



Four-loop form factor in $N = 4$

/w Tobias Huber and Gang Yang,
arXiv:1705.03444 & arXiv:1711.08449

Rutger Boels
University of Hamburg



in short

fun result: **a four loop form factor in N=4**

$$\dots + (0.0007 \pm 0.0186)\epsilon^{-3} + (1.60 \pm 0.19)\epsilon^{-2} + (-17.98 \pm 3.47)\epsilon^{-1}$$



in short

fun result: **a four loop form factor in N=4**

$$\dots + (0.0007 \pm 0.0186)\epsilon^{-3} + (1.60 \pm 0.19)\epsilon^{-2} + (-17.98 \pm 3.47)\epsilon^{-1}$$

- at -2,-1: $\neq 0$ \rightarrow speculation in literature
- can be computed **at all** \rightarrow **methods**



in short

fun result: a four loop form factor in $N=4$

our approach is almost certainly wrong...



in short

fun result: a four loop form factor in $N=4$

$$\dots + (0.0007 \pm 0.0186)\epsilon^{-3} + (1.60 \pm 0.19)\epsilon^{-2} + (-17.98 \pm 3.47)\epsilon^{-1}$$

our approach is almost certainly wrong...



two poles of research here at the workshop



two poles of research here at the workshop





two poles of research here at the workshop





perturbative QFT talks cheatsheet

“preprocessing”

generate
integrand

reduce to
simpler integrals

explicit integration

“postprocessing”



perturbative QFT talks cheatsheet

“preprocessing”



- what to compute and why?
- define ‘success’
- how badly do we want it?

generate
integrand

reduce to
simpler integrals

explicit integration

“postprocessing”



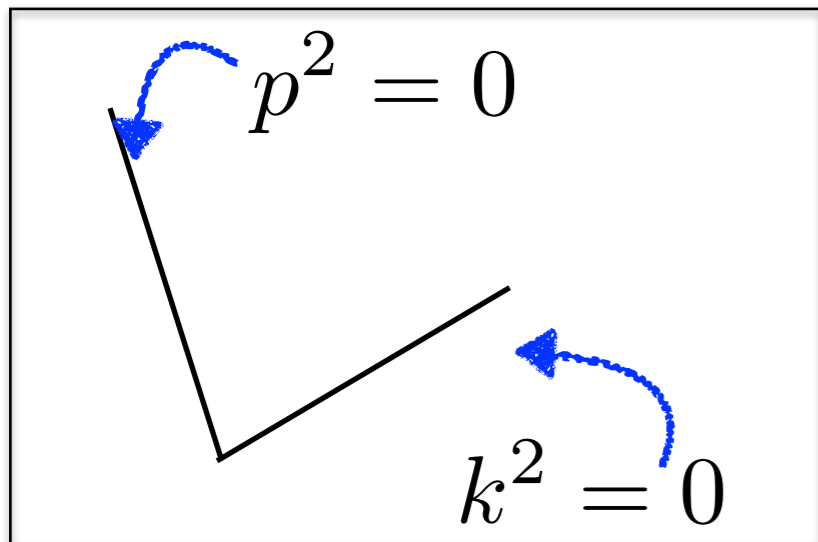
anomalous dimensions

- govern IR/UV divergences [long literature, ~70 - today]



anomalous dimensions

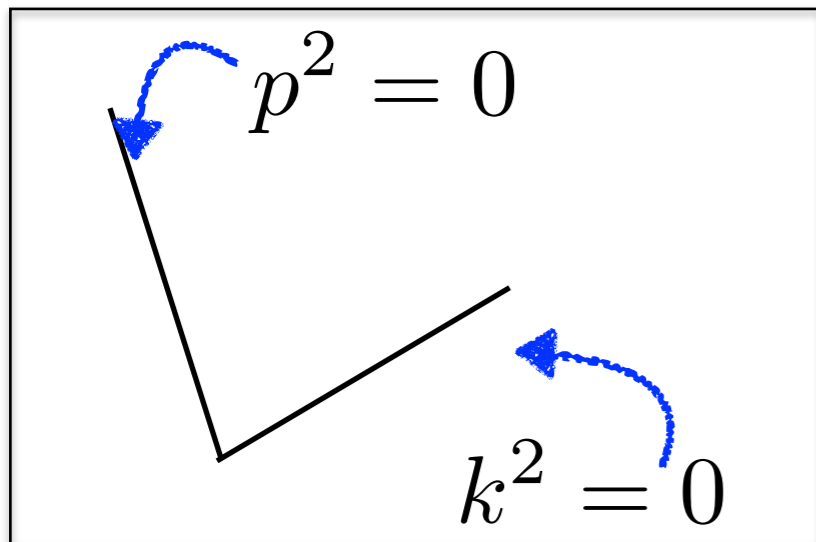
- govern IR/UV divergences [long literature, ~70 - today]
- example: light-like cusp anomalous dimension
 - I anomalous dimension of light like cusped Wilson line
 - II leading infrared divergence of amplitudes
 - III logarithmic growth of high-spin Wilson operators
 - IV related to gluon Regge trajectory
 - V appears in AdS/CFT (N=4)





anomalous dimensions

- govern IR/UV divergences [long literature, ~70 - today]
- example: light-like cusp anomalous dimension
 - I anomalous dimension of light like cusped Wilson line
 - II leading infrared divergence of amplitudes
 - III logarithmic growth of high-spin Wilson operators
 - IV related to gluon Regge trajectory
 - V appears in AdS/CFT (N=4)



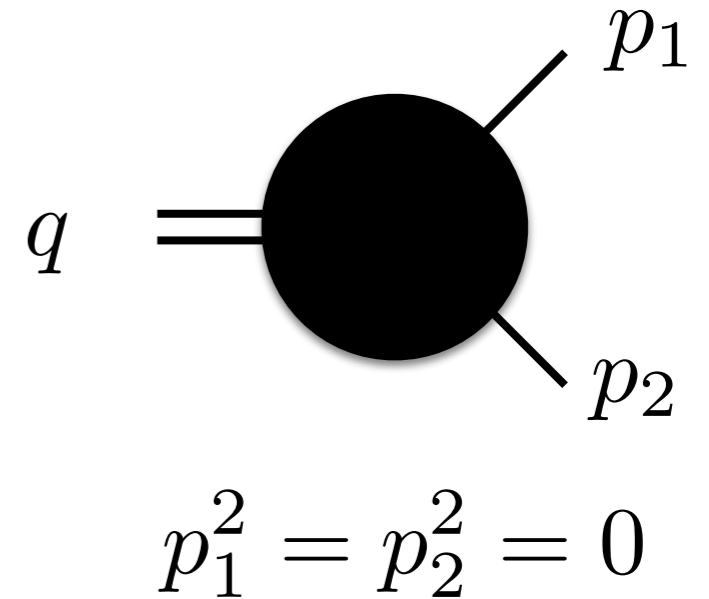
→ ample motivation to compute it!
(many approaches to compute it...)



anomalous dimensions

- here: two gluon + stress-tensor multiplet in N=4

$$\mathcal{F}^{(l)} = \mathcal{F}^{\text{tree}} g^{2l} (-q^2)^{-l\epsilon} F^{(l)}$$

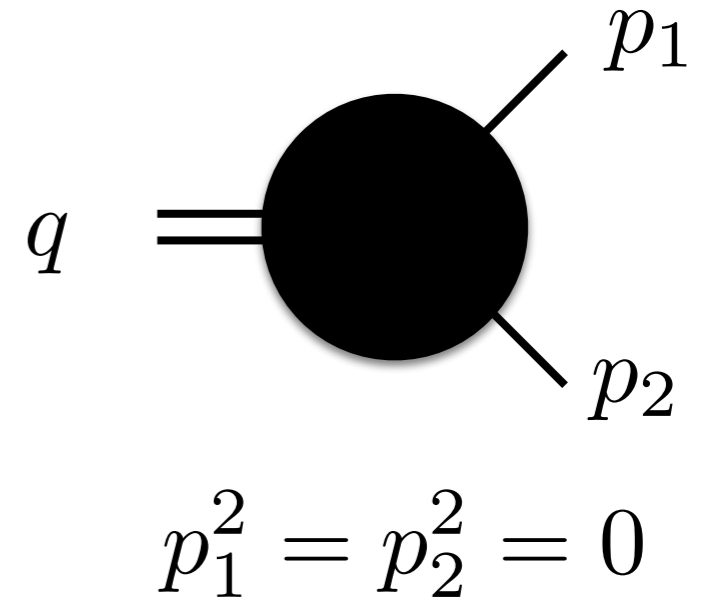




anomalous dimensions

- here: two gluon + stress-tensor multiplet in N=4

$$\mathcal{F}^{(l)} = \mathcal{F}^{\text{tree}} g^{2l} (-q^2)^{-l\epsilon} F^{(l)}$$



- general theory of IR divergences:

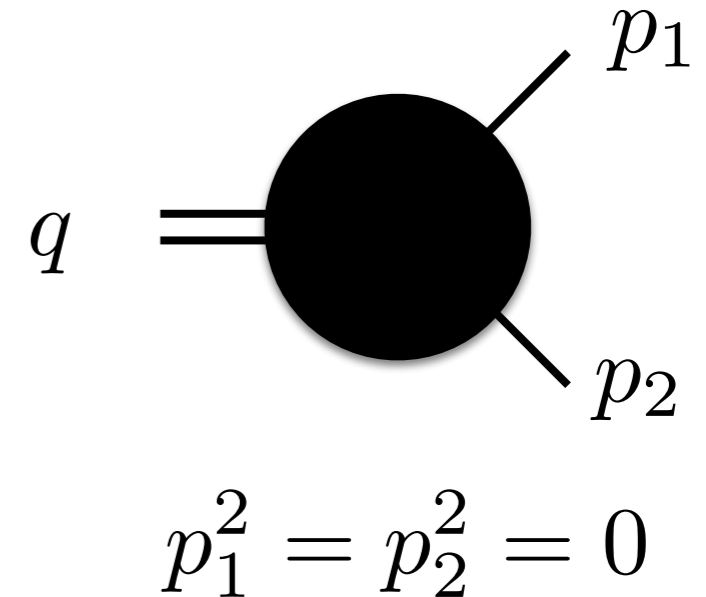
$$(\log F)^{(l)} = - \left[\frac{\gamma_{\text{cusp}}^{(l)}}{(2l\epsilon)^2} + \frac{\mathcal{G}_{\text{coll}}^{(l)}}{2l\epsilon} + \text{Fin}^{(l)} \right] + \mathcal{O}(\epsilon)$$



anomalous dimensions

- here: two gluon + stress-tensor multiplet in N=4

$$\mathcal{F}^{(l)} = \mathcal{F}^{\text{tree}} g^{2l} (-q^2)^{-l\epsilon} F^{(l)}$$



- general theory of IR divergences:

$$(\log F)^{(l)} = - \left[\frac{\gamma_{\text{cusp}}^{(l)}}{(2l\epsilon)^2} + \frac{\mathcal{G}_{\text{coll}}^{(l)}}{2l\epsilon} + \text{Fin}^{(l)} \right] + \mathcal{O}(\epsilon)$$

