



# FCNC, EFT, Anomalous couplings at ATLAS

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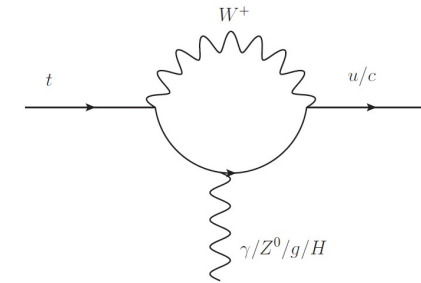
On Behalf of the ATLAS Collaboration



# Flavor Changing Neutral Currents

Not allowed at tree level in SM.

Allowed at loop-level, but GIM suppressed.  
Branching Ratio  $\sim 10^{-15} \rightarrow 10^{-12}$



BSM interactions could lead to much larger effects  
Branching Ratio up to  $10^{-4}$

There have been many searches for top FCNC in ATLAS.  
I will present the ones that currently provide our best limits



The Lagrangian for top FCNC can be written in terms of a set of anomalous couplings:

$$\begin{aligned}\mathcal{L} = \sum_{q=u,c} & \left[ \sqrt{2}g_s \frac{\kappa_{gqt}}{\Lambda} \bar{t} \sigma^{\mu\nu} T_a (f_{Gq}^L P_L + f_{Gq}^R P_R) q G_{\mu\nu}^a + \right. \\ & + \frac{g}{\sqrt{2}c_W} \frac{\kappa_{zqt}}{\Lambda} \bar{t} \sigma^{\mu\nu} (f_{Zq}^L P_L + f_{Zq}^R P_R) q Z_{\mu\nu} + \frac{g}{4c_W} \zeta_{zqt} \bar{t} \gamma^\mu (f_{Zq}^L P_L + f_{Zq}^R P_R) q Z_\mu - \\ & - e \frac{\kappa_{\gamma qt}}{\Lambda} \bar{t} \sigma^{\mu\nu} (f_{\gamma q}^L P_L + f_{\gamma q}^R P_R) q A_{\mu\nu} + \\ & \left. + \frac{g}{\sqrt{2}} \bar{t} \kappa_{Hqt} (f_{Hq}^L P_L + f_{Hq}^R P_R) q H \right] + h.c.\end{aligned}$$

These couplings can also be expressed in terms of EFT coefficients, or combinations of coefficients



# $t \rightarrow qH$

*JHEP 10 (2017) 129*  
[arXiv:1707.01404](https://arxiv.org/abs/1707.01404)

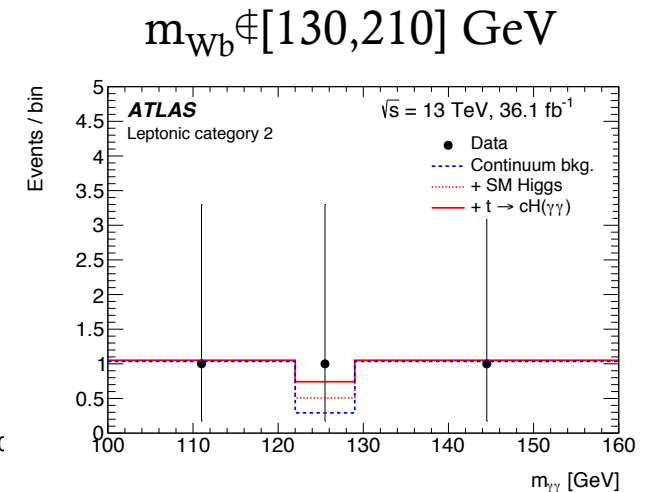
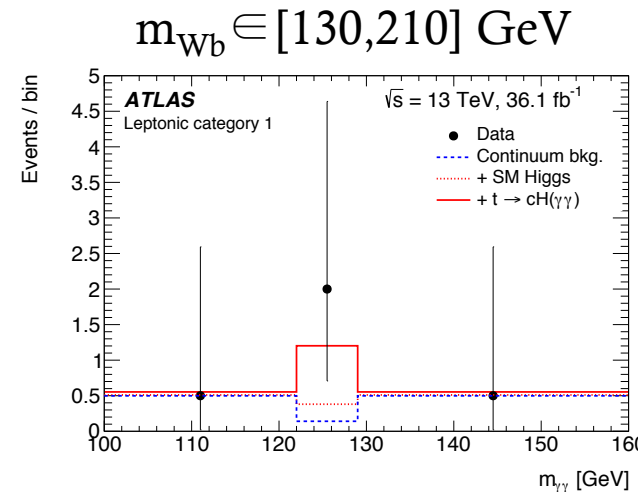
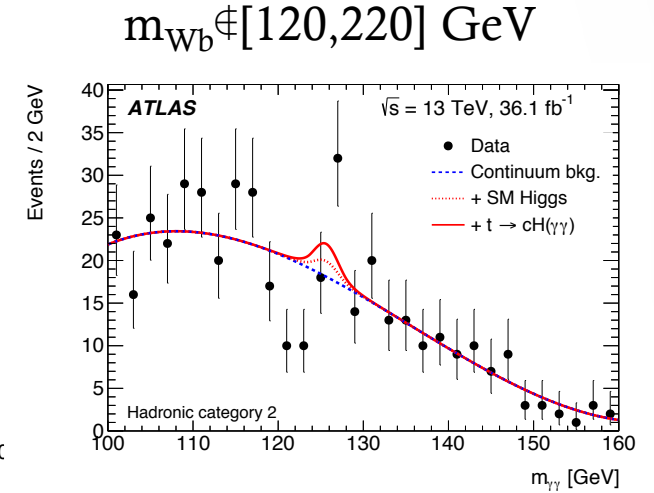
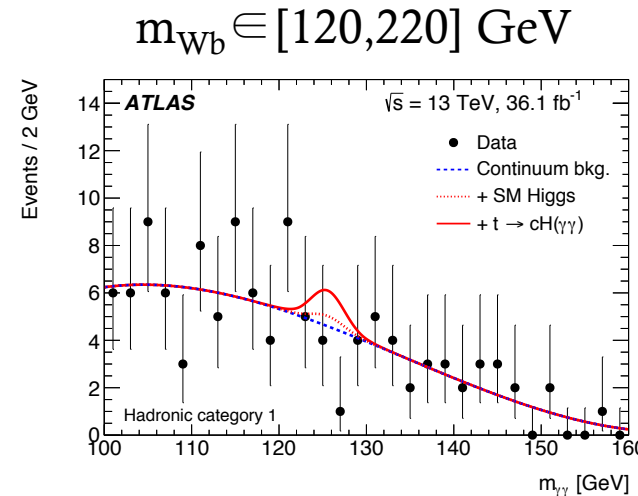


Search for FCNC in  $t\bar{t}$  events.  
One top decays to  $Wb$ ,  
the other to  $qH$ ,  $H \rightarrow \gamma\gamma$

$$m_{\gamma\gamma j} \in [152, 190]$$

Events categorized by whether

1.  $W$  decays to jets or leptons.
2.  $m_{Wb}$  inside top mass window or not.







