

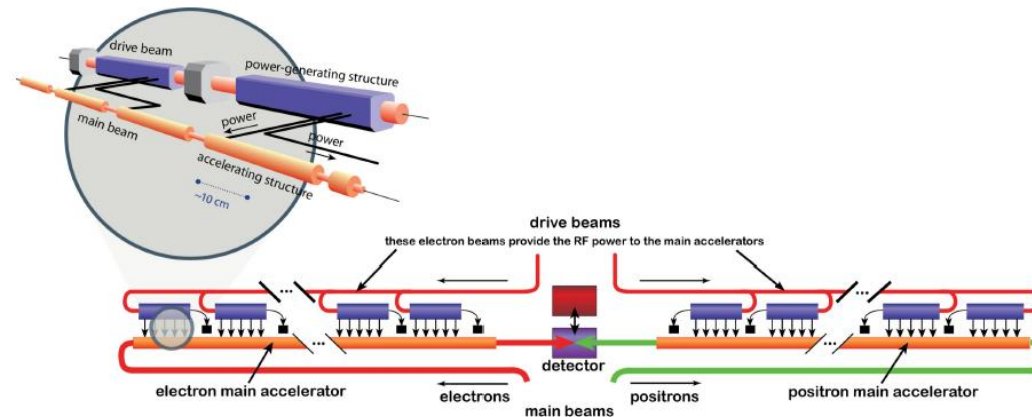
Linear collider physics at 380 GeV

Igor Boyko
(CLICdp/Dubna)

Projects of future e+e- colliders

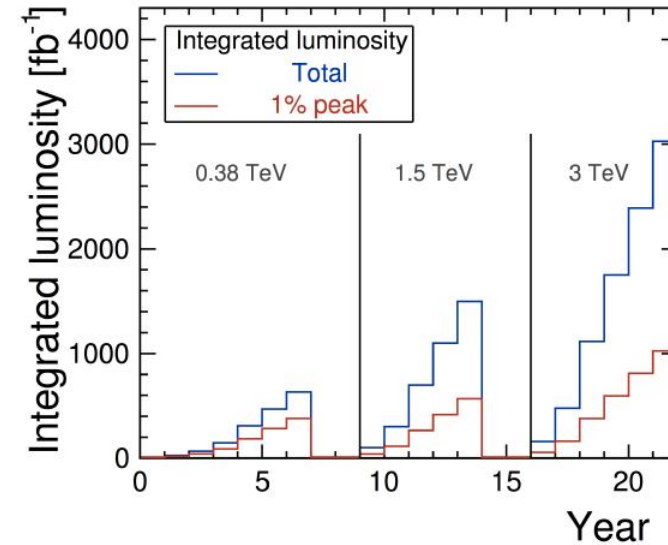
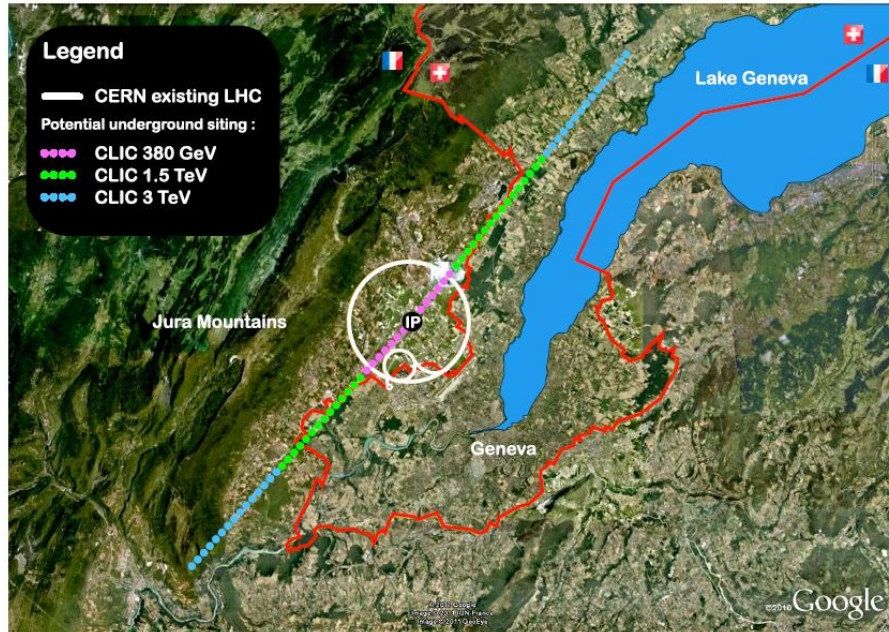
- Future Circular Collider (FCC-ee)
 - $\sqrt{s} = 91\text{-}350 \text{ GeV}$
- Circular Electron-Positron Collider (CEPC)
 - $\sqrt{s} = 90\text{-}240 \text{ GeV}$
- International Linear Collider (ILC)
 - $\sqrt{s} = 250\text{-}500 \text{ GeV (1 TeV)}$
- Compact Linear Collider (CLIC)
 - $\sqrt{s} = 380 \text{ GeV} / 1.5 \text{ TeV} / 3 \text{ TeV}$

CLIC (Compact Linear Collider)



- 2-beam acceleration scheme
 - Low energy: 2.4 GeV → 240 MeV, 100 A
 - High energy: 9 GeV → 1500 GeV, 1.2 A
- Gradient 100 MV/m, operated at room temperature
- Energy: 380-3000 GeV (staged)
- Length: 50 km (for 3 TeV)

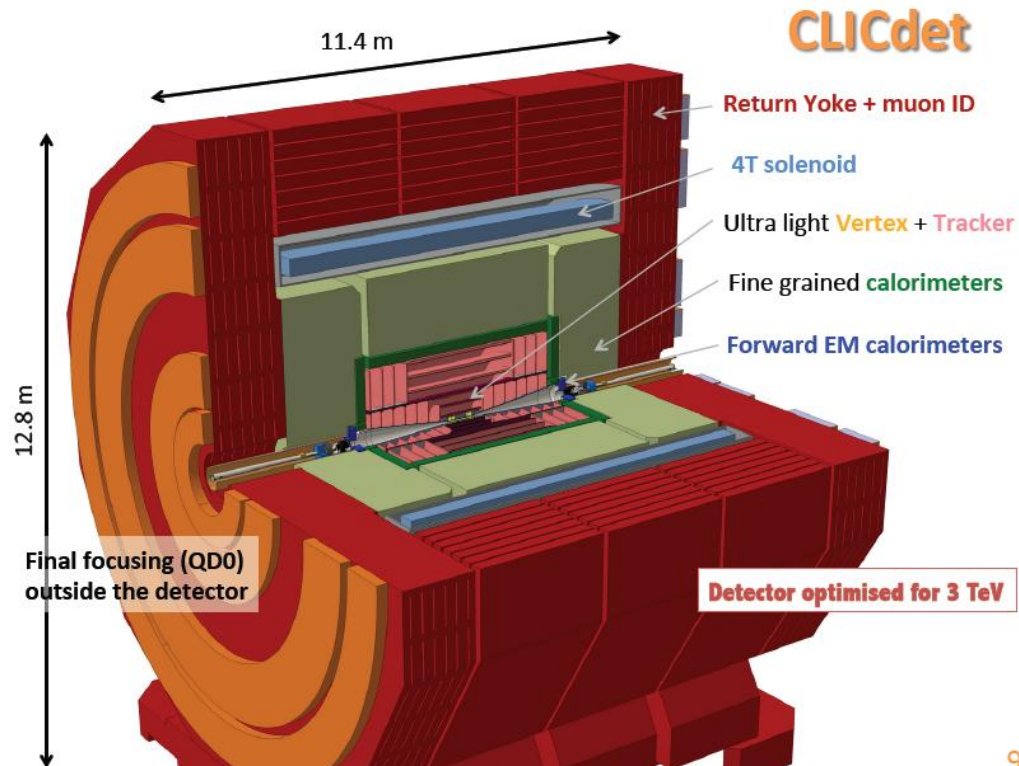
CLIC staging scenario



- Why 380 GeV?
 - Cost consideration (staging)
 - Unique and wonderful physics at 380 GeV!
 - Some benchmark results were obtained with “old” staging scenario, **350/1400/3000 GeV**

Stage	\sqrt{s} (GeV)	\mathcal{L}_{int} (fb^{-1})
1	380	500
	350	100
2	1500	1500
3	3000	3000