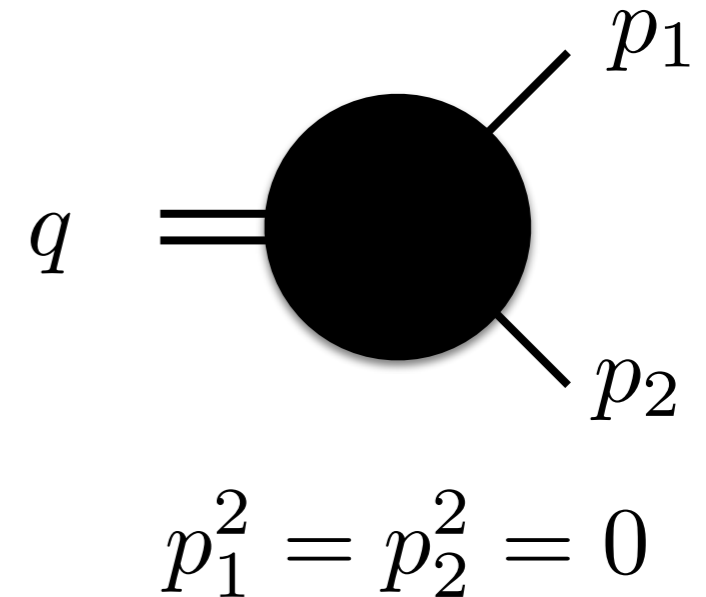
- in N=4 two loop form factor: [\[Van Neerven, 1986\]](#),
three loops [\[Gehrmann-Henn-Huber, 11\]](#)



anomalous dimensions

- here: two gluon + stress-tensor multiplet in N=4

$$\mathcal{F}^{(l)} = \mathcal{F}^{\text{tree}} g^{2l} (-q^2)^{-l\epsilon} F^{(l)}$$



- general theory of IR divergences:

$$(\log F)^{(l)} = - \left[\frac{\gamma_{\text{cusp}}^{(l)}}{(2l\epsilon)^2} + \frac{\mathcal{G}_{\text{coll}}^{(l)}}{2l\epsilon} + \text{Fin}^{(l)} \right] + \mathcal{O}(\epsilon)$$

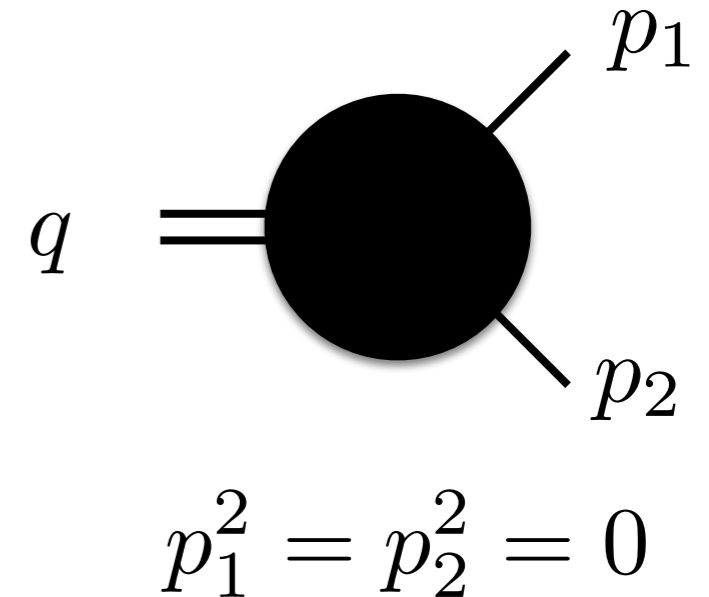
- in N=4 two loop form factor: [\[Van Neerven, 1986\]](#), three loops [\[Gehrmann-Henn-Huber, 11\]](#)
- in QCD, three loops [\[Gehrmann et.al, 06\]](#) - [\[Baikov et.al, 09\]](#)



anomalous dimensions

- here: two gluon + stress-tensor multiplet in N=4

$$\mathcal{F}^{(l)} = \mathcal{F}^{\text{tree}} g^{2l} (-q^2)^{-l\epsilon} F^{(l)}$$



- general theory of IR divergences:

$$(\log F)^{(l)} = - \left[\frac{\gamma_{\text{cusp}}^{(l)}}{(2l\epsilon)^2} + \frac{\mathcal{G}_{\text{coll}}^{(l)}}{2l\epsilon} + \text{Fin}^{(l)} \right] + \mathcal{O}(\epsilon)$$

- in N=4 two loop form factor: [\[Van Neerven, 1986\]](#), three loops [\[Gehrmann-Henn-Huber, 11\]](#)
- in QCD, three loops [\[Gehrmann et.al, 06\]](#) - [\[Baikov et.al, 09\]](#)
- planar limit known **exactly** [\[Beisert-Eden-Staudacher, 06\]](#)



4 loop Sudakov form factor

- function of coupling constant, group theory:

$$\gamma_{\text{cusp}} = \sum_l g^{2l} \gamma_{\text{cusp}}^{(l)} =$$

$$a_1 g^2 C_A + a_2 g^4 C_A^2 + a_3 g^6 C_A^3 + g^8 (a_4^{\text{P}} C_A^4 + a_4^{\text{NP}} d_{44}) + \mathcal{O}(g^9)$$



4 loop Sudakov form factor

- function of coupling constant, group theory:

$$\gamma_{\text{cusp}} = \sum_l g^{2l} \gamma_{\text{cusp}}^{(l)} = \sim N_c^4 \quad \sim \dots N_c^4 + \dots N_c^2$$
$$a_1 g^2 C_A + a_2 g^4 C_A^2 + a_3 g^6 C_A^3 + g^8 (a_4^{\text{P}} C_A^4 + a_4^{\text{NP}} d_{44}) + \mathcal{O}(g^9)$$

- first “non-planar” correction at **four** loops!



4 loop Sudakov form factor

- function of coupling constant, group theory:

$$\gamma_{\text{cusp}} = \sum_l g^{2l} \gamma_{\text{cusp}}^{(l)} = \sim N_c^4 \quad \sim \dots N_c^4 + \dots N_c^2$$
$$a_1 g^2 C_A + a_2 g^4 C_A^2 + a_3 g^6 C_A^3 + g^8 (a_4^{\text{P}} C_A^4 + a_4^{\text{NP}} d_{44}) + \mathcal{O}(g^9)$$

- first “non-planar” correction at **four** loops!
- a_{NP}^4 conjectured to vanish [Becher-Neubert, 09] in any QFT
cf. [(Dixon-)Gardi-Magnea, 09],[Ahrens-Neubert-Vernazza, 09]



4 loop Sudakov form factor

- function of coupling constant, group theory:

$$\gamma_{\text{cusp}} = \sum_l g^{2l} \gamma_{\text{cusp}}^{(l)} = \sim N_c^4 \quad \sim \dots N_c^4 + \dots N_c^2$$
$$a_1 g^2 C_A + a_2 g^4 C_A^2 + a_3 g^6 C_A^3 + g^8 (a_4^{\text{P}} C_A^4 + a_4^{\text{NP}} d_{44}) + \mathcal{O}(g^9)$$

- first “non-planar” correction at **four** loops!
- a_4^{NP} conjectured to vanish [Becher-Neubert, 09] in any QFT
cf. [(Dixon-)Gardi-Magnea, 09],[Ahrens-Neubert-Vernazza, 09]
- today: the first computation of nonplanar cusp in **any** QFT



4 loop Sudakov form factor

- function of coupling constant, group theory:

$$\gamma_{\text{cusp}} = \sum_l g^{2l} \gamma_{\text{cusp}}^{(l)} = \sim N_c^4 \quad \sim \dots N_c^4 + \dots N_c^2$$
$$a_1 g^2 C_A + a_2 g^4 C_A^2 + a_3 g^6 C_A^3 + g^8 (a_4^{\text{P}} C_A^4 + a_4^{\text{NP}} d_{44}) + \mathcal{O}(g^9)$$

- first “non-planar” correction at **four** loops!
- a_{NP}^4 conjectured to vanish [Becher-Neubert, 09] in any QFT
cf. [(Dixon-)Gardi-Magnea, 09],[Ahrens-Neubert-Vernazza, 09]
- today: the first computation of nonplanar cusp in **any** QFT
(see also more recent: [Moch-Ruijl-Ueda-Vermaseren-Vogt, 17],
[Grozin-Henn-Stahlhofen, 17])



perturbative QFT talks cheatsheet

“preprocessing”

- what to compute and why?
- define ‘success’
- how badly do we want it?



perturbative QFT talks cheatsheet

“preprocessing”

- what to compute and why?
- define ‘success’
- how badly do we want it?

- non-planar correction to the ‘Sudakov’ form factor in N=4 at four loops to at least leading divergent term: ϵ^{-2} , ϵ^{-1}
- is it zero? \rightarrow numerics (may) suffice
- quite... \rightarrow long-standing conjecture



perturbative QFT talks cheatsheet

“preprocessing”

generate
integrand

reduce to
simpler integrals

explicit integration

“postprocessing”



perturbative QFT talks cheatsheet

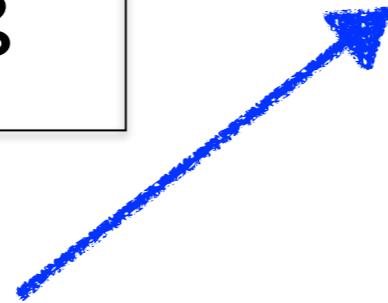
“preprocessing”

generate
integrand

reduce to
simpler integrals

explicit integration

“postprocessing”



- Feynman graphs
- **unitarity based approaches**
- (string theory)



Intermezzo: color kinematic duality

[Bern-Carrasco-Johannson, 08, 10]

- write a gauge theory tree **amplitude** as:

$$A_n = g_{ym}^{n-2} \sum_{\Gamma_i} \frac{n_i c_i}{s_i}$$



Intermezzo: color kinematic duality

[Bern-Carrasco-Johannson, 08, 10]

- write a gauge theory tree **amplitude** as:

$$A_n = g_{ym}^{n-2} \sum_{\Gamma_i} \frac{n_i c_i}{s_i}$$

sum over trivalent, connected tree graphs



Intermezzo: color kinematic duality [Bern-Carrasco-Johansson, 08, 10]

- write a gauge theory tree **amplitude** as:

color structure for each graph

$$A_n = g_{ym}^{n-2} \sum_{\Gamma_i} \frac{n_i c_i}{s_i}$$

sum over trivalent, connected tree graphs



Intermezzo: color kinematic duality

[Bern-Carrasco-Johansson, 08, 10]

- write a gauge theory tree **amplitude** as:

“kinematic numerator”

color structure for each graph

$$A_n = g_{ym}^{n-2} \sum_{\Gamma_i} \frac{n_i c_i}{s_i}$$

sum over trivalent, connected tree graphs



Intermezzo: color kinematic duality

[Bern-Carrasco-Johansson, 08, 10]

- write a gauge theory tree **amplitude** as:

“kinematic numerator”

color structure for each graph

$$A_n = g_{ym}^{n-2} \sum_{\Gamma_i} \frac{n_i c_i}{s_i} \text{ propagators}$$

sum over trivalent, connected tree graphs



Intermezzo: color kinematic duality

[Bern-Carrasco-Johansson, 08, 10]

- write a gauge theory tree **amplitude** as:

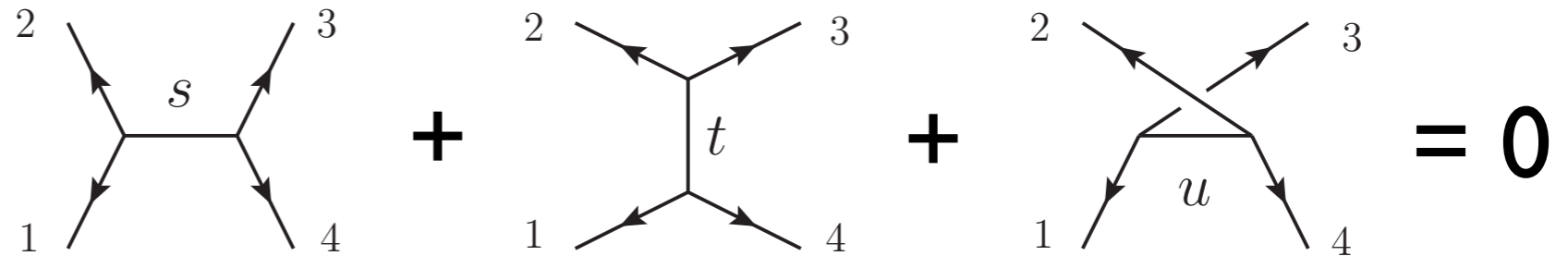
“kinematic numerator”

color structure for each graph

$$A_n = g_{ym}^{n-2} \sum_{\Gamma_i} \frac{n_i c_i}{s_i} \quad \text{propagators}$$

sum over trivalent, connected tree graphs

- \exists color Jacobi identities:





Intermezzo: color kinematic duality

[Bern-Carrasco-Johansson, 08, 10]

- write a gauge theory tree **amplitude** as:

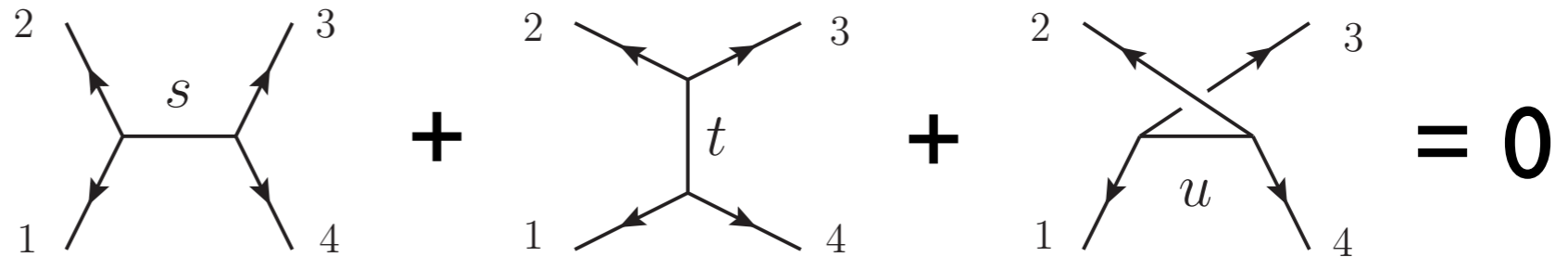
“kinematic numerator”

color structure for each graph

$$A_n = g_{ym}^{n-2} \sum_{\Gamma_i} \frac{n_i c_i}{s_i} \quad \text{propagators}$$

sum over trivalent, connected tree graphs

- \exists color Jacobi identities:



- demand that the kinematic numerators satisfy same Jacobi's:

$$\forall \{c_i = c_k - c_j\} \Rightarrow n_i = n_k - n_j$$



Intermezzo: color kinematic duality

[Bern-Carrasco-Johansson, 08, 10]

- write a gauge theory tree **amplitude** as:

“kinematic numerator”

color structure for each graph

$$A_n = g_{ym}^{n-2} \sum_{\Gamma_i} \frac{n_i c_i}{s_i}$$

propagators

sum over trivalent, connected tree graphs

- \exists color Jacobi identities:

$$\text{Diagram } s + \text{Diagram } t + \text{Diagram } u = 0$$

- demand that the kinematic numerators satisfy same Jacobi's:

$$\forall \{c_i = c_k - c_j\} \Rightarrow n_i = n_k - n_j$$

- always possible at tree level, very similar looking loop level conjecture, see review in [Isermann, 13]



suspicion of duality enough as “Ansatz-generator”

inspired by amplitude computation [[Bern-et.al, 12](#)]:



suspicion of duality enough as “Ansatz-generator”

inspired by amplitude computation [\[Bern-et.al, 12\]](#):

- draw all trivalent graphs, **relate numerators by duality**
- feed in expectations about answer:, e.g. UV divergences
- check Ansatz using cuts

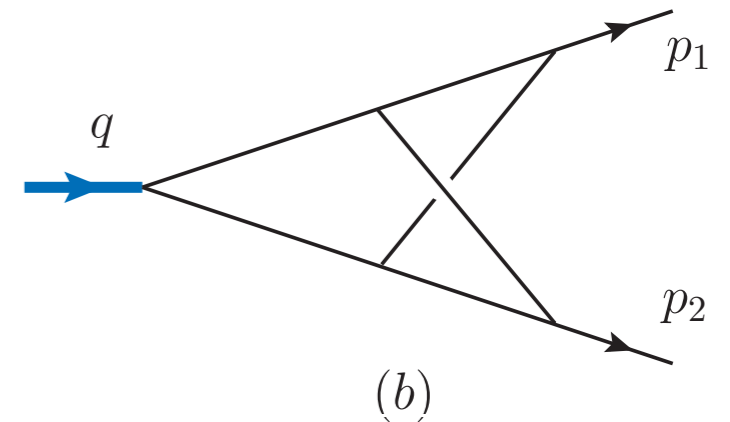
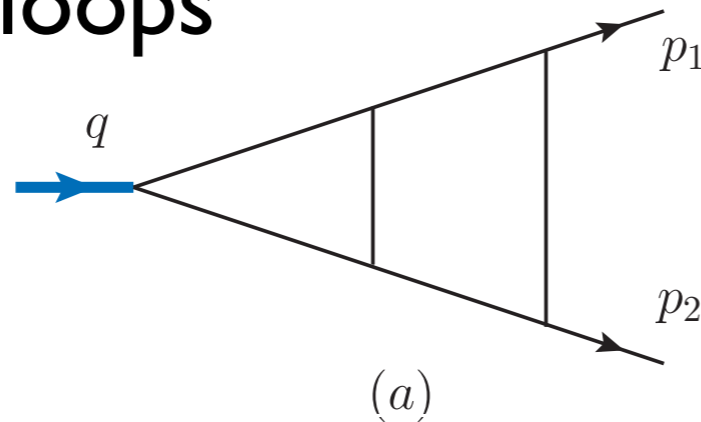


suspicion of duality enough as “Ansatz-generator”

inspired by amplitude computation [\[Bern-et.al, 12\]](#):

- draw all trivalent graphs, **relate numerators by duality**
- feed in expectations about answer:, e.g. UV divergences
- check Ansatz using cuts

e.g. $N=4$, 2pt @ 2 loops



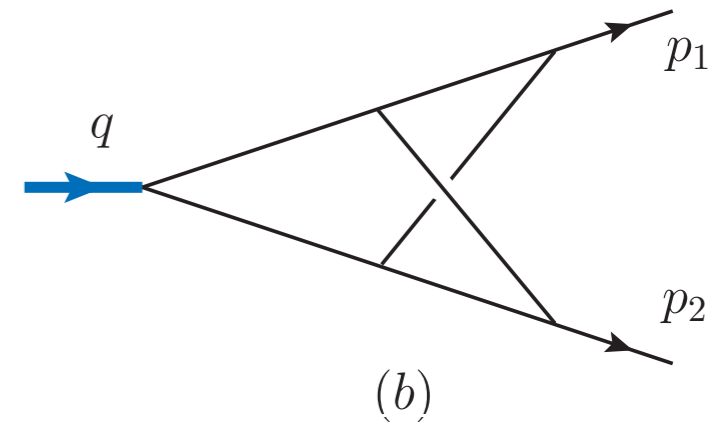
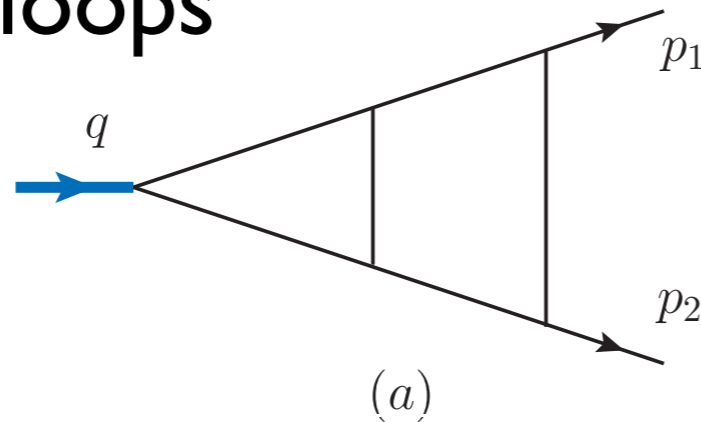


suspicion of duality enough as “Ansatz-generator”

inspired by amplitude computation [Bern-et.al, 12]:

- draw all trivalent graphs, **relate numerators by duality**
- feed in expectations about answer:, e.g. UV divergences
- check Ansatz using cuts

e.g. N=4, 2pt @ 2 loops



$$\begin{aligned}\mathcal{F}_2^{(2)} &= s_{12}^2 \mathcal{F}^{(0)} \sum_{\sigma_2} \sum_{i=a}^b \frac{1}{S_i} C_i I_i, \\ &= N_c^2 \delta^{a_1 a_2} s_{12}^2 (4 I_1 + I_2)\end{aligned}$$

