

# Prospects of Higgs portal DM models at future colliders

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Based on B. Dutta, T. Kamon, P. Ko, JL, arXiv: 1705.02149, 1712.05123

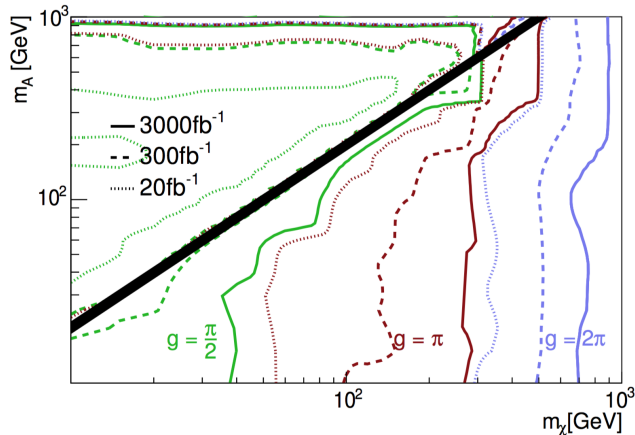
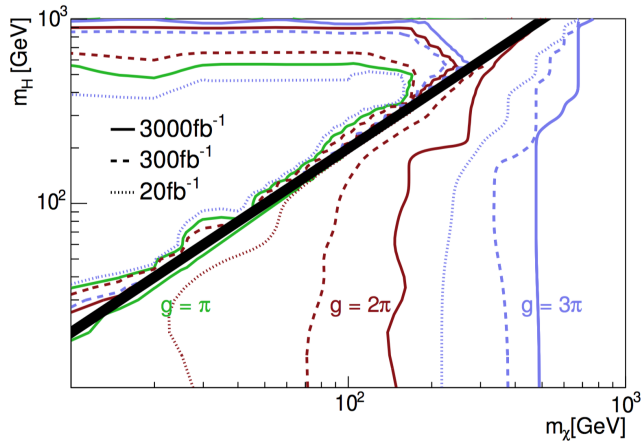
Apr. 5th, 2018

Dark Matter @ LHC 2018

Heidelberg University

- 1 Status of LHC searches for Higgs portal DM model
- 2 Prospects of searches at future colliders
  - 500 GeV ILC
  - 100 TeV FCC-pp/SppC
- 3 Conclusion

# Scalar mediated simplified model



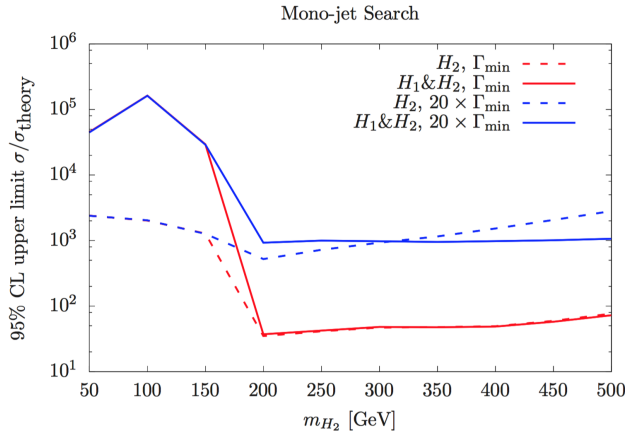
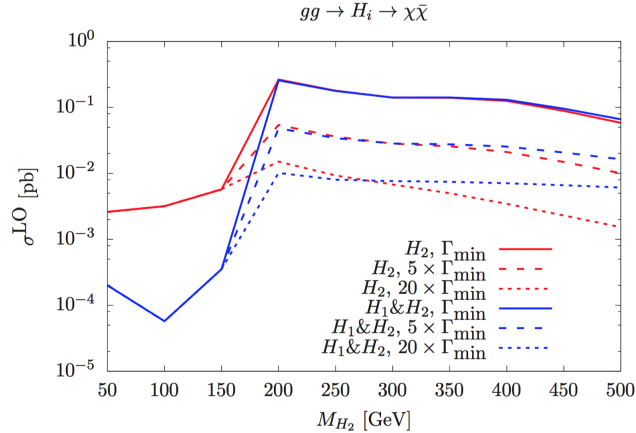
$$\mathcal{L}_H \ni -g_\chi H \bar{\chi} \chi - \sum_f \frac{g_v y_f}{\sqrt{2}} H \bar{f} f$$

$$\mathcal{L}_A \ni -ig_\chi A \bar{\chi} \gamma^5 \chi - \sum_f \frac{ig_v y_f}{\sqrt{2}} A \bar{f} \gamma^5 f$$

- Projected CMS invisible Higgs search in VBF
- Constraint on scalar is weaker than pseudoscalar, due to smaller cross section and a softer  $E_T^{\text{miss}}$ .
- Requires sizeable coupling for detection ( $g_v = g_\chi = g$ ), may not realistic, e.g.  $\Gamma > m_{H/A}$ .
- SM gauge symmetry is not conserved if  $H/A$  is an singlet.

