



Pelourinho, Salvador – colourful painting of the old colonial houses by street artist from yesterday tour

Measurement of cross sections and properties of the Higgs boson using the ATLAS detector

Luigi MARCHESE, Oxford university

on behalf of the ATLAS collaboration

LISHEP2018, 9-14 September, Salvador de Bahia



UNIVERSITY OF
OXFORD



OUTLINE

- Introduction to Higgs physics
- Higgs properties and cross sections
 - couplings
 - Higgs width
 - Higgs mass
- Conclusions

 Several new ATLAS results released in the past months



INTRODUCTION

- We have just recently celebrated the 6th Higgs birthday
- The Higgs boson discovery opened a new window into the sector of the SM Lagrangian responsible for EW symmetry breaking!
- Precision measurements are vital to test theoretical models
 - Initial results in Run 1
 - Many results in Run 2



THE HIGGS BOSON IN THE STANDARD MODEL

- By interacting with all the SM particles, the Higgs field gives them mass
 - Two different types of tree-level couplings

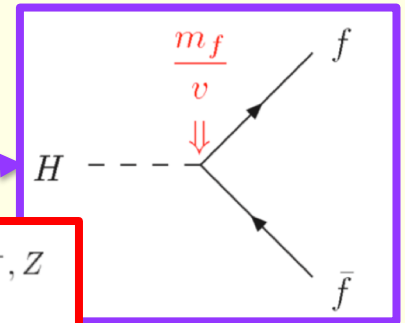
$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i \bar{\psi} \not{D} \psi + h.c.$$

$$+ \bar{\psi}_i y_{ij} \psi_j \phi + h.c.$$

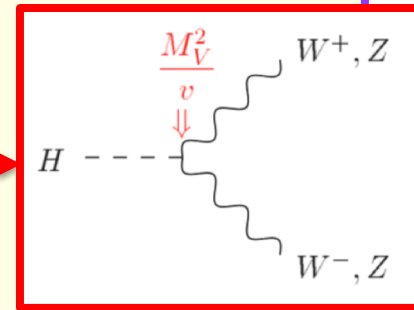
$$+ \left(\frac{1}{2} D_\mu \phi \right)^2 - V(\phi)$$



Fermions



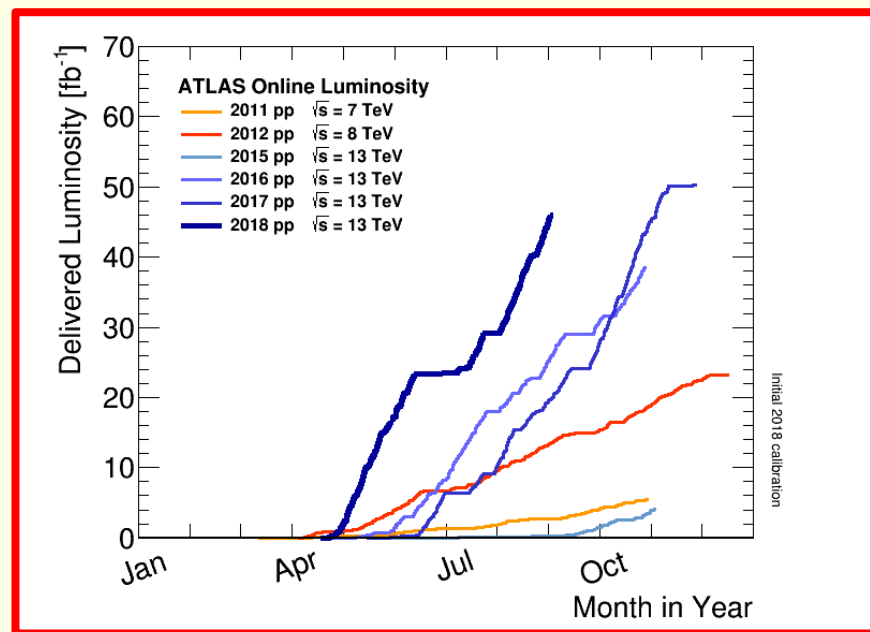
BOSONS





HIGGS PHYSICS AT LHC

- At the Large Hadron collider the delivered luminosity was:
 - 25 fb^{-1} at 7/8 TeV - Run 1 **Higgs discovery!**
 - 36 fb^{-1} (2015+2016), 44 fb^{-1} (2017) at 13 TeV - Run 2

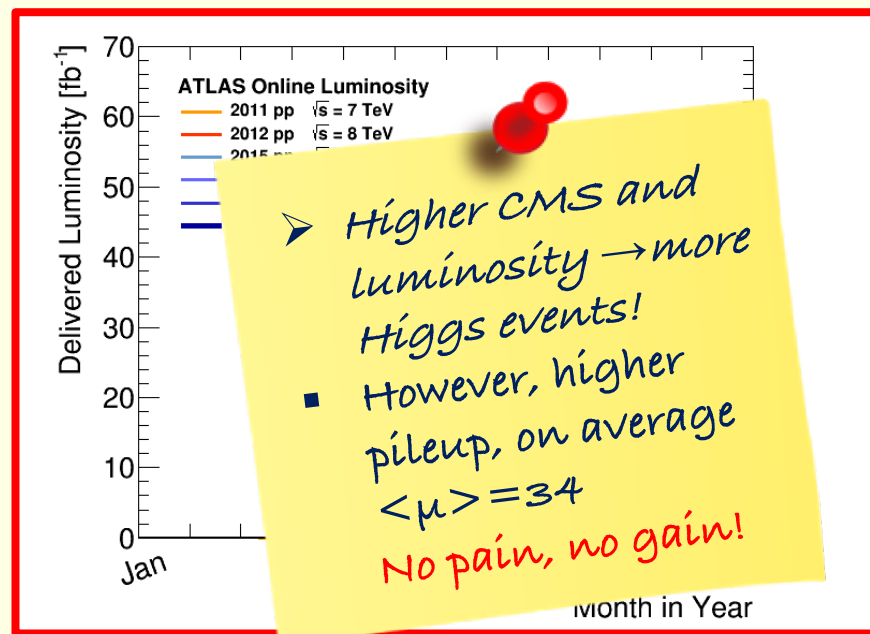


- 1 Higgs boson produced every 10^9 proton-proton collisions

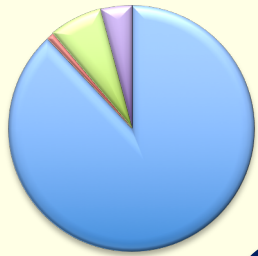


HIGGS PHYSICS AT LHC

- At the Large Hadron collider the delivered luminosity was:
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 - 36 fb^{-1} (2015+2016), 44 fb^{-1} (2017) at 13 TeV - Run 2

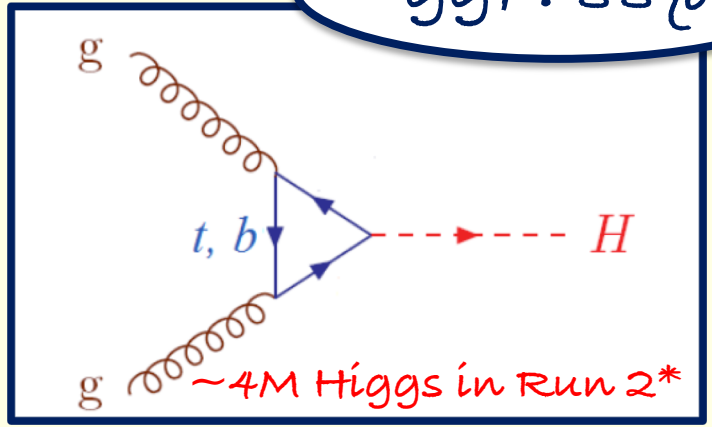


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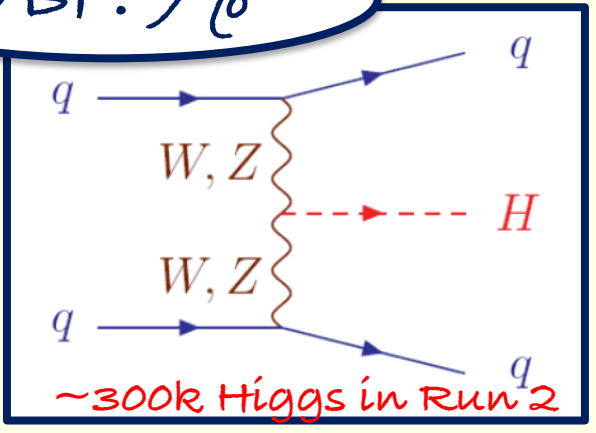


HIGGS PRODUCTION AT LHC

ggF: 88%



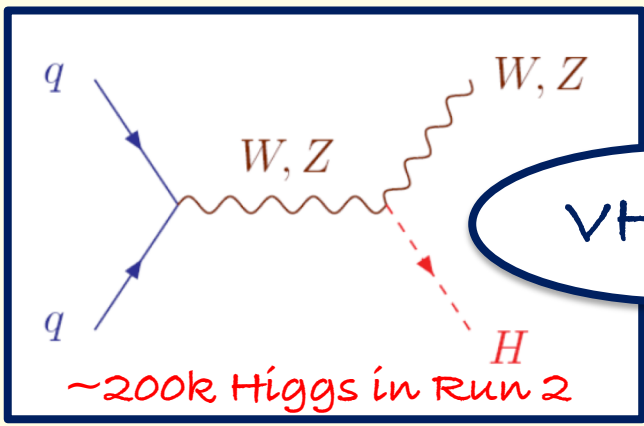
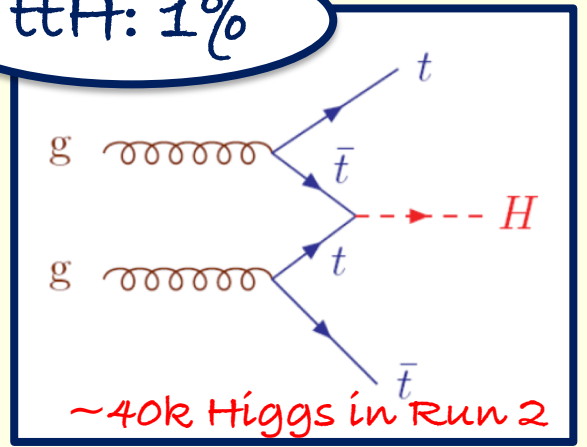
VBF: 7%



How should I be produced this time?



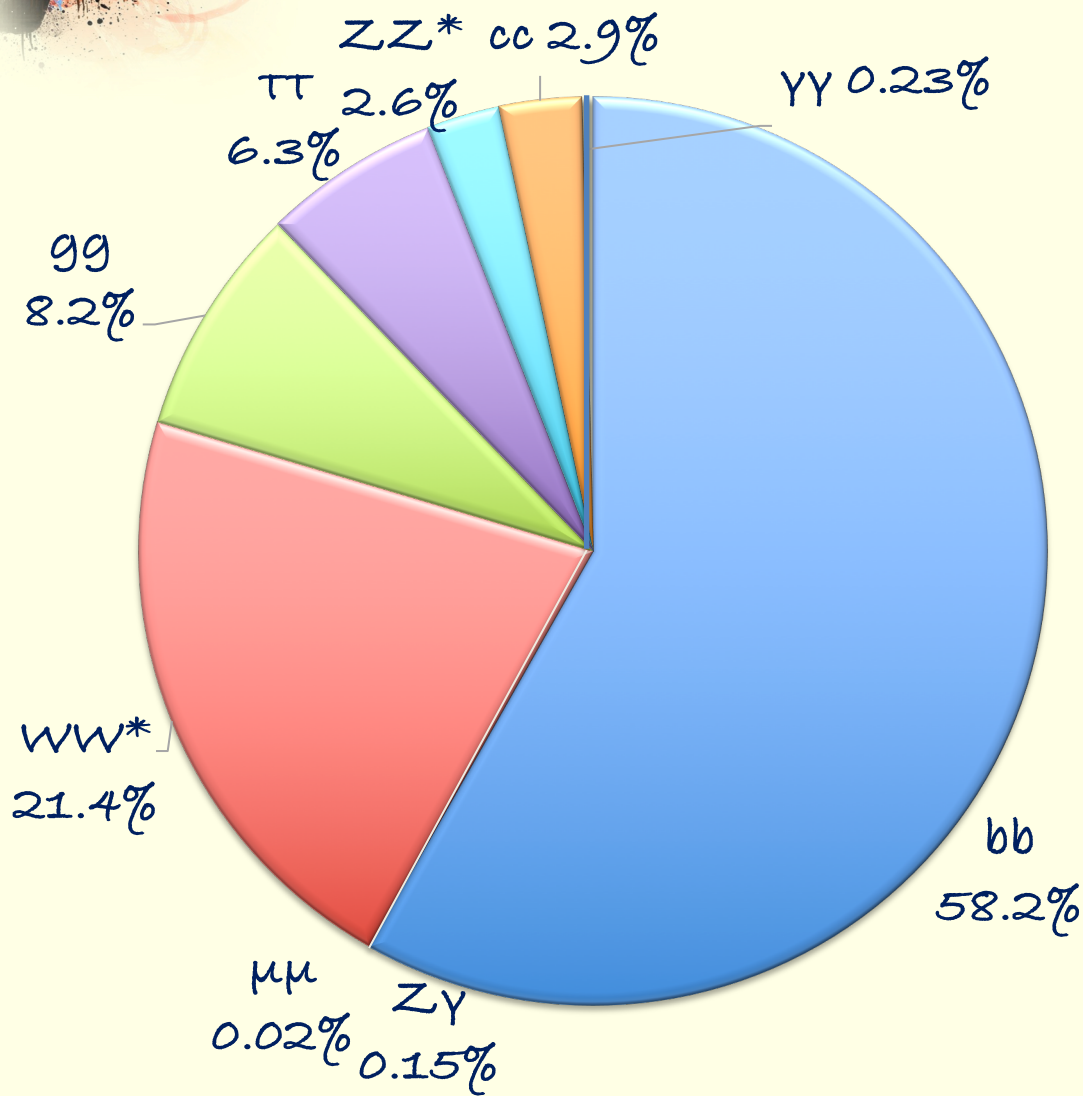
ttH: 1%



VH: 4%

*Run 2 refers to 2015-2017 at 13 TeV

HIGGS BOSON DECAY MODES



➤ Analyses ongoing in all the main channels, directly or indirectly
→ No couplings measurements without theory assumptions

➤ Combination of all the channels is crucial:

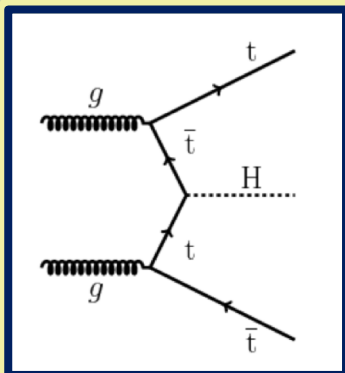
- to increase sensitivity
- to resolve ambiguity

HIGGS COUPLINGS

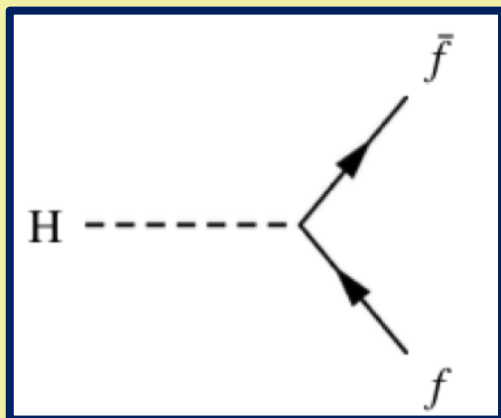


DIRECT PROBES

at tree level



ttH



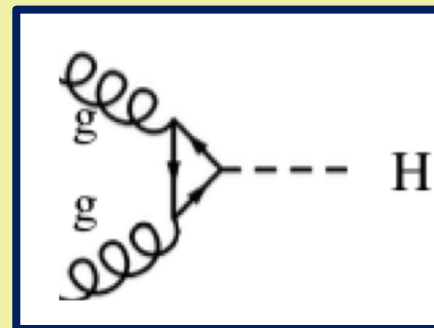
Different generations:

- Leptons
- Up-type quarks
- Down-type quarks

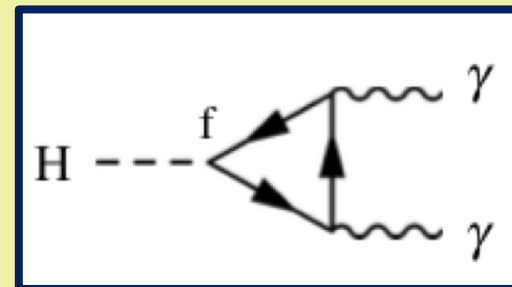


INDIRECT PROBES

via loop diagrams

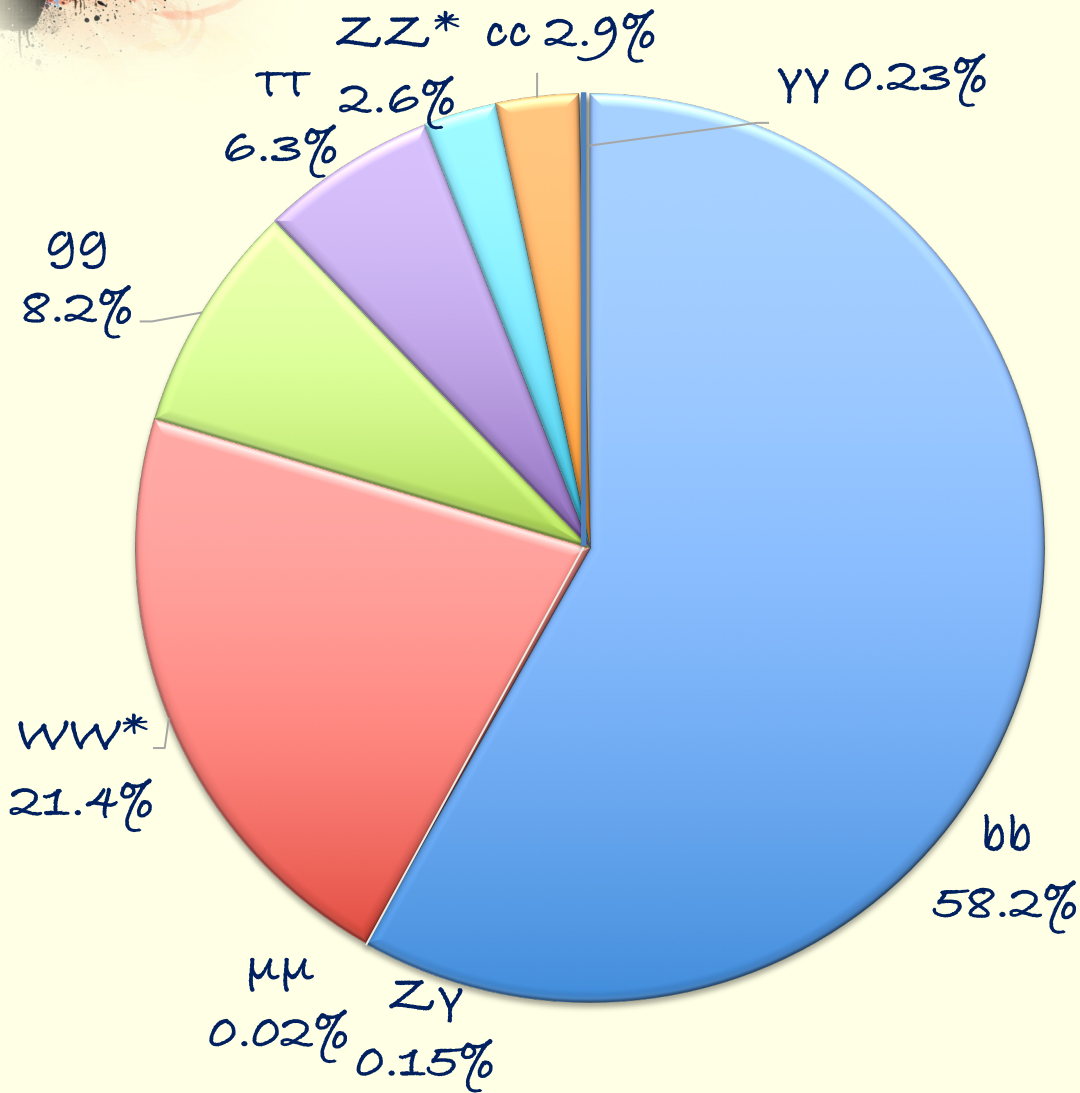


gluons



photons

HIGGS TO BOSONS



$ZZ^*, \gamma\gamma$



Good mass resolution

- Ideal for precision measurements since well modelled background and clear signatures



Low BR, especially ZZ^* , 0.012% in 4l

WW^*



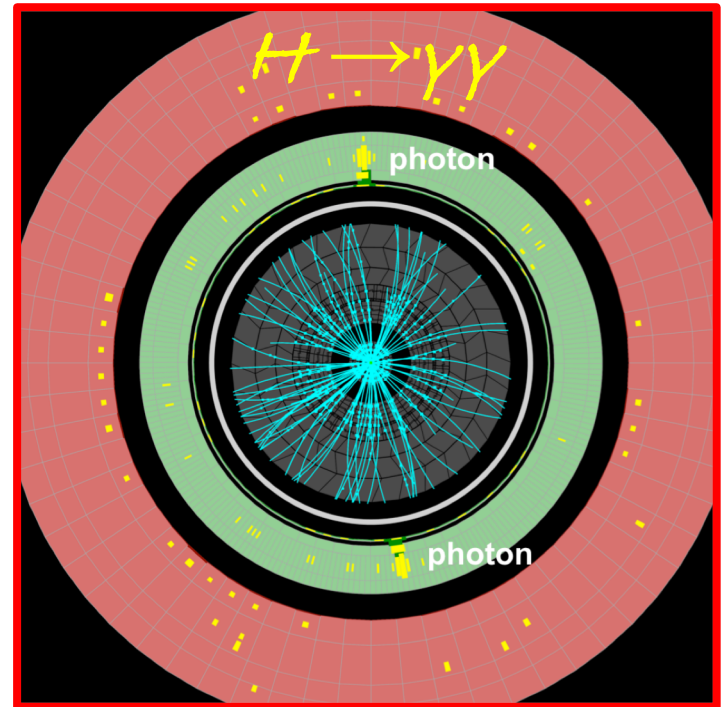
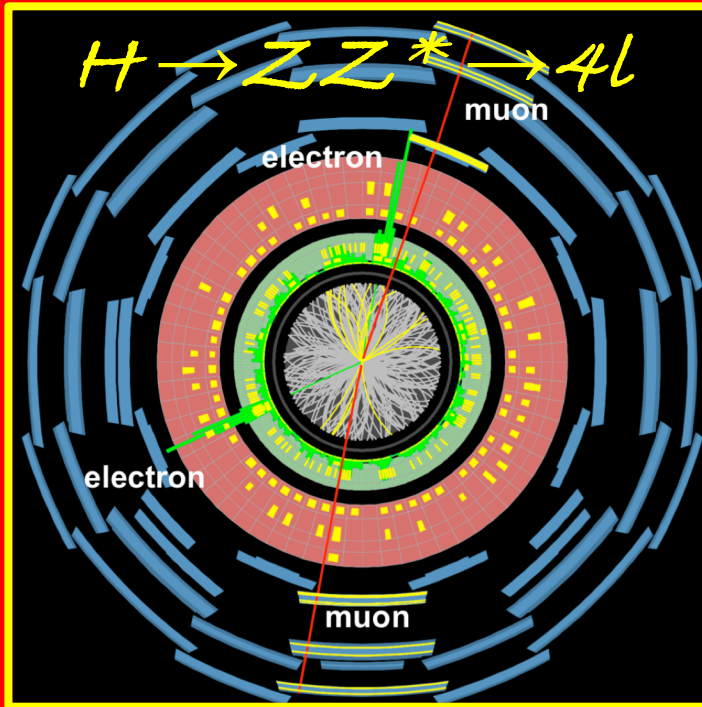
High BR, but reduced in the dilepton mode, 1.1%



Low mass resolution because of ν_s in final states



$H \rightarrow ZZ^*$ and $H \rightarrow \gamma\gamma$: DISCOVERY CHANNELS



Run-2 results (79.8 fb^{-1}):

[ATLAS-CONF-2018-018](#)

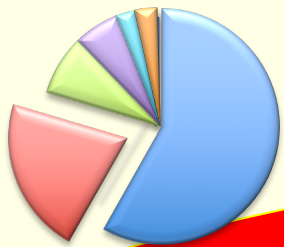
$$\sigma_{\text{ggF}} \mathcal{B}_{HZZ^*} = 1.22 \pm 0.18 \text{ pb}$$

$$\sigma_{\text{VBF}} \mathcal{B}_{HZZ^*} = 0.25 \pm 0.09$$

[ATLAS-CONF-2018-028](#)

$$\sigma_{\text{inclusive}} = 60.4 \pm 6.1 \text{ (stat.)}$$

$$\pm 0.3 \text{ (theo.) fb}$$



$$H \rightarrow WW^*$$

➤ Most precise cross-section measurements in Run 1

➤ Focus on leptonic decays

$$H \rightarrow WW^* \rightarrow l\nu l\nu$$

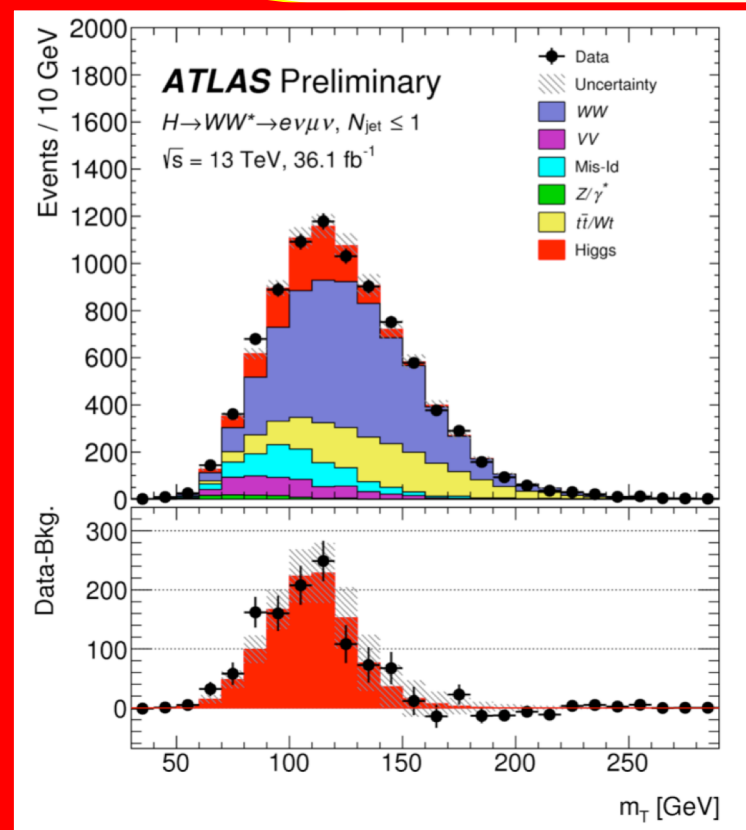
- Hard signature because of MET (i.e. Missing Transverse Energy)
- MET gets harder as pileup rises

Run-2 cross section measurements:

$$\sigma_{ggF} \mathcal{B}_{HWW^*} = 12.6^{+1.3}_{-1.2} \text{ (stat.) } ^{+1.9}_{-1.8} \text{ (syst.) pb and } \mu_{ggF} = 1.21^{+0.22}_{-0.21}$$

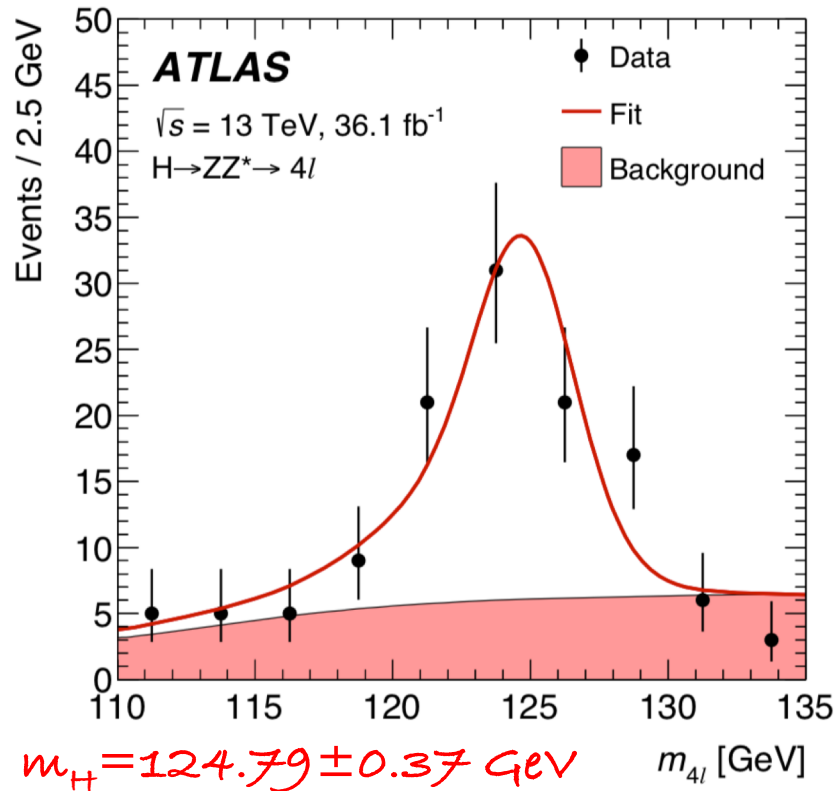
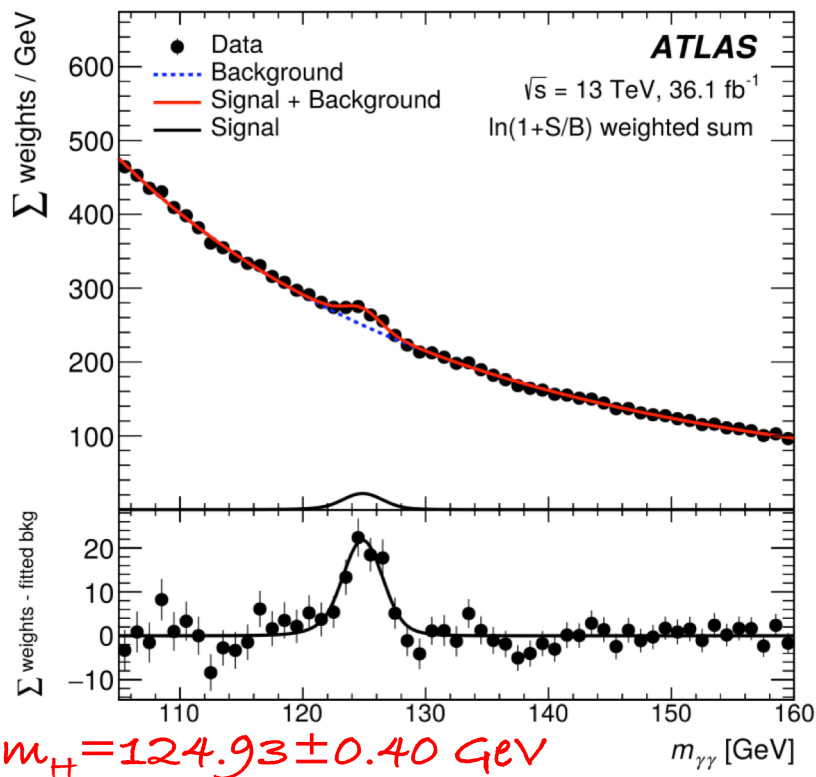
$$\sigma_{VBF} \mathcal{B}_{HWW^*} = 0.50^{+0.24}_{-0.23} \text{ (stat.) } + 0.18 \text{ (syst.) pb and } \mu_{VBF} = 0.62^{+0.37}_{-0.36}$$

ATLAS-CONF-2018-004



HIGGS BOSON MASS

➤ Extraction based on di-photon and 4l invariant mass



Combination with Run 1 : [arXiv:1806.00242](https://arxiv.org/abs/1806.00242)

$m_H = (124.97 \pm 0.24) \text{ GeV} - 1.9 \text{ permille}$



HIGGS BOSON WIDTH

- In contrast of LEP or ILC, at LHC only $\sigma \cdot BR$ can be measured
 - The measurement of Γ_H is extremely hard at LHC
 - Γ_H cannot be inferred from measurements of Higgs boson rates. The SM expectation is 4 MeV.
- **Direct** and indirect **strategies** have been considered:
 - From the on-shell mass peak – limited by mass resolution, ~ 1 GeV [arXiv:1806.00242](https://arxiv.org/abs/1806.00242)
 - From the lifetime



HIGGS BOSON WIDTH

- In contrast of LEP or ILC, at LHC only $\sigma \cdot BR$ can be measured
 - The measurement of Γ_H is extremely hard at LHC
 - Γ_H cannot be inferred from measurements of Higgs boson rates. The SM expectation is 4 MeV.
- Direct and **indirect strategies** have been considered:
 - From couplings
 - **From off-shell to on-shell production**
→ Best proxy to-date (under some assumptions)!