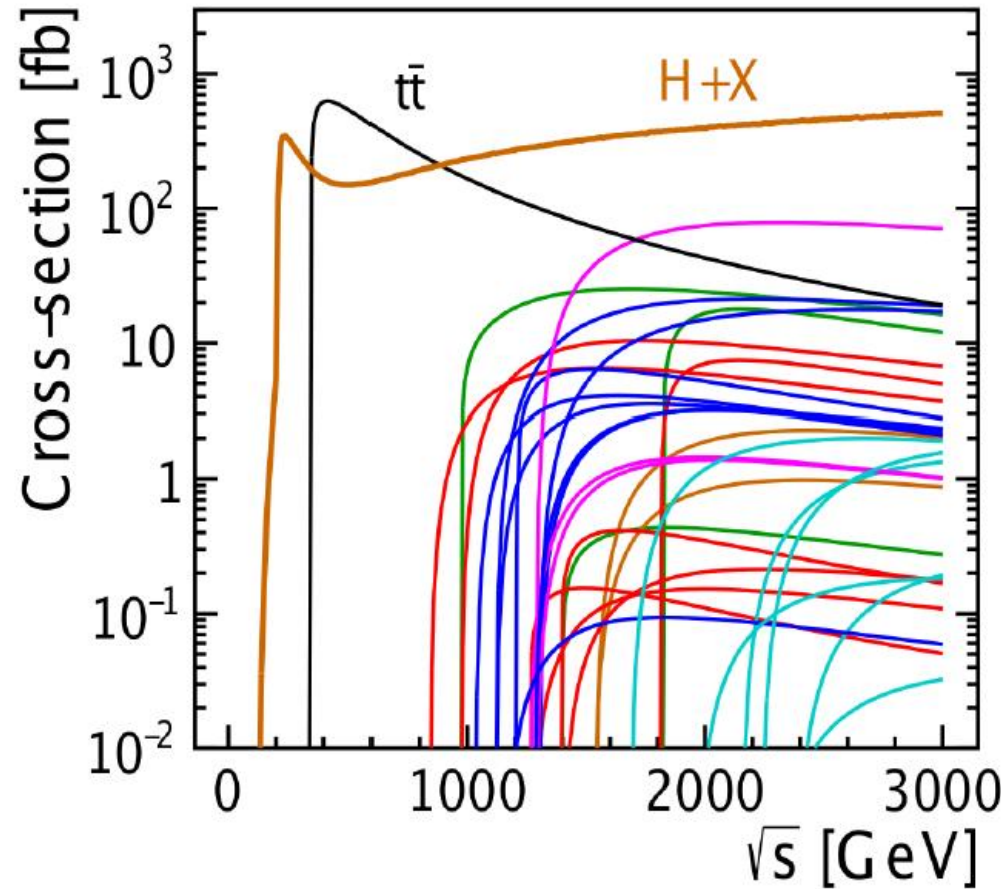
CLIC detector



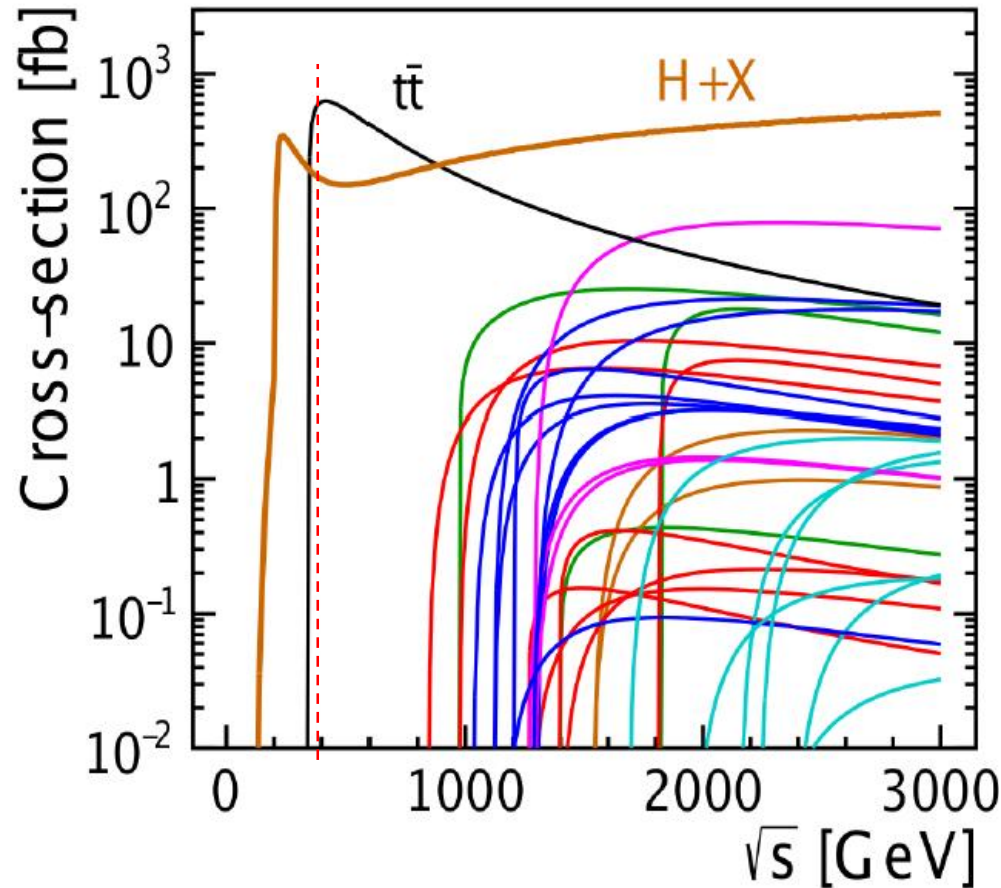
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- New detector model **CLICdet**
- Replaces old options **CLIC_ILD** and **CLIC_SiD**
- Many benchmark results are still based on the old options
- **See the detailed talk this afternoon!**

CLIC physics versus energy

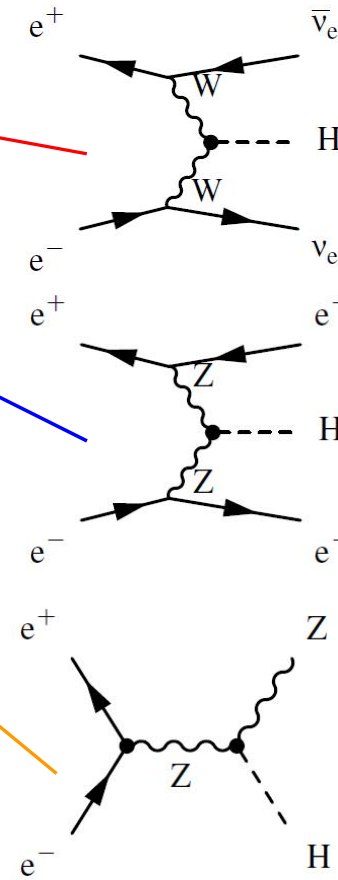
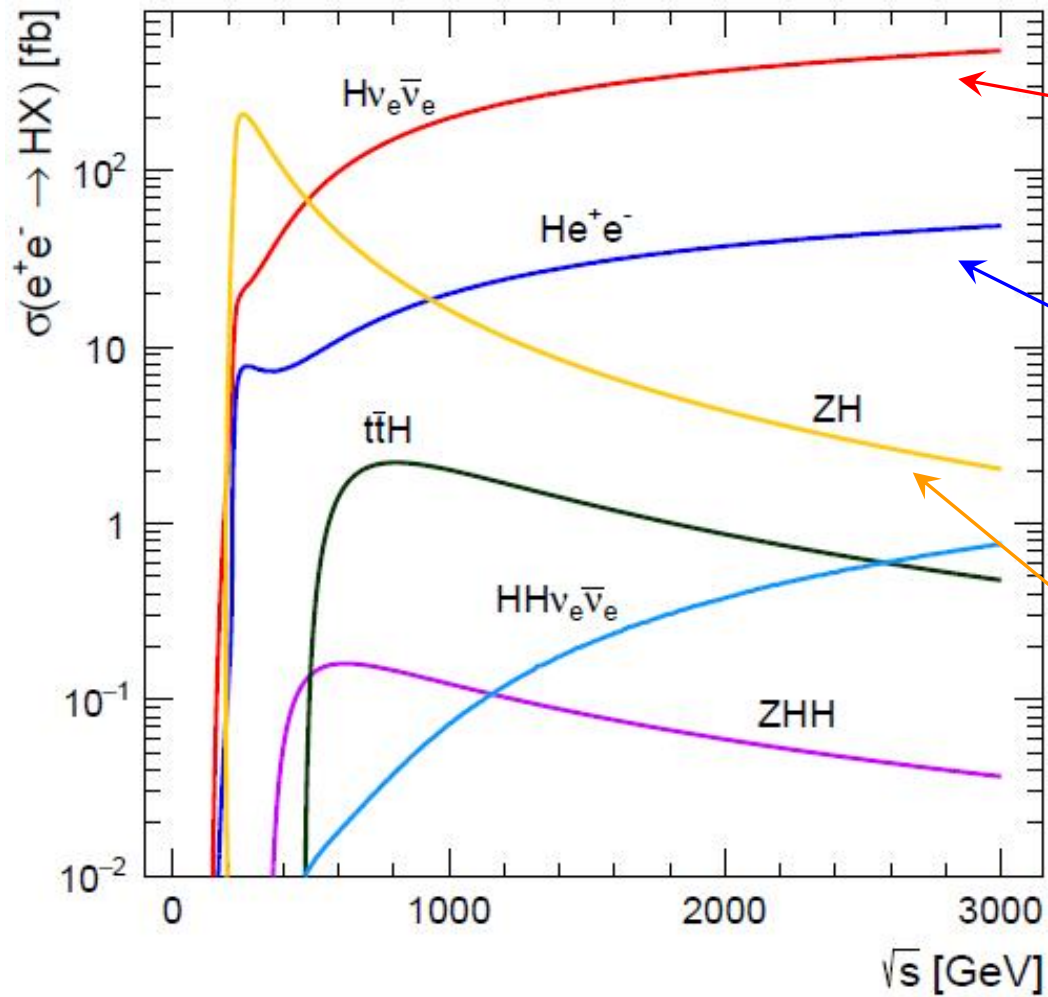


