



Neutron width statistics in a realistic resonance-reaction model



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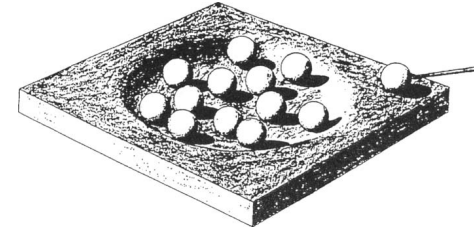
15th International Conference on Nuclear Reaction Mechanisms,
Varenna, Italy.

- Statistical model of compound nucleus reactions.
- Violation of statistical model in $n + \text{Pt}$ reaction and overview of proposed theoretical explanations.
- Novel realistic resonance-reaction model.
- Cross sections and neutron width fluctuations in the $n + {}^{194}\text{Pt}$ reaction.
- Summary and conclusions.

PF, G. F. Bertsch, and Y. Alhassid, arXiv:1710.00792, accepted in Phys. Rev. C

Statistical model of compound nucleus reactions

- Compound nucleus (CN): complex, equilibrated system of projectile + target.
- **Statistical model:** CN Hamiltonian is a random matrix from the Gaussian Orthogonal Ensemble (GOE).

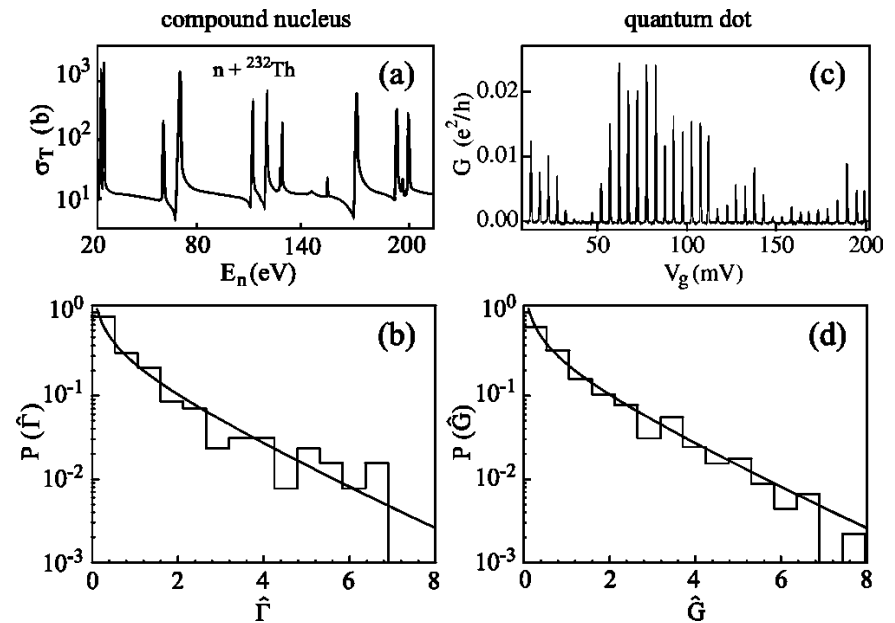


Bohr, Nature (1936)

- CN resonance energies and widths exhibit local fluctuations that are independent of system-specific details.

Mitchell, Weidenmüller, and Richter, RMP (2010).

- Modifies Hauser-Feshbach theory of CN reactions through **width fluctuation correction factor**.
 - Enhances elastic channels; decreases other open channels.
- Widely used in statistical reaction calculations.

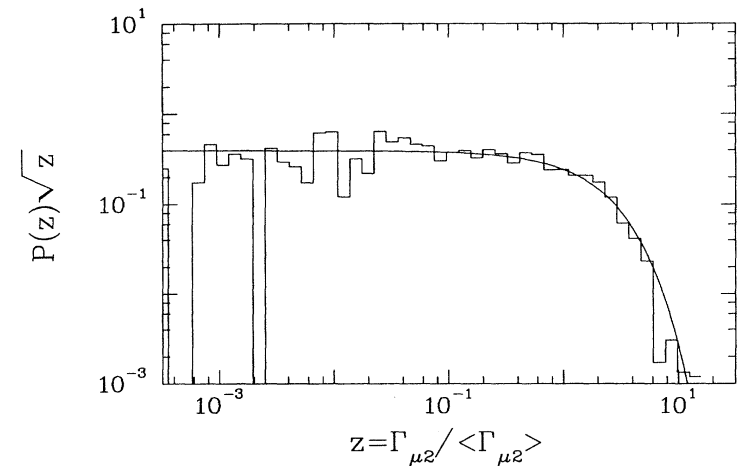


Alhassid, RMP (2000)

