



THE ELECTROWEAK PHASE TRANSITION, BEYOND THE STANDARD MODEL

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Based (in part) on:

**arXiv: 1711.09849
JHEP1703 (2017) 007**

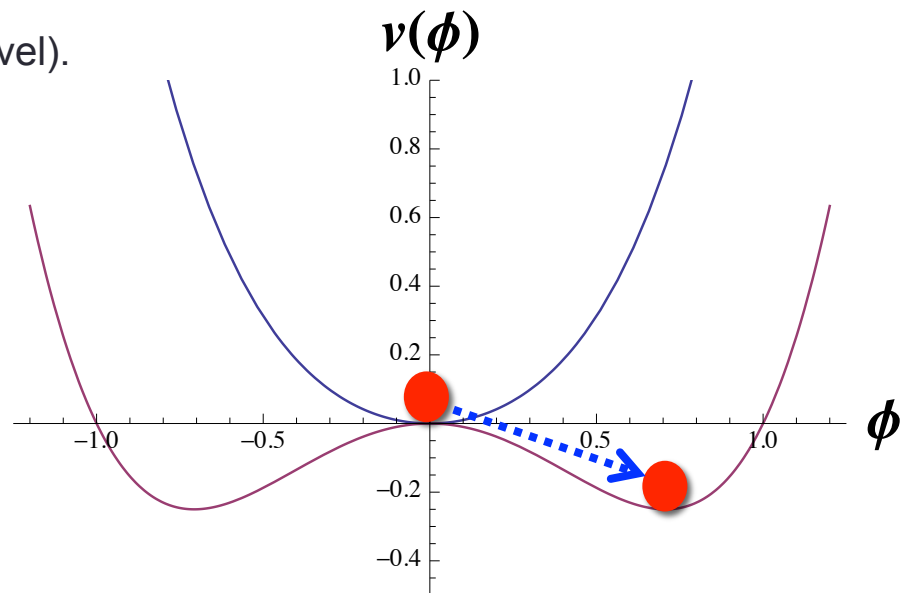
Electroweak phases

- The Higgs (effective) potential has a global minimum...
 - At "low" T : non-zero field value, $v = 246$ GeV.
 - Fermions have mass.
 - Gauge boson W, Z have mass.
 - Higgs has mass.
 - At "high" T : zero field value, $v = 0$.
 - Fermions are massless (at tree level).
 - Gauge bosons are massless (at tree level).

Phase transition at some "intermediate" T .

Context:

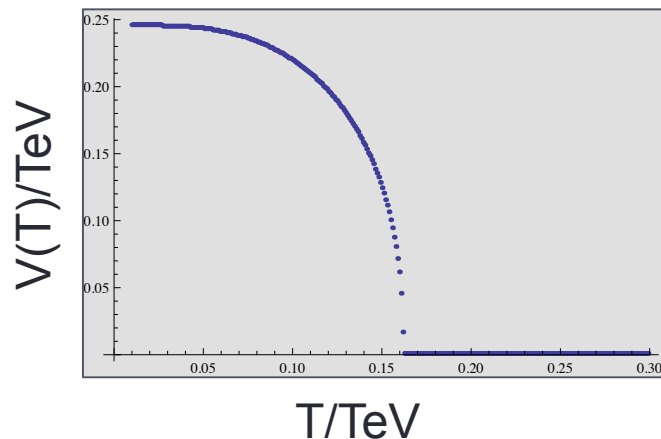
After the Big Bang \rightarrow "high" T .
 Expansion \rightarrow cool down \rightarrow "low" T .
 Phase transition in time, $T(t)$!



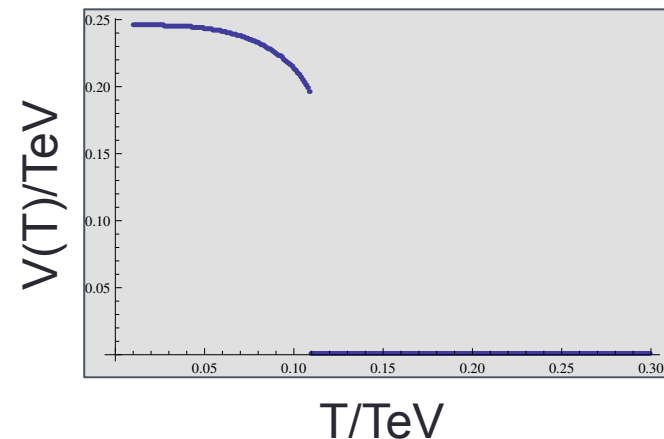
Finite temperature phase transitions

- The thermodynamic equilibrium system is described by a finite temperature effective potential (the free energy).
- A function of an "order" parameter.
- The value of the order parameter corresponds to the minimum of the free energy.

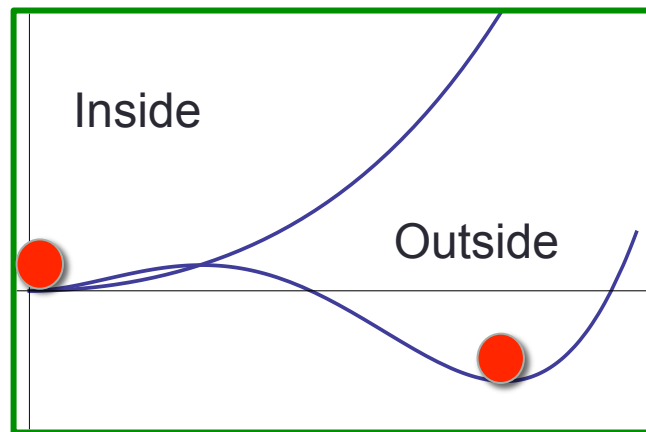
If the minimum evolves continuously as a function of $T \rightarrow 2.$ order phase transition. (Or higher...)



