



a common language

- o top quark probes SM in wildly different settings
 - o systematically limited high precision differential x-sec meas.
 - o rare-event type searches (4t, ...)
 - o everything in between $(T/TT+Z/W/H/\gamma, FCNC, ...)$
 - → uncertainties evolve differently with time
- o different energies 7/8/13 TeV, growing PU, evolving detector
 - Want to confront all that with plethora (and evolving) BSM predictions
 AND include Tevatron/LEP/EWPT results
 AND enjoy the results for some time demand longevity.
- o First approach: anomalous coupling to Lagrangian
 - o defined in the broken phase, simple, facilitates physics intuition
 - o can break symmetries, difficult beyond LO, no global hierarchy of effects
 - o this could all be solved, but relevance of results will decline with time



effective field theory

o generic extension of the Standard Model

Grzadkowski et.al. JHEP 1010 (2010) 085 https://arxiv.org/pdf/1008.4884.pdf

$$\mathcal{L}_{eff} = \mathcal{L}_{SM}^{(4)} + \sum \frac{C_x}{\Lambda^2} O_{6,x} + h.c.$$
 $\overset{\text{O}_{6,x}}{\subset_{\times}}$ 59 dim-6 gauge-invariant ops. $\overset{\text{O}_{6,x}}{\subset_{\times}}$ Wilson coefficients (complex)

scale of dim-6 interactions

- o defined in unbroken phase of SM → complex pattern after EWSB
- based on limited & well defined approximations
 - o global way to look for NP in SM measurements
 - o parameterizes deviations from higher-order SM predictions due to interference of NP with SM; even if new particles are too massive for LHC energies
- o EFT provides guidance to exp. searches

$$\sigma = \sigma_{\rm SM} + \sum_{i} \frac{1 \text{TeV}^2}{\Lambda^2} C_i \sigma_i + \sum_{i \le j} \frac{1 \text{TeV}^4}{\Lambda^4} C_i C_j \sigma_{ij}.$$

- o e.g. on combination strategy in TT+X
- e.g. to include 4-f ops in 2I FCNC states
- o etcetc.

new challenges for analysis design:

- understand operator hierarchy, understand EFT validity
- o longevity/unfolding → sensibly balance exp. uncertainties and longevity. MVA?



FCNC

- o FCNC suppressed to $10^{-12} 10^{-15}$ in SM by GIM mechanism
- o sensitive probe BSM models: 2HDM, SUSY, etc. arXiv:1311.2028
- o anomalous coupling Lagrangian:

$$\mathcal{L}_{FCNC} = \sum_{q=u,c} \left[\frac{\sqrt{2}}{2} g_s \, \frac{\kappa_{\mathrm{tgq}}}{\Lambda} \, \left(\bar{q} \sigma^{\mu\nu} T^a (f_{gq}^L P_L + f_{gq}^R P_R) t \right) G_{\mu\nu}^a \right. \quad 7 \, \& \, 8 \, \mathrm{TeV}^*$$

$$+ \frac{eQ_t}{\sqrt{2}} \, \frac{\kappa_{\mathrm{tq}\gamma}}{\Lambda} \, \left(\bar{q} \sigma^{\mu\nu} (f_{\gamma q}^L P_L + f_{\gamma q}^R P_R) t \right) F_{\mu\nu} \quad 8 \, \mathrm{TeV}^*$$

$$+ \frac{g}{4c_W} \, \frac{\kappa_{\mathrm{tq}Z}}{\Lambda} \, \left(\bar{q} \sigma^{\mu\nu} (\hat{f}_{Zq}^L P_L + f_{Hq}^R P_R) t \right) H \quad 8 \, \& \, 13 \, \mathrm{TeV}$$

$$+ \frac{g}{4c_W} \, \frac{\kappa_{\mathrm{tq}Z}}{\Lambda} \, \left(\bar{q} \sigma^{\mu\nu} (\hat{f}_{Zq}^L P_L + \hat{f}_{Zq}^R P_R) t \right) Z_{\mu\nu} \quad 8 \, \& \, 13 \, \mathrm{TeV}$$

$$+ \frac{g}{4c_W} \, \frac{\kappa_{\mathrm{tq}Z}}{\Lambda} \, \left(\bar{q} \sigma^{\mu\nu} (\hat{f}_{Zq}^L P_L + \hat{f}_{Zq}^R P_R) t \right) Z_{\mu} \right] + \mathrm{h.c.}$$

* see backup

- o often simplify chiral structure, e.g. $f^R = 1$.
- o q = u,c with more sensitivity to u (higher x-sec)



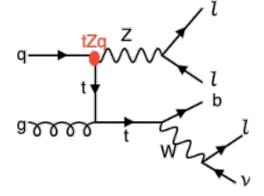
FCNC t/tt qZ

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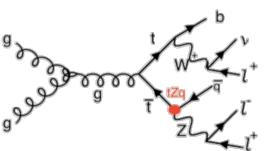
o combine t & tt FCNC channels

8 TeV TOP-12-039 JHEP 07 (2017) 003 https://arxiv.org/abs/1702.01404

FCNC signal	FCNC signal
single-top-quark	tī
1 jet, $ \eta < 2.4$	\geqslant 2 jets, $ \eta $ < 2.4
1 b tag	≥ 1 b tag
$m_{\mathrm{T}}^{\mathrm{W}} > 10 \mathrm{GeV}$	$m_{\mathrm{T}}^{\mathrm{W}} > 10 \mathrm{GeV}$
$p_{\mathrm{T}}^{\mathrm{miss}} > 40 \mathrm{GeV}$	$p_{\mathrm{T}}^{\mathrm{miss}} > 40 \mathrm{GeV}$



13 TeV TOP-17-017 http://cds.cern.ch/record/2292045

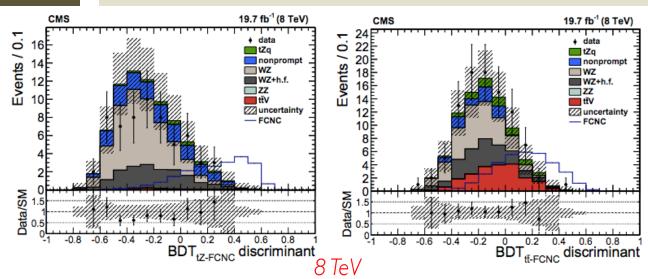


- o consider all flavor combinations eee/ μ ee/ $\mu\mu$ e/ $\mu\mu$ μ require same-flavor opposite-sign Z candidate
- o consider only tensor coupling \mathbf{K}_{tqZ}
- o three low njet/nbjet SB for (I) non-prompt leptons and W+Jets (separated by $m_T(W)$ and per flavor) and for NPL + (2) t and (3) tt
- o train BDTs to separate t and tt-FCNC signal, fit output discriminator in CR and SR simultaneous in t/tt



