

MODELIZATION OF RADIATION-INDUCED DAMAGE IN FLUKA AND MATERIAL DAMAGE ESTIMATES FOR CERN INJECTORS AND FUTURE FACILITIES

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FLUKA



FLUKA is a Monte Carlo code for calculations of particle transport and interactions with matter.

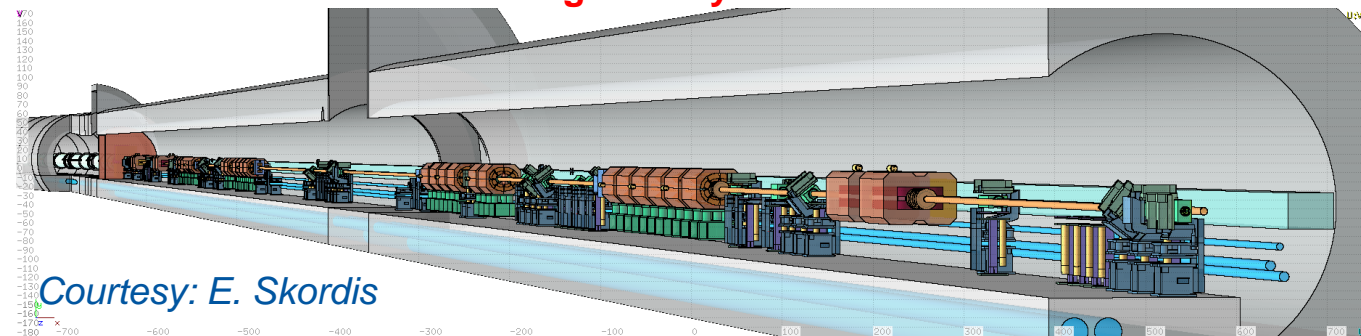
- Wide range of applications:

Proton and electron accelerator shielding, target design, calorimetry, activation, dosimetry, detector design, Accelerator Driven Systems, cosmic rays, neutrino physics, Radiotherapy, etc.

- **Extensively used at CERN for:**

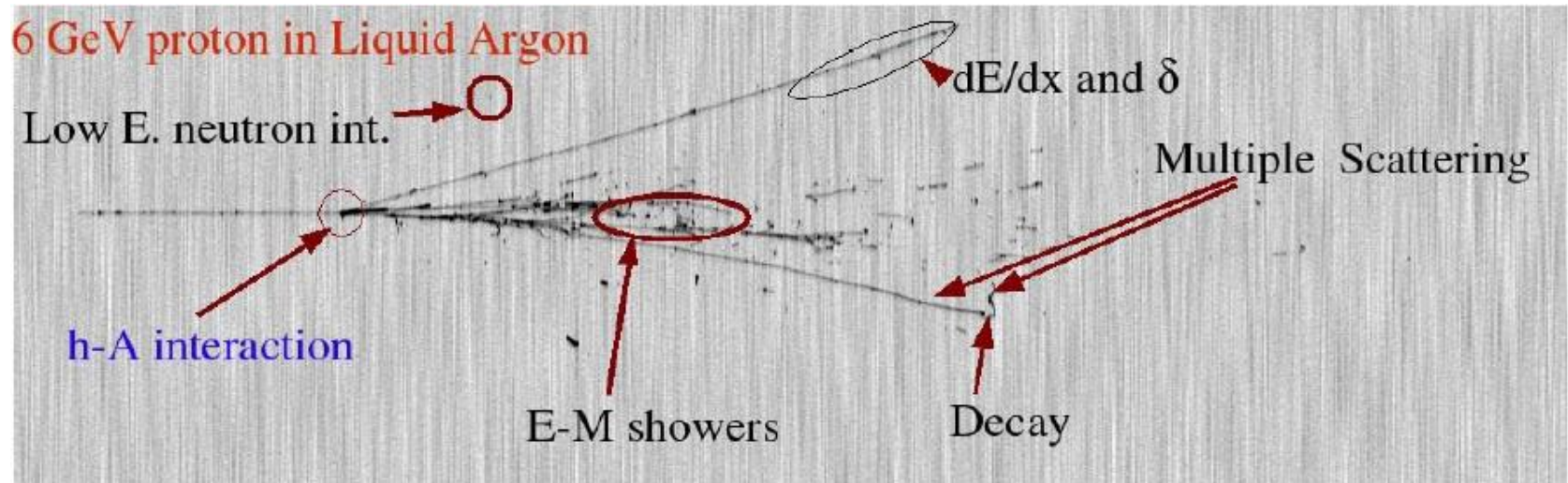
- Beam-machine interactions
- Radio-Protection calculations
- Facility design of future projects

FLUKA geometry of the LHC warm section of IR7



Courtesy: E. Skordis

Interaction and Transport Monte Carlo Code



- Hadron-nucleus interactions
- Nucleus-Nucleus interactions
- Electron interactions
- Photon interactions
- Muon interactions (inc. photonuclear)
- Neutrino interactions
- Decay
- Low energy neutrons
- Ionization
- Multiple scattering
- Combinatorial geometry
- Voxel geometry
- Magnetic field
- Analogue or biased
- On-line buildup and evolution of induced radioactivity and dose
- User-friendly GUI thanks to *Flair*

Info: <http://www.fluka.org>

Different kinds of damage

from Radiation-Matter interaction

Precious materials (healthy/tragic damage)

energy (dose) deposition, radioisotope production and decay & positron annihilation and photon pair detection

Oxidation

by generation of chemically active radicals (e.g. PVC de-hydrochlorination by X and g-rays, radiolysis,...)

Accidents

energy (power) deposition

Degradation

energy (dose) deposition, particle fluence, **DPA**

Gas production

residual nuclei production

Electronics

high energy hadron fluence, neutron fluence, energy (dose) deposition

Activation

residual activity and dose rate

F.Cerutti

DPA as Radiation Damage Estimator

The **dpa** quantity is a measure of the amount of radiation damage in irradiated materials. It means the **average number of displacements that every atom in the crystal structure has suffered.**

$$dpa = \frac{A}{N_A \rho} N_F$$

A is the mass number

N_A is the Avogadro number

ρ is the density

N_F is the number of defects or **Frenkel pairs**.

Experimentally: no direct determination. Indirect through study of macroscopic effects (electric and thermal conductivities, radiation hardening, swelling...)

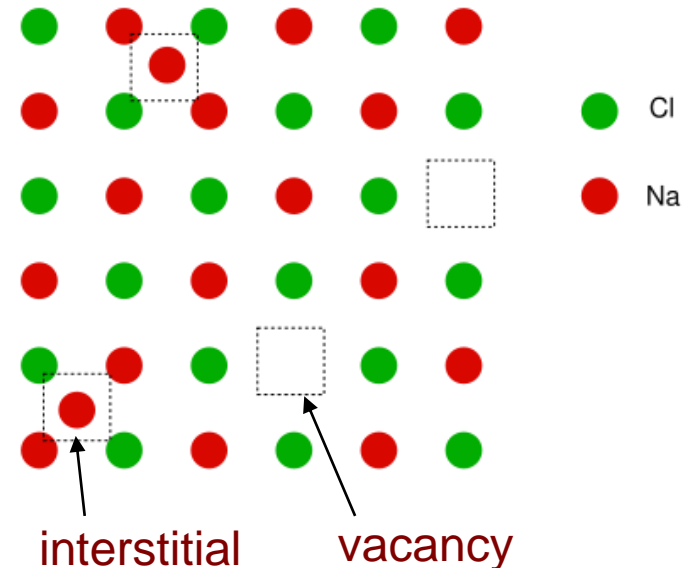
Amount of dpa \longleftrightarrow Macroscopic effects

Frenkel pairs

- Frenkel pair N_F (defect or disorder), is a compound crystallographic defect formed when an atom or ion leaves its place in the lattice (leaving a **vacancy**), and lodges nearby in the crystal (becoming an **interstitial**)

$$N_{NRT} \equiv N_F = \kappa \frac{\xi(T)T}{2E_{th}}$$

- N_{NRT} Defects by Norgert, Robinson and Torrens
- $\kappa=0.8$ is the displacement efficiency
- T kinetic energy of the primary knock-on atom (PKA)
- $\xi(T)$ partition function (LSS theory)
- $\xi(T) T$ directly related to the **NIEL**(non ionizing energy loss)
- E_{th} damage threshold energy

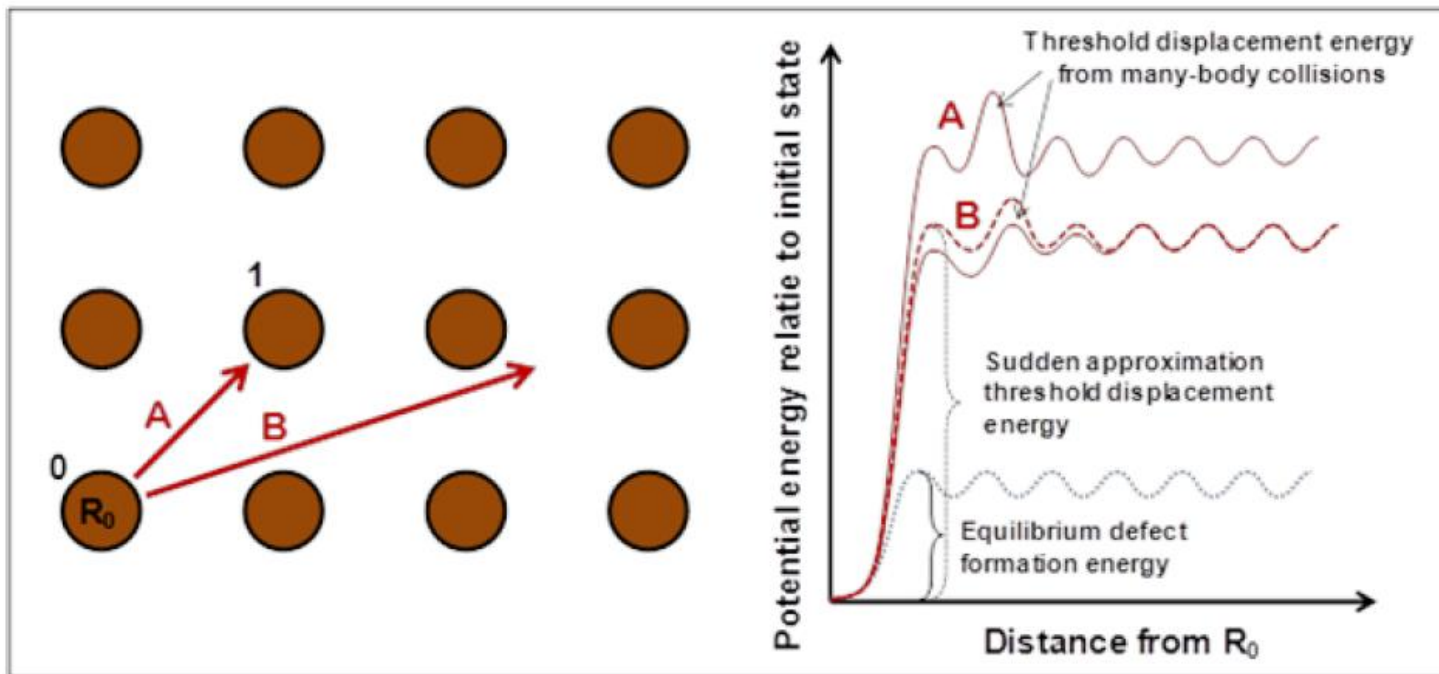


Damage Threshold

$$N_{NRT} \equiv N_F = \kappa \frac{\xi(T)T}{2E_{th}}$$

From: NEA/NSC/DOC(2015)9

- Damage threshold **depends on the direction of the recoil** in the crystal lattice.
- Also depends on the **compound combination**: e.g. NaCl: $E_{th}(\text{Na-Na})$, $E_{th}(\text{Na-Cl})$, $E_{th}(\text{Cl-Na})$, $E_{th}(\text{Cl-Cl})$
- FLUKA use the “**average**” threshold over all crystallographic directions (user defined)
- **Sensitivity studies** using different E_{th} values can provide upper and lower limits on dpa