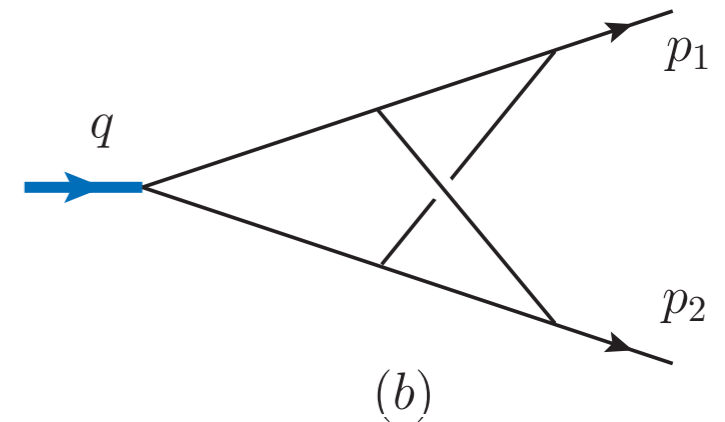
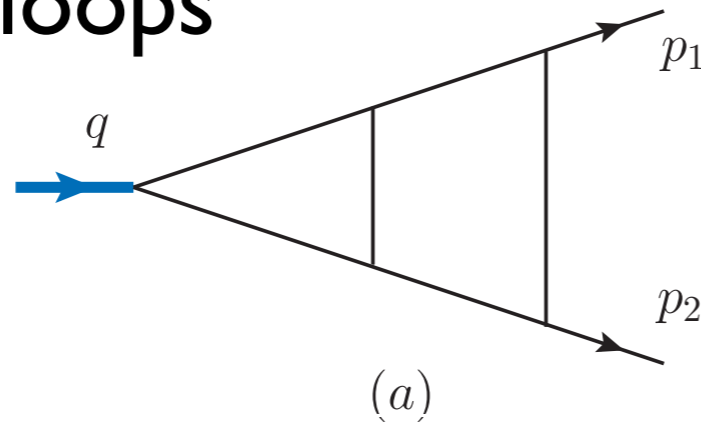


suspicion of duality enough as “Ansatz-generator”

inspired by amplitude computation [Bern-et.al, 12]:

- draw all trivalent graphs, **relate numerators by duality**
- feed in expectations about answer:, e.g. UV divergences
- check Ansatz using cuts

e.g. N=4, 2pt @ 2 loops



$$\begin{aligned}\mathcal{F}_2^{(2)} &= s_{12}^2 \mathcal{F}^{(0)} \sum_{\sigma_2} \sum_{i=a}^b \frac{1}{S_i} C_i I_i, \\ &= N_c^2 \delta^{a_1 a_2} s_{12}^2 (4 I_1 + I_2) \checkmark\end{aligned}$$



Application: form factors

[RB-Kniehl-Tarasov-Yang]

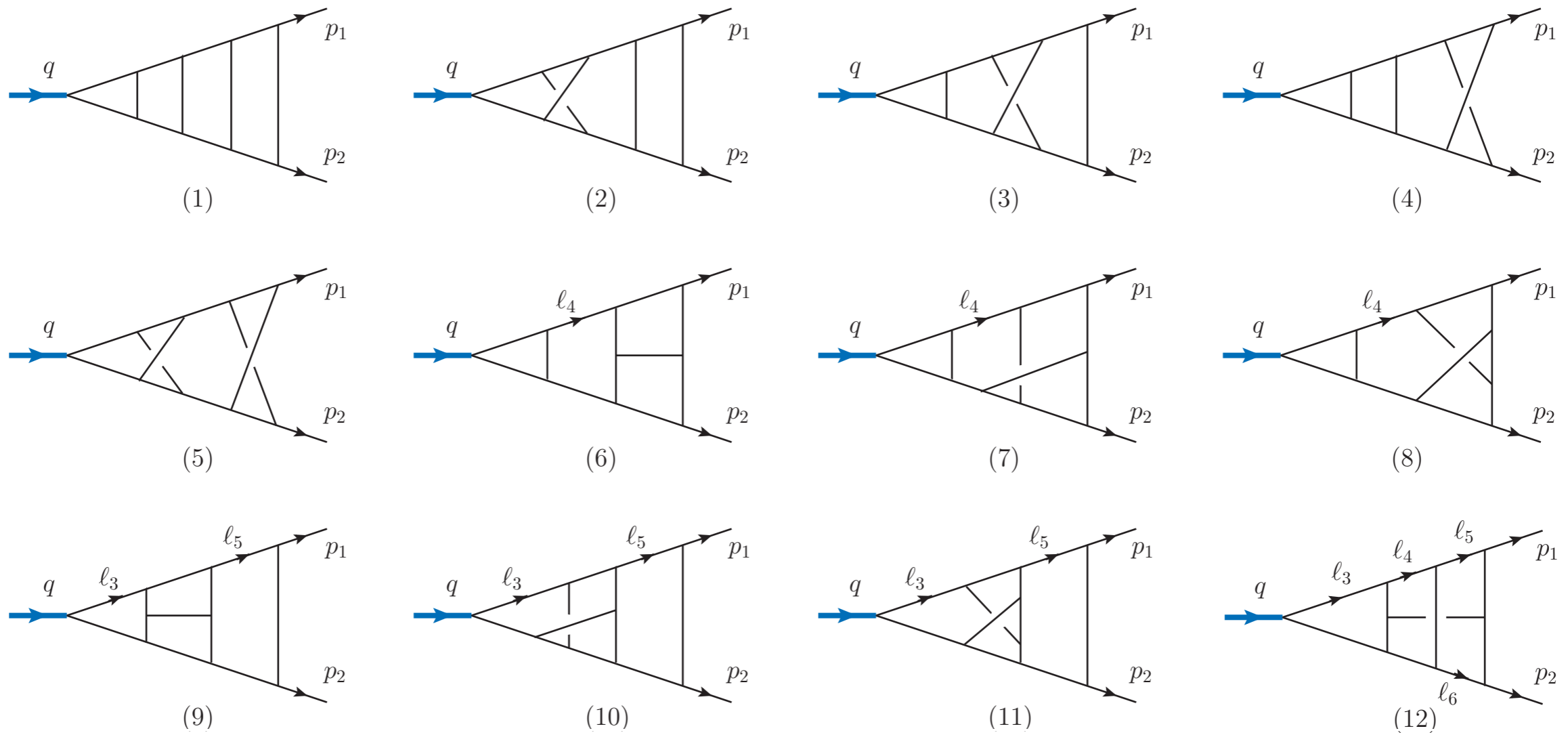
- checked 3 loop-2 point, 2 loop-3 point results (simple!)
- result for 4 loop-2 point:



Application: form factors

[RB-Kniehl-Tarasov-Yang]

- checked 3 loop-2 point, 2 loop-3 point results (simple!)
- result for 4 loop-2 point:

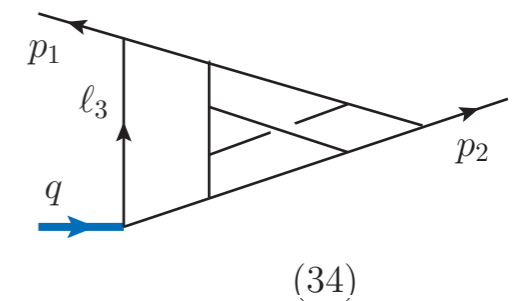
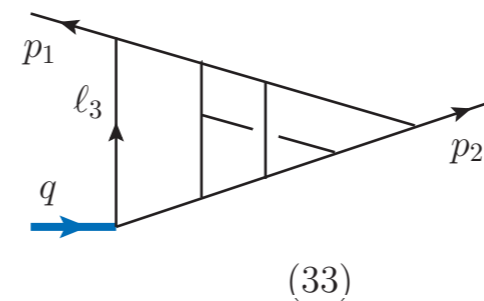
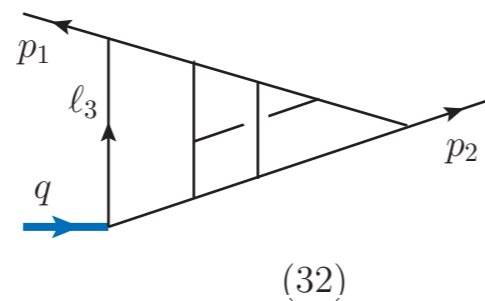
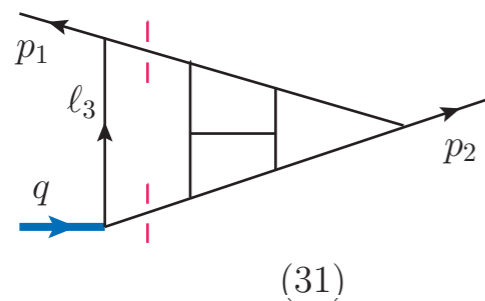
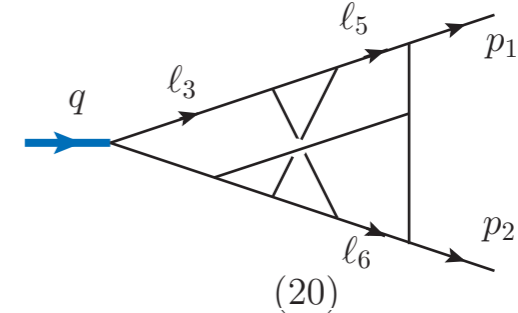
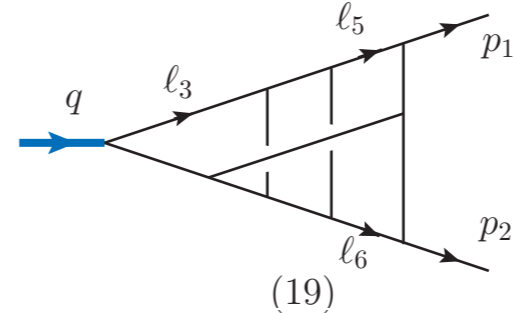
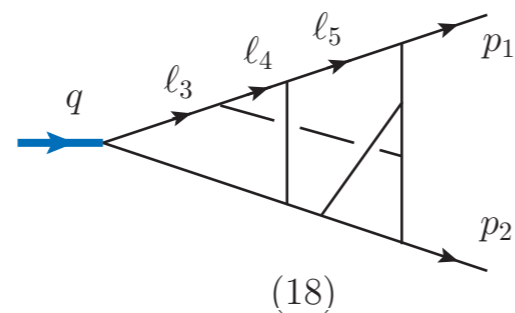
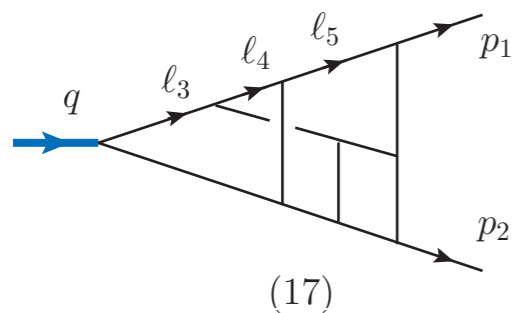
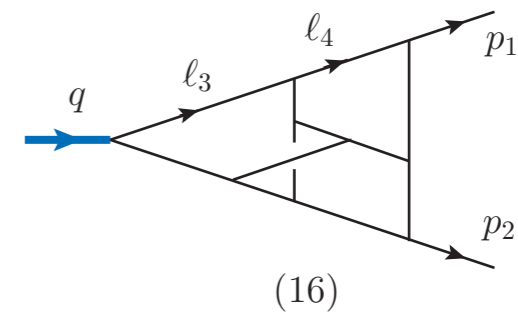
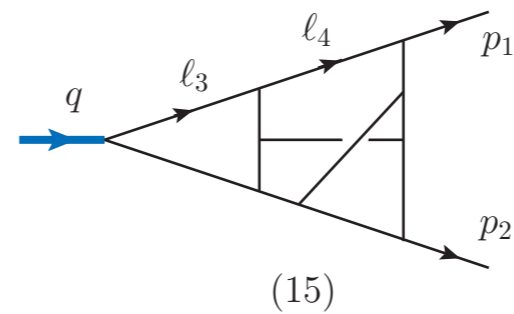
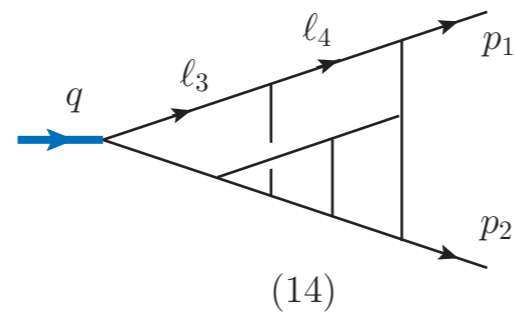
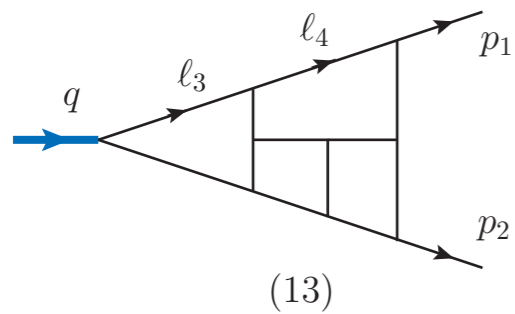




