

# Phenomenology of Right-Handed Neutrinos in the Littlest Seesaw

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Advances in High Energy Physics and Cosmology  
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22 Mar 2018

ongoing work with S. Molina-Sedgwick and Prof. S.F. King

- Background
  - ▶ Type-I Seesaw Mechanism
  - ▶ Leptogenesis
  - ▶ Minimal Model: *The Littlest Seesaw*
- This Work: *Phenomenology of RH Neutrinos*
- Preliminary Results
- Outlook

- Standard Model cannot explain **neutrino masses and oscillations** or the observed **BAU**

$$Y_B \equiv \frac{n_B - n_{\bar{B}}}{s} > 0$$

- Possible solution: models that utilise a Type-I Seesaw mechanism

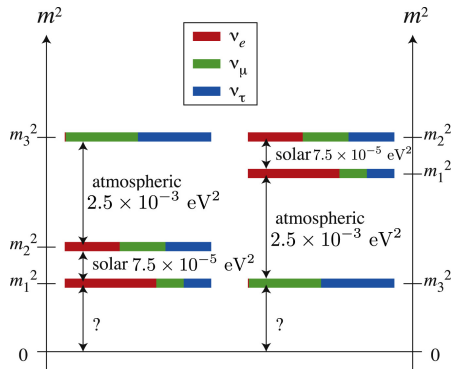


Figure: [King:1701.04413]

# Type-I Seesaw Mechanism

- Extend the SM by a number of right-handed neutrinos -  $\nu_R$
- Give rise to additional terms in the Lagrangian:

$$\mathcal{L}_m^D = -Y_\nu \bar{\ell}_L \tilde{H} \nu_R + h.c. \quad \xrightarrow{\text{EWSB}} \quad -m_D \bar{\nu}_L \nu_R + h.c.$$

$$\mathcal{L}_m^M = -\frac{1}{2} M_R \bar{\nu}_R^c \nu_R + h.c.$$

- Collect such terms  $\implies$  neutrino mass matrix:

$$\begin{pmatrix} \bar{\nu}_L & \bar{\nu}_R^c \end{pmatrix} \begin{pmatrix} 0 & m_D \\ m_D^T & M_R \end{pmatrix} \begin{pmatrix} \nu_L^c \\ \nu_R \end{pmatrix}$$

- In the limit  $M_R \gg m_D$ , diagonalise to obtain light neutrino masses:

$$m_\nu = m_D M_R^{-1} m_D^T = v^2 Y_\nu M_R^{-1} Y_\nu^T$$

- Method of creating baryon asymmetry
- Generate a lepton asymmetry through the decay of lightest  $RH\nu$  into each lepton flavour

$$Y_{\Delta\alpha} = \eta_\alpha \epsilon_\alpha Y_{\nu_1}^{eq}$$

- Lepton asymmetry converted to baryon excess through sphaleron processes in the SM

$$Y_B = \frac{12}{37} \sum_{\alpha} Y_{\Delta\alpha}$$

Extend the SM by **two** RH  $\nu$  singlets:  $\nu_R = \begin{pmatrix} \nu_R^{atm} \\ \nu_R^{sol} \end{pmatrix}$

$$-\mathcal{L}_{LS} = -\mathcal{L}_{SM} + (Y_\nu \bar{\ell}_L \tilde{H} \nu_R + \frac{1}{2} M_R \overline{\nu_R^c} \nu_R + h.c.)$$

N.B.  $Y_\nu$  is a  $2 \times 3$  matrix and  $M_R$  is  $2 \times 2$

*Constrained Sequential Dominance*  $\implies$  **heaviest** RH $\nu$  gives dominant contribution to the **heaviest** LH neutrino mass

Two RH $\nu$ s  $\implies$  lightest LH neutrino is **massless**,  $m_1 = 0$

# The Littlest Seesaw

Fixes the absolute scale of neutrino masses

Constraints on  $\Delta m_{12}$  and  $\Delta m_{13} \rightarrow$  constraints on  $m_2$  and  $m_3$