$$\phi \cdot \bar{f} f = \phi \cdot (\bar{f}_L f_R + \bar{f}_R f_L)$$

# Higgs portal dark matter model



$$\begin{aligned}
 \mathcal{L}_{\text{FDM}} \ni & -g_\chi S \bar{\chi} \chi - \lambda_{HS} H^\dagger H S^2 - \mu_0^3 S \\
 & - \mu_1 S H^\dagger H - \frac{\mu_2}{3!} S^3 - \frac{\lambda_S}{4!} S^4 \\
 \mathcal{L}_{\text{FDM}}^{\text{int}} \ni & -(H_1 \cos \alpha + H_2 \sin \alpha) \sum_f \frac{m_f}{v_h} \bar{f} f \\
 & + g_\chi (H_1 \sin \alpha - H_2 \cos \alpha) \bar{\chi} \chi
 \end{aligned}$$

- Recast of the CMS invisible Higgs search in Mono-jet.
- $g_\chi = 3$ ,  $\sin \alpha = 0.3$ ,  $m_\chi = 80$  GeV (Throughout the talk.)
- Production limited by SM Higgs-fermion couplings.
- Interference between two mediators (reduce the sensitivity at low  $M_{H_2}$ ).
- All points way below the current LHC sensitivity.

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# Minimal gauge invariant models: FDM, VDM and SDM

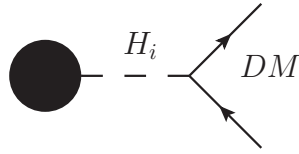
## Lagrangians:

- $\mathcal{L}_{\text{FDM}} = \mathcal{L}_{\text{SM}} + \bar{\chi}(i\not{\partial} - m_\chi - g_\chi S)\chi + \frac{1}{2}\partial_\mu S\partial^\mu S - \frac{1}{2}m_0^2 S^2 - \lambda_{HS}H^\dagger HS^2 - \mu_0^3 S - \mu_1 SH^\dagger H - \frac{\mu_2^2}{3!}S^3 - \frac{\lambda_S}{4!}S^4$
- $\mathcal{L}_{\text{VDM}} = -\frac{1}{4}V_{\mu\nu}V^{\mu\nu} + D_\mu\Phi^\dagger D^\mu\Phi - \lambda_\Phi(\Phi^\dagger\Phi - \frac{v_\phi^2}{2})^2 - \lambda_{H\Phi}(H^\dagger H - \frac{v_h^2}{2})(\Phi^\dagger\Phi - \frac{v_\phi^2}{2})$
- $\mathcal{L}_{\text{SDM}} = \frac{1}{2}\partial_\mu S\partial^\mu S - \frac{1}{2}m_0^2 S^2 - \lambda_{HS}H^\dagger HS^2 - \frac{\lambda_S}{4!}S^4$

## Interaction Lagrangians:

- $\mathcal{L}_{\text{FDM}}^{\text{int}} = -(H_1 \cos \alpha + H_2 \sin \alpha) \left[ \sum_f \frac{m_f}{v_h} \bar{f}f - \frac{2m_W^2}{v_h} W_\mu^+ W^{-\mu} - \frac{m_Z^2}{v_h} Z_\mu Z^\mu \right] + g_\chi(H_1 \sin \alpha - H_2 \cos \alpha) \bar{\chi}\chi$
- $\mathcal{L}_{\text{VDM}}^{\text{int}} = -(H_1 \cos \alpha + H_2 \sin \alpha) \left[ \sum_f \frac{m_f}{v_h} \bar{f}f - \frac{2m_W^2}{v_h} W_\mu^+ W^{-\mu} - \frac{m_Z^2}{v_h} Z_\mu Z^\mu \right] - \frac{1}{2}g_V m_V (H_1 \sin \alpha - H_2 \cos \alpha) V_\mu V^\mu$
- $\mathcal{L}_{\text{SDM}}^{\text{int}} = h \left[ \sum_f \frac{m_f}{v_h} \bar{f}f - \frac{2m_W^2}{v_h} W_\mu^+ W^{-\mu} - \frac{m_Z^2}{v_h} Z_\mu Z^\mu \right] + \lambda_{HS} v_h h S^2$

# General feature

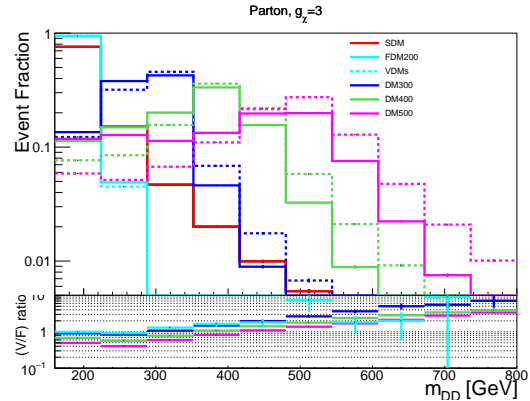
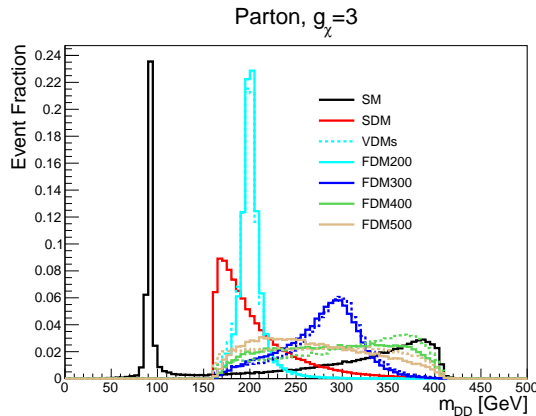


