

Pion Scattering and Lattice QCD

UKQCD

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Plan for this talk

- 1 Why is Lattice QCD interesting & useful?
- 2 Pion Scattering
- 3 Why is any of this useful to you?
- 4 Frontiers in Lattice QCD

Why is low energy physics interesting?

- We know the SM is not the full picture: Dark matter, CP-violation ... we are missing something.
- After a many years of searching we have not detected a signal of BSM Physics.
- "So its time to get the shovel and start digging."-David Newbold (CMS)

However a number of resonances have started to appear in the hadronic sector.

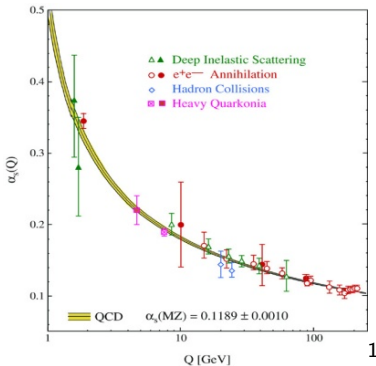


Where do we have large errors?

- Hadronic (QCD) uncertainties are often the dominant errors in theoretical predictions.

Why is low energy difficult?

- Perturbation theory breaks down when applied to QCD at low energies.
- Due to perturbations in the strong force becoming large at low energy. i.e. coupling increases as energy decreases.

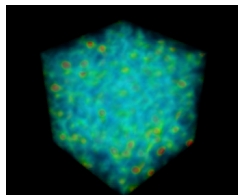


¹<https://arxiv.org/pdf/hep-ex/0606035.pdf>

How does Lattice QCD help?

Fermion action:

$$\mathcal{S}_F = \int d^4x \sum_f \bar{\psi}_f (i\gamma^\mu D_\mu - m_f) \psi_f$$



Lattice Fermion action:

$$\mathcal{S}_F = a^4 \sum_{n \in \Lambda} \sum_f \bar{\psi}_f(n) \left(\gamma_\mu \frac{\psi(n + \hat{\mu}) - \psi(n - \hat{\mu})}{2a} + m \psi_f(n) \right)$$

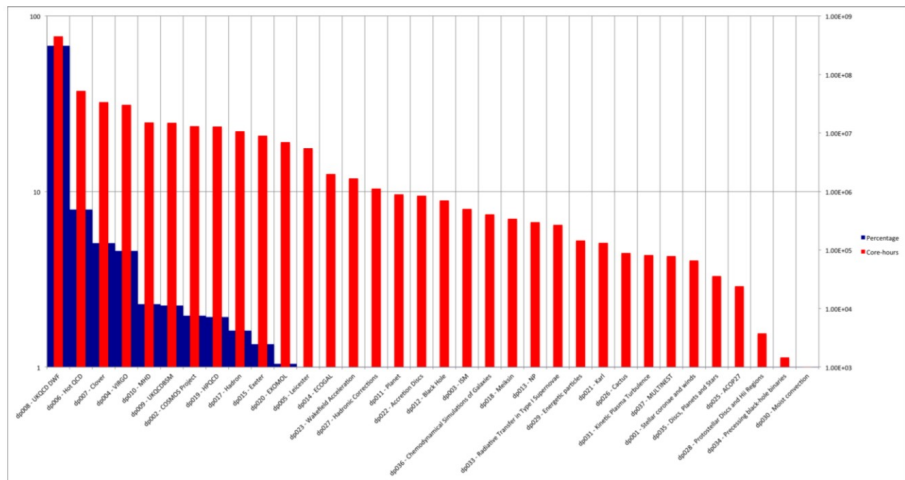
How do you solve a path integral without perturbation theory?

$$\langle \hat{\mathcal{O}} \rangle = \frac{1}{Z} \int \mathcal{D}[U, \psi, \bar{\psi}] e^{-S_{Lat}[U, \psi, \bar{\psi}]} \mathcal{O}[\psi, \bar{\psi}]$$

$$\langle \mathcal{O} \rangle = \lim_{N \rightarrow \infty} \frac{1}{N} \sum_{n=1}^N \mathcal{O}_n$$



Trade Off: Computing time



2

²DiRAC-3 Technical Case

Trade Off: Computationally Expensive

Discretised path integrals can be evaluated by means of Hybrid Monte Carlo methods.