Application: form factors

[RB-Kniehl-Tarasov-Yang]

- checked 3 loop-2 point, 2 loop-3 point results (simple!)
- result for 4 loop-2 point:

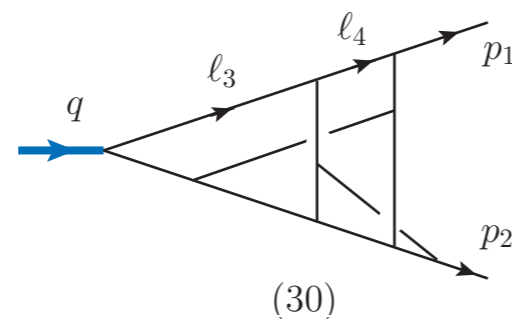
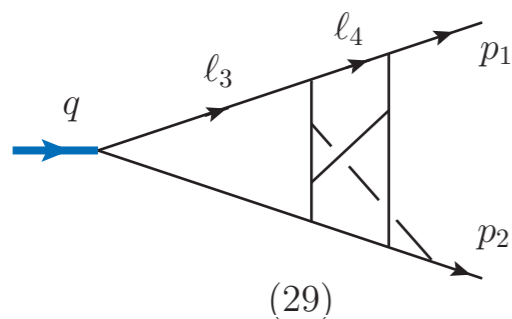
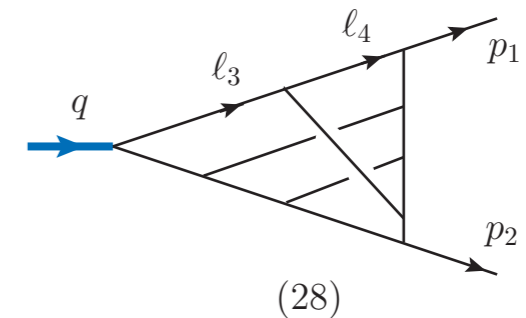
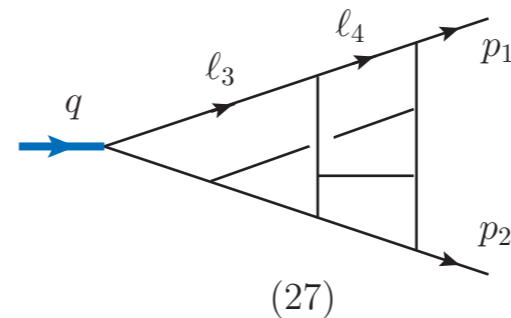
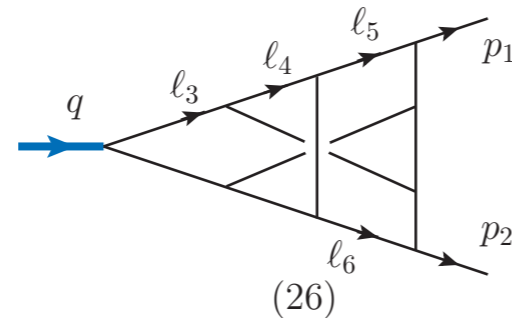
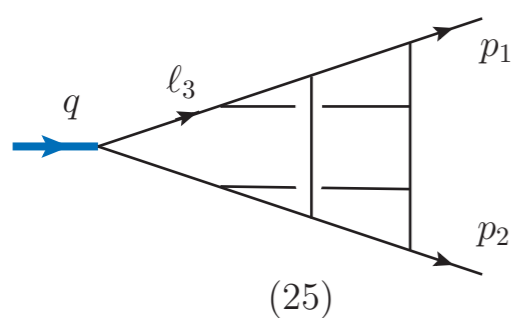
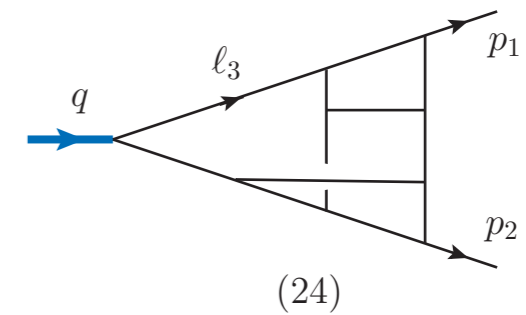
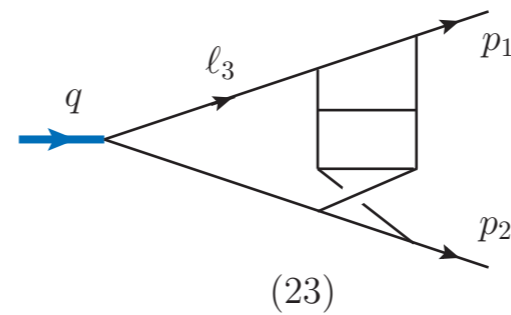
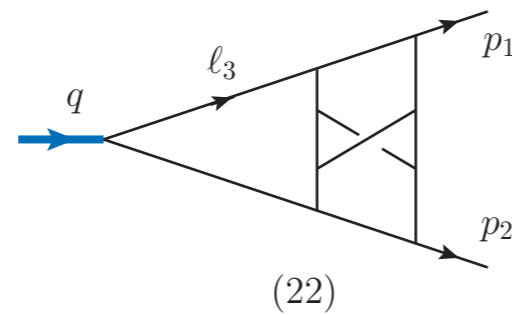
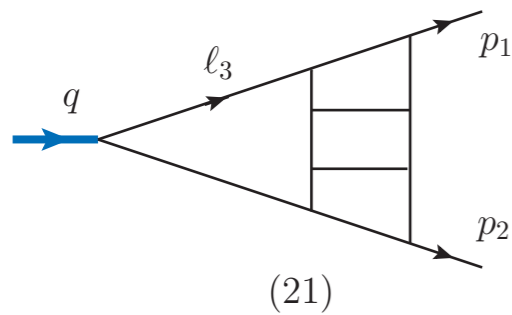




Application: form factors

[RB-Kniehl-Tarasov-Yang]

- checked 3 loop-2 point, 2 loop-3 point results (simple!)
- result for 4 loop-2 point:





Application: form factors

[RB-Kniehl-Tarasov-Yang]

- checked 3 loop-2 point, 2 loop-3 point results (simple!)
- result for 4 loop-2 point:

$$\mathcal{F}_2^{(4)} = s_{12}^2 \mathcal{F}_2^{(0)} \sum_{\sigma_2} \sum_{i=1}^{34} \frac{1}{S_i} C_i I_i .$$



- checked 3 loop-2 point, 2 loop-3 point results (simple!)
- result for 4 loop-2 point:

$$\mathcal{F}_2^{(4)} = s_{12}^2 \mathcal{F}_2^{(0)} \sum_{\sigma_2} \sum_{i=1}^{34} \frac{1}{S_i} C_i I_i .$$

- 34 graphs, 2 “master” graphs.
- first true non-planar corrections
- Ansatz constructed, most unitarity cuts checked
- 1 truly free parameter left



- checked 3 loop-2 point, 2 loop-3 point results (simple!)
- result for 4 loop-2 point:

$$\mathcal{F}_2^{(4)} = s_{12}^2 \mathcal{F}_2^{(0)} \sum_{\sigma_2} \sum_{i=1}^{34} \frac{1}{S_i} C_i I_i .$$

- 34 graphs, 2 “master” graphs.
- first true non-planar corrections
- Ansatz constructed, most unitarity cuts checked
- 1 truly free parameter left

→ color-kinematic duality exists up to four loops for (some) form factors



- checked 3 loop-2 point, 2 loop-3 point results (simple!)
- result for 4 loop-2 point:

$$\mathcal{F}_2^{(4)} = s_{12}^2 \mathcal{F}_2^{(0)} \sum_{\sigma_2} \sum_{i=1}^{34} \frac{1}{S_i} C_i I_i .$$

- 34 graphs, 2 “master” graphs.
- first true non-planar corrections
- Ansatz constructed, most unitarity cuts checked
- 1 truly free parameter left

→ color-kinematic duality exists up to four loops for (some) form factors

[Yang, 16]: five loop case



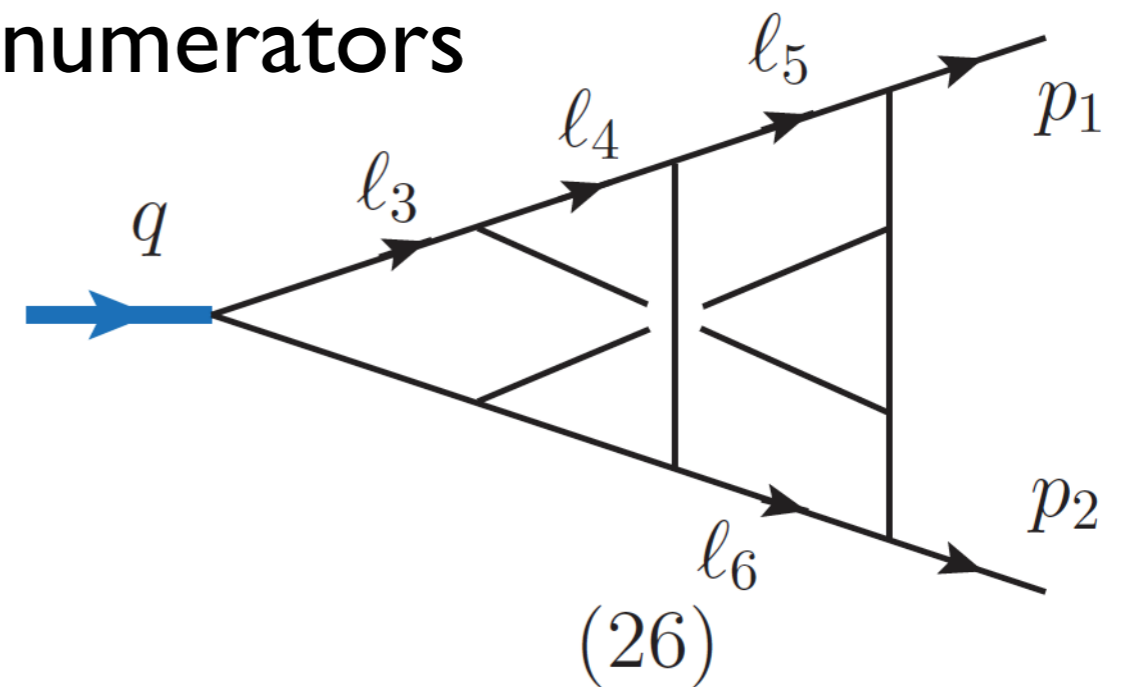
constructing the form factor integrand



constructing the form factor integrand

Integral statistics after generation:

- 34 integrals
- 13 have a non-planar color part
- 10 are purely non-planar color
- mostly quadratic in 6 irreducible numerators
- topology 26: no internal boxes



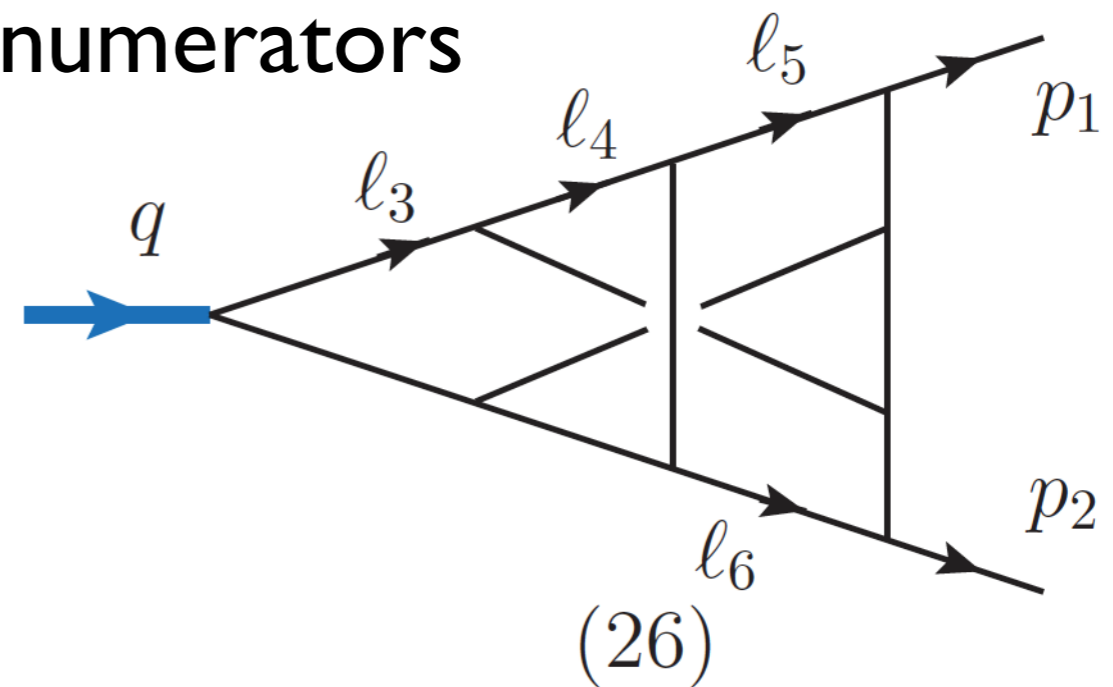


constructing the form factor integrand

Integral statistics after generation:

- 34 integrals
- 13 have a non-planar color part
- 10 are purely non-planar color
- mostly quadratic in 6 irreducible numerators
- topology 26: no internal boxes

→ way to complicated!





constructing the form factor integrand

Inte

recent progress planar, Nf dependent pieces:

- 34
- 13
- 10
- m
- to

[Henn-Smirnov²-Steinhauser, 16]

[Henn-Lee-Smirnov²-Steinhauser, 16]

[Von Manteuffel-Schabinger, 16]

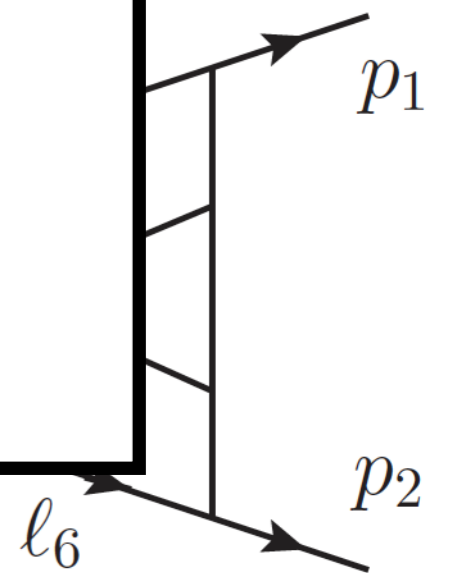
[Lee-Smirnov²-Steinhauser, 17]

[Ahmed-Henn-Steinhauser, 17]

[Grozin-Henn-Stahlhofen, 17]

→ Volodya Smirnov, Peter Marquard's talks

→ way to complicated!



(26)



perturbative QFT talks cheatsheet

“preprocessing”

generate
integrand

reduce to
simpler integrals

explicit integration

“postprocessing”



perturbative QFT talks cheatsheet

“preprocessing”

generate
integrand

reduce to
simpler integrals

explicit integration

“postprocessing”



- integration-by-parts identities
- graph symmetries
- dimensional shifts
- choice of basis



simplification?



simplification?

- integration-by-parts identities

$$\int d^D l_1 \dots d^D l_L \frac{\partial}{\partial l_i^\mu} (\text{integrand}) = 0$$

- LARGE systems of linear equations, solution in terms of **choice** of master integrals
- Laporta algorithm, implemented in e.g. LiteRed, FIRE, Reduze, Kira, Air, private → Volodya's talk



simplification?

- integration-by-parts identities

$$\int d^D l_1 \dots d^D l_L \frac{\partial}{\partial l_i^\mu} (\text{integrand}) = 0$$

- LARGE systems of linear equations, solution in terms of **choice** of master integrals
- Laporta algorithm, implemented in e.g. LiteRed, FIRE, Reduze, Kira, Air, private → Volodya's talk
- here: Reduze ([\[Von Manteuffel-Studerus, 12\]](#)) in [\[Boels-Kniehl-Yang, 15\]](#)

two problems:

- too many hard master integrals
- epsilon dependent coefficients

$$FF = \dots + \left(\frac{\sim 1}{\epsilon^4} + \frac{\sim 10}{\epsilon^3} + \dots \right) I_{\text{master}}$$



good master integrals?

key idea: uniformly transcendental integrals are good



good master integrals?

key idea: uniformly transcendental integrals are good

in expansions of Feynman integrals certain constants
always appear: multiple zeta values

e.g. $\zeta(n)$

not that many constants:

$$\{1\}, \{ \}, \{\pi^2\}, \{\zeta_3\}, \{\pi^4\}, \{\pi^2 \zeta_3, \zeta_5\}, \{\pi^6, \zeta_3^2\}, \dots$$



good master integrals?

key idea: uniformly transcendental integrals are good

in expansions of Feynman integrals certain constants always appear: multiple zeta values

e.g. $\zeta(n)$

not that many constants:

$$\{1\}, \{ \}, \{\pi^2\}, \{\zeta_3\}, \{\pi^4\}, \{\pi^2 \zeta_3, \zeta_5\}, \{\pi^6, \zeta_3^2\}, \dots$$

- every order in epsilon expansion typically has a maximal ‘weight’
 - in N=4, **only maximal terms observed**
 - maximal transcendent part of QCD \leftrightarrow N=4



good master integrals?

key idea: uniformly transcendental integrals are good

in expansions of Feynman integrals certain constants always appear: multiple zeta values

e.g. $\zeta(n)$

not that many constants:

$$\{1\}, \{ \}, \{\pi^2\}, \{\zeta_3\}, \{\pi^4\}, \{\pi^2 \zeta_3, \zeta_5\}, \{\pi^6, \zeta_3^2\}, \dots$$

- every order in epsilon expansion typically has a maximal ‘weight’
 - in N=4, **only maximal terms observed**
 - maximal transcendent part of QCD \leftrightarrow N=4
- idea: find integrals only have maximal terms



good master integrals?

key idea: uniformly transcendental integrals are good

in expansions of Feynman integrals certain constants always appear: multiple zeta values

e.g. $\zeta(n)$

not that many constants:

$$\{1\}, \{ \}, \{\pi^2\}, \{\zeta_3\}, \{\pi^4\}, \{\pi^2 \zeta_3, \zeta_5\}, \{\pi^6, \zeta_3^2\}, \dots$$

- every order in epsilon expansion typically has a maximal ‘weight’
 - in N=4, **only maximal terms observed**
 - maximal transcendent part of QCD \leftrightarrow N=4
- idea: find integrals that are uniformly transcendental



good master integrals?

key idea: uniformly transcendental integrals are good

in expansions of Feynman integrals certain constants always appear: multiple zeta values

e.g. $\zeta(n)$

not that many constants:

$$\{1\}, \{ \}, \{\pi^2\}, \{\zeta_3\}, \{\pi^4\}, \{\pi^2 \zeta_3, \zeta_5\}, \{\pi^6, \zeta_3^2\}, \dots$$

- every order in epsilon expansion typically has a maximal ‘weight’
 - in N=4, **only maximal terms observed**
 - maximal transcendent part of QCD \leftrightarrow N=4
- idea: find integrals that are uniformly transcendental
- cf. [\[Gehrmann-Henn-Huber, 11\]](#) at 3 loops



when is an integral uniformly transcendental?

example from [\[Henn-Smirnov-Smirnov-Steinhauser, 16\]](#):

$$I_{12 \text{ prop}} = \frac{1}{576} + \epsilon^2 \frac{1}{216} \pi^2 + \epsilon^3 \frac{151}{864} \zeta_3 + \epsilon^4 \frac{173}{10368} \pi^4 + \epsilon^5 \left[\frac{505}{1296} \pi^2 \zeta_3 + \frac{5503}{1440} \zeta_5 \right] + \\ + \epsilon^6 \left[\frac{6317}{155520} \pi^6 + \frac{9895}{2592} \zeta_3^2 \right] + \epsilon^7 \left[\frac{89593}{77760} \pi^4 \zeta_3 + \frac{3419}{270} \pi^2 \zeta_5 - \frac{169789}{4032} \zeta_7 \right]$$



when is an integral uniformly transcendental?