CLIC physics versus energy



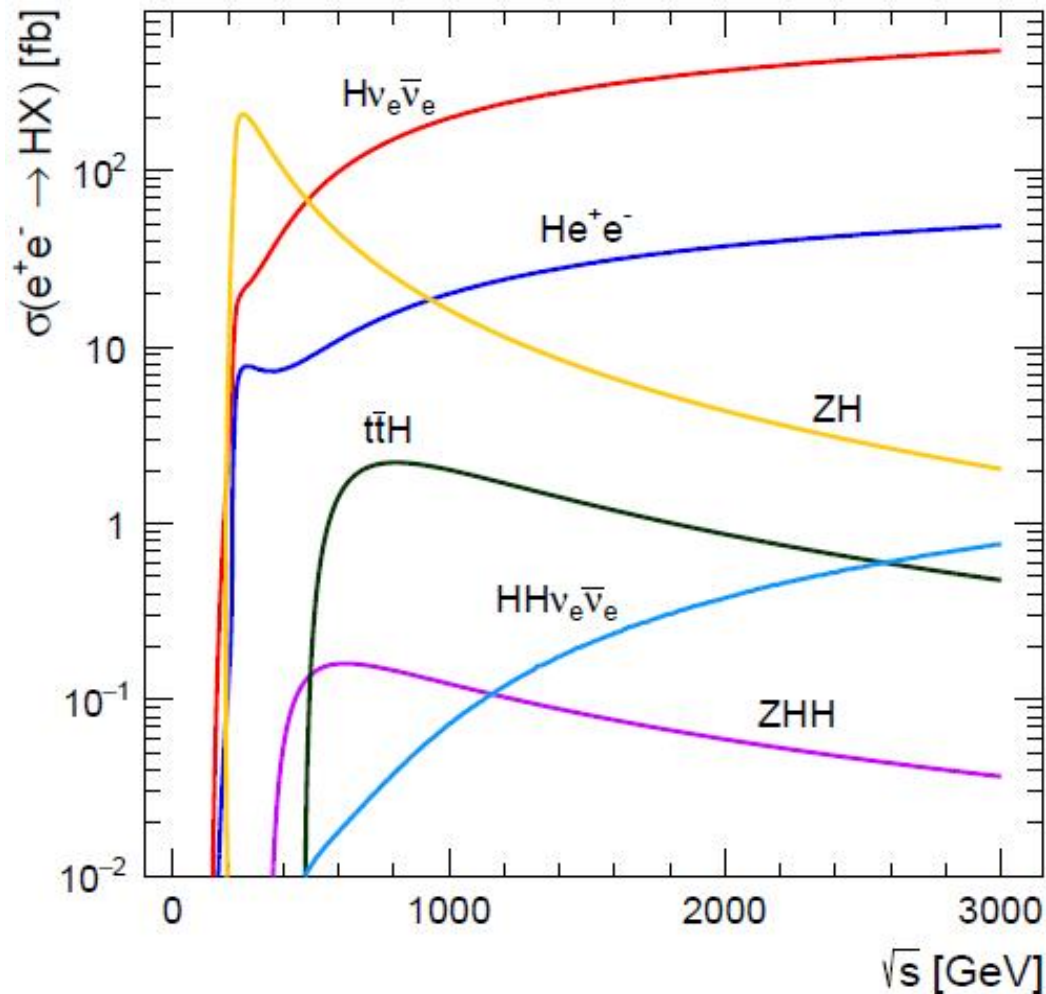
- BSM physics: direct searches require the highest possible energies
 - Note: deviations from the Standard Model can be observed in lower energy precision measurements
- Initial stage at **380 GeV** is optimized for Higgs and Top physics (including $t\bar{t}$ -threshold scan at **350 GeV**)

Higgs production at CLIC energies



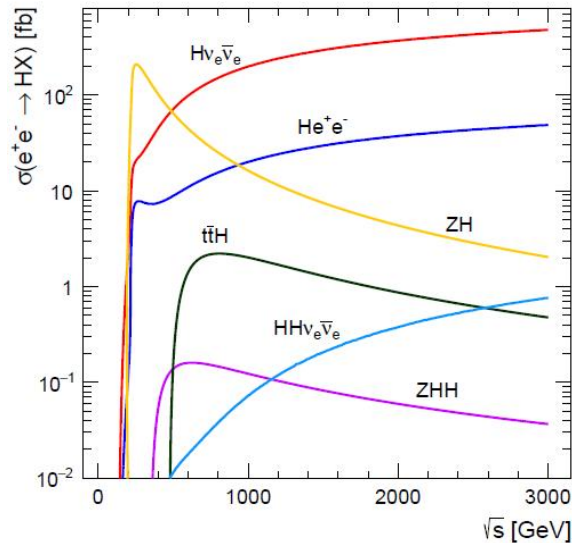
Higgsstrahlung!

Higgs production at CLIC energies



- For CLIC energies, the maximum Higgs cross-section is at **3 TeV**, dominated by $ee \rightarrow H\nu\nu$ mode
- At **250-500 GeV** Higgs production is dominated by Higgsstrahlung $ee \rightarrow ZH$
- It offers the unique opportunity of model-independent measurement, where Higgs is predicted as a system recoiling against the Z^0
- Although ZH cross-section peaks at ~ 250 GeV, the optimal precision is reached around 350 GeV