$$t \equiv m_{DD}^2$$

$$\frac{d\sigma_{SDM}}{dt} \propto \sigma_{SDM}^{h^*} \times \left| \frac{1}{t - m_h^2 + im_h\Gamma_h} \right|^2,$$

$$\frac{d\sigma_{FDM}}{dt} \propto \sigma_{FDM}^{h^*} \times \left| \frac{1}{t - m_{H_1}^2 + im_{H_1}\Gamma_{H_1}} - \frac{1}{t - m_{H_2}^2 + im_{H_2}\Gamma_{H_2}} \right|^2 \cdot (2t - 8m_\chi^2),$$

$$\frac{d\sigma_{VDM}}{dt} \propto \sigma_{VDM}^{h^*} \times \left| \frac{1}{t - m_{H_1}^2 + im_{H_1}\Gamma_{H_1}} - \frac{1}{t - m_{H_2}^2 + im_{H_2}\Gamma_{H_2}} \right|^2 \cdot \left( 2 + \frac{(t - 2m_D^2)^2}{4m_V^4} \right).$$



The dominant DM production process:

$$e^+e^- \rightarrow Z(\rightarrow ff) H_{1,2}^{(*)}(\rightarrow DD)$$

DM pair four-momentum:

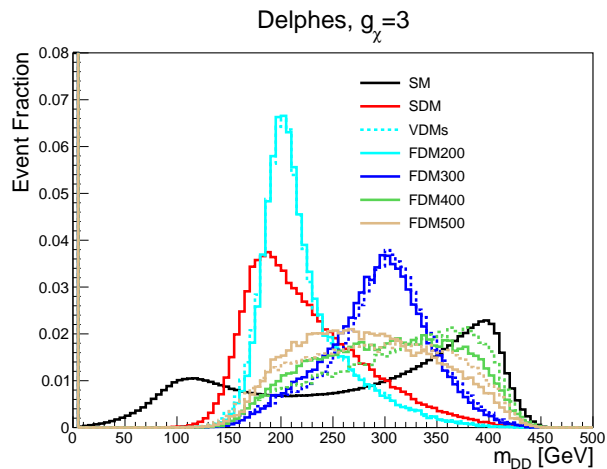
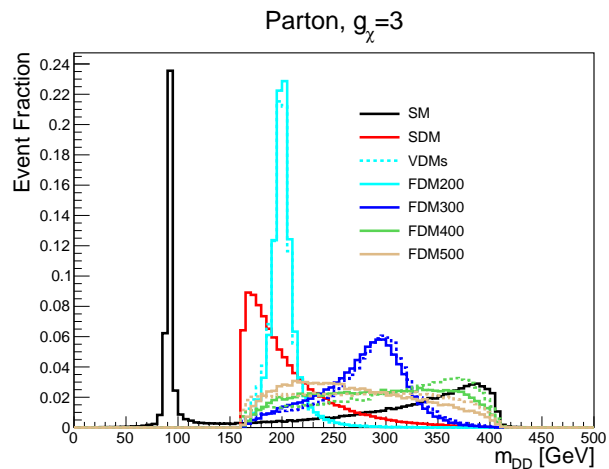
$$P_{DD}^\mu = P_{e^+}^\mu + P_{e^-}^\mu - P_Z^\mu = (\sqrt{s} - E_Z, -\vec{p}_Z)$$

DM pair invariant mass:

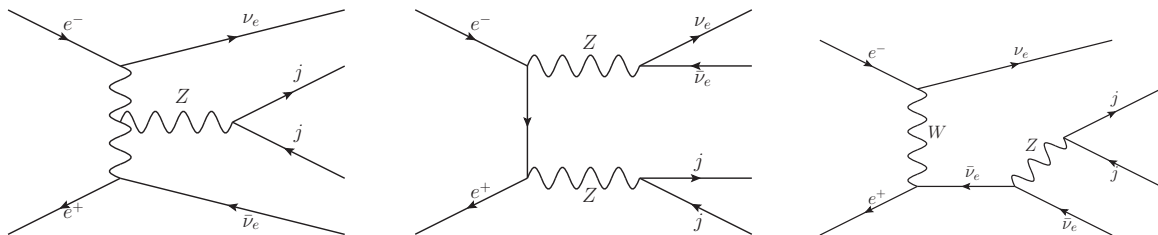
$$m_{DD}^2 = s + m_Z^2 - 2E_Z\sqrt{s}$$



# Features of DM production at the ILC



Dominant background processes:



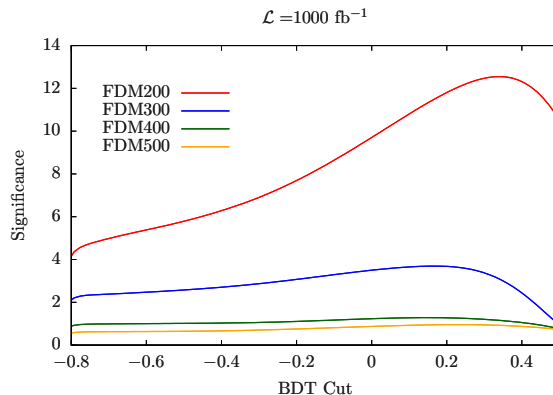
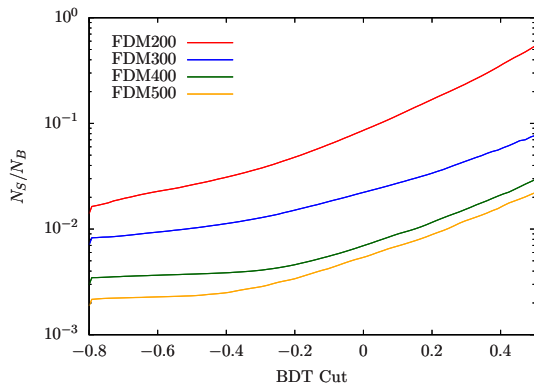
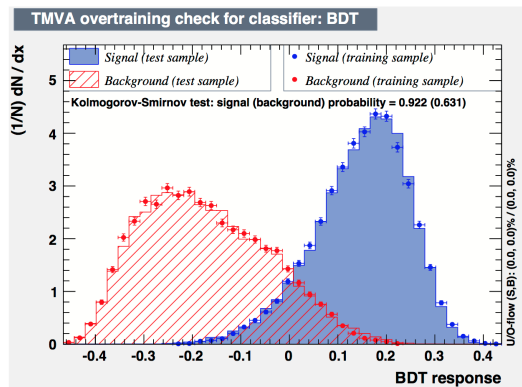
# Discovery prospects of the hadronic channel (FDM)

Preselection cuts:

- Lepton veto
- Exactly two jets
- $E_T^{\text{miss}} > 50 \text{ GeV}$

Boosted decision tree analysis with inputs:

$$m_{DD}, p_T(j_1), p_T(Z), E_T^{\text{miss}}, \Delta\phi^{\text{min}}, p_T(j_2), m_{jj}$$

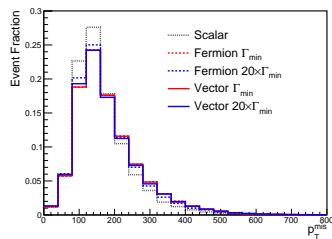
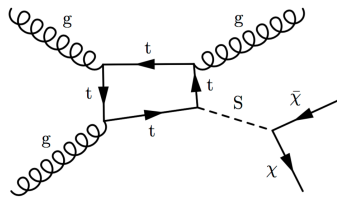


# Discovery prospects of the hadronic channel (FDM)

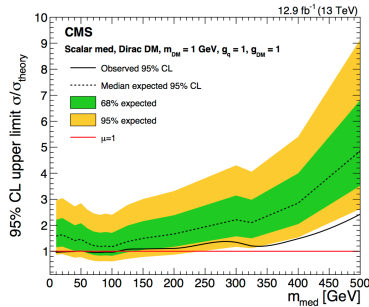
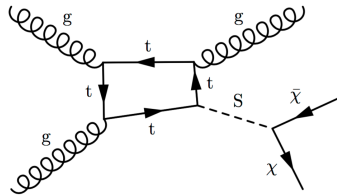
	<b>FDM200</b>	<b>FDM300</b>	FDM400	FDM500
$\sigma^0$ [fb]	1.643	0.9214	0.4221	0.2526
$\epsilon^{\text{pre}}$	0.796	0.717	0.655	0.698
BDT	0.3615	0.2132	0.1929	0.2129
$N_S/1000 \text{ fb}^{-1}$	697.8	410.5	148	102
$N_B/1000 \text{ fb}^{-1}$	2248.5	11453.5	12736	10898
$N_S/\sqrt{N_S + N_B}$	<b>12.85</b>	<b>3.769</b>	1.31	0.97

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# Signal: mono-jet?

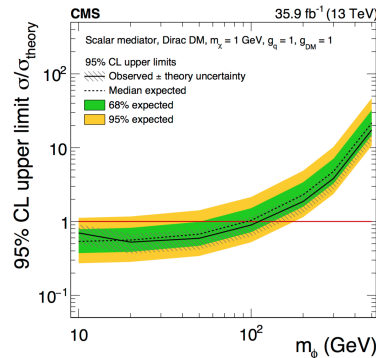
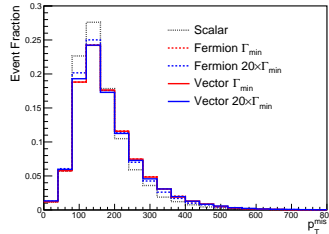


# Signal: mono-jet?



CMS mono-jet search,

[arXiv:1703.01651](https://arxiv.org/abs/1703.01651)