Violation of Porter-Thomas Distribution in Platinum

- For isolated resonances, widths for any reaction channel should fluctuate according to the Porter-Thomas distribution (PTD).
- Neutron resonance widths are considered compelling evidence for PTD, especially the Nuclear Data Ensemble.
[Haq, Bohigas, and Pandey, PRL \(1982\).](#)
- Evidence in other quantum chaotic systems described by the GOE, e.g. microwave stadium billiard.
- **Koehler *et al.* PRL (2010): neutron width fluctuations from neutron scattering off Pt isotopes are much broader than the PTD.**
- Other statistical model violations observed in reanalysis of the NDE and in total gamma width distributions.

[Porter and Thomas, Phys. Rev. \(1956\).](#)



[Alt *et al.* PRL \(1995\)](#)

[Koehler, PRC \(2011\).](#)
[Koehler *et al.* PRC \(2013\).](#)

Statistical Model Explanations

1. Energy dependence of neutron strength function changed by near-threshold bound or virtual state of neutron channel potential [Weidenmüller PRL 2010].

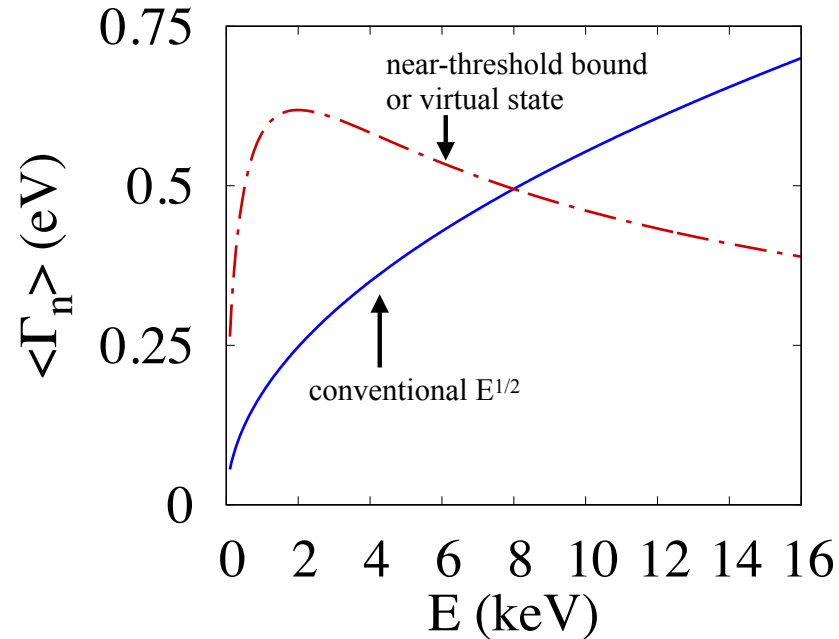
- Evidence: peak in the NSF near Pt.

2. Non-statistical interactions of the resonances due to coupling to the neutron channel perturbs the GOE [Celardo *et al.* PRL (2011), Volya *et al.* PRL (2015), Bogomolny PRL (2017)].

$$H^{\text{eff}} = H^{\text{GOE}} + \Delta - i\pi W W^T$$

↑ effective Hamiltonian of CN resonances ↑ GOE Hamiltonian of bound CN states ↙ ↘ real (imaginary) shift due to off (on)-shell couplings to channels