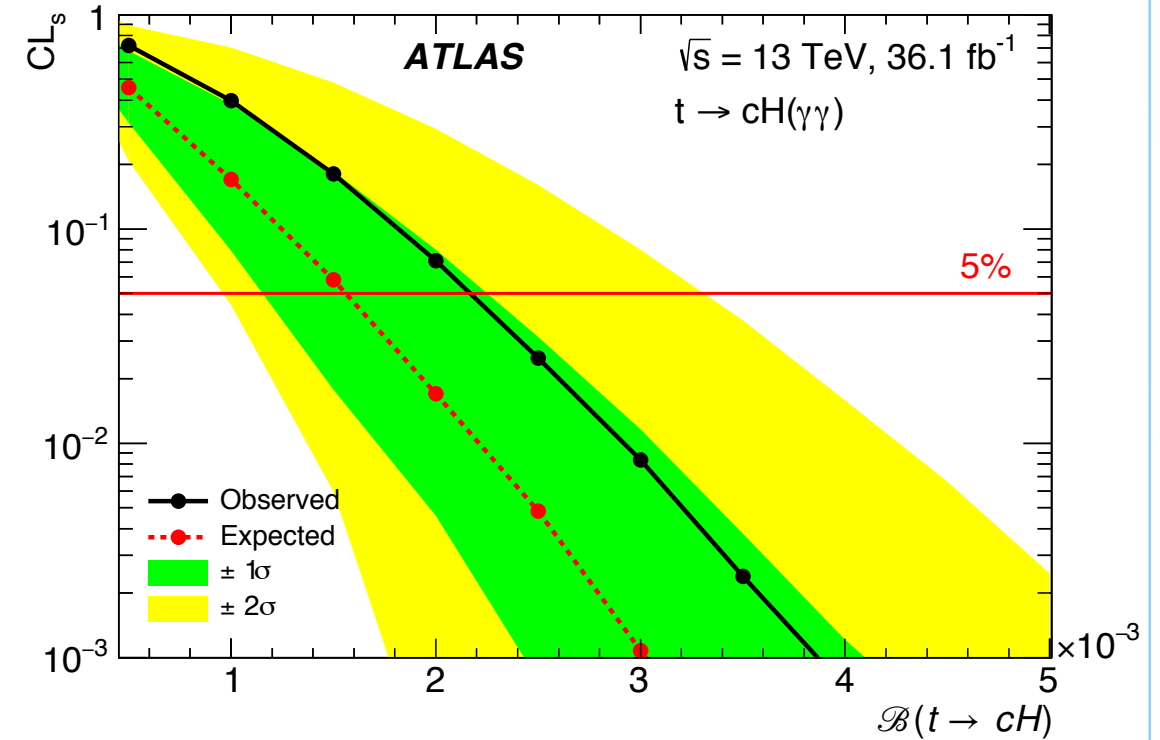
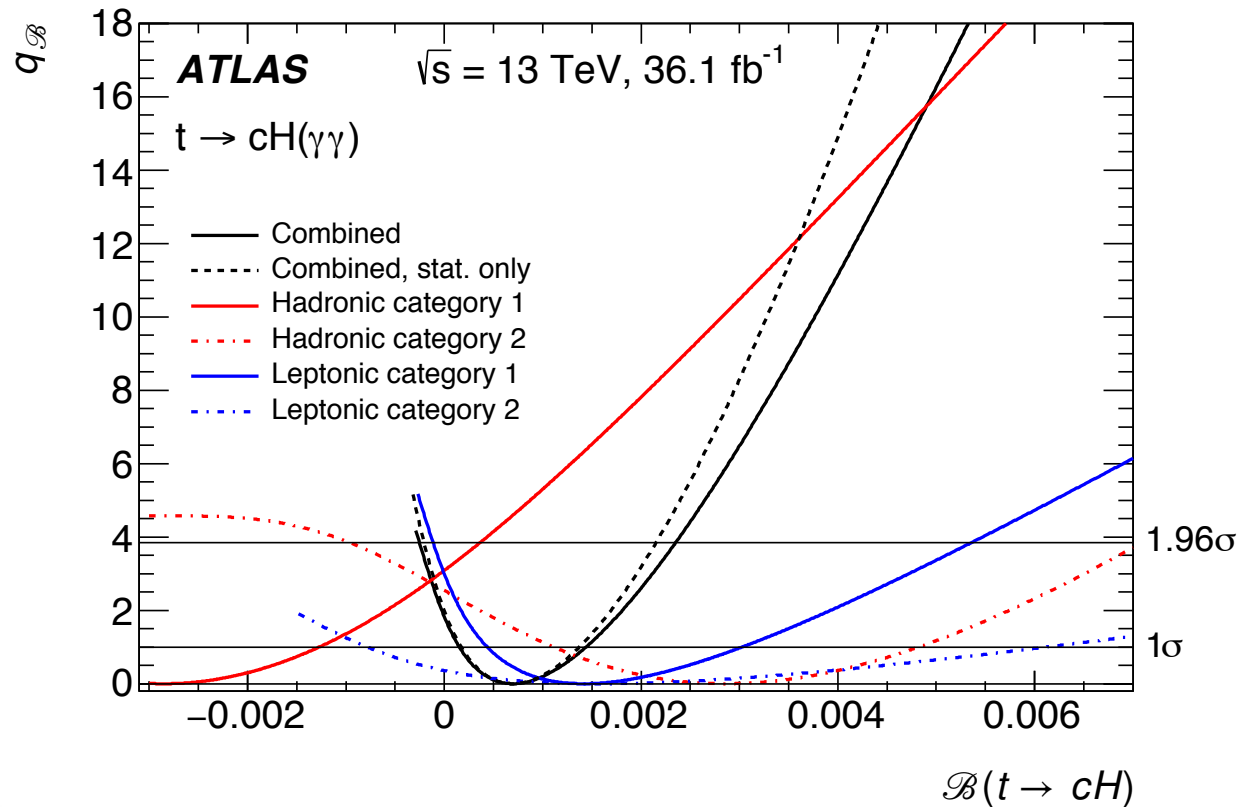
# $t \rightarrow qH$

*JHEP 10 (2017) 129*  
*arXiv:1707.01404*



$$q_B = -2(\ln L(\mathcal{B}) - \ln L(\hat{\mathcal{B}}))$$

Shown for each individual category and the combination.



$B(t \rightarrow Hu) < 0.24\%$  (obs)  
 $< 0.17\%$  (exp)  
 $B(t \rightarrow Hc) < 0.22\%$  (obs)  
 $< 0.16\%$  (exp)



# limits extracted from $t \rightarrow qH$



*[JHEP 10 \(2017\) 129](#)*  
*[arXiv:1707.01404](#)*

The branching ratio limits have been used to obtain limits on the off-diagonal Yukawa couplings

$$\sqrt{|\lambda_{tcH}|^2 + 0.92|\lambda_{tuH}|^2} < 0.090$$

Where each  $\lambda$  is the quadratic sum of the two possible chirality couplings.

This could be recast as a limit on a quadratic combination of all the relevant EFT coefficients  $C_{u\phi}^{13}$ ,  $C_{u\phi}^{31*}$ ,  $C_{u\phi}^{23}$ , and  $C_{u\phi}^{32*}$ .