FCNC t/tt qZ





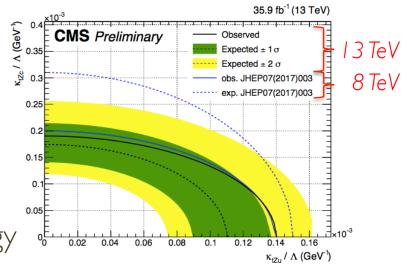
8 TeV TOP-12-039 JHEP 07 (2017) 003 https://arxiv.org/abs/1702.01404

13 TeV TOP-17-017 http://cds.cern.ch/record/2292045

o BDT output for t- and tt FCNC

Branching fraction	Expected	68% CL range	Observed		
$\mathcal{B}(t o Zu)$ (%)	0.027	0.018 - 0.042	0.022		
$\mathcal{B}(t o Zc)$ (%)	0.118	0.071 - 0.222	0.049		
Branching fraction	Expected	Observed			
$\mathcal{B}(t o Zu)$ (%)	0.015	0.024			
$\mathcal{B}(t o Zc)$ (%)	0.037	0.045			
13 ToV					

o statistics dominated; profit from energy and lumi; excluded BR $\sim O(10^{-4})$



Top physics at the precision frontier, FNAL, 18-01-18



FCNC tqH

Data

tī+cc

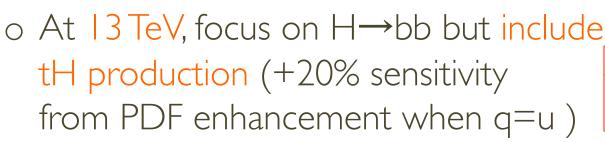
tt+bb

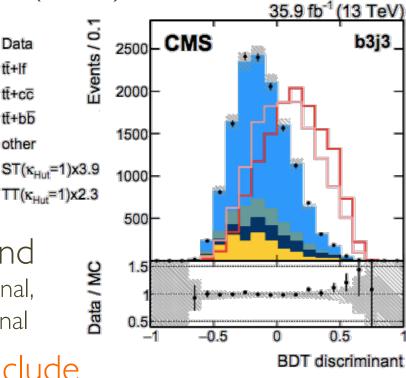
other

o Combine 8 TeV results from top quark pair production in $H \rightarrow bb/\gamma\gamma/WW+\tau\tau(+ZZ)$

13 TeV TOP-17-003 https://arxiv.org/abs/1712.02399 8 TeV TOP-13-017 JHEP 02 (2017) 079 https://arxiv.org/abs/1610.04857

- o H→yy most sensitive bkg estimation from m_{vv} SB fit
- o For $H \rightarrow WW + TT(+ZZ)$ combine SS and multi-lepton channels
- $TT(\kappa_{Hut}=1)x2.3$ o H→bb has largest branching but large combinatorial background BDT to select correct assignment in FCNC signal, ANN (8 TeV) or BDT (13 TeV) to selecting signal





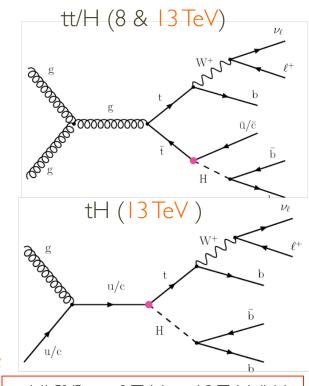
UL[%] 13 TeV (bb) 8 TeV BR(t→Hu) 0.55(0.40) 0.47(0.34) BR(t→Hc) 0.40(0.43) 0.47(0.44)



FCNC tqH

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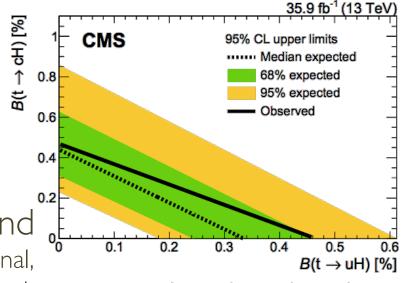
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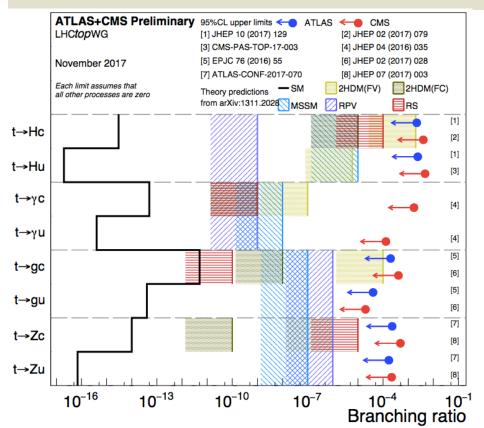


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FCNC summary – ATLAS & CMS

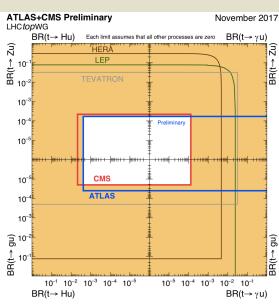
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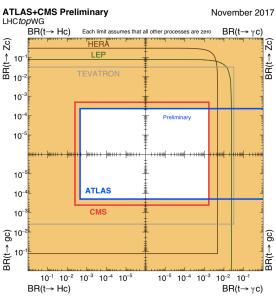


- common language pragmatic choice: BR

EFT approach in: G. Durieux, F. Maltoni, C. Zhang Phys. Rev. D 91, 074017 (2015) https://arxiv.org/pdf/1412.7166.pdf

o points out few missed contributions







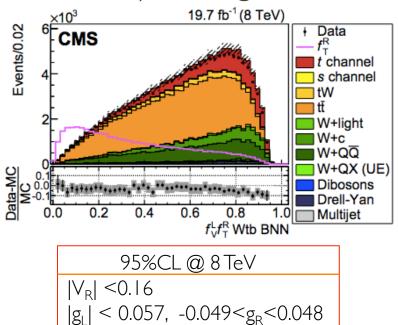
anomalous charged current

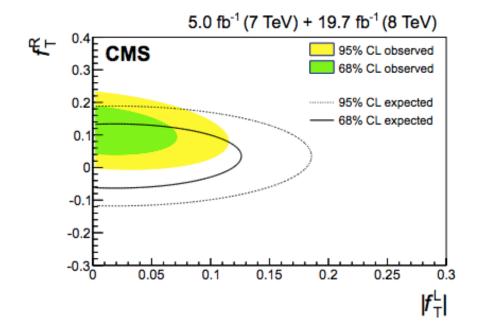
o charged current anomalous interactions