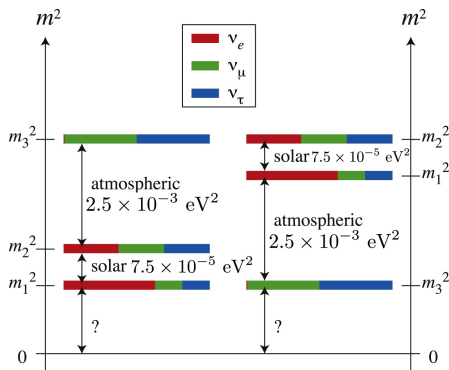


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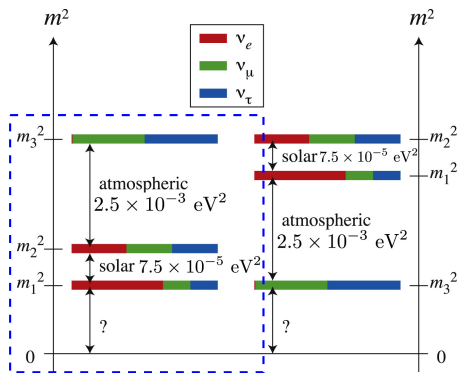


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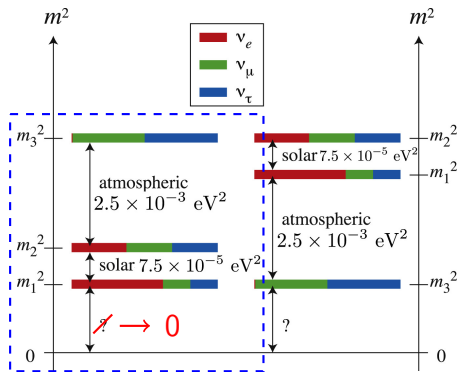


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Four distinct cases of the LS, two of which relevant for leptogenesis.

- Case A:

$$Y_\nu^A = \begin{pmatrix} 0 & be^{i\eta/2} \\ a & nbe^{i\eta/2} \\ a & (n-2)be^{i\eta/2} \end{pmatrix}, \quad M_R^A = \begin{pmatrix} M_{atm} & 0 \\ 0 & M_{sol} \end{pmatrix}$$

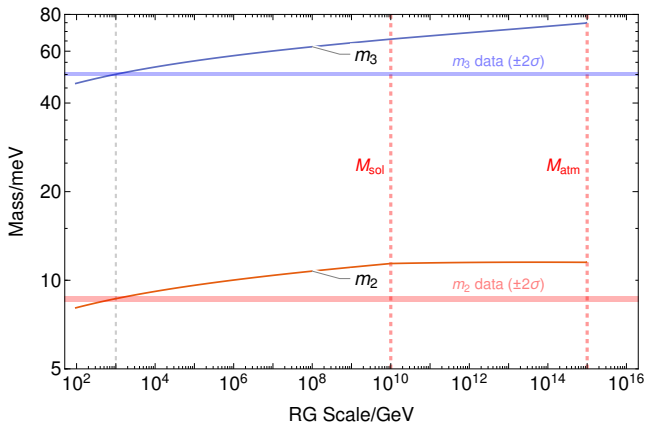
- Case D:

$$Y_\nu^D = \begin{pmatrix} be^{i\eta/2} & 0 \\ (n-2)be^{i\eta/2} & a \\ nbe^{i\eta/2} & a \end{pmatrix}, \quad M_R^D = \begin{pmatrix} M_{sol} & 0 \\ 0 & M_{atm} \end{pmatrix}$$

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- Evolve observables to low scales through **RG running**

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Running of Light  $\nu$  Mass Eigenstates