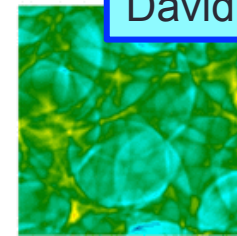
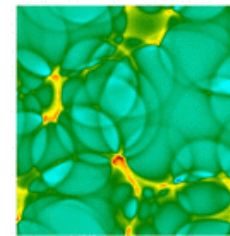
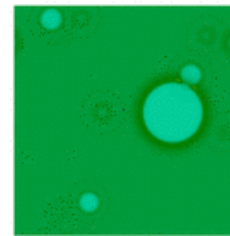
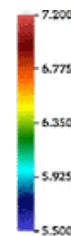
If the minimum evolves discontinuously as a function of $T \rightarrow 1.$ order phase transition.



What do I need a first order transition for?



- If 1. order \rightarrow nucleation of vacuum bubbles.
- Outside: high-T phase, $v = 0$.
 - Inside: low-T phase, $v > 0$.
 - Bubbles expand at up to the speed of light.



David Weir

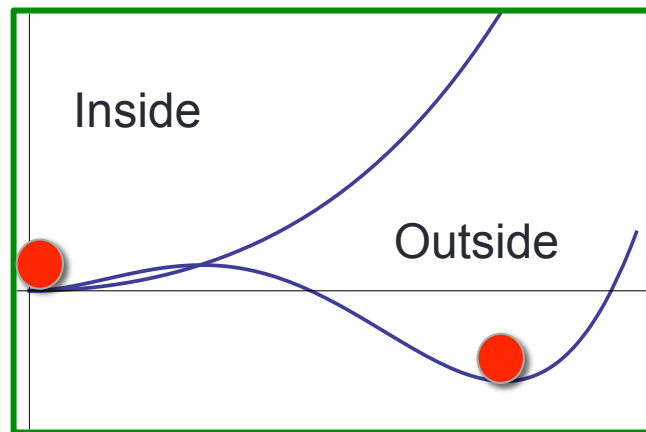
If particle physics model has:

- C- and CP-violation (like in the Standard Model)
 - Baryon number violation (like in the Standard Model)
- \rightarrow A baryon asymmetry is created as the bubble walls sweep through the plasma.

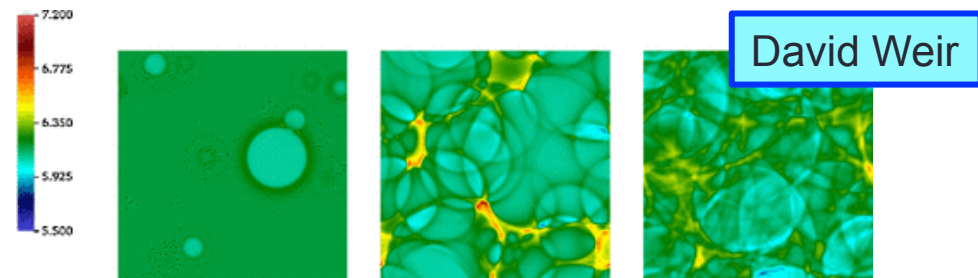


Electroweak
Baryogenesis

What do I need a first order transition for?

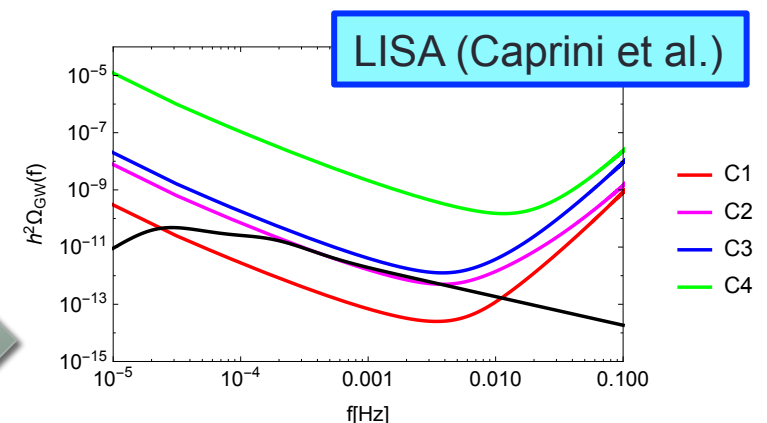


- If 1. order \rightarrow nucleation of vacuum bubbles.
- Outside: high-T phase, $v = 0$.
 - Inside: low-T phase, $v > 0$.
 - Bubbles expand at up to the speed of light.



- As the bubble walls push around the plasma
- \rightarrow gravitational waves are created.
- As the bubble walls collide
- \rightarrow gravitational waves are created.

