Dark Matter Searches in the Z(ll) +MET channel

Christopher Anelli
on Behalf of the ATLAS Experiment

University of Victoria

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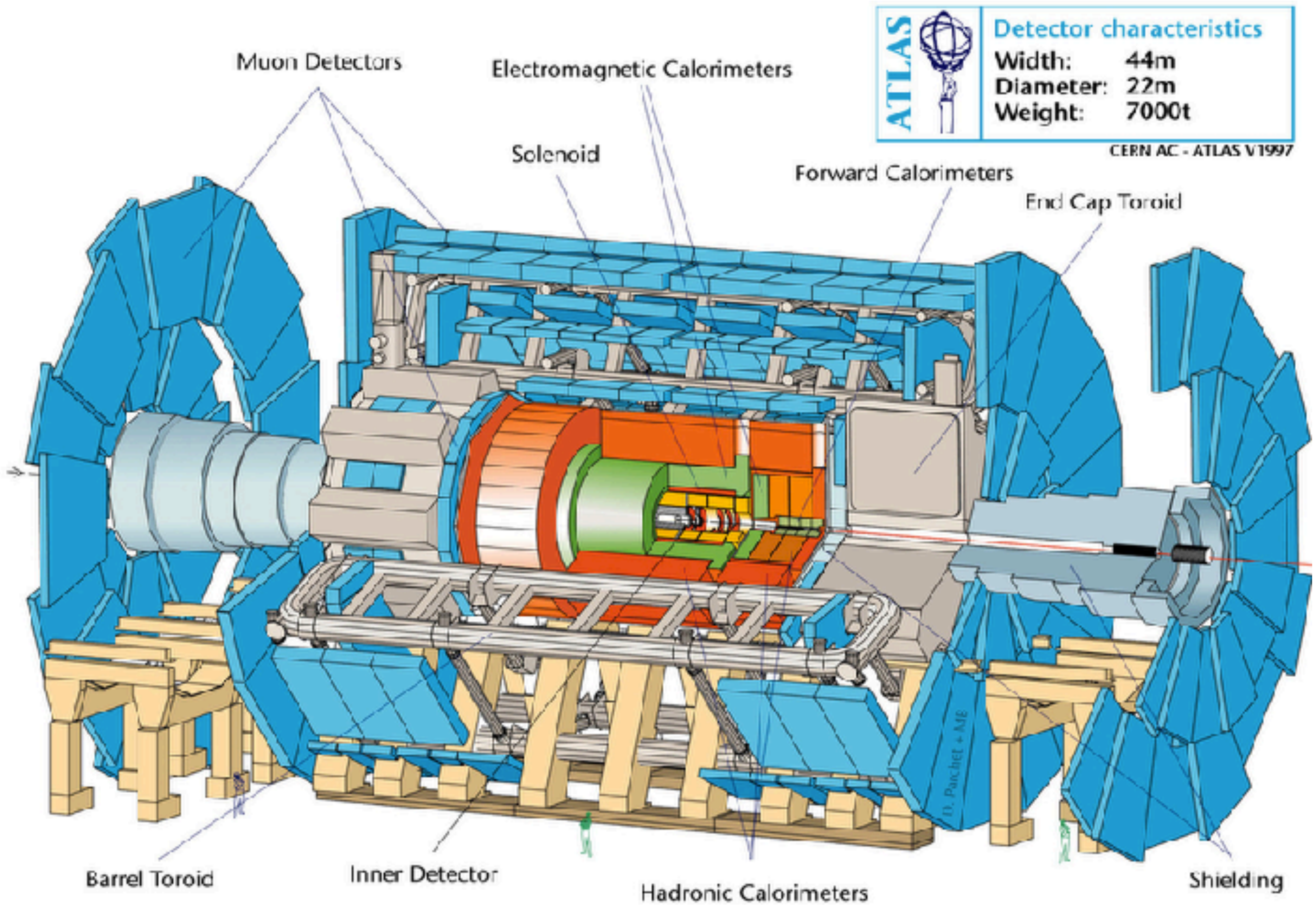
Dark Matter @ LHC 2018

ATLAS Experiment



- Analyze **13 TeV** proton-proton collisions collected by the ATLAS experiment during 2015 and 2016 of LHC Run-II.
- Collected events correspond to an integrated luminosity of **36.1 fb^{-1}**.

Z(ll) + MET

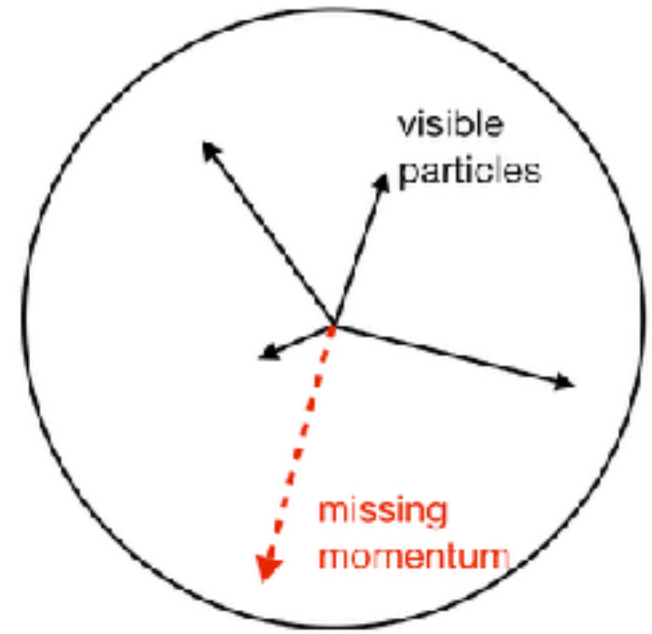
Search for DM through the presence of missing transverse momentum in the detector:

$E_T^{miss} = -\Sigma\ \vec{p}_T$ All Reconstructed Objects

†ATLAS definition also includes soft term for tracks not associated with a particle.

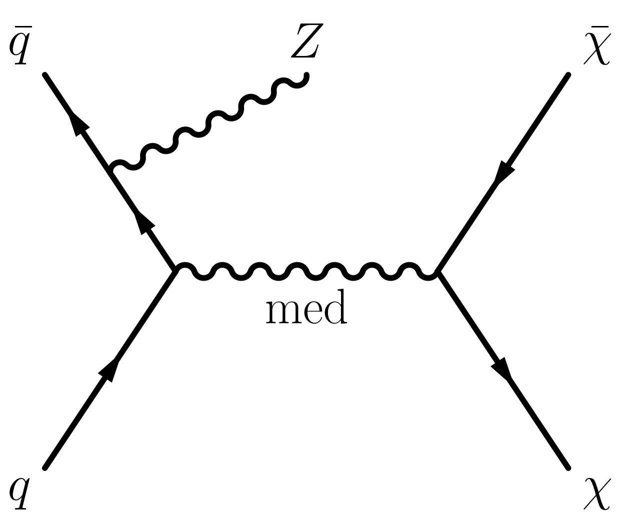
Missing Transverse Momentum Performance Paper (ATLAS 2018).

Neutrinos also produce E_T^{miss} and are a main source of background.

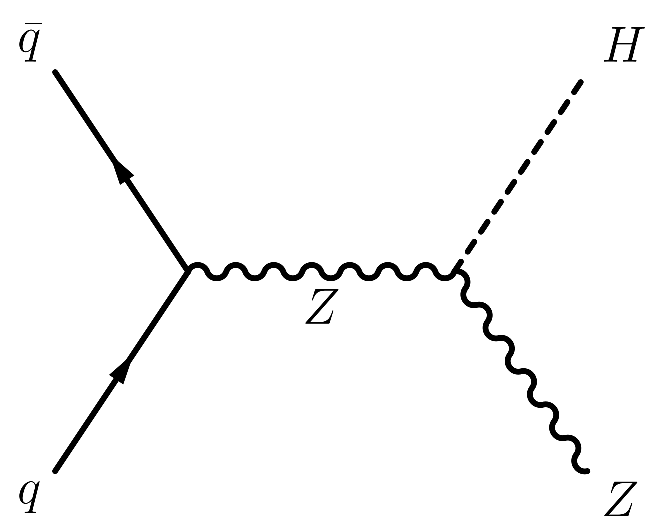


Sensitive to DM models where p_T imbalance is created from DM production recoiled against a Z boson.

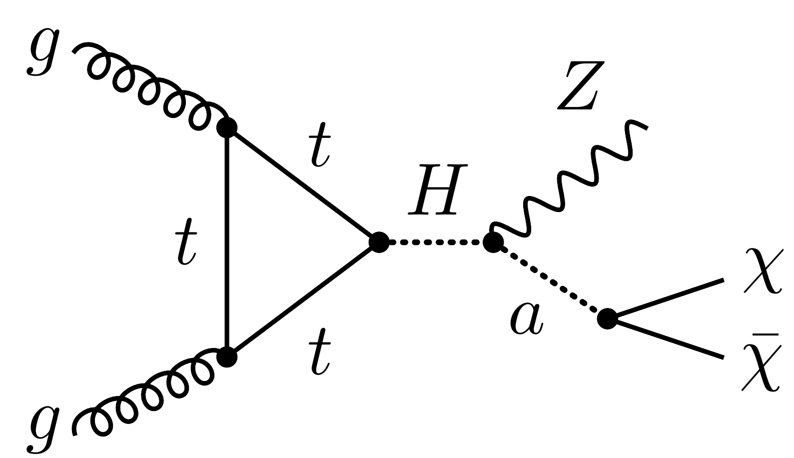
Simplified Models



Higgs Portal



Pseudoscalar 2HDM



Event Selection

Triggers:

- Low pT Triggers for isolated electrons (muons). Thresholds at 24 (20)† GeV.
- And high pT Triggers without an isolation requirement, 50 (60) GeV.

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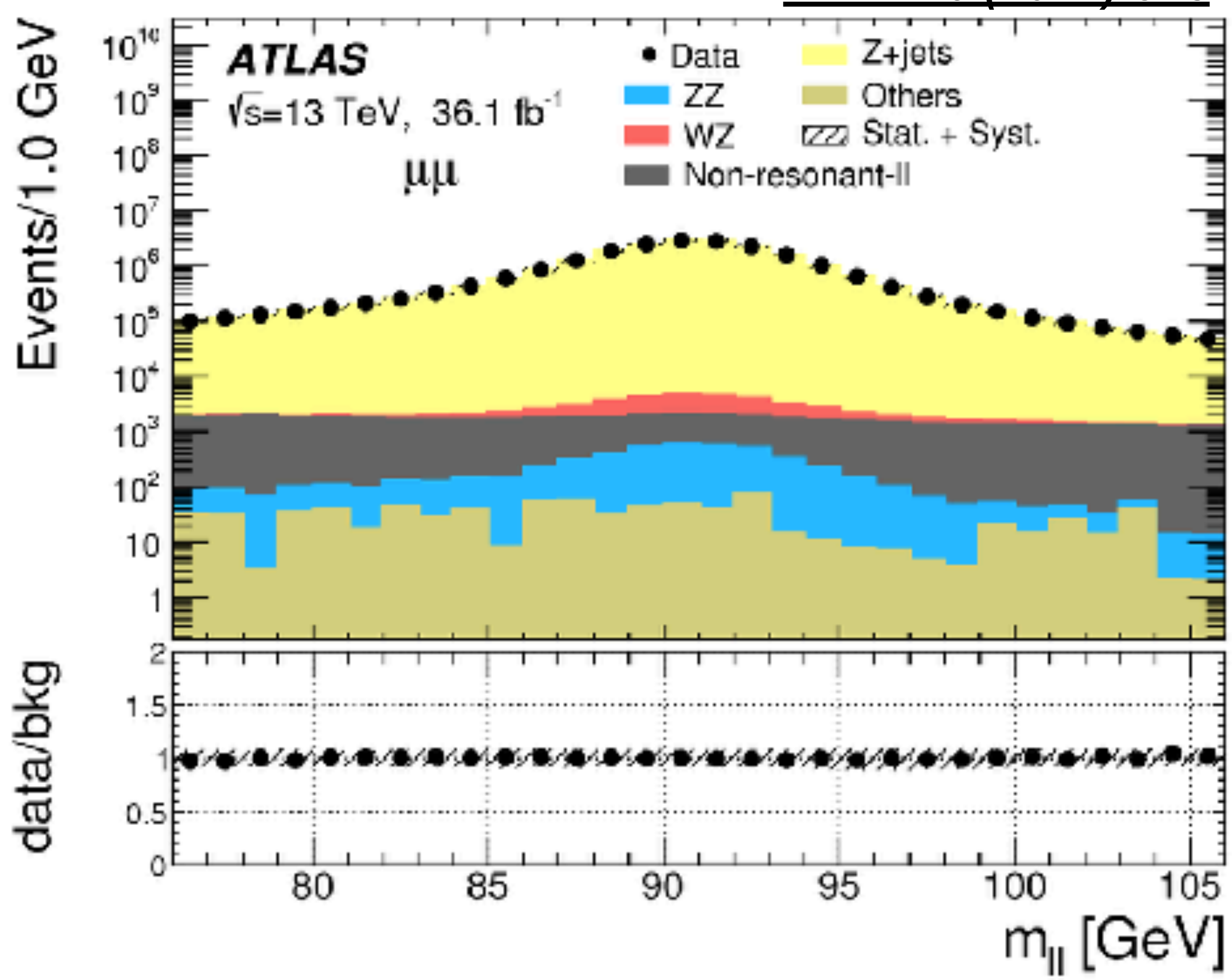

Distribution of data and background estimates following Z-window cut.

