Transverse Effects in the Hollow Channel Plasma Accelerator Carl A. Lindstrøm, University of Oslo

CLIC Novel Accelerator Working Group, CERN – Oct 13, 2017

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# Transverse Effects in the Hollow Channel Plasma Accelerator

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## Why use a hollow channel?

- Bigger picture:
  - If you want an electron-positron collider, you need positron acceleration (!)
  - You therefore need to charge-symmetrize your technology

(or make another different, but equally good technology)

• Main reasons for using a hollow plasma channel are



• However, we soon discover that unless you are perfectly aligned, there is a strong transverse wakefield.



## A useful analogy

#### Born with a fatal illness



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## PART 1: Conception

Hollow plasma channel transverse wakefield theory



## Linear model

- Assume:
  - plasma electrons are not moving (just slightly perturbed)
  - that the plasma behaves as a dielectric medium
- The most significant longitudinal wakefield (for acceleration):

$$W_{z0}(z) = \frac{ek_p \chi_{\parallel}^2}{2\pi\epsilon_0 a} \frac{B_{00}(a,b)}{B_{10}(a,b)} \cos(\chi_{\parallel} k_p z) \Theta(z)$$

• The most significant transverse wakefield is a dipole mode:

$$W_{x1}(z) = -\frac{e\Delta x\chi_{\perp}}{\pi\epsilon_0 a^3} \frac{B_{11}(a,b)}{B_{21}(a,b)} \sin(\chi_{\perp}k_p z)\Theta(z)$$

• When the response is non-linear, we turn to particle-incell (PIC) simulations.



$$\begin{split} \chi_{\parallel} &= \sqrt{\frac{2B_{10}(a,b)}{2B_{10}(a,b) - k_p a B_{00}(a,b)}} \\ \chi_{\perp} &= \sqrt{\frac{2B_{21}(a,b)}{4B_{21}(a,b) - k_p a B_{11}(a,b)}} \end{split}$$

$$B_{ij}(a,b) = I_i(k_p a) K_j(k_p b) + (-1)^{i-j+1} I_j(k_p b) K_i(k_p a)$$



Comparing numbers to conventional wakefields

- Typical CLIC transverse wakefields per offset in structures: ~1-100 V/pC/m/mm
- Hollow channel (500 µm diameter at 3 x 10<sup>15</sup> cm<sup>-3</sup>):
  ~150 V/m^2/particle = ~1 000 000 V/pC/m/mm





Figure 8: The envelope of the transverse wakefields for both planes is shown. The CLIC bunch spacing is 0.15 m.



## PART 2: Diagnosis

## Experiments at FACET, SLAC



## The E225 experiment

- One of these experiments was the E225 Hollow Channel experiment, lead by Spencer Gessner.
- E225 successfully demonstrated acceleration of a trailing positron bunch in a positron driven hollow plasma channel.
- Transverse wakefields were also measured in the E225 experiment.





Spencer Gessner (left) and Sebastien Corde (right) at FACET tunnel, SLAC. Image source: SLAC National Accelerator Laboratory



## E225 – Experimental setup



- The SLAC linac provided two 20 GeV bunches, made from one bunch using a beam notching device.
- The FACET laser (up to 10 TW, 60 fs pulses) was adjusted down to ensure no ionisation in the channel.
- A lithium oven was set to give a neutral gas density of 3x10<sup>16</sup> cm<sup>-3</sup> (but was necessarily fully ionized).



## The experiment

- Our goal was to measure the how the transverse wakefield varied longitudinally.
- The probe bunch observing the wakefield is deflected angularly (kicked) when the channel and the drive bunch are relatively offset.
- The experiment performed was:
  Transverse channel offsets
  for various bunch separations
  - The channel (250 µm radius) was offset by transverse laser jitter (20-40 µm rms)
  - The bunch separation was varied by stretching the bunch and adjusting the notching device.
- Diagnostics:
  - Laser offset imaged downstream (laser cameras).
  - Probe kick measured on a spectrometer (in the non-dispersed plane).
  - Bunch separation measured using an **electro-optical sampler**.