CMS  $tt+DM$  search in dileptonic

channel,

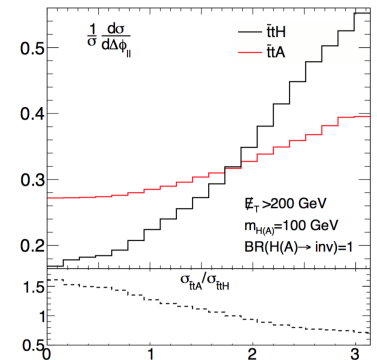
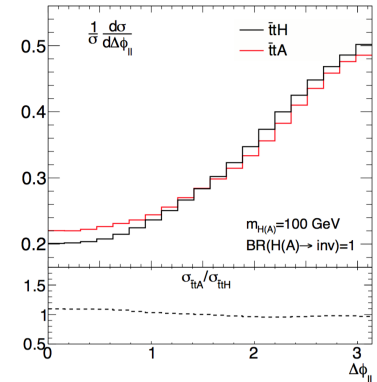
[arXiv:1711.00752](https://arxiv.org/abs/1711.00752)

*M. R Buckley et al.;*

*U. Haisch et al.;*

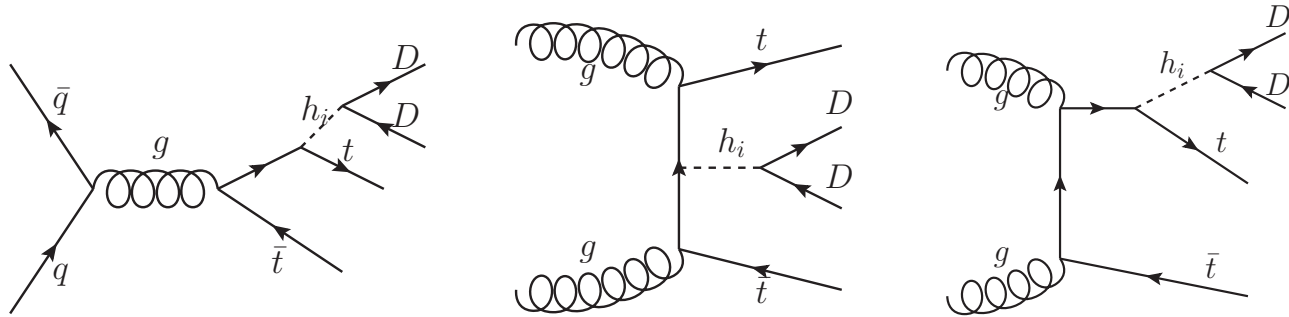
*F. Boudjema et al.;*

*J. Ellis et al. ...*



# Signal and background

The dominant DM production process:

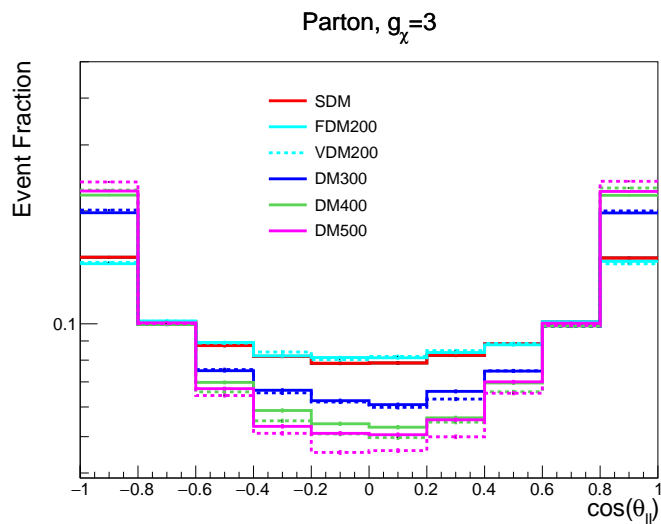
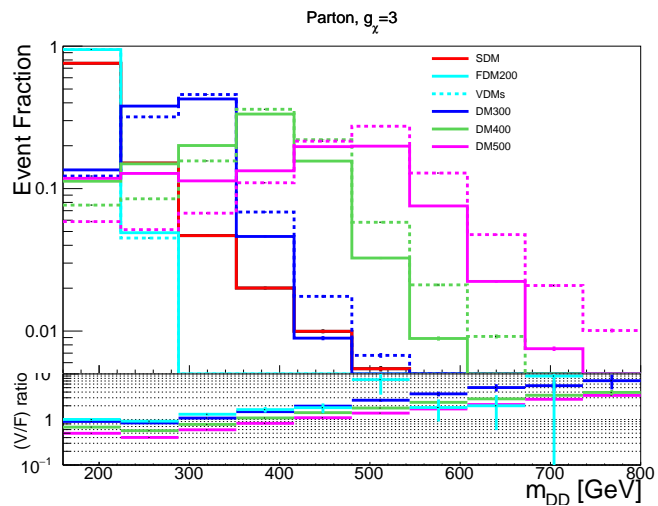


Dominant background processes:

	Cross section (NLO)
$\bar{t}t$	1316.5 pb
$\bar{t}tW$	20.5 pb
$\bar{t}tZ$	64.2 pb

	$t(\rightarrow bl\nu)t(\rightarrow bl\nu) + \text{DM}$
FDM200	34.2 fb
FDM300	18.7 fb
FDM400	14.8 fb
FDM500	12.5 fb