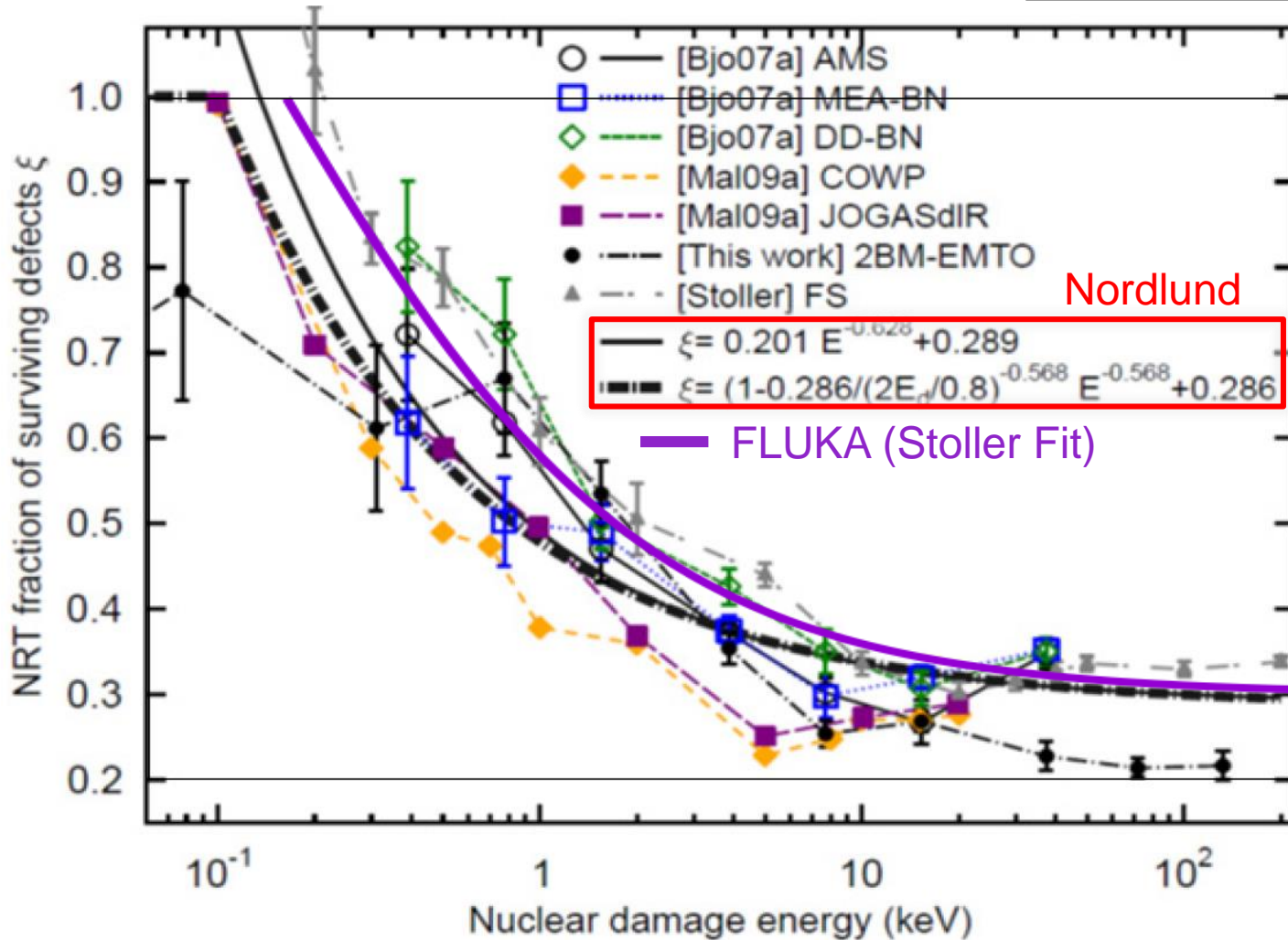


Typical values used in NJOY99 code

Element	E _{th} (eV)	Element	E _{th} (eV)
Lithium	10	Co	40
C in SiC	20	Ni	40
Graphite	30..35	Cu	40
Al	27	Nb	40
Si	25	Mo	60
Mn	40	W	90
Fe	40	Pb	25

Displacement efficiency κ Stoller vs Nordlund

$$N_{NRT} \equiv N_F = \kappa \frac{\xi(T)T}{2E_{th}}$$



Recombination from overlap of collision cascades, “athermal recombination-corrected” dpa: Arc-dpa

Stoller results Comes from Molecular Dynamics simulations

Thermal recombination of defects not considered

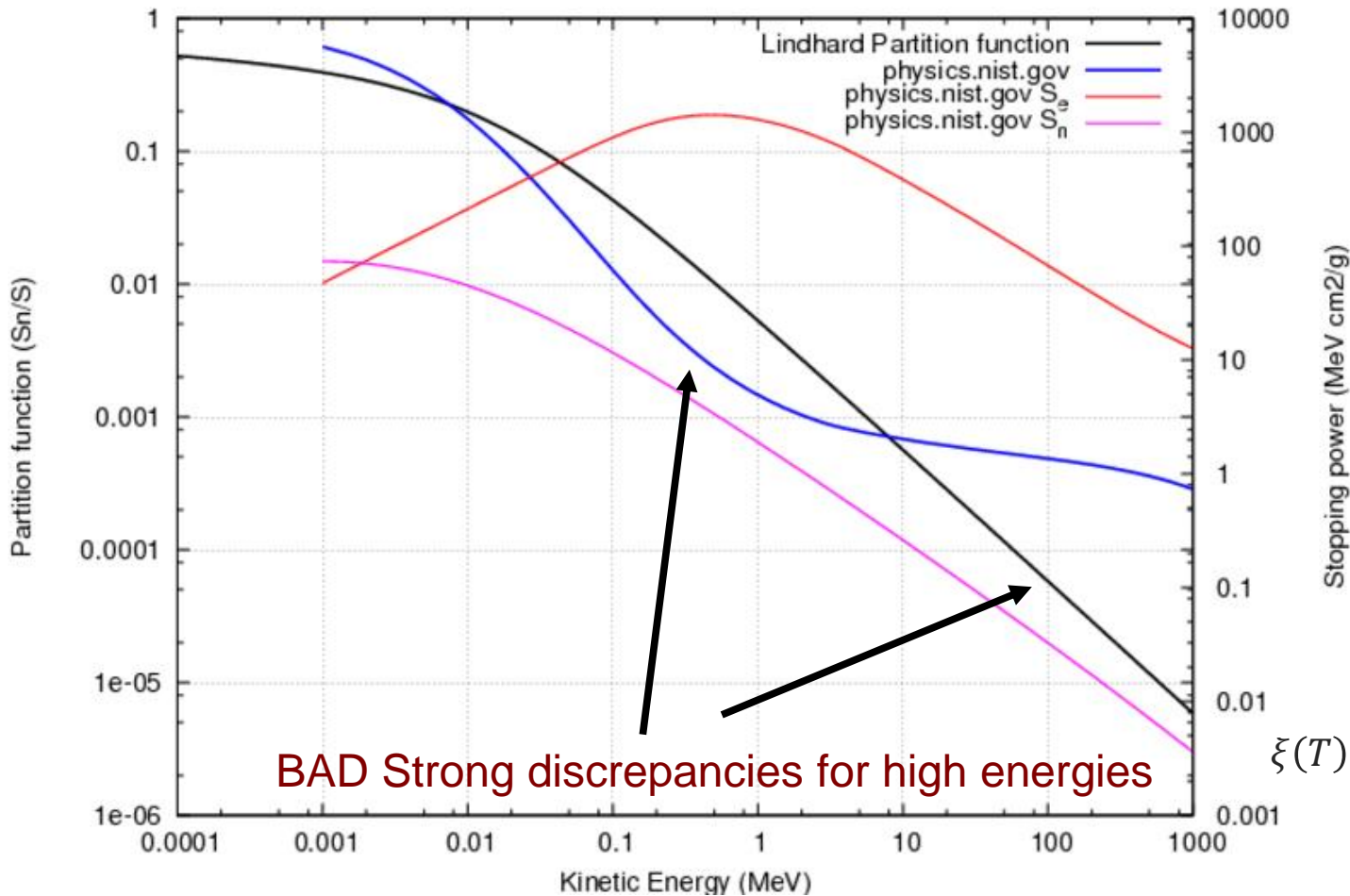
↓ becomes relevant for high temperatures

↓ Overestimation of dpa's

Lindhard partition function ξ

$$N_F = \kappa \frac{\xi(T) T}{2E_{th}}$$

Partition function α on Silicon



Stopping power (MeV cm²/g)

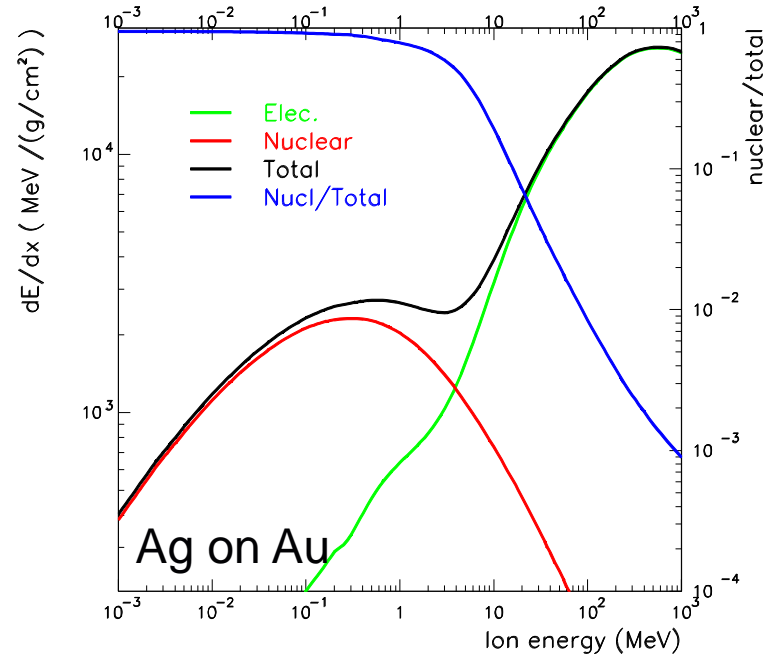
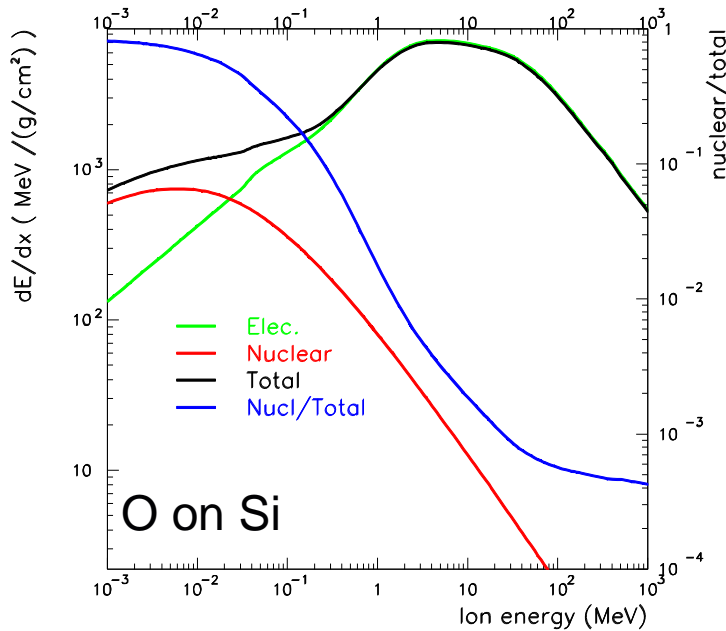
Fraction of stopping power $S(T)$ going into NIEL (non-ionizing energy loss)

$$\xi(T) = \frac{S_n(T)}{S(T)} = \frac{S_n(T)}{S_n(T) + S_e(T)}$$

Nuclear Stopping Power

$$\xi(T) = \frac{S_n(T)}{S(T)}$$

$$N_F = \kappa \frac{\xi(T) T}{2E_{th}}$$



The total (S), nuclear (S_n) and electronic (S_e) stopping power. The **partition function** $S_n/(S_n+S_e)$ is also plotted. The abscissa is the ion total kinetic energy

Partition function decreases with energy
And increases with charge



NIEL/DPA are dominated by
Low energy (heavy) recoils

F.Cerutti

Restricted Nuclear Stopping Power

$$N_F = \kappa \frac{\xi(T) T}{2E_{th}}$$

- Lindhard approximation uses the **unrestricted NIEL**. Including all the energy losses also those below the threshold $E_{th} \rightarrow$ **Overestimation of DPAs**
- FLUKA is using a more accurate way by employing the **restricted nuclear losses**

$$S(E, E_{th}) = N \int_{E_{th}}^{\gamma E} T \left(\frac{d\sigma}{dT} \right) dT$$

where:

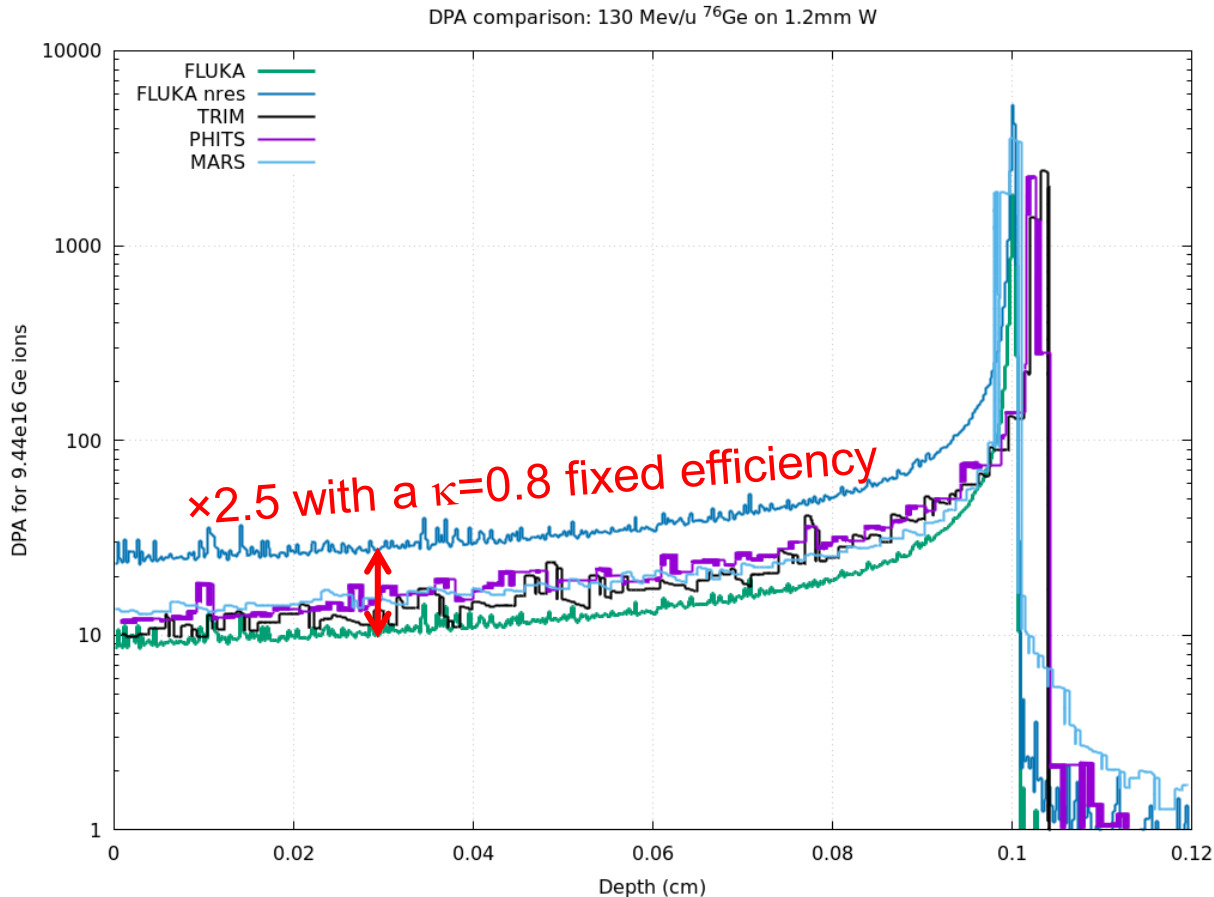
$S(E, E_{th})$ is the restricted energy loss
 N atomic density
 T energy transfer during ion-solid interaction
 $d\sigma/dT$ differential scattering cross section

$$\gamma = \frac{2E(2m + E)}{M + \frac{m^2}{M} + 2(m + E)}$$

maximum fraction of energy transfer during collision

Comparison with other simulation codes

^{76}Ge ion pencil beam of 130 MeV/A uniform in W target a disc of R=0.3568 mm, 1.2 mm thickness



Non-restricted provides higher values if we consider fixed Displ. Effic.



