



# Srážky těžkých iontů v detektoru ATLAS

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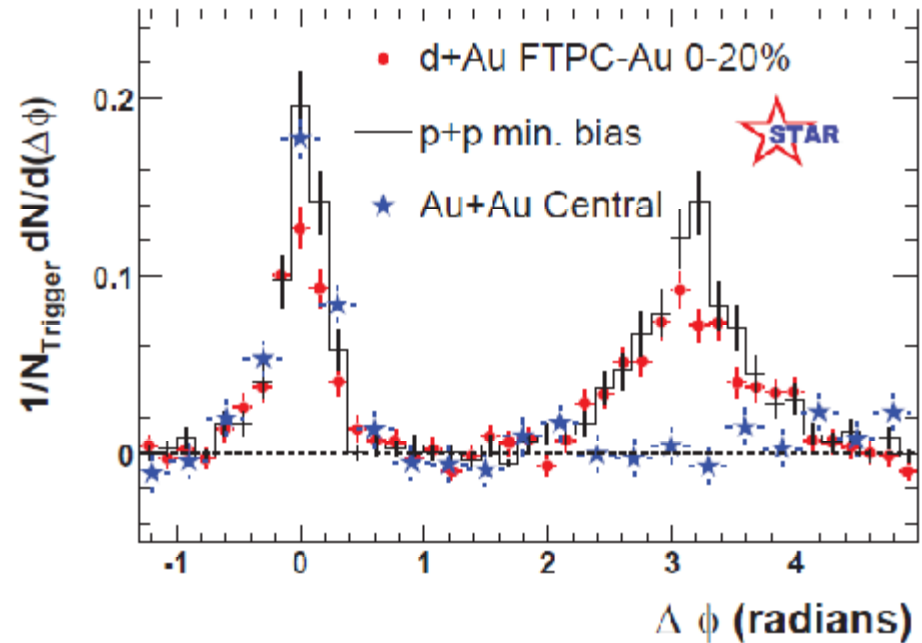
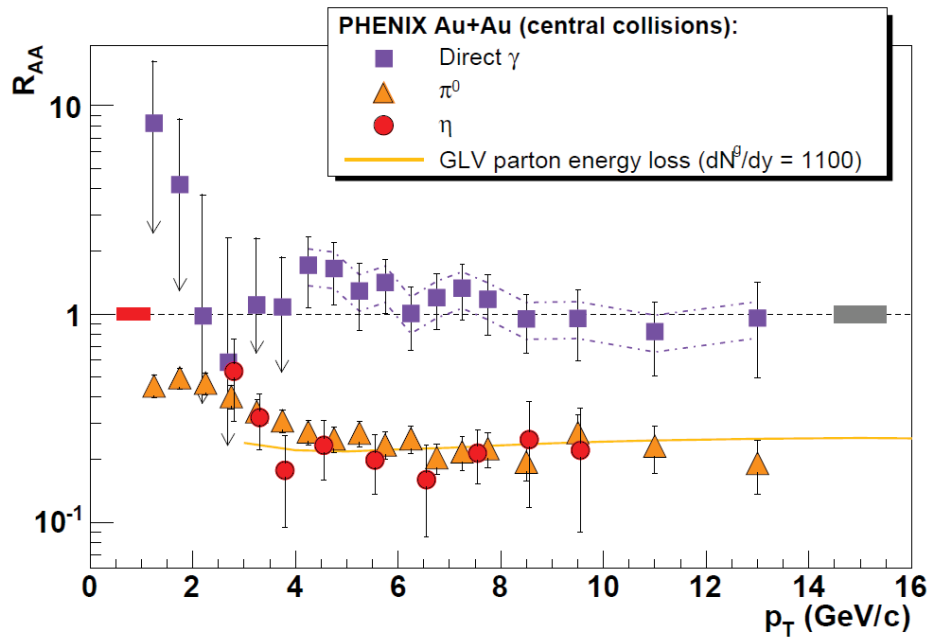
Martin Spousta, UČJF, MFF UK



# Motivation



- Main task is to better understand QCD at high temperature where quarks and gluons are deconfined.
- In this limit strong interactions can possibly exhibit unexpected characteristics that do not follow from perturbative considerations.
- One of useful tools are jets - jet quenching predicted in early 90's and observed at RHIC experiments.



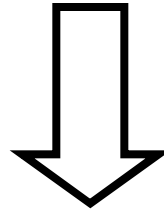


# Motivation



The **main results** concerning high- $p_T$  probes from RHIC experiments:

- High  $p_T$  hadrons are suppressed by a factor of  $\sim 5$  compared to binary scaled  $p+p$ .
- Away side peak is completely absent in di-jet correlation measurements.
- Measured spectra obey power-law behavior characteristic for pQCD processes.
- ...



