

Physics at future e^+e^- colliders

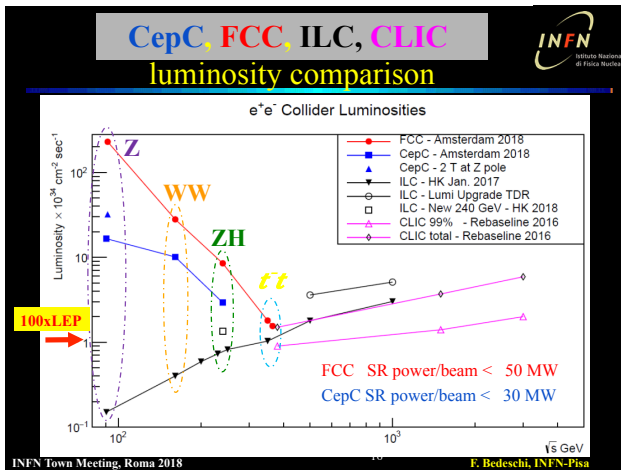
Fulvio Piccinini



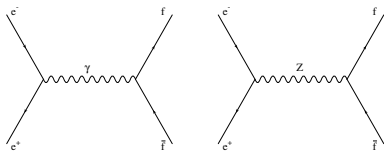
Pisa School on future colliders, Pisa

Pisa, 17-21 September 2018

luminosity comparison among lepton colliders



F. Bedeschi, INFN EU Strategy Meeting, 6-7 September 2018, Roma, Italy



Primary observables

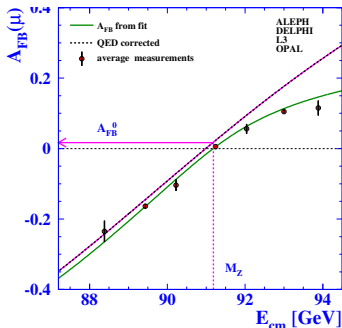
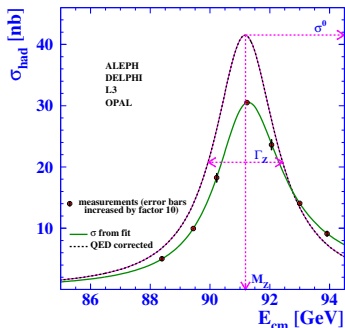
- absolute cross sections for different species of fermions $\sigma_f(s)$
- Forward-backward asymmetries $A_{FB}(s)$ (sensitive to parity violation)

$$\sigma = \sigma_F + \sigma_B$$

$$A_{FB} = \frac{\sigma_F - \sigma_B}{\sigma_F + \sigma_B}$$

$$\sigma_F = 2\pi \int_0^1 d \cos \vartheta \frac{d\sigma}{d\Omega}, \quad \sigma_B = 2\pi \int_{-1}^0 d \cos \vartheta \frac{d\sigma}{d\Omega}$$

measurements @ Z peak at LEP



Large differences due to QED initial state radiation

$$\sigma_{\text{T}}(s) = \int_{z_0}^1 dz H(z; s) \hat{\sigma}_{\text{T}}(zs)$$

$$A_{\text{FB}}(s) = \frac{\pi \alpha^2 Q_e^2 Q_f^2}{\sigma_{\text{tot}}} \int_{z_0}^1 dz \frac{1}{(1+z)^2} H_{\text{FB}}(z; s) \hat{\sigma}_{\text{FB}}(zs)$$

Pseudo-observables

- Idea: extract as much as possible information on the Z properties independently of the event selection details

Observable	Symbol
hadronic peak cross-section	σ_h
partial leptonic and hadronic widths	Γ_l ($l = e, \mu, \tau$), Γ_c, Γ_b
total width	Γ_Z
hadronic width	Γ_h
invisible width	Γ_{inv}
ratios	R_l, R_b, R_c
forward-backward asymmetries	$A_{FB}^l, A_{FB}^b, A_{FB}^c$
polarization asymmetries	P^τ, P^b
left-right asymmetry (SLC)	A_{LR}^e
effective sine	$\sin^2 \vartheta_{eff}^l, \sin^2 \vartheta_{eff}^b$