Gravitational
waves



How to compute it...perturbatively?

$$V_{\text{eff}}(\phi, T) = V_{\text{tree}}(\phi) + V_{\text{CW}}(\phi) + V_{\text{c.t.}}(\phi) + V_T^1(\phi, T) + V_T^2(\phi, T)$$

$$V_T^1(\phi, T) = \frac{T^4}{2\pi^2} \sum_i N_i \int_0^\infty dx x^2 \log \left[1 \pm e^{-\sqrt{x^2 + \frac{M_i^2(\phi)}{T^2}}} \right]$$

$$V_T^2(\phi, T) = \frac{T}{12\pi} \sum_i N_i \text{Tr} [M_i^3(\phi) - M_i^3(\phi, T)]$$

Procedure:

Scanning through your parameter space:

- Compute the 1-loop, daisy-resummed, finite temperature effective potential $V(\Phi, T)$.
- Find the minimum.
- Plot the minimum as a function of T .
- Does it go to zero continuously as T increases?
 - Yes? No bubbles \rightarrow Game Over.
 - No? Bubbles \rightarrow How big is the jump?

Lots and lots of people:

Carrington, Kainulainen, Cline, Huber, Seniuch, Konstandin, Servant, Kozaczuk, Laine, Nardini, Ramsey-Musolf, Damgaard, AT, Petersen, O'Connell, Haarr, Kvellestad, Dorsch, No, Harmann, Mimasu, Profumo, Shaughnessy, ...

How to compute it...non-perturbatively?

- Lattice Monte Carlo in full 4D (expensive).
- Lattice Monte Carlo in effective 3D model (less expensive).

Procedure:

Scanning through your parameter space:

- Analytically, match your 4D theory to a 3D effective theory.
- In the 3D theory, numerically compute the full, nonperturbative effective potential as a function of T .
- Find the minimum.
- Plot the minimum as a function of T .
- Does it go to zero continuously as T increases?
 - Yes? No bubbles → Game Over.
 - No? Bubbles → How big is the jump?

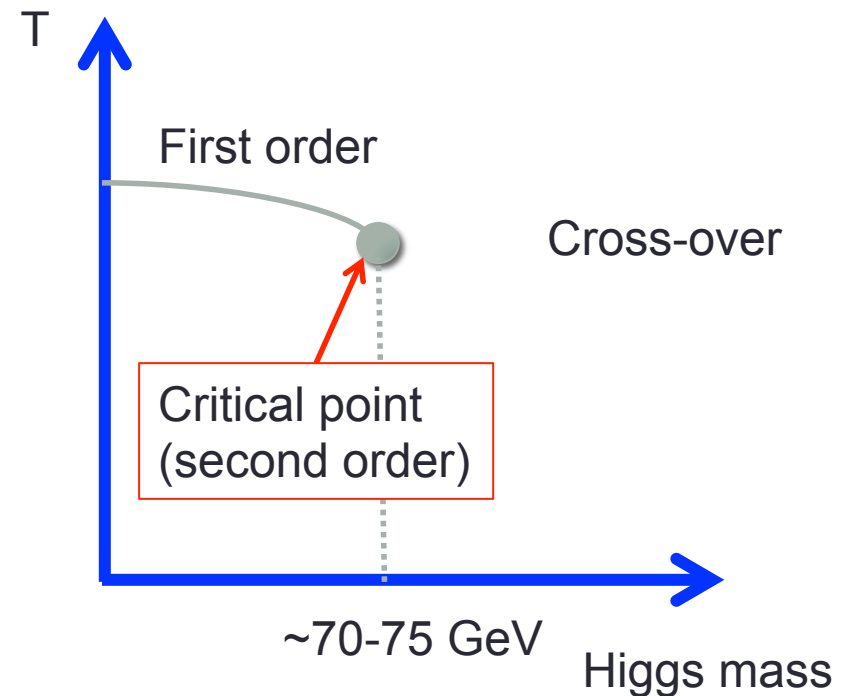
Quite few people:

Appelquist, Pisarski, Kajantie, Laine, Rummukainen, Shaposhnikov, Farakos, Csikor, Fodor, Heitger.

The phase transition in the Standard Model

Higgs mass unknown in 1998.
Phase diagram and T_c as a
function of m_H .

4D-endpoint $\rightarrow 72$ GeV.
3D-endpoint $\rightarrow \sim 80$ GeV.



Higgs mass = 125 GeV!
No first order transition.