$t \rightarrow qZ$

ATLAS-CONF-2017-070

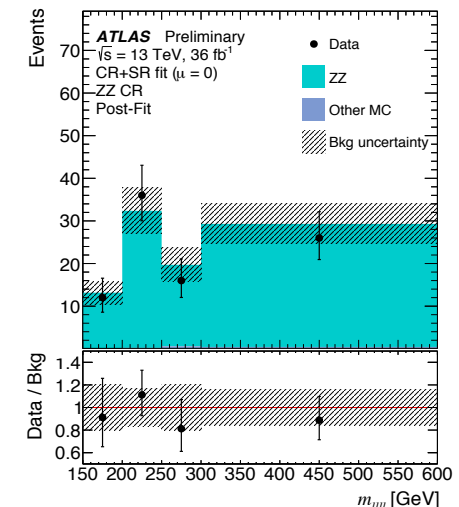
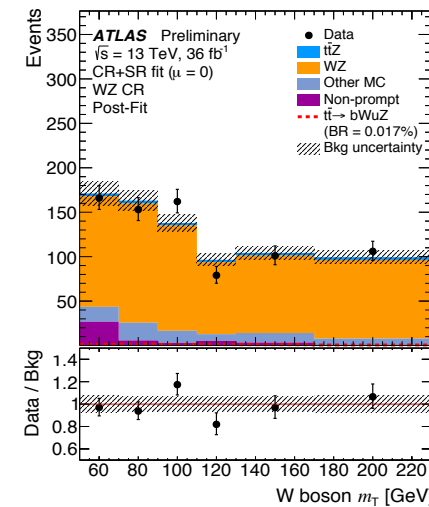
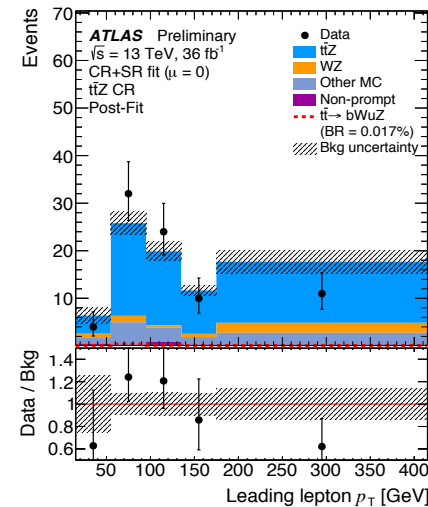
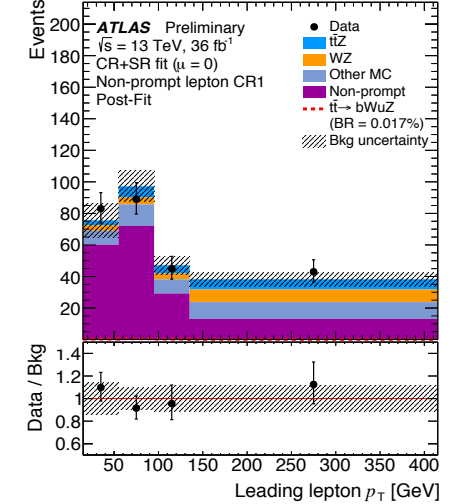
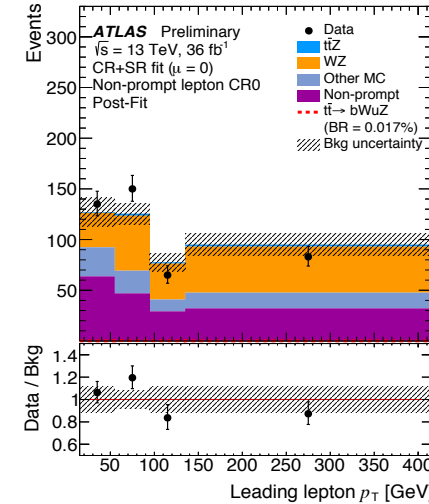
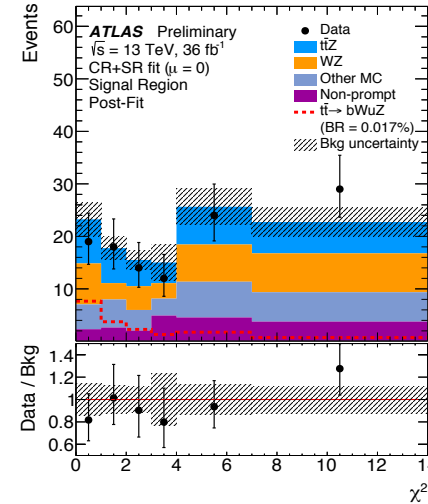


Search for FCNC in  $t\bar{t}$  events.  
One top decays to  $Wb$ ,  $W \rightarrow l\nu$   
the other to  $qZ$ ,  $Z \rightarrow l^+l^-$   
3 leptons plus 2 jets; 1 b-tagged,  
plus MET.

Form  $\chi^2$  based on reconstructed  
 $m_{jll}$ ,  $m_{bl\nu}$ , and  $m_{l\nu}$ .

Simultaneous fit to 6 distributions  
in Signal and Control Regions

Main syst due to background  
Modeling



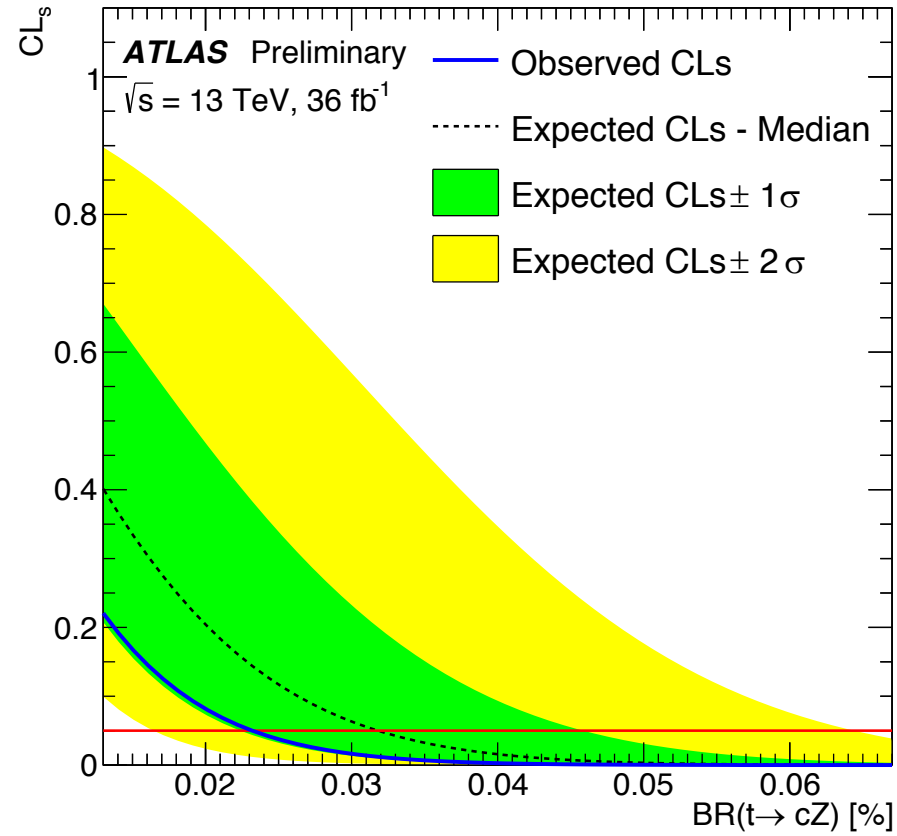
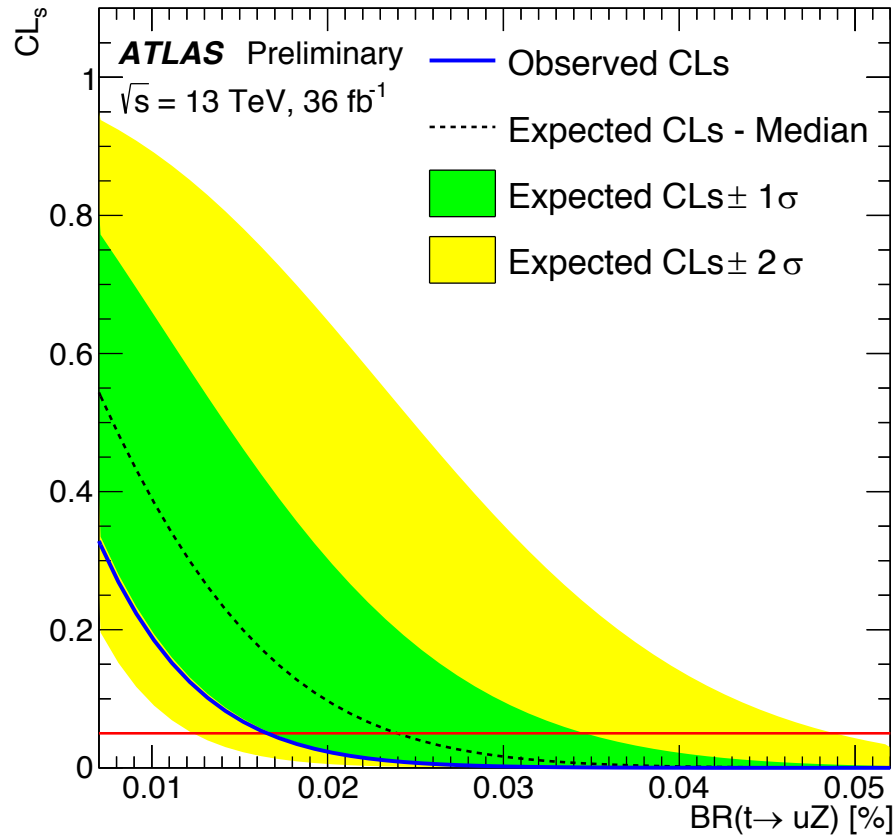


# $t \rightarrow qZ$

ATLAS-CONF-2017-070



## Results



$B(t \rightarrow Zu) < 0.017\%$  (obs)  
 $< 0.024\%$  (exp)  
 $B(t \rightarrow Zc) < 0.023\%$  (obs)  
 $< 0.032\%$  (exp)



# EFT limits from $t \rightarrow qZ$



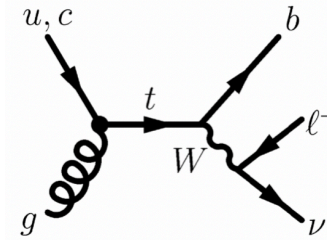
*Eur. Phys. J. C (2016) 76:55*  
[arXiv:1509.00294](https://arxiv.org/abs/1509.00294)

- Limits on EFT coefficients related to the tensor  $Ztq$  couplings have also been extracted from the branching ratio limits.
- Each limit assumes all other coefficients are 0.
  - Since the tensor couplings depend on linear combinations of  $C_{uW}$  and  $C_{uB}$ , it would be best to show these correlations.
  - The  $\gamma tq$  couplings depend on a different combination of couplings, so a combined analysis would be more powerful.
- This also assumes the vector  $Ztq$  couplings are zero. How does this affect limits obtained on the tensor couplings?