Overestimation of DPA

$$N_F = \kappa \frac{\xi(T)T}{2E_{th}}$$

FLUKA Implementation

Charged particles and heavy ions

- **During transport**
 - Calculate the **restricted non ionizing energy loss**
- **Below threshold**
 - Calculate the **integrated nuclear stopping power** with the **Lindhard** partition function
- **At (elastic and inelastic) interactions**
 - Calculate **the recoil**, to be transported or treated as below threshold

Neutrons:

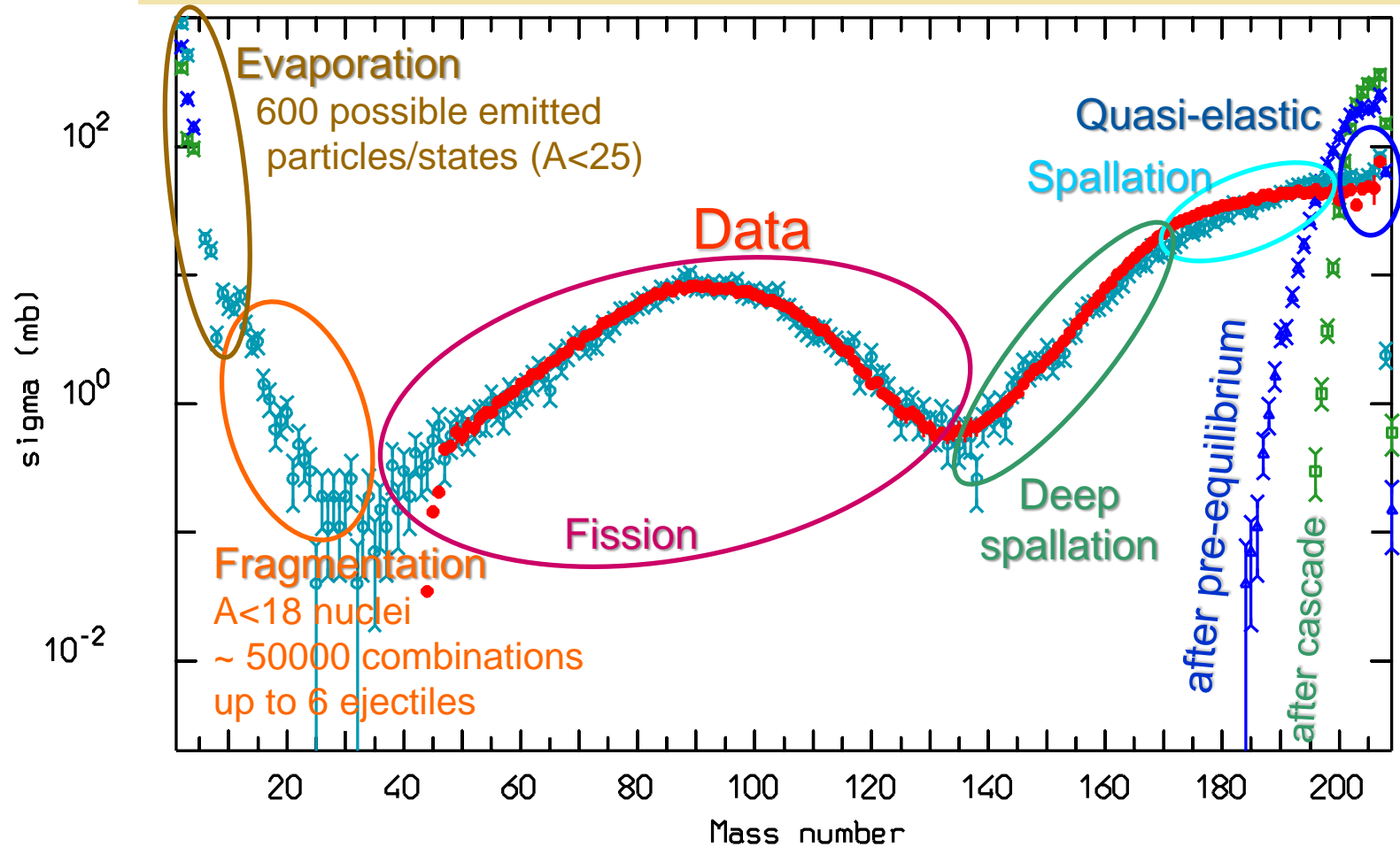
- **High energy $E_n > 20$ MeV**
 - Calculate the **recoils** after interaction
 - Treat recoil as a “normal” charged particle/ion
- **Low energy $E_n \leq 20$ MeV (group-wise)**
 - Calculate the NIEL from NJOY
- **Low energy $E_n \leq 20$ MeV (point-wise)**
 - Calculate the recoil if possible
 - Treat recoil as a “normal” charged particle/ion

Implementation in FLUKA: A. Fasso et al. Prog. In Nucl. Science and Technology, Vol. 2, p769-775 (2011)

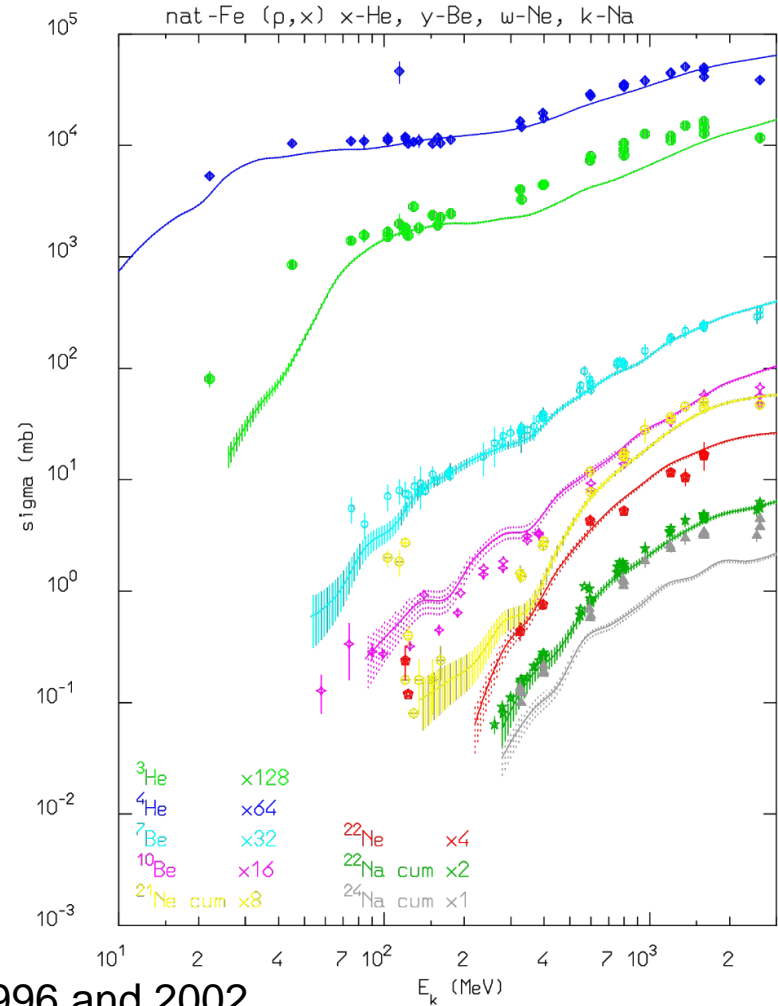
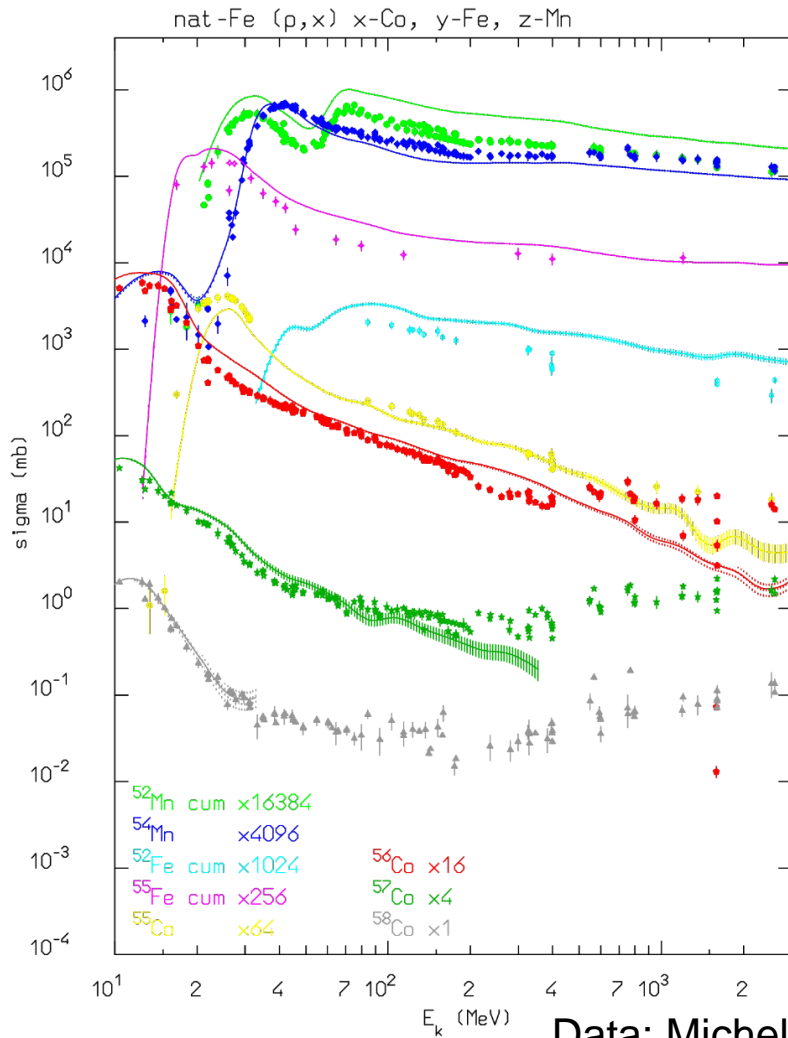


Example of fission/evaporation

1 A GeV $^{208}\text{Pb} + \text{p}$ reactions Nucl. Phys. A 686 (2001) 481-524



Isotope production for $^{nat}\text{Fe}(p,x)$:



Data: Michel et al. 1996 and 2002

Estimates for CERN injectors and future facilities

- Beam Dump Facility
- BLIP capsule
- PS Internal Beam Dumps

Beam Dump Facility (BDF)

Beam:

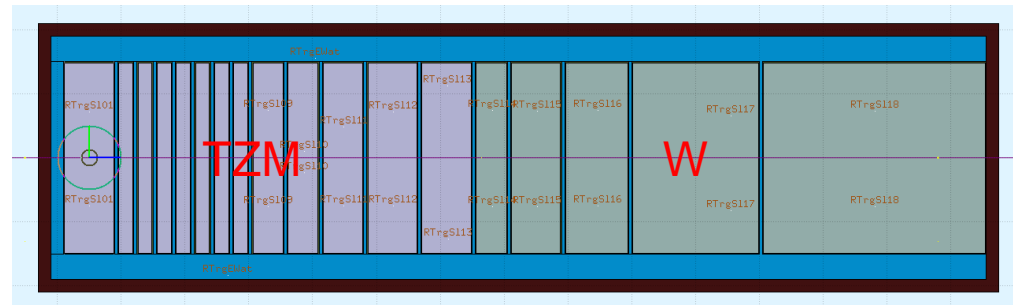
- Protons: 400 GeV/c
- Sweep pattern:
 - radius 3 cm
 - 1σ 0.6 cm

Geometry:

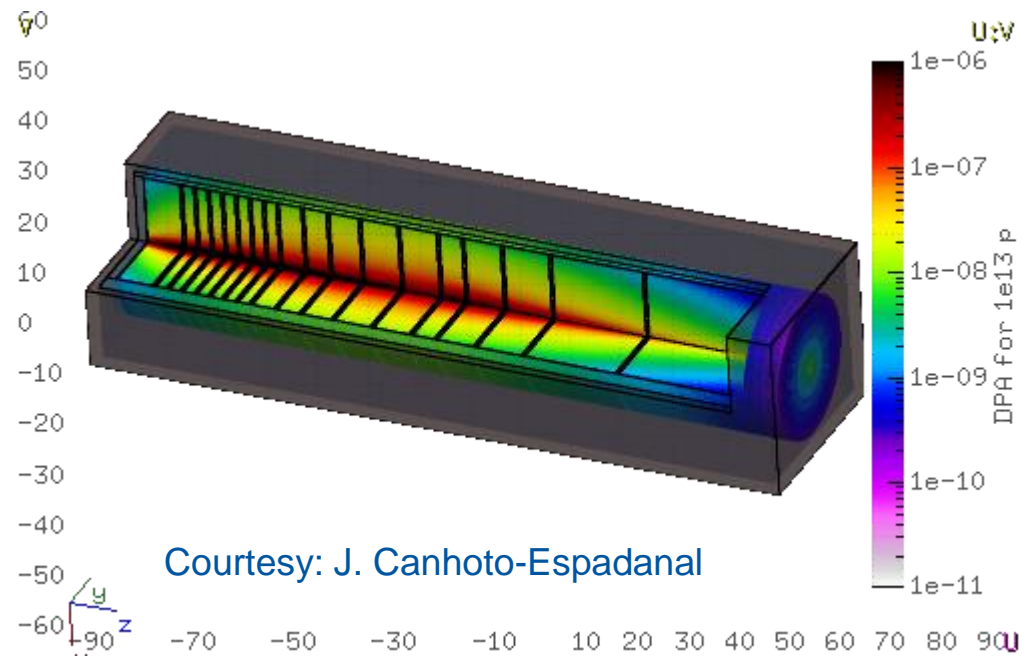
- 1.4 m long cylinder discs of TZM enclosed in Ta
- W enclosed in Ta
- 1.5 mm Ta cladding and 5 mm water gaps