CONSTRAINING THE HIGGS BOSON WIDTH

- Using the relative on-shell and off-shell production, we can indirectly constrain the Higgs boson total width

$$\mu_{\text{off-shell}} = \frac{\sigma_{\text{off-shell}}}{\sigma_{\text{off-shell,SM}}} = k_{g,\text{off-shell}}^2 k_{v,\text{off-shell}}^2$$

$$\mu_{\text{on-shell}} = \frac{\sigma_{\text{on-shell}}}{\sigma_{\text{on-shell,SM}}} = \frac{k_{g,\text{off-shell}}^2 k_{v,\text{off-shell}}^2}{\Gamma_H / \Gamma_{H,SM}}$$

$$\frac{\mu_{\text{off-shell}}}{\mu_{\text{on-shell}}} = \frac{\Gamma_H}{\Gamma_{H,SM}}$$

From an independent analysis

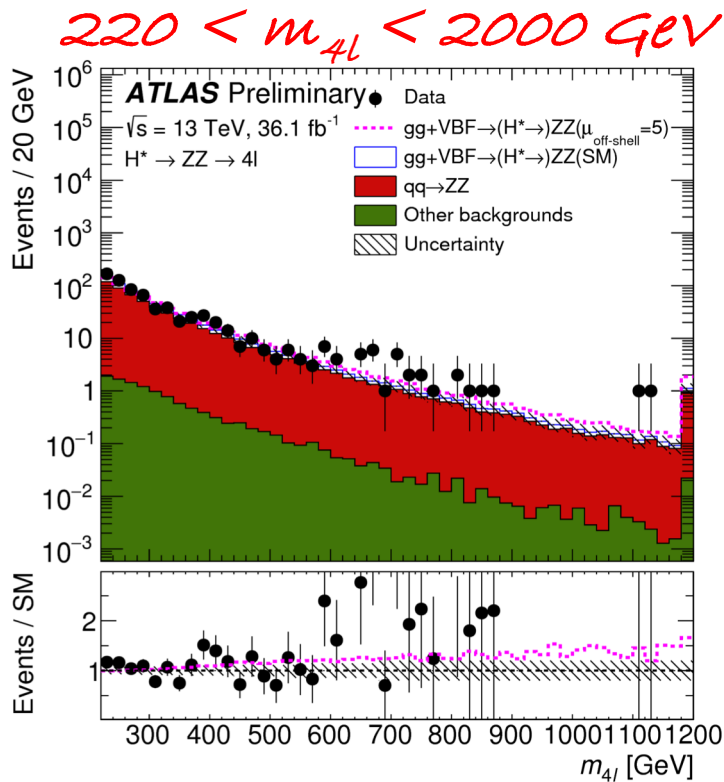


- This strategy is assuming identical on-shell and off-shell couplings
- No new physics alters the Higgs couplings in the off-shell regime

STRATEGY AND RESULTS

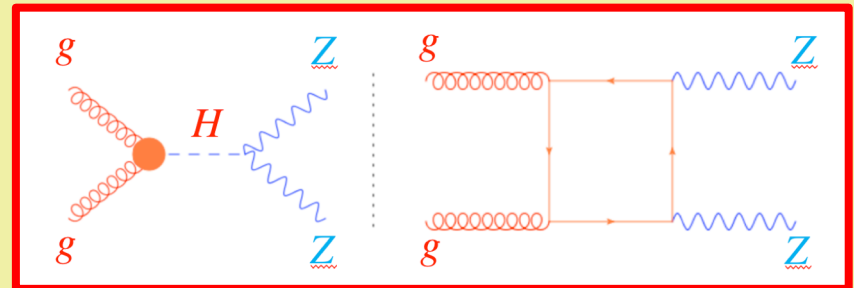
- Two decay channels, $H^* \rightarrow ZZ \rightarrow 4l$ and $H^* \rightarrow ZZ \rightarrow 2l2\nu$
- Discriminants: transverse mass ($2l2\nu$) and Matrix-Element ($4l$)

4l invariant mass



INTERFERENCE

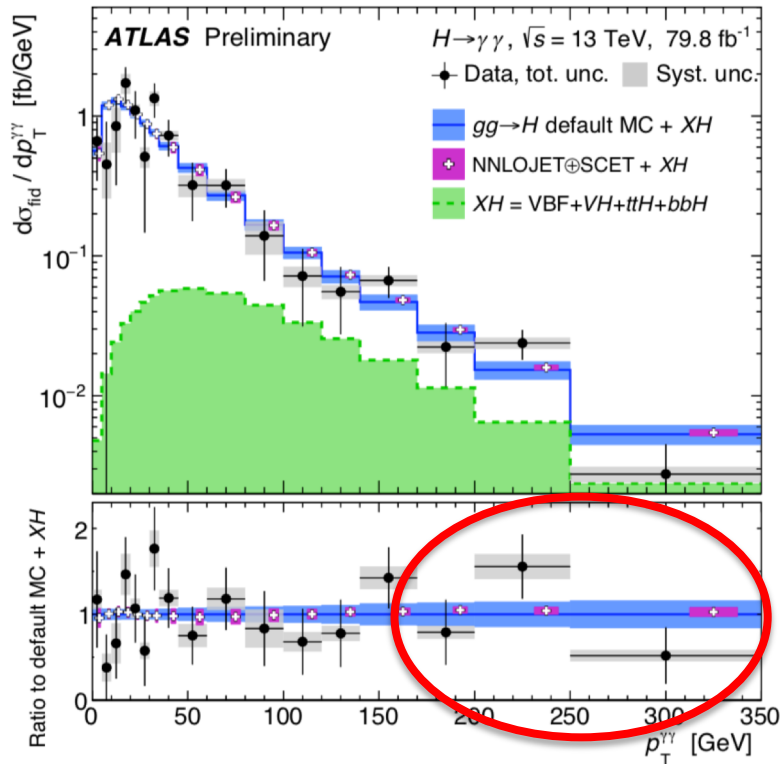
- Production of two Z bosons in fusion of two gluons can occur either directly or through the Higgs bosons
- The two amplitudes interfere destructively in the SM



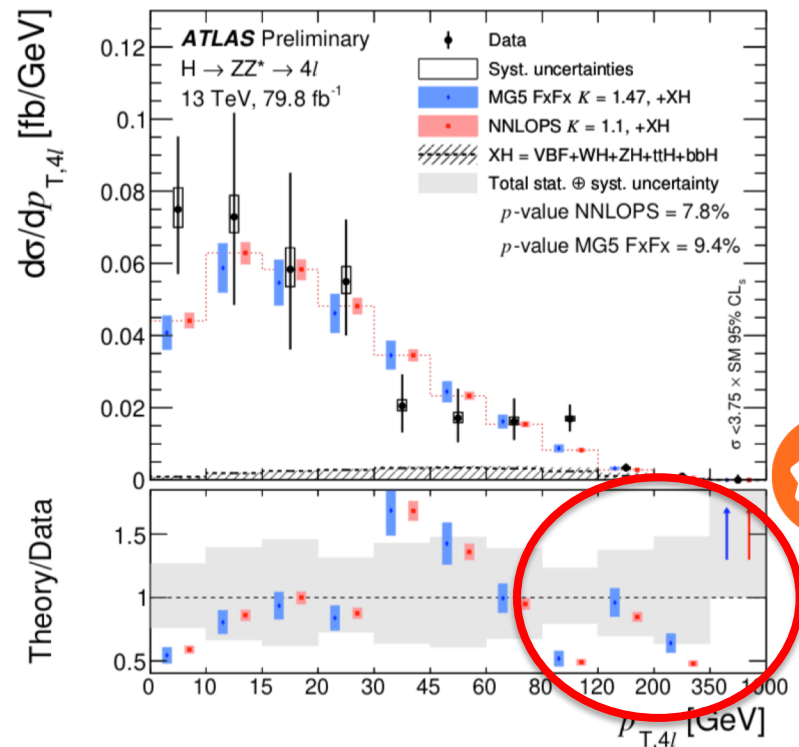
Run 2 result: $\Gamma_H < 14.4 \text{ MeV obs. (15.2 MeV exp.)}$ [arXiv:1808.0119](https://arxiv.org/abs/1808.0119)

DIFFERENTIAL MEASUREMENTS

- New precision on the Higgs- p_T spectrum
- Distribution tails sensitive to new physics



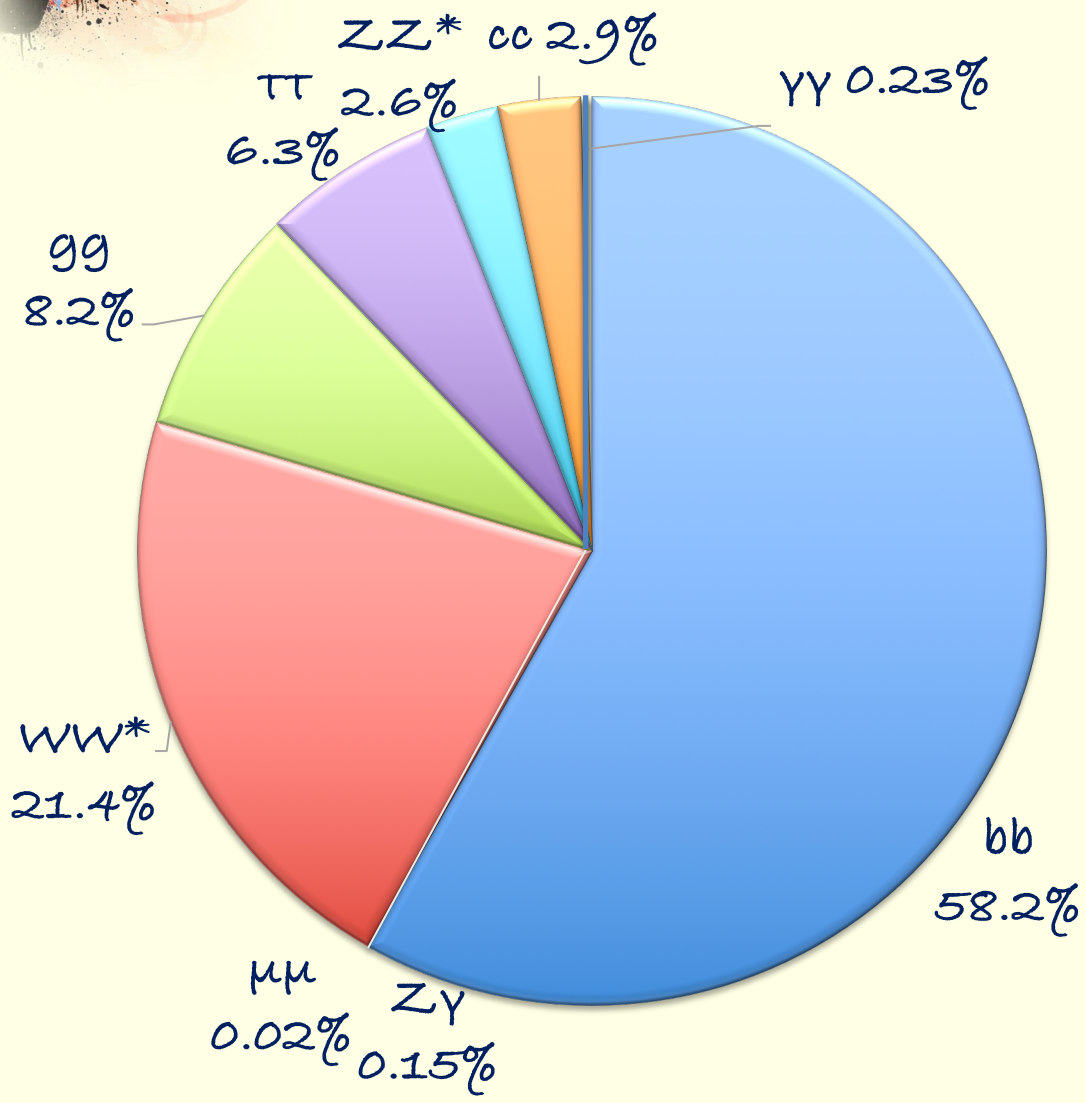
ATLAS-CONF-2018-004




ATLAS-CONF-2018-018

No signs of new physics!


HIGGS TO FERMIONS




$bb, \tau\tau$


 Significant BR

- Allow direct probe to fermions

 Low S/B, challenging measurement

$\mu\mu$

 It allows couplings measurements to 2nd generation

 very small BR



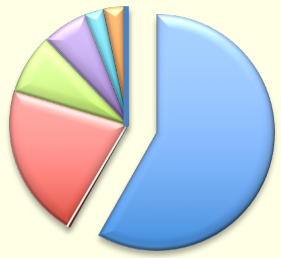
$$H \rightarrow bb$$

- Direct probe to coupling to b-type quarks
- Because of the large BR, the uncertainty on the total decay width (\rightarrow measurement of the absolute couplings) relies on this channel

First Evidence
from TEVATRON in 2012
2.80 Obs. (1.50 Exp.)
3.10 Global
[Phys. Rev. Lett. 109 \(2012\) 071804](#)

LHC
Evidence in 2017
3.60 Obs. (4.00 Exp.) ATLAS
3.80 Obs. (3.80 Exp.) CMS
[JHEP 12 \(2017\) 024](#), [Phys. Lett. B 780 \(2018\) 501](#)

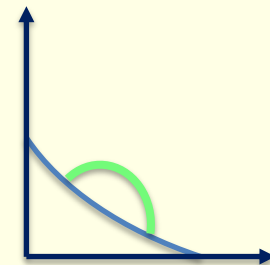
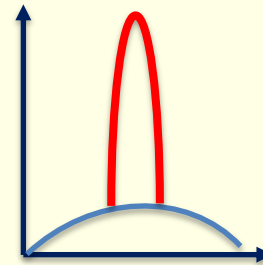
- VH production is the most sensitive mode to search for $H \rightarrow bb$ at LHC



CHALLENGES OF $H \rightarrow bb$

- Let's compare it with one of the discovery channels
 - values in the table before the analyses strategies (discriminants, BDT...)

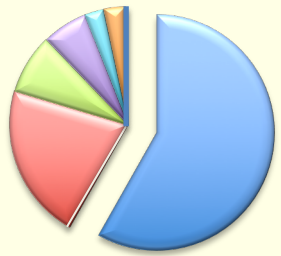
	$H \rightarrow ZZ^* \rightarrow 4l$	$H \rightarrow bb$
Branching Ratio	0.012%	58%
Mass resolution	1%	10%
S/B	2	0.02



→ More is not always better!

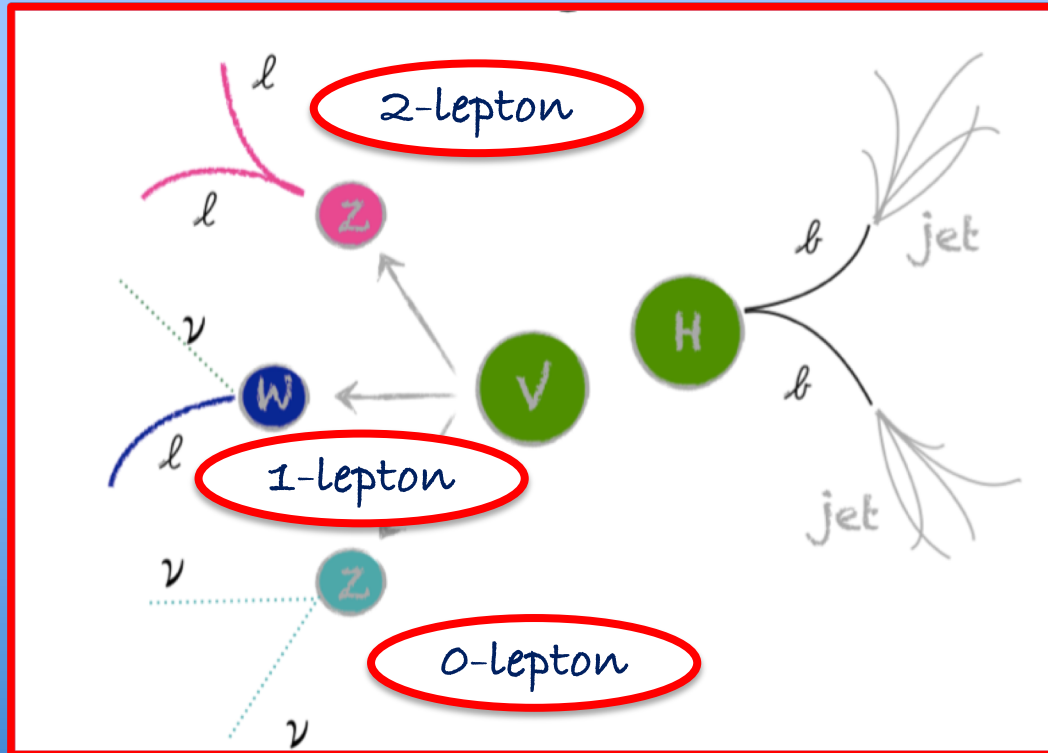
RECIPE FOR THE PERFECT SEARCH

1. Great b-jet identification performance
2. Best possible $m(bb)$ resolution
3. Leverage on all the event information



VH($H \rightarrow bb$) TOPOLOGY

SIGNAL



CATEGORIZATION

- 3 decay channels
- Classification based on the number of leptons from $V=W/Z$ bosons



VH($H \rightarrow bb$) TOPOLOGY

SIGNAL



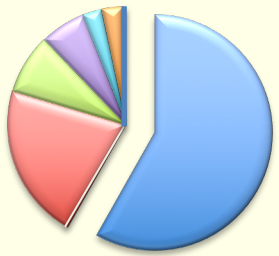
B-taggers:
selecting
the 2b-jets
we need



b-jet efficiency
70%

SELECTION

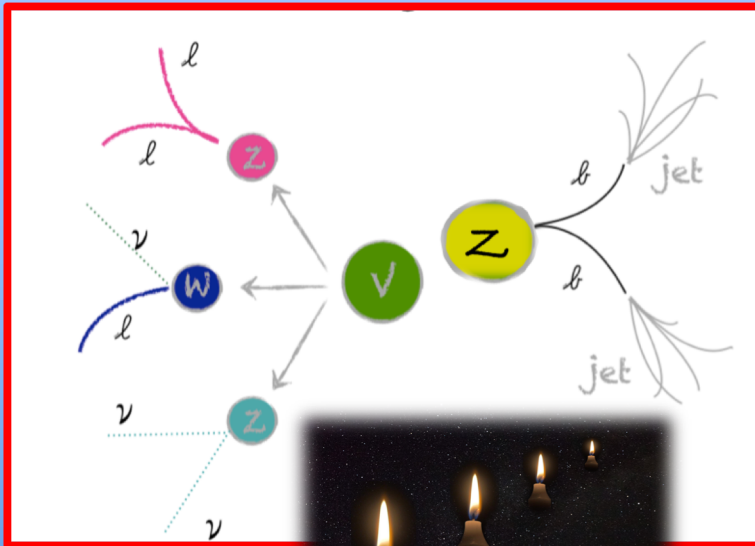
- 2 b-jets from Higgs + leptons from ν decay
- $p_T(\nu) > 75$ or 150 GeV
- b-jets identification based on multivariate techniques exploiting the long b-hadron lifetime