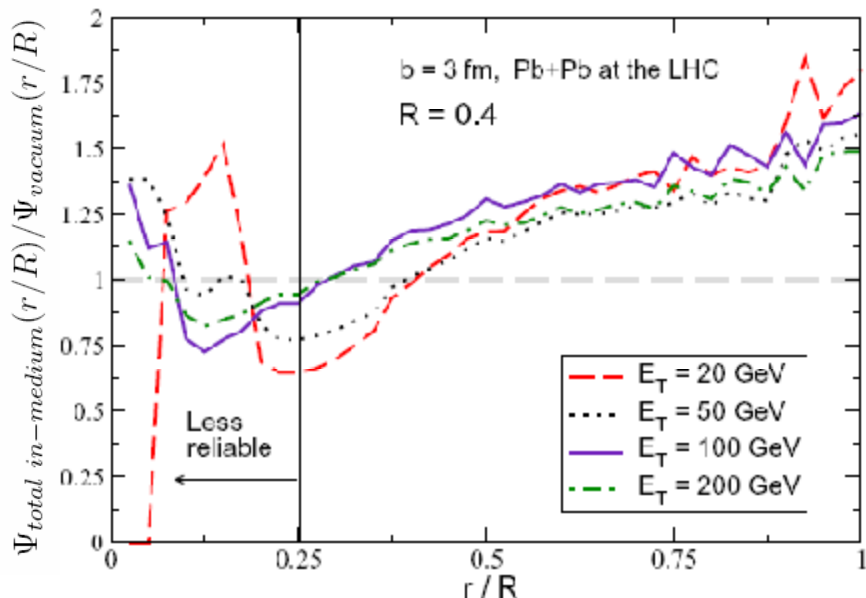
- Energy loss is a final state effect, not an initial state effect.
- Energy loss is at the parton level, not at the hadron level.
- Partons lose their energy dominantly by the radiative energy loss, collisional energy loss may be important only for heavy quarks.
- Partons lose  $1-10 \text{ GeV}/\text{fm}^3$  in the created medium. This energy loss seems to depend linearly on the path-length traveled through the expanding medium.



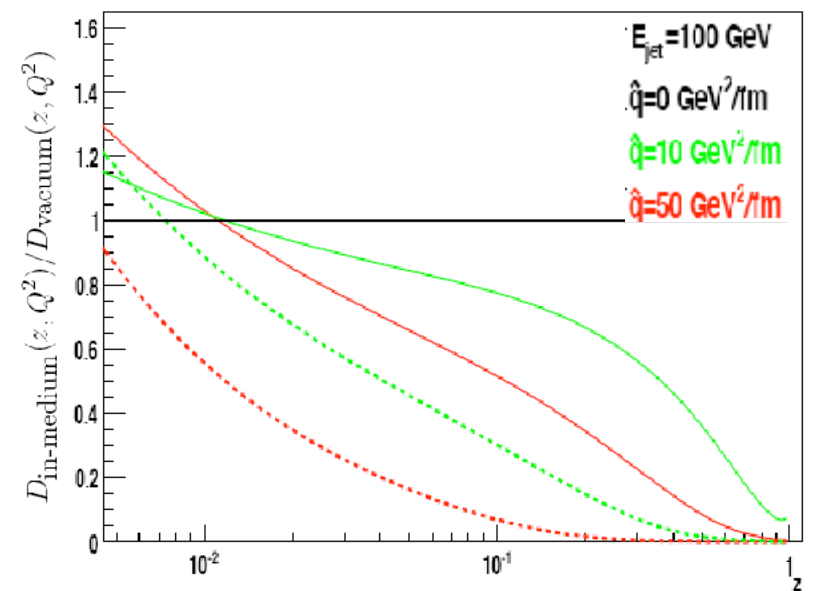
# Motivation



- Jets copiously produced at LHC energies  $\rightarrow$  possibility to study details of energy loss mechanism in QGP medium
- Models and predictions for LHC energies exist but the details of QCD energy loss mechanisms in medium not well understood