$$\Gamma_f = 4N_c \Gamma_0 [(g_V^f)^2 R_V^f + (g_A^f)^2 R_A^f]$$

$$\Gamma_{inv} = \Gamma_Z - \Gamma_e - \Gamma_\mu - \Gamma_\tau - \Gamma_h$$

$$R_l = \frac{\Gamma_h}{\Gamma_l}$$

$$R_{b,c} = \frac{\Gamma_{b,c}}{\Gamma_h}$$

$$\sigma_{had}^0 = 12\pi \frac{\Gamma_e \Gamma_h}{M_Z^2 \Gamma_Z^2}$$

$$A_f = \frac{2g_V^f g_A^f}{(g_V^f)^2 + (g_A^f)^2}$$

$$A_{FB}^f = \frac{3}{4} A_e A_f$$

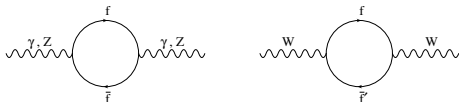
$$A_{LR}^e = A_e$$

$$P^f = -A_f$$

$$4|Q_f| \sin^2 \vartheta_{eff}^f = 1 - \frac{g_V^f}{g_A^f}$$

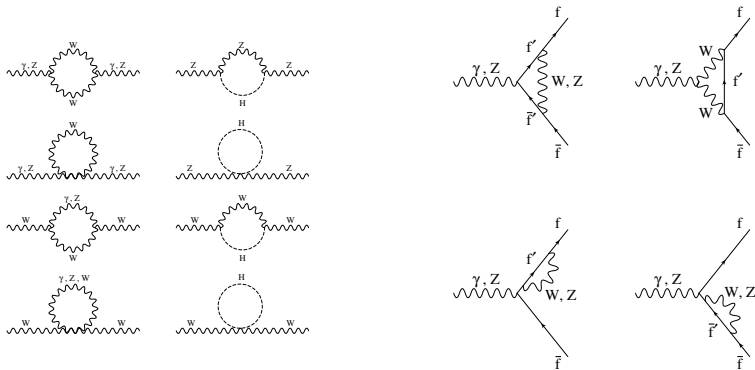
High exp. precision and th. precision

- Exp. precision at the level of 0.01% or better
- How accurate are theoretical predictions? The one-loop example
 - we have to rely on perturbation theory
 - the “theoretical accuracy” is given by the size of the next (not calculated) perturbative order
 - the tree-level approximation has an uncertainty of the order of several %, as proved by the calculated effects at one-loop approximation \implies we need to include higher order effects



- in the loop we have to sum over all possible (also heavy) particles which couple to γ and/or Z

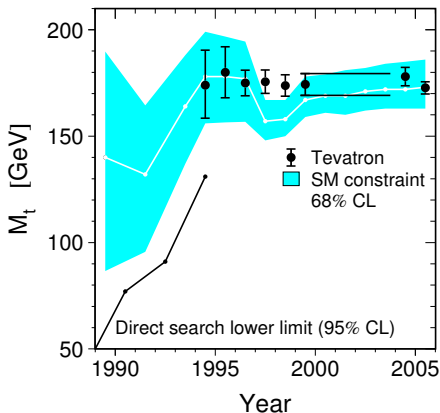
Not only fermions but also bosons...



- Theoretical predictions become sensitive to all the spectrum and structure of the theory: m_t , number of light neutrinos, m_H , non abelian γWW and ZWW vertices
- m_t experimental uncertainty becomes a source of th. uncertainty
 \implies **very important the run at the $t\bar{t}$ threshold**

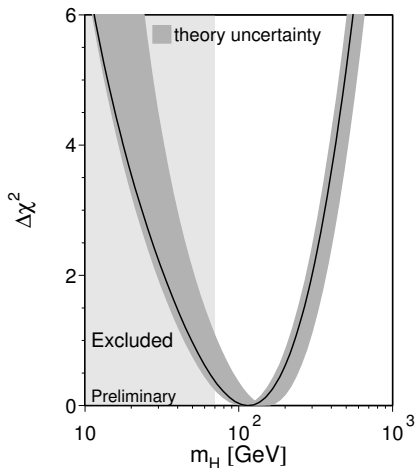
The “indirect” discovery of the quark top at LEP

- before 1995 the dependence on m_t could be used to “measure” m_t from a best-fit to Z -peak data (χ^2 depends quadratically on $G_F(m_t^2 - m_b^2)$ in a gauge theory with SSB)