Materials:

- Tungsten $E_d=90$ eV
- SS 316N $E_d=40$ eV
- Tantalum $E_d=53$ eV
- TZM (Mo,Zr,Ti...) $E_d=60$ eV

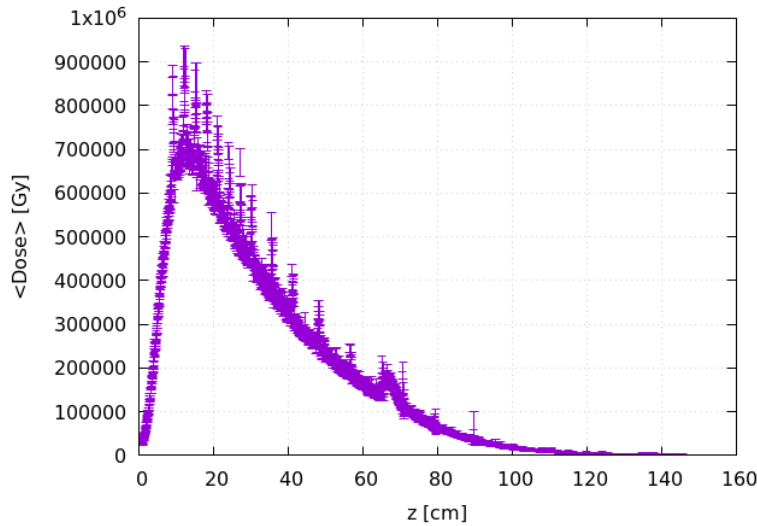


Ta cladding Water cooling

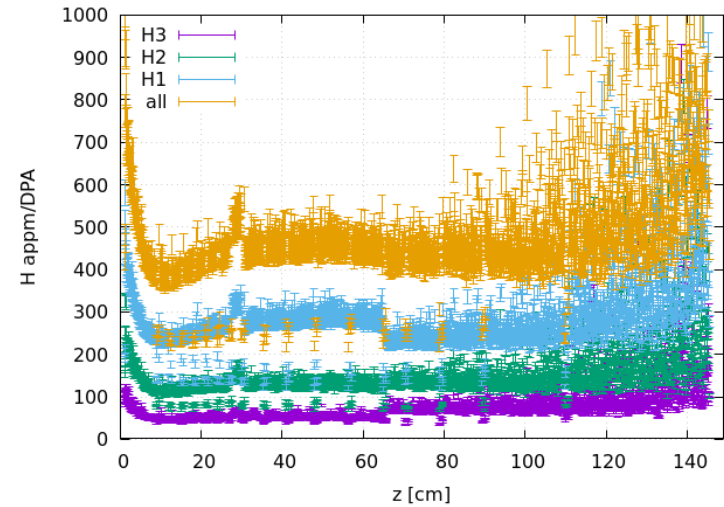


BDF Results: H/He[appm] vs DPA

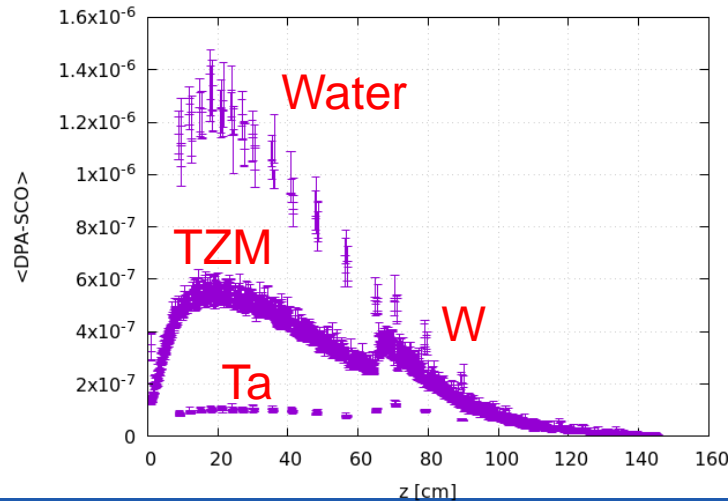
Dose (energy deposited per unit mass), $R < 0.25\text{cm}$, for $3e13$ POT



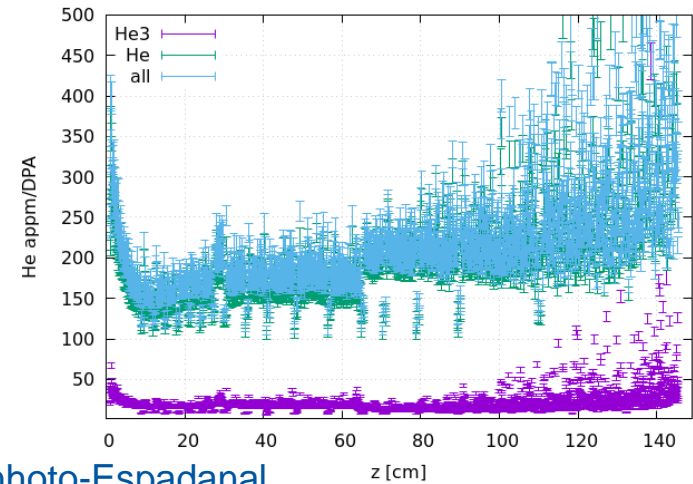
H appm/ DPA



Displacements per atoms, $R < 0.25\text{cm}$, for $3e13$ POT



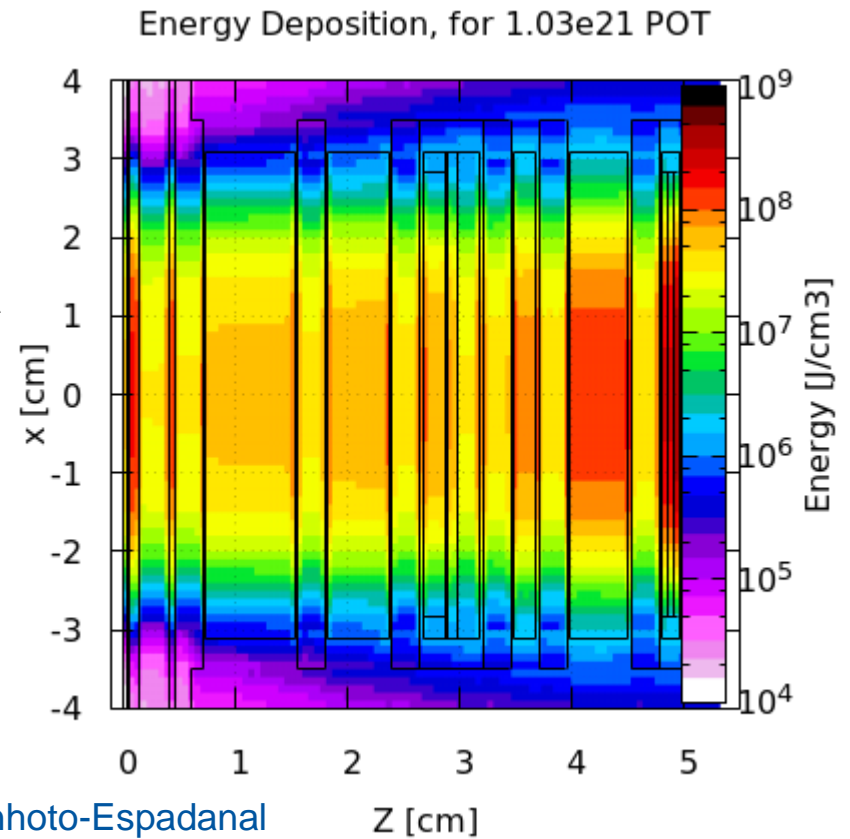
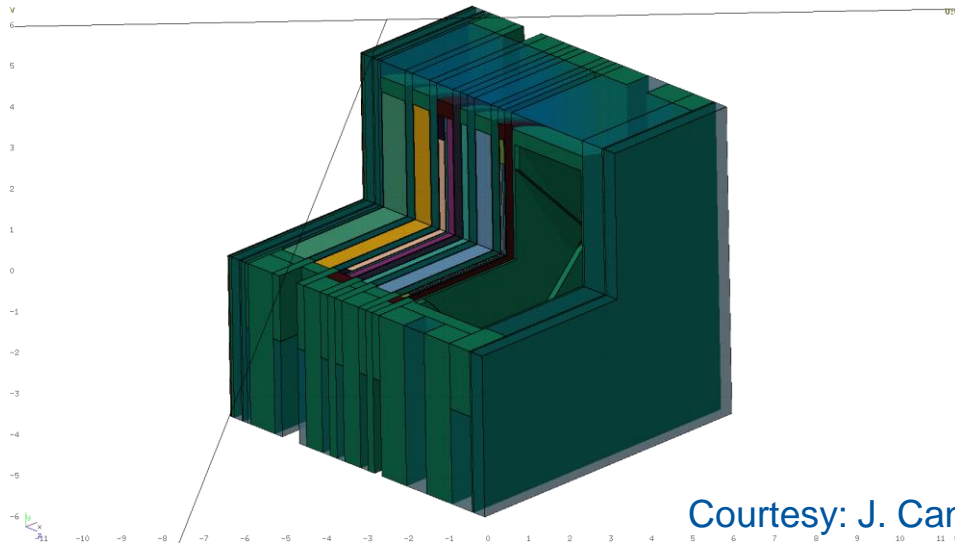
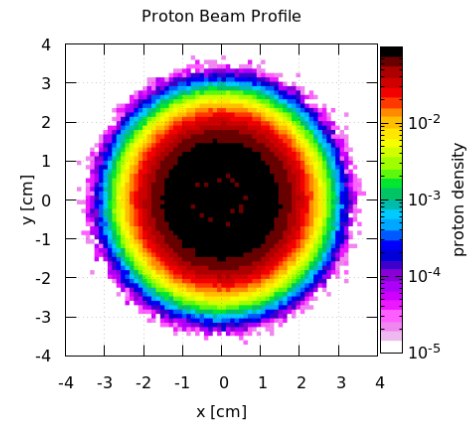
He appm/ DPA



Courtesy: J. Canhoto-Espadanal

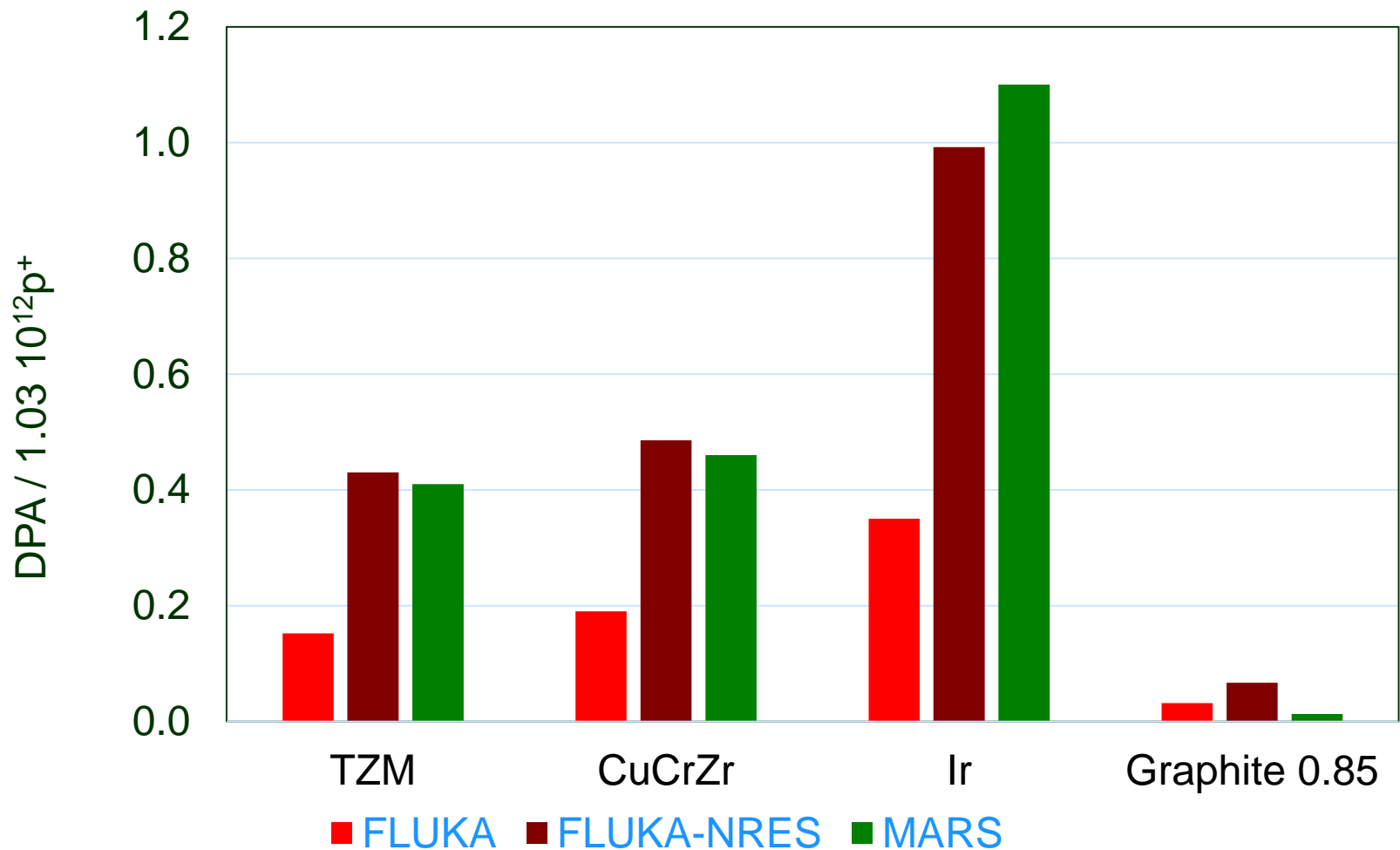
BLIP capsule

- Beam
 - Proton E=181 MeV
 - $\sigma_{x,y} = 5.1$ mm
- Geometry: Layers of
 - Window SS304L 0.3mm
 - TZM 0.5mm
 - CuCrZr 0.5mm
 - Ir 0.5mm
 - Graphite(0.85g/cm³) 0.1mm