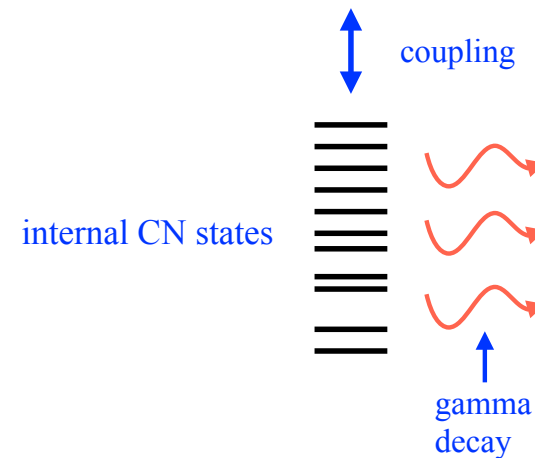
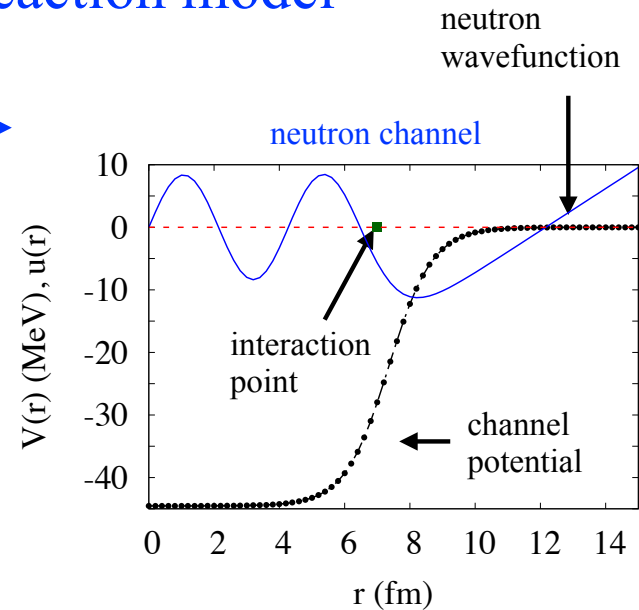
- Explanations have not been assessed together in a realistic model. Origin of observed PTD violation still an open question



Realistic resonance-reaction model

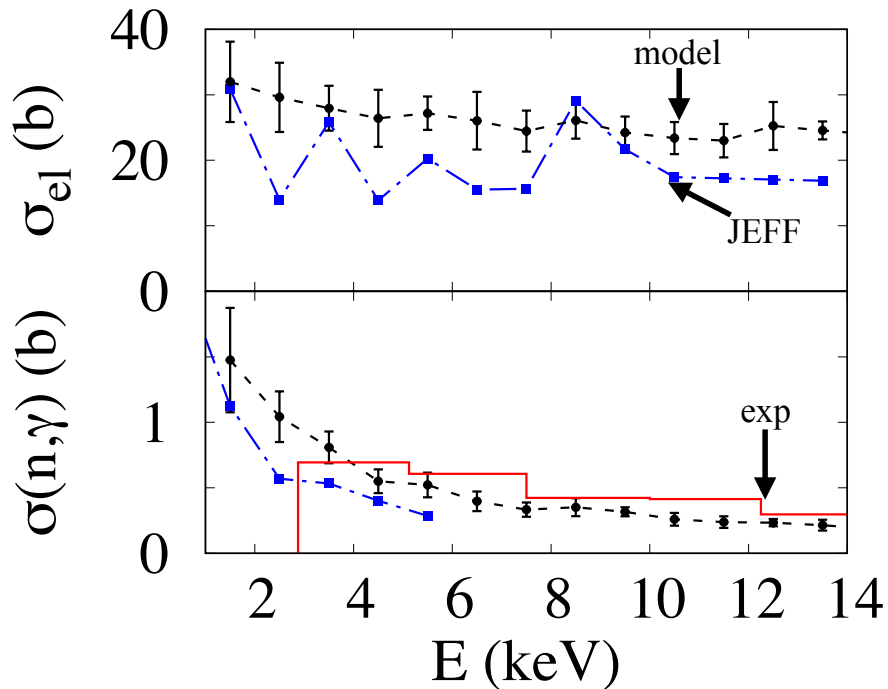
$$\mathbf{H} = \begin{pmatrix} \mathbf{H}_n & \mathbf{V} \\ \mathbf{V}^T & \mathbf{H}_c \end{pmatrix} \longleftrightarrow$$