8 TeV TOP-14-007 JHEP 02 (2017) 028 https://arxiv.org/pdf/1610.03545.pdf

$$\mathfrak{L} = \frac{g}{\sqrt{2}}\bar{b}\gamma^{\mu}\left(f_{\mathrm{V}}^{\mathrm{L}}P_{\mathrm{L}} + f_{\mathrm{V}}^{\mathrm{R}}P_{\mathrm{R}}\right)\mathsf{t}W_{\mu}^{-} - \frac{g}{\sqrt{2}}\bar{b}\frac{\sigma^{\mu\nu}\partial_{\nu}W_{\mu}^{-}}{M_{\mathrm{W}}}\left(f_{\mathrm{T}}^{\mathrm{L}}P_{\mathrm{L}} + f_{\mathrm{T}}^{\mathrm{R}}P_{\mathrm{R}}\right)\mathsf{t} + \mathrm{h.c.}$$

- o constrained in single-top t-channel (I μ) at 7 & 8 TeV
 - o 3 BNNs to separate V_L contributions from V_R , g_L , g_R
 - o 3D analysis in right handed vector- and the L&R tensor couplings







helicity fractions

o charged current anomalous interactions

$$\mathfrak{L} = rac{g}{\sqrt{2}} ar{b} \gamma^{\mu} \left(f_{\mathrm{V}}^{\mathrm{L}} P_{\mathrm{L}} + f_{\mathrm{V}}^{\mathrm{R}} P_{\mathrm{R}} \right) \mathrm{t} W_{\mu}^{-} - rac{g}{\sqrt{2}} ar{b} rac{\sigma^{\mu \nu} \partial_{\nu} W_{\mu}^{-}}{M_{\mathrm{W}}} \left(f_{\mathrm{T}}^{\mathrm{L}} P_{\mathrm{L}} + f_{\mathrm{T}}^{\mathrm{R}} P_{\mathrm{R}} \right) \mathrm{t} + \mathrm{h.c.}$$

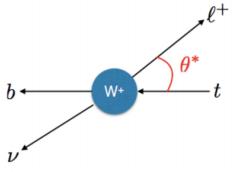
o also modifies helicity fractions and diff. decay w.

$$\frac{1}{\Gamma} \frac{d\Gamma}{d\cos\theta_{\ell}^{*}} = \frac{3}{8} (1 + \cos\theta_{\ell}^{*})^{2} F_{R} + \frac{3}{8} (1 - \cos\theta_{\ell}^{*})^{2} F_{L} + \frac{3}{4} \sin^{2}\theta_{\ell}^{*} F_{0}$$

o semi-lep tt+jets with full kinematic reco

angular I/j

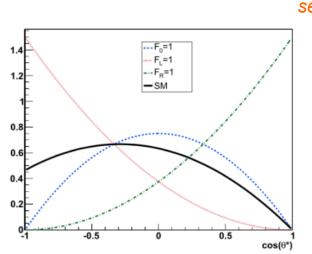
8 TeV TOP-13-008 Phys.Lett. B762 (2016) 512-534 https://arxiv.org/abs/1605.09047

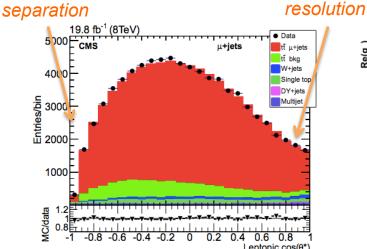


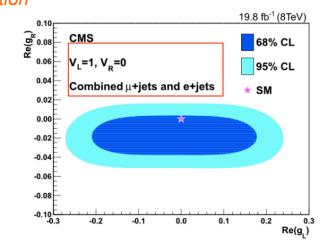
 $F_0 = 0.681 \pm 0.012 \text{ (stat.)} \pm 0.023 \text{ (syst.)}$

 $F_L = 0.323 \pm 0.008 \text{ (stat.)} \pm 0.014 \text{ (syst.)}$

 $F_R = 0.004 \pm 0.005 \text{ (stat.)} \pm 0.014 \text{ (syst.)}$









spin correlation

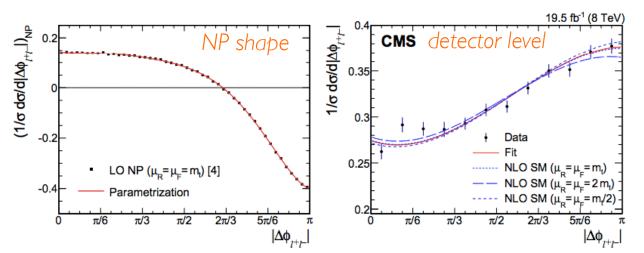
top decay products are a probe of the ttbar spin correlation

$$\frac{1}{m_{
m t}} < \frac{1}{\Gamma_{
m t}} < \frac{1}{\Lambda_{
m QCD}} < \frac{m_{
m t}}{\Lambda^2}$$

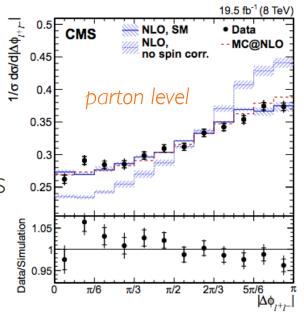
production lifetime hadronization spin-flip 10^{-27} s 10^{-21} s

- o $|\Delta \Phi(II)|$ in 2I-ttbar unfolded to parton-level
- o sensitive to chromo magn. & el. dipole moments

$$\mathcal{L}_{ ext{eff}} = -rac{\widetilde{\mu}_{ ext{t}}}{2}ar{ ext{t}}\sigma^{\mu
u}T^{a} ext{t}G^{a}_{\mu
u} - rac{\widetilde{d}_{ ext{t}}}{2}ar{ ext{t}}i\sigma^{\mu
u}\gamma_{5}T^{a} ext{t}G^{a}_{\mu
u}$$