Courtesy: J. Canhoto-Espadanal

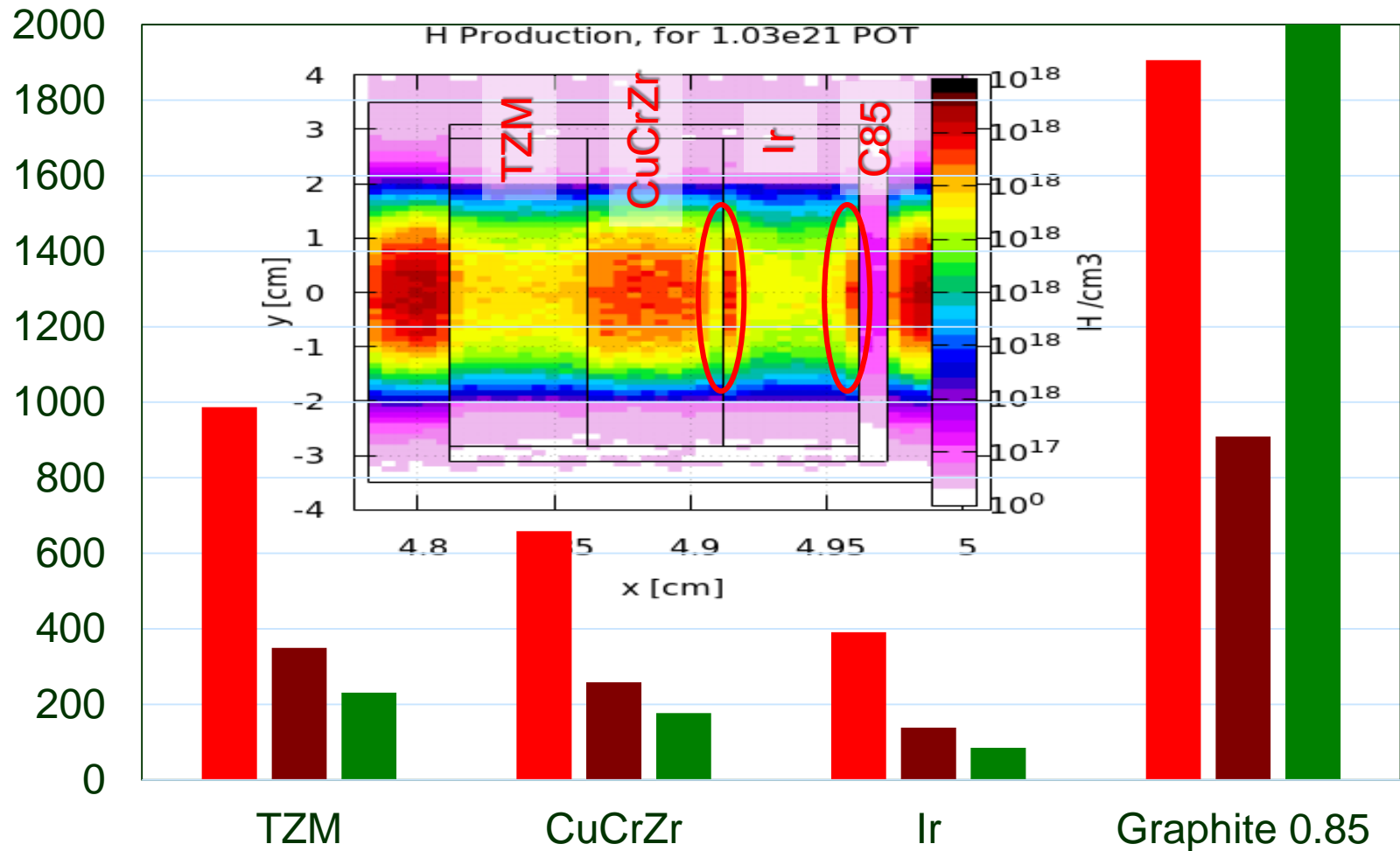
DPA High-Z BLIP [FLUKA vs MARS]



Courtesy: J. Canhoto-Espadanal

Note: NRT model of MARS

H appm/DPA High-Z BLIP

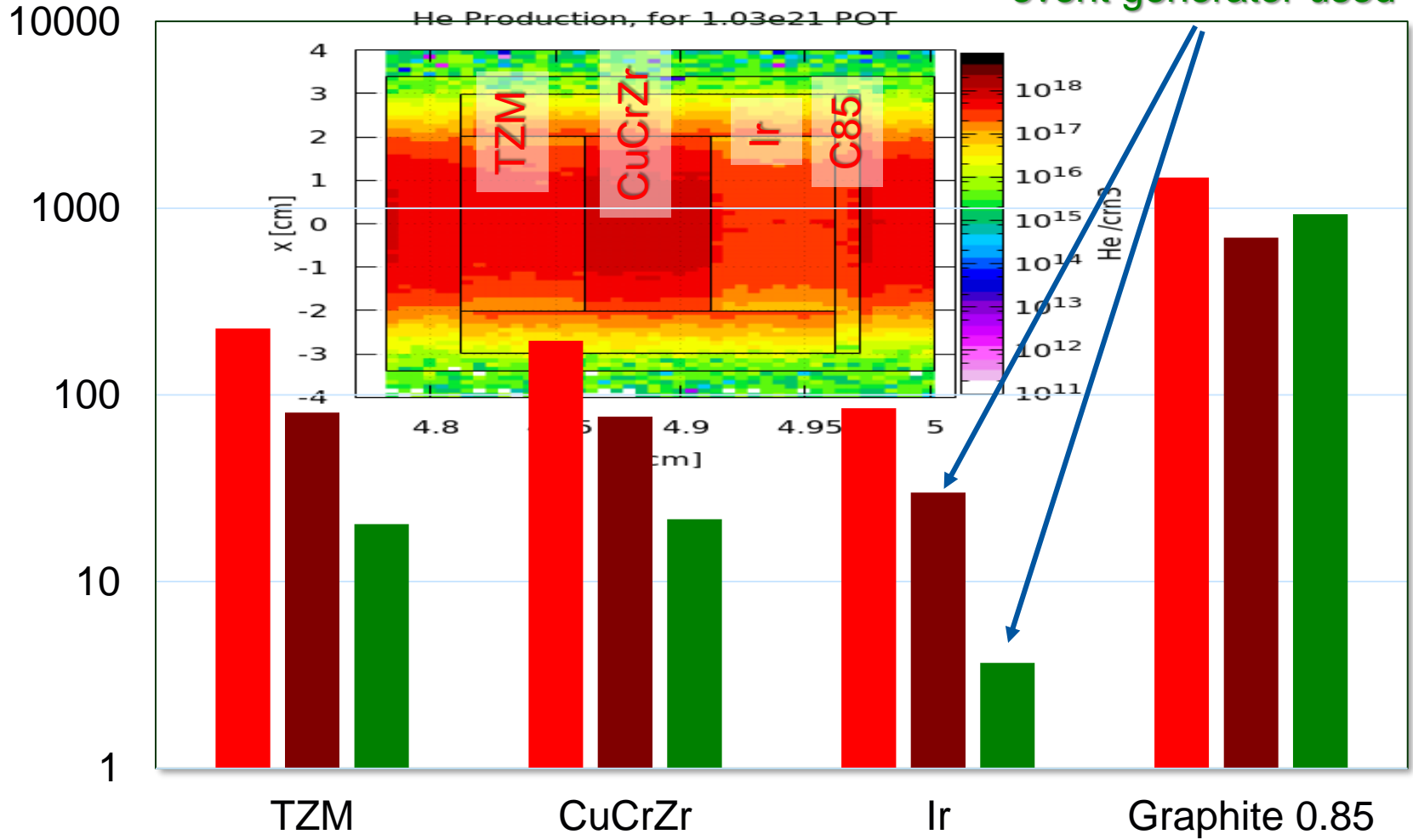


Courtesy: J. Canhoto-Espadanal ■ FLUKA ■ FLUKA-NRES ■ MARS

He appm/DPA High-Z BLIP

Probably due to "Old" MARS event generator used

Warning:
Log-scale



Courtesy: J. Canhoto-Espadanal ■ FLUKA ■ FLUKA-NRES ■ MARS



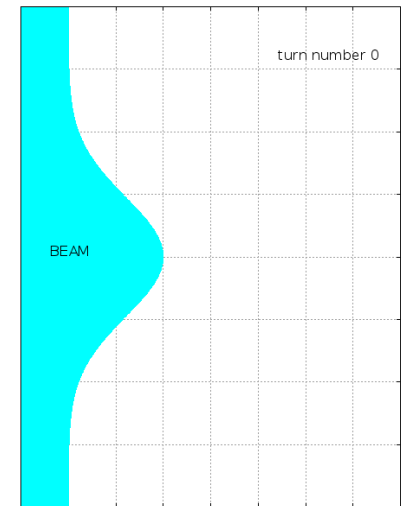
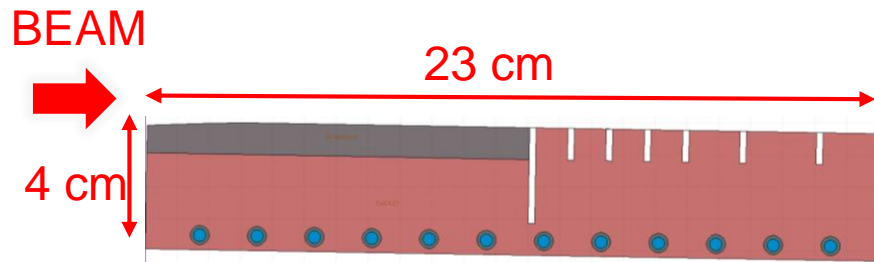
27-28/11/2017

1st Workshop of ARIES WP17 Power Mat - Politecnico de Torino

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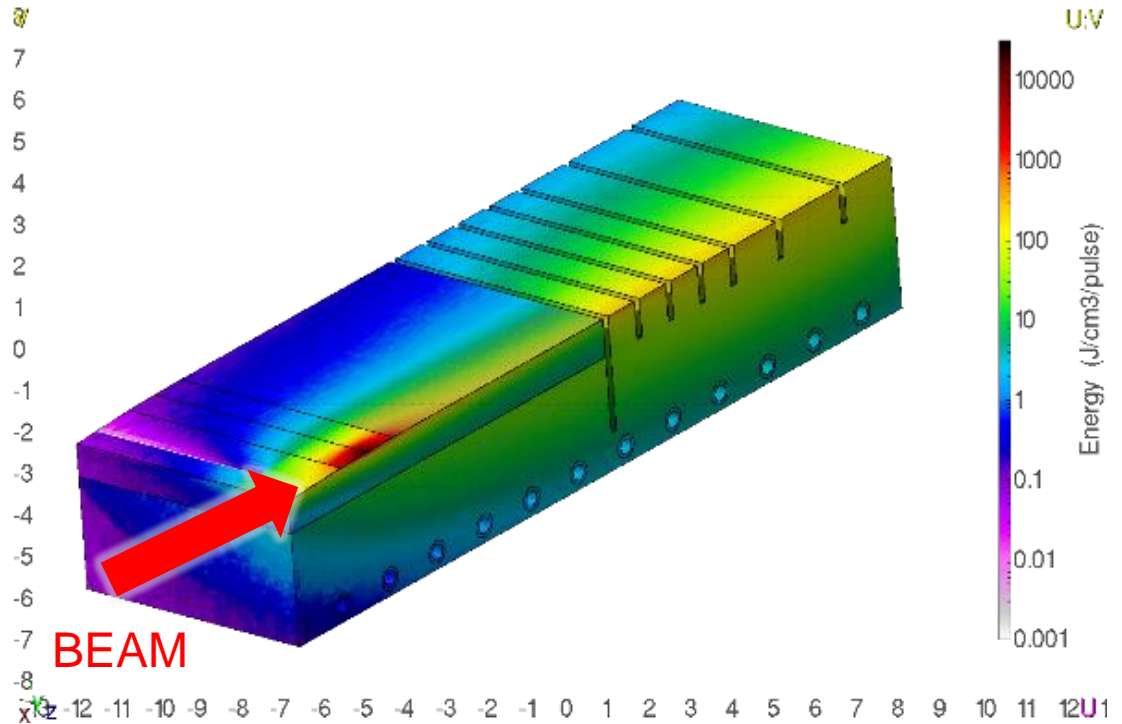
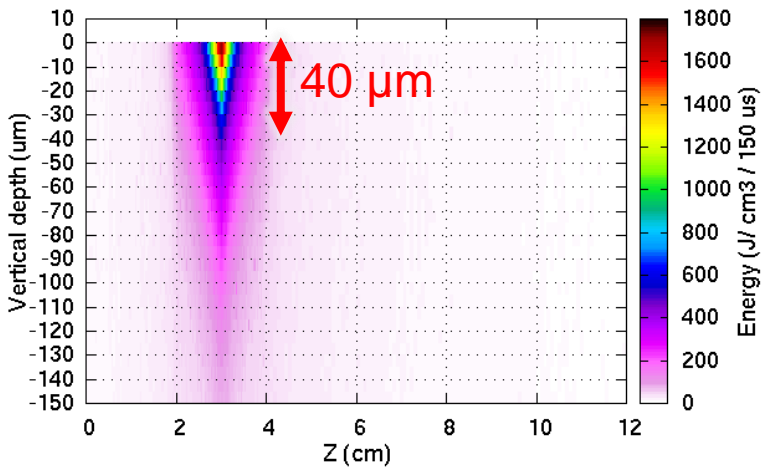