- Neutron channel represented on a spatial mesh with a Woods-Saxon potential.
 - CN states have GOE spectrum with a constant imaginary shift to describe gamma decay.
 - Channel coupled to CN states at one interaction point. Coupling strengths have a Gaussian distribution.
 - Able to calculate resonance energies and widths, as well as cross sections, within the same framework.
 - Based on [Mazama](#) code of G. F. Bertsch.
- [Bertsch, Brown, and Davis, arXiv:1804.00364 \(2018\).](#)

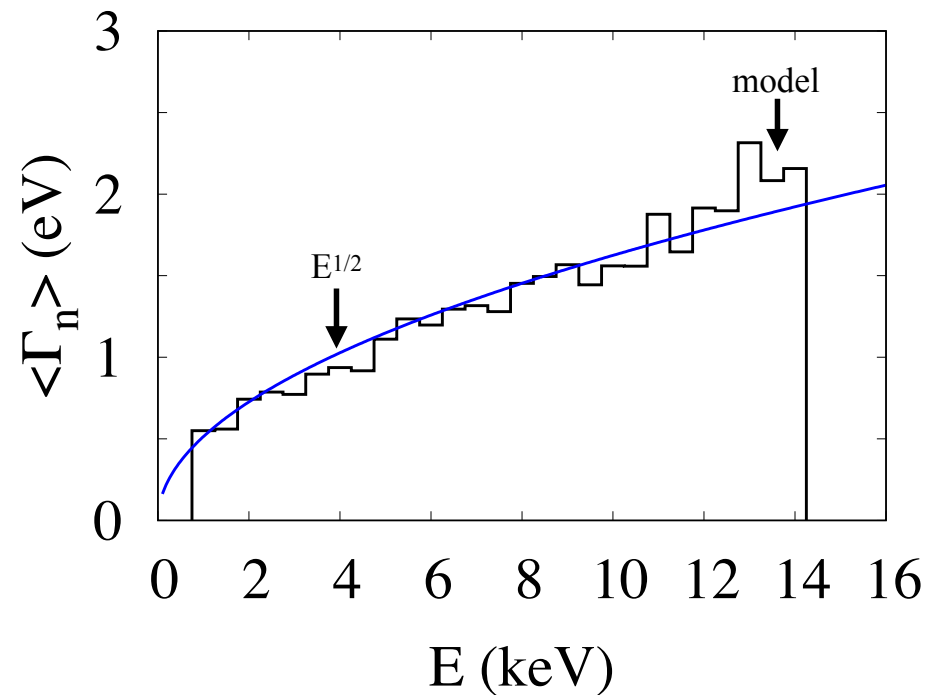


n + ¹⁹⁴Pt: Baseline Model

- Woods-Saxon: $V_0 = -44.54$ MeV, $r_0 = 1.27$ fm, $a_0 = 0.67$ fm from Bohr and Mottelson.
- Average resonance spacing $D = 82$ eV, gamma width $\Gamma_\gamma = 72$ meV from RIPL-3.
- Average coupling strength parameter $v_0 = 11$ keV-fm^{1/2} tuned to match RIPL-3 neutron strength function at 8 keV neutron energy.