8 TeV TOP-14-023 PRD 93 (2016) 052007 https://arxiv.org/pdf/1601.01107.pdf



 $|\Delta \Phi(II)|$ & other asymmetries:

$$-0.053 < \text{Re}(\hat{\mu}_t) < 0.026$$

 $-0.068 < \text{Im}(\hat{d}_t) < 0.067$

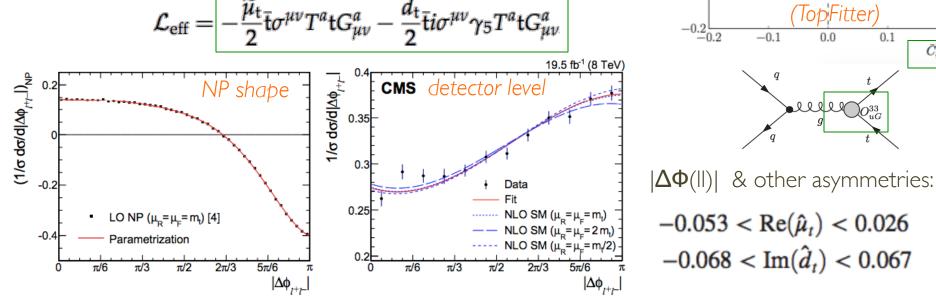


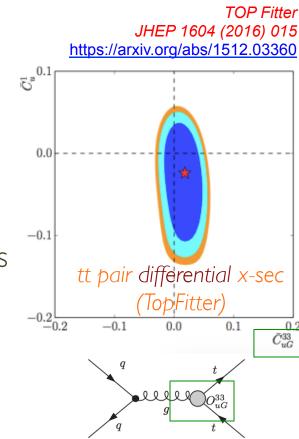
spin correlation

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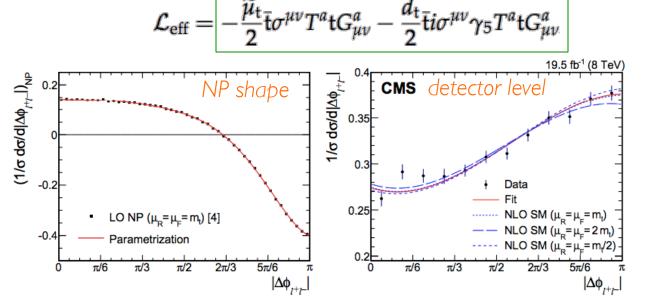


spin correlation

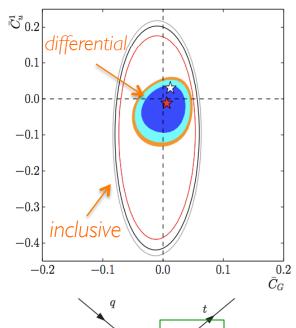
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- o sensitive to chromo magn. & el. dipole moments



TOP Fitter JHEP 1604 (2016) 015 https://arxiv.org/abs/1512.03360





$$-0.053 < \text{Re}(\hat{\mu}_t) < 0.026$$

 $-0.068 < \text{Im}(\hat{d}_t) < 0.067$

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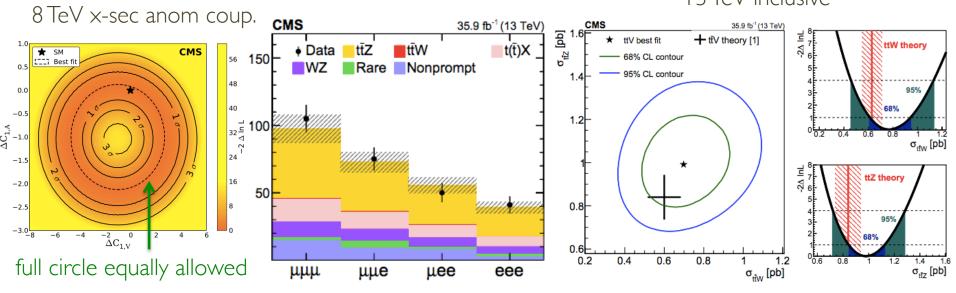
o measured inclusive tt+ $Z/W/\gamma$ x-sec at 8 & 13 TeV

$$\mathcal{L} = e^{\bar{u}(p_t)} \left(\gamma^{\mu} \left(\underline{C_{1,V} + \gamma^5 C_{1,A}} \right) + \frac{i \sigma^{\mu\nu} q_{\nu}}{m_Z} \left(\underline{C_{2,V}^Z + i \gamma^5 C_{2,A}^Z} \right) \right) v(p_{\bar{t}}) Z_{\mu}$$
 vector tensor \rightarrow EWK dipole moments

13 TeV TOP-17-005 submitted to JHEP https://arxiv.org/abs/1711.02547 8 TeV TOP-14-021 JHEP 01 (2016) 096 https://arxiv.org/abs/1510.01131 A. Alloul, B. Fuks, V. Sanz JHEP 1404 (2014) 110 https://arxiv.org/abs/1310.5150

- o 8 TeV: limit on current couplings C_{IV,A}
 - o also used tt+Z inclusive x-sec to constrain "HEL" & anomalous couplings
- o 13 TeV: Combine tt+W/Z inclusive x-sec to constrain 7 W.C.

13 TeV inclusive







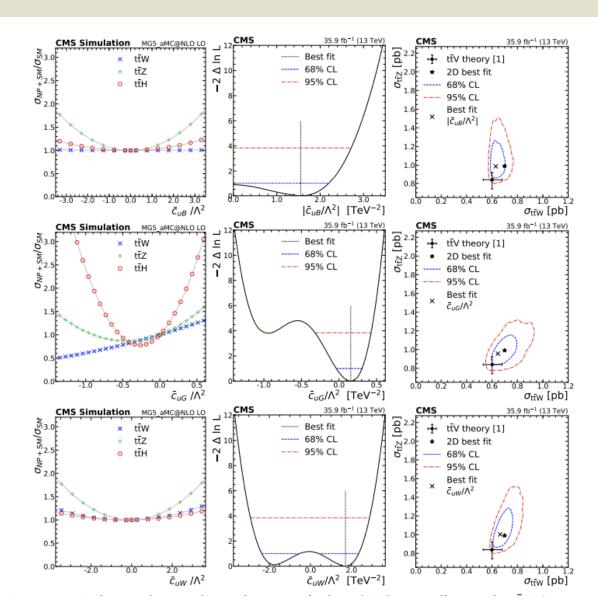
constraints from tt+Z/W

95% CL [TeV⁻²]