Preselection

- Require <u>exactly</u> 2 SF OS leptons
- Lead (sublead) lepton $p_T > 30$ (20) GeV
- Z window cut $| m_{\ell\ell} - m_Z | < 15$ GeV

Highlighted Selection Cuts

- $E_T^{miss} > 90$ GeV and $E_T^{miss} / H_T > 0.6$
- $\Delta\varphi(p_T^{\ell\ell}, E_T^{miss}) > 2.7$
- $\Delta R_{\ell\ell} < 1.8$
- $|p_T^{\ell\ell} - p_T^{miss,jets})| / p_T^{\ell\ell} < 0.2$

†Depending on luminosity the electron trigger threshold is increased to 26 GeV.

Backgrounds

Background Estimates for Signal Region:

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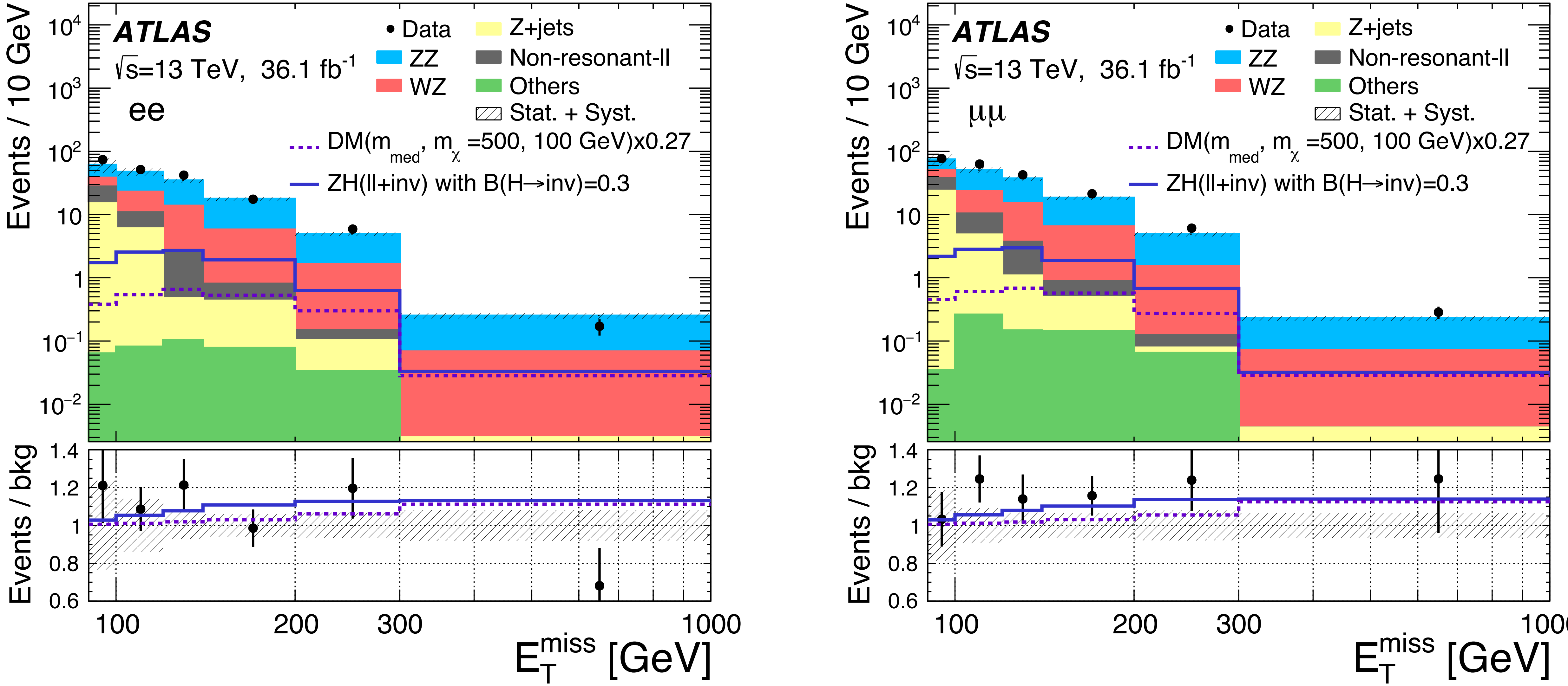

- Main background is ZZ production, with one Z→νν. Estimated with Monte Carlo.
- WZ background, (W→ℓν), where the lepton is lost or, in tau case hadronically decays. Shape modeled with MC, but normalized to data using WZ enriched control region.
- Non-resonant→ℓℓ background estimated with opposite sign eμ control region.
- Z+jets background, mis-reconstructed jets fake leptons, Z+jets background. Estimated using data driven ABCD method (E_T^{miss} , E_T^{miss}/H_T).

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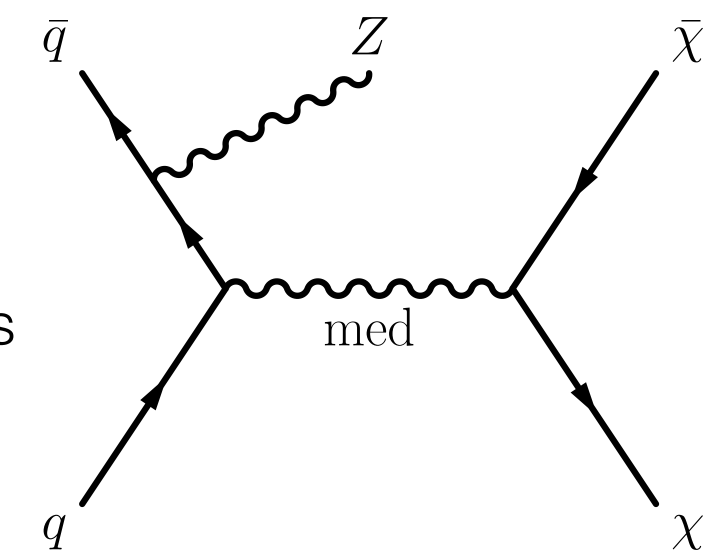
DM Simplified Models

DM simplified models are distinguished by mediator type and coupling strengths:

Mediator Types - Vector, Axial-vector, Scalar, Pseudoscalar.

Simplified models have 6 free parameters:

- g_q - mediator coupling to quarks
- g_ℓ - mediator coupling to leptons
- g_χ - mediator coupling to DM
- M_{DM} - DM mass
- M_{med} - Mediator mass
- Γ_{med} - Mediator width



Following the recommendations of the DM Working Group, ATLAS studies two benchmark models:

- **A1** - Axial-vector model with g_q=0.25, g_ℓ=0, g_χ =1
- **V1** - Vector model with g_q=0.25, g_ℓ=0, g_χ =1

- Γ_{med} is set using the minimal width formula.
- Results are shown as 2D exclusion plots in M_{med} : M_{DM}.

Axial-vector Model (A1)

95% confidence-level exclusion limits for the axial-vector model (A1):

A1 - Axial-vector model with g_q=0.25, g_ℓ=0, g_χ =1

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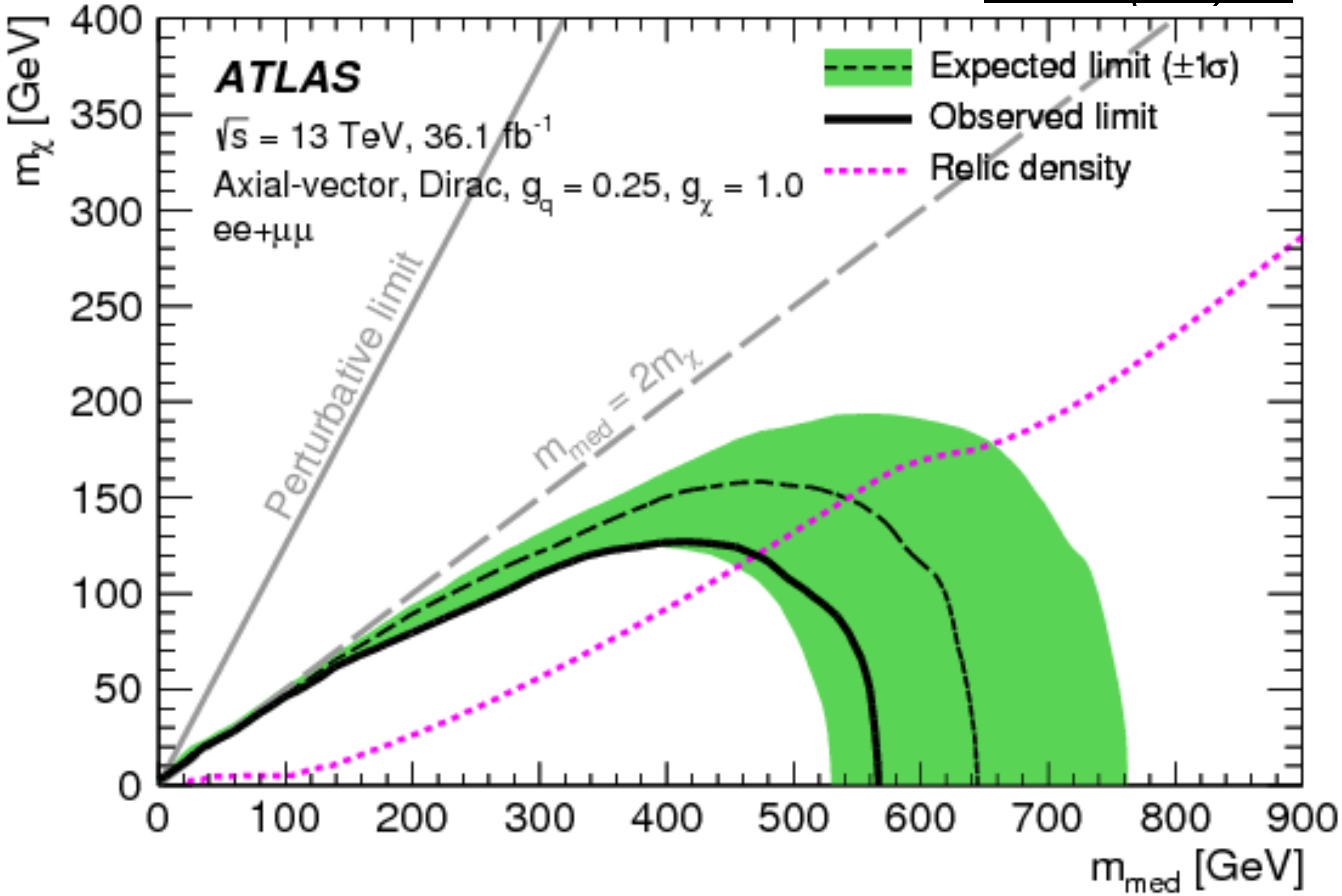

Vector Model (V1)

95% confidence-level exclusion limits for the vector model (V1):

V1 - Vector model with g_q=0.25, g_ℓ=0, g_χ =1

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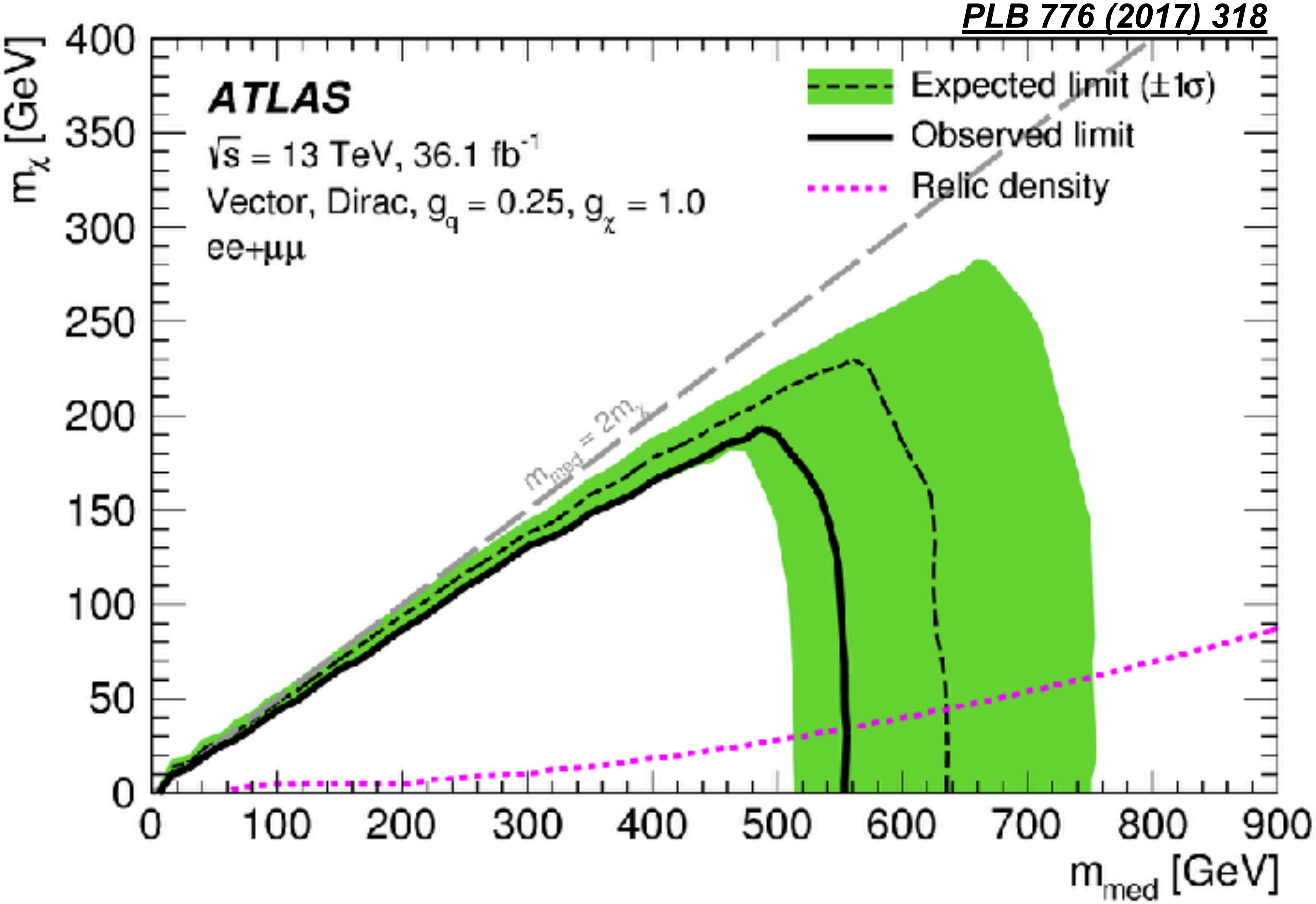

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Comparison to Direct Detection

Simplified models allow for comparison between direct detection experiments and colliders.