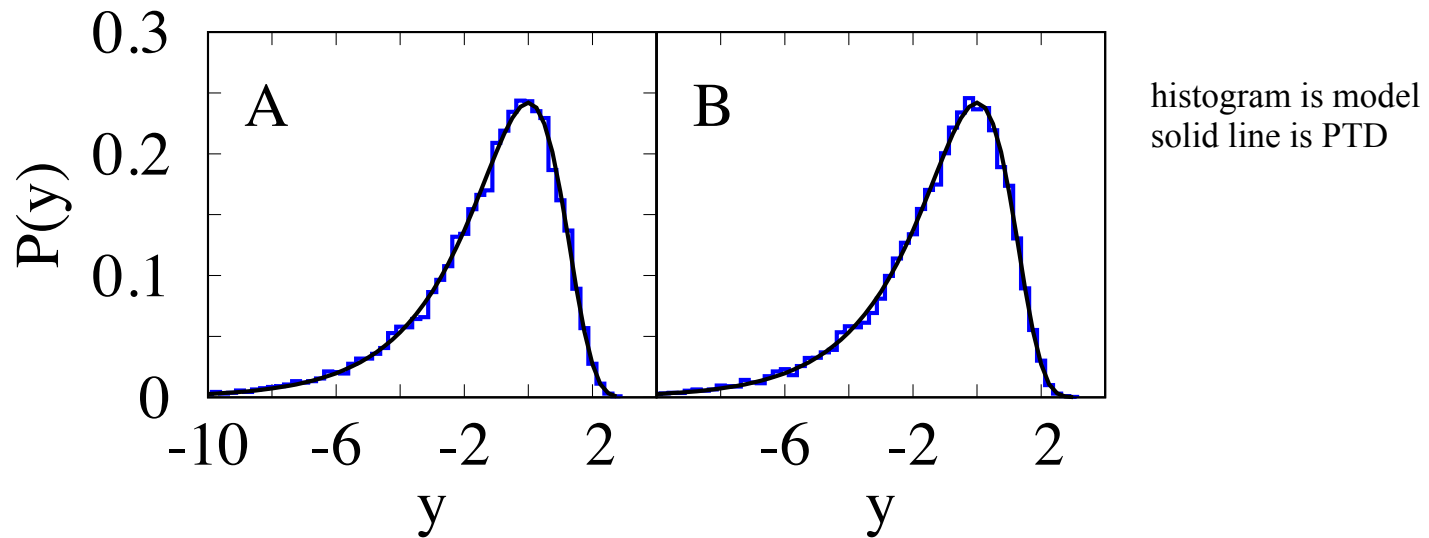


Model cross sections agree with JEFF-3.2 library and experimental cross section [Koehler and Guber, PRC (2012)].



Average neutron width matches the $E^{1/2}$ expectation.