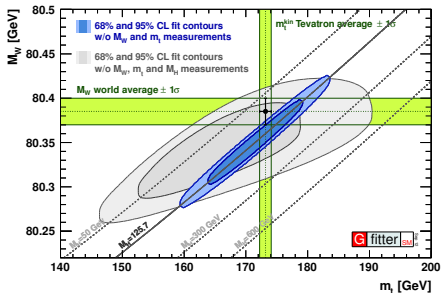
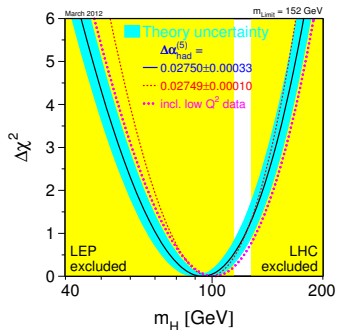


the same could be said about m_H

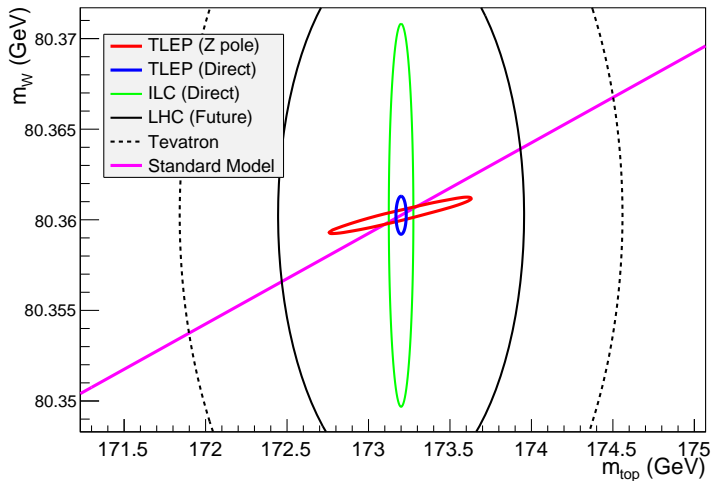
- however, dependence on m_H is only logarithmic
- at the end of LEP1



some years later, after switching on LHC



possible contour 68% plot after FCC-ee running



M. Bicer et al., arXiv:1308.6176

N_ν from Z invisible width

$$R_{\text{inv}}^0 = \frac{\Gamma_{\text{inv}}}{\Gamma_{ll}} = \sqrt{\frac{12\pi R_l^0}{\sigma_{\text{had}}^0 m_Z^2}} - R_l^0 - (3 + \delta_\tau)$$

- assuming lepton universality

$$(R_{\text{inv}}^0)_{\text{exp}} = N_\nu \left(\frac{\Gamma_{\nu\bar{\nu}}}{\Gamma_{ll}} \right)_{\text{SM}}$$

- from LEP Z -peak measurements

$$N_\nu = 2.9840 \pm 0.0082$$
$$\delta N_\nu \simeq 10.5 \frac{\delta n_{\text{had}}}{n_{\text{had}}} \oplus 3.0 \frac{\delta n_{\text{lept}}}{n_{\text{lept}}} \oplus 7.5 \frac{\delta \mathcal{L}}{\mathcal{L}}$$
$$\frac{\delta \mathcal{L}}{\mathcal{L}} = 0.061\% \implies \delta N_\nu = 0.0046$$

ADLO, SLD and LEPEWWG, Phys. Rept. 427 (2006) 257, hep-ex/0509008

- δN_ν severely affected by luminosity uncertainty through σ_0

QED processes: candles for e^+e^- collider luminosity

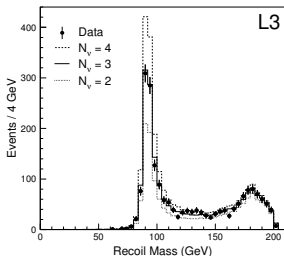
- Luminosity is a machine (process independent) parameter entering every experimental cross-section

$$\frac{N_{obs}}{\mathcal{L}} = \sigma \rightarrow \mathcal{L} = \frac{N_{obs}^X}{\sigma_{theory}^X}$$

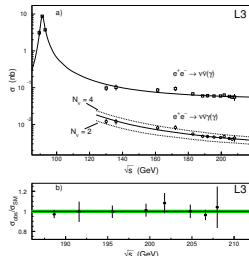
- in order to minimize $\delta\mathcal{L}$, the reference process X has to **have large statistics, be well calculable theoretically, be cleanly detectable** (small systematics)
- the best choice are QED processes, in particular **Bhabha**
 - ★ at small angles
 - huge statistics, by far dominated by photon t -channel contribution (only QED, no “ Z contamination”)
- to reduce the **theoretical error on σ** all the **relevant radiative corrections (RC)** must be included
- at present the theoretical ingredients to reach 0.01% are available