Expected Higgs statistics



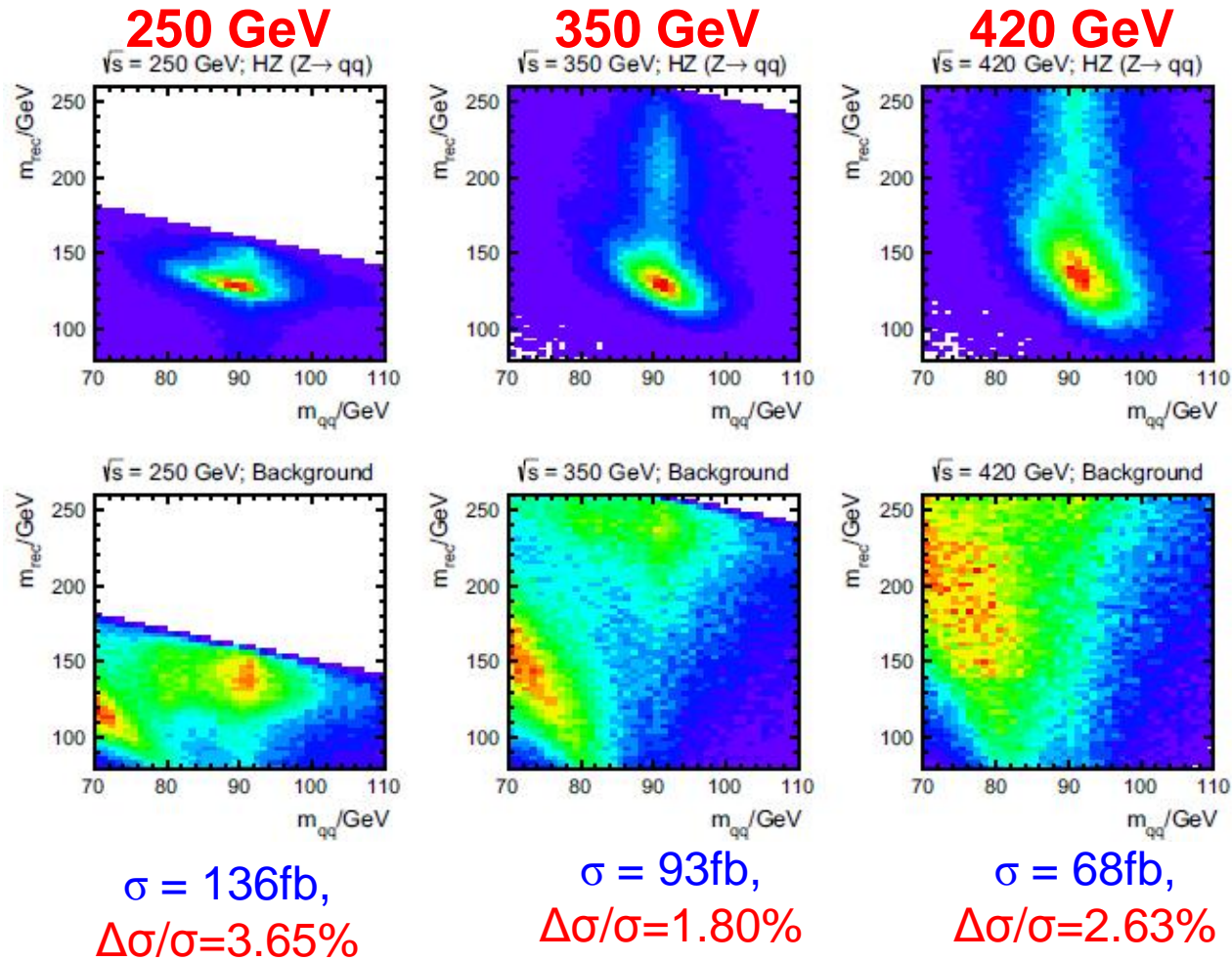
\sqrt{s}	350 GeV	1.4 TeV	3 TeV
\mathcal{L}_{int}	500 fb^{-1}	1.5 ab^{-1}	2 ab^{-1}
$\sigma(e^+e^- \rightarrow ZH)$	133 fb	8 fb	2 fb
$\sigma(e^+e^- \rightarrow H\nu_e\bar{\nu}_e)$	34 fb	276 fb	477 fb
$\sigma(e^+e^- \rightarrow H e^+e^-)$	7 fb	28 fb	48 fb
# HZ events	68,000	20,000	11,000
# $H\nu_e\bar{\nu}_e$ events	17,000	370,000	830,000
# $H e^+e^-$ events	3,700	37,000	84,000

$P(e^-/e^+) = -80/+30\%$: (HZ)x1.4, (H $\nu\nu$)x2.34

- The really big statistics will be provided by $ee \rightarrow H\nu\nu$ at high energy running
- However, the ZH events at 350/380 GeV will provide a unique sample of “tagged” Higgses
 - Possibility of model-independent measurements