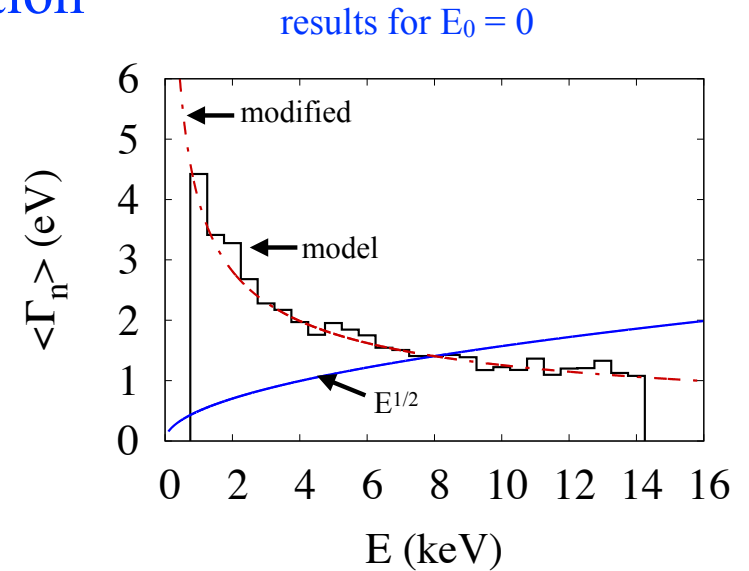
Baseline neutron width fluctuations



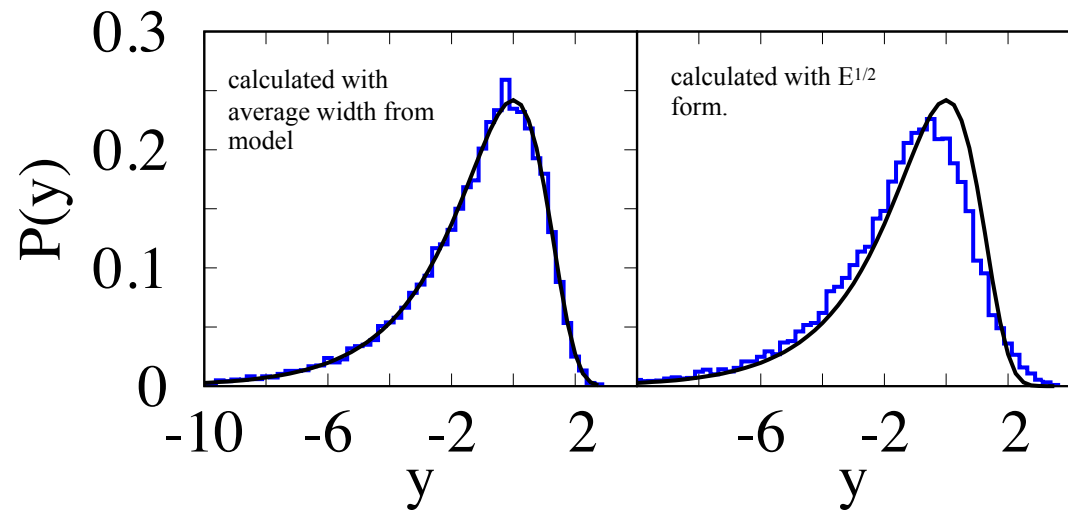
- Reduced neutron width: $\gamma_{n,r} = \Gamma_{n,r} / \langle \Gamma_n \rangle (E_r)$
- **Reduction A**: average neutron width taken from model.
- **Reduction B**: assumes $\langle \Gamma_n \rangle (E) \propto E^{1/2}$
- Figure shows distribution of $y = \ln(x)$, $x = \gamma / \langle \gamma \rangle$

Parameter variation

- Reducing neutron channel potential depth by a few MeV creates a near-threshold bound or virtual state of the channel.
- Modifies the $E^{1/2}$ behavior of the average width, agrees with the formula of Weidenmüller: $\langle \Gamma_n \rangle(E) \propto \frac{\sqrt{E}}{E + |E_0|}$
- The reduced width distribution extracted with the $E^{1/2}$ form is broader than PTD.
- Results insensitive to physically large changes in average channel-CN coupling.



histogram is model
solid line is PTD



Summary and Conclusions

- In our realistic model, the PTD describes the neutron width fluctuations well if the energy dependence of the average neutron width is correctly described.
- This study excludes explanations for the observed PTD violation based on coupling to the neutron channel.
- Within a reasonable parameter range, there can be a near-threshold bound or virtual state of the neutron channel. In this case, using the $E^{1/2}$ form will produce PTD violation. Near-threshold state must be within a few keV to have a significant effect.
- A reanalysis by the experimentalists [Koehler *et al.* arXiv:1101.4533 (2011)] found that using this state did not improve their agreement with the PTD. Problem with multilevel R-matrix analysis?
- The elastic and total cross sections are quite sensitive to the existence of such a state.
- More experimental investigation into resonance width statistics would be useful to explore possible breakdowns of statistical model.

Thank you for your attention!

Extra Slides

Resonance determination

- To find the complex resonance wavenumbers k_r , solve Schrödinger equation with appropriate boundary conditions

$$u(0) = 0$$

$$u(r) \rightarrow B(k)e^{ikr} \quad \text{for large } r$$

- neutron wavefunction is regular at origin.
- neutron wavefunction is purely outgoing.

$$\Rightarrow u(N_n + 1) = u(N_n)e^{ik\Delta r}$$

- With discretized approach, obtain a **nonlinear eigenvalue problem (NEVP)**.

$$\mathbf{M}(k)\vec{u} = [\mathbf{H} - E - te^{ik\Delta r}\mathbf{C}]\vec{u}$$

- Solve NEVP with an iterative method to find resonance wavenumbers k_r .

$$t = \hbar^2/2m\Delta r^2, \quad \mathbf{C}_{ij} = \delta_{i,N_n}\delta_{ij}$$

- Find resonance energies, total widths, and neutron widths from wavenumbers.

$$E_r - \frac{i}{2}\Gamma_r = \frac{\hbar^2 k_r^2}{2m}$$

- Can calculate elastic and capture cross sections [details in additional slides].

$$\Gamma_{n,r} = \Gamma_r - \Gamma_\gamma$$

Non-statistical interactions

$$H^{\text{eff}} = H^{\text{GOE}} + \Delta - i\pi W W^T$$

↑ effective Hamiltonian of CN resonances

↑ GOE Hamiltonian of bound CN states

← real (imaginary) shift due to off (on)-shell couplings to channels

↑