VH($H \rightarrow bb$) TOPOLOGY

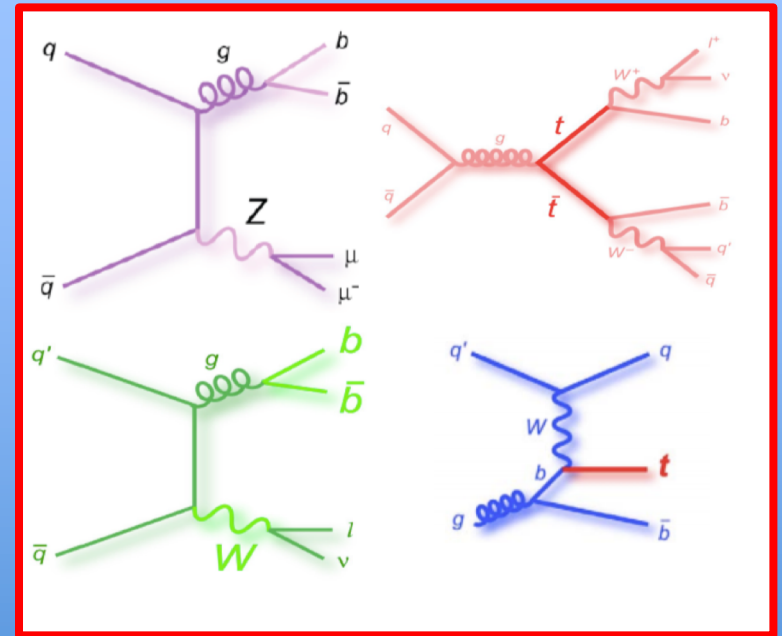
BACKGROUND

DIBOSON



VZ, a perfect standard candle to validate the analysis strategy

OTHERS IRREDUCIBLE



Normalization from data (CR), shapes from MC



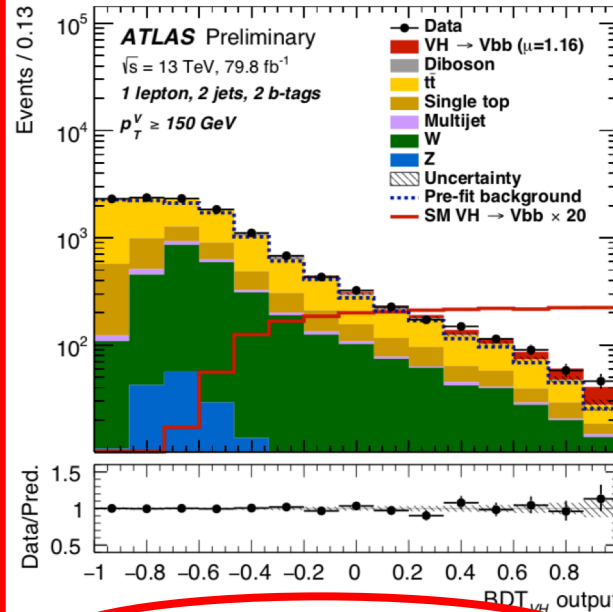
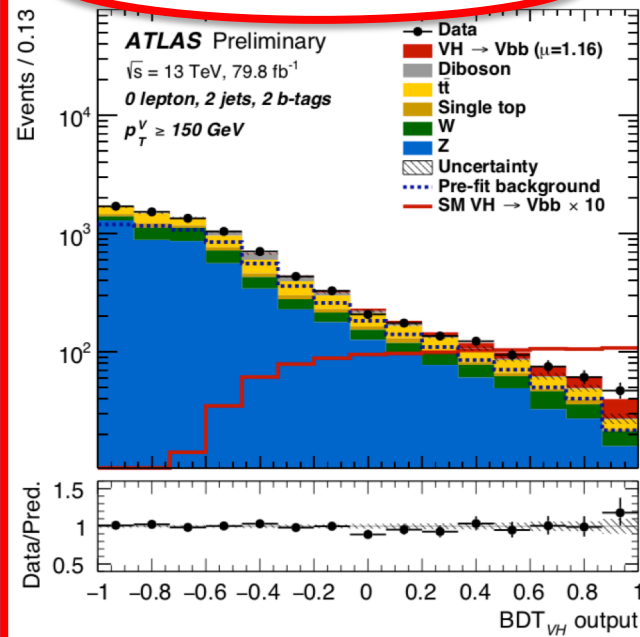
VH($H \rightarrow bb$) STRATEGY AND RUN-2 RESULTS

➤ Main discriminants m_{bb} , $p_T(\nu)$ and ΔR_{bb} combined into a Boosted Decision Tree (BDT)

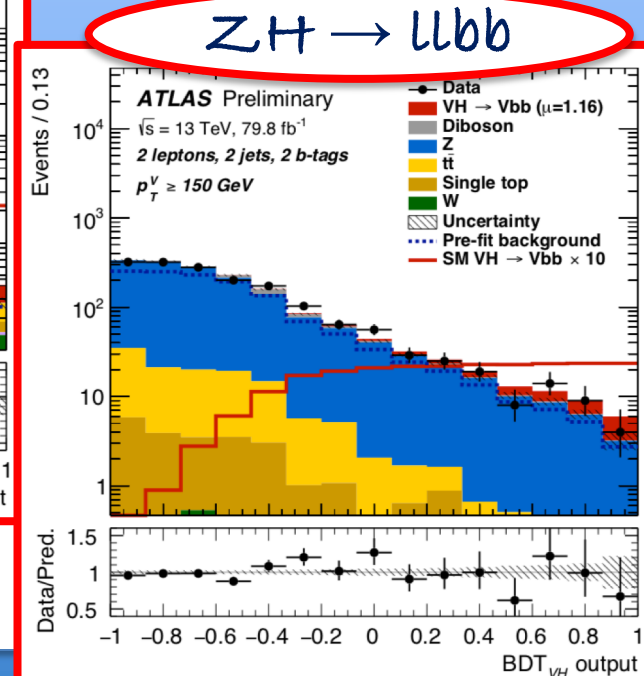
combined into a

NEW

$ZH \rightarrow \nu b\bar{b}$



$WH \rightarrow l\nu b\bar{b}$



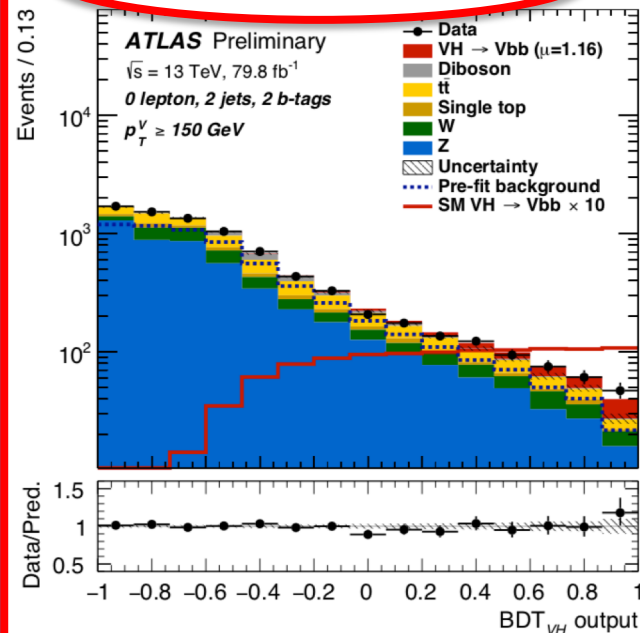


VH(H → bb) STRATEGY AND RUN-2 RESULTS

- Main discriminants m_{bb} , $p_T(V)$ and ΔR_{bb} combined in a Boosted Decision Tree (BDT)

NEW

ZH → vbb



Events / 0.13

Data/Pred.

BDT_{VH} output

RESULTS

- Fit result with 79.8 fb^{-1} (Run 2)
 $\mu = \sigma_{\text{meas}} / \sigma_{\text{SM}} = 1.16^{+0.27}_{-0.25}$
 significance:
 4.90 obs. (4.30 exp.)
- combination with Run 1
 $\mu = 0.98 \pm 0.14$ (stat) $^{+0.17}_{-0.16}$ (syst)
 significance:
 4.90 obs. (5.10 exp.)

Data/Pred.

BDT_{VH} output



VH and $H \rightarrow bb$ RESULTS

$H \rightarrow bb$ combination

Run-1 and Run-2 analyses:

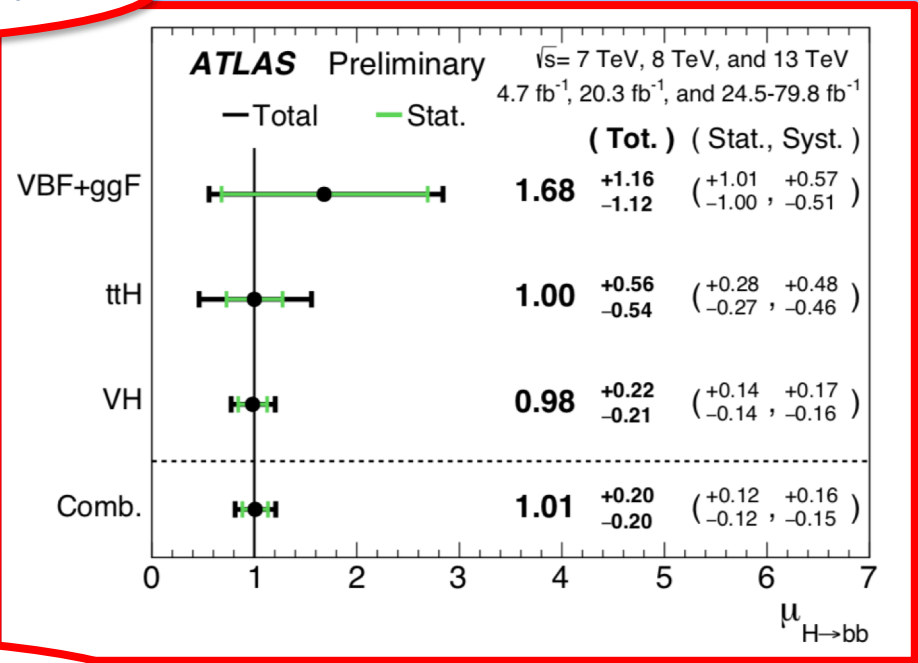
- $VH, H \rightarrow bb$
- $VBF(+ggF), H \rightarrow bb$
- $ttH, H \rightarrow bb$



Significance: [arXiv:1808.08238](https://arxiv.org/abs/1808.08238)

5.40 obs. (5.50 exp.)

Observation of $H \rightarrow bb$!



VH combination

Run-2 analyses:

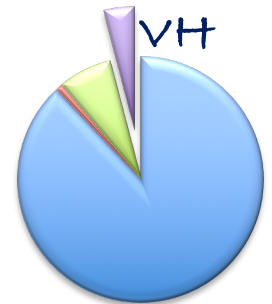
- $VH, H \rightarrow bb$
- $VH, H \rightarrow \gamma\gamma$
- $VH, H \rightarrow ZZ^*$

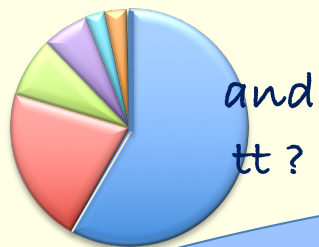


Significance: [arXiv:1808.08238](https://arxiv.org/abs/1808.08238)

5.30 obs. (4.80 exp.)

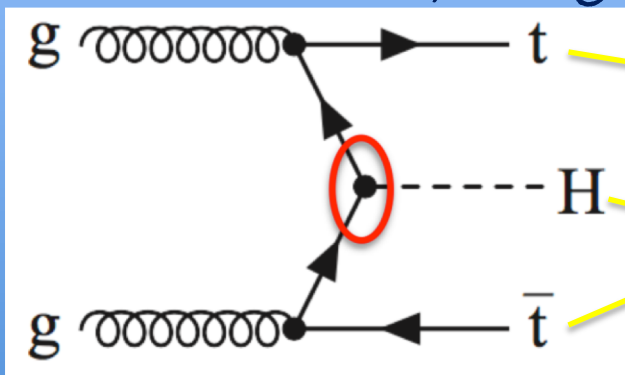
Observation of VH production!





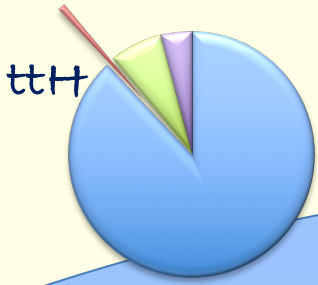
COUPLING TO TOP-QUARK

- The largest coupling to top quark
 - can only be studied directly in $t\bar{t}H$ production mode
 - can be studied indirectly using ggF, being dominated by the top quark in the SM
- $t\bar{t}H$: high multiplicity of final state objects → complex analysis



$2\text{ }b\text{-jets} + \begin{cases} l\nu l\nu \\ l\nu + 2\text{jets} \\ 4\text{jets} \end{cases}$
 jets, leptons, MET, γ

- Several categories analysed with multivariate algorithms

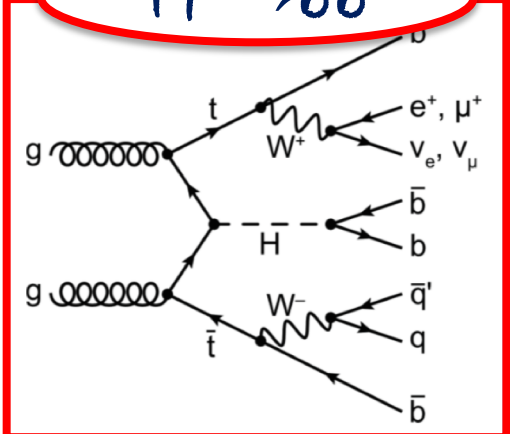
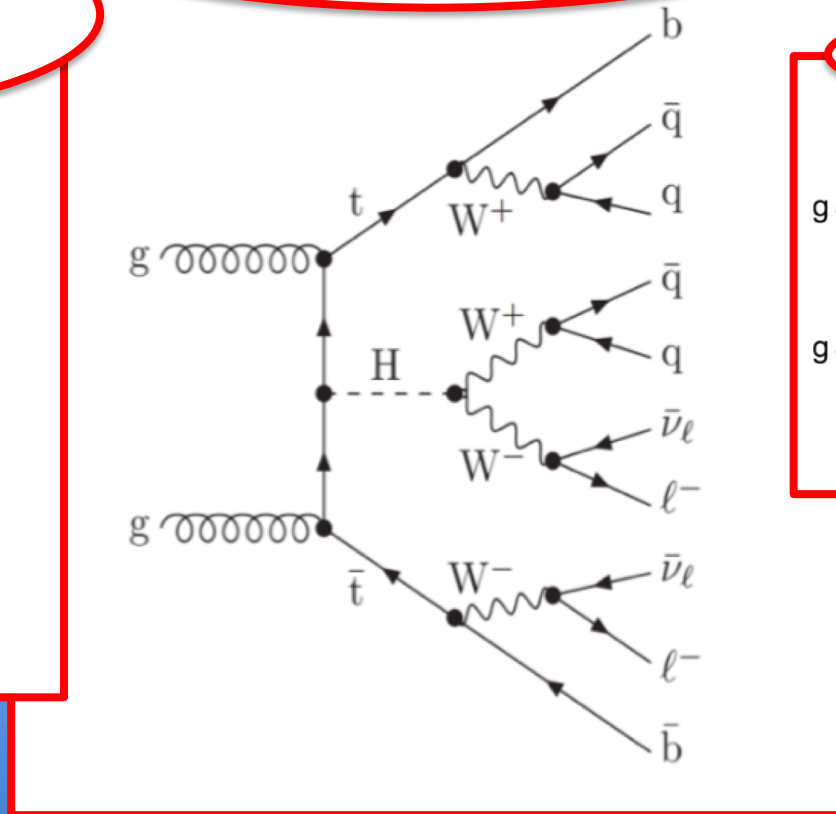
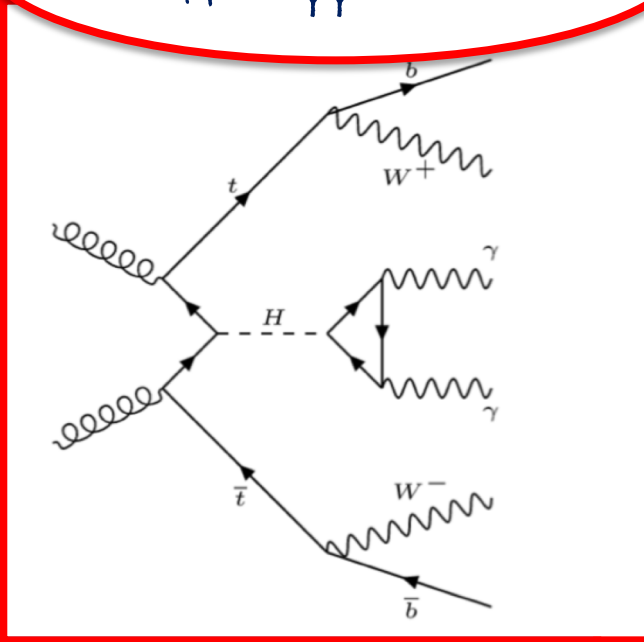


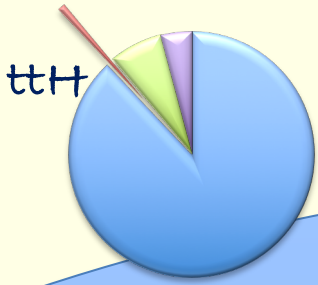
ttH CATEGORIES

- $H \rightarrow WW^* \rightarrow l\nu l\nu$
- $H \rightarrow \tau\tau$ and
- $H \rightarrow ZZ^*$ without $ZZ^* \rightarrow 4l$
Multi-leptons