Axial-vector model:

$$\sigma^0_{\mathrm{SD}} \approx 4.6 \times 10^{-41}\ \mathrm{cm}^2 \cdot \left(\frac{g_{\mathrm{DM}}\, g_q}{1}\right)^2 \left(\frac{1\ \mathrm{TeV}}{M_{\mathrm{med}}}\right)^4 \left(\frac{\mu_{n\chi}}{1\ \mathrm{GeV}}\right)^2$$

Observed limits for Mono-Z(ll) search

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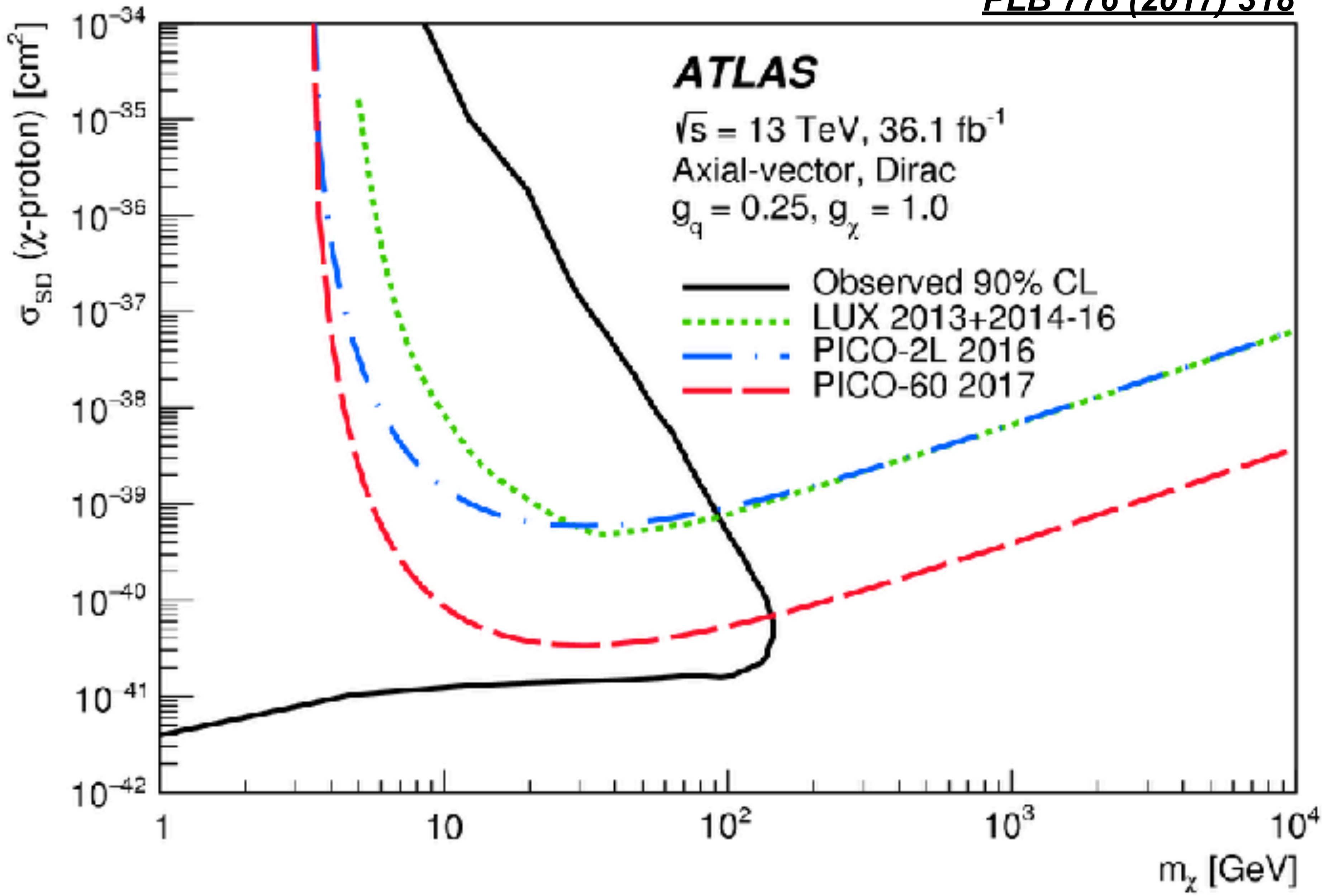


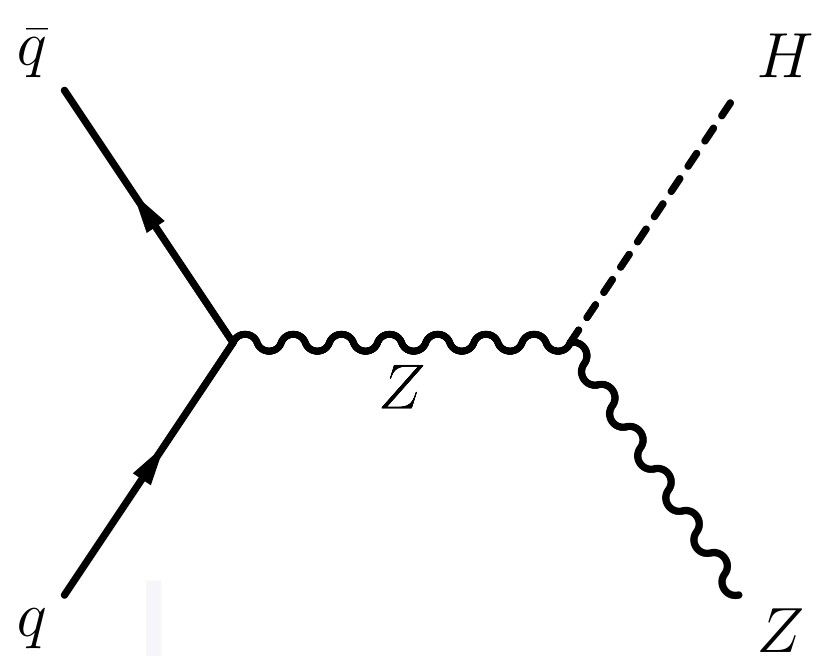
N.B. Collider limits are model dependent, and different couplings result in different limits.

Invisible Higgs Decay

The Z(ll) + MET channel is sensitive to invisible decays of the Higgs boson. Assuming SM ZH production, ucan measure the branching fraction, B(H→inv).

The SM H→ZZ→νννν branching fraction is 1.06 $\cdot 10^{-3}$, and observing a larger branching fraction is clear evidence of Beyond the SM physics.

In particular, larger branching fractions are consistent with Higgs portal dark matter or Higgs decays to light neutralinos. (To be sensitive, must have $m_\chi < 0.5\, m_H$)



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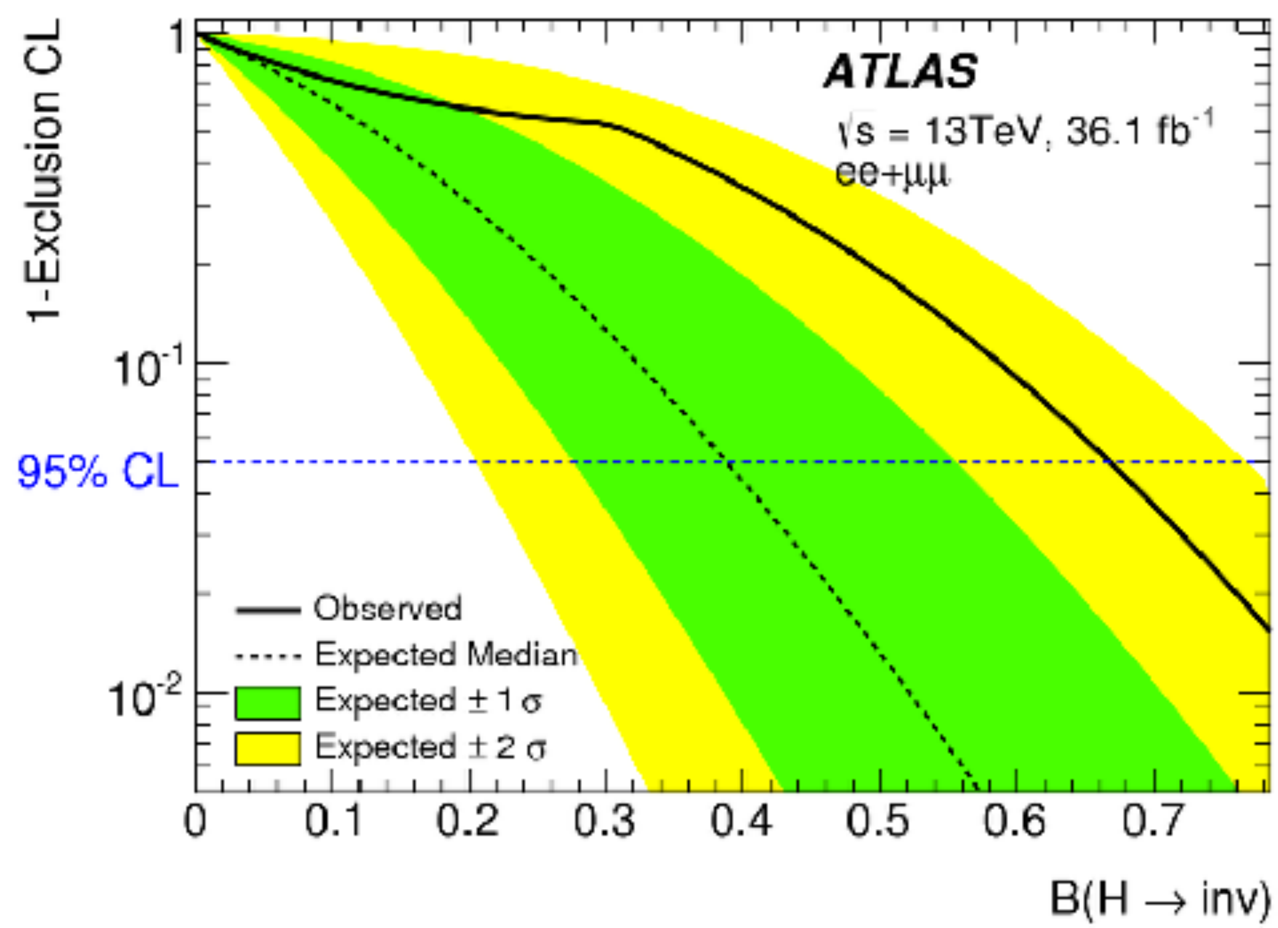
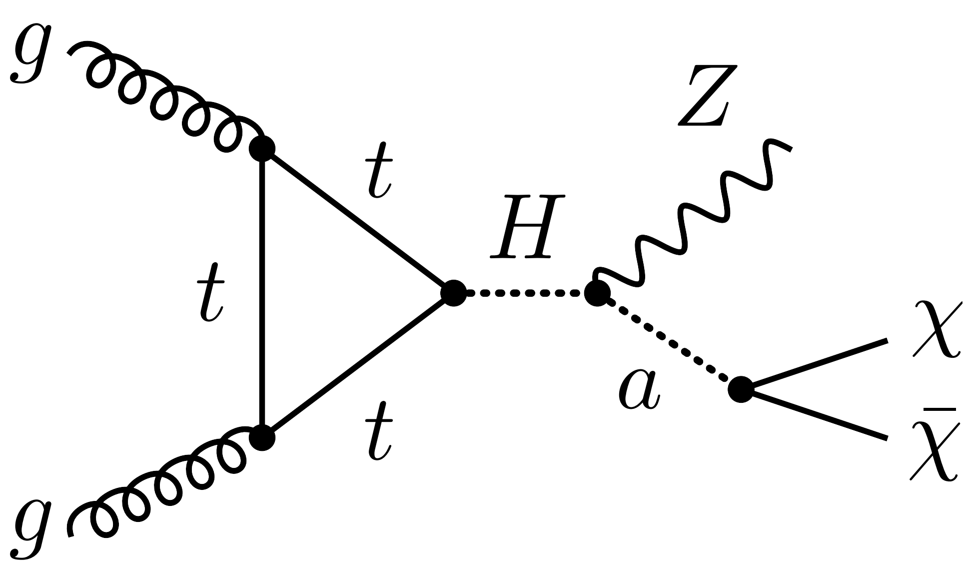
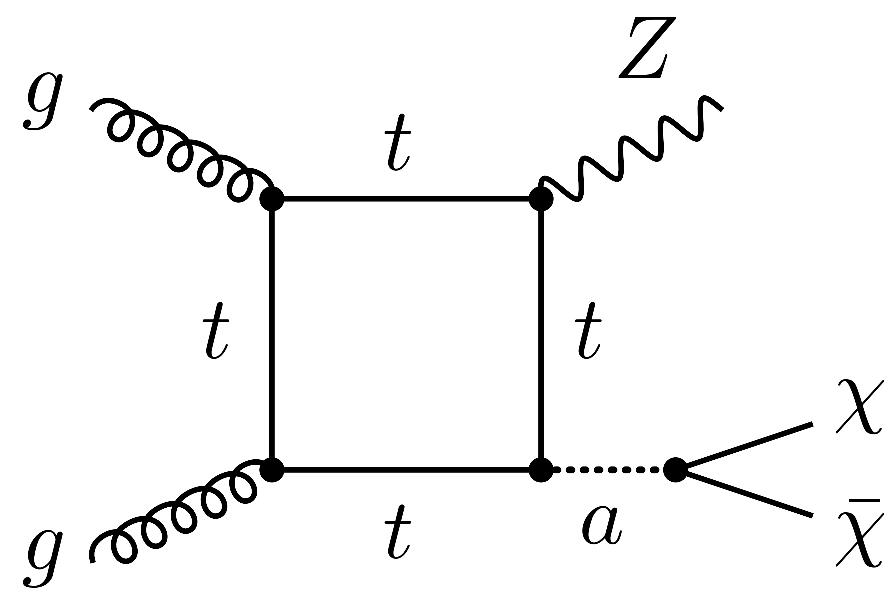


ATLAS EXPERIMENT

Pseudoscalar Two Higgs Doublet Model

Feynman Diagrams for Mono-Z + MET production.