I.Vitev, *et al.* – modification of jet shapes



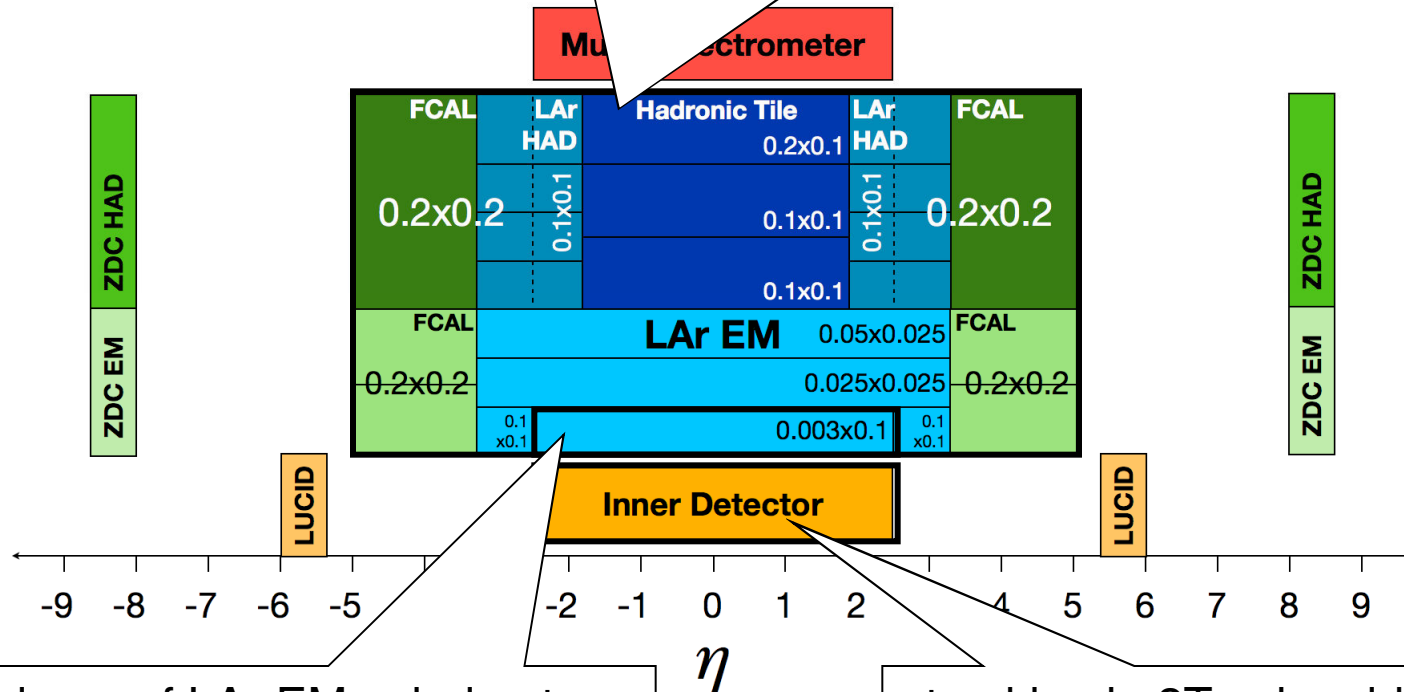
N.Armeo, *et al.* – modification of fragmentation functions



# Why ATLAS?



jet reconstruction using calorimeter,  
full azimuth, 10 units of pseudorapidity



first layer of LAr EM calorimeter  
excellent for photon isolation,  
other layers also well segmented

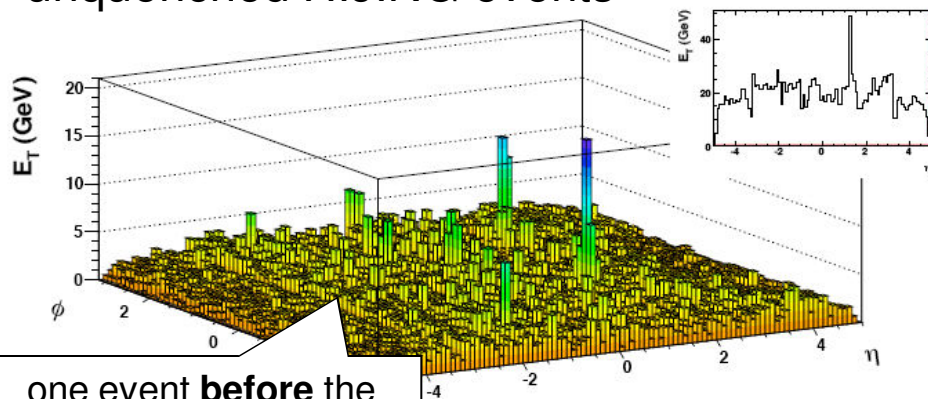
tracking in 2T solenoid –  
fragmentation studies