PS Internal Beam Dumps

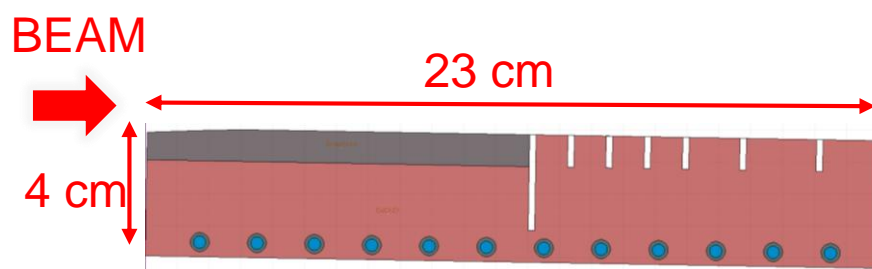


Challenging Energy density

Superficial energy deposition



PS Internal Beam Dumps: Damage



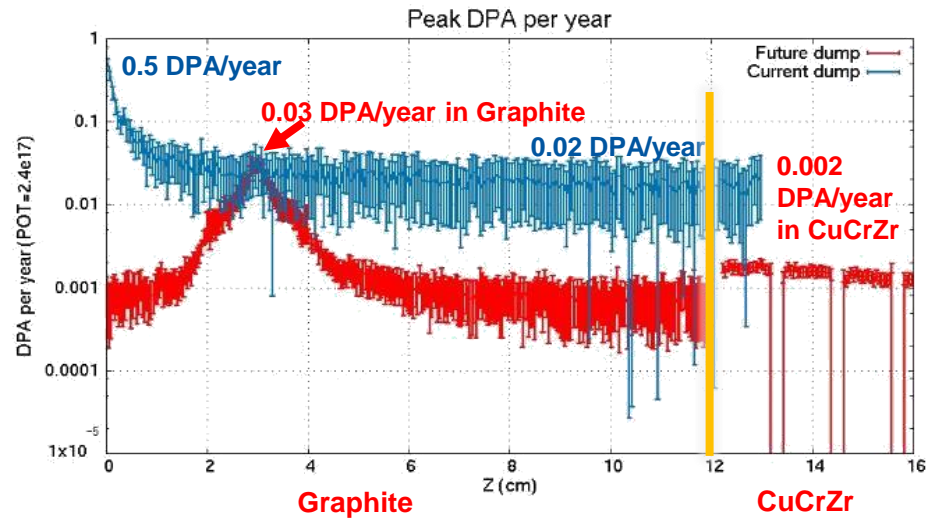
Beam properties:

- 26 GeV/c proton beam
- 2.4×10^{17} POT per year
- $\sigma_h = 1.74$ mm $\sigma_v = 0.87$ mm
- Beam is shaved in a thin top most layer

Graphite:

Experience at CERN with **CNGS** air cooled graphite target (SPS beam). About 1200°C reached for each pulse. At the end of operation: 1.5 DPA

→ No problem observed on graphite



Damage threshold energies considered:
 $E_{th}(\text{Graphite}) = 30$ eV - typical value 30-35 eV
 $E_{th}(\text{CuCrZr and SS304L}) = 40$ eV

CuCrZr:

Literature on neutron irradiation indicated for **similar dpa damage** some possible radiation hardening and thermal conductivity **degradation but not dramatic effects**

Summary

- FLUKA dpa model uses a **restricted NIEL** computed during initialization and run time.
- **Not based on Lindhard** but reworked all formulas
- The **only free parameter** for the user is the **damage threshold**. It depends on the direction of the recoil. Simple averaging is not correct
- Uniform treatment from the transport threshold up to the highest energies
- Use of **Stoller displacement efficiency instead of a fixed 0.8** as NRT suggests
- Not considered thermal recombination of defects → **overestimation of dpa**. Improving the estimate would require Molecular Dynamics simulations

Summary

- **FLUKA is employed to evaluate radiation-induced damage** on new dumps facilities (BDF, BLIP) and elements of the CERN injector chain (PS internal dumps)
- Simulations provided a way of quantifying the damage through estimation of dpa and gas production (H, He)
- Comparison of **estimations of dpa (simulations)** and modifications of macroscopic quantities (experimental: thermal, electric, etc..., properties) helps to extract conclusions on radiation-induced damage

Possible Future improvements:

- Implementation of the Nordlund arc-dpa
- More accurate recoil momentum cross section for **pair production** and **Bremsstrahlung**
- **Point wise** treatment of **low energy neutrons** will provide correct recoil information
- Multiple damage thresholds for compounds



Thank you for your attention!

Any question?

Extra slides

$$N_F = \kappa \frac{\xi(T)T}{2E_{th}}$$

E_{th} Damage Threshold Energy

- E_{th} is the value of the threshold displacement energy averaged over all crystallographic directions or a minimum energy to produce a defect

Element	E_{th} (eV)	Element	E_{th} (eV)
Lithium	10	Co	40
C in SiC	20	Ni	40
Graphite	30..35	Cu	40
Al	27	Nb	40
Si	25	Mo	60
Mn	40	W	90
Fe	40	Pb	25

Typical values used in NJOY99 code

- The only variable requested for FLUKA
 - MAT-PROP** *WHAT(1)* = E_{th} (eV)
 - WHAT(4,5,6)* = Material range
 - SDUM* = **DPA-ENER**

Damage Threshold in Compounds

- NJOY (MT=444) sums up the cross section multiplied by the damage energies, which is the damage production cross section representing the effective kinetic energy of recoiled atom for reaction types i at neutron energy E_n

$$(E\sigma)_{DPA} = \sum_i E_{th,i} \sigma_i(E_n)$$

Problematic:

- Damage threshold depends on the lattice structure
- Damage threshold can be quite different for each combination for **the specific compound**
e.g. NaCl: $E_{th}(\text{Na-Na})$, $E_{th}(\text{Na-Cl})$, $E_{th}(\text{Cl-Na})$, $E_{th}(\text{Cl-Cl})$
- Simple weighting with the atom/mass fraction doesn't work
- FLUKA's approximation is using a unique average damage threshold E_{th} for the compounds as well → **A sensitivity study can be performed**

Only free parameter for the FLUKA user is E_{th}

$$N_F = \kappa \frac{\xi(T)T}{2E_{th}}$$

κ displacement efficiency

- $\kappa=0.8$ value deviates from the **hard sphere model** (K&P), and compensates for the forward scattering in the displacement cascade
- The displacement efficiency κ can be considered as independent of T only in the range of $T \leq 1-2$ keV. At higher energies, the development of collision cascades results in **defect migration** and **recombination of Frenkel pairs** due to overlapping of different branches of a cascade which translates into decay of $\kappa(T)$.
- From molecular dynamics (MD*) simulations of the primary cascade the number of surviving displacements, N_{MD} , normalized to the number of those from NRT model, N_{NRT} , decreases down to the values about 0.2–0.3 at $T \approx 20-100$ keV. The efficiency in question only slightly depends on atomic number Z and the temperature. $N_{MD}/N_{NRT} = 0.3-1.3$

- $$N_{MD} / N_{NRT} = 0.3 - 1.3 \left(-\frac{9.57}{X} + \frac{17.1}{X^{4/3}} - \frac{8.81}{X^{5/3}} \right)$$

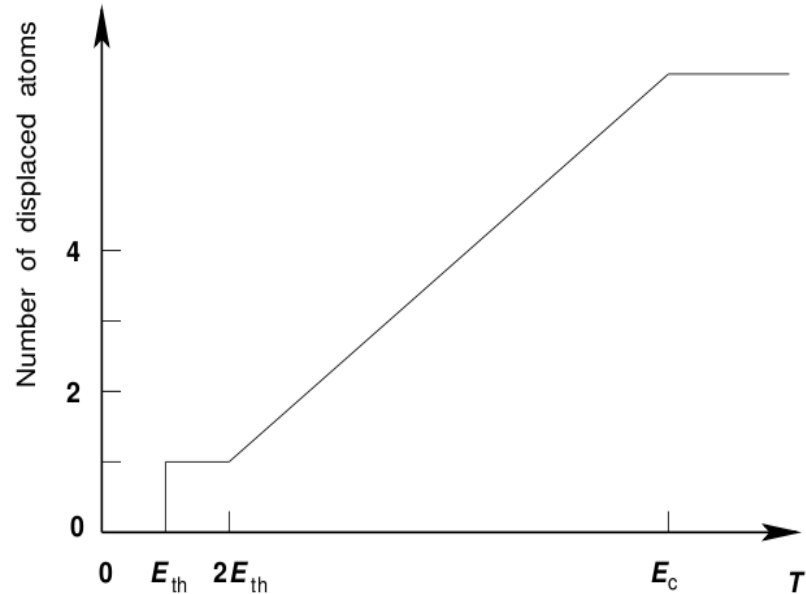
where $X \equiv 20 T$ (in keV).

$$N_F = \kappa \frac{\xi(T)T}{2E_{th}}$$

Factor of 2 (Kinchin & Pease)

- The cascade is created by a sequence of two-body elastic collisions between atoms
- In the collision process, the energy transferred to the lattice is zero
- For all energies $T < E_c$ electronic stopping is ignored and only atomic collisions take place. No additional displacement occur above the cut-off energy E_c
- The energy transfer cross section is given by the **hard-sphere** model.

$$\begin{aligned}
 v(T) &= 0 && \text{for } 0 < T < E_{th} \text{ (phonons)} \\
 v(T) &= 1 && \text{for } E_{th} < T < 2E_{th} \\
 v(T) &= T/2E_{th} && \text{for } 2E_{th} < T < E_c \\
 v(T) &= E_c/2E_{th} && \text{for } T > E_c
 \end{aligned}$$



Schematic relation between the number of displaced atoms in the cascade and the kinetic energy T of the primary knock-on atom

Energy is equally shared between two atoms after the first collision
Compensates for the energy lost to sub threshold reactions

$$N_F = \kappa \frac{\xi(T)T}{2E_{th}}$$

Lindhard partition function ξ [1/2]

- The partition function gives the fraction of **stopping power S** that goes to NIEL
- Approximations used: Electrons do not produce recoil nuclei with appreciable energy, lattice binding energy is neglected, etc...

where

$$(S_n + S_e)E'_n(E) = \int E_n(T) \frac{d\sigma_n}{dT} dT$$

- approximated to

$$S_{n\epsilon}(E) = \int T_{n\epsilon} d\sigma_{n\epsilon}$$

$$\xi(T) = \frac{1}{1 + F_L \cdot (3.4008 \cdot \epsilon(T)^{1/6} + 0.40244 \cdot \epsilon(T)^{3/4} + \epsilon(T))}$$

$$F_L = 30.724 \cdot Z_1 \cdot Z_2 \sqrt{Z_1^{2/3} + Z_2^{2/3}}$$

$$\epsilon(T) = \frac{T}{0.0793 \frac{Z_1^{2/3} \cdot \sqrt{Z_2}}{(Z_1^{2/3} + Z_2^{2/3})^{3/4}} \cdot \frac{(A_1 + A_2)^{3/2}}{A_1^{3/2} \sqrt{A_2}}}$$

Z,A	charge and mass
1	projectile
2	medium
T	recoil energy (eV)

Nice feature: It can handle any projectile Z_1, A_1 whichever charged particle



Nuclear Stopping power

- Nuclear stopping power (unrestricted)

$$\frac{1}{\rho} S_n(E, E_{th}) = -2\pi N \int_0^{b_{\max}} b \frac{db}{d\theta} W(\theta, E) db$$

- Energy transferred to recoil atom

$$W(\theta, T) = \gamma T \sin^2(\theta / 2)$$

- Deflection angle, by integrating over all impact parameters b

$$\theta = \pi - 2 \int \frac{b dr}{r^2 \sqrt{1 - \frac{V(r)}{E_{cms}} - \frac{b^2}{r^2}}}$$

- Universal potential

$$V(r) = \frac{Z_1 Z_2 e^2}{r} F_s \left(\frac{r}{r_s} \right)$$

where:

$$F_s(x) = \sum a_i \exp(-c_i x)$$

$$r_s = 0.88534 r_B / (Z_1^{0.23} + Z_2^{0.23})$$

$$r_s = 0.88534 r_B Z_1^{-1/3}$$

screening function
screening length
in case of particle

Ziegler approximation

- Reduced kinetic energy ε (T in keV)

$$\varepsilon = \frac{32.536 T}{(Z_1^{0.23} + Z_2^{0.23}) \left(1 + \frac{M_1}{M_2}\right) Z_1 Z_2}$$

- Reduced stopping power

$$\text{if } \varepsilon < 30 \quad \hat{S}_n(\varepsilon) = \frac{0.5 \ln(1 + 1.1383 \varepsilon)}{\varepsilon + 0.01321 \varepsilon^{0.21226} + 0.19593 \sqrt{\varepsilon}}$$

$$\text{if } \varepsilon \geq 30 \quad \hat{S}_n(\varepsilon) = \frac{\ln(\varepsilon)}{2\varepsilon}$$

Important features of Reduced Stopping Power

- Independent from the **projectile** and **target** combination
- Accurate within **1%** for $\varepsilon < 1$ and to within **5%** or better for $\varepsilon > 3$
- Stopping power (MeV/g/cm²)

$$\frac{1}{\rho} S_n(T) = \frac{5105.3 Z_1 Z_2 \hat{S}_n(\varepsilon)}{(Z_1^{0.23} + Z_2^{0.23}) \left(1 + \frac{M_2}{M_1}\right) A}$$

Restricted Stopping Power

- The restricted nuclear stopping power is calculated the same way only integrating from 0 impact parameter up to a maximum b_{max} which corresponds to a transfer of energy equal to the

$$E_{th} = W_{min}(\theta_{min}, T)$$

$$\frac{1}{\rho} S_n(E, E_{th}) = -2\pi N \int_0^{b_{max}} b \frac{db}{d\theta} W(\theta, E) d\theta$$