2HDMa model introduces 6 bosons:

h - light scalar, identified as SM Higgs

H - heavy scalar

$H^{\pm}$ - two heavy, charged scalars

A - heavy pseudoscalar

a - light pseudo scalar

} Mixture of 2HDM (A_0) and DM Mediator (a_0) pseudoscalars

Couples to DM and SM particles

Parameters:

mH, $mH^{\pm}$, mA, ma, m_{χ} - masses

sin(θ) - a, A mixing angle

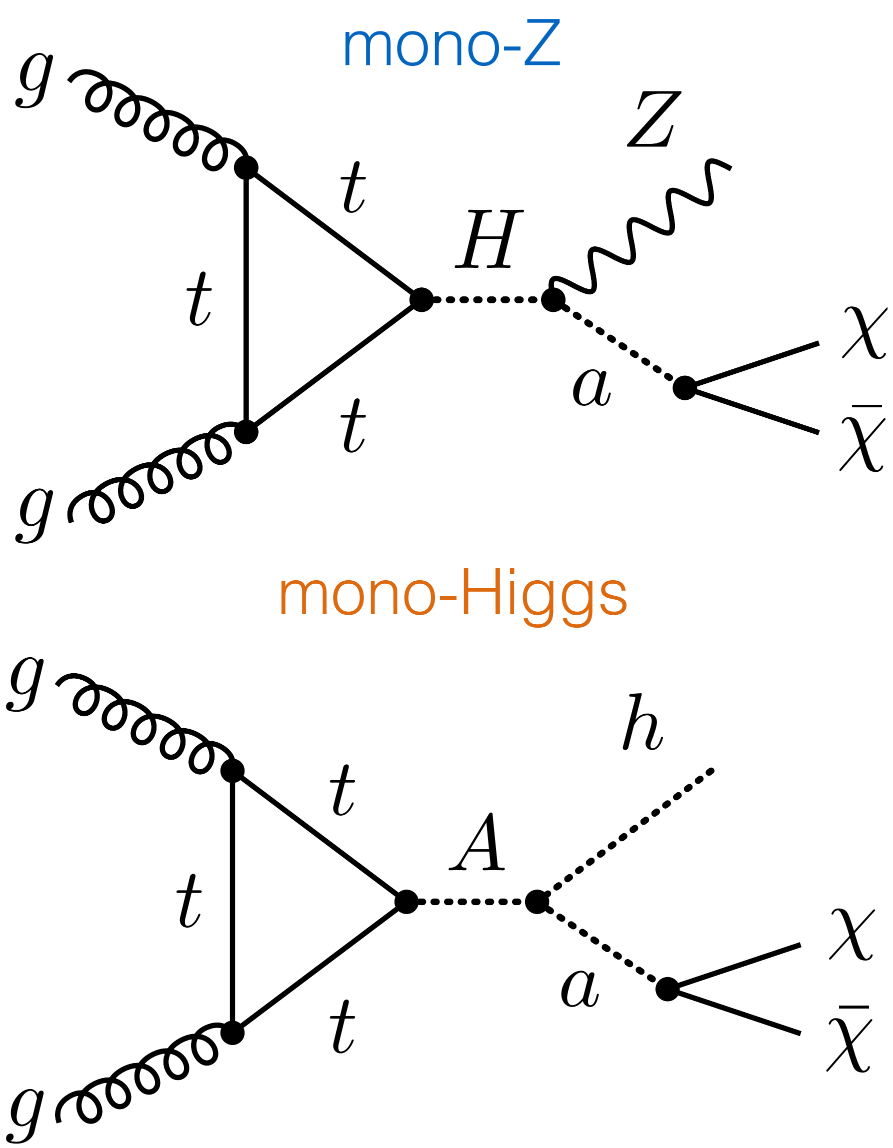
tan(β) - ratio of vacuum expectation values

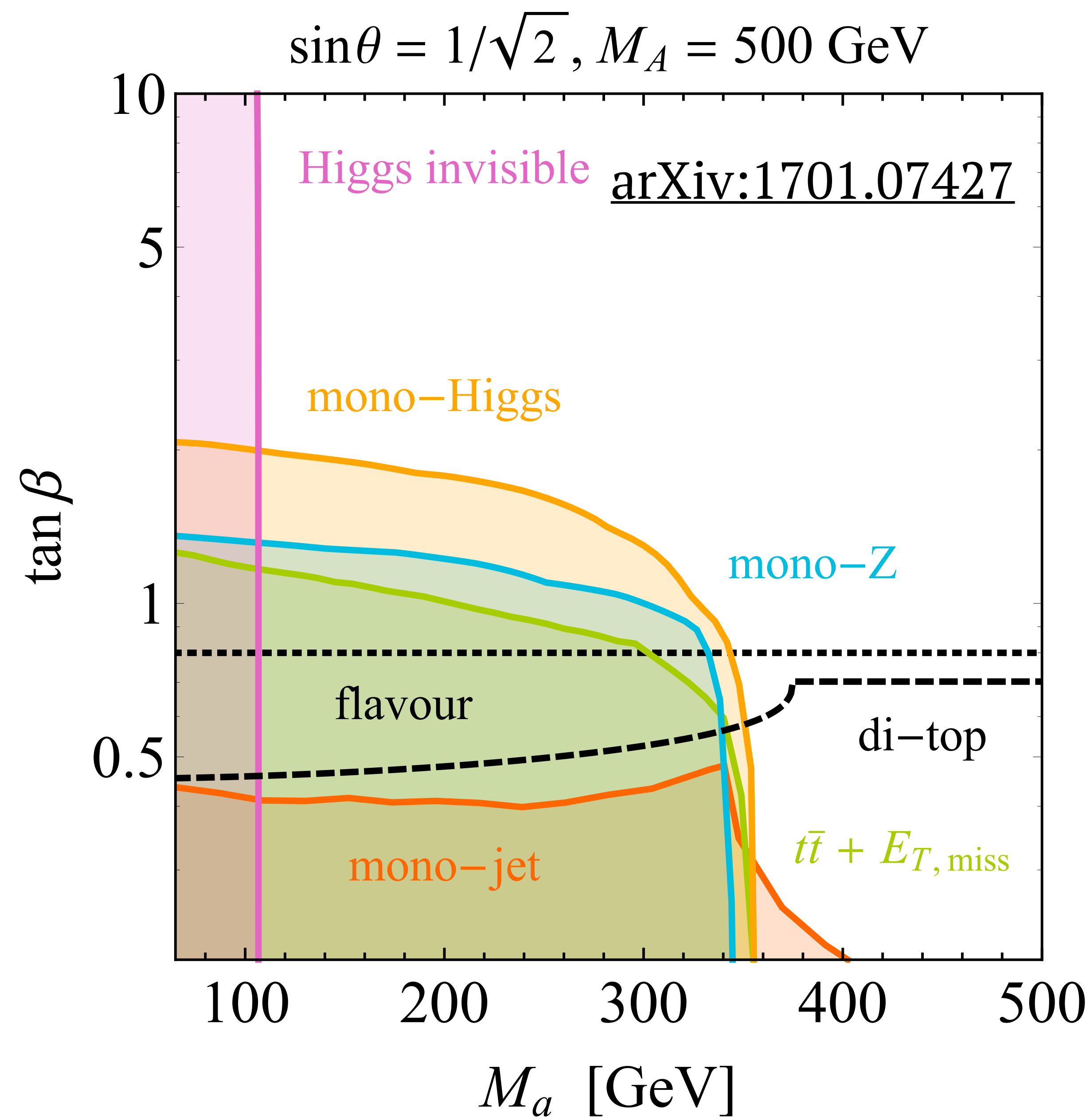
λ_3, λ_{1p} , λ_{2p} - quartic scalar couplings

y_{χ} - DM coupling

Pseudoscalar Two Higgs Doublet Model

For 2HDMa, Mono-Z and Mono-Higgs channels are enhanced due to resonant production of heavy scalar (H) or heavy pseudoscalar (A) particles, can be more sensitive than Mono-Jet.



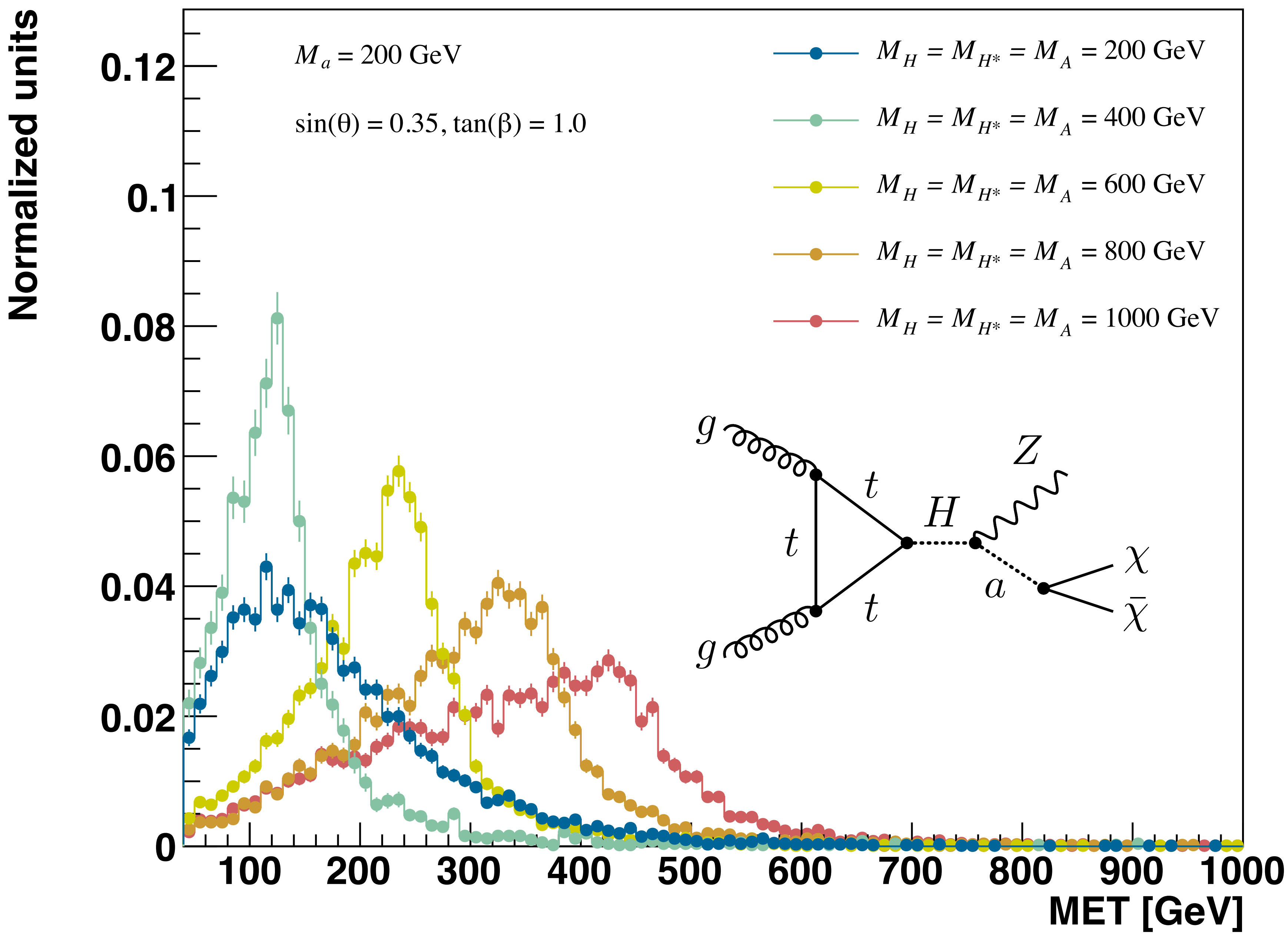


Study a subset of 2HDMa models with **mA =mH = mH±** (Mono-Z and Mono-Higgs channels complement each other.)