Independent way for ν count: $\nu\bar{\nu}\gamma$ at higher energies

- radiative return to the Z peak through emission of a hard photon
- provided large enough luminosity is available to be competitive with Γ_{inv} method (not a problem at future e^+e^- colliders!)
at LEP2, $190 \text{ GeV} \leq \sqrt{s} \leq 208 \text{ GeV}$, $\mathcal{L} \sim 600 \text{ pb}^{-1}$



L3 Collab., P. Achard et al., CERN-EP/2003-068 (2003)

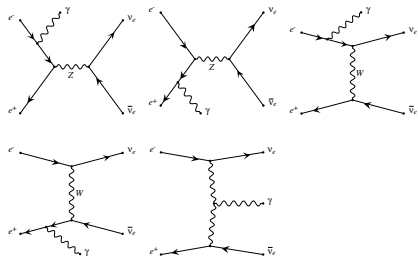


- agreement of data with SM predictions at % level
- $N_\nu = 2.98 \pm 0.05 \pm 0.04$ (L3) (important but not competitive with the Γ_{inv} method)
- similar results for ALEPH, DELPHI and OPAL

- a factor $10^3/10^4$ of improvement in luminosity w.r.t. LEP allows to exploit the ratios

$$\frac{d\sigma(e^+e^- \rightarrow \nu\bar{\nu}\gamma)}{d\sigma(e^+e^- \rightarrow \mu^+\mu^-\gamma)}$$

in order to cancel common systematics (such as luminosity)

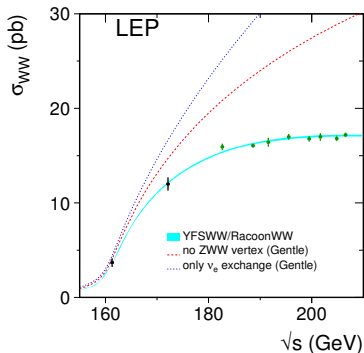
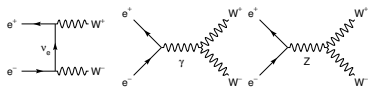


- $\mu^+\mu^-$ only s -channel but ISR and FSR
- ν_μ and ν_τ f.s.: only s -channel ISR
- ν_e f.s.: ISR with t -channel
- ν_e f.s.: also W radiation

- preliminary investigations show that QED effects are very small

talk by S. Jadach at FCC-ee physics Workshop, Paris, 27-29 October 2014

Direct M_W measurement at WW threshold

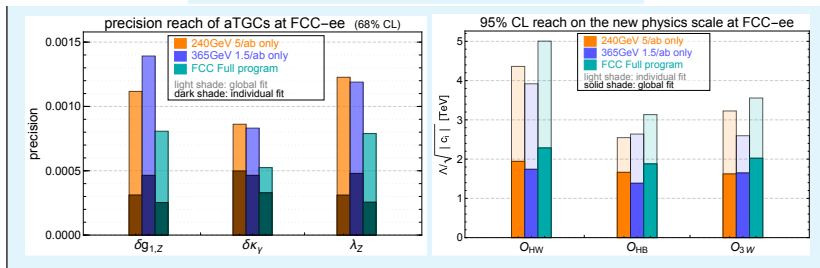


- by using the lineshape dependence on M_W at threshold
 $\implies M_W$ measurement with uncertainty $\Delta M_W \sim 0.6$ MeV
(precision 200 times higher than the best present M_W direct measurement)

talk by P. Azzurri at FCC Week 2018, Amsterdam

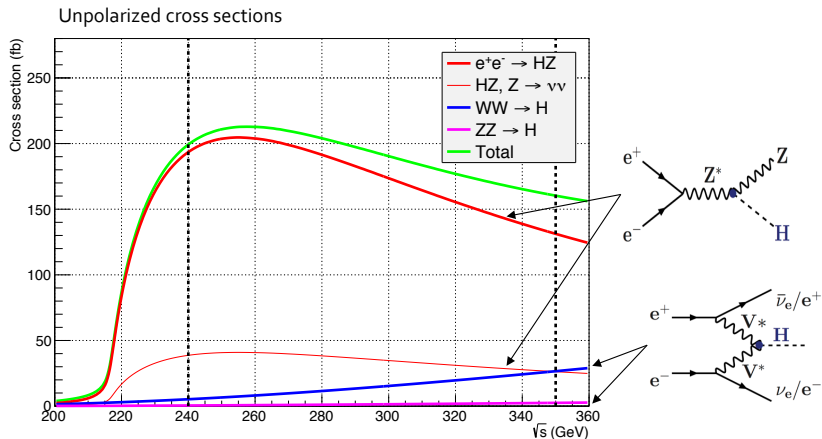
Anomalous triple gauge couplings

\sqrt{s}	Luminosity	# of events
161 GeV	10 ab^{-1}	$\sim 3 \times 10^7$
240 GeV	5 ab^{-1}	$\sim 8 \times 10^7$
350 GeV	0.2 ab^{-1}	$\sim 0.2 \times 10^7$
365 GeV	1.5 ab^{-1}	$\sim 1.5 \times 10^7$