Case D:

$$M_{atm} = 1.0 \times 10^{15} \text{ GeV}$$

$$M_{sol} = 1.0 \times 10^{10} \text{ GeV}$$

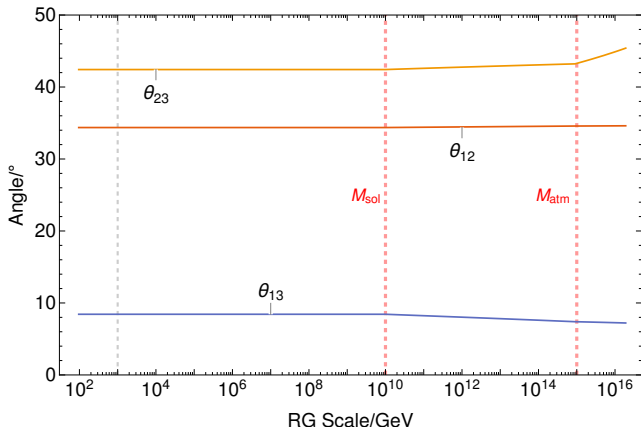
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$$b = 0.00114306$$

$$n = 3$$

$$\eta = -2\pi/3$$

Running of PMNS Angles



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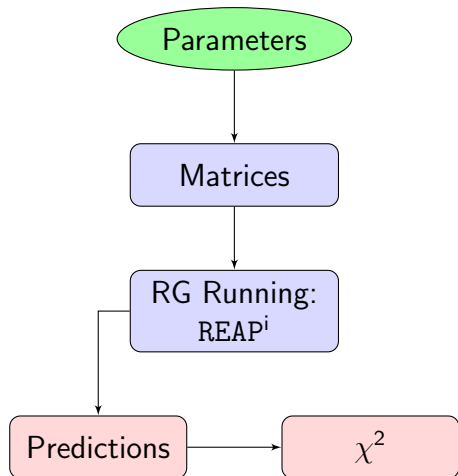
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- Use precision neutrino data and **BAU from Leptogenesis**

<i>Observable</i>	<i>Measured Value</i>
$\sin^2(\theta_{12})$	$0.307^{+0.013}_{-0.012}$
$\sin^2(\theta_{13})$	$0.02206^{+0.00075}_{-0.00075}$
$\sin^2(\theta_{23})$	$0.538^{+0.033}_{-0.069}$
$\delta$	$-126^{\circ+43^{\circ}}_{-31^{\circ}}$
$\Delta m_{12}^2$	$7.40^{+0.21}_{-0.20} \times 10^{-5} \text{eV}^2$
$\Delta m_{13}^2$	$2.494^{+0.033}_{-0.031} \times 10^{-3} \text{eV}^2$
$Y_B$	$0.87^{+0.01}_{-0.01} \times 10^{-10}$

[Esteban et. al.(NuFit 3.2),  
arXiv:1611.01514]

[Plack Collaboration:  
arXiv:1303.5076]



Scan over RH neutrino masses:

$$1.0 \times 10^9 \leq M_1 \leq 5.0 \times 10^{12}$$

$$5M_1 \leq M_2 \leq 1.0 \times 10^{16}$$

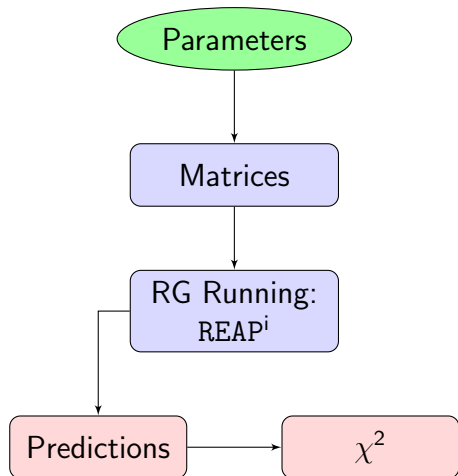
$a$  and  $b$  are left as free parameters.

$n$  fixed to 3 [Geib et. al.:1709.07425]

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<sup>i</sup>Antusch et. al.:hep-ph/0501272.