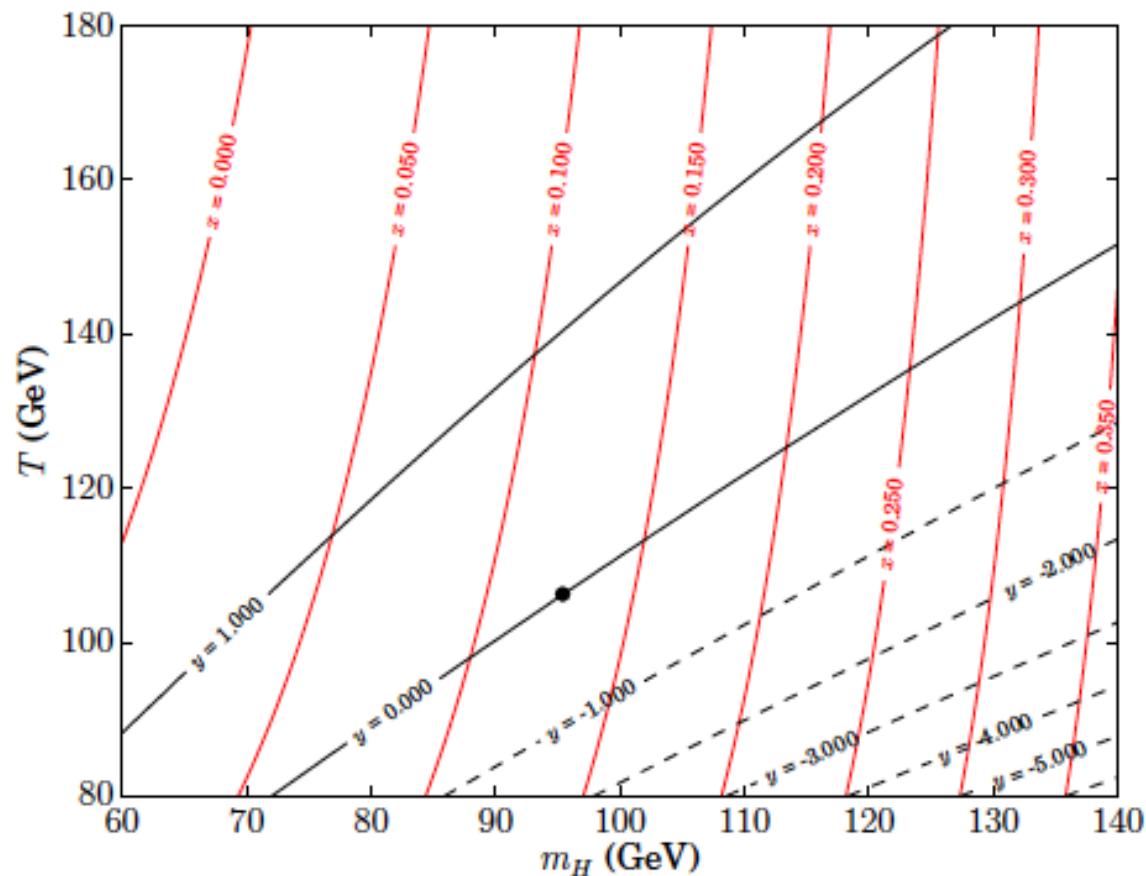
\rightarrow No bubbles.
 \rightarrow Game over!

Our existence \rightarrow proof of BSM physics.

Kajantie, Laine, Rummukainen,
Shaposhnikov: 1995-8
Csikor, Fodor, Heitger: 1998

The phase transition in the Standard Model

3D-model:
$$S = \int d^3x \left[\frac{1}{4} F_{ij}^a F^{ij,a} + (D_i \phi)^\dagger D^i \phi + m_3^2 \phi^\dagger \phi + \lambda_3 (\phi^\dagger \phi)^2 \right]$$



Redone at 1-loop level.
Different DimRed procedure.

$$x = \frac{\lambda_3}{g_3^2},$$

$$y = \frac{m_3^2}{g_3^4}$$

Phase transition at $y \sim 0$.
1. Order at $x < x_c = 1/8$.

Brauner, Tenkanen, AT,
Vuorinen, Weir: 2017.

BSM: Add a scalar

Existence of matter requires BSM physics!

If electroweak baryogenesis is responsible, need 1. order phase transition.
 → Add new degrees of freedom: Try with a scalar coupled to the Higgs.

SM

$$\begin{aligned}
 &+ \frac{1}{2} \partial_\mu \sigma \partial^\mu \sigma + \frac{1}{2} \mu_\sigma^2 \sigma^2 \\
 &+ \mu_1 \sigma + \frac{1}{3} \mu_3 \sigma^3 + \frac{1}{4} \lambda_\sigma \sigma^4 \\
 &+ \frac{1}{2} \mu_m \sigma \phi^\dagger \phi + \frac{1}{2} \lambda_m \sigma^2 \phi^\dagger \phi
 \end{aligned}$$