Operator	Observed	Expected
$ C_{uB}^{(31)} $	0.25	0.30
$ C_{uW}^{(31)} $	0.25	0.30
$ C_{uB}^{(32)} $	0.30	0.34
$ C_{uW}^{(32)} $	0.30	0.34

$t \rightarrow qg$

Search for FCNC in exclusive single-top events.  
isolated lepton, MET, exactly 1 b-tagged jet.

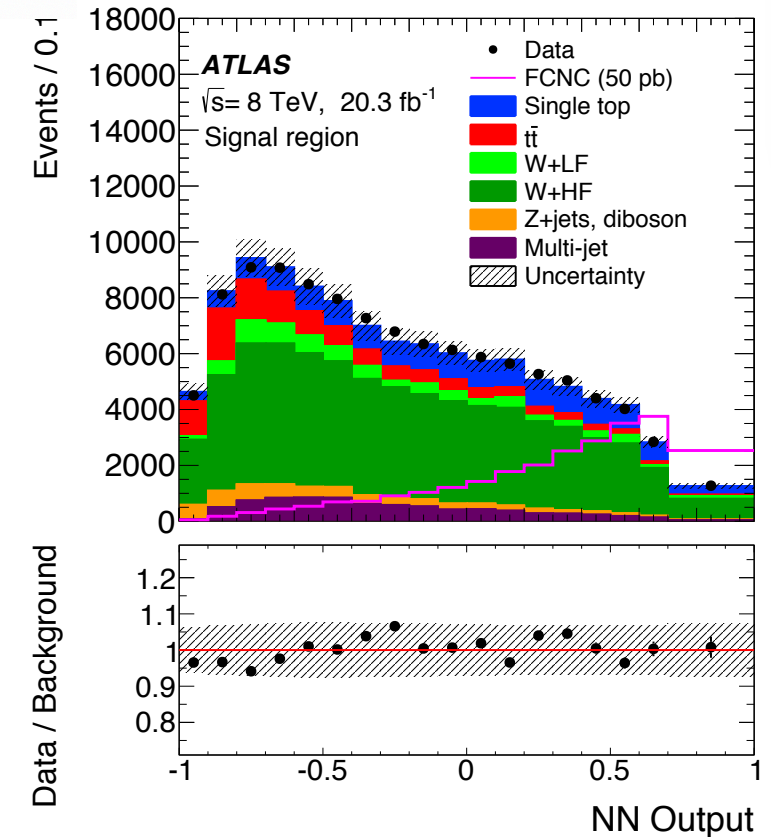
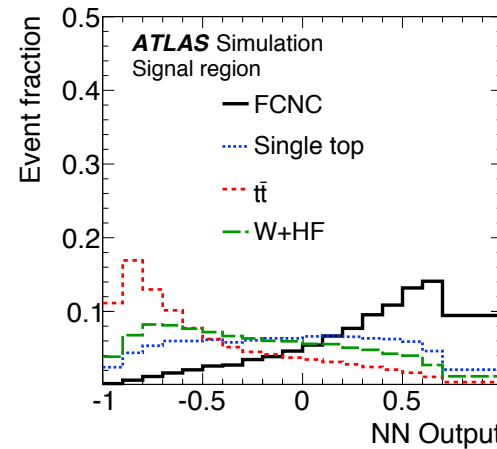
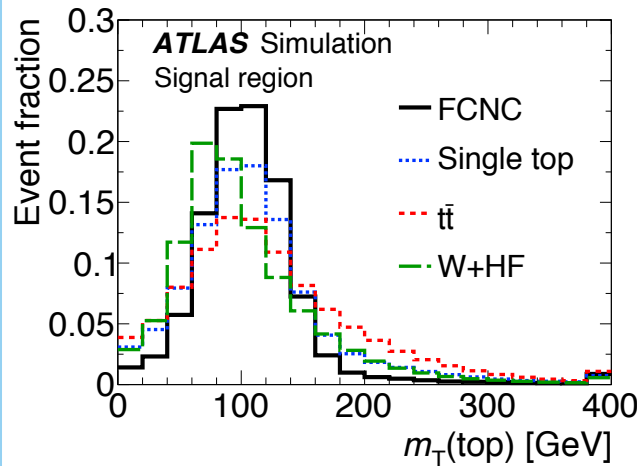


*Eur. Phys. J. C (2016) 76:55*  
[arXiv:1509.00294](https://arxiv.org/abs/1509.00294)

W+jets suppressed with Neural Net.

13 variables, most important:

Transverse mass of reconstructed top





$t \rightarrow qg$

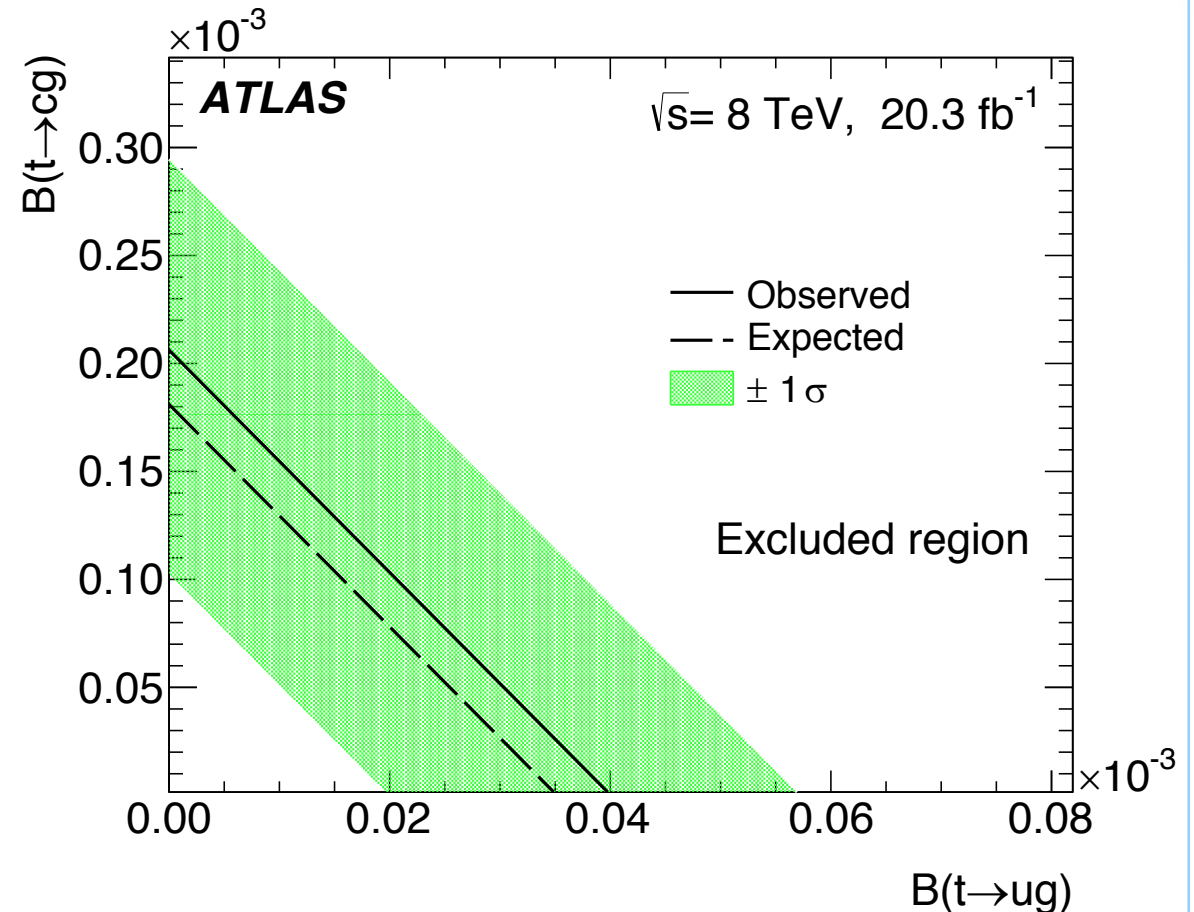


Result shows excluded region in plane of Branching Ratios.  
Shaded region shows one-sigma variation of expected limit.