# Jet reconstruction strategy

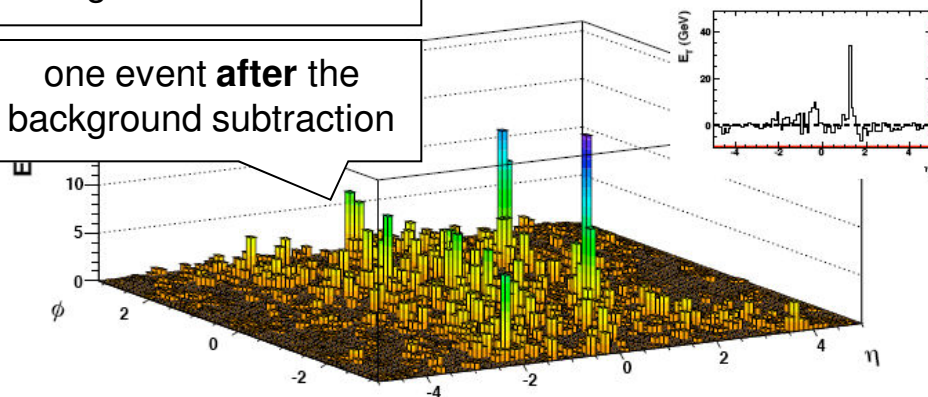


PYTHIA dijets embedded to the unquenched HIJING events



one event **before** the background subtraction

one event **after** the background subtraction



## Cone jet reconstruction:

- regions of interest found (seed regions) – fast sliding window algorithm used
- background computed
  - excluding the seed-regions
  - vs.  $\eta$ , vs. *layer*
- background subtracted
- standard p+p jet finding algorithm used (seeded iterative  $R=0.4$  cone algorithm)

$$R = \sqrt{(\Delta\phi)^2 + (\Delta\eta)^2}$$

An alternative:  **$k_T$ -algorithm** based reconstruction strategy – also studied



# Jet energy resolution

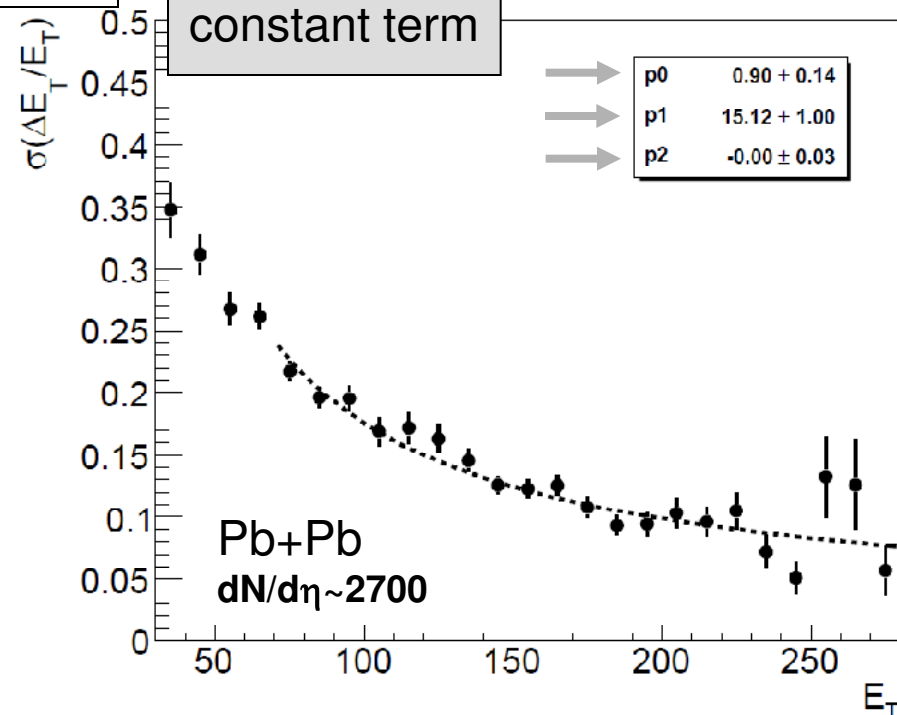
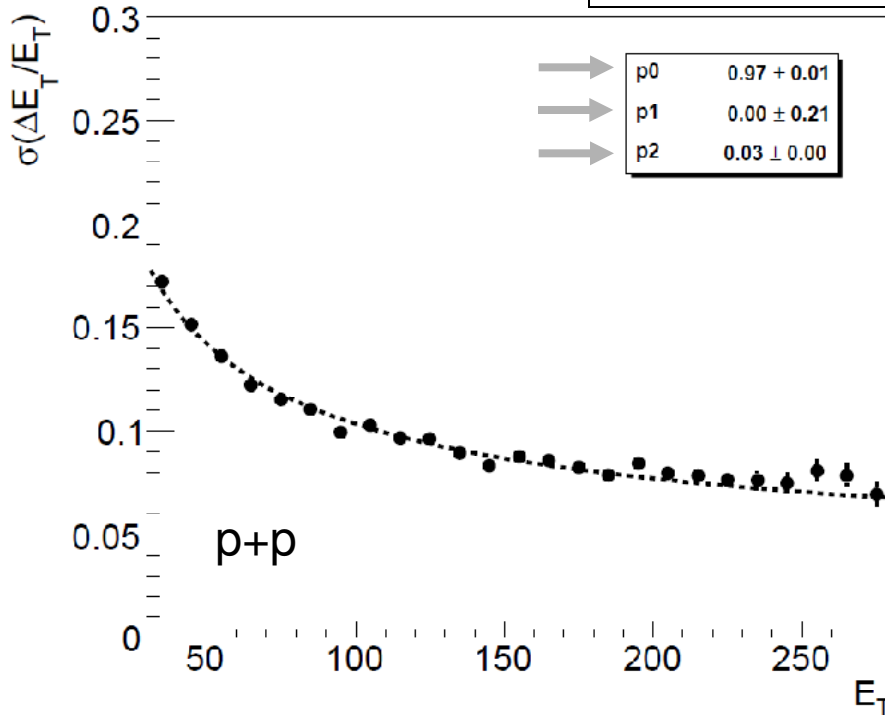
stochastic term