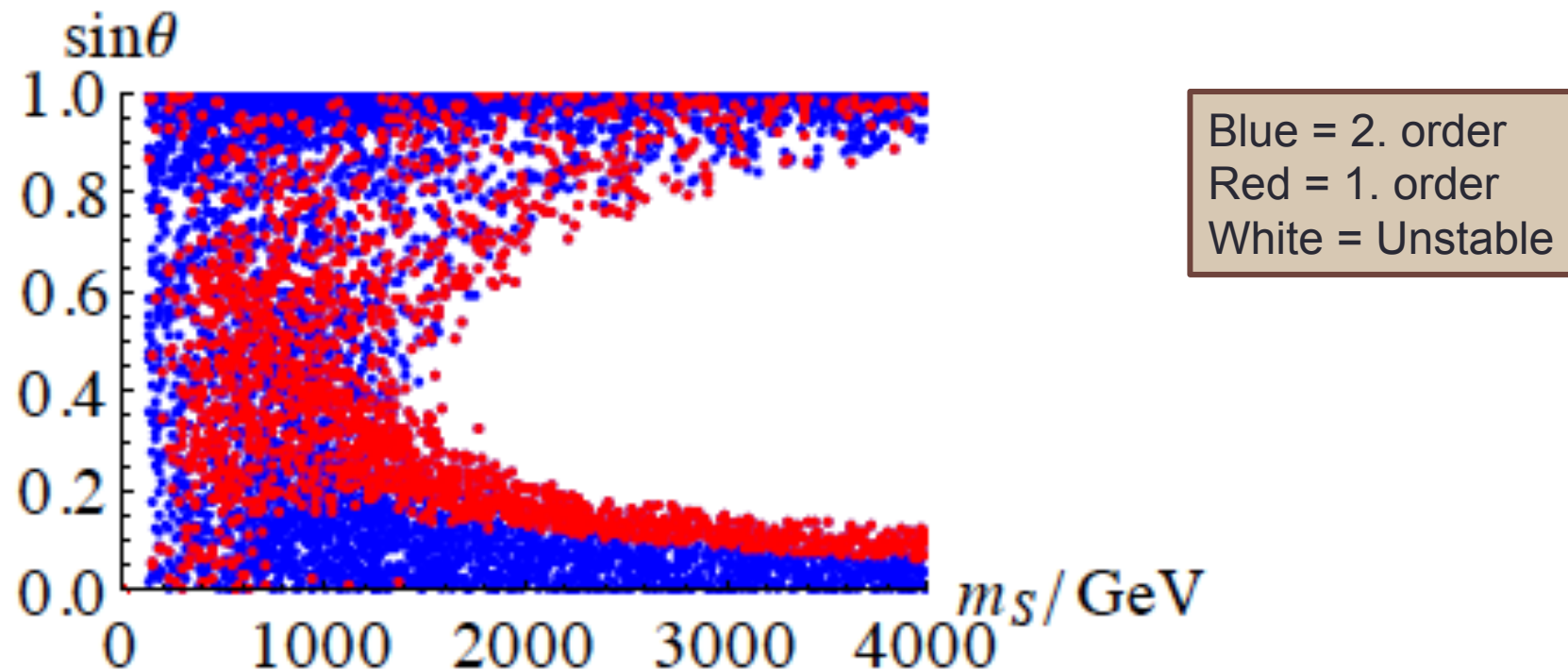
Singlet

SM – Higgs

$$\begin{aligned}
 &+ (D_\mu \phi_1)^\dagger D^\mu \phi_1 + (D_\mu \phi_2)^\dagger D^\mu \phi_2 + \\
 &\frac{1}{2} m_{11}^2 \phi_1^\dagger \phi_1 + \frac{1}{2} m_{22}^2 \phi_2^\dagger \phi_2 + \\
 &\frac{1}{2} m_{12}^2 \phi_1^\dagger \phi_2 + \text{h.c.} + \\
 &\frac{1}{2} \lambda_1 (\phi_1^\dagger \phi_1)^2 + \frac{1}{2} \lambda_2 (\phi_2^\dagger \phi_2)^2 + \\
 &\lambda_3 (\phi_1^\dagger \phi_1) (\phi_2^\dagger \phi_2) + \lambda_4 (\phi_1^\dagger \phi_2) (\phi_2^\dagger \phi_1) + \\
 &\frac{1}{2} \lambda_5 (\phi_1^\dagger \phi_2)^2 + \text{h.c.}
 \end{aligned}$$