Kinematic Dependence on Masses

For resonant production MET distribution is characterized by a Jacobian peak. Shape depends strongly on M_H and M_a.

Fixed parameters: sin(θ)=0.35, tan(β)=1.0, m_χ=10 GeV

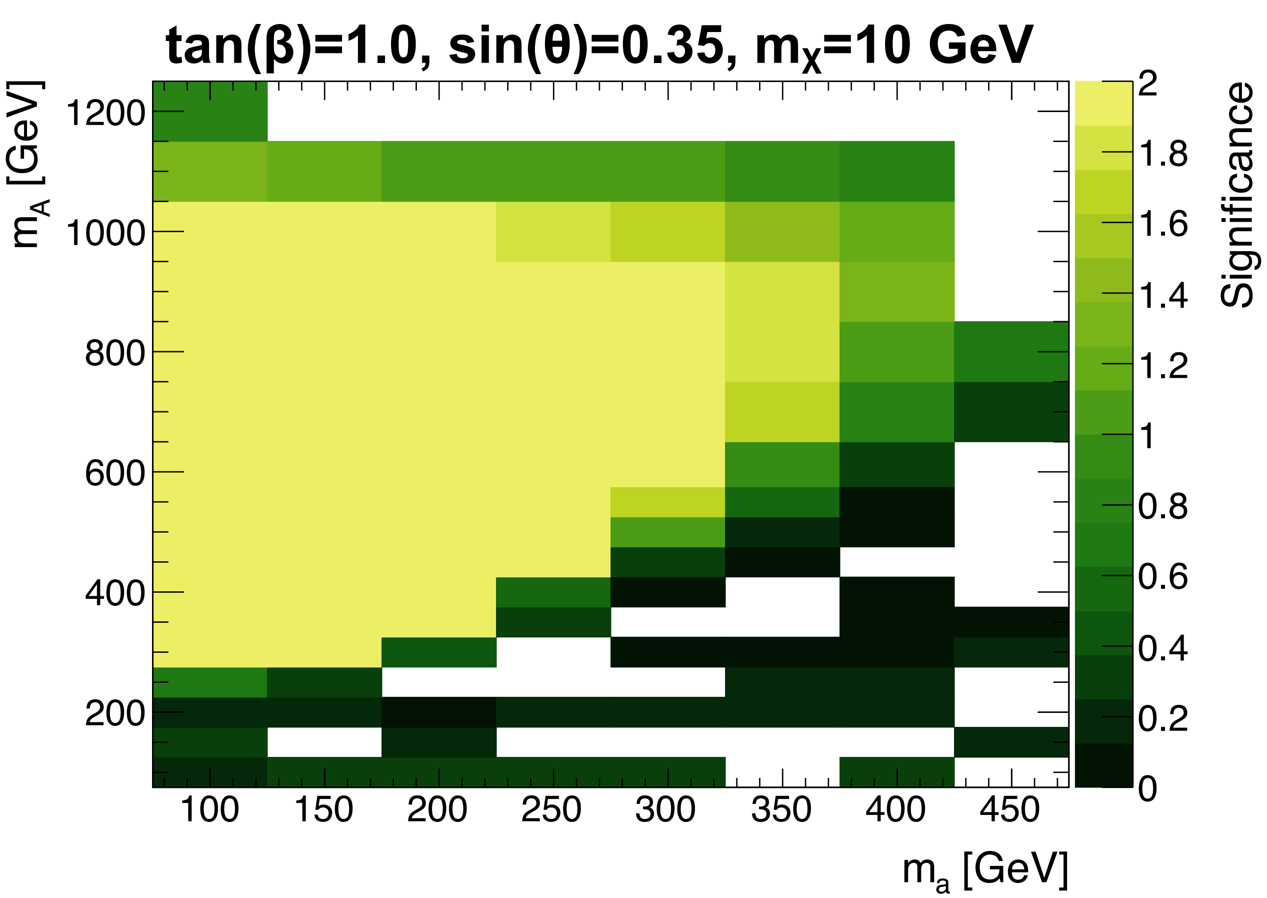


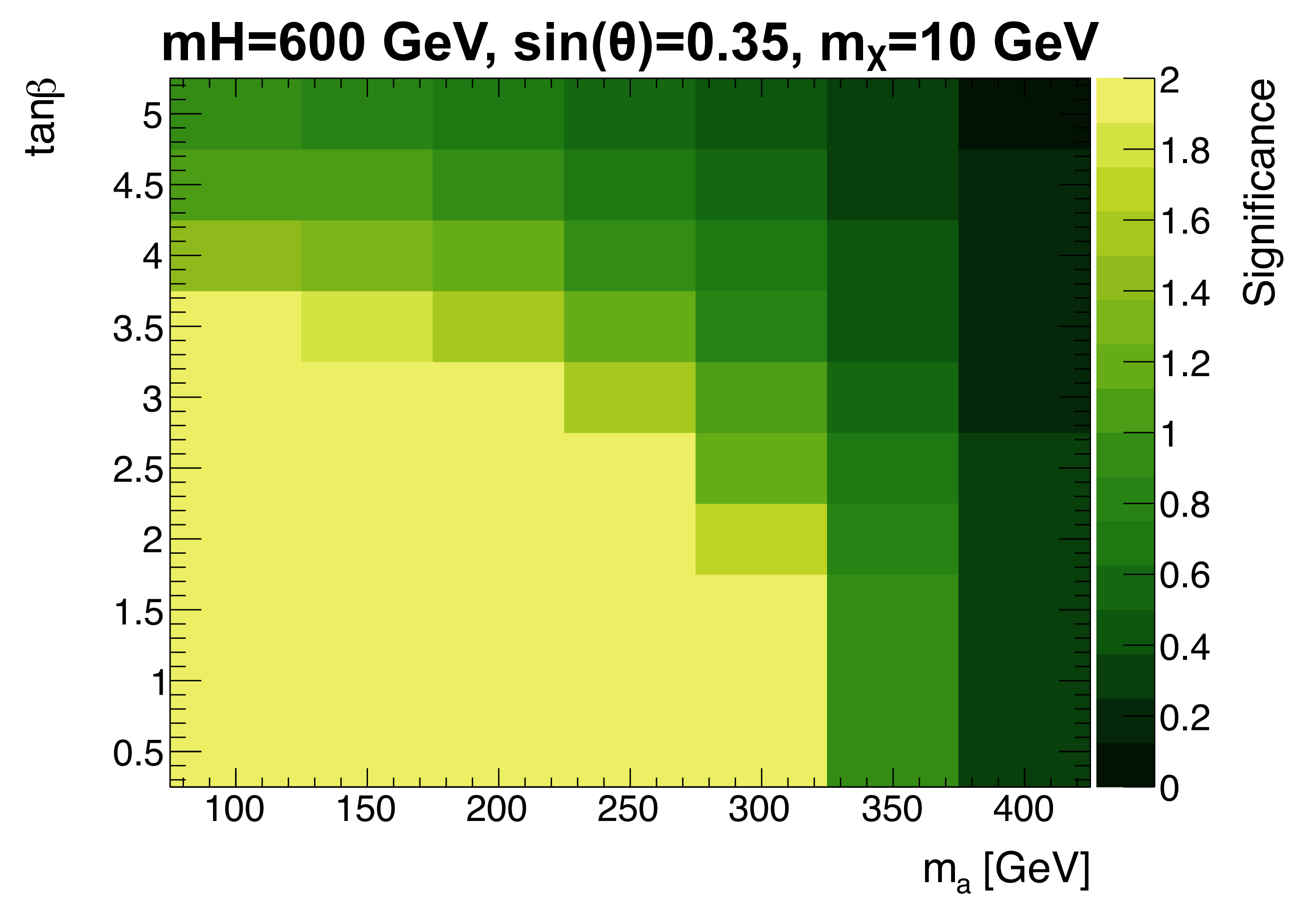
$$E_T^{miss,max} \approx \frac{\sqrt{(M_A^2 - M_a^2 - M_h^2)^2 - 4M_a^2 M_h^2}}{2M_A}$$

Expected Significance (2HDMa)

Mono-Z(ll): Truth-Level Significance

- Calculate significance using Asimov approximation for profile likelihood ratio (Cowan Paper). Per-bin significances summed in quadrature.
- Truth level cuts mirror Mono-Z(ll) (ATLAS) selection, background estimates taken from analysis, and conservative 10-20% background systematic per MET bin applied.
- Assume reconstruction efficiency of 75%.
- Exclude Phase Space with Significance > 2.





DMWG Report : https://indico.cern.ch/event/665524/contributions/2929794/attachments/1623601/2584712/DMWG_2HDM.pdf

Conclusions

- Current ATLAS measurements of the Z(ll) + MET channel are consistent with SM predictions.
- Upper limits are placed on **Simplified Models** of DM, the **Invisible Branching Fraction** of the Higgs, and can be reinterpreted for the **pseudoscalar Two Higgs Doublet Model**.
- As LHC collects more data through Run-2, exclusion limits as well as the potential for discovery should continue to improve.

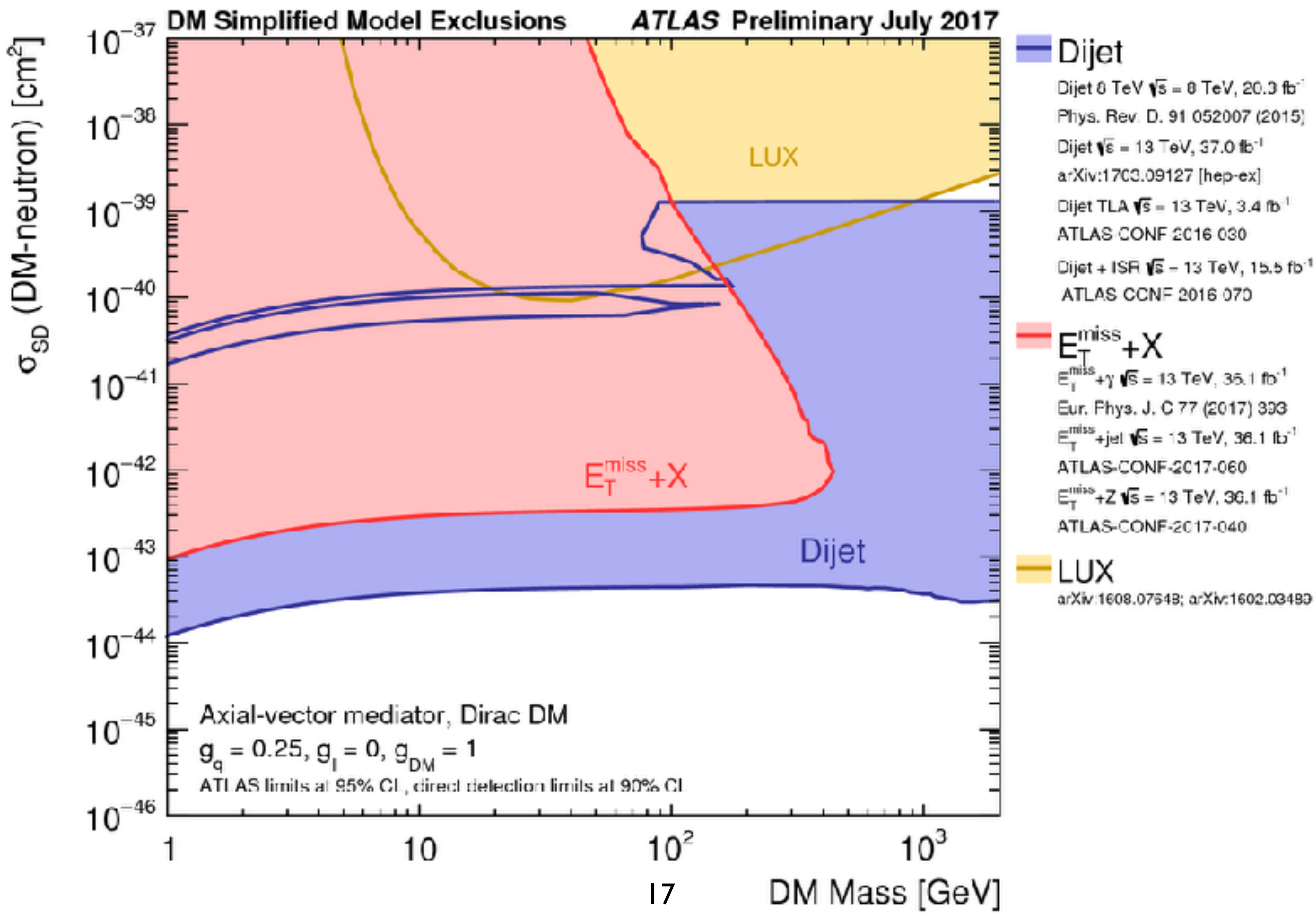
Backup Slides

Comparison to Direct Detection

Simplified models allow for comparison between direct detection experiments and colliders.