noise term

Fitting well known formula:

$$\frac{\sigma(E)}{E} = \frac{a}{\sqrt{E}} \oplus \frac{b}{E} \oplus c$$

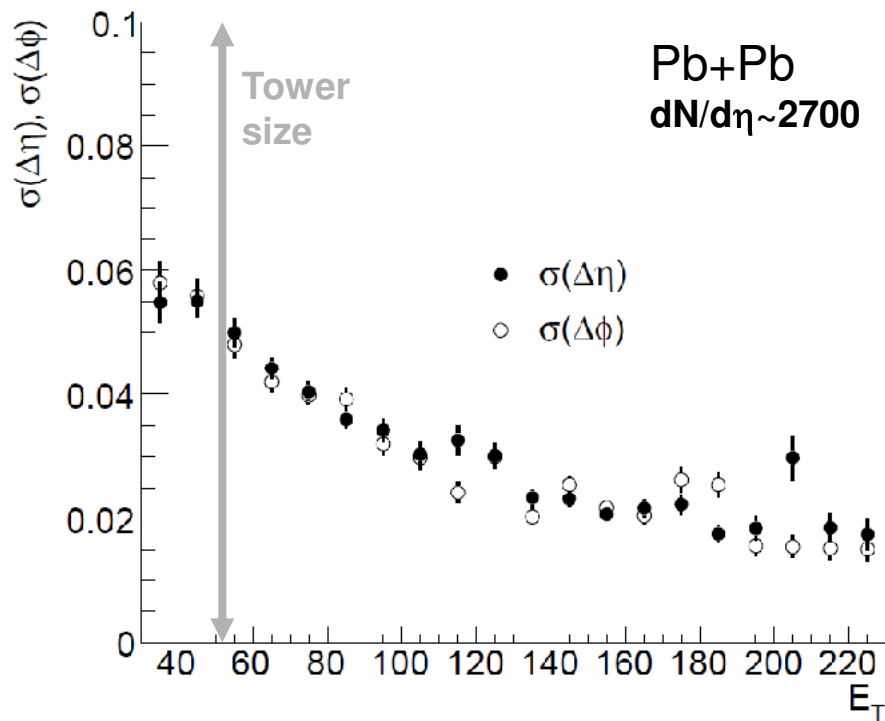
to p+p and Pb+Pb energy resolution



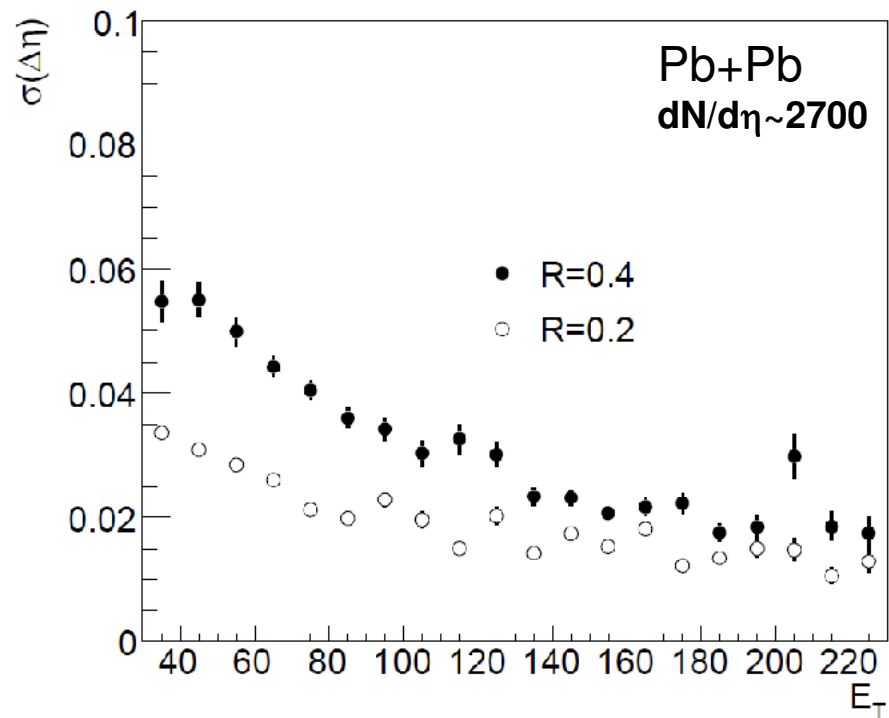
- Jet energy resolution below 25% for 70 GeV jets in the most central collisions (dN/dη ~ 2700 ⇔ b=2 fm, unquenched HIJING)