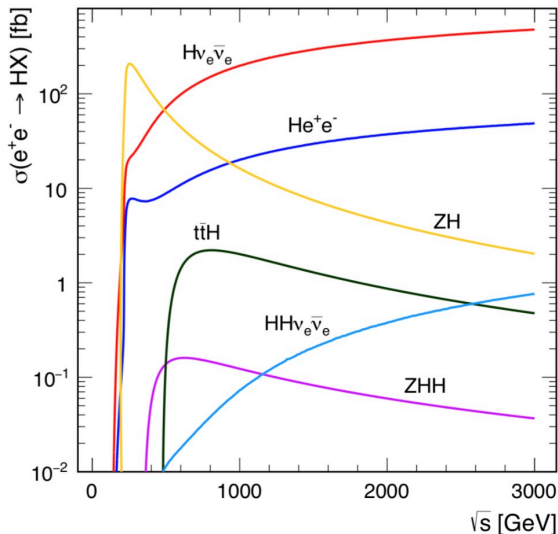
poster by J. Gu at FCC Week 2018, Amsterdam

Higgs production at circular e^+e^- colliders

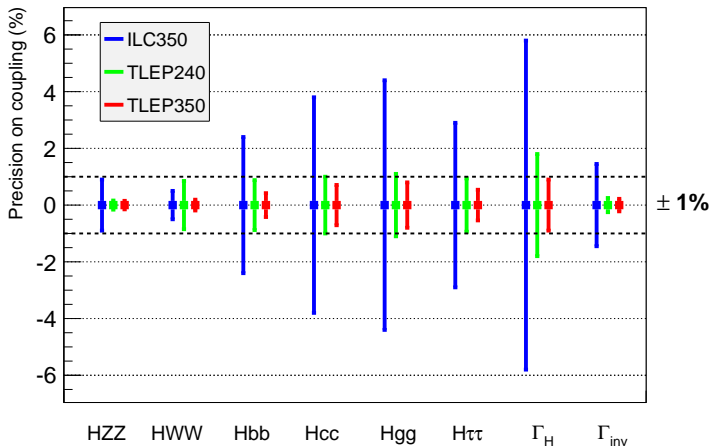


M. Bicer et al., arXiv:1308.6176

Higgs production at e^+e^- linear colliders

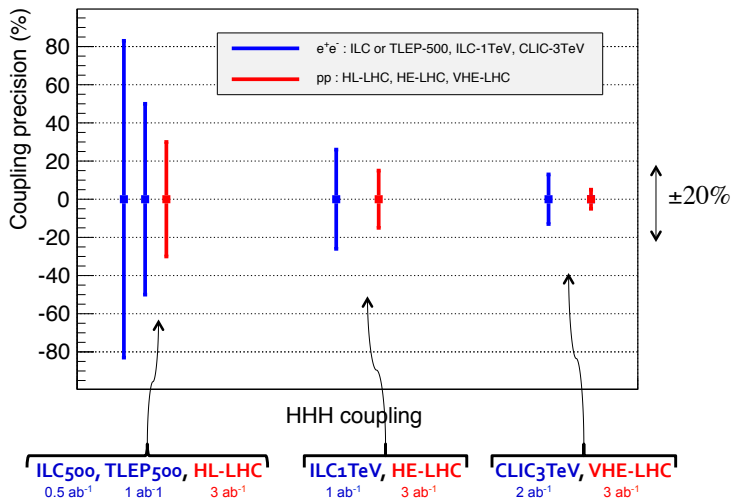


after five years of running FCC-ee



M. Bicer et al., arXiv:1308.6176

Comparison on λ_{HHH} between linear e^+e^- and pp



M. Bicer et al., arXiv:1308.6176

- rich physics potential of e^+e^- colliders
- very high precision measurements at Z peak give the strongest internal consistency checks of the SM
- challenge for precision in theory calculations
- direct measurements of M_W and m_t at the respective thresholds give strong input to the internal consistency checks
- Higgs couplings can be measured at the % level
 - except for Htt , λ_{HHH} where pp colliders display more sensitivity, as well as rare channels like $HZ\gamma$