*Eur. Phys. J. C (2016) 76:55*  
[arXiv:1509.00294](https://arxiv.org/abs/1509.00294)

$$B(t \rightarrow ug) < 4 \times 10^{-5} \text{ (obs)}$$
$$B(t \rightarrow cg) < 20 \times 10^{-5} \text{ (obs)}$$

Largest systematic uncertainties due to  
Missing  $E_T$  modeling  
Jet Energy resolution







# coupling limits extracted from $t \rightarrow qg$



*Eur. Phys. J. C (2016) 76:55*

[arXiv:1509.00294](https://arxiv.org/abs/1509.00294)

Assumes only left-handed couplings.

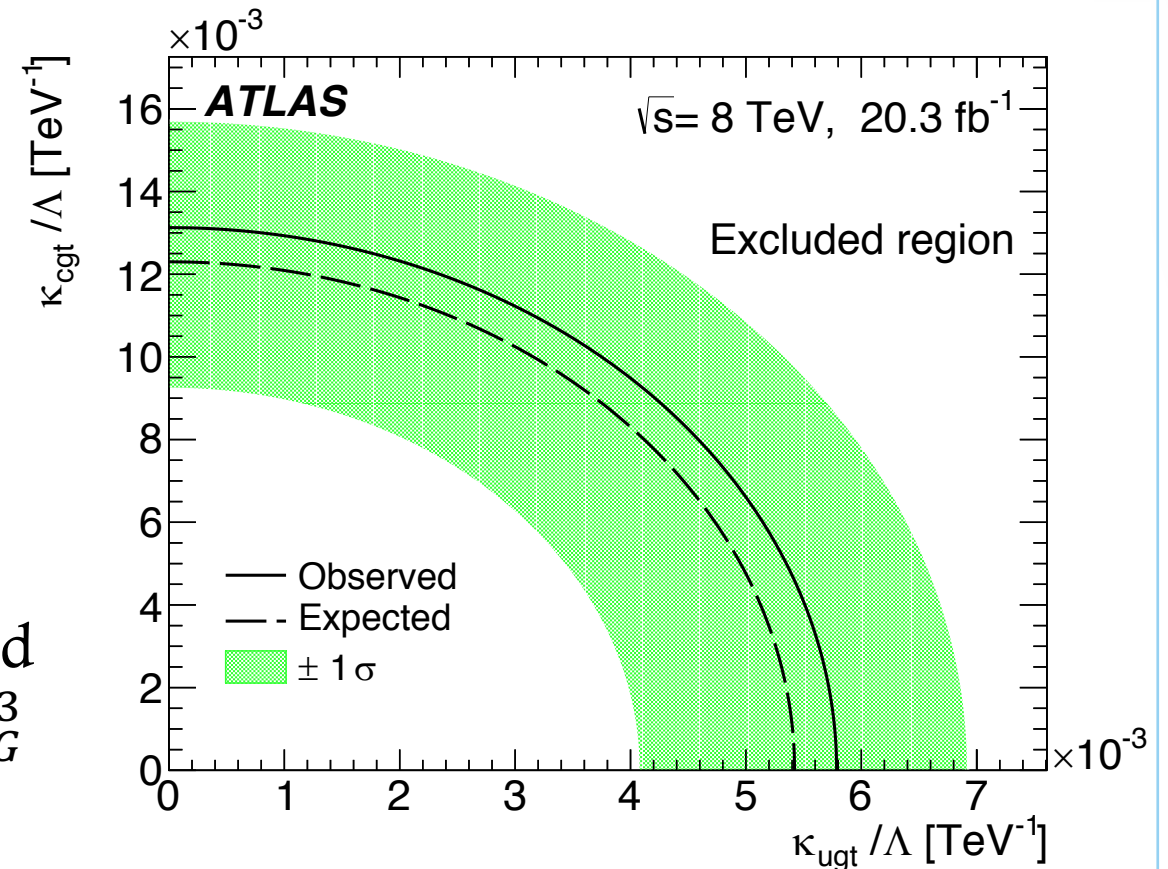
Conservative limit as right-handed couplings would be easier to distinguish from SM top production.

Assuming all other couplings are SM, one can quote

$$\kappa_{ugt}/\Lambda < 5.8 \times 10^{-3}, \kappa_{cgt}/\Lambda < 13 \times 10^{-3}$$

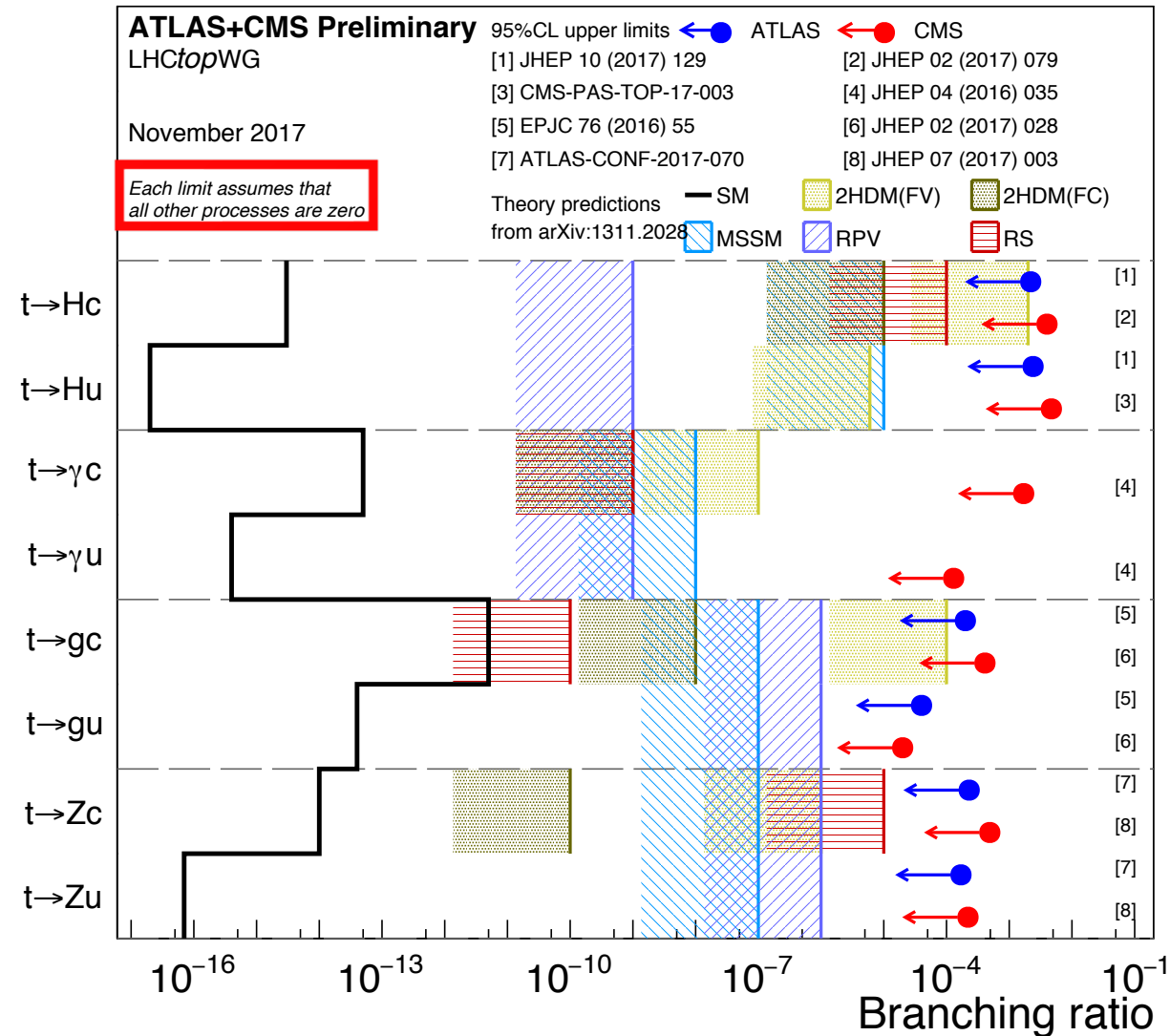
In future, perhaps we can quote how limits vary with assumed right- vs left-handed couplings.