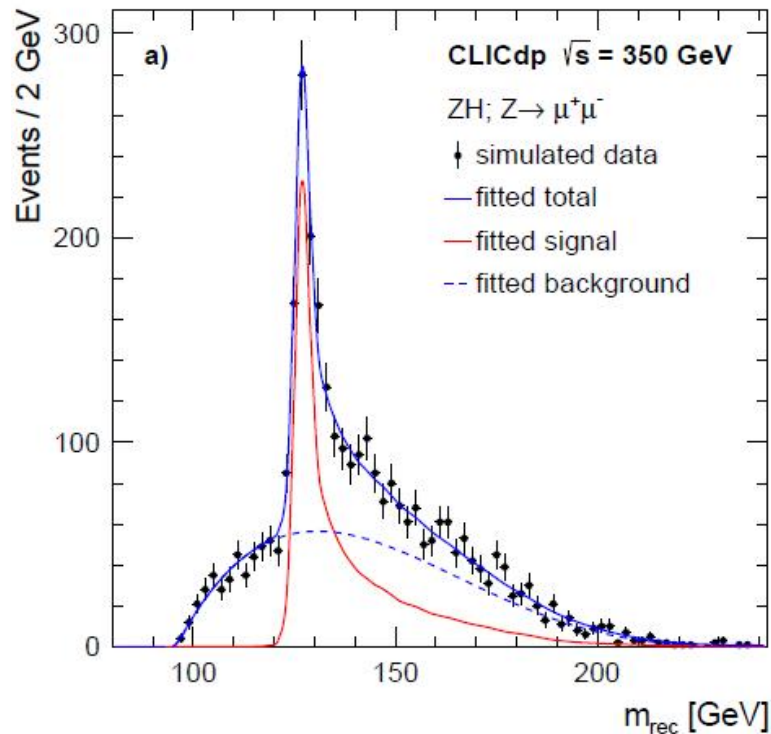
ee \rightarrow HZ: energy optimization

Thomson, EPJ C76:72



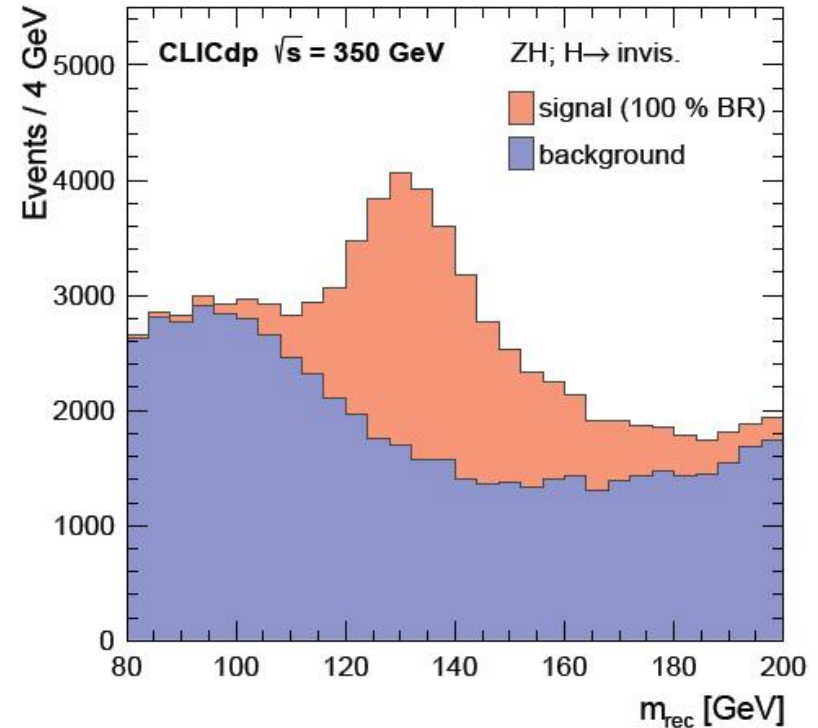
ZH with Z^0 reconstruction

M_H from recoil mass



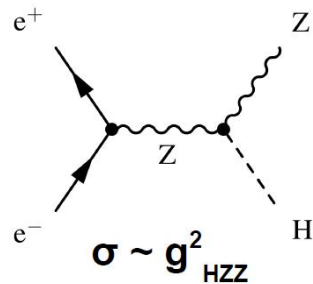
$\sigma(M_H) = 110$ MeV

Invisible Higgs decays



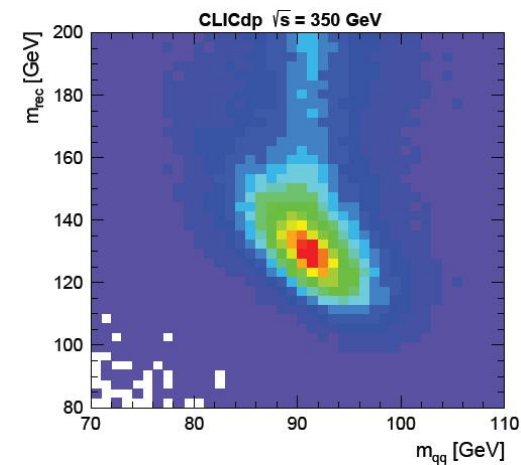
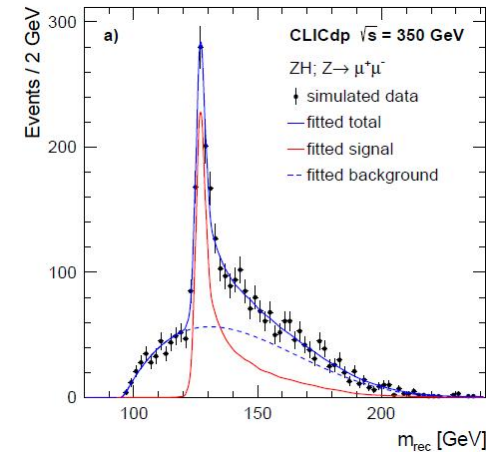
$BR(H \rightarrow \text{invis}) < 0.97\%$ (90%CL)
[SM: 0.1%]

Model-independent measurement of HZZ coupling

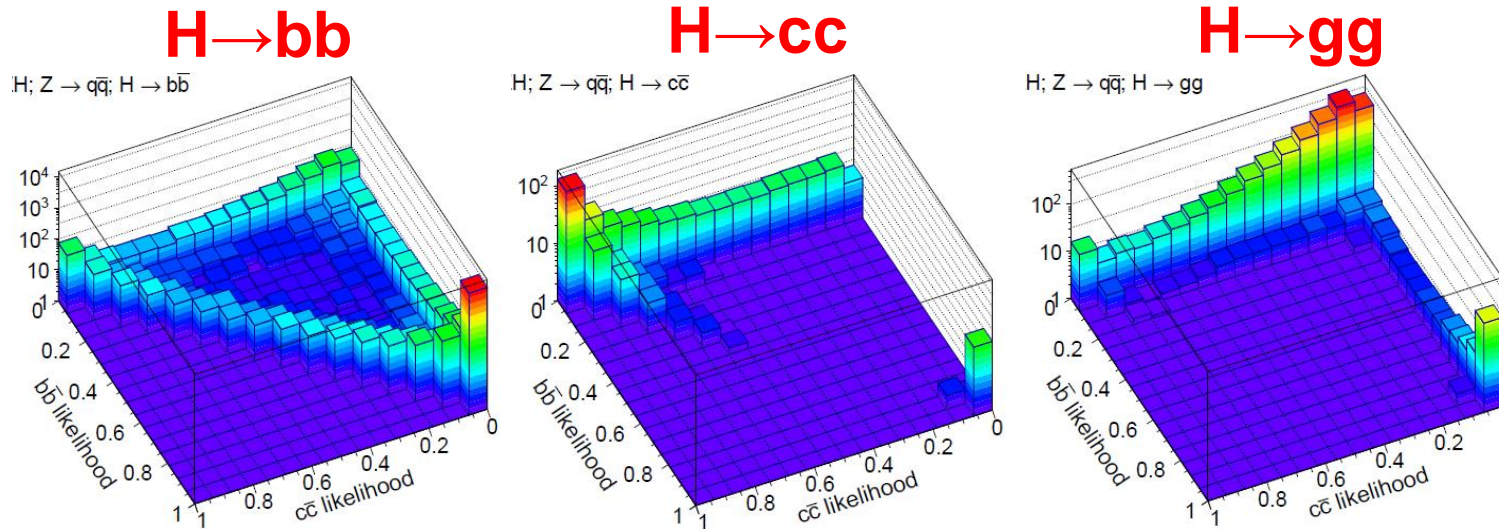


- Higgs is “tagged” rather than reconstructed
- Selection efficiency nearly independent from the Higgs decay mode
- Combining $Z \rightarrow \ell\ell$ and $Z \rightarrow qq$:

$$\frac{\Delta\sigma(ZH)}{\sigma(ZH)} = 1.65\% \quad \frac{\Delta g_{HZZ}}{g_{HZZ}} = 0.8\%$$



Higgs hadronic branchings



- Hadronic decay preselection
 - $ZH \rightarrow 4\text{jets}$
 - $(ZH, H\nu\nu) \rightarrow 2\text{jets} + \text{missing } p_T$
- $bb/cc/gg$ jets are separated using the MVA technique based on flavour tagging

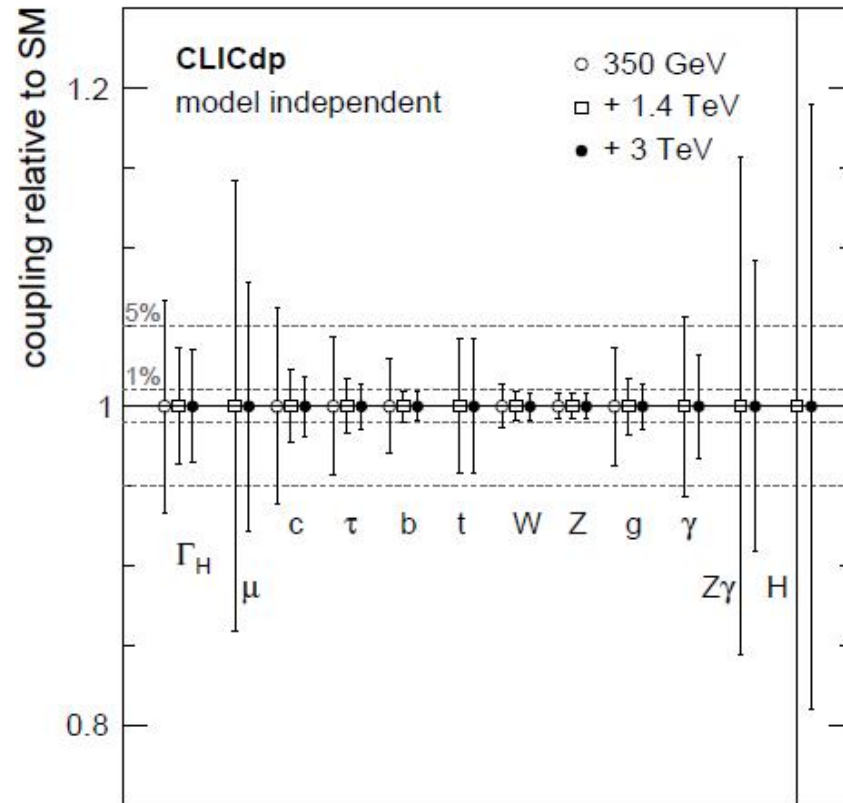
Decay	Statistical uncertainty	
	Higgsstrahlung	WW-fusion
$H \rightarrow b\bar{b}$	0.86 %	1.9 %
$H \rightarrow c\bar{c}$	14 %	26 %
$H \rightarrow gg$	6.1 %	10 %