- Irreducible background fluctuations: ~ 15 GeV / jet



# Jet position resolution



- Jet position resolution in  $\phi$  similar to that in  $\eta$  (full field simulated)
- It improves with increasing jet energy, in the whole energy range jet position resolution is within a half of a tower

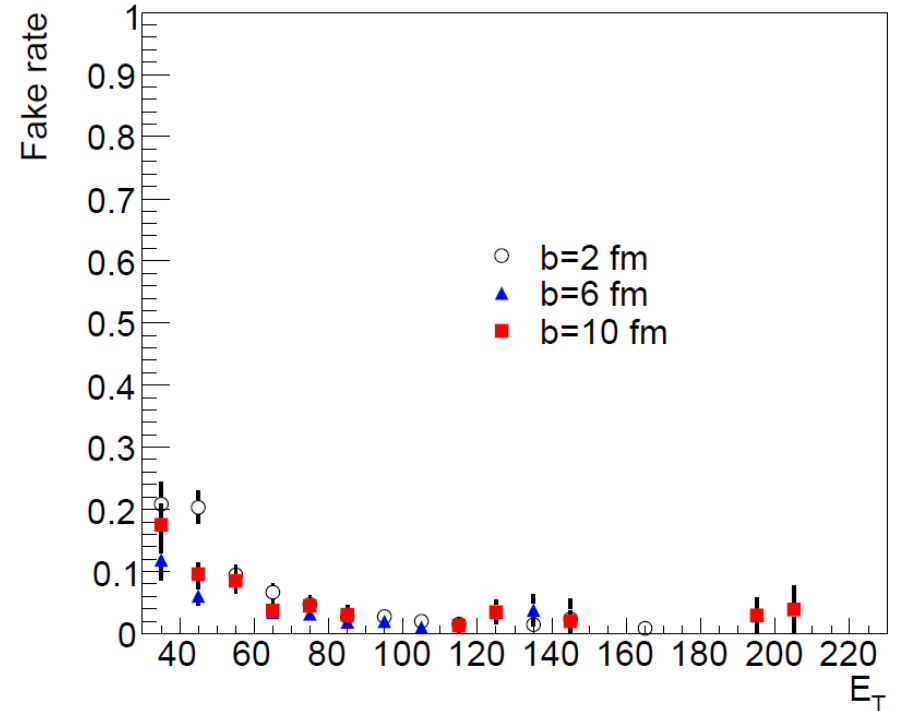
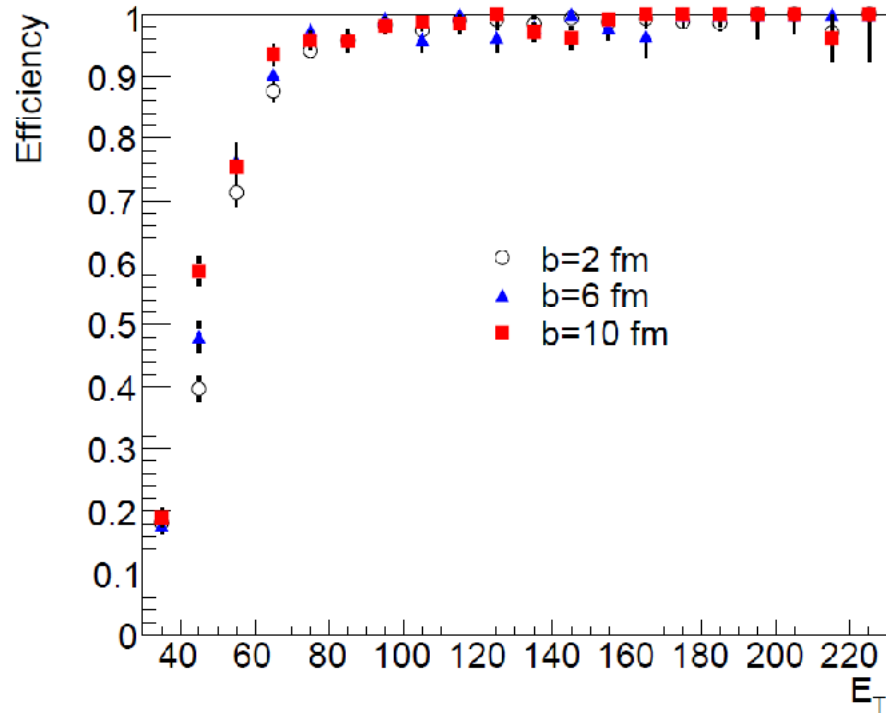