- $H \rightarrow ZZ^* \rightarrow 4l$
- $H \rightarrow \gamma\gamma$

H → bb





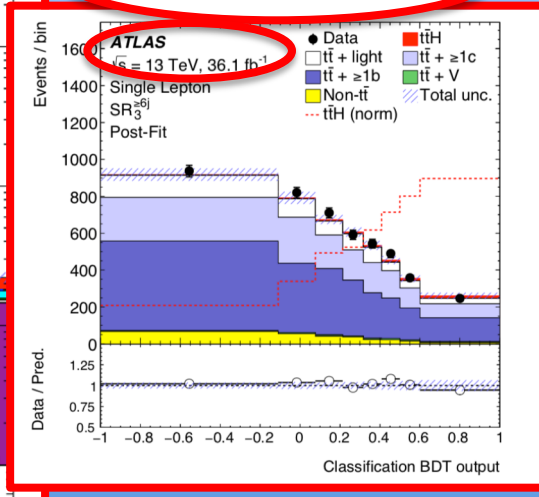
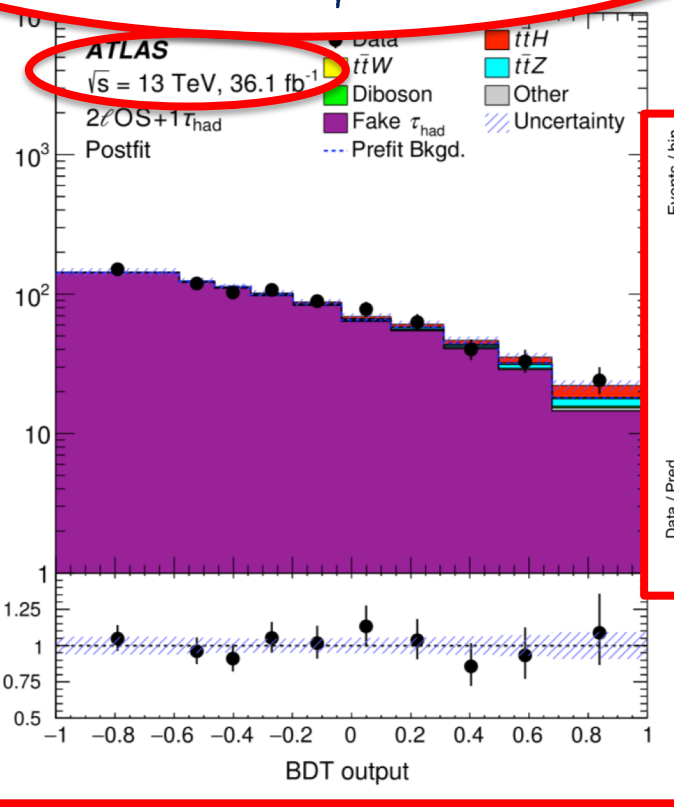
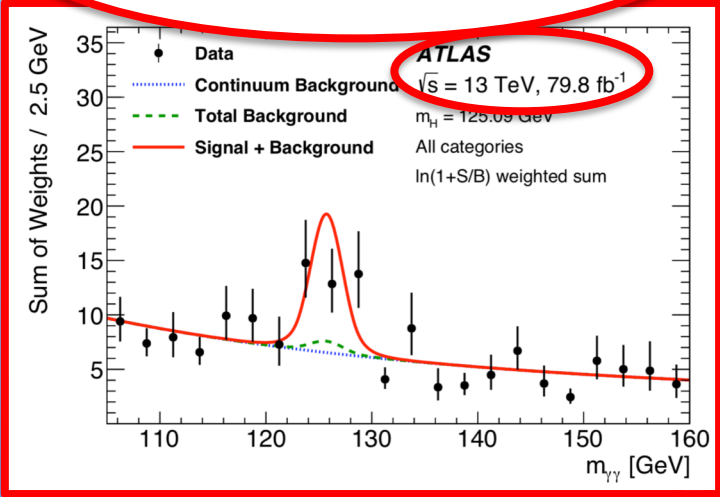
ttH CATEGORIES

- $H \rightarrow WW^* \rightarrow l\nu l\nu$
- $H \rightarrow \tau\tau$ and
- $H \rightarrow ZZ^*$ without $ZZ^* \rightarrow 4l$

Multi-leptons

- $H \rightarrow ZZ^* \rightarrow 4l$
- $H \rightarrow \gamma\gamma$

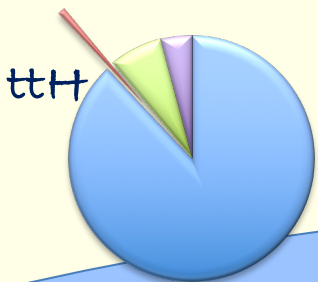
$H \rightarrow bb$



$tt(H \rightarrow \gamma\gamma)$ 4.10 obs. (3.70 exp.)
 $tt(H \rightarrow ZZ^*)$ 0.00 obs. (1.20 exp.)

$tt(H \rightarrow bb)$ 1.40 obs. (1.60 exp.)

$tt(H \rightarrow \text{multi-lepton})$ 4.10 obs. (2.80 exp.)



ttH RESULTS

ttH combination

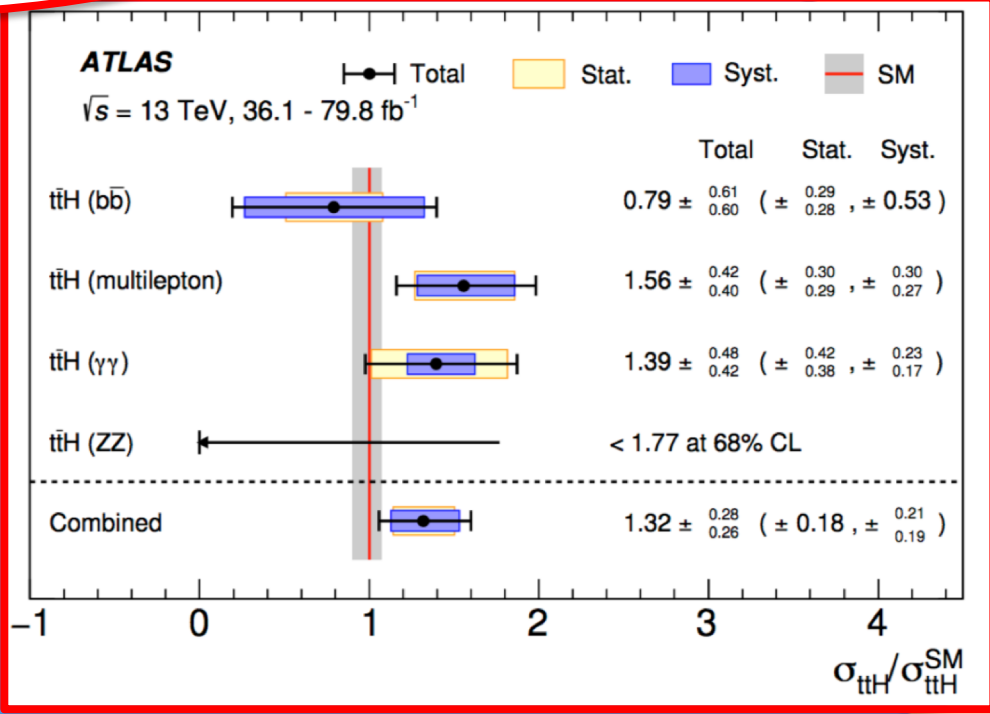
Run-2 analyses:

- ttH, H → bb
- ttH, multilepton
- ttH, H → γγ
- ttH, H → ZZ*

Significance: [arXiv:1806.00425](https://arxiv.org/abs/1806.00425)

5.8σ obs. (4.9σ exp.)

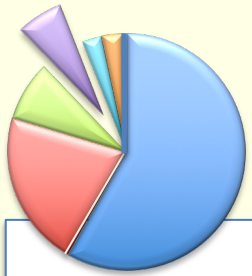
Observation of ttH production!



Run-2 + Run-1 analyses: Significance 6.3σ obs. (5.1σ exp.)

Cross section: $\sigma_{ttH} = 670 \pm 90$ (stat.) $^{+110}_{-100}$ (syst.) fb

with $\sigma_{ttH, SM} = 507^{+35}_{-50}$ fb [arXiv:1806.00425](https://arxiv.org/abs/1806.00425)

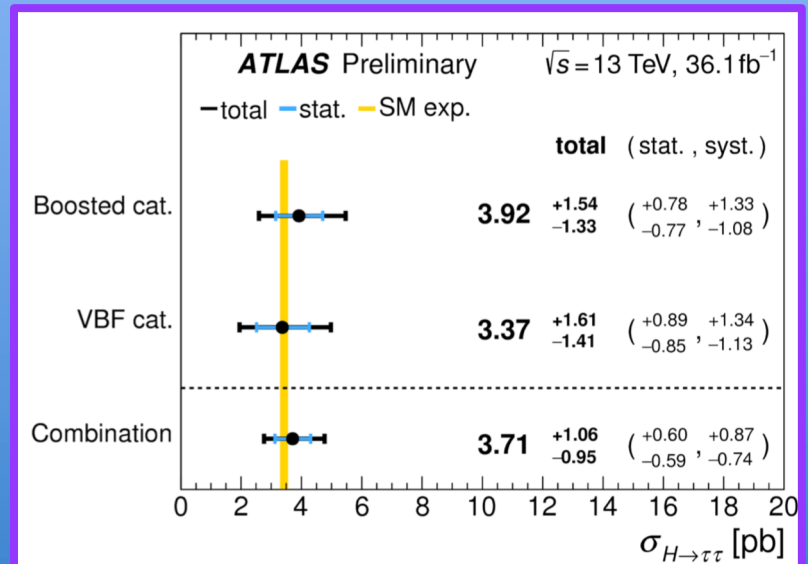
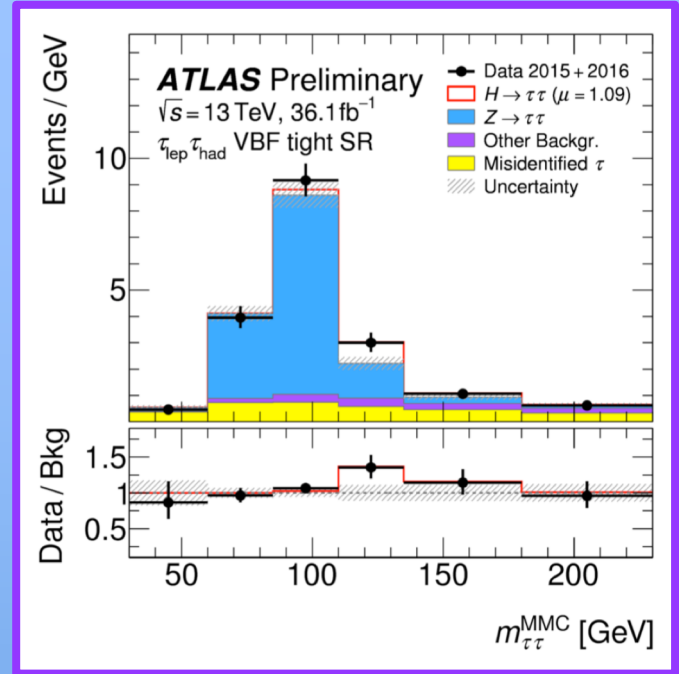


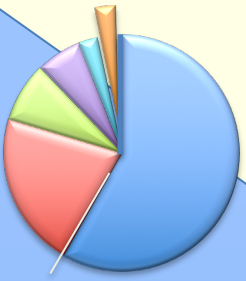
$$H \rightarrow \tau\tau$$

- Analyses use all the τ leptonic and hadronic decay modes: $\tau \rightarrow \text{hadrons} + \nu$ and $\tau \rightarrow l\nu$
- Categories considered: VBF and ggF in the “boosted regime”
- Main background $Z \rightarrow \tau\tau$
 - Normalization from data and shapes from MC
- Main discriminant $m_{\tau\tau}$

*Significance Run 1 + Run 2:
6.40 obs. (5.40 exp.)*

ATLAS-CONF-2018-021

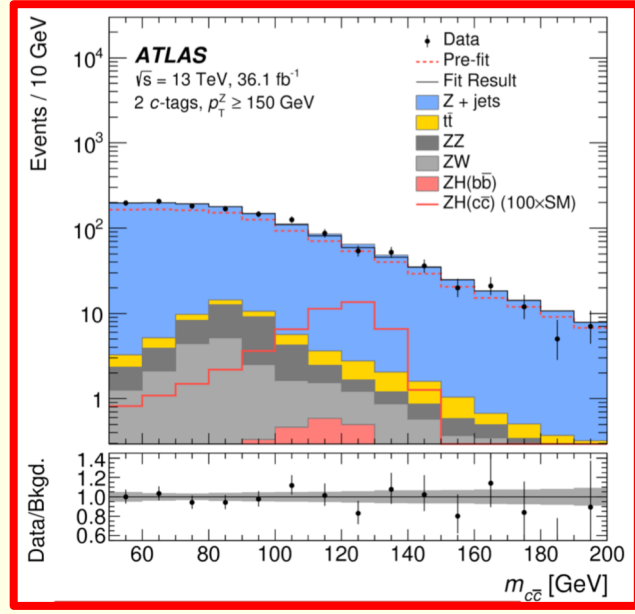
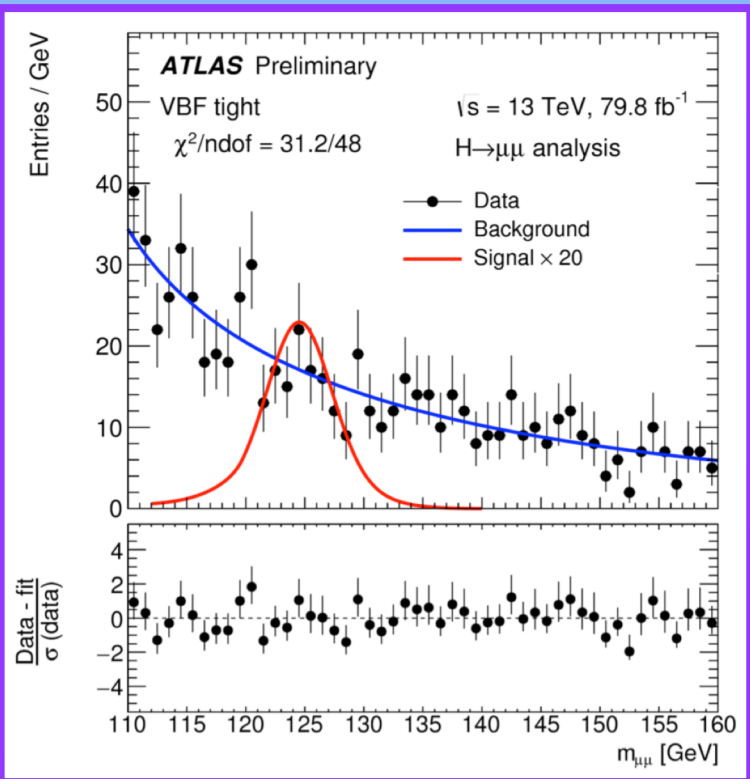




2nd GENERATION: $H \rightarrow cc$ and $H \rightarrow \mu\mu$

$H \rightarrow \mu\mu$
 Selection: events with 2 isolated and opposite-sign muons

$H \rightarrow cc$
 Selection: similar to $H \rightarrow bb$, but based on specific c-tagging techniques



[arXiv:1802.04329](https://arxiv.org/abs/1802.04329)

[ATLAS-CONF-2018-026](#)

Run 2 (36.1 fb⁻¹)

Run 2 (80 fb⁻¹) getting close to SM sensitivity

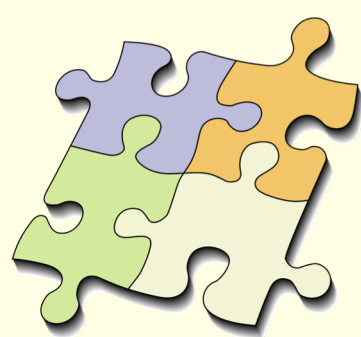
$\mu_{cc} < 110$

w.r.t. SM

exp.

$$\mu_{\mu\mu} = 0.1^{+1.0}_{-1.1}$$

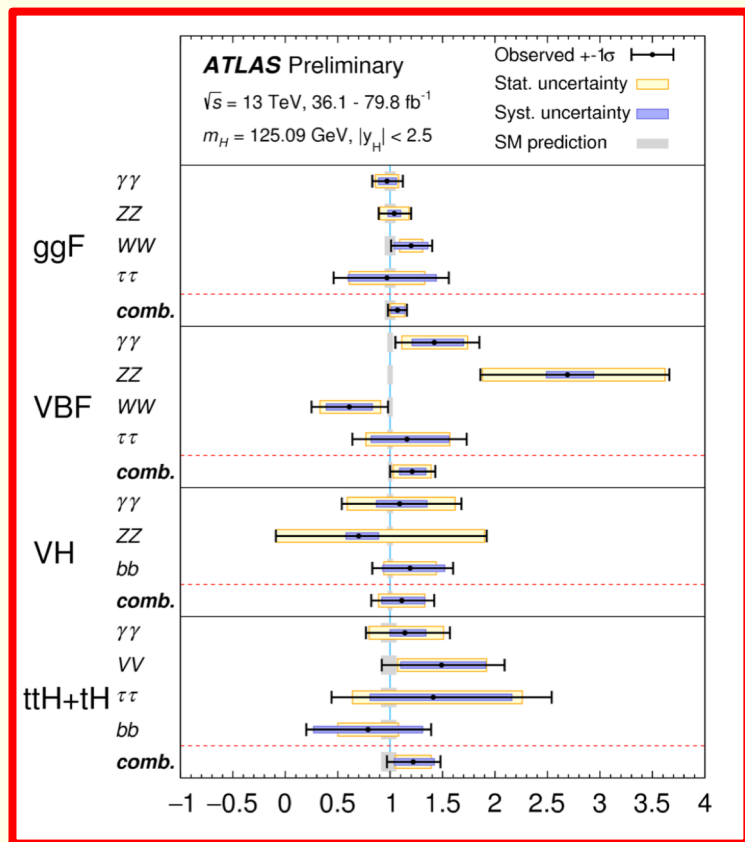
$$\mu_{\mu\mu} < 2.1 \text{ obs. (2.0 exp.)}$$