Left handed couplings could also be expressed in terms of  $C_{uG}^{31*}$ ,  $C_{uG}^{32*}$ ; right-handed:  $C_{uG}^{13}$ ,  $C_{uG}^{23}$





# Summary of all Branching Ratio limits





# Other constraints on couplings

The angular correlations, mentioned by Ben yesterday, have also been used to extract anomalous couplings associated with the  $Wtb$  vertex.

$$\mathcal{L}_{Wtb} = -\frac{g}{\sqrt{2}} \bar{b} \gamma^\mu (V_L P_L + V_R P_R) t W_\mu^- - \frac{g}{\sqrt{2}} \bar{b} \frac{i \sigma^{\mu\nu} q_\nu}{m_W} (g_L P_L + g_R P_R) t W_\mu^- + \text{h.c.},$$

These can be associated with the  $V_{tb}$  from the SM plus four EFT coefficients.

$$V_L = V_{tb} + C_{\phi q}^{(3,3+3)} \frac{v^2}{\Lambda^2}; \quad V_R = \frac{1}{2} C_{\phi\phi}^{3,3*} \frac{v^2}{\Lambda^2}; \quad g_L = \sqrt{2} C_{dW}^{3,3*} \frac{v^2}{\Lambda^2}; \quad g_R = \sqrt{2} C_{uW}^{3,3} \frac{v^2}{\Lambda^2}$$



# W Helicity Fractions

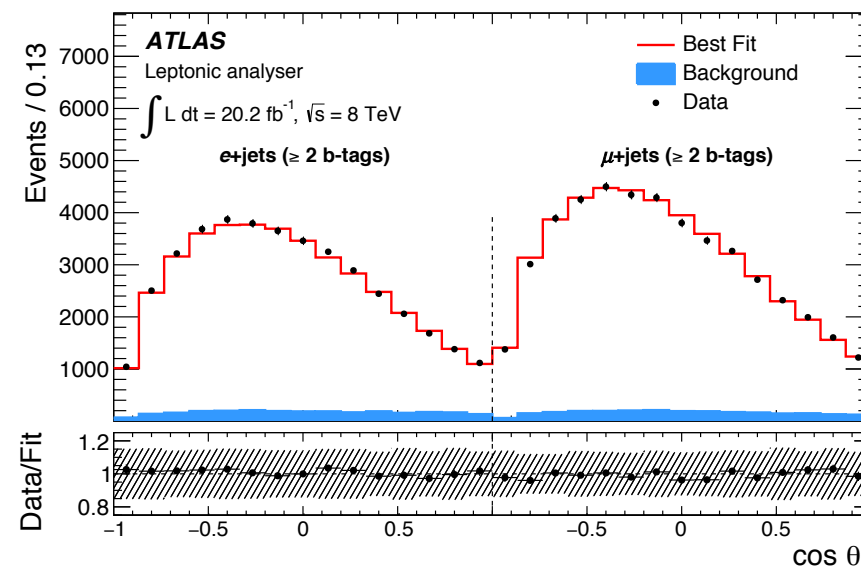
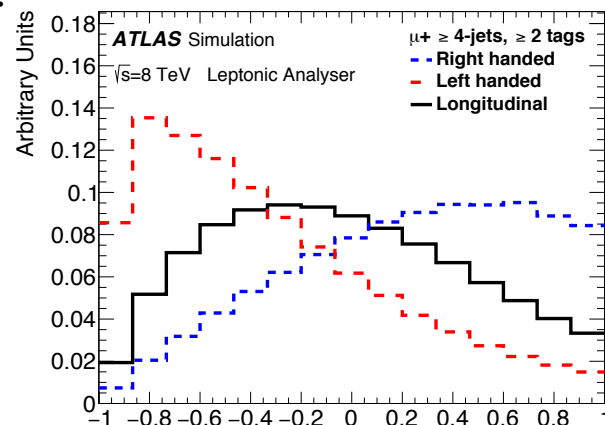


*Eur. Phys. J. C 77 (2017) 264*  
[arXiv:1612.02577](https://arxiv.org/abs/1612.02577)

The standard W Helicity fractions are extracted in  $t\bar{t}$  events.  
 1 lepton plus 4 jets (2 b-tagged) plus missing  $E_T$ .

Template fit to distribution in angle between lepton and helicity axis in W rest frame.

Systematic uncertainty somewhat larger than statistical.  
 Largest systematic for  $F_0$  is MC template statistics.



$$F_0 = 0.709 \pm 0.012 \text{ (stat.+bkg. norm.)}^{+0.015}_{-0.014} \text{ (syst.)}$$

$$F_L = 0.299 \pm 0.008 \text{ (stat.+bkg. norm.)}^{+0.013}_{-0.012} \text{ (syst.)}$$

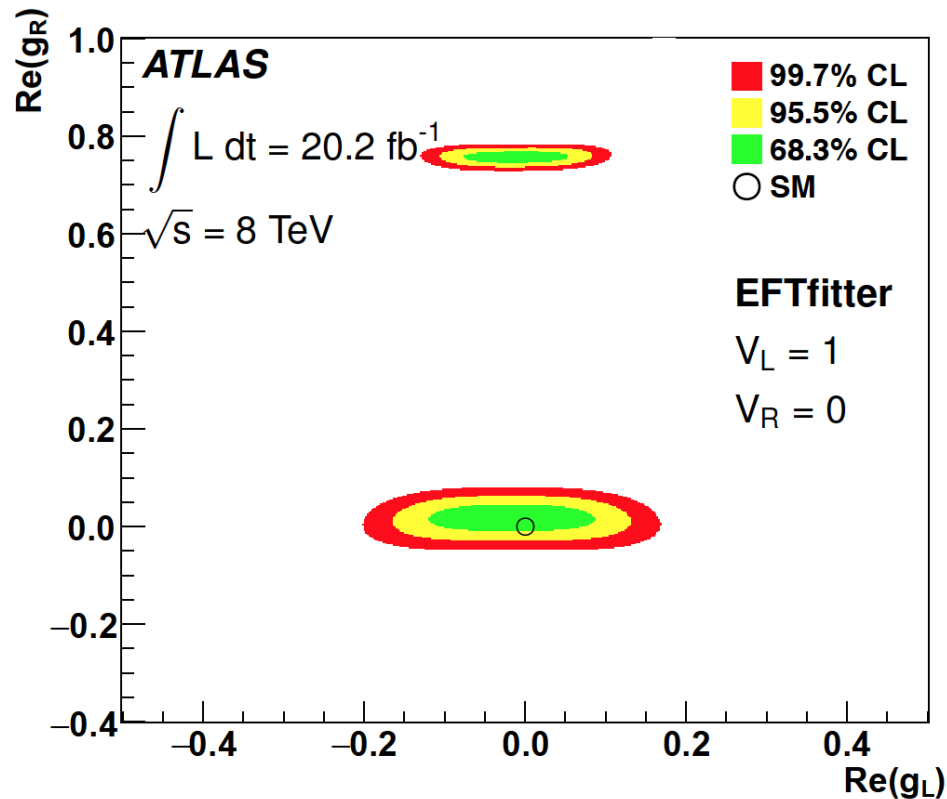
$$F_R = -0.008 \pm 0.006 \text{ (stat.+bkg. norm.)} \pm 0.012 \text{ (syst.)}$$