	Wilson coefficient	$68\% \text{ CL } [\text{TeV}^{-2}]$	95% CL [TeV
$(ar{q}_p\sigma^{\mu u}u_r) au^I\widetilde{arphi}W^I_{\mu u}$	$\bar{c}_{\mathrm{uW}}/\Lambda^2$	[-1.6, 1.5]	[-2.2, 2.2]
$\partial^{\mu}igl[\Phi^{\dagger}\Phiigr]\partial_{\mu}igl[\Phi^{\dagger}\Phiigr]$	$ \bar{c}_{\rm H}/\Lambda^2 - 16.8{\rm TeV}^{-2} $	[3.7, 23.4]	[0,28.7]
$f^{abc}G_{\mu}^{a u}G_{ u}^{b ho} ilde{G}_{ ho}^{c\mu}$	$\widetilde{c}_{3\mathrm{G}}/\Lambda^2$	[-0.5, 0.5]	[-0.7, 0.7]
$f^{abc}G^{a u}_{\mu}G^{b ho}_{ u}G^{c\mu}_{ ho}$	$ar{c}_{ m 3G}/\Lambda^2$	[-0.3, 0.7]	[-0.5, 0.9]
$(ar q_p \sigma^{\mu u} T^A u_r) \widetilde arphi G^A_{\mu u}$	$\bar{c}_{\mathrm{uG}}/\Lambda^2$	[-0.9, -0.8] and $[-0.3, 0.2]$	[-1.1, 0.3]
$(ar q_p \sigma^{\mu u} u_r) \widetilde arphi B_{\mu u}$	$ \bar{c}_{\mathrm{uB}}/\Lambda^2 $	[0, 1.5]	[0, 2.1]
$(ar{u}_R \gamma^\mu u_R) \left(H^\dagger \overleftrightarrow{D}_\mu H ight)$	$\bar{c}_{\mathrm{Hu}}/\Lambda^2$	[-9.2, -6.5] and $[-1.6, 1.1]$	[-10.1, 2.0]
$\left(D^{\mu}G_{\mu u} ight)^{a}\left(D_{ ho}G^{ ho u} ight)^{a}$	$\bar{c}_{\mathrm{2G}}/\Lambda^{2}$	[-0.7, 0.4]	[-0.9, 0.6]



constraints from tt+Z/W





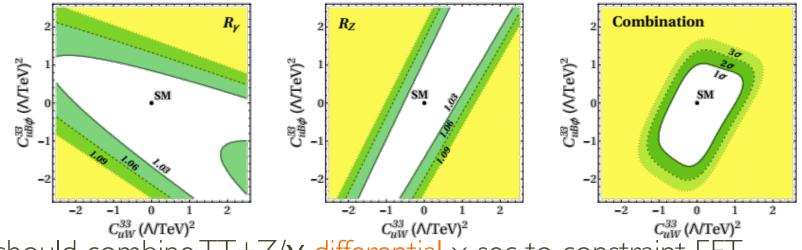
lessons on TT+V from EFT

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$$egin{aligned} \mathcal{L}_{ extstyle ex$$

M. Schulze, Y. Soreg Eur.Phys.J. C76 (2016) no.8, 466 https://arxiv.org/abs/1603.08911

- o gauge symmetry relate EWK dipole moments
- o manifest in EFT \rightarrow Same W.C. enter $C_2^{\mathbb{Z}/Y}$, but in different lin. comb.



- o should combine TT+Z/Y differential x-sec to constraint EFT
- o gauge symmetries relate $C_2^{Z/Y}$ also with the Wtb vertex
 - o new "dim6_top" EFT from TopLHCWG disentangles tt+V from Wtb



new model from TopLHCWG

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- o global strategy provides conservative, well-defined & enduring results
- o "dim6 top EFT" from TopLHCWG http://desy.de/~durieux/topbasis/eft_note-submitted-to-readers.pdf

6 A simple analysis strategy

The simplest strategy to constrain the standard-model effective field theory with measurements in the top sector could rely on fiducial observables defined at, and unfolded to, the particle level. It is meant to be simultaneously practical and useful on a long-term basis, being for instance able to

unfolding	advantages	disadvantages
None (detector level)	high precision, MVAs feasible	hard to integrate in global effort difficult to profit from e.g. theory & MC developments
particle level unfolding (fiducial)	scalable, can use cutting edge theory, active role for theorists	less precise data, limits on e.g. MVA
parton level unfolding (full phase space)	scalable, can use cutting edge theory	least precise, frozen unc. e.g. from acceptance, correlation pattern, strong limits on MVA usage



EFT is not a simple BSM model

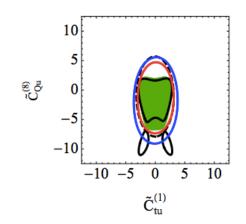
o recent tttt x-sec constraints C_{tu} $\mathcal{O}_{tu}^{(8)} = (\bar{t}_R \gamma_\mu T^a t_R) (\bar{u}_R \gamma^\mu T^a u_R)$

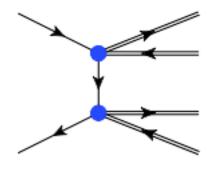
$$\sigma_{\text{LO}}(4t) = 6.1 + 0.10 \tilde{C}_{tu}^{(8)} + 0.081 \tilde{C}_{tu}^{(8)2} + 0.016 \tilde{C}_{tu}^{(8)3} + 0.0048 \tilde{C}_{tu}^{(8)4}$$

o 4th power in W.C., two dim-6 insertions can neglect dim-8 for wide class of BSM 4t xsec: $-8.8 < \tilde{C}_{tu}^{(8)} < 7.1$

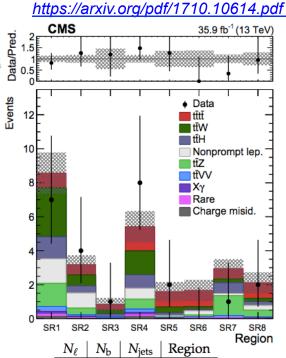
inclusive tt x-sec: $-11.8 < \tilde{C}_{tu}^{(8)} < 4.6$

- □ tt̄ inclusive
- ĭ tī m_{tt}
- tī global
- □ tttt M_{cut}=3 TeV
- \square tttt M_{cut} =4 TeV





C. Zhang, 2017 https://arxiv.org/pdf/1708.05928.pdf TOP-17-009



N_ℓ	$N_{\rm b}$	$N_{ m jets}$	Region	
		<u>≤</u> 5	CRW	
		6	SR1	
	2	7	SR2	
2		≥8	SR3	
	3	5,6	SR4	
	3		≥7	SR5
	≥4	≥7 ≥5	SR6	
≥3	2	≥5	SR7	
≥3	≥3	≥4	SR8	
Inve	Inverted Z veto		CRZ	



final remarks

- o we'll always have anomalous coupling interpretations
 - o comprehensible, quickly doable, a selling point
- o Run-II endgame: we should have (carefully planned) EFT interpretations using (multi-) differential distributions
 - \circ TT+W/Z/ γ /H:

profit from independent linear combinations

o ttbar x-sec + spin corr

constrain tt pair production & chromo moments

o Wtb/single-t + helicity

constraint the Wtb vertex

- o FCNC
- o similar to e.g. SUSY-pMSSM efforts
 - o EFT is not a simple BSM model
- o particle-level unfolding is a good compromise & defines a suitable interface to theorists