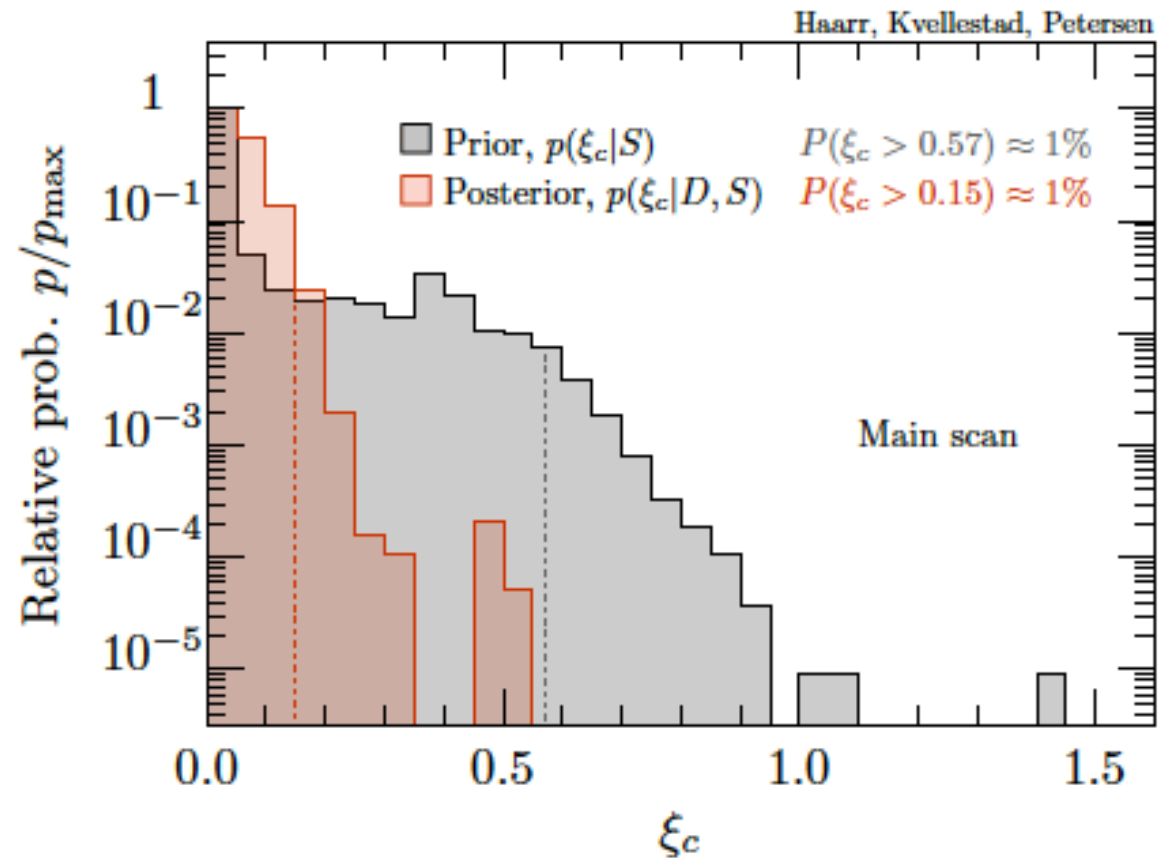
Doublet

Perturbatively: SM + singlet



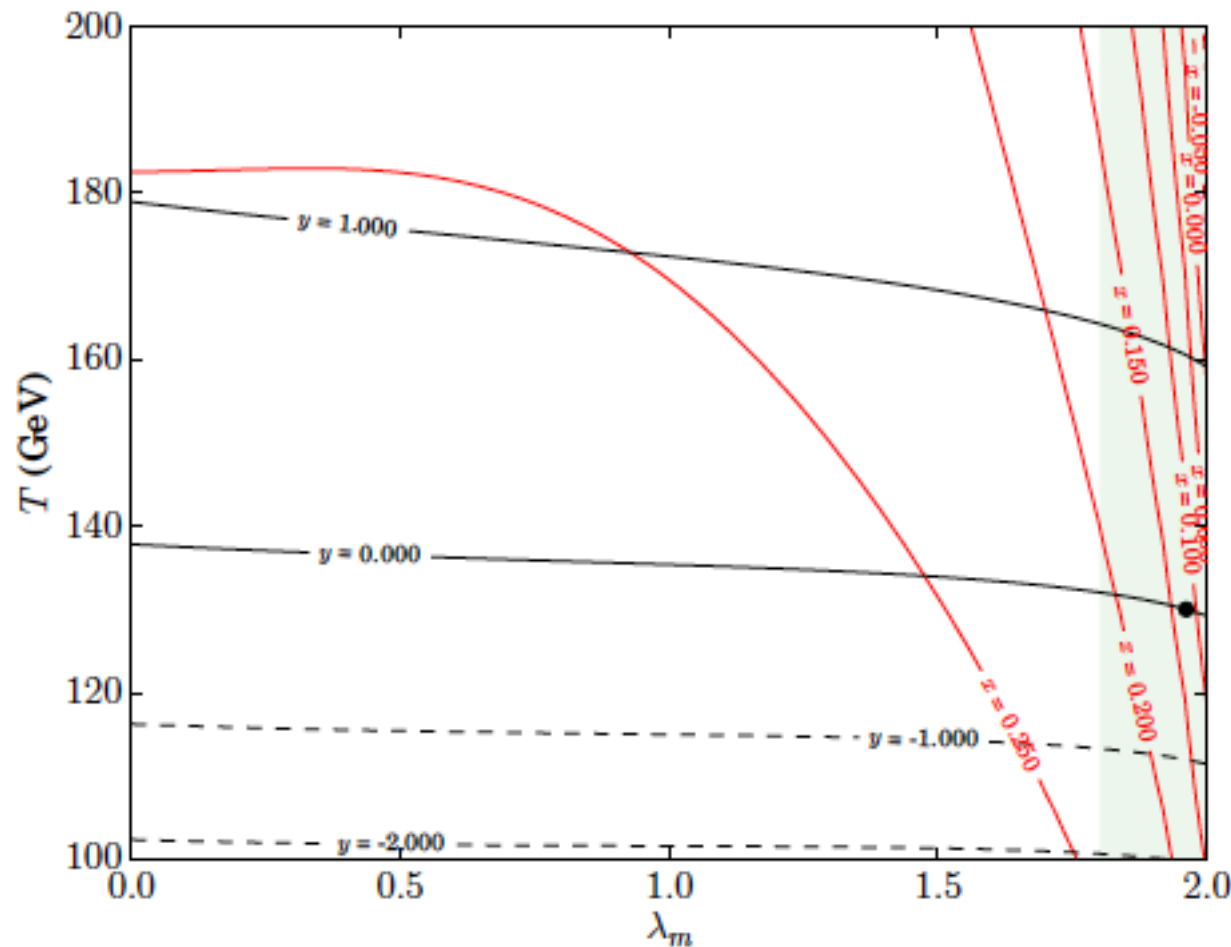
Absolutely biased selection: Example only!
Damgaard, Haarr, O'Connell, AT: 2015.

Perturbatively: SM + Doublet



Absolutely biased selection: Example only!
 Haarr, Kvellestad, Petersen: 2017.