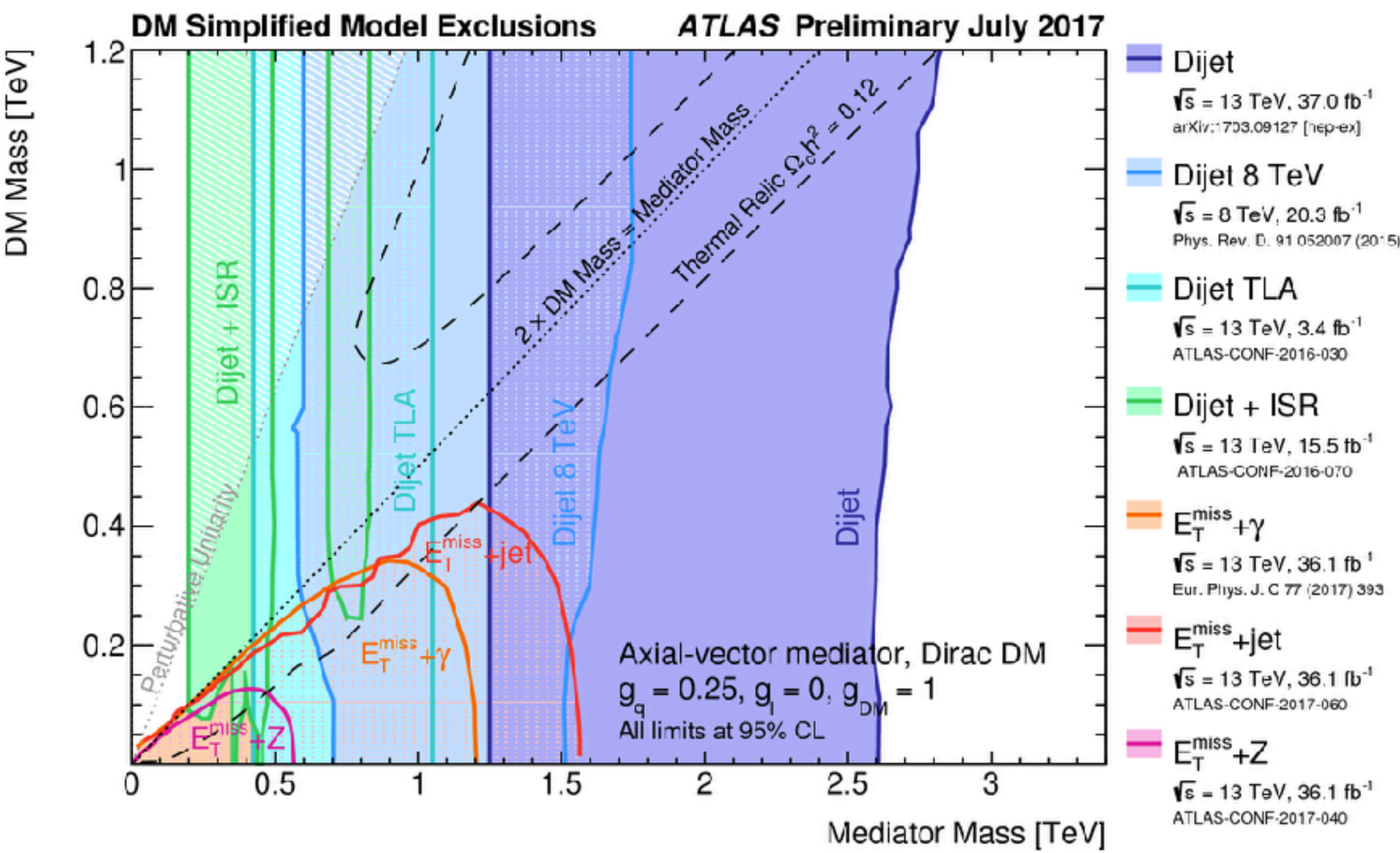
Axial-vector model:

$$\sigma^0_{\mathrm{SD}} \approx 4.6 \times 10^{-41}\ \mathrm{cm}^2 \cdot \left(\frac{g_{\mathrm{DM}}\, g_q}{1}\right)^2 \left(\frac{1\ \mathrm{TeV}}{M_{\mathrm{med}}}\right)^4 \left(\frac{\mu_{n\chi}}{1\ \mathrm{GeV}}\right)^2$$



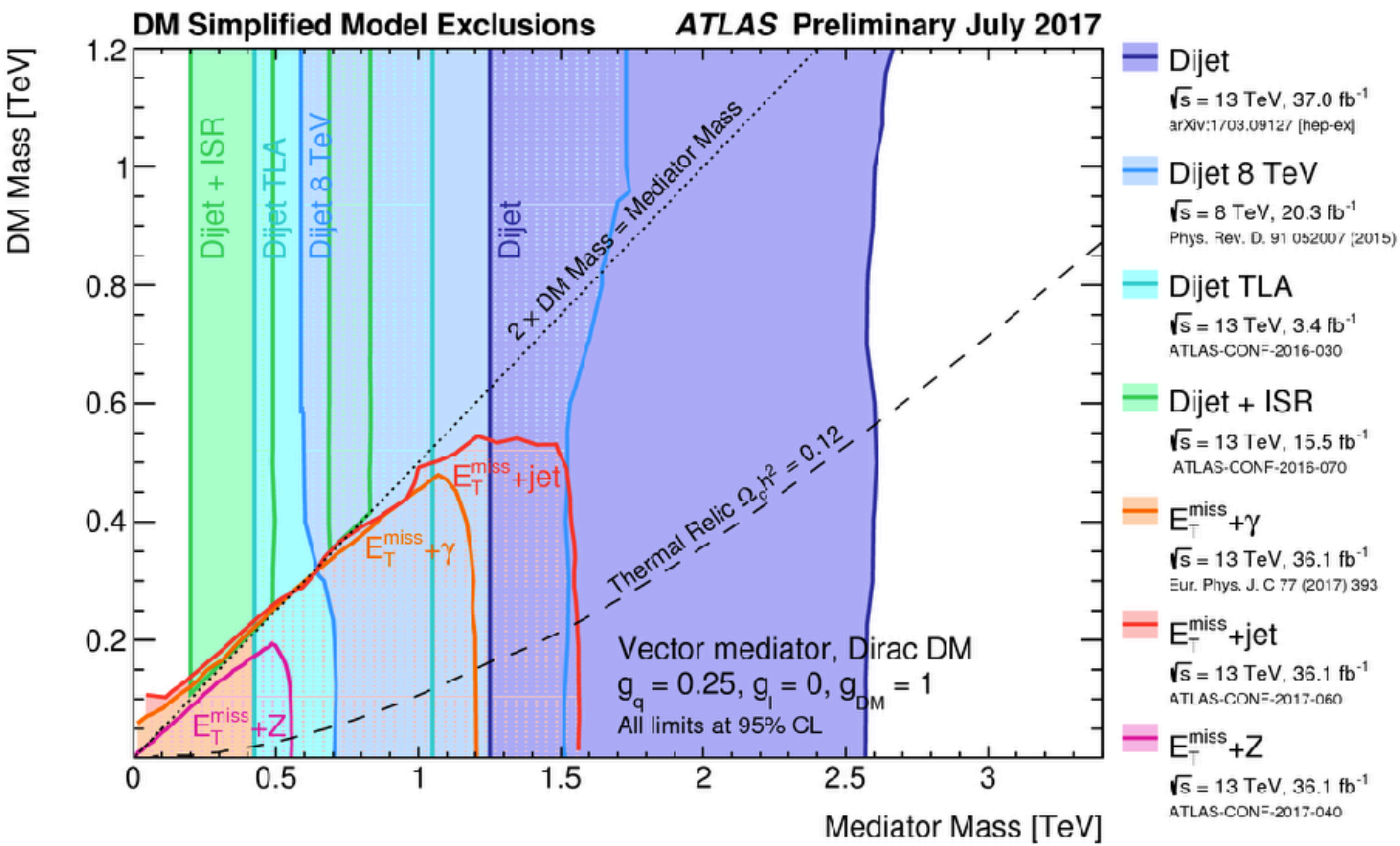
Exclusion Limits

Axial-vector mediator (A1):



Exclusion Limits

Vector mediator (V1):



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Observed limits for Mono-Z(ll) channel.

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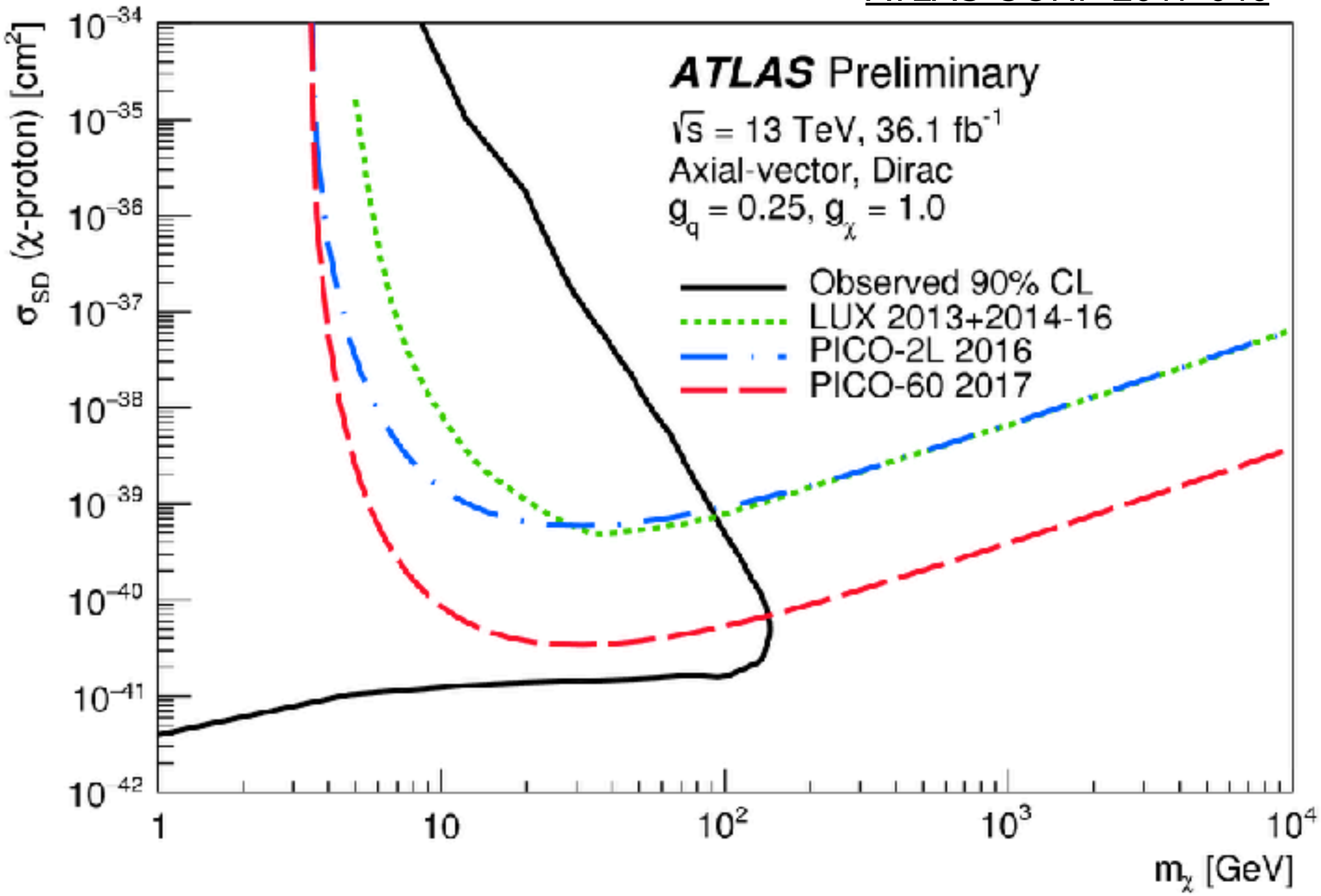