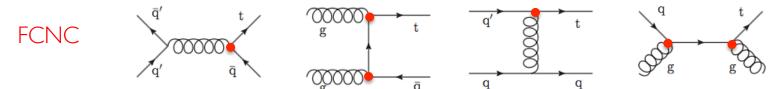


FCNC tgq (q=u/c)

23

o single top t-channel events (I µ final state)

8 TeV TOP-14-007 JHEP 02 (2017) 028 https://arxiv.org/pdf/1610.03545.pdf

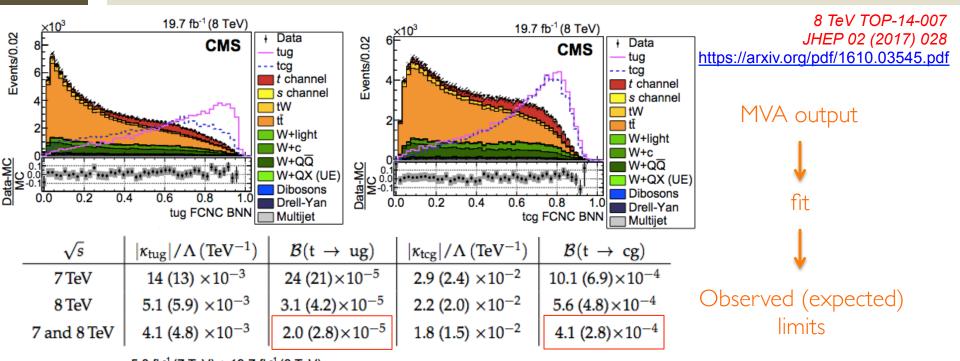


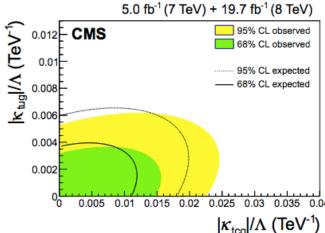
- o require I bjet, I light forward jet as predicted in t-channel
- o BNN to reject multijet BG (fit to normalize)
- o In FCNC, light forward jet is softer + more subtle effects
- o top reconstruction (W mass → neutrino z momentum)
- o BNN to distinguish FCNC from SM
 - o combine kinematic and angular observables (FCNC top softer)
 - o lepton charge & PDF effect in u & c → different BNN



FCNC tgq (q=u/c)

24





- o fit procedure to extract individual limits on \mathbf{K}_{tgq} and BR for q=u/c
- o BR(t \rightarrow ug) < 2.0 10⁻⁵ BR(t \rightarrow cg) < 4.1 10⁻⁴



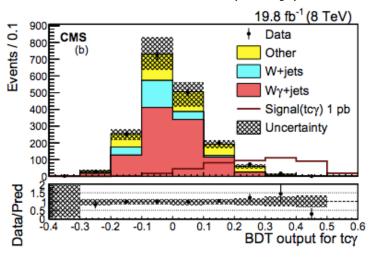
FCNC tqY

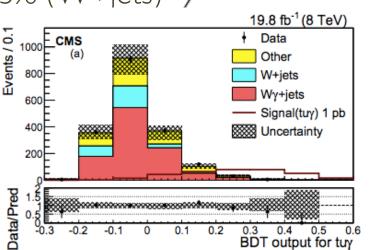
25

o single top t-channel (\rightarrow single μ final state)

8 TeV TOP-14-003 JHEP 04 (2016) 035 https://arxiv.org/abs/1511.03951

- o $p_T(\gamma) > 50$ GeV, $\leq 1b$ jet
- o top reconstruction ($p_T(\mathbf{v})$ from m_W)
- o main bkg W+jets with real or fake Y
 - o extract normalisation from NN shape fit in data: $\pm 17\%$ (WY) and $\pm 23\%$ (W+jets)





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u/c

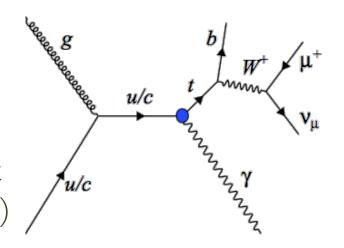


FCNC tqY

o single top t-channel (\rightarrow single μ final state)

8 TeV TOP-14-003 JHEP 04 (2016) 035 https://arxiv.org/abs/1511.03951

- o $p_T(\gamma) > 50$ GeV, $\leq 1b$ jet
- o top reconstruction ($p_T(\mathbf{V})$ from m_W)
- o main bkg W+jets with real or fake γ
 - o extract normalisation from NN shape fit in data: $\pm 17\%$ (W γ) and $\pm 23\%$ (W+jets)



	Exp. limit (LO)	$\pm 1\sigma$ (exp. limit)	$\pm 2\sigma$ (exp. limit)	Obs. limit (LO)
$\sigma_{tu\gamma}\mathcal{B}$ (fb)	40	30–56	23–78	25
$\sigma_{ m tc\gamma}\mathcal{B}$ (fb)	39	30-55	24–76	34
$\kappa_{ m tu\gamma}$	0.036	0.032 - 0.043	0.028 - 0.051	0.029
$\kappa_{ m tc\gamma}$	0.111	0.098-0.132	0.087-0.16	0.10
$\mathcal{B}(t o u \gamma)$	2.7×10^{-4}	$(2.0 - 3.8) \times 10^{-4}$	$(1.6-5.4)\times10^{-4}$	1.7×10^{-4}
$\mathcal{B}(t o c \gamma)$	2.5×10^{-3}	$(1.9 - 3.6) \times 10^{-3}$	$(1.5-4.9)\times 10^{-3}$	2.2×10^{-3}