- Jet position resolution can be improved using a method of smaller cones: jet axis of a reconstructed jet is substituted by the jet axis from jet reconstructed with  $R < R_{\text{orig}}$





# Efficiency and fake-rate



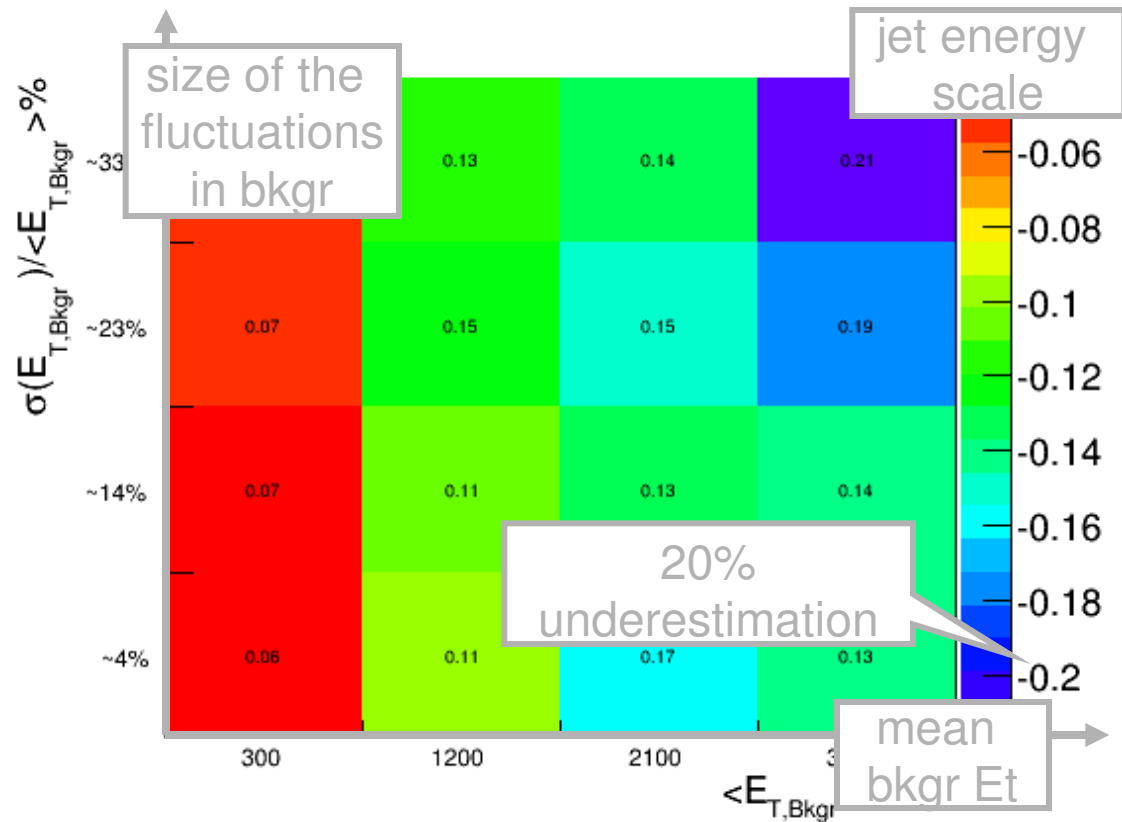
- Efficiency is almost centrality independent – easier interpretation of jet properties vs. centrality
- Above 70 GeV the efficiency is above 90%
- Above 70 GeV very low fake rate  $< 5\%$  (without any fake rejection)