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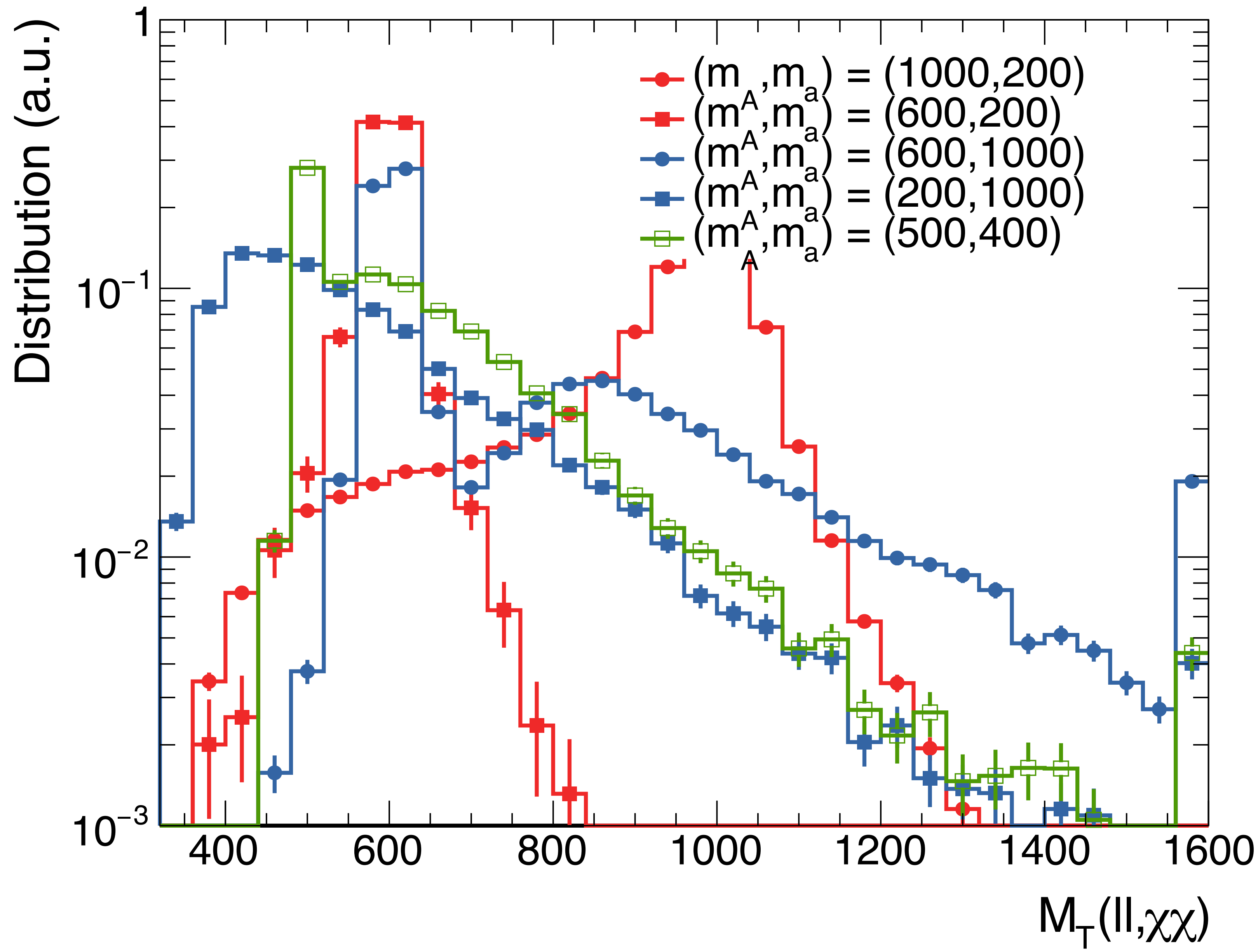
Transverse Mass Distibution for 2HDMa

Transverse mass distribution for 2HMDa model. For resonant production, has characteristic shape.

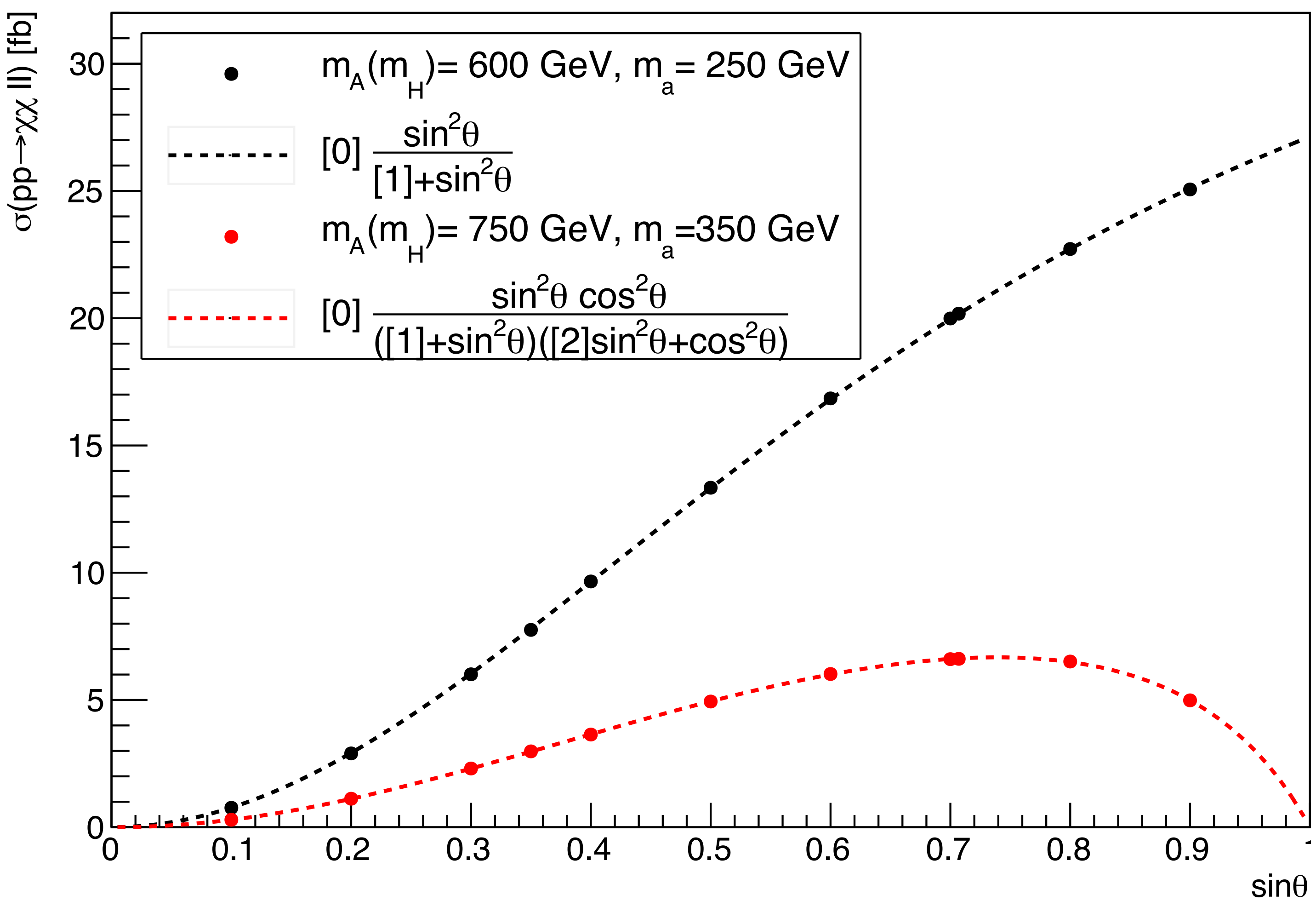


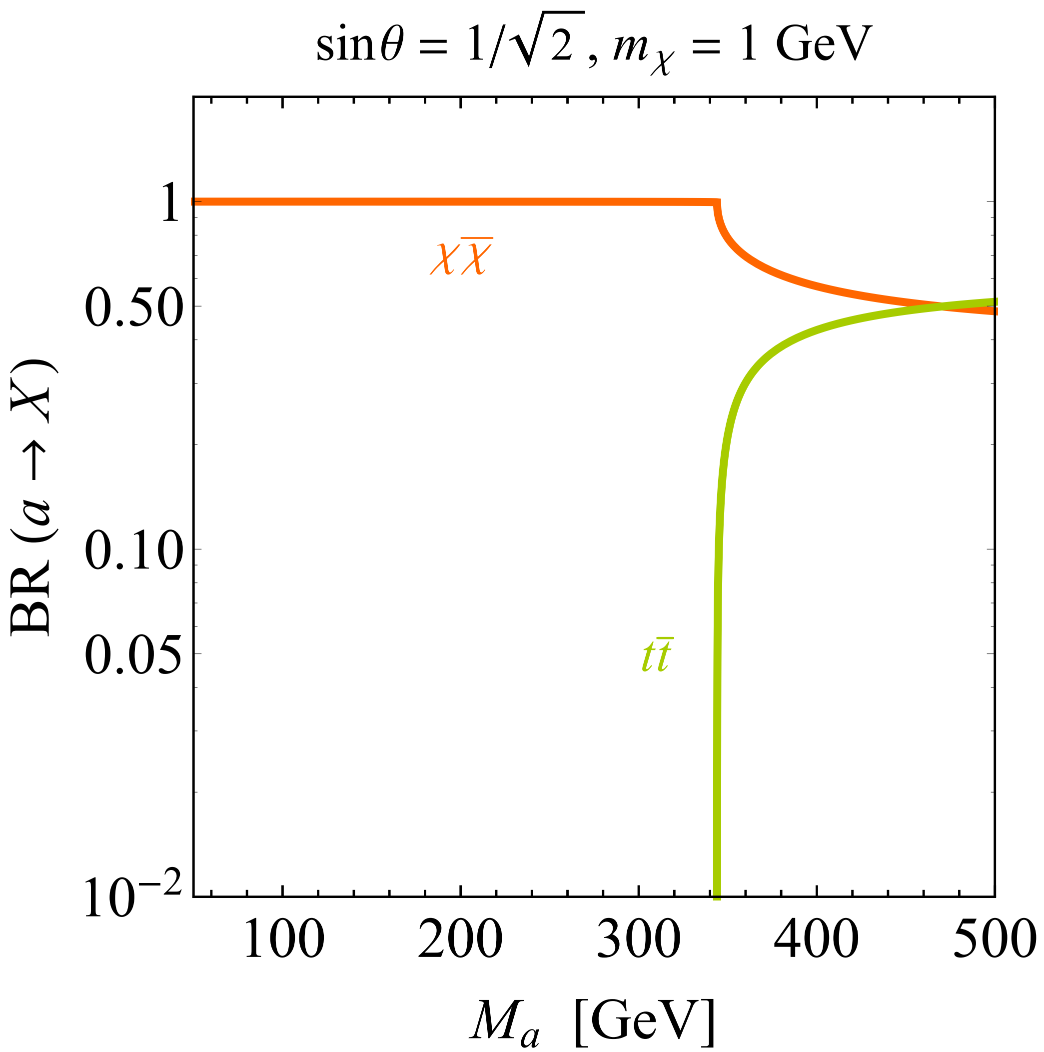
$$M_T^2 = m_1^2 + m_2^2 + 2\left(E_{T,1}E_{T,2} - \vec{p}_{T,1}\cdot\vec{p}_{T,2}\right)$$

Cross Sections

Dependence of cross section on tan(β), M_χ , and sin(θ):

- Cross Section decreases with increasing tan(β).
- Is flat as a function of M_χ, but drops steeply for $2M_\chi > Ma$.
- sin(θ) dependence is interesting. For Ma < 350 GeV, mixing angle only impacts branching fraction of H→aZ and is strictly increasing. For Ma > 350 GeV, decay of the mediator to tt becomes accessible (increasing with sin(θ)) leading to a turnover point.





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ATLAS EXPERIMENT

Object and Event Selection

Single Lepton Triggers

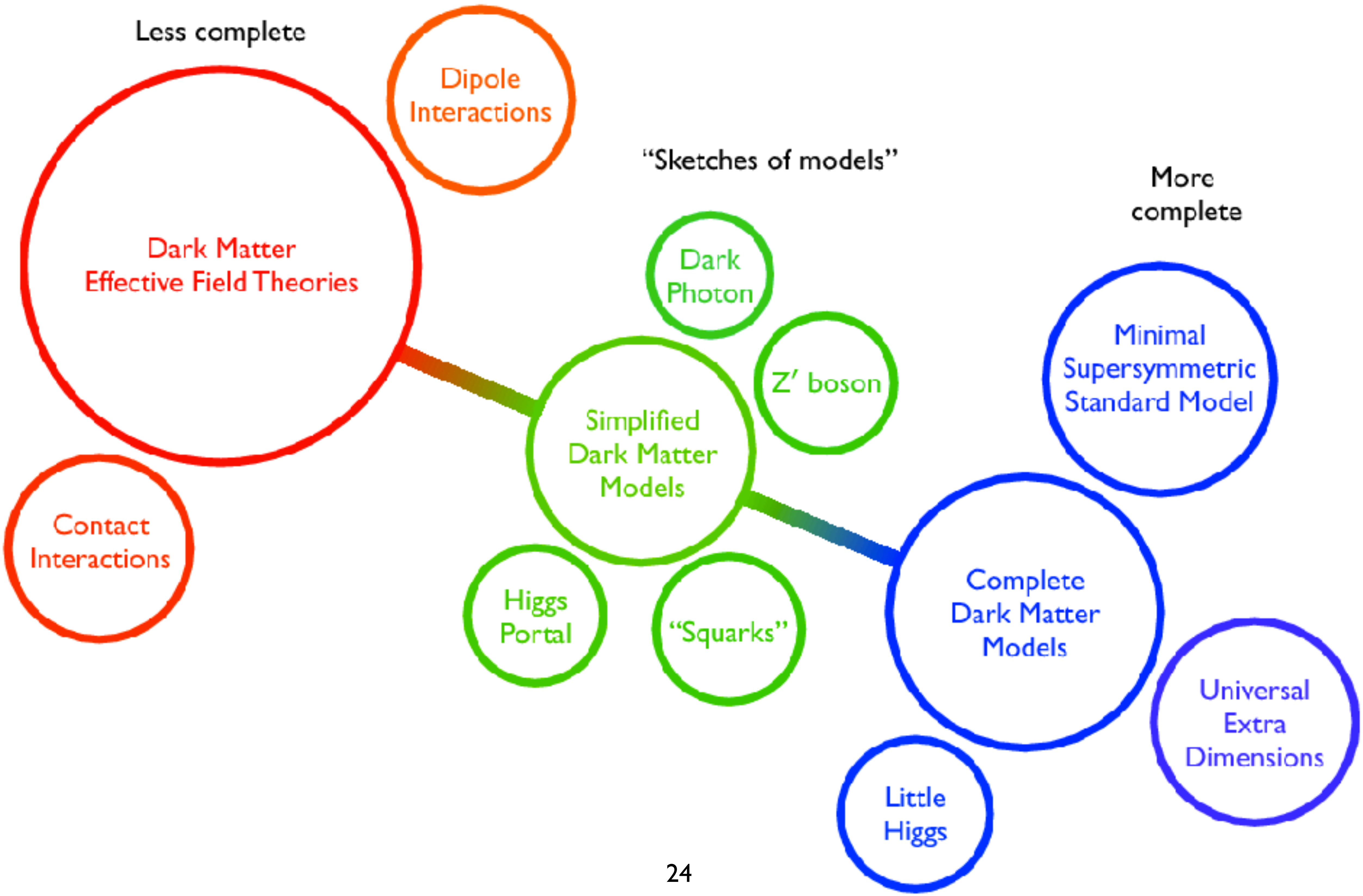
- Low, pT Triggers for isolated electrons (muons). Thresholds at 24 (20)- 26 GeV.
- And high, pT Triggers without an isolation requirement, 50 (60) GeV.