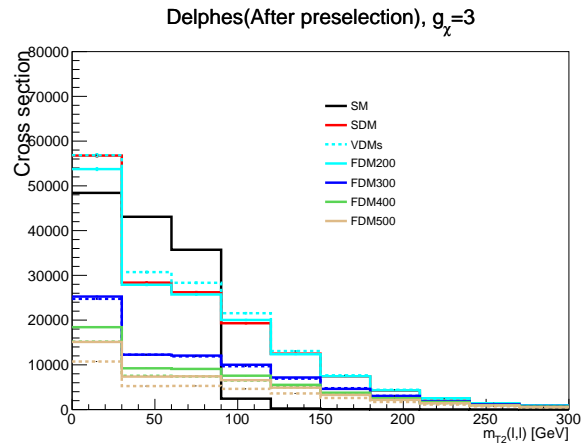
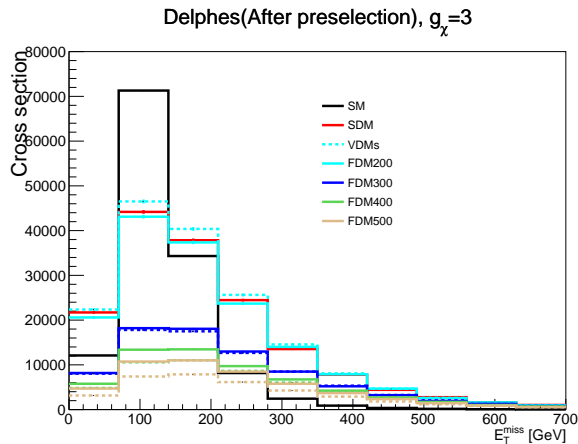
# Features of DM spin





# Analysis strategy

- Preselection: Exactly two opposite sign lepton and at least one b jet in the final state.
- $m_{\ell\ell} \notin [85, 95] \text{ GeV}$ .



- $E_T^{\text{miss}} > 150 \text{ GeV}$ .
- $m_{T_2}(l,l) > 150 \text{ GeV}$ .

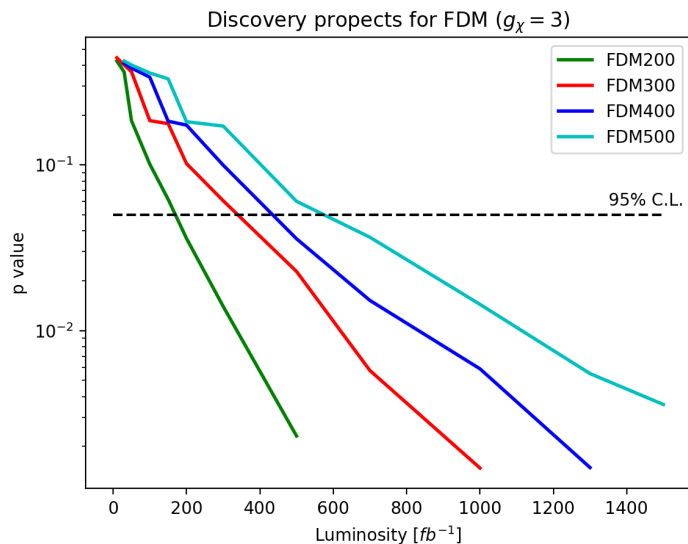
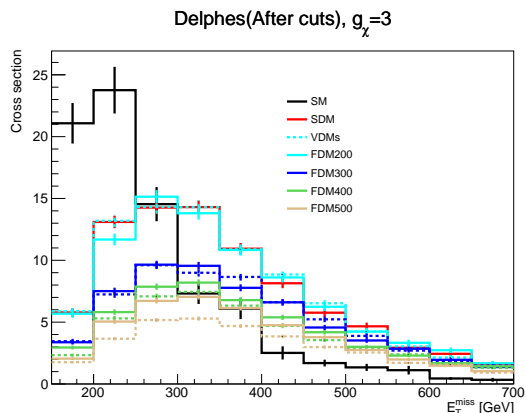
# Cuts flow for SM processes and Signals

Backgrounds	$\bar{t}t$	$\bar{t}tW$	$\bar{t}tZ$
Cross section	1316.5 pb	20.5 pb	64.2 pb
Preselections	63.76 pb	351.8 fb	1.9 pb
$m_{\ell\ell} \notin [85, 95] \text{ GeV}$	59.8 pb	330.4 fb	1.05 pb
$E_T^{miss} > 150 \text{ GeV}$	17.76 pb	69.61 fb	261.14 fb
$m_{T_2}(l, l) > 150 \text{ GeV}$	23.83 fb	1.92 fb	32.1 fb

Signals	FDM200	FDM300	FDM400	FDM500
Cross section	34.2 fb	18.7 fb	14.8 fb	12.5 fb
Preselections	7.86 fb	3.99 fb	3.05 fb	2.55 fb
$m_{\ell\ell} \notin [85, 95] \text{ GeV}$	7.47 fb	3.82 fb	2.92 fb	2.44 fb
$E_T^{miss} > 150 \text{ GeV}$	4.17 fb	2.44 fb	1.93 fb	1.63 fb
$m_{T_2}(l, l) > 150 \text{ GeV}$	0.87 fb	0.62 fb	0.54 fb	0.47 fb
$\mathcal{L}^{95\%} [\text{fb}^{-1}]$	305	602	793	1047