Non-perturbatively: SM + singlet



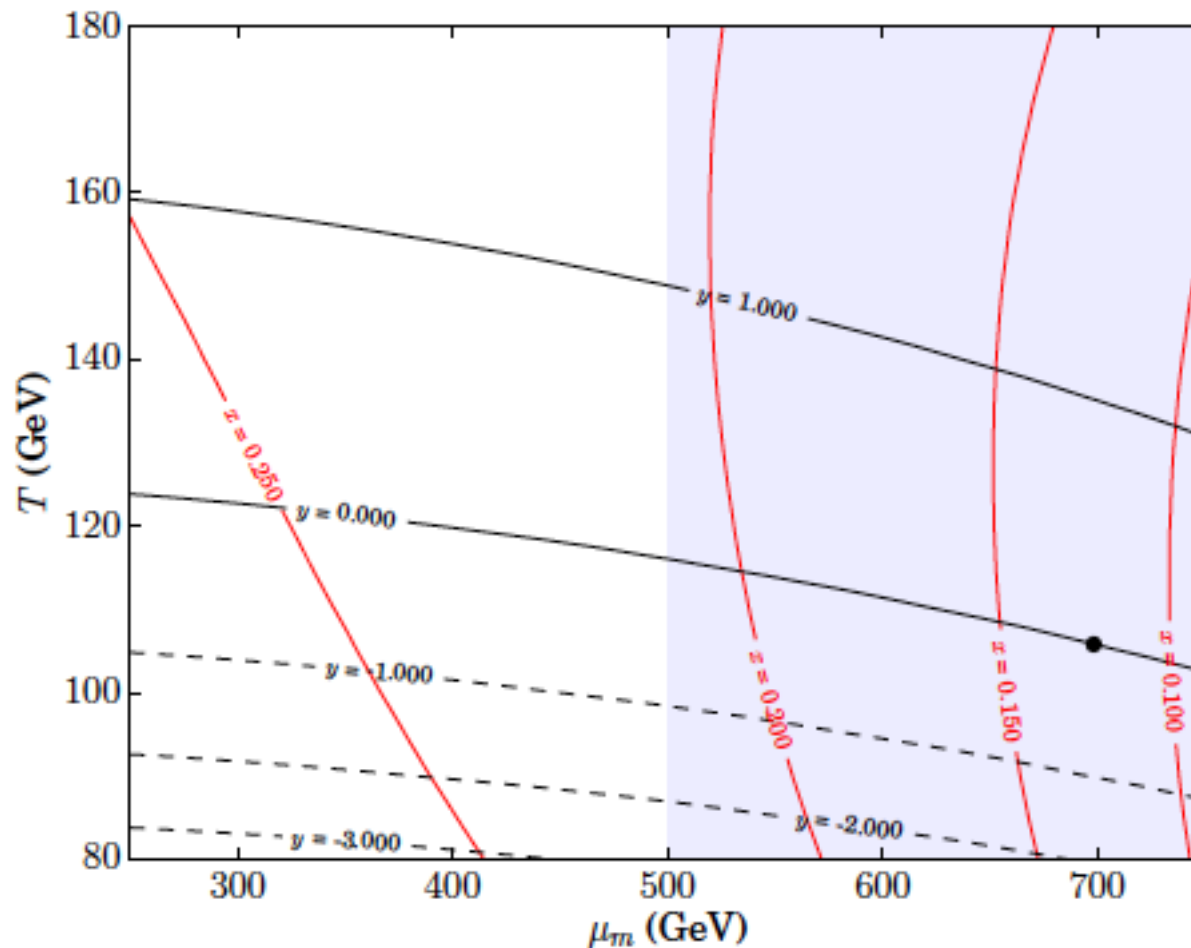
Singlet mass: 500 GeV

SM

$$\begin{aligned}
 & + \frac{1}{2} \partial_\mu \sigma \partial^\mu \sigma + \frac{1}{2} \mu_\sigma^2 \sigma^2 \\
 & + \cancel{\mu_1 \sigma} + \cancel{\frac{1}{3} \mu_3 \sigma^3} + \frac{1}{4} \lambda_\sigma \sigma^4 \\
 & + \cancel{\frac{1}{2} \mu_m \sigma \phi^\dagger \phi} + \frac{1}{2} \lambda_m \sigma^2 \phi^\dagger \phi
 \end{aligned}$$

