



30 GeV proton

 <20 GeV/c unseparated 2ndary beams (mostly pions), ~10⁷/spill 105 m

~10⁷/spill K, pbars

Summary

- Drell-Yan process is a powerful tool to explore the partonic structures of nucleons. We
 are moving ahead from 1D to 3D imaging of nucleons
- Measurement of GPDs through the exclusive π -induced Drell-Yan process will offer
 - Test of universality of GPD in space-like and time-like processes.
 - Test of factorization of exclusive Drell-Yan process.
 - Pion DA, Pion timelike FF (at small |t|), Transition DA (at large |t|), Transition GPD
 - Transition of Inclusive DY \rightarrow Semi-inclusive DY \rightarrow Exclusive DY

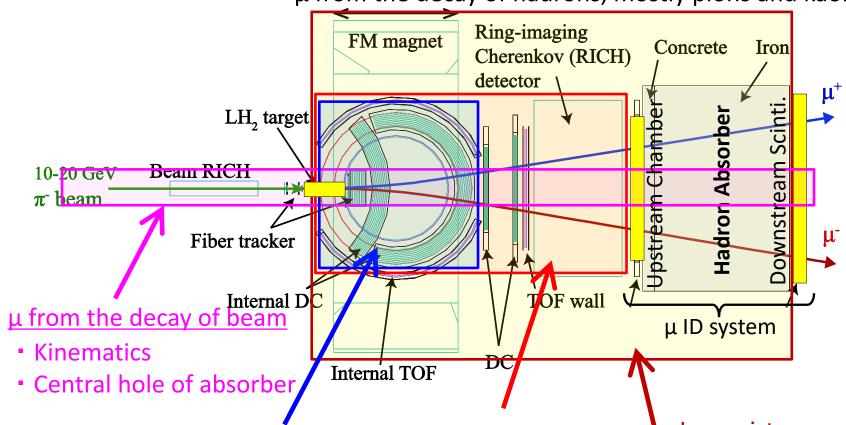
with an increase of beam time (50 days in simulation \rightarrow 100 days or more) or beam luminosity, optimization of setup:

- GPD at large-Q'² region
- QCD-evolution properties of GPD
- The preliminary study shows that such measurement is feasible with E-50 spectrometers in the coming high momentum beamline at J-PARC.

Spares

Background μ rejection (Offline)

μ from the decay of hadrons, mostly pions and kaons



- Bad χ² probability in the reconstructions
- Vertex position
- Kink of the decay vertex

•Inconsistency of the trajectory between the spectrometer and the upstream chamber in µID system

•Inconsistency of the hit location between the upstream chamber and downstream Scinti. in µID system