the 59 SM EFT operators

	X^3	$arphi^6$ and $arphi^4 D^2$ $\psi^2 arphi^3$		$\psi^2 arphi^3$	
Q_G	$f^{ABC}G^{A u}_{\mu}G^{B ho}_{ u}G^{C\mu}_{ ho}$	Q_{arphi}	$(arphi^\daggerarphi)^3$	Q_{earphi}	$(arphi^\daggerarphi)(ar{l}_p e_r arphi)$
$Q_{\widetilde{G}}$	$f^{ABC}\widetilde{G}^{A u}_{\mu}G^{B ho}_{ u}G^{C\mu}_{ ho}$	$Q_{arphi\square}$	$(arphi^\daggerarphi)\Box(arphi^\daggerarphi)$	Q_{uarphi}	$(arphi^\daggerarphi)(ar q_p u_r \widetildearphi)$
Q_W	$arepsilon^{IJK}W_{\mu}^{I u}W_{ u}^{J ho}W_{ ho}^{K\mu}$	$Q_{arphi D}$	$\left(arphi^\dagger D^\mu arphi ight)^\star \left(arphi^\dagger D_\mu arphi ight)$	Q_{darphi}	$(\varphi^\dagger\varphi)(\bar{q}_pd_r\varphi)$
$Q_{\widetilde{W}}$	$arepsilon^{IJK}\widetilde{W}_{\mu}^{I u}W_{ u}^{J ho}W_{ ho}^{K\mu}$				
	$X^2 arphi^2$		$\psi^2 X \varphi$		$\psi^2 arphi^2 D$
$Q_{arphi G}$	$arphi^\dagger arphi G^A_{\mu u} G^{A\mu u}$	Q_{eW}	$(\bar{l}_p \sigma^{\mu\nu} e_r) au^I arphi W^I_{\mu u}$	$Q_{arphi l}^{(1)}$	$(arphi^\dagger i \overset{\leftrightarrow}{D}_\mu arphi) (ar{l}_p \gamma^\mu l_r)$
$Q_{arphi\widetilde{G}}$	$arphi^\dagger arphi \widetilde{G}^A_{\mu u} G^{A\mu u}$	Q_{eB}	$(ar{l}_p\sigma^{\mu u}e_r)arphi B_{\mu u}$	$Q_{arphi l}^{(3)}$	$(arphi^\dagger i \overset{\leftrightarrow}{D}_{\mu}^I arphi) (ar{l}_p au^I \gamma^\mu l_r)$
$Q_{arphi W}$	$arphi^\dagger arphi W^I_{\mu u} W^{I\mu u}$	Q_{uG}	$(ar{q}_p \sigma^{\mu u} T^A u_r) \widetilde{arphi} G^A_{\mu u}$	$Q_{arphi e}$	$(arphi^\dagger i \overleftrightarrow{D}_\mu arphi) (ar{e}_p \gamma^\mu e_r)$
$Q_{arphi\widetilde{W}}$	$arphi^\dagger arphi \widetilde{W}^I_{\mu u} W^{I \mu u}$	Q_{uW}	$(ar{q}_p \sigma^{\mu u} u_r) au^I \widetilde{arphi} W^I_{\mu u}$	$Q_{arphi q}^{(1)}$	$(arphi^\dagger i \overleftrightarrow{D}_\mu arphi) (ar{q}_p \gamma^\mu q_r)$
$Q_{arphi B}$	$arphi^\dagger arphi B_{\mu u} B^{\mu u}$	Q_{uB}	$(ar q_p \sigma^{\mu u} u_r) \widetilde arphi B_{\mu u}$	$Q_{arphi q}^{(3)}$	$(arphi^\dagger i \overset{\leftrightarrow}{D}_\mu^I arphi) (ar{q}_p au^I \gamma^\mu q_r)$
$Q_{arphi\widetilde{B}}$	$arphi^\dagger arphi \widetilde{B}_{\mu u} B^{\mu u}$	Q_{dG}	$(ar{q}_p \sigma^{\mu u} T^A d_r) arphi G^A_{\mu u}$	$Q_{arphi u}$	$(arphi^\dagger i \overleftrightarrow{D}_\mu arphi) (ar{u}_p \gamma^\mu u_r)$
$Q_{arphi WB}$	$arphi^\dagger au^I arphi W^I_{\mu u} B^{\mu u}$	Q_{dW}	$(ar{q}_p \sigma^{\mu u} d_r) au^I arphi W^I_{\mu u}$	$Q_{arphi d}$	$(arphi^\dagger i \overleftrightarrow{D}_\mu arphi) (ar{d}_p \gamma^\mu d_r)$
$Q_{arphi\widetilde{W}B}$	$arphi^\dagger au^I arphi \widetilde{W}^I_{\mu u} B^{\mu u}$	Q_{dB}	$(ar q_p \sigma^{\mu u} d_r) arphi B_{\mu u}$	$Q_{arphi ud}$	$i(\widetilde{arphi}^\dagger D_\mu arphi)(ar{u}_p \gamma^\mu d_r)$