Model-independent fit to Higgs couplings

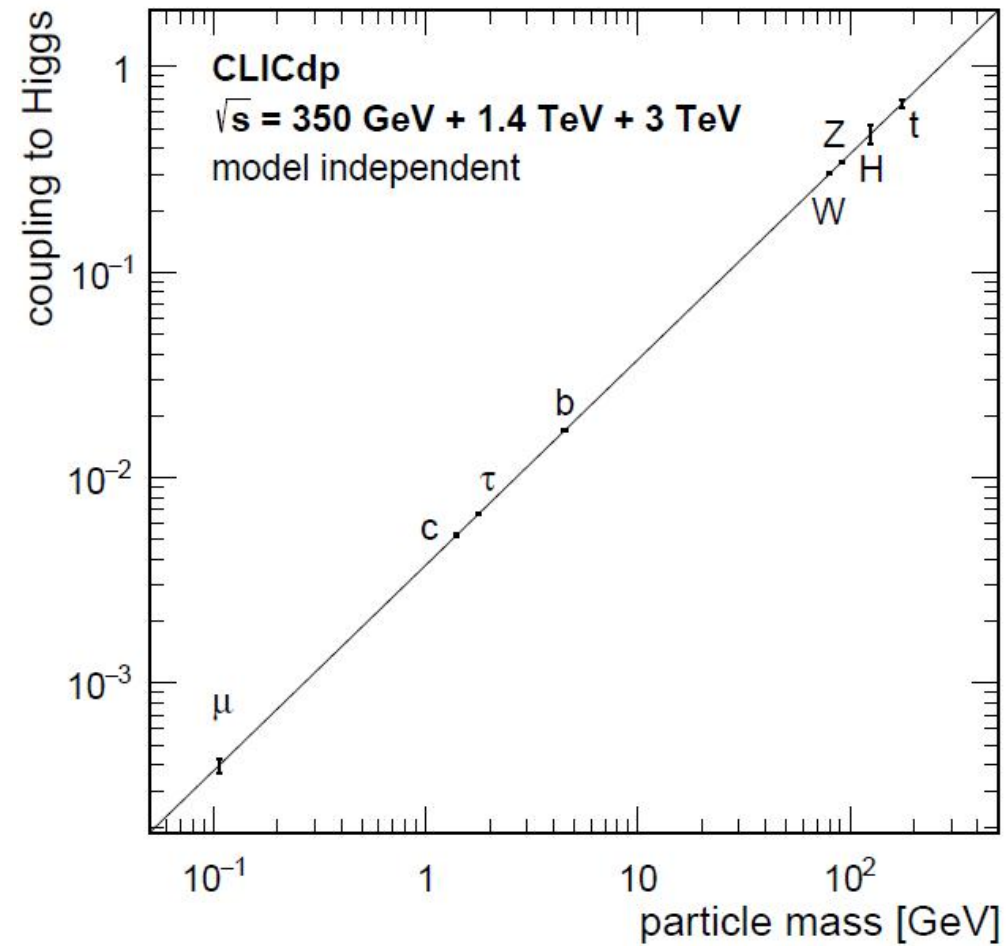
$$\frac{\sigma(e^+e^- \rightarrow ZH) \times BR(H \rightarrow b\bar{b})}{\sigma(e^+e^- \rightarrow \nu_e\bar{\nu}_e H) \times BR(H \rightarrow b\bar{b})} \propto \left(\frac{g_{HZZ}}{g_{HWW}} \right)^2$$

$$\sigma(H\nu_e\bar{\nu}_e) \times BR(H \rightarrow WW^*) \propto \frac{g_{HWW}^4}{\Gamma_H}$$

Parameter	Relative precision		
	350 GeV 500 fb ⁻¹	+ 1.4 TeV + 1.5 ab ⁻¹	+ 3 TeV + 2 ab ⁻¹
g_{HZZ}	0.8 %	0.8 %	0.8 %
g_{HWW}	1.4 %	0.9 %	0.9 %
g_{Hbb}	3.0 %	1.0 %	0.9 %
g_{Hcc}	6.2 %	2.3 %	1.9 %
$g_{H\tau\tau}$	4.3 %	1.7 %	1.4 %
$g_{H\mu\mu}$	—	14.1 %	7.8 %
g_{Htt}	—	4.2 %	4.2 %
g_{Hgg}^\dagger	3.7 %	1.8 %	1.4 %
$g_{H\gamma\gamma}^\dagger$	—	5.7 %	3.2 %
$g_{HZ\gamma}^\dagger$	—	15.6 %	9.1 %
Γ_H	6.7 %	3.7 %	3.5 %



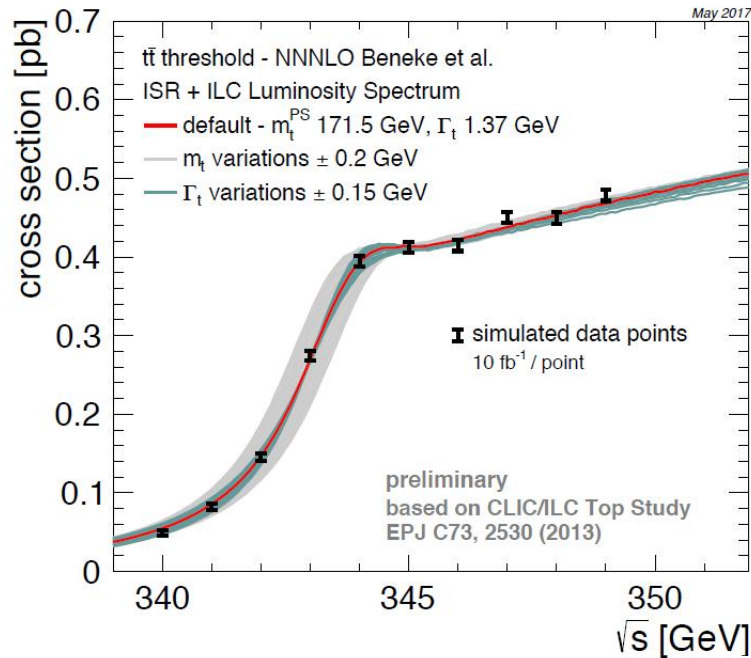
Higgs coupling precision



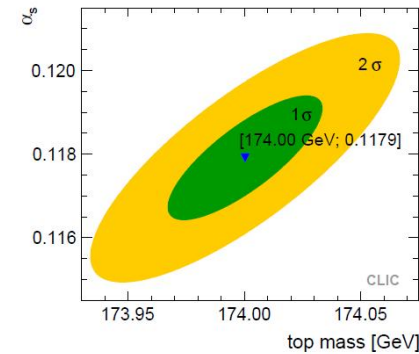
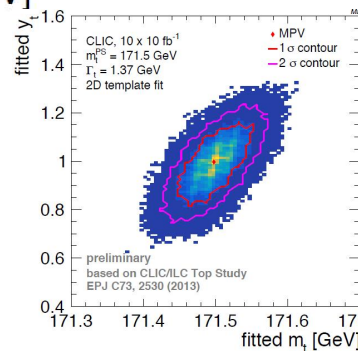
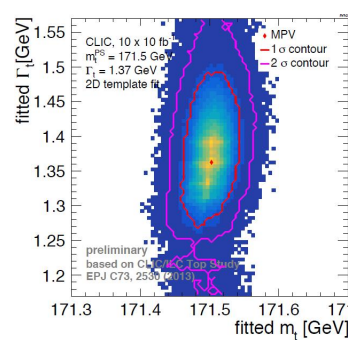
Top physics at 350-380 GeV

- Top mass (and width) from threshold scan
- Top mass from jet reconstruction
- Rare Top decays
- Precision electroweak couplings

Top threshold scan

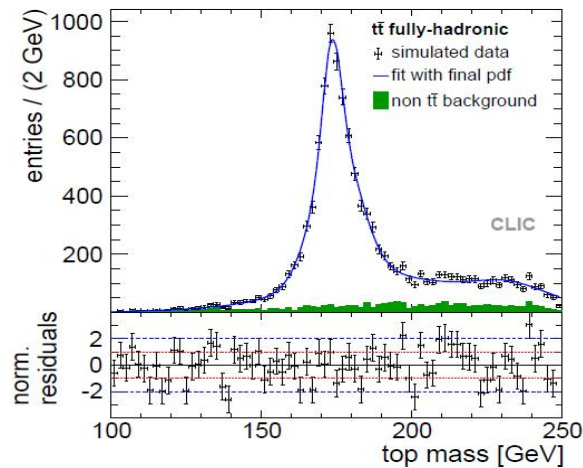


- 10 points, 10fb $^{-1}$ each (1 year running at 350 GeV)
- Cross-section curve directly sensitive to the top mass, width, Yukawa coupling, strong coupling constant
- $\delta(m_t) = \pm 20$ MeV(stat)
 ± 40 MeV(syst) ± 40 MeV(scale)
- $\delta(\Gamma_t) = \pm 45$ MeV(stat)
 ± 60 MeV(scale)



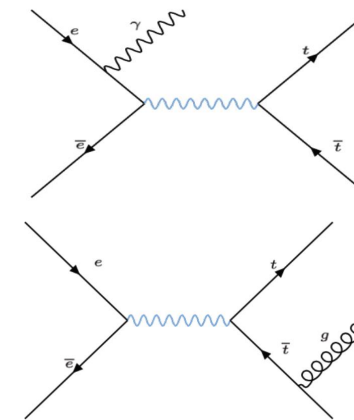
Top mass above threshold

Invariant mass reconstruction



- Systematics-dominated measurement
- Experimental systematic error on m_t about ± 80 MeV, dominated by Jet Energy Scale
- Similar statistical error reached already with 100 fb^{-1}
- Additional theoretical errors from scale and colour reconnection uncertainties