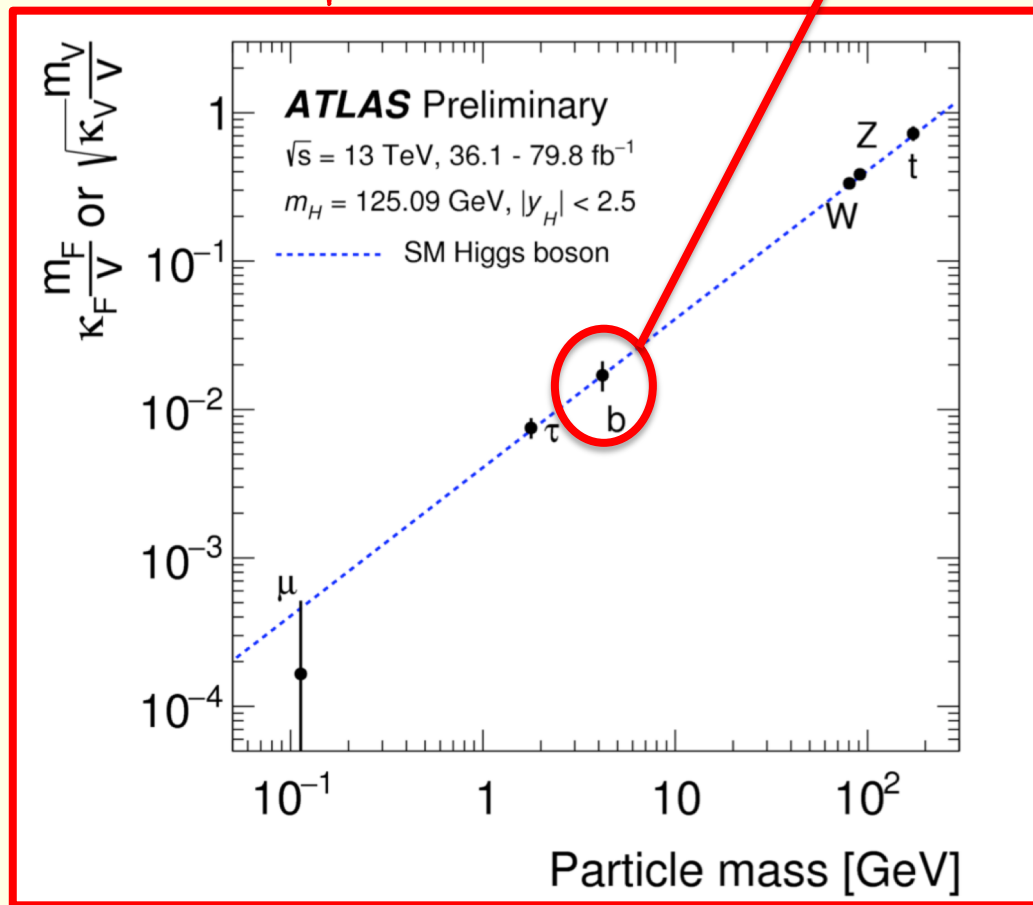
COMBINING ALL THE MEASUREMENTS

To be updated with the Hbb observation

PRODUCTION



COUPLINGS



Global signal strength: $\mu_H = 1.13^{+0.09}_{-0.08}$ [ATLAS-CONF-2018-031](#)



CONCLUSIONS

- The LHC Run 2 has improved the precision of Higgs physics and allowed new couplings measurements
 - Direct observation for all the main production modes: ggF, VBF, VH and ttH
 - All five decay modes foreseen for light Higgs (WW^* , ZZ^* , $\gamma\gamma$, bb and $\tau\tau$) observed
 - New precision for the mass and better (indirect) constraints for the width:
 $m_H = (124.97 \pm 0.24) \text{ GeV}$ and $\Gamma_H < 14.4 \text{ MeV}$
 - Couplings to muons look promising
 - SM consistency reduces BSM possibilities

We have shed light on 5 orixas-couplings ...
more to come!



Lagoa dos Orixas, Dique do Tororo, Salvador

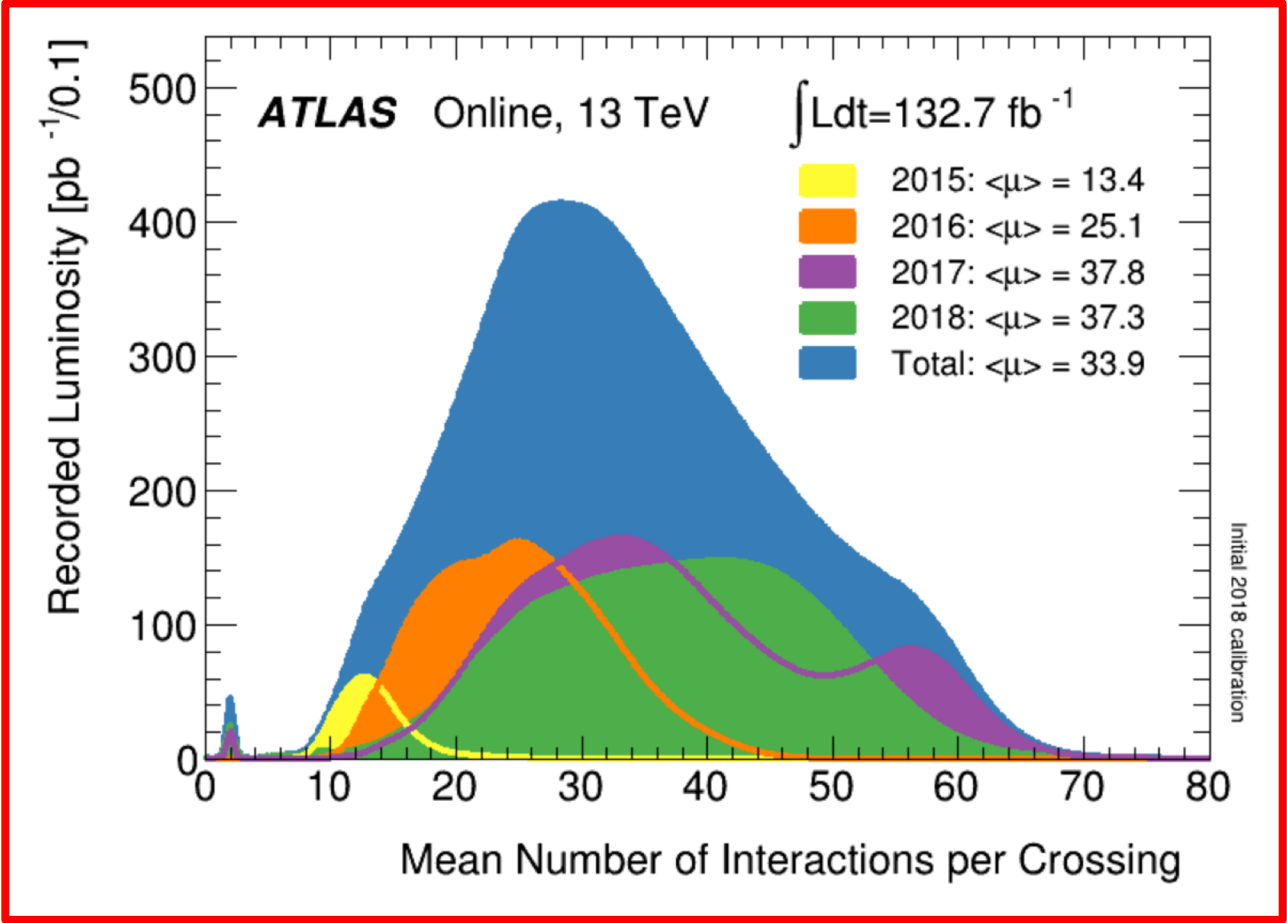
MUITO OBRIGADO!



BACK-UP



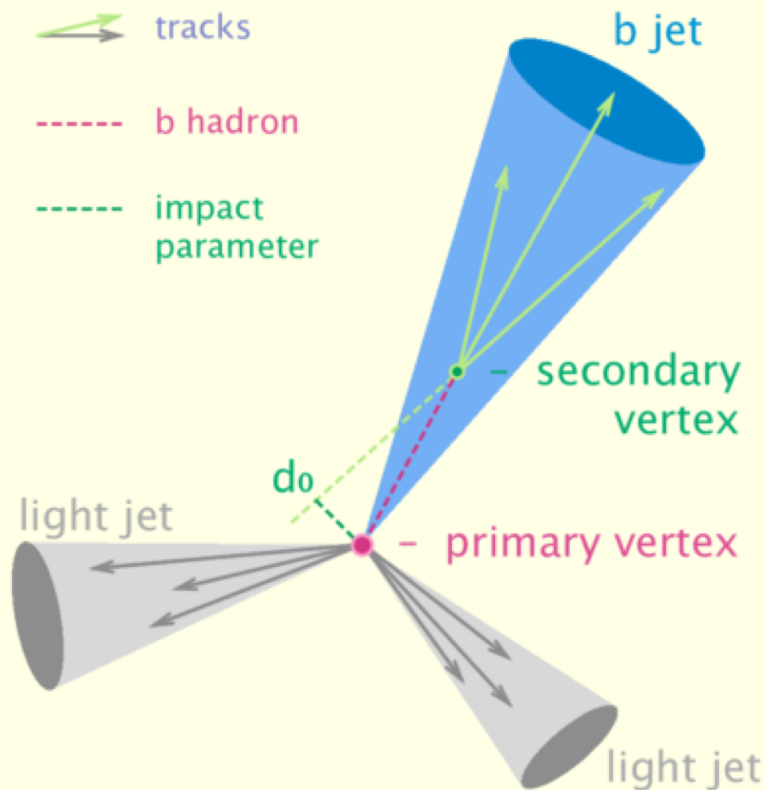
PILEUP





b-TAGGING

- Specific algorithms developed to identify (tag) jets produced by b-quarks
- Separation of b-jets from light (u, d, s, g) jets using specific b-hadrons properties:
 - Long b-hadron lifetime $\tau \sim 1.5$ ps
→ 20 GeV b-hadrons decays after ~ 2 mm
 - Search for tracks or vertexes displaced w.r.t. primary vertex and large impact parameters for displaced tracks
 - Search leptons from semileptonic b decays, with large transverse momentum w.r.t. jet axis

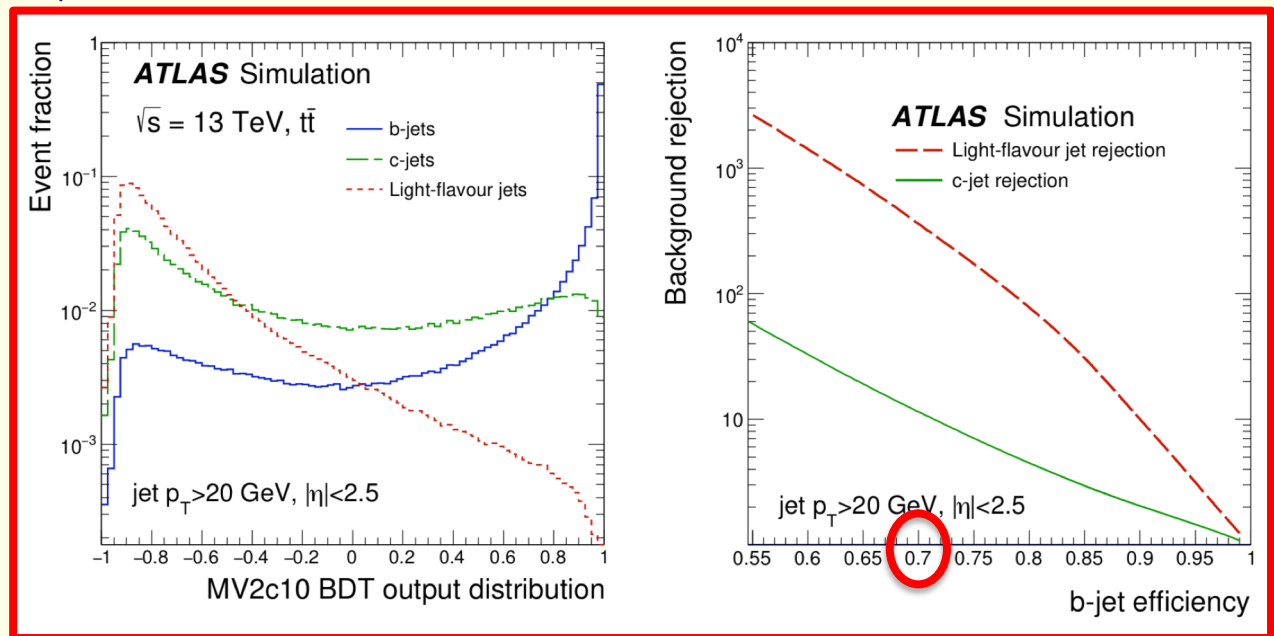




b-TAGGING PERFORMANCES

- A new multivariate b-tagging algorithm, MV2c10, was developed for Run2 and it's based on a BDT: [arXiv:1805.01845](https://arxiv.org/abs/1805.01845)
 - BDT trained on a events from a tt sample
 - Selection tuned to produce an average efficiency of 70% for b-jets
 - Jets are tagged as containing b-hadrons by requiring a large MV2c10 BDT output value

- ❖ b-jet purity is:
 - $e\mu$ SAMPLE**
 - 72% - 2j sample
 - 53% - 3j sample
 - $ee/\mu\mu$ SAMPLE**
 - 62% - 2j sample
 - 48% - 3j sample





VH(Hbb) SYSTEMATIC UNCERTAINTIES

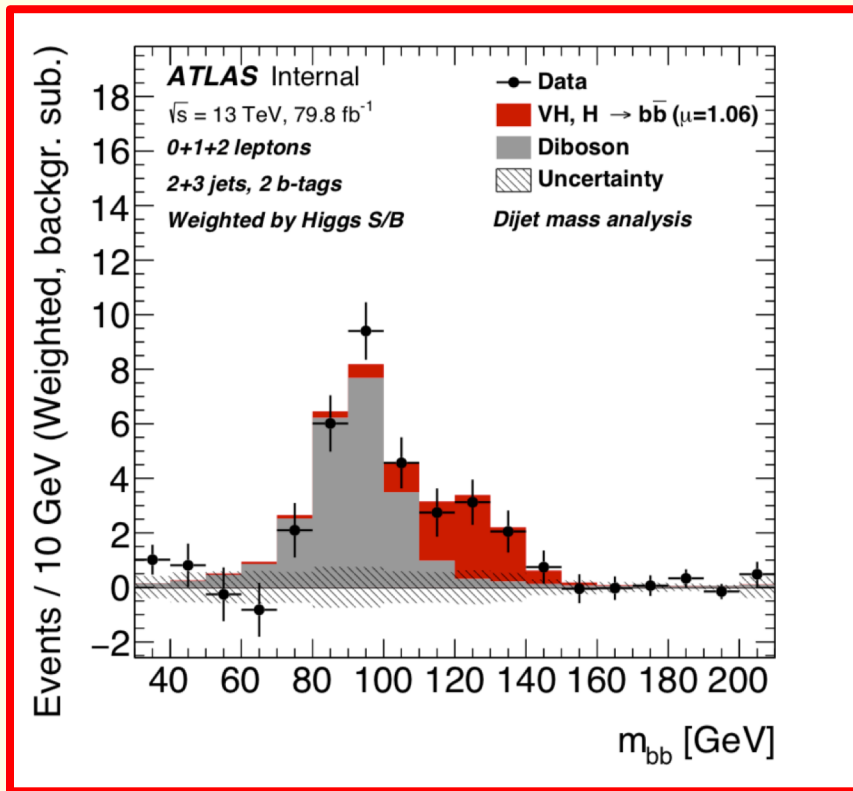
- Breakdown of the contributions to the uncertainty in $\mu_{\text{VH(Hbb)}}$
 - Leading contributions:
 - Modelling of the signal
 - Limited size of simulated samples
 - Modelling of background
 - b-tagging efficiency

Source of uncertainty	σ_μ	
Total	0.259	
Statistical	0.161	
Systematic	0.203	
Experimental uncertainties		
Jets	0.035	
E_T^{miss}	0.014	
Leptons	0.009	
b-tagging	b-jets	0.061
	c-jets	0.042
	light jets	0.009
	extrapolation	0.008
Pile-up	0.007	
Luminosity	0.023	
Theoretical and modelling uncertainties		
Signal	0.094	
Floating normalisations	0.035	
Z + jets	0.055	
W + jets	0.060	
$t\bar{t}$	0.050	
Single top quark	0.028	
Diboson	0.054	
Multijet	0.005	
MC statistical	0.070	



VH(Hbb): VALIDATION

- A validation of the multivariate analysis is performed using the m_{bb} variable as a discriminant



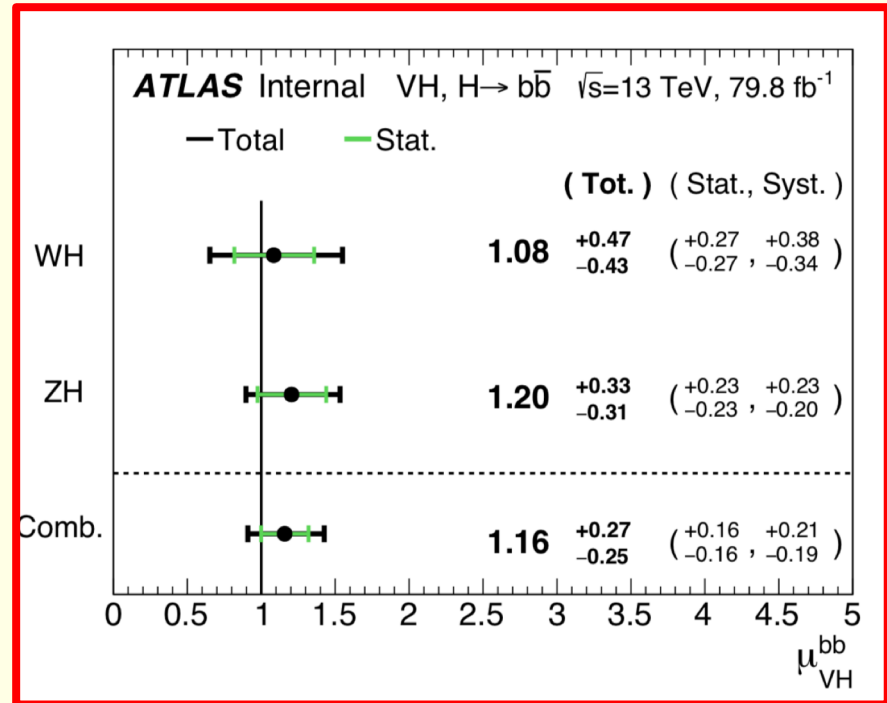
- RESULTS**
- Fit result with 79.8 fb^{-1} (Run 2)
 $\mu = 1.06^{+0.36}_{-0.33}$
Good agreement with the multivariate analysis!
Significance:
 3.6σ obs. (3.5σ exp.)