# Constraints from W-Helicity Fractions



*Eur. Phys. J. C 77 (2017) 264*  
[arXiv:1612.02577](https://arxiv.org/abs/1612.02577)



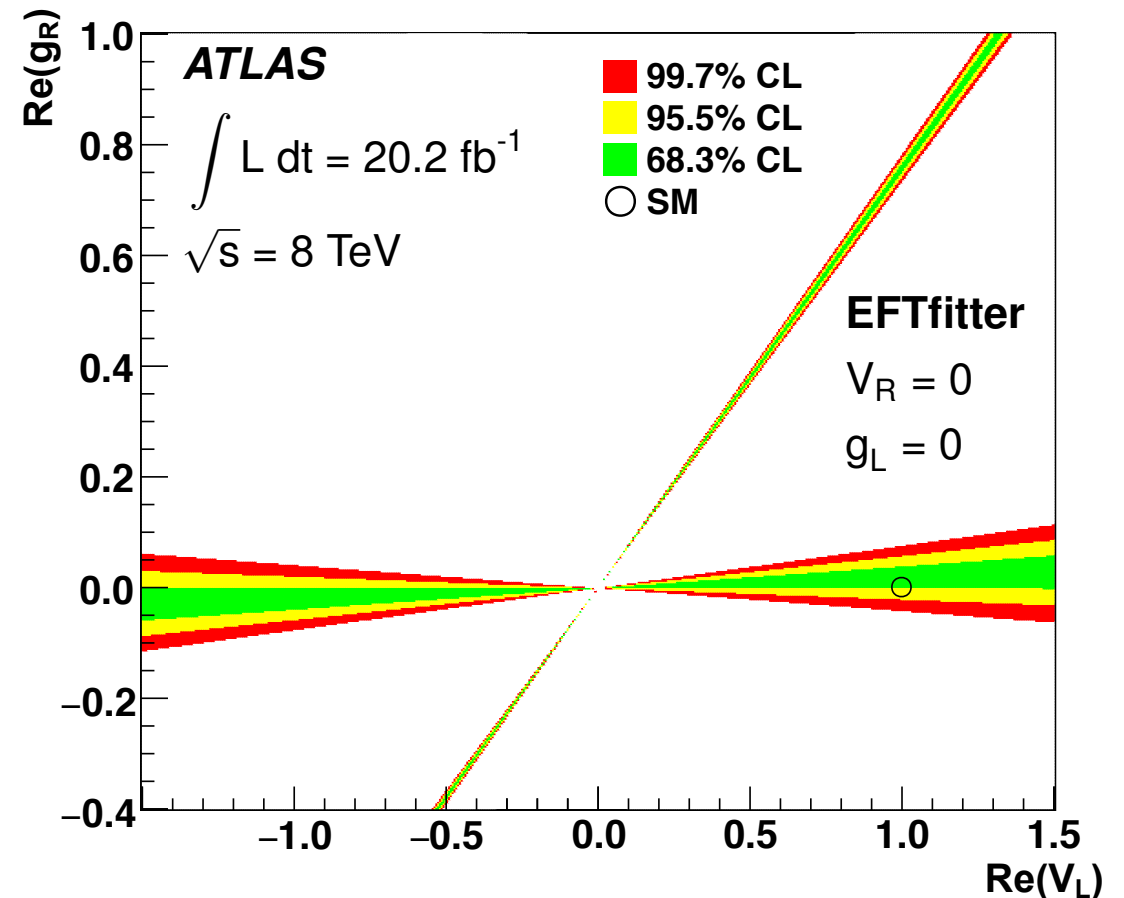
W-helicity measurements can be used to constrain anomalous couplings. This plot shows tight constraints in  $g_R$  vs  $g_L$ .

*There are assumptions that the two coupling not shown are fixed to their SM values, and the ones shown are purely real. If these assumptions are removed, most of the plane is allowed.*

If one examines a different set, of couplings one see that helicity *fractions* are sensitive to ratios of couplings. They very precisely determine two allowed regions in  $g_R/V_L$ .

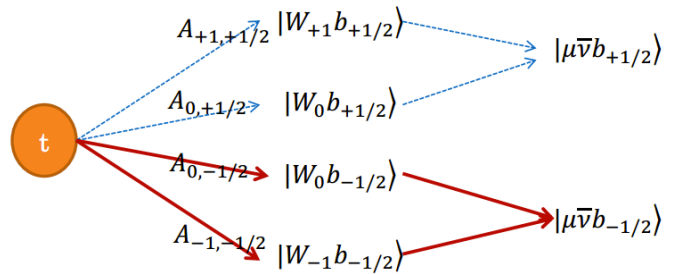
Again, if one releases the assumptions on the other two couplings, the interpretation becomes more complicated.

*Eur. Phys. J. C 77 (2017) 264*  
[arXiv:1612.02577](https://arxiv.org/abs/1612.02577)



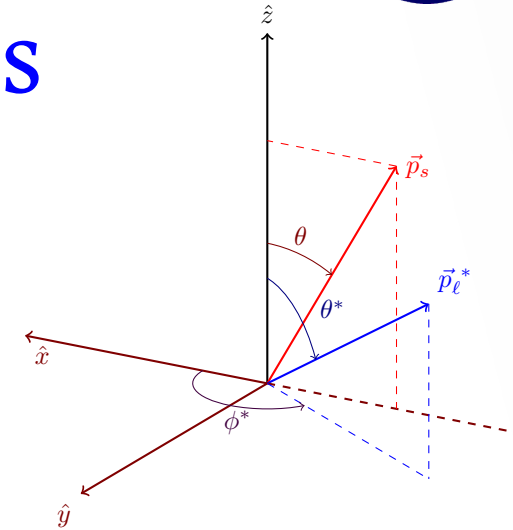


# Top decay distribution can be fully expressed in terms of 4 Helicity amplitudes



$$A_{\lambda_W, \lambda_b} = \begin{cases} \lambda_b = 1/2, \lambda_W = 0, 1. \\ \lambda_b = -1/2, \lambda_W = 0, -1. \end{cases}$$

$\lambda$  represents the helicity fractions of  $b$  quark and  $W$  boson.  $A_{\lambda_W, \lambda_b}$  is the decay amplitude.



$$\begin{aligned} \frac{1}{N} \frac{dN}{d\Omega d\Omega^*} &= \frac{1}{(4\pi)^2} \left[ \frac{3}{4} |A_{1, \frac{1}{2}}|^2 (1 + P \cos \theta) (1 + \cos \theta^*)^2 \right. \\ &+ \frac{3}{4} |A_{-1, -\frac{1}{2}}|^2 (1 - P \cos \theta) (1 - \cos \theta^*)^2 \\ &+ \frac{3}{2} \left( |A_{0, \frac{1}{2}}|^2 (1 - P \cos \theta) + |A_{0, -\frac{1}{2}}|^2 (1 + P \cos \theta) \right) \sin^2 \theta^* \\ &- \frac{3\sqrt{2}}{2} P \sin \theta \sin \theta^* (1 + \cos \theta^*) \Re \left( e^{i\phi^*} A_{1, \frac{1}{2}} A_{0, \frac{1}{2}}^* \right) \\ &\left. - \frac{3\sqrt{2}}{2} P \sin \theta \sin \theta^* (1 - \cos \theta^*) \Re \left( e^{-i\phi^*} A_{-1, -\frac{1}{2}} A_{0, -\frac{1}{2}}^* \right) \right], \end{aligned}$$

The full triple differential decay distribution is determined in t-channel single-top events by measuring coefficients of orthogonal functions of  $\theta$ ,  $\theta^*$ , and  $\phi^*$ .