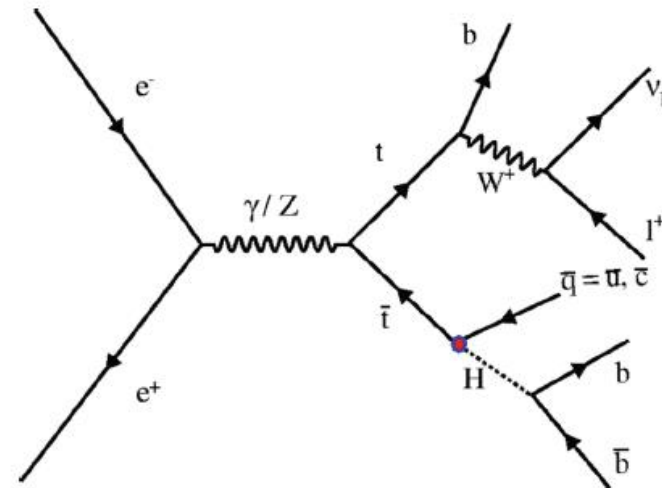
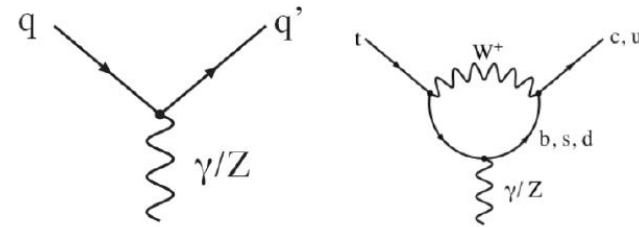
Radiative events



- Idea: measure differential cross-section of ISR (extra photon) or FSR (extra jet)
- Work in progress
- Very preliminary estimate: expect m_t precision of the order of ± 100 MeV

Rare top decays

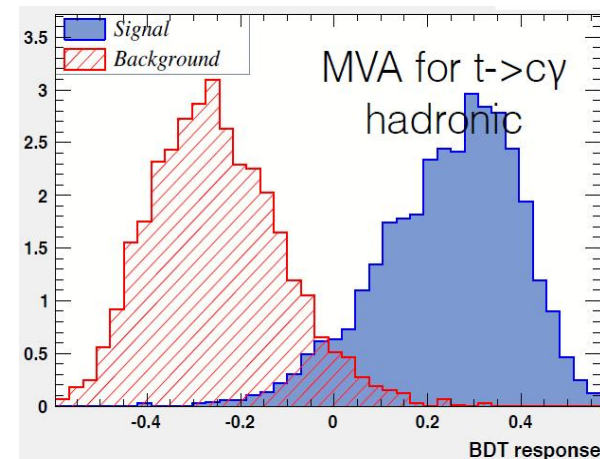
- The FCNC decays $t \rightarrow c\gamma/cZ/cg/cH$ have negligible branchings in the Standard Model (10^{-12} - 10^{-14})
- Currently 2 channels are under study: $t \rightarrow c\gamma$ and $t \rightarrow cH$
- Signal: for $ee \rightarrow tt$ one top decays anomalously, another decay is standard, $t \rightarrow Wb$



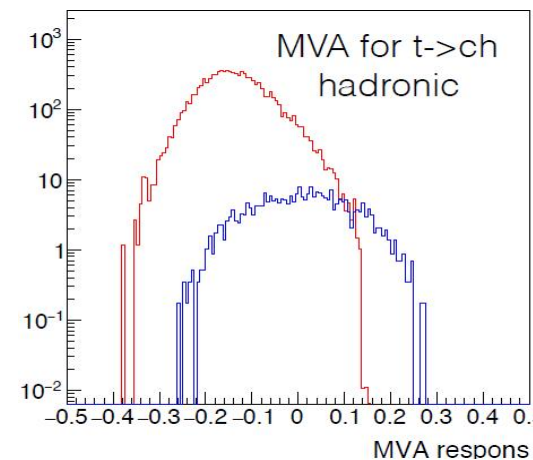
Rare top decays

very preliminary results

- $t \rightarrow c\gamma$
 - CLIC sensitivity:
 $BR(t \rightarrow c\gamma) < 0.5 \cdot 10^{-4}$ (95%CL)
 - Expected HL-LHC: $2 \cdot 10^{-4}$



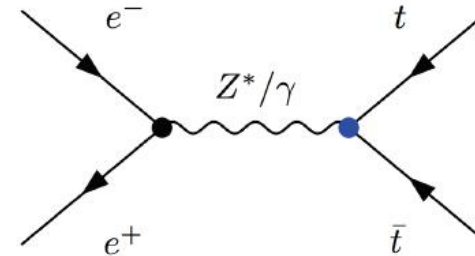
- $t \rightarrow cH$
 - CLIC sensitivity:
 $BR(t \rightarrow cH) < 1.6 \cdot 10^{-4}$ (95%CL)
 - Expected HL-LHC: $2 \cdot 10^{-4}$



Top electroweak couplings

$$\Gamma^{t\bar{t}X}(k^2, q, \bar{q}) = ie \left\{ \underbrace{\gamma_\mu (F_{1V}^X(k^2) + \gamma_5 F_{1A}^X(k^2))}_{\text{vector axial}} - \frac{\sigma_{\mu\nu}}{2m_t} (q + \bar{q})^\nu \underbrace{(iF_{2V}^X(k^2) + \gamma_5 F_{2A}^X(k^2))}_{\text{tensor CPV}} \right\}$$

- Observables to distinguish Z and γ couplings:
 - Energy dependence of the cross-section
 - Forward-backward charge asymmetry
 - Beam polarization asymmetry ($P_{e^-} = \pm 80\%$)
 - Top quark polarization
 - Spin correlation
- Deviations from SM can be parameterized in terms of New Physics, e.g. in the EFT language



$$F_{1,V}^Z - F_{1,V}^{Z,SM} = \frac{1}{2} \left(\underline{C_{\varphi Q}^{(3)}} - \underline{C_{\varphi Q}^{(1)}} - C_{\varphi t} \right) \frac{m_t^2}{\Lambda^2 s_W c_W} = -\frac{1}{2} \underline{C_{\varphi q}^V} \frac{m_t^2}{\Lambda^2 s_W c_W} \quad \varepsilon$$

$$F_{1,A}^Z - F_{1,A}^{Z,SM} = \frac{1}{2} \left(-\underline{C_{\varphi Q}^{(3)}} + \underline{C_{\varphi Q}^{(1)}} - C_{\varphi t} \right) \frac{m_t^2}{\Lambda^2 s_W c_W} = -\frac{1}{2} \underline{C_{\varphi q}^A} \frac{m_t^2}{\Lambda^2 s_W c_W}$$