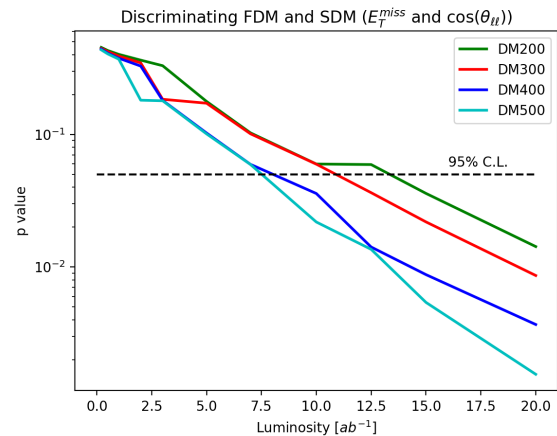
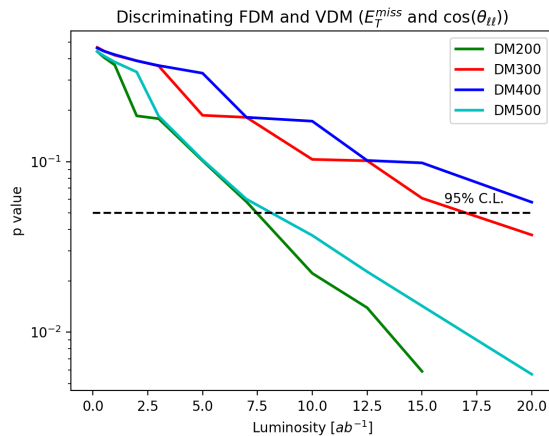
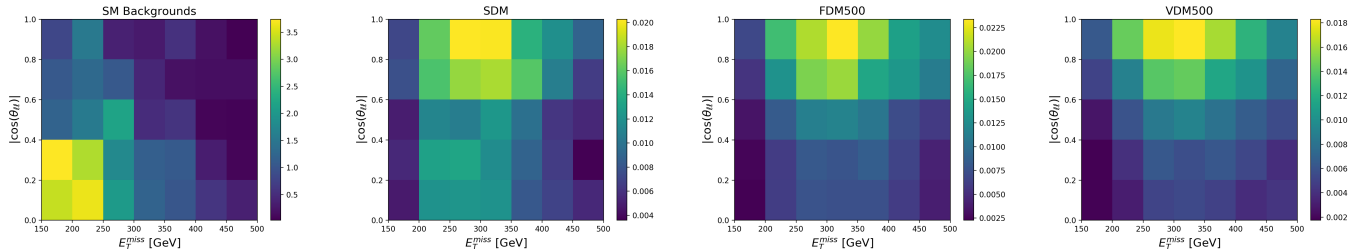
Binned log-likelihood analysis:

$$\mathcal{L}(\text{data}|\mathcal{H}_\alpha) = \prod_i \frac{t_i^{n_i} e^{-t_i}}{n_i!}, \quad \mathcal{Q} = -2 \log \left( \frac{\mathcal{L}(\text{data}|\mathcal{H}_\alpha)}{\mathcal{L}(\text{data}|\mathcal{H}_0)} \right).$$



# Spin characterization ( $E_T^{miss}$ and $\cos(\theta_{\ell\ell})$ )

Two dimensional binned log-likelihood test:  $\mathcal{L}(\text{data}|\mathcal{H}_\alpha) = \prod_{i,j} \frac{n_{ij} e^{-t_{ij}}}{n_{ij}!}$



# Conclusion

- The gauge invariant Higgs portal DM models for FDM and VDM require at least two mediators, while that of SDM only need one.
- At the ILC,  $m_{H_2} \lesssim 300$  GeV can be probed at more than  $3\text{-}\sigma$  level.
- At the 100 TeV  $p\text{-}p$  collider, (1) All benchmark points should be probable at integrated luminosity of  $\mathcal{O}(100)$  fb $^{-1}$  at 100 TeV p-p collider; (2) The spin discriminations for our benchmark points are possible at  $\mathcal{O}(10)$  ab $^{-1}$ . (3) Those values are all below the targets luminosity of FCC-hh, which is  $\sim 20$  ab $^{-1}$ .

Backup.

# Benchmark points

- The relevant parameters in FDM for collider search:  
 $g_\chi = 3$ ,  $\sin \alpha = 0.3$ ,  $m_\chi = 80$  GeV and  $m_{H_2} = (200, 300, 400, 500)$  GeV.

$$\begin{aligned} \Gamma_{\min}^{\text{FDM}}(H_2) &= \Gamma(H_2 \rightarrow \chi\chi) + \Gamma(H_2 \rightarrow WW/ZZ) + \Gamma(H_2 \rightarrow ff) \\ &= \cos^2 \alpha \cdot g_\chi^2 \frac{m_{H_2}}{8\pi} \left(1 - \frac{4m_\chi^2}{m_{H_2}^2}\right)^{3/2} + \sin^2 \alpha \cdot \frac{G_\mu m_{H_2}^3}{16\sqrt{2}\pi} \delta_V \sqrt{1 - 4\frac{m_V^2}{m_{H_2}^2}} \left(1 - 4\frac{m_V^2}{m_{H_2}^2} + 12\frac{m_V^4}{m_{H_2}^4}\right) \\ &\quad + \sin^2 \alpha \cdot \left(\frac{m_f}{v}\right)^2 \frac{3m_{H_2}}{8\pi} \left(1 - \frac{4m_f^2}{m_{H_2}^2}\right)^{3/2}, \end{aligned}$$

where  $f$  is the SM fermion,  $V = Z, W$  and  $\delta_V = 1(2)$  for  $Z(W^\pm)$ .