example from [\[Henn-Smirnov-Smirnov-Steinhauser, 16\]](#):

$$I_{12 \text{ prop}} = \frac{1}{576} + \epsilon^2 \frac{1}{216} \pi^2 + \epsilon^3 \frac{151}{864} \zeta_3 + \epsilon^4 \frac{173}{10368} \pi^4 + \epsilon^5 \left[\frac{505}{1296} \pi^2 \zeta_3 + \frac{5503}{1440} \zeta_5 \right] + \\ + \epsilon^6 \left[\frac{6317}{155520} \pi^6 + \frac{9895}{2592} \zeta_3^2 \right] + \epsilon^7 \left[\frac{89593}{77760} \pi^4 \zeta_3 + \frac{3419}{270} \pi^2 \zeta_5 - \frac{169789}{4032} \zeta_7 \right]$$

can you tell an integral is UT **without integrating it?**



when is an integral uniformly transcendental?

example from [\[Henn-Smirnov-Smirnov-Steinhauser, 16\]](#):

$$I_{12 \text{ prop}} = \frac{1}{576} + \epsilon^2 \frac{1}{216} \pi^2 + \epsilon^3 \frac{151}{864} \zeta_3 + \epsilon^4 \frac{173}{10368} \pi^4 + \epsilon^5 \left[\frac{505}{1296} \pi^2 \zeta_3 + \frac{5503}{1440} \zeta_5 \right] + \\ + \epsilon^6 \left[\frac{6317}{155520} \pi^6 + \frac{9895}{2592} \zeta_3^2 \right] + \epsilon^7 \left[\frac{89593}{77760} \pi^4 \zeta_3 + \frac{3419}{270} \pi^2 \zeta_5 - \frac{169789}{4032} \zeta_7 \right]$$

can you tell an integral is UT **without integrating it?**

(differential equations)



when is an integral uniformly transcendental?

example from [\[Henn-Smirnov-Smirnov-Steinhauser, 16\]](#):

$$I_{12 \text{ prop}} = \frac{1}{576} + \epsilon^2 \frac{1}{216} \pi^2 + \epsilon^3 \frac{151}{864} \zeta_3 + \epsilon^4 \frac{173}{10368} \pi^4 + \epsilon^5 \left[\frac{505}{1296} \pi^2 \zeta_3 + \frac{5503}{1440} \zeta_5 \right] + \\ + \epsilon^6 \left[\frac{6317}{155520} \pi^6 + \frac{9895}{2592} \zeta_3^2 \right] + \epsilon^7 \left[\frac{89593}{77760} \pi^4 \zeta_3 + \frac{3419}{270} \pi^2 \zeta_5 - \frac{169789}{4032} \zeta_7 \right]$$

can you tell an integral is UT **without integrating it?**



when is an integral uniformly transcendental?

example from [\[Henn-Smirnov-Smirnov-Steinhauser, 16\]](#):

$$I_{12 \text{ prop}} = \frac{1}{576} + \epsilon^2 \frac{1}{216} \pi^2 + \epsilon^3 \frac{151}{864} \zeta_3 + \epsilon^4 \frac{173}{10368} \pi^4 + \epsilon^5 \left[\frac{505}{1296} \pi^2 \zeta_3 + \frac{5503}{1440} \zeta_5 \right] + \\ + \epsilon^6 \left[\frac{6317}{155520} \pi^6 + \frac{9895}{2592} \zeta_3^2 \right] + \epsilon^7 \left[\frac{89593}{77760} \pi^4 \zeta_3 + \frac{3419}{270} \pi^2 \zeta_5 - \frac{169789}{4032} \zeta_7 \right]$$

can you tell an integral is UT **without integrating it?**

- dLog form exist: certainly UT
 - conjecture: constant 'leading singularity' integrals are UT
- [\[Bern-Hermann-Litsey-Stankowicz-Trnka, 14\]](#)
[\[Henn-Smirnov-Smirnov-Steinhauser, 16\]](#)



finding UT integrals

express all loop momenta in a four D basis:

$$l^i = \alpha_1^i p_1 + \alpha_2^i p_2 + \alpha_3^i q_1 + \alpha_4^i q_2$$

consider integrand $I(\vec{\alpha}^i)$ in $D=4$

‘constant leading singularity’ \rightarrow simple poles in all variables



finding UT integrals

express all loop momenta in a four D basis:

$$l^i = \alpha_1^i p_1 + \alpha_2^i p_2 + \alpha_3^i q_1 + \alpha_4^i q_2$$

consider integrand $I(\vec{\alpha}^i)$ in D=4

‘constant leading singularity’ \rightarrow simple poles in all variables

if non-simple pole appears in taking multi-residues:
integral **not** UT

 **integral** property from **integrand**



finding UT integrals

express all loop momenta in a four D basis:

$$l^i = \alpha_1^i p_1 + \alpha_2^i p_2 + \alpha_3^i q_1 + \alpha_4^i q_2$$

consider integrand $I(\vec{\alpha}^i)$ in D=4

‘constant leading singularity’ → simple poles in all variables

if non-simple pole appears in taking multi-residues:
integral **not** UT

 **integral** property from **integrand**

- many multi residues possible (4*4=16 variables)
- pick random sequences: non-UT integrals tend to fail quickly
→ double or higher poles



finding UT integrals: algorithm

if non-simple pole appears in taking multi-residues:
integral certainly not UT

- take a set of integrals
- find any higher-pole-generating sequence of residues
- derive constraint on set of integrals to evade higher order residue → smaller set of integrals
- repeat



finding UT integrals: algorithm

if non-simple pole appears in taking multi-residues:
integral certainly not UT

- take a set of integrals
- find any higher-pole-generating sequence of residues
- derive constraint on set of integrals to evade higher order residue → smaller set of integrals
- repeat

→ output is a set of integrals that pass checks: UT candidates



finding UT integrals: algorithm

if non-simple pole appears in taking multi-residues:
integral certainly not UT

- take a set of integrals
- find any higher-pole-generating sequence of residues
- derive constraint on set of integrals to evade higher order residue → smaller set of integrals
- repeat

→ output is a set of integrals that pass checks: UT candidates

maximal initial set of integrals from dimensional analysis: here quadratic numerator integrals (190)



finding UT integrals: algorithm

if non-simple pole appears in taking multi-residues:
integral certainly not UT

- take a set of integrals
- find any higher-pole-generating sequence of residues
- derive constraint on set of integrals to evade higher order residue → smaller set of integrals
- repeat

→ output is a set of integrals that pass checks: UT candidates

maximal initial set of integrals from dimensional analysis: here quadratic numerator integrals (190)

(some topologies have no candidates!)



express form factor in terms of UT integrals

- now known: form factor & a set of candidate UT integrals & IBP solution → in principle enough information

result: full form factor expressed in UT-candidate integrals



express form factor in terms of UT integrals

- now known: form factor & a set of candidate UT integrals & IBP solution → in principle enough information
- refinement: IBP relations without epsilon dependence
- can be obtained directly, but here from IBP-subreduction, [Boels-Kniehl-Yang, 16]
- output is a **minimal set of rational IBP relations** for given set of integrals
 - advantage: fits easily in laptop memory!
 - disadvantage: less powerful

result: full form factor expressed in UT-candidate integrals



express form factor in terms of UT integrals

- now known: form factor & a set of candidate UT integrals & IBP solution → in principle enough information
- refinement: IBP relations without epsilon dependence
- can be obtained directly, but here from IBP-subreduction, [Boels-Kniehl-Yang, 16]
- output is a **minimal set of rational IBP relations** for given set of integrals
 - advantage: fits easily in laptop memory!
 - disadvantage: less powerful

result: full form factor expressed in UT-candidate integrals

form factor is (likely) maximally transcendental



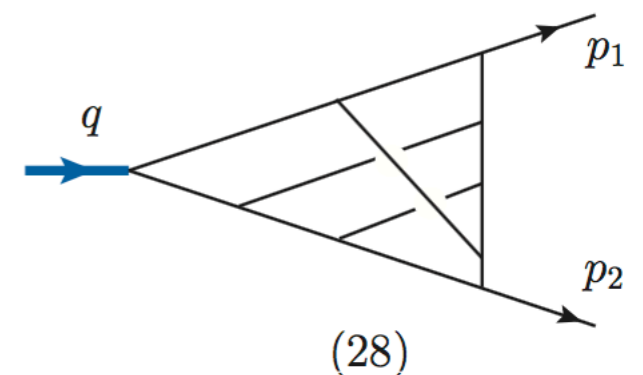
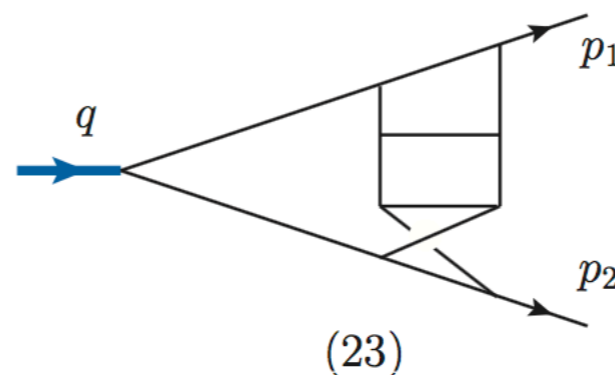
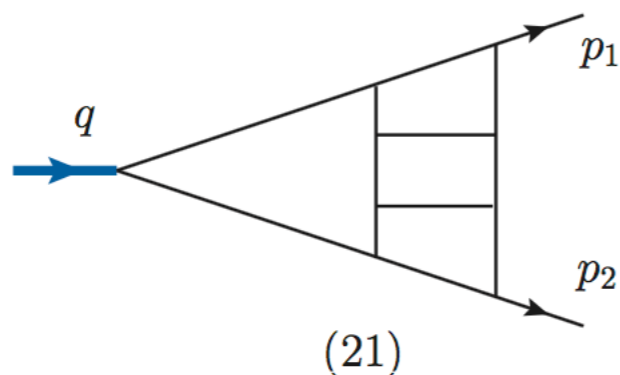
express form factor in terms of UT integrals

- want to put UT integrals into product form for easy input into FIESTA / MB

$$\sum c_i \text{UTC}_i = (\text{quadratic in } l_i) (\text{quadratic in } l_i)$$

- brute force using Mathematica
- aim to minimise number of integrals for form factor
- found choice of 23 / 34 UT integrals non-planar/planar
- all passing > 10.000 random residue checks separately

- dLog:





perturbative QFT talks cheatsheet

“preprocessing”

generate
integrand

reduce to
simpler integrals

explicit integration

“postprocessing”



perturbative QFT talks cheatsheet

“preprocessing”

generate
integrand

reduce to
simpler integrals

explicit integration

“postprocessing”



- Mellin-Barnes representation
- dimensional recurrences
- sector decomposition
- otherwise



numerical integration, **non-planar**

important observation: UT integrals are simpler to integrate numerically than non-UT ones!



numerical integration, **non-planar**

important observation: UT integrals are simpler to integrate numerically than non-UT ones!

- for some integrals derived low-dimensional valid MB representation by hand & inspection → precise results
- automated tools used include: [Czakov, 05], [Smirnov², 07], [Gluza, Kajda, Riemann, 07 / 11], [Blümlein et.al, 14]
- one integral known analytically



numerical integration, **non-planar**

important observation: UT integrals are simpler to integrate numerically than non-UT ones!