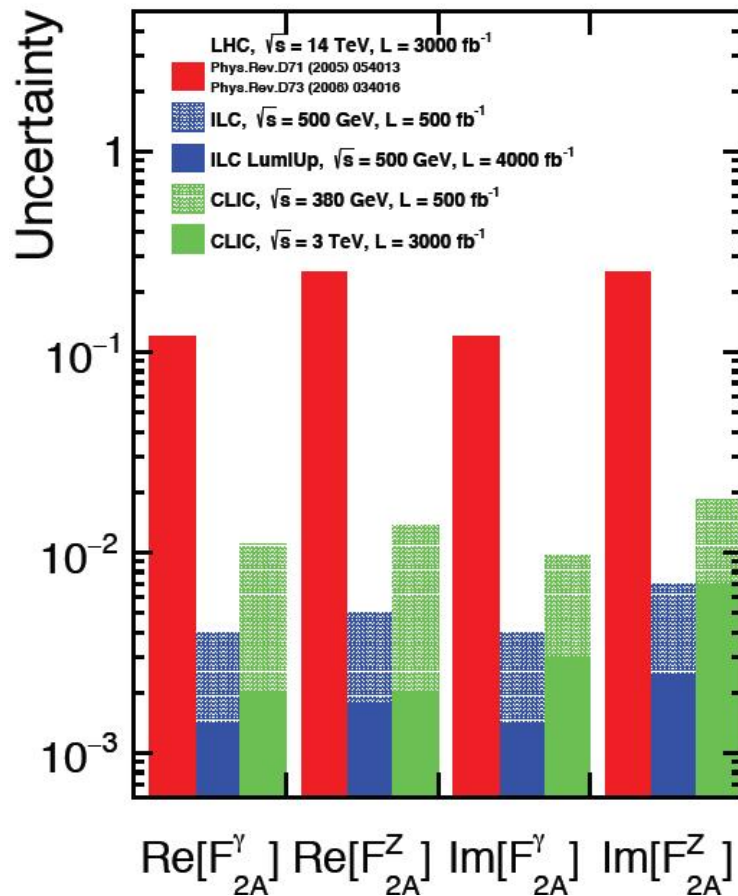
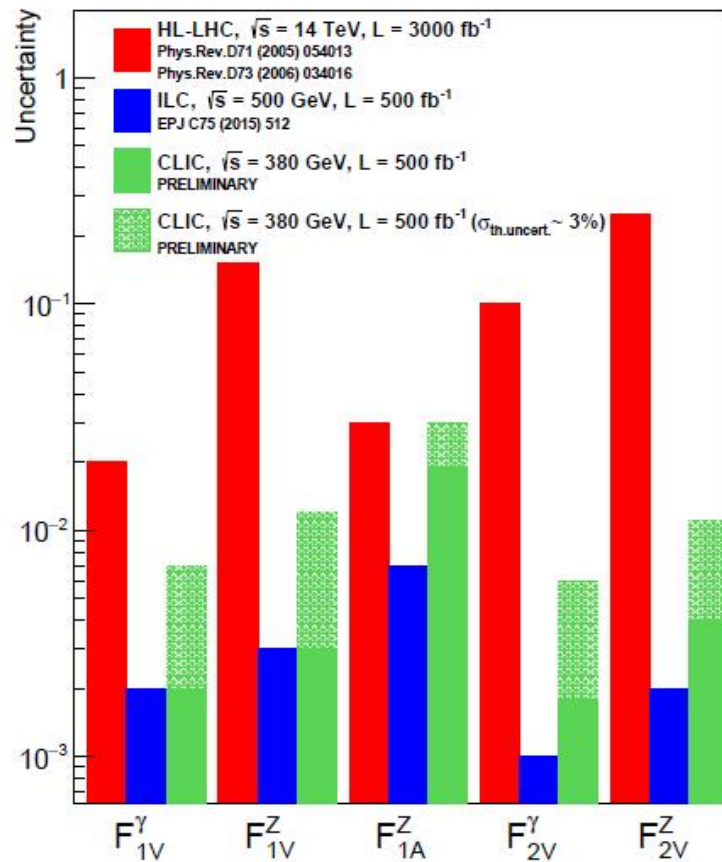
$$F_{2,V}^Z = \left(\underline{\text{Re}\{C_{tW}\}} c_W^2 - \text{Re}\{C_{tB}\} s_W^2 \right) \frac{4m_t^2}{\Lambda^2 s_W c_W} = \text{Re}\{\underline{C_{uZ}}\} \frac{4m_t^2}{\Lambda^2}$$

$$F_{2,V}^\gamma = \left(\underline{\text{Re}\{C_{tW}\}} + \text{Re}\{C_{tB}\} \right) \frac{4m_t^2}{\Lambda^2} = \text{Re}\{\underline{C_{uA}}\} \frac{4m_t^2}{\Lambda^2}$$

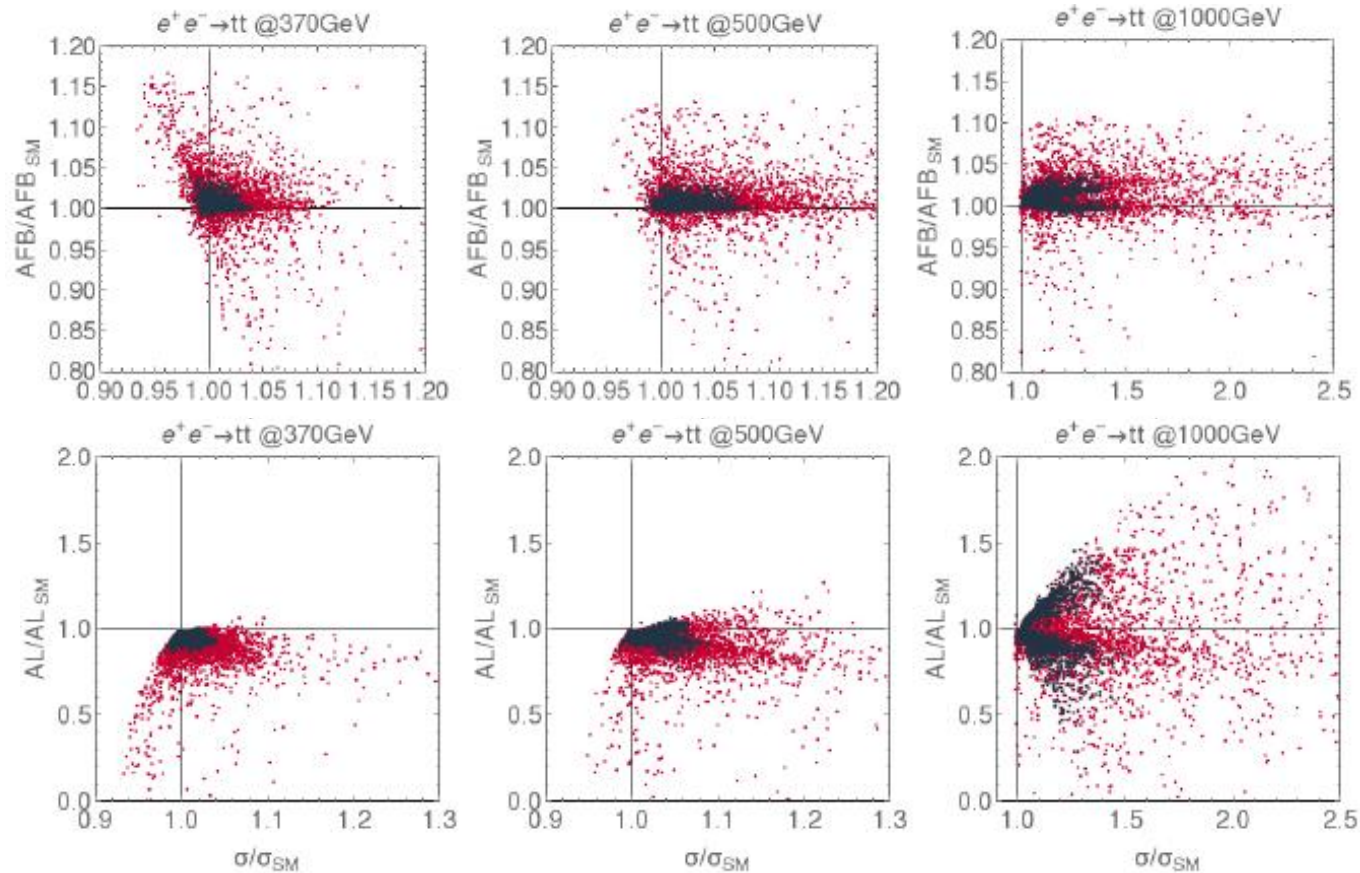
$$[F_{2,A}^Z, F_{2,A}^\gamma] \propto [\text{Im}\{C_{tW}\}, \text{Im}\{C_{tB}\}]$$

Top coupling sensitivity

$$\Gamma^{t\bar{t}X}(k^2, q, \bar{q}) = ie \left\{ \underbrace{\gamma_\mu (F_{1V}^X(k^2) + \gamma_5 F_{1A}^X(k^2))}_{\text{vector axial}} - \frac{\sigma_{\mu\nu}}{2m_t} (q + \bar{q})^\nu \underbrace{(iF_{2V}^X(k^2) + \gamma_5 F_{2A}^X(k^2))}_{\text{tensor CPV}} \right\}$$



Top couplings: composite Higgs scenario



Points: 4DCHM scan with $f=0.75-1.5$ TeV, $g_p=1.5-3$

Barducci et al, 1504.05407

Summary

- CLIC-380 will be the first step on the way to the ultimate energy frontier at 3 TeV
- However, this stage on its own will provide invaluable data for precision measurements
 - Model-independent, low-systematics measurements with ZH sample
 - Top mass with unprecedented precision from the threshold scan
 - Precision top-quark measurements with the huge sample of tt pairs
 - Although direct BSM discoveries require higher energy, the hints of New Physics may be observed as small deviations in precision measurements