***JHEP 12 (2017) 017***  
***arXiv:1707.05393***

Terms in Black: measured by W-Helicity analyses

Terms in Green: also measure relative phase between amplitudes

Terms in Red: measures polarization and different combinations of amplitudes





# Extracted Observables



*JHEP* 12 (2017) 017  
arXiv:1707.05393

Likelihood is calculated numerically in full 6-d space of parameters

Two 2-d and one 1-d profiles shown here

$$f_1 = \frac{|A_{-1,-\frac{1}{2}}|^2 + |A_{1,\frac{1}{2}}|^2}{|A_{-1,-\frac{1}{2}}|^2 + |A_{0,-\frac{1}{2}}|^2 + |A_{0,\frac{1}{2}}|^2 + |A_{1,\frac{1}{2}}|^2} = F_R + F_L = 1 - F_0$$

$$f_1^+ = \frac{|A_{1,\frac{1}{2}}|^2}{|A_{-1,-\frac{1}{2}}|^2 + |A_{1,\frac{1}{2}}|^2} = \frac{F_R}{F_R + F_L}$$

$$f_0^+ = \frac{|A_{0,\frac{1}{2}}|^2}{|A_{0,-\frac{1}{2}}|^2 + |A_{0,\frac{1}{2}}|^2}$$

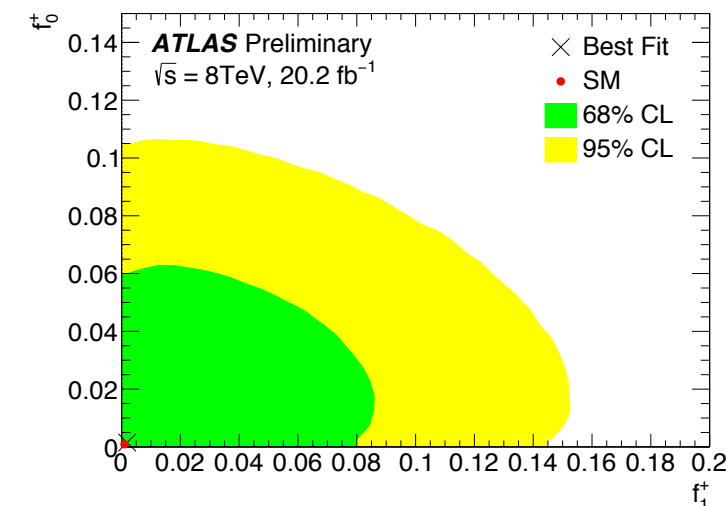
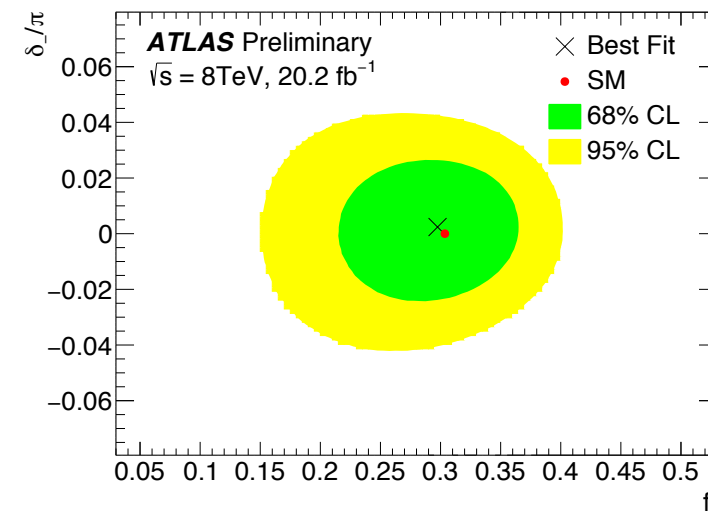
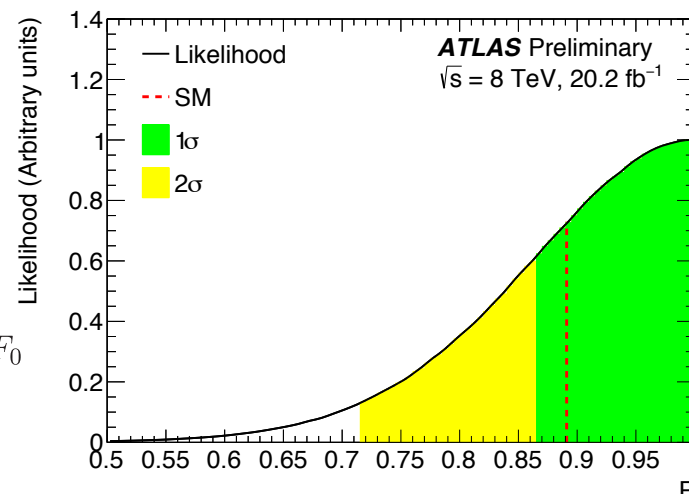
$$\delta_- = \arg A_{-1,-\frac{1}{2}} A_{0,-\frac{1}{2}}^*$$

$$\delta_+ = \arg A_{1,-\frac{1}{2}} A_{0,\frac{1}{2}}^*$$

$P$  = top polarization

Not previously measured

Non-zero phase could imply CP-violation

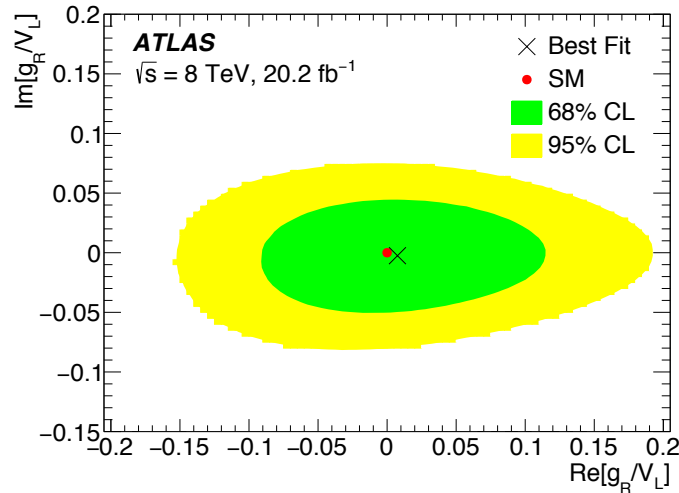




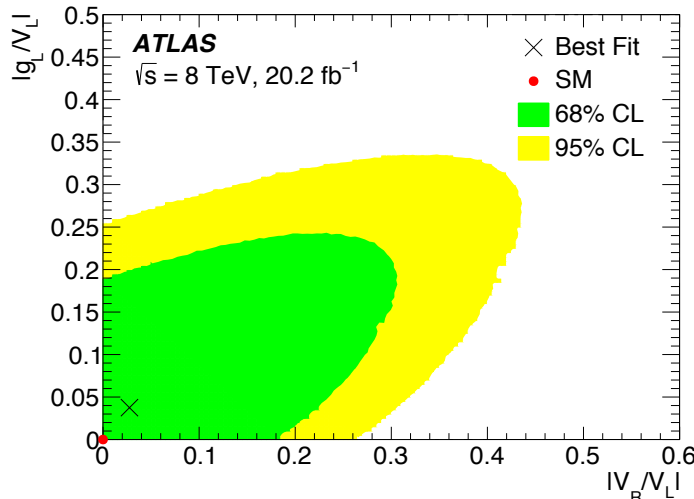
# Constraints from triple-differential top decay



*JHEP 12 (2017) 017*  
[arXiv:1707.05393](https://arxiv.org/abs/1707.05393)



In the analysis with single-top, all assumptions on the couplings are removed, we explicitly state the result as ratios of couplings, and  $V_R$ ,  $g_L$ , and  $g_R$  are allowed to be complex.



$f_1^+$  and  $f_0^+$  are determined by different linear combinations of  $g_L$  and  $V_R$ . In this case it is not conservative to fix one to its SM value when setting limits on the other. One needs the full correlation.



# Summary



- ATLAS has many top FCNC results.
  - The full Run 2 data set will allow for more sensitivity to BSM physics.
    - Will need to keep control of systematics for  $t \rightarrow Zq$  and  $t \rightarrow gq$
- The FCNC results have been used to provide constraints on couplings/EFT coefficients.
- Anomalous couplings/EFT also extracted from W Helicity measurements in  $t\bar{t}$  events and from triple differential decay distributions in single-top events.
  - Correlations are important.
- A top sub-group has been formed in ATLAS to more fully examine ways to extract EFT coefficients from the wealth of top data to be extracted in Run 2 and beyond.