Brauner, Tenkanen, AT, Vuorinen, Weir: 2017

Non-perturbatively: SM + singlet

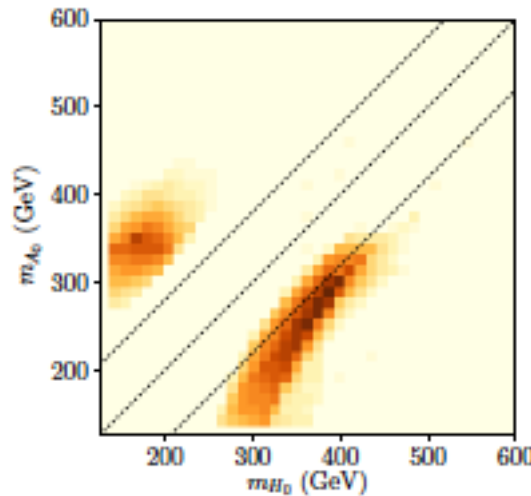
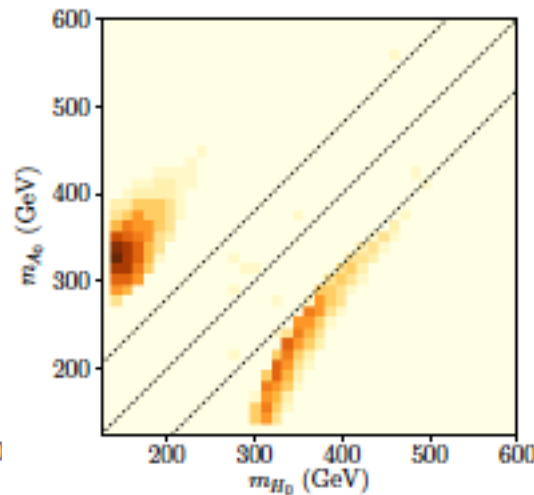
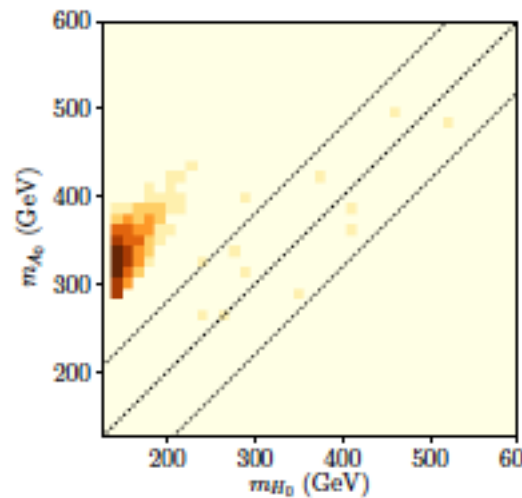
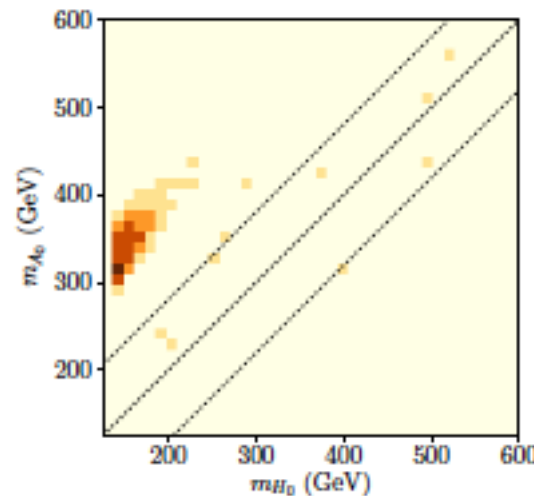


Singlet mass: 250 GeV

SM

$$\begin{aligned}
 & + \frac{1}{2} \partial_\mu \sigma \partial^\mu \sigma + \frac{1}{2} \mu_\sigma^2 \sigma^2 \\
 & + \cancel{\mu_1 \sigma} + \cancel{\frac{1}{3} \mu_3 \sigma^3} + \frac{1}{4} \lambda_\sigma \sigma^4 \\
 & + \frac{1}{2} \mu_m \sigma \phi^\dagger \phi + \frac{1}{2} \cancel{\lambda_\sigma \sigma^2 \phi^\dagger \phi}
 \end{aligned}$$

Non-perturbatively: SM + doublet

(a) $\tan(\beta) = 1.5$ (b) $\tan(\beta) = 2.0$ (c) $\tan(\beta) = 2.5$ (d) $\tan(\beta) = 3.0$

SM – Higgs

$$\begin{aligned}
 &+(D_\mu \phi_1)^\dagger D^\mu \phi_1 + (D_\mu \phi_2)^\dagger D^\mu \phi_2 + \\
 &\frac{1}{2}m_{11}^2 \phi_1^\dagger \phi_1 + \frac{1}{2}m_{22}^2 \phi_2^\dagger \phi_2 + \\
 &\frac{1}{2}m_{12}^2 \phi_1^\dagger \phi_2 + \text{h.c.} + \\
 &\frac{1}{2}\lambda_1(\phi_1^\dagger \phi_1)^2 + \frac{1}{2}\lambda_2(\phi_2^\dagger \phi_2)^2 + \\
 &\lambda_3(\phi_1^\dagger \phi_1)(\phi_2^\dagger \phi_2) + \lambda_4(\phi_1^\dagger \phi_2)(\phi_2^\dagger \phi_1) + \\
 &\frac{1}{2}\lambda_5(\phi_1^\dagger \phi_2)^2 + \text{h.c.}
 \end{aligned}$$

$$m_{H^\pm} = m_{A_0},$$

$$\tan(\beta) = \frac{v_2}{v_1}, \quad v_1^2 + v_2^2 = (246 \text{ GeV})^2$$

$$\cos(\beta - \alpha) = 0,$$

$$\mu_{12} < 400 \text{ GeV},$$

All real \rightarrow no CP-violation

Andersen, Gorda, Niemi, Weir,
Tenkanen, AT, Vuorinen: 2017

Outlook

A first order electroweak phase transition would provide:

1. A viable baryogenesis mechanism.
2. A detectable cosmological GW source.

Where the SM model fails, extended Higgs sectors may succeed:

Collider constraints are complementary:

- LHC suggests extensions give small corrections to SM.
- A 1. order phase transition requires large corrections to SM.

Combine theoretical computation, collider constraints, BG and measurements of GW to rule in or out TeV-scale extensions of the SM.

Perturbative computations allow broad parameter scans.

But non-perturbatively, transition may be weaker or go away!

To do: In parallel compute:

- Numerically, phase diagrams of more general 3D theories.
- Analytically, matching relations from 4D theories to those 3D theories.

Thank you!