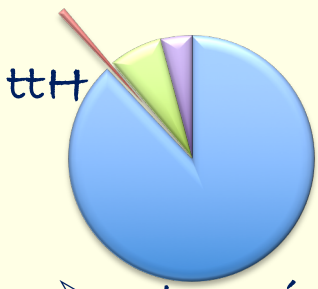
VH(Hbb) INDIVIDUAL RESULTS

➤ Individual-category results: 80% compatibility

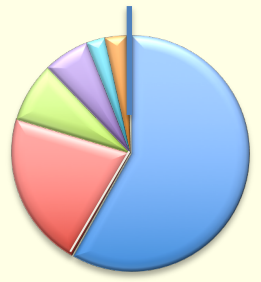
➤ WH and ZH results: 84% compatibility

Signal strength parameter	Signal strength
0-lepton	$1.04^{+0.34}_{-0.32}$
1-lepton	$1.09^{+0.46}_{-0.42}$
2-lepton	$1.38^{+0.46}_{-0.42}$
VH, H → b \bar{b} combination	$1.16^{+0.27}_{-0.25}$

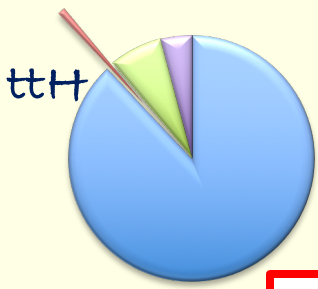




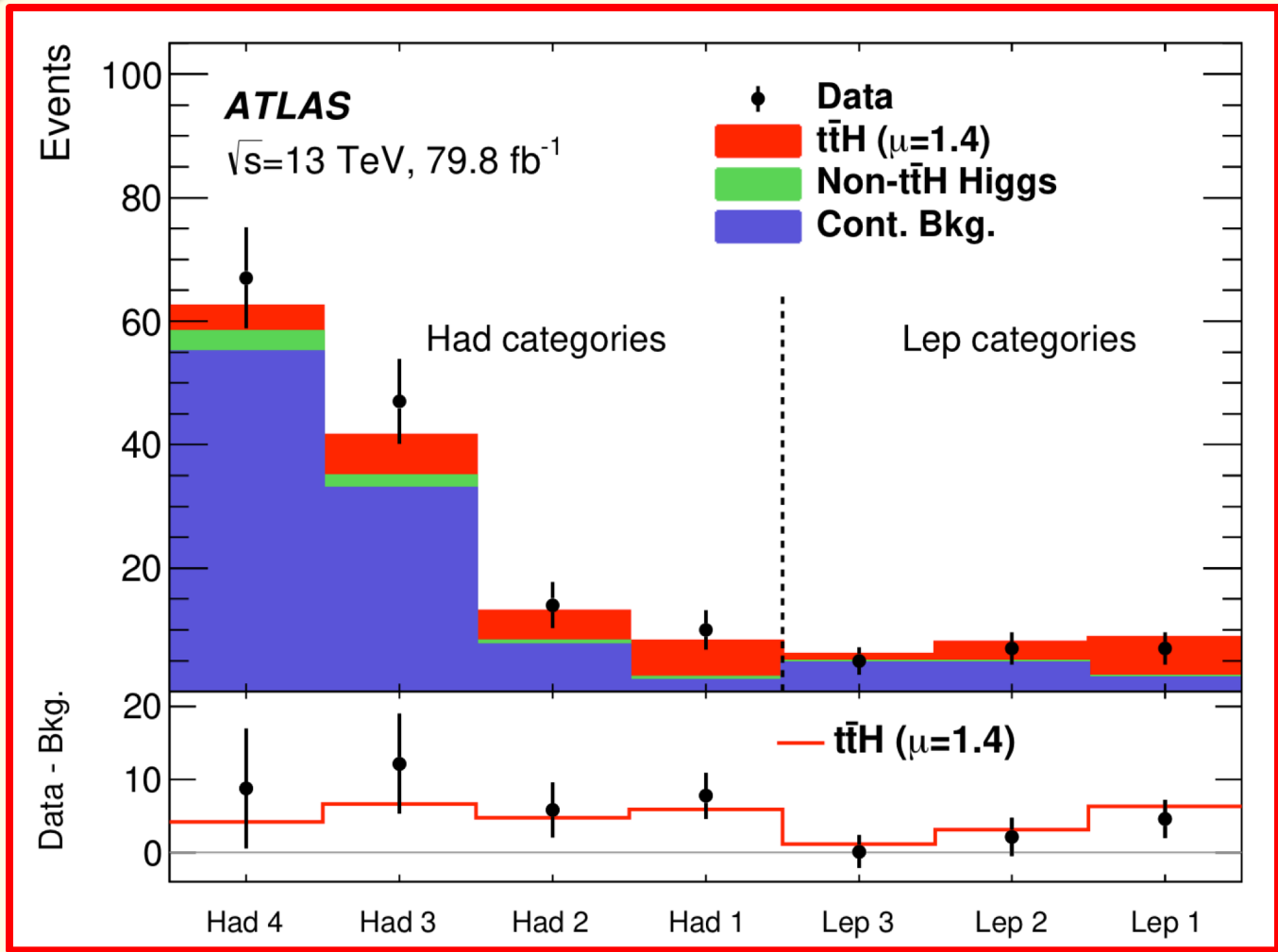
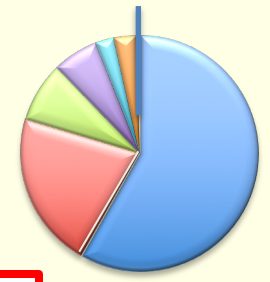
$$ttH \rightarrow \gamma\gamma$$

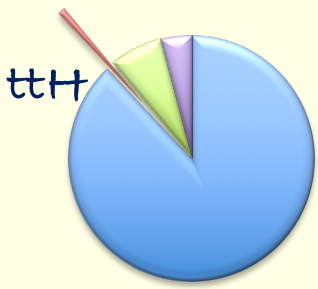


- Clear signature: 2 isolated photons + at least one b-tagged jet
- 2 SRS {
 - Had: at least 2 additional jets and no isolated leptons
Enriched in hadronic top-decays
 - Lep: at least one isolated leptons
Enriched in semi-leptonic top-decays
- Background: non-resonant diphoton processes, including $tt+2\gamma$. Also non- ttH processes, i.e. tH and ggF in 'Had' and tH and VH in 'Lep'
- Strategy: 2 dedicated BDTs whose input variables are: $[p_T, \eta, \phi$ and E of the jets (leptons)] + $(p_T/m_{\gamma\gamma}, MET$ info, η and ϕ of the photons)
 - BDT bins chosen to optimize the expected ttH sensitivity
 - In each BDT bin the ttH signal yield is extracted by performing an unbinned likelihood fit to the diphoton invariant mass

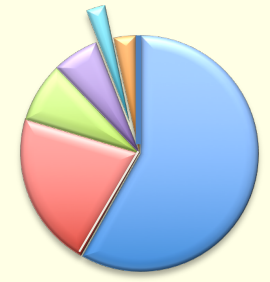


$$ttH \quad H \rightarrow \gamma\gamma$$

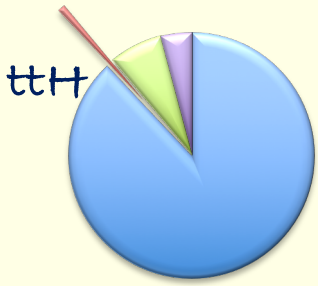




$$ttH \rightarrow ZZ^* \rightarrow 4l$$



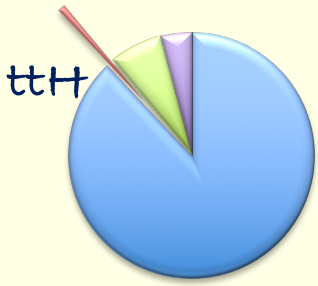
- Clear signature: at least 4 isolated leptons (sfos) + at least one b-tagged jet
- 2 SRs: 'Had' (at least 3 additional jets and no isolated leptons) and 'Lep' (at least one additional isolated lepton and one jet) similarly to $ttH \rightarrow \gamma\gamma$
- Background: ttW , ttZ and non- ttH processes (ggF and tH for 'Had' and tH for 'Lep')
- Strategy: BDT in 'Had' whose input variables are kinematic variables of jets and leptons and a single region for 'Lep'
 - 2 BDT bins chosen to optimize the expected ttH sensitivity in the 'Had'
 - The ttH observed signal and expected background yields extracted in the 'Had' BDT bins and in the 'Lep' regions are used as input for a likelihood fit to the diphoton invariant mass to extract the ttH yield.



ttH COMBINATION

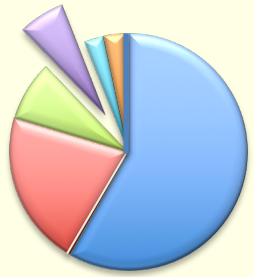
- The final combination is performed using the profile likelihood method based on simultaneous fits to the signal regions and CRs of the individual analyses
- Syst. uncertainties effects taken into account through nuisance parameters
- Non-ttH processes fixed to SM predictions in the final fit

Analysis	Integrated luminosity [fb^{-1}]	$t\bar{t}H$ cross section [fb]	Obs. sign.	Exp. sign.
$H \rightarrow \gamma\gamma$	79.8	710^{+210}_{-190} (stat.) $^{+120}_{-90}$ (syst.)	4.1σ	3.7σ
$H \rightarrow \text{multilepton}$	36.1	790 ± 150 (stat.) $^{+150}_{-140}$ (syst.)	4.1σ	2.8σ
$H \rightarrow b\bar{b}$	36.1	400^{+150}_{-140} (stat.) ± 270 (syst.)	1.4σ	1.6σ
$H \rightarrow ZZ^* \rightarrow 4\ell$	79.8	<900 (68% CL)	0σ	1.2σ
Combined (13 TeV)	36.1–79.8	670 ± 90 (stat.) $^{+110}_{-100}$ (syst.)	5.8σ	4.9σ
Combined (7, 8, 13 TeV)	4.5, 20.3, 36.1–79.8	–	6.3σ	5.1σ



ttH SYSTEMATIC UNCERTAINTIES

Uncertainty source	$\Delta\sigma_{t\bar{t}H}/\sigma_{t\bar{t}H}$ [%]
Theory uncertainties (modelling)	11.9
<i>t\bar{t}</i> + heavy flavour	9.9
<i>t\bar{t}H</i>	6.0
Non- <i>t\bar{t}H</i> Higgs boson production	1.5
Other background processes	2.2
Experimental uncertainties	9.3
Fake leptons	5.2
Jets, E_T^{miss}	4.9
Electrons, photons	3.2
Luminosity	3.0
τ -leptons	2.5
Flavour tagging	1.8
MC statistical uncertainties	4.4



- Three analyses channels: $\tau_{lep}\tau_{lep\tau}$, $\tau_{lep}\tau_{had}$ and $\tau_{had}\tau_{had}$
- Background: $Z \rightarrow \tau\tau$ dominant in all the channels, contributions from top and vector bosons decays, as well as from misidentified leptonic or hadronic τ decays depend on the channel considered
- Two categories: 'boosted' ggF (additional recoiling jet) and VBF
- Higgs is reconstructed from the visible decay products of τ_s and from MET
- Di-tau invariant mass derived using the Missing Mass Calculator, MMC
- Strategy: maximum likelihood fit performed on data using distributions of the di-tau mass in SRs simultaneously with event yields from CRs (included to constrain normalization of major backgrounds estimated from simulation)



NEW PHYSICS: THE EFT APPROACH

- The SM can be supplemented by possible new physics effects
3 dim6 operators modifying the t- and b- quark couplings + point-like Hg coupling

$$\mathcal{L} = \mathcal{L}_{SM} + \sum_i \frac{c_i}{\Lambda^2} \mathcal{O}_i + \dots$$

New physics at the the scale Λ
[arXiv:1612.00283](https://arxiv.org/abs/1612.00283)

With c_i dimensionless coefficients and \mathcal{O}_i operators of dim6 from the SM

- Multi-TeV scale can be:

- tested with sub-percent level measurements

$$\delta O \sim \left(\frac{v}{\Lambda}\right)^2 \sim 6\% \left(\frac{\text{TeV}}{\Lambda}\right)^2 - 1\% \text{ effect on coupling for } \Lambda \sim 2.5 \text{ TeV}$$

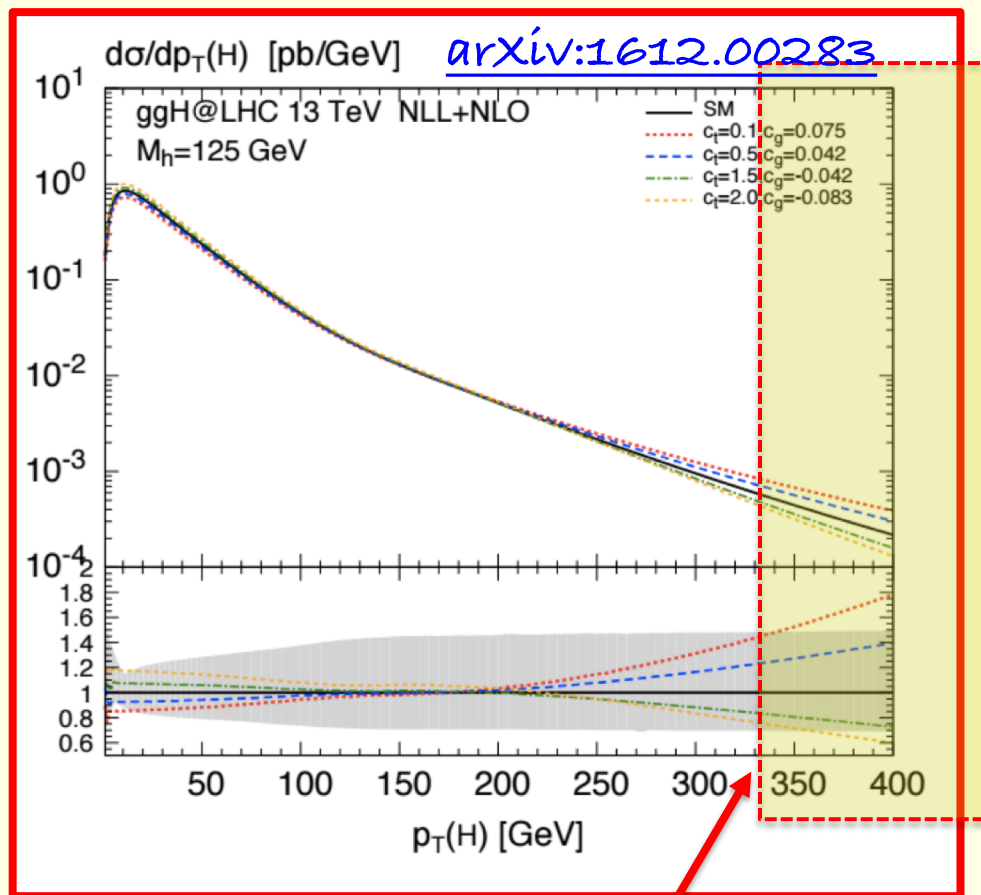
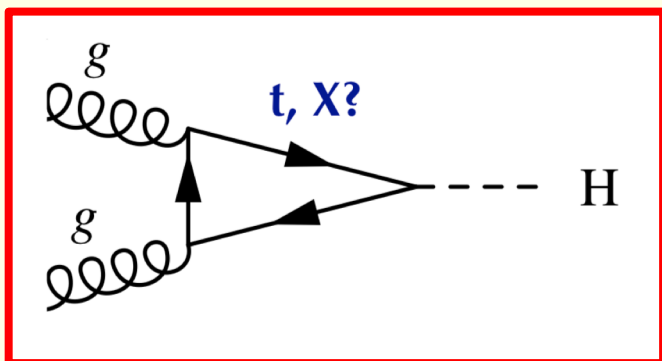
- tested with high- p_T measurements even if at low precision

$$\delta O \sim \left(\frac{Q}{\Lambda}\right)^2 - 16\% \text{ effect on coupling for } \Lambda \sim 2.5 \text{ TeV}$$