- To find b_{max} we have to approximately solve the previous θ integral using an iterative approach for

$$\theta_{min} = 2 \arcsin \left(\sqrt{\frac{E_{th}}{\gamma T}} \right)$$

This can be done either by integrating numerically for θ or using the magic scattering formula from Biersack-Haggmark that gives a fitting to $\sin^2(\theta/2)$

Implementation: Charged Particles

- During the transport of all charged particles and heavy ions the dpa estimation is based on the restricted nuclear stopping power while for NIEL on the unrestricted one.
- For every charged particle above the transport threshold and for every Monte Carlo step, the number of defects is calculated based on a modified multiple integral
- Taking into account also the **second level of sub-cascades initiated by the projectile**

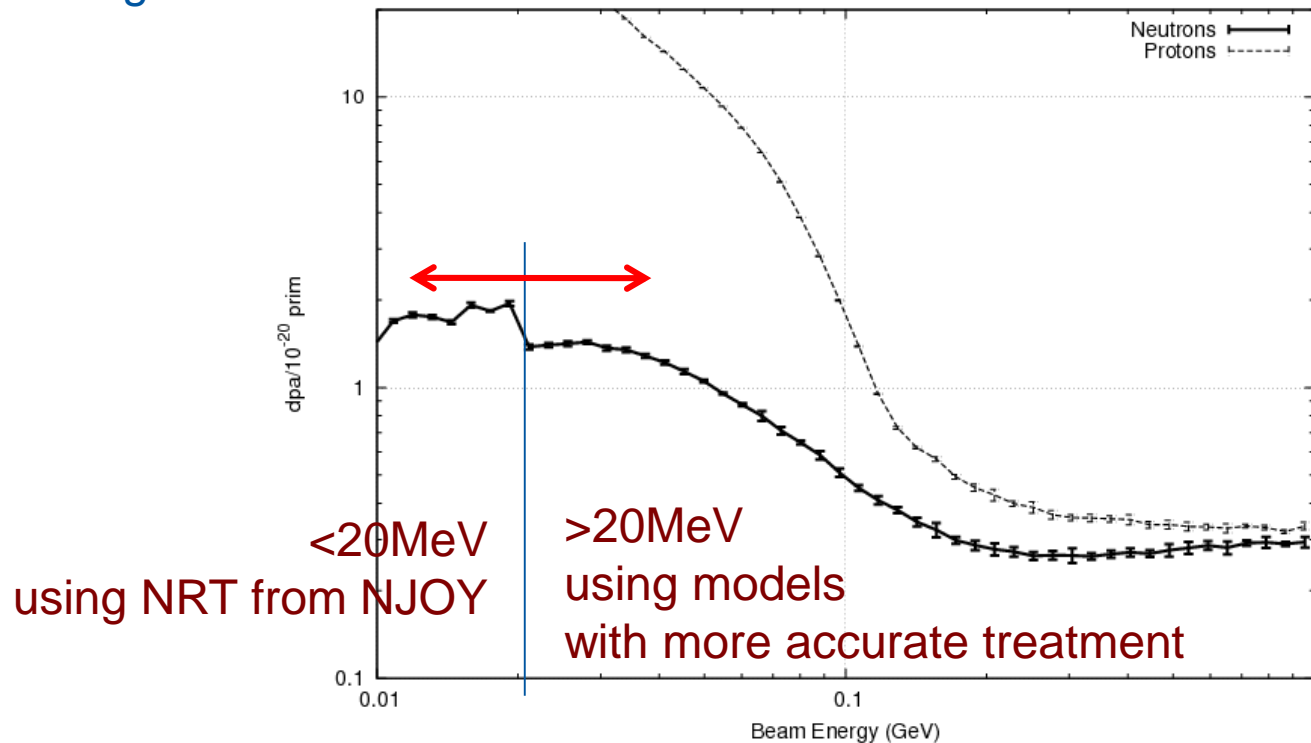
$$N(E) = \int_{E_{th}}^{\gamma E} \left[\xi_r(T, E_{th}) \left(\frac{d\sigma}{dT} \right)_E \int_{E_{th}}^{\gamma T} \kappa(T') \xi(T') T' \left(\frac{d\sigma}{dT'} \right)_{T'} dT' \right] dT$$

↑
restricted partition function
↑
Lindhard partition function

- Below the transport threshold (1 keV) it employs the Lindhard approximation

Group Wise Neutron Artifacts

- Due to the group treatment of low-energy neutrons, there is no direct way to calculate properly the recoils.
- Therefore the evaluation is based on the KERMA factors calculated by NJOY, which in turn is based on the Unrestricted Nuclear losses from using the NRT model.



Implementation: others

For Bremsstrahlung and pair production the recoil is sampled randomly from an approximation of the recoil momentum cross section

Bremsstrahlung

$$\frac{d\sigma}{dp_{\perp}} = \frac{32a(Za)^2}{kp_{\perp}^3} \left[1 - \frac{k}{E} + \frac{1}{2} \left(\frac{k}{E} \right)^2 \right] \ln \left(\frac{p_{\perp}}{m_e} \right)$$

Pair production

$$\frac{d\sigma}{dp} = \frac{0.183 \cdot 10^{-2} Z^2}{p^3} (\ln(p) + 0.5)$$

both can be written in the same approximate way as

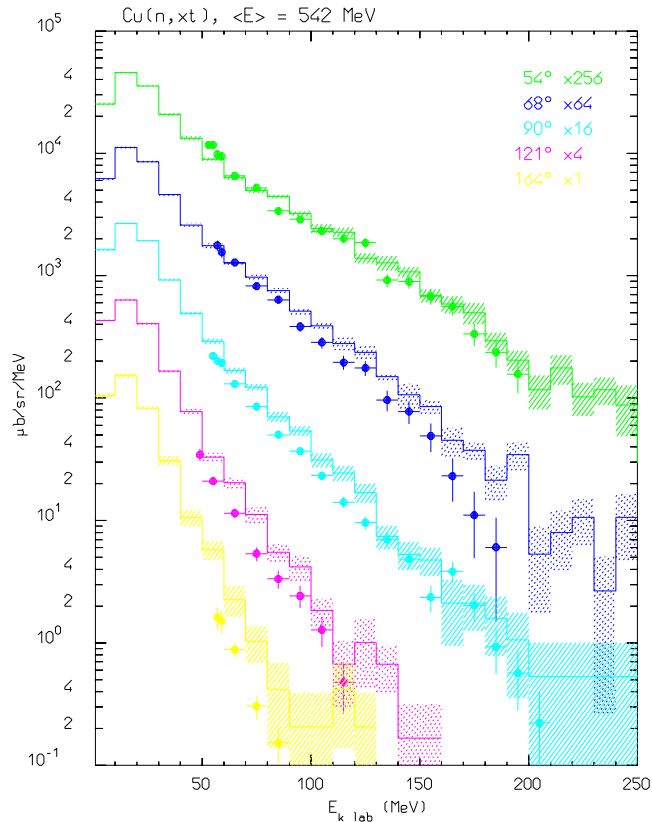
$$\frac{d\sigma}{dp} \propto \frac{\ln(p/c)}{p^3}$$

where the recoil momentum is sampled randomly by rejection from a similar function

Coalescence:

- d, t, ^3He , and alpha's generated during the (G)INC and preequilibrium stage
- All possible combinations of (unbound) nucleons and/or light fragments checked at each stage of system evolution
- FOM evaluation based on phase space "closeness" used to decide whether a light fragment is formed rather than not
 - ❑ FOM evaluated in the CMS of the candidate fragment at the time of minimum distance
 - ❑ Naively a momentum or position FOM should be used, but not both due to quantum non commutation
 - ❑ ... however the best results are obtained with a Wigner transform FOM (assuming gaussian wave packets) which should be the correct way of considering together positions and momenta
- Binding energy redistributed between the emitted fragment and residual excitation (exact conservation of 4-momenta)

Coalescence



High energy light fragments are emitted through the coalescence mechanism: “put together” emitted nucleons that are near in phase space.

Example : double differential t production from 542 MeV neutrons on Copper

Warning: coalescence is OFF by default
Can be important, ex for . residual nuclei.

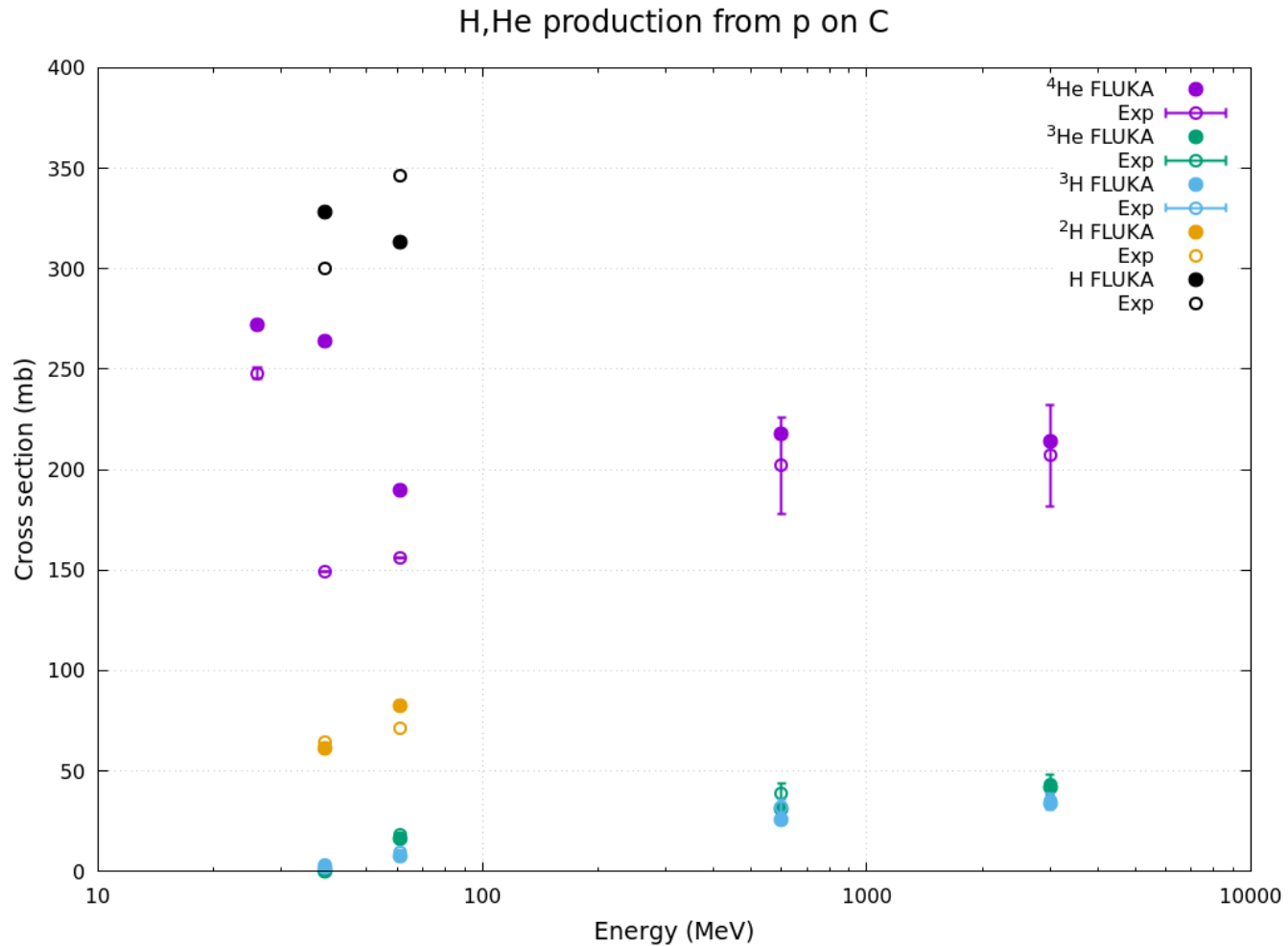
To activate it:

PHYSICS 1.

COALESCE

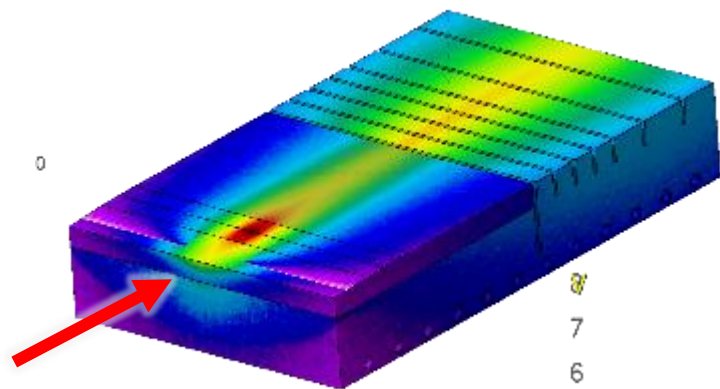
If coalescence is on, switch on Heavy ion transport and interactions (see later)