- Theoretically possible.
- Practically hazardous...

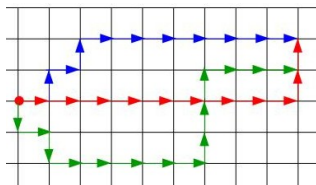
Scale of a propagator:

$$4_{spin} \times 3_{colour} = 12_{dof/site}$$

Lattice sizes vary but $\approx (96_{spacial})^3 \times 192_{time}$

$$12 \times 96^3 \times 192 = 2038431744_{floats}$$

$$(2038431744_{floats} \times 8_{bytes} \times 2_{complex})^2 = Large$$



The propagator is the sum over paths of link variables.

What can we do?

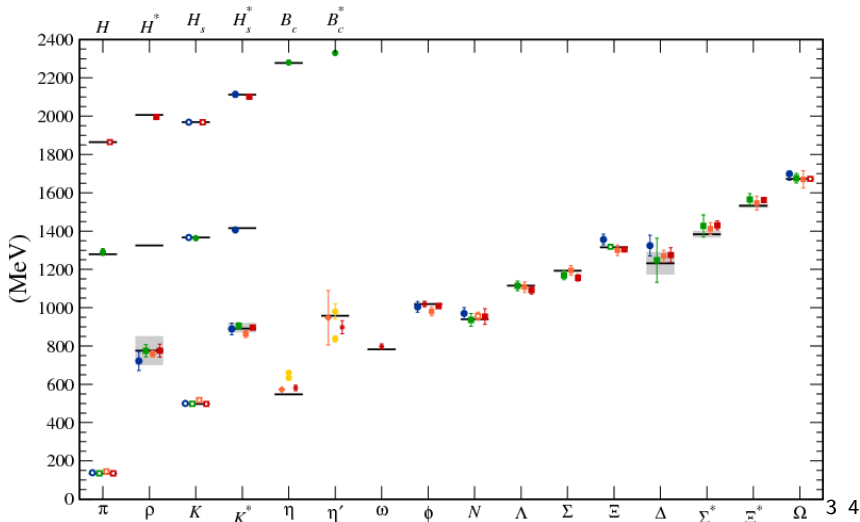
- Simulations of QCD with dynamic (sea) u,d,s,c quarks with physical masses.
- Inverse lattice spacing $a^{-1} \leq 4\text{GeV}$
- Volume $L \leq 6\text{fm}$
- Standard: Hadron Spectroscopy
- Difficult: Interactions of pairs of hadrons.
- "Under construction": multi channel final states.

Parameter tuning

- Tune light quark mass am_l such that: $\frac{am_\pi}{am_p} = \frac{am_\pi^{PDG}}{am_p^{PDG}}$
- Tune strange quark mass: $\frac{am_\pi}{am_K} = \frac{am_\pi^{PDG}}{am_K^{PDG}}$
- Fix lattice spacing $a = \frac{af_\pi}{f_\pi^{PDG}}$

Once these parameters are "tuned" theory is fully predictive! No more input required.

Lattice QCD and the Hadronic spectrum



³b-flavored meson masses are offset by -4000MeV.

⁴<http://inspirehep.net/record/1092992/plots>

Regualisation

IR regulator

- The finite size of the box.

UV regulator

- The inverse lattice spacing.

Discretisation errors

$$\frac{d\psi(x)}{dx} = \lim_{a \rightarrow 0} \frac{\psi(x+a) - \psi(x-a)}{2a} + O(a^2)$$

- We often run a number of simulations of the same physics with different lattice spacings.
- Extrapolation techniques can be used to estimate the continuum limit.
 $a \rightarrow 0$

Finite volume errors

- We have to settle for a reasonable box size and deal with the finite volume errors.

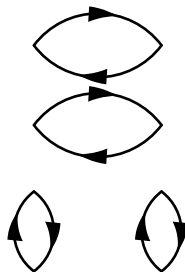
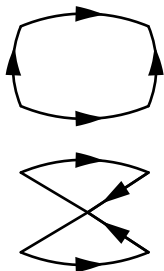
Pion Scattering: $\pi\pi \rightarrow \pi\pi$

- Two possible states $I = 0, 2$

$$|00\rangle = \frac{1}{\sqrt{3}} (|\pi^+\pi^-\rangle - |\pi^0\pi^0\rangle + |\pi^-\pi^+\rangle)$$

$$|22\rangle = |\pi^+\pi^+\rangle$$

- We have to be careful to constrain our simulation so that we stay below the inelastic limit.