Mono-Z(ll) Selection

- Electrons: isolated, pT > 7 GeV , $|\eta| < 2.47$
- Muons: isolated pT > 7 GeV, $|\eta| < 2.5$
- $E_T^{miss} = - \Sigma$ All reconstructed particles†
- Require <u>exactly</u> 2 opposite sign leptons with lead pT > 30 GeV and sublead pT > 20 GeV.
- Z-window cut ($76 < m_{\ell\ell} < 106$ GeV)
- $E_T^{miss} > 90$ GeV and $E_T^{miss} / H_T > 0.6$.
- $\Delta\varphi(p_T^{\ell\ell}, E_T^{miss}) > 2.7$
- $\Delta R_{\ell\ell} < 1.8$
- Fractional p_T difference : $|p_T^{\ell\ell} - p_T^{miss,jets}) > 2.7 | / p_T^{\ell\ell} < 0.2$
- B jet veto

†And soft tracks not associated with a particle.

Dark Matter Models

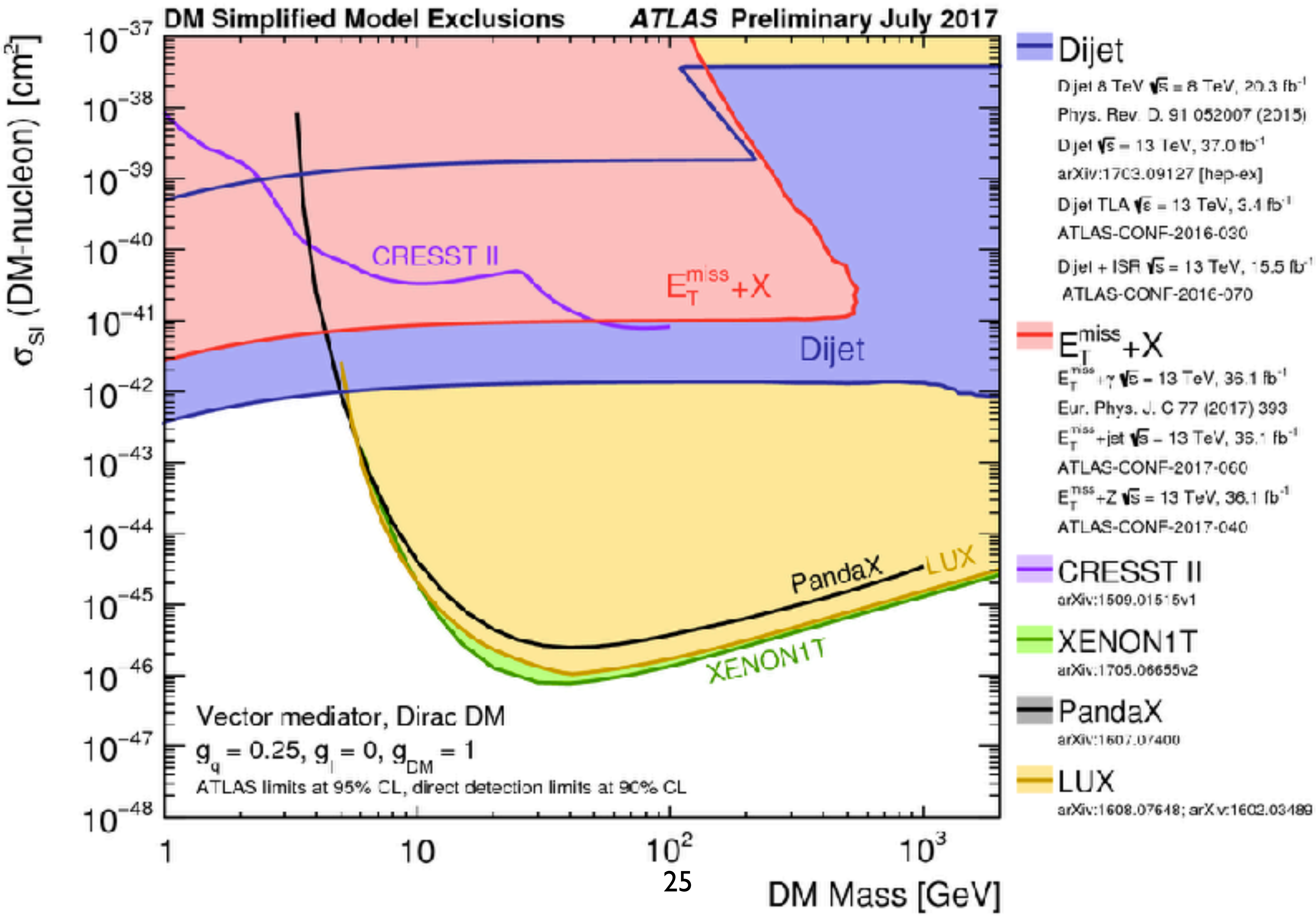


Comparing to Direct Detection

Simplified models allow for comparison between direct detection and collider results:

Vector model:

$$\sigma^0_{\rm SI} \approx 1.1 \times 10^{-39}\ {\rm cm}^2 \cdot \left(\frac{g_{\rm DM}\, g_q}{1}\right)^2 \left(\frac{1\ {\rm TeV}}{M_{\rm med}}\right)^4 \left(\frac{\mu_{n\chi}}{1\ {\rm GeV}}\right)^2$$



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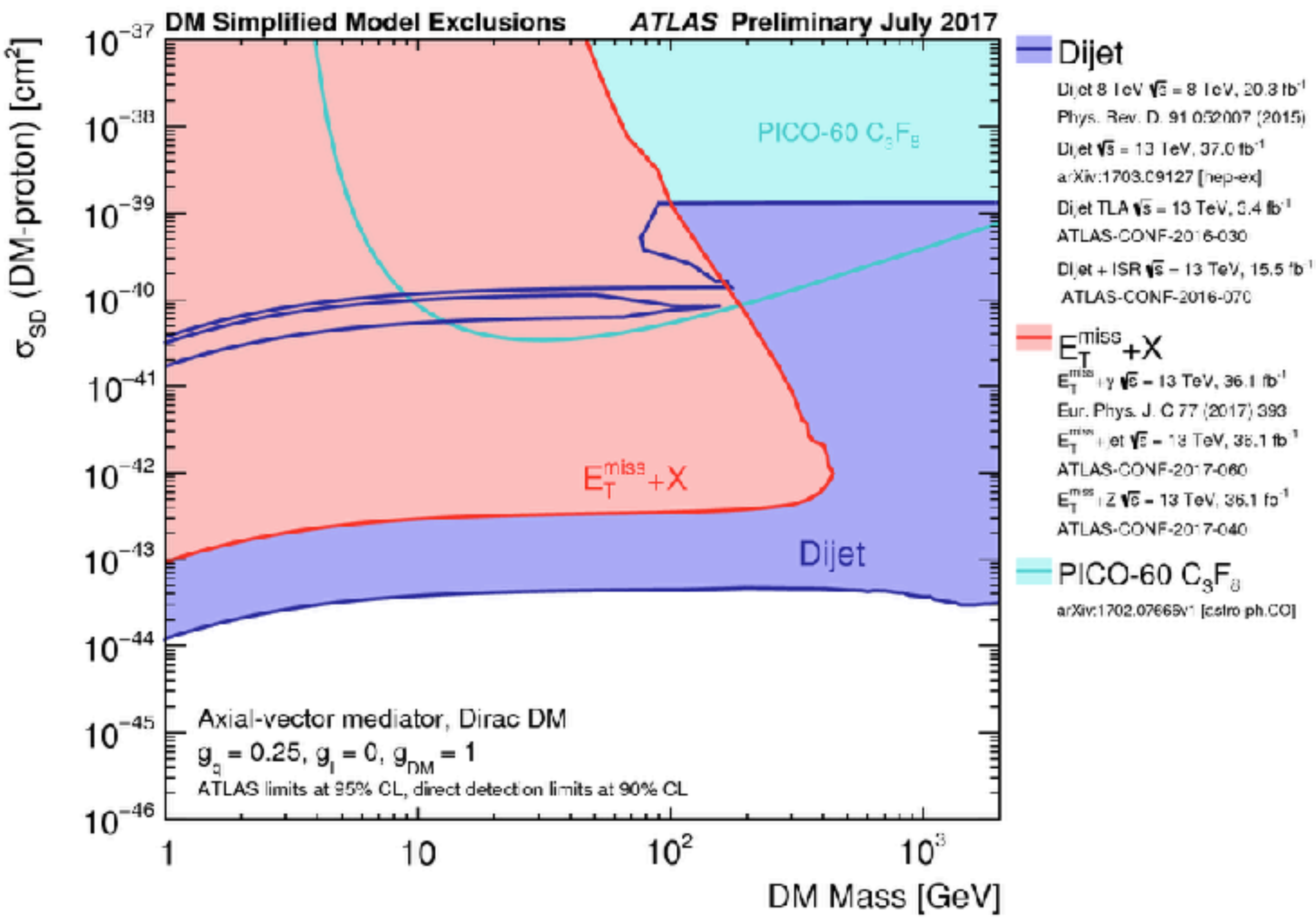


Significance Caluculation

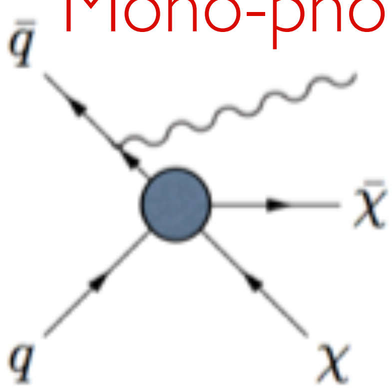
- (Cowan Paper)

$$Z'_{bin} = \sqrt{2 \cdot \left((s+b)\ln[\frac{(s+b)(b+\sigma_b^2)}{b^2+(s+b)\sigma_b^2}] - \frac{b^2}{\sigma_b^2}\ln[1+\frac{\sigma_b^2 s}{b(b+\sigma_b^2)}] \right)}$$

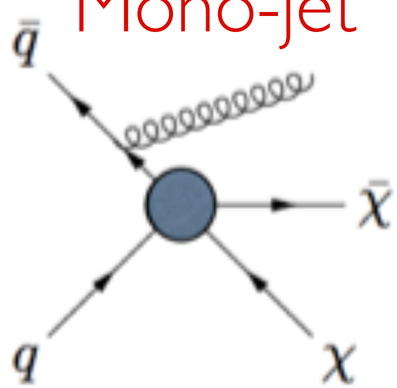
Mono-X Searches

Mono-X searches look for p_T imbalances created from DM production recoiled against a detectable particle.

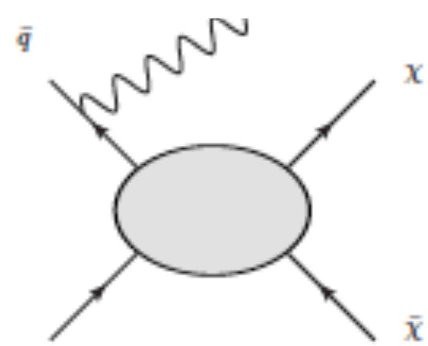
Mono-photon



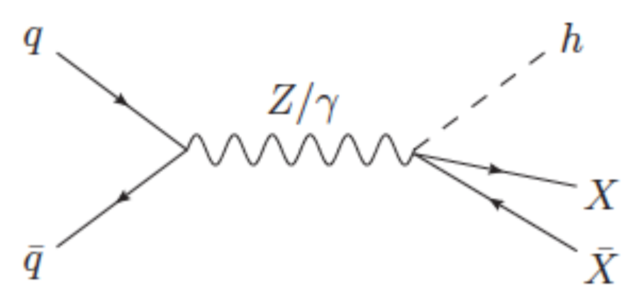
Mono-jet



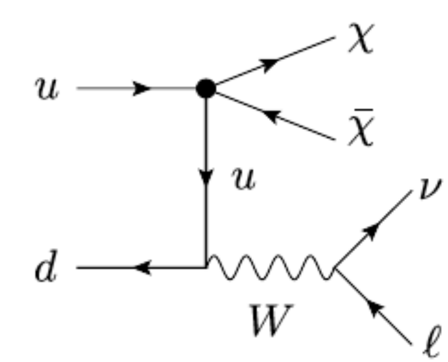
Mono-Z



Mono-Higgs



Mono-W



†And soft tracks not associated with a particle.

Cross Sections

MadGraph + Pythia generators were used to simulate 2HDMa events and calculate cross sections at LO, g g > xd xd~ l+ l-

Dependence of cross section on m_H and m_a:

- Destructive interference for $m_H = m_a$.
- For m_A, m_a > 350 GeV, cross sections decrease as pseudoscalars can also decay to ttbar. Impact on m_a is minimized for small sin(θ).