Particle production in C(p,x) reaction

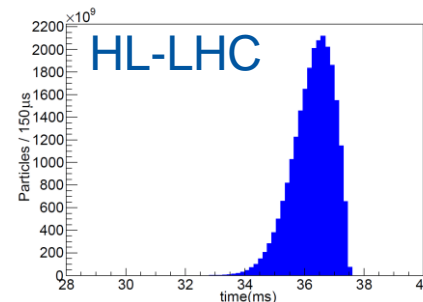


Data: JNST36 313 1999, PRC7 2179 1973

Energy Density Distribution from FLUKA

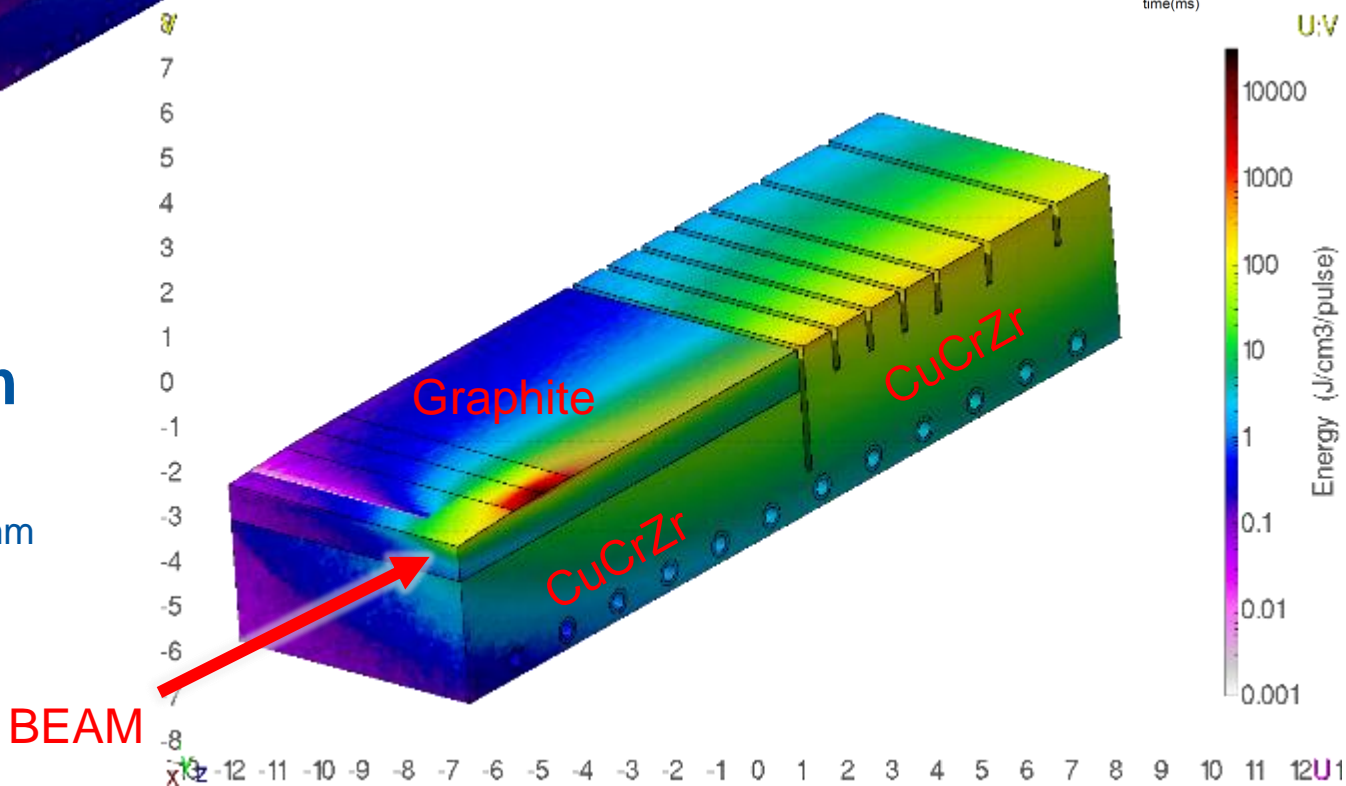


Values shown
are accumulated
per pulse

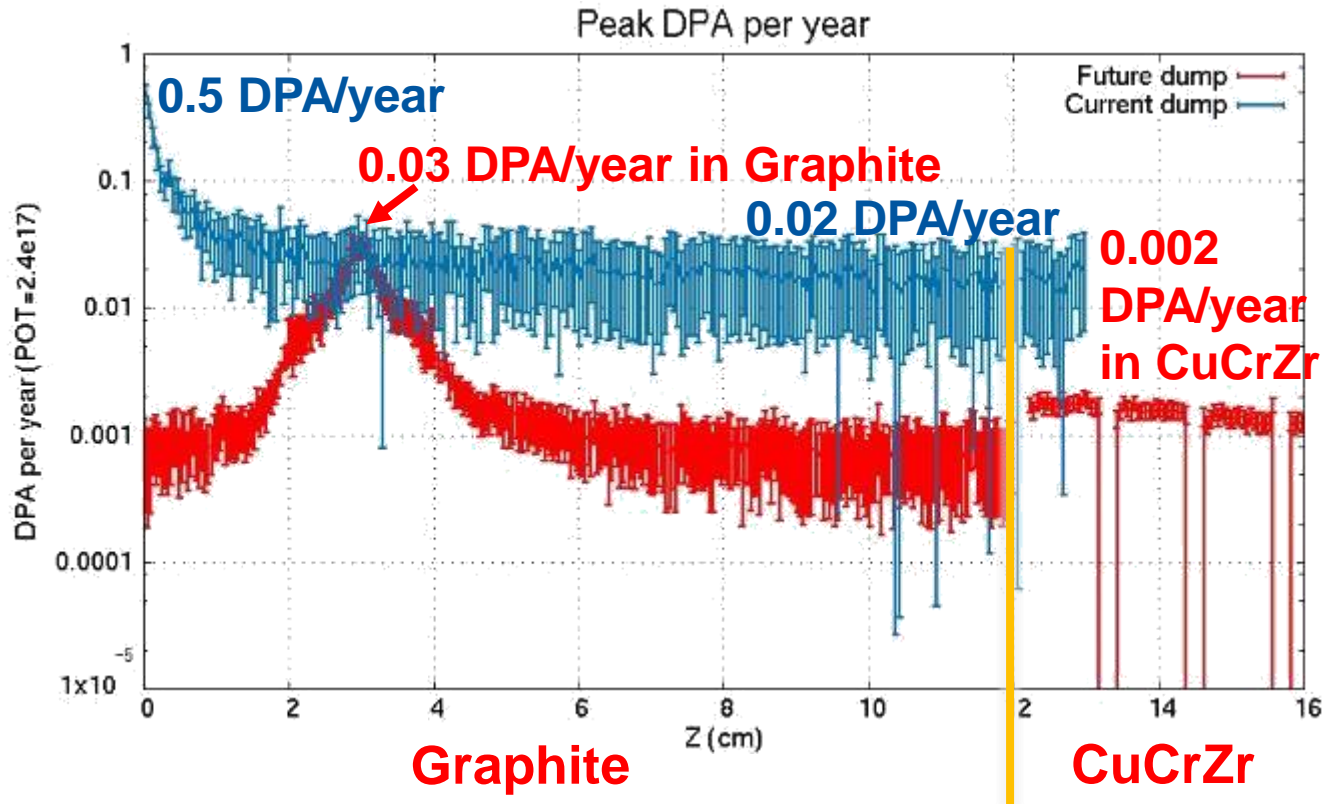


HL-LHC beam

$p=26$ GeV/c
 2.4×10^{13} ppp
 $\sigma_x \times \sigma_y = 1.74$ mm x 0.87 mm
 $\epsilon_x = 1.8$ mm mrad
 $\epsilon_y = 1.8$ mm mrad



Structural Damage on Dump Core



Lower general damage than current dump
(~1 order of magnitude)

Lower peak DPA on copper region than current dump
(~2 orders of magnitude)

Graphite protects the copper block from structural damage

Damage threshold energies considered:
 $E_{th}(\text{Graphite}) = 30 \text{ eV}$ - typical value 30-35 eV
 $E_{th}(\text{CuCrZr and SS304L}) = 40 \text{ eV}$

POT=2.4e17 (assumed same POT for current and future dumps)

Irradiation on Graphite

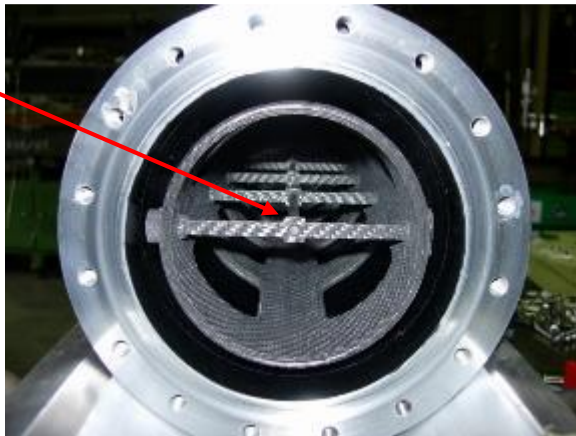
0.03 DPA/year estimated in Graphite block of PS dump

Experience at CERN: **CNGS air cooled graphite target (SPS beam)**

- About 1200°C reached for each pulse
- At the end of operation: **1.5 DPA**
- **No problem observed on graphite**

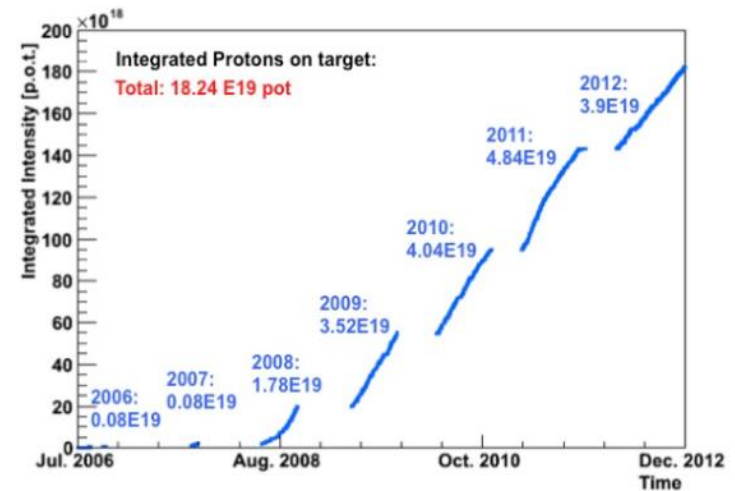
- 3.5×10^{13} protons per pulse,
- **10.5 μs pulse length < 1 mm spot size**
- 2 extractions per cycle separated by 50 ms, **occurring every 6 s**
- 2 000 000 extractions achieved by end of 2009
- 4.5×10^{19} protons at 400 GeV/c on CNGS target **per year**

Graphite rods
2020PT
(Mersen)



[Ref]: Spallation materials R&D for CERN's fixed target program, M. Calviani et al. IWSMT, Oct. 2014, Austria

CNGS 2006 - 2012
Total Integrated Protons on Target: 18.24 E19 pot



PS dump graphite irradiation shall not be a concern

Irradiation on CuCrZr

- Peak value obtained of 0.002 DPA per year (**0.04 DPA in 20 years**) in CuCrZr block of PS Dump
- **Localized peak DPA**
- Information for **neutron irradiation** found in literature
- CuCrZr shows **radiation hardening** until saturation values around 0.1 – 0.5 DPA [1][3] → Some hardening may occur
- CuCrZr is **void swelling resistant** [1][2] (below 2% density change for up to 150 DPA [1])
- Some **thermal conductivity** degradation may occur (5 – 10 % reduction for doses > 0.1 DPA at < 150 °C [2])

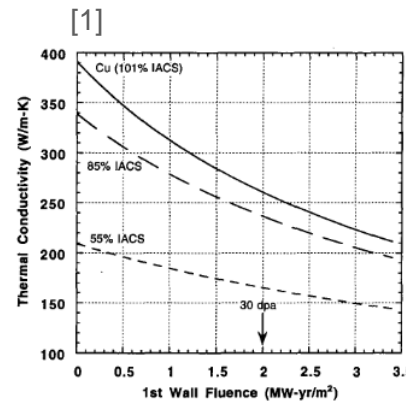
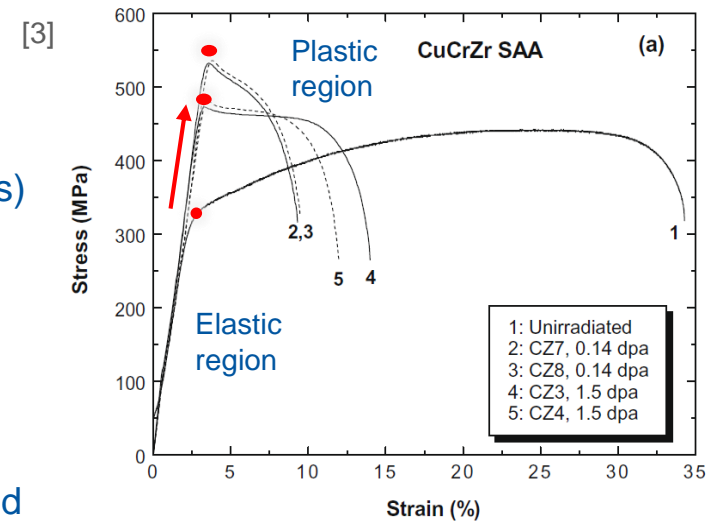


FIG. 21. Effect of solid transmutations on the thermal conductivity of copper at a fusion reactor first wall. The calculations were performed for three different initial conductivities, expressed in terms of percent International Annealed Copper Standard (IACS).

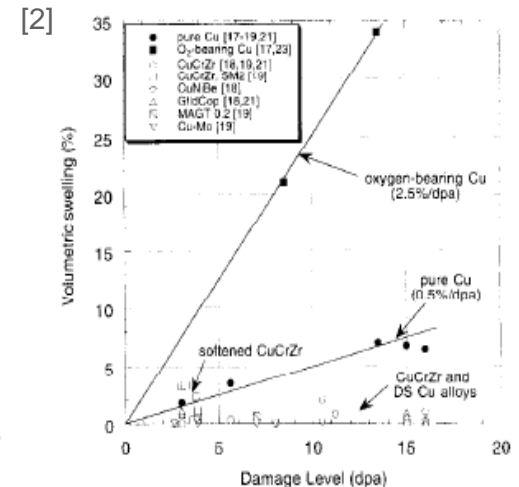


Fig. 3. Irradiation-induced swelling of neutron-irradiated copper and copper alloys. (a) Temperature dependence of void swelling in copper [20]. (b) Dose dependence of volumetric swelling in pure copper and copper alloys at ~ 400°C [17–19,21,23].

[1] C. Bobeldijk (ed.). (1994). Atomic and Plasma-Material Interaction Data for Fusion. Vol. 5. Supplement to the Journal Nuclear Fusion

[2] S.A. Fabritsiev & S.J. Zinkle & B.N. Singh. (1996). Evaluation of copper alloys for fusion reactor divertor and first wall components. Journal of Nuclear Materials. Vol. 233-237. pp. 127-137.

[3] M. Li & M.A. Sokolov & S.J. Zinkle. (2009). Tensile and fracture toughness properties of neutron-irradiated CuCrZr. Journal of Nuclear Materials. Vol. 393. pp. 36-46.