Pion Scattering: The Correlator

Correlator of a pion created, propagating and then annihilated:

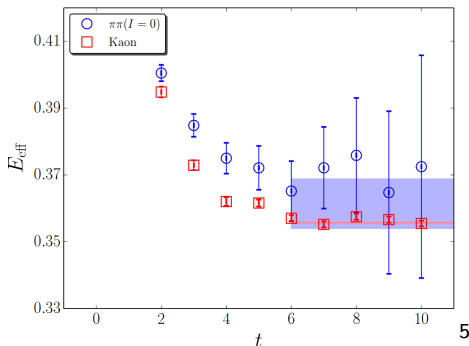
$$\begin{aligned}\langle 0 | \mathcal{O}_\pi(t) \mathcal{O}_\pi(0)^\dagger | 0 \rangle &= \frac{1}{Z} \int \mathcal{D}[U, \psi, \bar{\psi}] \mathcal{O}_\pi(t) \mathcal{O}_\pi(0)^\dagger e^{-S[U, \psi, \bar{\psi}]} \\ \langle 0 | \mathcal{O}_\pi(t) \mathcal{O}_\pi(0)^\dagger | 0 \rangle &= \sum_n \langle 0 | \mathcal{O}_\pi(t) | n \rangle \langle n | \mathcal{O}_\pi(0)^\dagger | 0 \rangle \\ &= \sum_n | \langle 0 | \mathcal{O}_\pi(0) | n \rangle |^2 e^{-E_n t} \\ &= |A_\pi|^2 e^{-E_\pi t} (1 + \mathcal{O}(e^{-\Delta E t}))\end{aligned}$$

What is the form of the pion correlator that we want to fit to our data?

$$C_{\pi\pi} = \langle \mathcal{O}_{\pi\pi}(t) \mathcal{O}_{\pi\pi}(0)^\dagger \rangle = |A_{\pi\pi}|^2 \left(e^{-E_{\pi\pi} t} + e^{-E_{\pi\pi}(T-t)} \right) + K$$

Pion Scattering: Analysis

$$E_{\pi\pi} = \ln\left(\frac{C_{\pi\pi}(t+2) - C_{\pi\pi}(t+1)}{C_{\pi\pi}(t+1) - C_{\pi\pi}(t)}\right)$$



Result: $E_{\pi\pi}^{I=0} \approx 0.36$ in lattice units, $aE_{\pi\pi}^{I=0} = 498\text{MeV}$

⁵ <https://arxiv.org/pdf/1505.07863.pdf>

Pion Scattering: Phase shift

Lüscher quantisation condition:

$$n\pi - \delta(k) = \phi(q)$$

We calculate the pion momentum:

$$k = \sqrt{\frac{E_{\pi\pi}^2}{4} - m_\pi^2}, \quad q = \frac{Lk}{2\pi}$$

We have the inputs required to solve the equation for the phase shift:

$$\delta(k) = -\arctan\left(\frac{-q\pi^{3/2}}{Z_{00}(1; q^2)}\right)$$

The result is for the S wave scattering is:

$$\delta_0(k)_{lat} = 23.8^\circ(4.9^\circ) \quad 6$$

$$\delta_0(k)_{ch.sym/exp} \approx 30 - 35^\circ \quad 7$$

⁶<https://arxiv.org/pdf/1505.07863.pdf>

⁷<https://arxiv.org/pdf/hep-ph/0103088.pdf>

CP Violation and Kaon Decays

The experimental measure of CP violation:

$$\text{Re}\left(\frac{\epsilon'}{\epsilon}\right) = \text{Re}\left(\frac{i\omega e^{i(\delta_2 - \delta_0)}}{\sqrt{2}\epsilon} \left[\frac{\text{Im}A_2}{\text{Re}A_2} - \frac{\text{Im}A_0}{\text{Re}A_0} \right]\right)$$

This is dependent on the pion scattering phases δ_0 we have just evaluated.