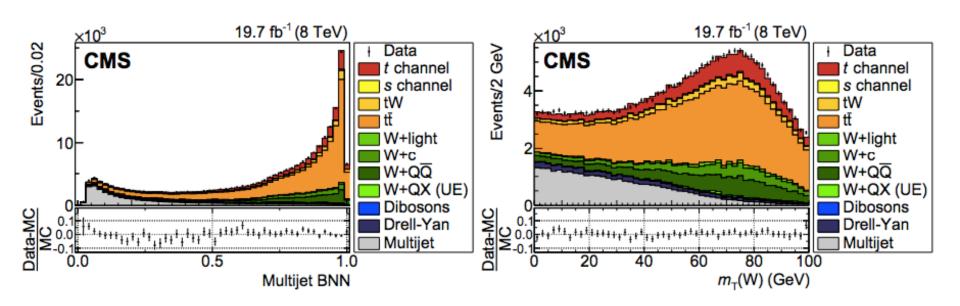
	$(\bar{L}L)(\bar{L}L)$		$(\bar{R}R)(\bar{R}R)$		$(\bar{L}L)(\bar{R}R)$
Q_{ll}	$(\bar{l}_p \gamma_\mu l_r)(\bar{l}_s \gamma^\mu l_t)$	Q_{ee}	$(ar{e}_p\gamma_\mu e_r)(ar{e}_s\gamma^\mu e_t)$	Q_{le}	$(\bar{l}_p \gamma_\mu l_r)(\bar{e}_s \gamma^\mu e_t)$
$Q_{qq}^{(1)}$	$(ar q_p \gamma_\mu q_r) (ar q_s \gamma^\mu q_t)$	Q_{uu}	$(ar{u}_p\gamma_\mu u_r)(ar{u}_s\gamma^\mu u_t)$	Q_{lu}	$(ar{l}_p\gamma_\mu l_r)(ar{u}_s\gamma^\mu u_t)$
$Q_{qq}^{(3)}$	$(ar q_p \gamma_\mu au^I q_r) (ar q_s \gamma^\mu au^I q_t)$	Q_{dd}	$(ar{d}_p\gamma_\mu d_r)(ar{d}_s\gamma^\mu d_t)$	Q_{ld}	$(ar{l}_p\gamma_\mu l_r)(ar{d}_s\gamma^\mu d_t)$
$Q_{lq}^{\left(1 ight) }$	$(ar{l}_p\gamma_\mu l_r)(ar{q}_s\gamma^\mu q_t)$	Q_{eu}	$(ar{e}_p\gamma_\mu e_r)(ar{u}_s\gamma^\mu u_t)$	Q_{qe}	$(ar{q}_p\gamma_\mu q_r)(ar{e}_s\gamma^\mu e_t)$
$Q_{lq}^{(3)}$	$(ar{l}_p \gamma_\mu au^I l_r) (ar{q}_s \gamma^\mu au^I q_t)$	Q_{ed}	$(ar{e}_p\gamma_\mu e_r)(ar{d}_s\gamma^\mu d_t)$	$Q_{qu}^{(1)}$	$(ar q_p \gamma_\mu q_r) (ar u_s \gamma^\mu u_t)$
		$Q_{ud}^{(1)}$	$(ar{u}_p\gamma_\mu u_r)(ar{d}_s\gamma^\mu d_t)$	$Q_{qu}^{(8)}$	$(ar{q}_p \gamma_\mu T^A q_r) (ar{u}_s \gamma^\mu T^A u_t)$
		$Q_{ud}^{(8)}$	$(ar{u}_p \gamma_\mu T^A u_r) (ar{d}_s \gamma^\mu T^A d_t)$	$Q_{qd}^{(1)}$	$(ar{q}_p\gamma_\mu q_r)(ar{d}_s\gamma^\mu d_t)$
				$Q_{qd}^{(8)}$	$(ar{q}_p \gamma_\mu T^A q_r) (ar{d}_s \gamma^\mu T^A d_t)$
$(\bar{L}R)$	$(\bar{R}L)$ and $(\bar{L}R)(\bar{L}R)$		B-viol	ating	
Q_{ledq}	$(ar{l}_p^j e_r) (ar{d}_s q_t^j)$	Q_{duq}	$\left[q - rac{arepsilon^{lphaeta\gamma}arepsilon_{jk}\left[(d_p^lpha)^TCu_r^eta ight]\left[(q_s^{\gamma j})^TCl_t^k ight]}{\left[(d_p^lpha)^TCu_r^eta ight]}$		
$Q_{quqd}^{(1)}$	$(ar{q}_p^j u_r) arepsilon_{jk} (ar{q}_s^k d_t)$	Q_{qqu}	$arepsilon^{lphaeta\gamma}arepsilon_{jk}\left[(q_p^{lpha j})^TCq_r^{eta k} ight]\left[(u_s^{\gamma})^TCe_t ight]$		
$Q_{quqd}^{(8)}$	$(ar{q}_p^j T^A u_r) arepsilon_{jk} (ar{q}_s^k T^A d_t)$	Q_{qqq}	$arepsilon^{lphaeta\gamma}arepsilon_{jn}arepsilon_{km}\left[(q_p^{lpha j})^TCq_r^{eta k} ight]\left[(q_s^{\gamma m})^TCl_t^n ight]$		
$Q_{lequ}^{(1)}$	$(ar{l}_p^j e_r) arepsilon_{jk} (ar{q}_s^k u_t)$	Q_{duu}	$arepsilon^{lphaeta\gamma}\left[(d_p^lpha)^TCu_r^eta ight]\left[(u_s^\gamma)^TCe_t ight]$		
$Q_{lequ}^{(3)}$	$(ar{l}_p^j\sigma_{\mu u}e_r)arepsilon_{jk}(ar{q}_s^k\sigma^{\mu u}u_t)$				

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FCNC tgq

o multijet BNN discriminator



FCNC tHq, $H \rightarrow bb$



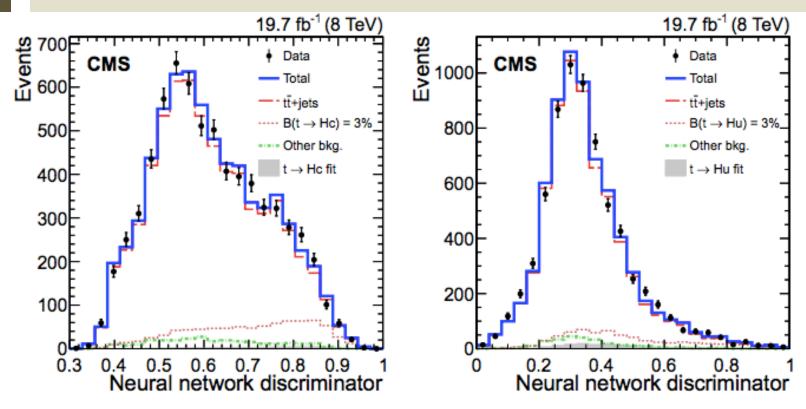


Figure 6: The output distributions from the ANN discriminator for data (points) and simulated background (lines) where the ANN was trained to discriminate the backgrounds from either $t \to Hc$ (left) or $t \to Hu$ (right) decays. The solid line shows the result of the fit of the signal and background templates to data. The dotted line gives the predicted signal distribution from simulation for $\mathcal{B}(t \to Hc) = 3\%$ and the filled histogram shows the proportion of signal estimated from the fit.



tZq individual limits

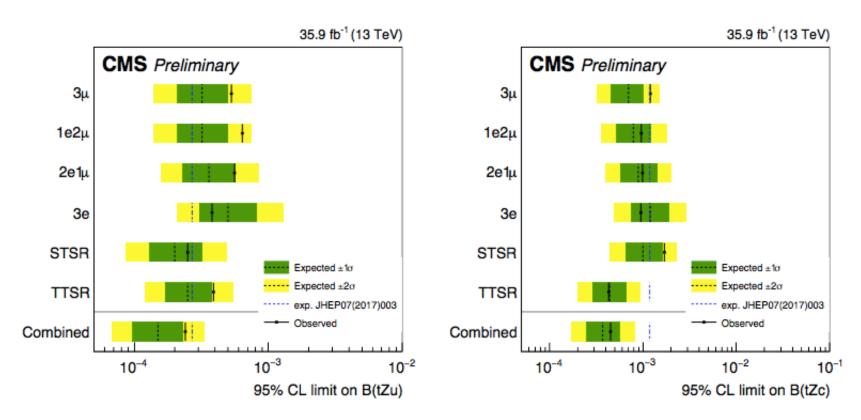


Figure 6: Exclusion limits at 95% CL for each leptonic channel and signal region on the FCNC tZu (left) and tZc (right) branching fractions considering one non-vanishing coupling at a time. The CMS 8 TeV [7] observed (expected) limit of is given with a blue dashed line).