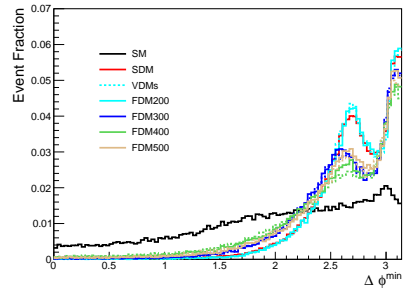
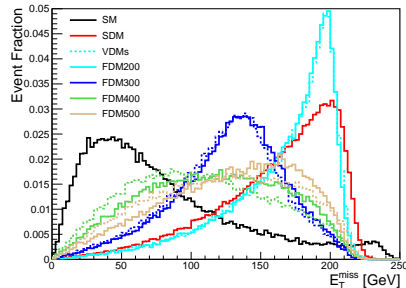
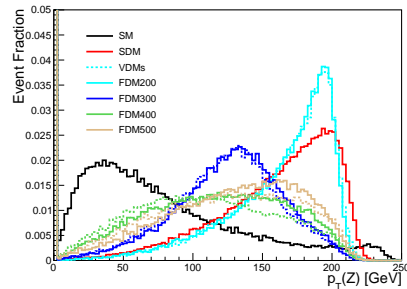
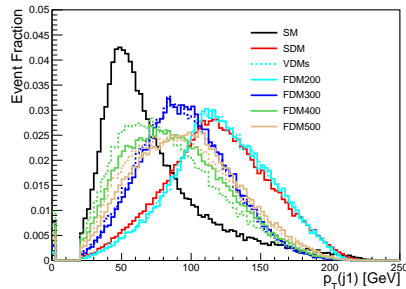
- Parameters for the vector DM production are chosen accordingly:  
 $\sin \alpha = 0.3$ ,  $m_V = 80$  GeV and  $g_V$  is chosen such that the total decay width of  $H_2$  is the same as benchmark points of FDM.

$m_{H_2}$ [GeV]	200	300	400	500
$\Gamma_{\min}(H_2)$ [GeV]	14.2	60.1	103.0	144.5
$g_V$	3.53	3.07	2.37	1.91

- Fix  $m_S = 80$  GeV and take appropriate  $\lambda_{HS}$  such that the production cross section of the signal process is the same with that in the FDM.

# Discovery prospects of the hadronic channel

Kinematic distributions:



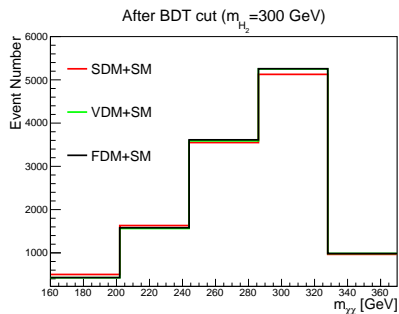
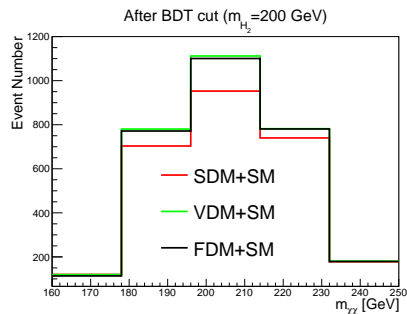
$$\Delta\phi^{\min} = \min_{i=1,2} \Delta\phi(p_T^{\text{miss}}, p(j_i))$$



# Spin characterization

The same preselection and BDT cuts as used for FDM the benchmark point FDM200 (FDM300) are applied to the corresponding benchmark point SDM200 (SDM300) and VDM200 (VDM300).

	SDM200	SDM300	VDM200	VDM300
$\sigma^0$ [fb]	1.643	0.9214	1.734	0.8674
$\epsilon^{\text{pre}}$	0.7875	0.7875	0.801	0.711
$N_S/1000 \text{ fb}^{-1}$	447	322.3	726	363.5
$\mathcal{S}$	4.4	3.3	0.59	0.44



$$\text{SDM: } \delta\chi^2 = \sum_{i=1}^5 \left( \frac{N_i^{\text{FDM+SM}} - N_i^{\text{SDM+SM}}}{\sqrt{N_i^{\text{FDM+SM}}}} \right)^2$$

$$\text{VDM: } \mathcal{S} = |N_S^{\text{FDM}} - N_S^{\text{VDM}}| / \sqrt{N_B}$$

# Spin characterization

Using only the distribution of  $E_T^{miss}$ :

$\mathcal{H}_0$  is the FDM + SM,  $\mathcal{H}_\alpha$  can be VDM/SDM + SM