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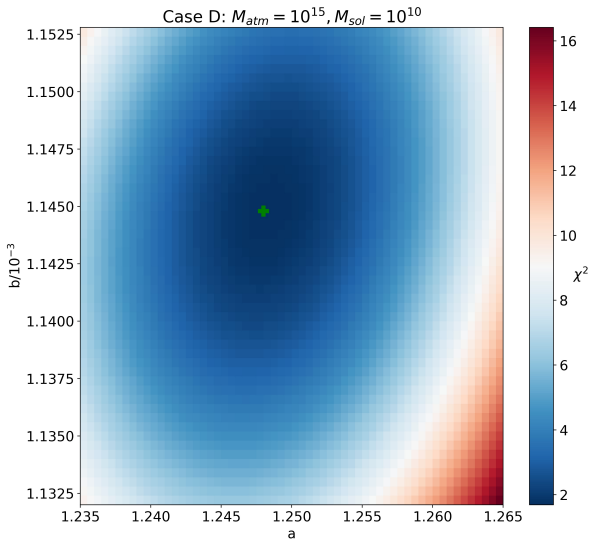
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Determine the best fit point for each case by minimising  $\chi^2$

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- Minimising  $\chi^2$  only useful if there is a **stable global minimum** in the parameter space
- Checked a/b parameter space (fix  $M_{atm}$  and  $M_{sol}$ ) to ensure minimal  $\chi^2$  is stable



Set of current best-fit points after a single coarse scan over  $M_{atm}, M_{sol}, a, b$  (**no leptogenesis**)

	Parameters				$\chi^2$	$\chi_\nu^2$
	$M_{atm}/\text{GeV}$	$M_{sol}/\text{GeV}$	$a$	$b$		
A	$1.0 \times 10^{11}$	$1.0 \times 10^{15}$	0.0113	0.370534	7.31	3.66
D	$1.0 \times 10^{15}$	$1.0 \times 10^{10}$	1.24691	0.00114	<b>1.79</b>	<b>0.90</b>

Concentrate on Case D, look at variation of  $\chi^2$  with RH $\nu$  masses.  
Varying  $M_{atm}$ :

Parameters				$\chi^2$
$M_{atm}/\text{GeV}$	$M_{sol}/\text{GeV}$	$a$	$b$	
$5.0 \times 10^{12}$	$1.0 \times 10^{10}$	0.0782107	0.00108097	10.0119
$1.0 \times 10^{13}$	$1.0 \times 10^{10}$	0.110607	0.00108097	13.0470
$5.0 \times 10^{13}$	$1.0 \times 10^{10}$	0.251836	0.00109801	8.56636
$1.0 \times 10^{14}$	$1.0 \times 10^{10}$	0.35615	0.00109801	11.0418
$5.0 \times 10^{14}$	$1.0 \times 10^{10}$	0.839184	0.00111505	3.50215
$1.0 \times 10^{15}$	$1.0 \times 10^{10}$	0.836253	0.00114306	1.78567

Concentrate on Case D, look at variation of  $\chi^2$  with RH $\nu$  masses.  
Varying  $M_{sol}$ :

Parameters				$\chi^2$
$M_{atm}/\text{GeV}$	$M_{sol}/\text{GeV}$	$a$	$b$	
$1.0 \times 10^{15}$	$1.0 \times 10^{10}$	0.836253	0.00114306	1.78567
$1.0 \times 10^{15}$	$5.0 \times 10^{10}$	0.836253	0.00257733	1.79572
$1.0 \times 10^{15}$	$1.0 \times 10^{11}$	1.24691	0.00367512	2.17086
$1.0 \times 10^{15}$	$5.0 \times 10^{11}$	1.24691	0.00828538	2.03957
$1.0 \times 10^{15}$	$1.0 \times 10^{12}$	1.24691	0.0117173	1.78793
$1.0 \times 10^{15}$	$5.0 \times 10^{12}$	1.24691	0.0264144	1.98070

- Neutrino data seems to be fixing heaviest RH neutrino mass
  - Hints that BAU constraint acting to fix lightest RH neutrino mass
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**Thank you for your attention!**