$$\text{Re}(\epsilon'/\epsilon)_{\text{Lat}} = 1.38(5.15)(4.59) \times 10^{-4} \quad 8$$

$$\text{Re}(\epsilon'/\epsilon)_{\text{Exp}} = 16.6(2.3) \times 10^{-4}$$

Improve by increasing statistics:

$$\frac{1}{\sqrt{N}} = \frac{1}{\sqrt{216}} = \frac{1}{14.7} = 0.068$$

$$\frac{1}{\sqrt{N}} = \frac{1}{\sqrt{1000}} = \frac{1}{31.6} = 0.032$$

⁸<https://arxiv.org/pdf/1505.07863.pdf>

- QED effects
 - QED not currently included due to finite volume effects. Photons interact over long range.
 - Large volumes not currently available, so use effective theory to subtract finite volume effect.
 - Analytically compute the difference between the finite volume and infinite volume self energies.
 - QED in meson decay: Assume photons are soft and do not resolve the hadronic structure.
- Isospin breaking effects
 - Currently assume $d = u = l$ and have equal mass.

- Lattice QCD is a method for making predictions of the SM in a low energy regime.
- Introducing previously challenging effects and using increased statistics we can drive down errors
- This means we can tightly constrain the Standard Model and hopefully break it in the hadronic sector.

Thank you

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Thank you

Back up Slides

Mesons on the Lattice

Now we have a way of calculating the correlators we need operators:

$$\mathcal{O}_M(n) = \bar{\psi}^{(f_1)}(n)\Gamma\psi(n)^{(f_2)}$$

Operator product of a meson propagating from a point n to m :

(supressed intergal and indices)

$$\begin{aligned}\langle \mathcal{O}_T(n)\bar{\mathcal{O}}_T(m) \rangle &= \langle \bar{d}(n)\Gamma u(n)\bar{u}(m)\Gamma d(m) \rangle \\ &= \Gamma\Gamma \langle \bar{d}(n)u(n)\bar{u}(m)d(m) \rangle \\ &= \Gamma\Gamma \langle u(n)\bar{u}(m) \rangle \langle d(m)\bar{d}(n) \rangle \\ &= \Gamma\Gamma D_u^{-1}(n|m)D_d^{-1}(n|m) \\ &= \text{tr}[\Gamma D_u^{-1}(n|m)\Gamma D_d^{-1}(n|m)]\end{aligned}$$

Final form of a hadron correlator:

$$\langle \mathcal{O}_\pi(n)\bar{\mathcal{O}}_\pi(m) \rangle = \frac{1}{Z} \int \mathcal{D}e^{-S_{Lat}} \det[D_u] \det[D_d] \text{tr}[\Gamma D_u^{-1}(n|m)\Gamma D_d^{-1}(n|m)]$$

The form of the zeta function is:

$$Z_{00}(s; q^2) = \frac{1}{\sqrt{4\pi}} \sum_{\vec{n} \in \mathbb{Z}^3} (\vec{n}^2 - q^2)^{-s}$$

Phase shift of Pion Scattering

Luscher quantisation condition:

$$n\pi - \delta(k) = \phi(q)$$

we relate k and q in the following way:

$$q = \frac{Lk}{2\pi}$$

We calculate the pion momentum:

$$k = \sqrt{\frac{E_{\pi\pi}^2}{4} - m_\pi^2}$$

but we don't calculate the mass we have E_π . To calculate the mass we can use the relation:

$$m_\pi = \sqrt{E_\pi^2 - 3\left(\frac{\pi}{L}\right)^2}$$

$$E_{\pi\pi} = \ln\left(\frac{C_{\pi\pi}(t)}{C_{\pi\pi}(t+1)}\right) \quad (1)$$

$$E_{\pi\pi} = \ln\left(\frac{C_{\pi\pi}(t+2) - C_{\pi\pi}(t+1)}{C_{\pi\pi}(t+1) - C_{\pi\pi}(t)}\right) \quad (2)$$

Pion Scattering: Analysis Schematic

- 1 Raw data from simulation output.
- 2 Form of the data is a $C \times T$ matrix (Configurations, Timeslices)
- 3 Perform Jackknife or Bootstrap Resampling.
- 4 Calculate the covariance matrix.
- 5 Fit functional form of the correlator to data.
- 6 The result is one value for each of the fitted free parameters this is done for each configuration.
- 7 Average to get central value and take standard deviation to get the error.