HIGH- p_T HIGGS

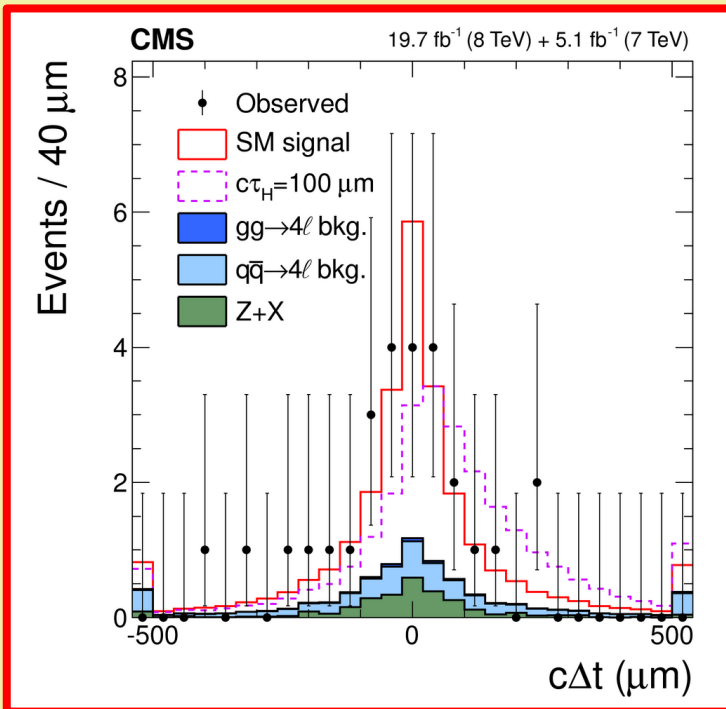
➤ At high- p_T Higgs we can probe modifications in top couplings



ggF Hbb: powerful test

HIGGS WIDTH: DIRECT STRATEGIES

From the lifetime



➤ using the Higgs lifetime we can set a direct lower bound

$$\Delta t = \frac{m_{4l}}{p_T} (\Delta \vec{r}_t \cdot \widehat{p}_T) \rightarrow$$

$$\langle \Delta t \rangle = \tau_H = \hbar / \Gamma_H$$

Lifetime of each H candidate

- $\Delta \vec{r}_t$ Displacement between the production and decay vertices in the transverse plane

➤ Observables:

Δt and $D_{bkg}(m_{4l}$ and $D^{kin})$



HIGGS WIDTH: FROM COUPLINGS

➤ Using the coupling analysis framework we can constraint Γ_H :

$$\Gamma_i = \Gamma_i^{SM} \cdot k_i^2 \text{ and so } \Gamma_H = \frac{k_H^2 \cdot \Gamma_H^{SM}}{1 - B_{BSM}}$$

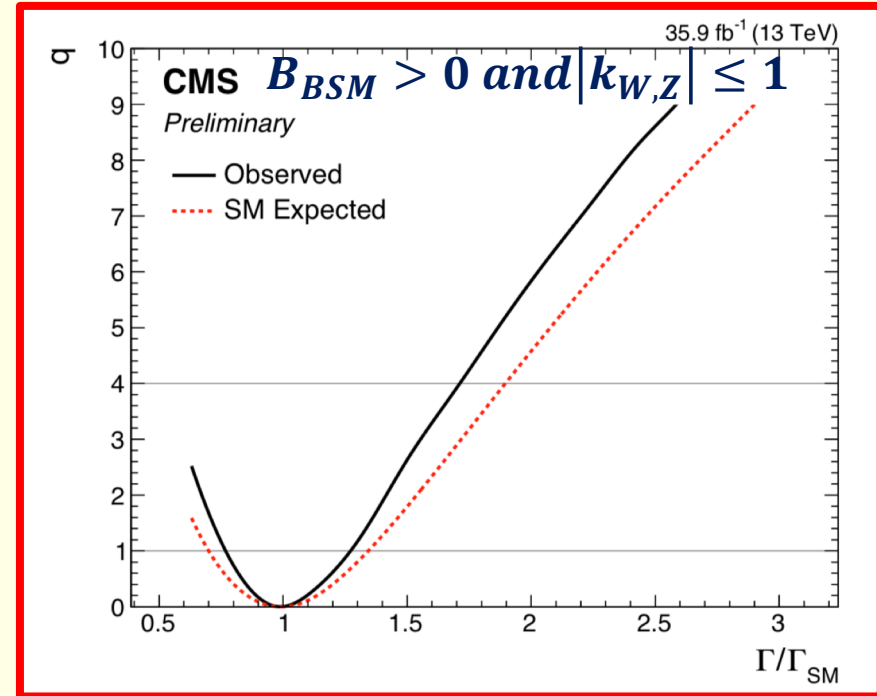
Two possible interpretations:

- $B_{BSM} = 0$
- $B_{BSM} > 0 \text{ and } |k_{W,Z}| \leq 1$

Production	Loops	Interference	Effective scaling factor	Resolved scaling factor
$\sigma(\text{ggH})$	✓	b - t	κ_g^2	$1.04 \cdot \kappa_t^2 + 0.002 \cdot \kappa_b^2 - 0.038 \cdot \kappa_t \kappa_b$
$\sigma(\text{VBF})$	-	-	-	$0.73 \cdot \kappa_W^2 + 0.27 \cdot \kappa_Z^2$
$\sigma(\text{WH})$	-	-	-	κ_W^2
$\sigma(\text{qq/qg} \rightarrow \text{ZH})$	-	-	-	κ_Z^2
$\sigma(\text{gg} \rightarrow \text{ZH})$	✓	Z - t	-	$2.46 \cdot \kappa_Z^2 + 0.47 \cdot \kappa_t^2 - 1.94 \cdot \kappa_Z \kappa_t$
$\sigma(\text{ttH})$	-	-	-	κ_t^2
$\sigma(\text{gb} \rightarrow \text{WtH})$	-	W - t	-	$2.91 \cdot \kappa_t^2 + 2.40 \cdot \kappa_W^2 - 4.22 \cdot \kappa_t \kappa_W$
$\sigma(\text{qb} \rightarrow \text{tHq})$	-	W - t	-	$2.63 \cdot \kappa_t^2 + 3.58 \cdot \kappa_W^2 - 5.21 \cdot \kappa_t \kappa_W$
$\sigma(\text{bbH})$	-	-	-	κ_b^2
Partial decay width				
Γ^{ZZ}	-	-	-	κ_Z^2
Γ^{WW}	-	-	-	κ_W^2
$\Gamma^{\gamma\gamma}$	✓	W - t	κ_γ^2	$1.59 \cdot \kappa_W^2 + 0.07 \cdot \kappa_t^2 - 0.67 \cdot \kappa_W \kappa_t$
$\Gamma^{\tau\tau}$	-	-	-	κ_τ^2
Γ^{bb}	-	-	-	κ_b^2
$\Gamma^{\mu\mu}$	-	-	-	κ_μ^2

Total width for $BR_{BSM} = 0$

Γ_H	✓	-	κ_H^2	$0.58 \cdot \kappa_b^2 + 0.22 \cdot \kappa_W^2 + 0.08 \cdot \kappa_g^2 + 0.06 \cdot \kappa_t^2 + 0.026 \cdot \kappa_Z^2 + 0.029 \cdot \kappa_c^2 + 0.0023 \cdot \kappa_\tau^2 + 0.0015 \cdot \kappa_{Z\gamma}^2 + 0.00025 \cdot \kappa_s^2 + 0.00022 \cdot \kappa_\mu^2$
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HIGGS WIDTH:

OFF-SHELL STRATEGY

- Maximum likelihood fit to the Matrix-Element (ME) based discriminant distribution (4l) and the transverse-mass, $m_T(ZZ)$, distribution (2l2v)

1

Off-shell signal strength constraints

- Combination of the 2l2v and 4l channel fixing the ratio of the signal strength in ggF and VBF to the SM

prediction:
$$\frac{\mu_{off-shell}^{ggF}}{\mu_{off-shell}^{VBF}} = 1$$

2

Higgs boson total width constraints

- Combination with the on-shell result assuming the same

- on-shell signal strength in VBF and ggF:
$$\frac{\mu_{on-shell}^{ggF}}{\mu_{on-shell}^{VBF}} = 1$$
- on-shell and off-shell couplings

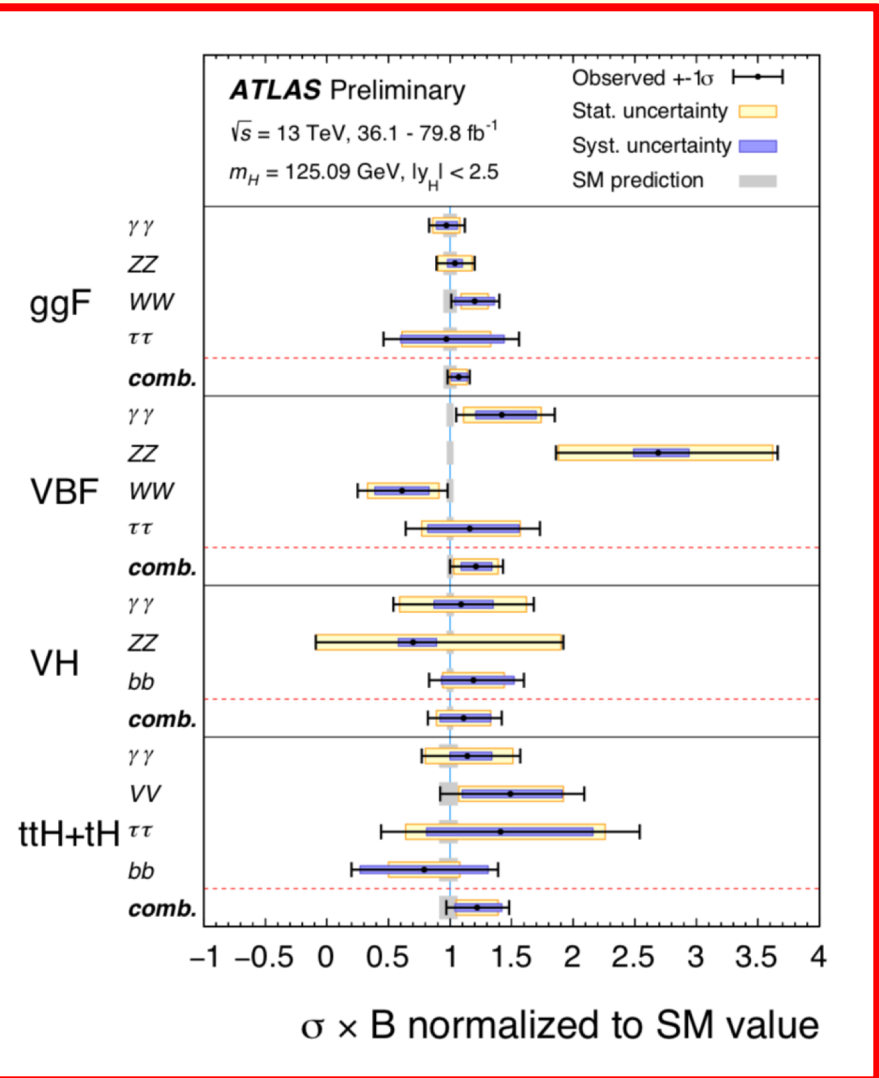
CROSS-SECTIONS IN RUN 2

➤ No absolute couplings measurements at LHC

➤ For the observable $\sigma(AA \rightarrow H)BR(H \rightarrow BB)$ the quantity measurable is proportional to:

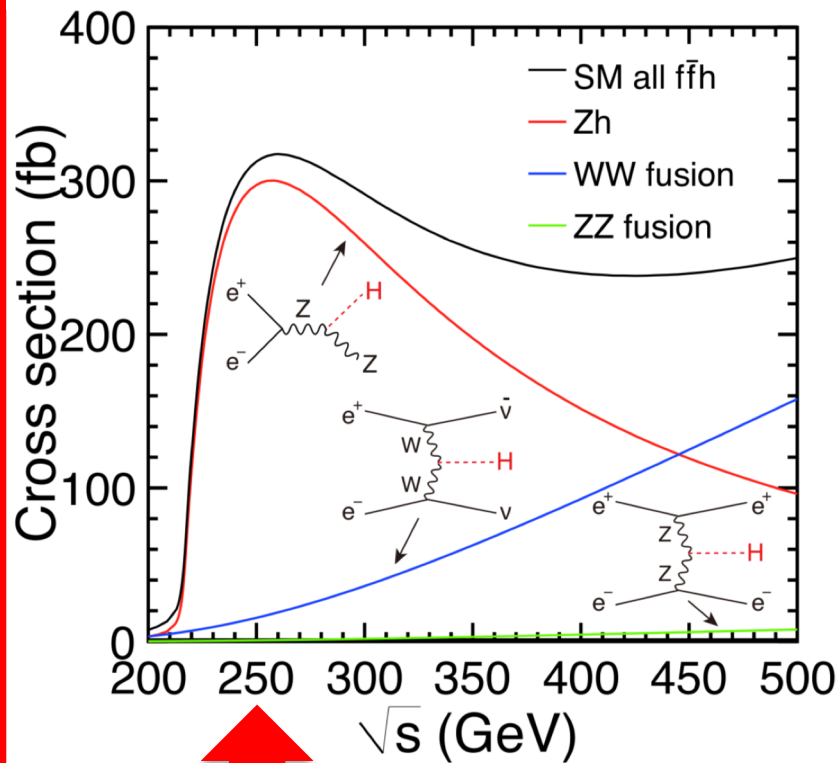
$$g^2(HAA)g^2(HBB)/\Gamma_H$$

The Higgs boson total width must be known to have absolute magnitudes of the couplings!



HIGGS PHYSICS AT ILC

- At ILC the total Higgs production cross section could be measured ➡ measurement of Γ_H
- Depending on \sqrt{s} different production modes

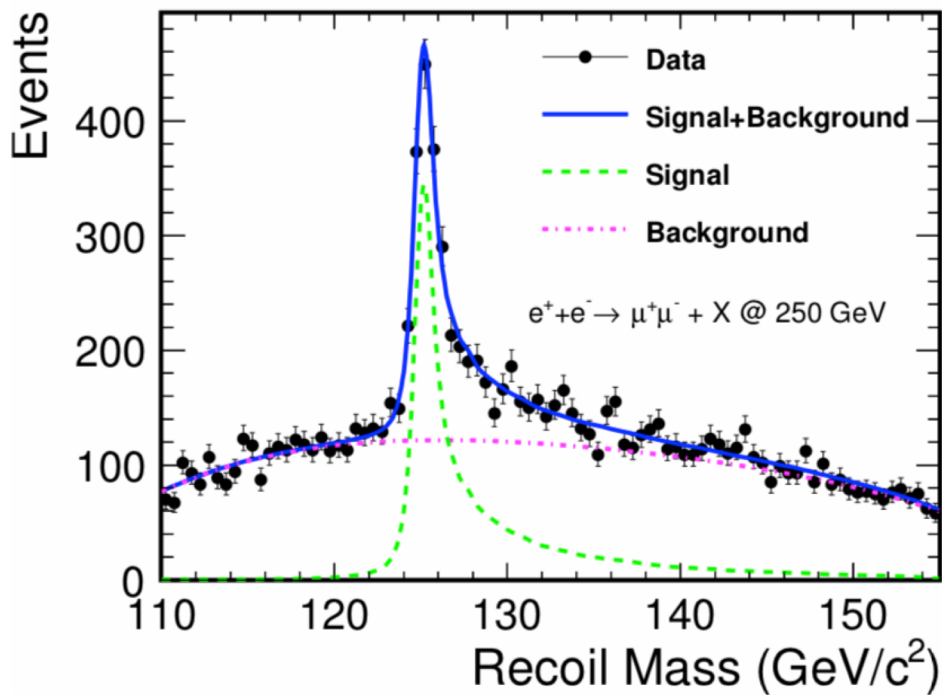


- The Higgs-strahlung production is maximum at 250 GeV
- 2000 fb⁻¹ in 20 years of data acquisition (H20 program):
 - ZH ➡ ~500 K Higgs
 - WW-fusion ➡ ~15 K Higgs

- ZH cross section measurable at 1.0%
- From the HZ sample, measurement of g_{HZZ} : $\sigma(e^+e^- \rightarrow ZH) \propto g_{HZZ}^2$

HZ COUPLING AT ILC

- Unique opportunity for a model-independent measurement of the HZ coupling from the recoil mass distribution in $e^+e^- \rightarrow ZH$



$$M_{rec}^2 = (\sqrt{s} - E_{ll})^2 - |\vec{p}_{ll}|^2$$

- Higgs events are tagged with the Z boson decays, independently of the Higgs decay mode
- From the HZ sample, measurement of g_{HZZ} :

$$\sigma(e^+e^- \rightarrow ZH) \propto g_{HZZ}^2$$

g_{HZZ} : KEY TO THE ILC PROGRAM

- From the ratio of the Higgs-strahlung and WW-fusion cross sections for the same exclusive Higgs boson final-state $H \rightarrow X\bar{X}$:

$$\frac{\text{Diagram 1}}{\text{Diagram 2}} = \frac{\frac{g_{HZZ}^2 \cdot g_{HXX}^2}{\Gamma_H}}{\frac{g_{HWW}^2 \cdot g_{HXX}^2}{\Gamma_H}} \Rightarrow g_{HWW}^2$$

Diagram 1: $e^+e^- \rightarrow ZH$ (Higgs-strahlung)

Diagram 2: $e^+e^- \rightarrow W^+W^-H$ (WW-fusion)

Measuring $\sigma(e^+e^- \rightarrow ZH) \times BR(H \rightarrow WW^*) \propto \frac{g_{HZZ}^2 \cdot g_{HWW}^2}{\Gamma_H}$

Γ_H Accuracy achievable 1.7% (ILC250+ILC500)

COUPLINGS PRECISION IN FUTURE

ATLAS Simulation Preliminary

$\sqrt{s} = 14 \text{ TeV}$: $\int \mathcal{L} dt = 300 \text{ fb}^{-1}$; $\int \mathcal{L} dt = 3000 \text{ fb}^{-1}$