- for some integrals derived low-dimensional valid MB representation by hand & inspection → precise results
- automated tools used include: [Czakon, 05], [Smirnov², 07], [Gluza, Kajda, Riemann, 07 / 11], [Blümlein et.al, 14]
- one integral known analytically
- rest: FIESTA + CUBA (mostly vegas) + complete cluster [Smirnov-Tentyukov,08][Smirnov²-Tentyukov,09] [Smirnov, 13, 15] + [Hahn, 04]



numerical integration, **non-planar**

important observation: UT integrals are simpler to integrate numerically than non-UT ones!

- for some integrals derived low-dimensional valid MB representation by hand & inspection → precise results
- automated tools used include: [Czakon, 05], [Smirnov², 07], [Gluza, Kajda, Riemann, 07 / 11], [Blümlein et.al, 14]
- one integral known analytically
- rest: FIESTA + CUBA (mostly vegas) + complete cluster [Smirnov-Tentyukov,08][Smirnov²-Tentyukov,09] [Smirnov, 13, 15] + [Hahn, 04]
- FIESTA uses sector decomposition (cf Sophia Borowska's talk)



numerical integration, **non-planar**

important observation: UT integrals are simpler to integrate numerically than non-UT ones!

- for some integrals derived low-dimensional valid MB representation by hand & inspection \rightarrow precise results
- automated tools used include: [Czakon, 05], [Smirnov², 07], [Gluza, Kajda, Riemann, 07 / 11], [Blümlein et.al, 14]
- one integral known analytically
- rest: FIESTA + CUBA (mostly vegas) + complete cluster [Smirnov-Tentyukov,08][Smirnov²-Tentyukov,09] [Smirnov, 13, 15] + [Hahn, 04]
- FIESTA uses sector decomposition (cf Sophia Borowska's talk)
- (some cross-checks for simple integrals)



numerical integration and results

limiting factor: improvements of accuracy scale as $\sqrt{\text{maxeval}}$
integration time scales as maxeval



numerical integration and results

limiting factor: improvements of accuracy scale as $\sqrt{\text{maxeval}}$
integration time scales as maxeval

other users of the local cluster after some time:





numerical integration and results

- integrals diverge as $\sim \frac{1}{\epsilon^8}$
- non-planar cusp is at $\sim \frac{1}{\epsilon^2}$ (transcendentality 6)

seven orders of expansion, first six should **cancel**



numerical integration and results

- integrals diverge as $\sim \frac{1}{\epsilon^8}$
- non-planar cusp is at $\sim \frac{1}{\epsilon^2}$ (transcendentality 6)

seven orders of expansion, first six should **cancel**

numerics for first five orders good enough to apply “PSLQ”
to convert to “**small** rational * zeta value”

(mathematica: “FindIntegerNullVector”)



numerical integration and results

- integrals diverge as $\sim \frac{1}{\epsilon^8}$
- non-planar cusp is at $\sim \frac{1}{\epsilon^2}$ (transcendentality 6)

seven orders of expansion, first six should **cancel**

numerics for first five orders good enough to apply “PSLQ”
to convert to “**small** rational * zeta value”

(mathematica: “FindIntegerNullVector”)

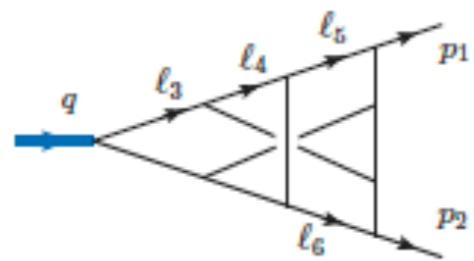
→ non-planar form factor cancels analytically down to $\sim \frac{1}{\epsilon^4}$

$$\dots + (0.0007 \pm 0.0186)\epsilon^{-3} + (1.60 \pm 0.19)\epsilon^{-2} + (-17.98 \pm 3.47)\epsilon^{-1}$$



some sample result...

$$I_7^{(26)} = \text{Diagram} \times \left\{ 4 [(\ell_4 - \ell_5)(\ell_3 - \ell_4 + \ell_5 - p_1)] [(\ell_4 - \ell_6)(\ell_3 - \ell_4 + \ell_6 - p_2)] \right.$$



$$\begin{aligned} & - \ell_5^2 (\ell_6 - p_2)^2 - 4 (\ell_4 - \ell_5)^2 (\ell_3 - \ell_4 + \ell_6 - p_2)^2 - \ell_6^2 (\ell_5 - p_1)^2 \\ & - (\ell_3 - \ell_4)^2 (\ell_5 + \ell_6 - \ell_4)^2 - \ell_4^2 (\ell_3 - \ell_4 + \ell_5 + \ell_6 - p_1 - p_2)^2 \end{aligned} \Big\}$$

$$= \frac{0.00347222}{\epsilon^8} - \frac{0.00000000013}{\epsilon^7} + \frac{0.0114231(17)}{\epsilon^6} + \frac{1.1631(3)}{\epsilon^5} + \frac{2.90880(35)}{\epsilon^4} - \frac{12.2720(43)}{\epsilon^3} + \frac{29.708(57)}{\epsilon^2} + \frac{3185.60 \pm 2.63}{\epsilon},$$

$$I_{7,\text{PSLQ}}^{(26)} = \frac{1}{288\epsilon^8} + \frac{\zeta_2}{144\epsilon^6} + \frac{209\zeta_3}{216\epsilon^5} + \frac{43\zeta_4}{16\epsilon^4} + \mathcal{O}(\epsilon^{-3}).$$





$$\dots + (0.0007 \pm 0.0186)\epsilon^{-3} + (1.60 \pm 0.19)\epsilon^{-2} + (-17.98 \pm 3.47)\epsilon^{-1}$$

integration error is somewhat naive

(cf. [\[Marquard-Smirnov²-Steinhauser-Wellmann, 16\]](#))

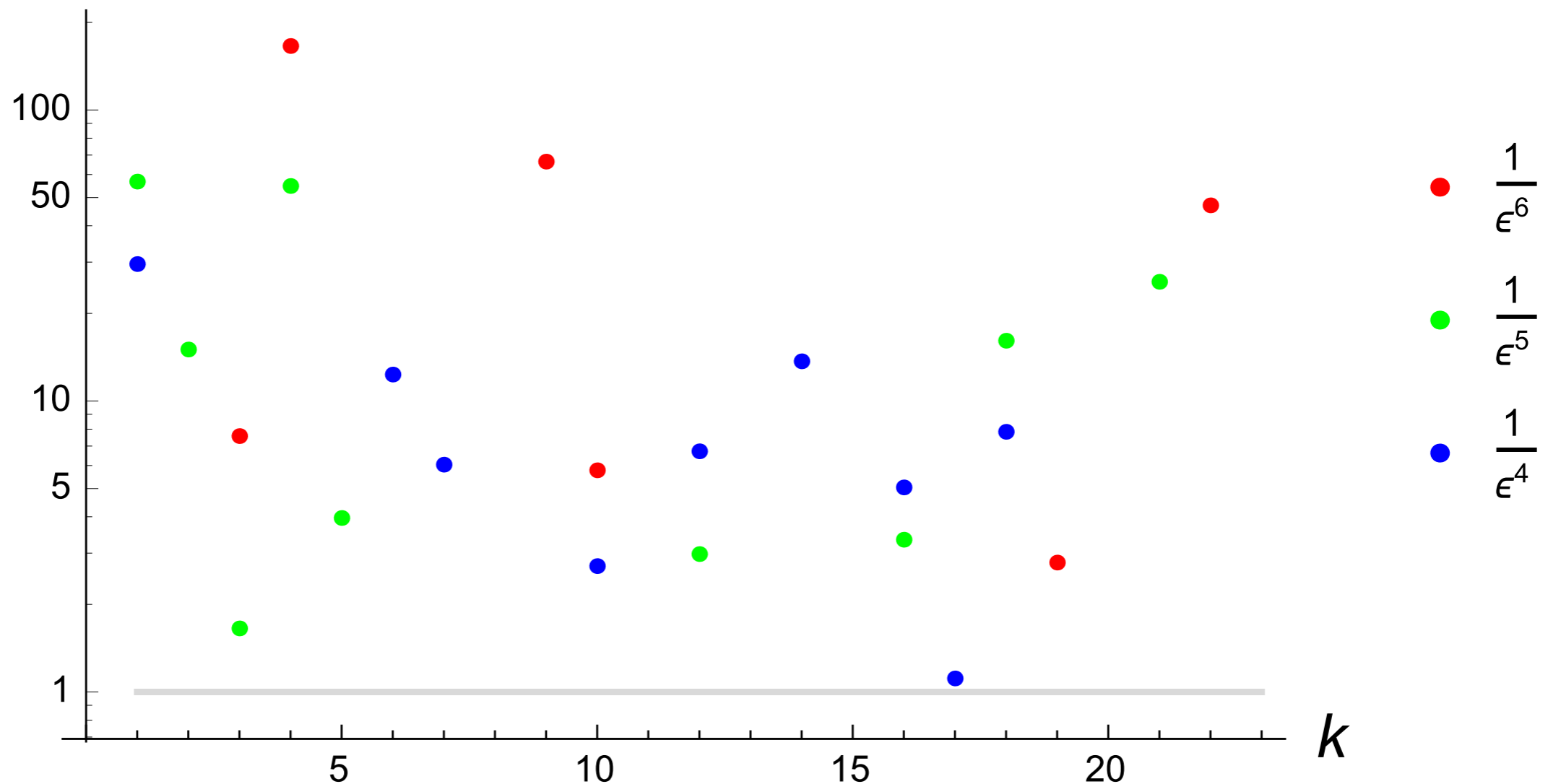
- MB, exact cross-checks
- PSLQ possible
- ϵ^{-3} coefficient central value
- checked stability of central value with increasing points
- error dominated by very few integrals



the rug: example check

- take PSLQ result as exact result, and study numerical deviation

$$\frac{|FIESTAError_k|}{|I_{k,FIESTA} - I_{k,PSLQ}|} > 0$$



- uses number theory to check numerical computations!



in short

fun result: a four loop form factor in $N=4$

- at $-2, -1$: $\neq 0$ → speculation in literature
- can be computed **at all** → **methods**



in short

fun result: **a four loop form factor in N=4**

$$\dots + (0.0007 \pm 0.0186)\epsilon^{-3} + (1.60 \pm 0.19)\epsilon^{-2} + (-17.98 \pm 3.47)\epsilon^{-1}$$

- at -2,-1: $\neq 0$ → speculation in literature
- can be computed **at all** → **methods**



in short

fun result: **a four loop form factor in N=4**

$$\dots + (0.0007 \pm 0.0186)\epsilon^{-3} + (1.60 \pm 0.19)\epsilon^{-2} + (-17.98 \pm 3.47)\epsilon^{-1}$$

- at -2,-1: $\neq 0$ → speculation in literature
- can be computed **at all** → **methods**

-
- extend UT finding to other integrals (five loops!)
 - analytical results for integrals needed...
 - QCD applications: nice choice of basis
 - input for non-planar Beisert-Eden-Staudacher



THANKS FOR A NICE
WORKSHOP!



Your Question
Here?