#### Prediction:



#### Experiment (2D "scan"):

#### Varying bunch separations (scanned)





## Another independent measurement

- An independent measurement is beneficial (due to high complexity).
- It is possible to estimate the transverse wakefield per offset from the measured longitudinal wakefield, via the Panofsky-Wenzel theorem and the linear model.

Estimate of transverse from longitudinal wakefield:

$$\frac{\partial W_x}{\partial z} = \frac{\partial W_z}{\partial x} \quad \stackrel{\text{Integrate (++)}}{\longrightarrow} \quad \frac{W_x(z)}{\Delta x} \approx -\frac{\kappa(a,b)}{a^2} \int_0^z W_z(z') dz' \quad \text{where} \quad \kappa(a,b) = \frac{4\chi_{\perp}^2 - 2}{\chi_{\parallel}^2 - 1}$$

- Not perfect: Assumes linear model, breaks down far behind the drive bunch.
- Provides verification of numerical calibrations, etc.

Panofsky-Wenzel theorem:

• The longitudinal wakefield was measured by the energy change of the probe bunch (on a spectrometer).



## Observed data (deflection vs. channel offset)



- For each bunch separation, a correlation between channel offset and probe bunch angular deflection was observed.
- The slope of this correlation is proportional to the transverse wakefield per offset at the z-location of the probe bunch.



## Final experimental results



- Plasma density determined by a wavelength fit (10% ionization =  $3 \times 10^{15}$  cm<sup>-3</sup>)
- Good fit, largely consistent with theory. Some discrepancy at larger separations.



### Implications

- Overall, the measurement agrees with the theoretical models.
- Simulation-based parameter scans indicate that the <u>discrepancy</u> at large separations can possibly be <u>explained by using a more complex radial</u> <u>plasma shape</u> (not possible to exclude with our diagnostics).
- Implication: There is indeed a strong transverse wakefield.



• Submitting these results to Physical Review Letter.



## PART 3: Finding a cure

## New ideas for fixing the hollow channel



## Fundamental problem: transverse wakefields



Ideas for transverse wakefield suppression

- Conventional BNS damping
  => Requires very large energy spreads
- External focusing  ${}^{\bullet}$ => Requires very strong focusing => Solenoids do not scale well with energy (1/ $y^2$ , but kick scales with 1/y) => Does not help unless perfectly aligned <sup>~l</sup> center Near hollow channels (low density plas => Breaks the child e symmetry Tailoring locally c Transverse Effects in the Hollow Channel Plasma Accel



## Clarification on what needs to be damped

- Two different problems
  - Transverse instability of the drive beam (feedback loop)
  - Transverse deflection of accelerated beam caused by drive beam wakefield

Transverse instability of the accelerated beam (from its own transverse wakefield)

- The drive beam only needs to **survive one stage**. It does not need to preserve its emittance.
- The accelerated beam can in principle be placed on the zero crossing of the transverse wakefield from the driver.
- The self-wakefield can in principle be used to transversely beam load the transverse wakefield of the driver (at the zero crossing).
- Now we just need to make a hollow channel with all the above + high accelerating gradient + high efficiency PEASY!





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## Tailored radial plasma profiles

- Currently just at the speculative stage: Hollow plasmas are naturally not high-Q cavities: can we use this to our advantage?
- A strength of plasma channels compared to structure-based channels: Many degrees of freedom to play with (radial distribution of plasma density)
- This work will benefit from a theoretical model for arbitrary radial plasma profiles n(r) (as an alternatively a brute force PIC simulation search)
- If it actually works, the next step is to somehow make this plasma profile (with the kinoform of all kinoforms?)





## Conclusions

- Strong wakefields is a big challenge for the hollow channel
- We did an experiment at FACET to diagnose the problem
- Experiments correspond well with theory (but still some discrepancies)
- We have ideas on how to fix it (with radially tailored plasma profiles?)
- Will be pursued at FACET-II

### A Haiku of Hope

We do not de-clare

the ho-llow cha-nnel dead yet. First try all id-eas.



## Thank you for your attention!