# Comparison among jet finding algorithms



- $k_T$  algorithm exhibits serious problems with the jet energy scale: for more severe background,  $k_T$  underestimates the jet energy. This is due to the fact that  $k_T$  preferably clusters the softer part of a jet with the background  $k_T$
- Problematic algorithms:
  - $k_T$  algorithm,
  - Cambridge/Aachen
- Non-problematic:
  - anti- $k_T$  algorithm,
  - (ATLAS) cone algorithm
- Difference among algorithms should be studied with real data (real jets & real noise) **x** favoured strategy of ATLAS is to rely on one (max. two) jet algorithms!

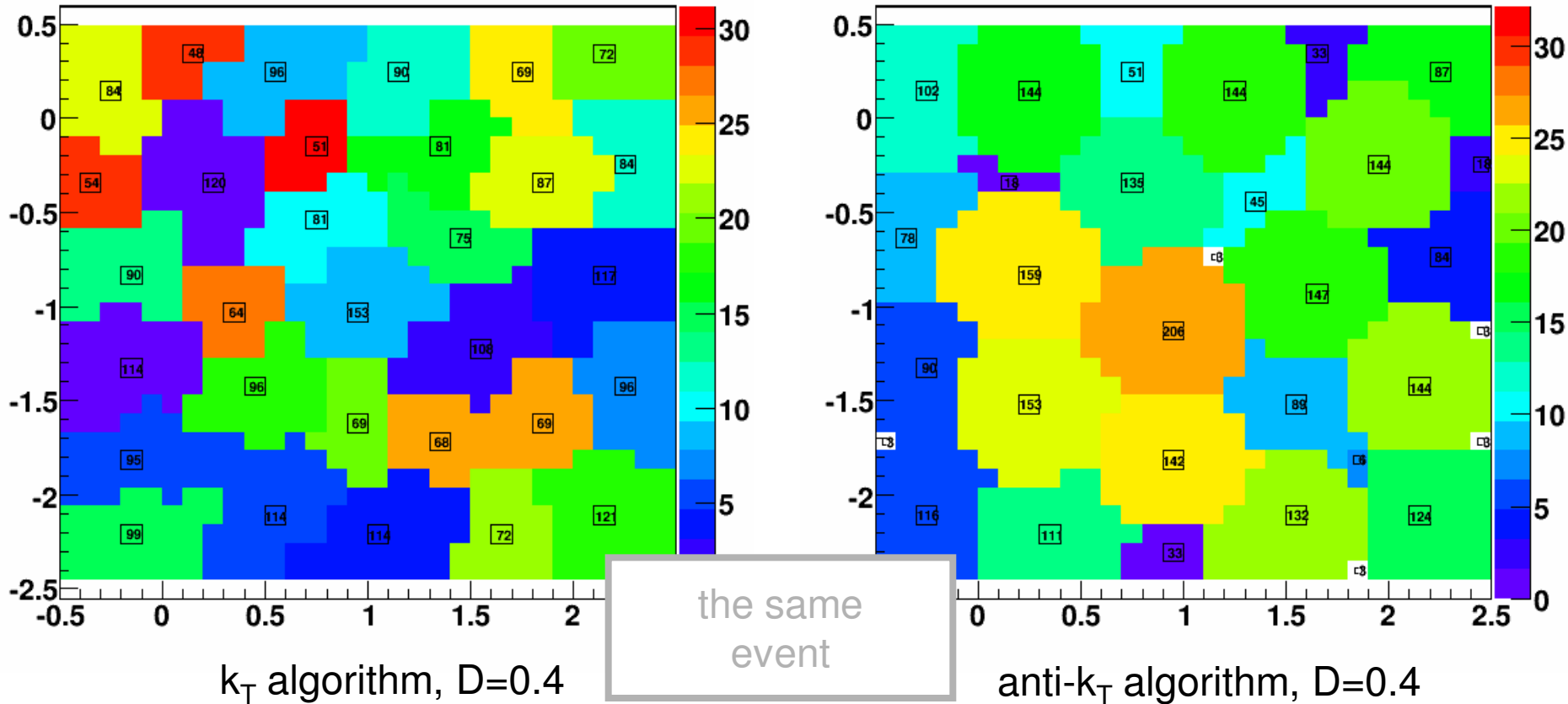




# Comparison among jet finding algorithms

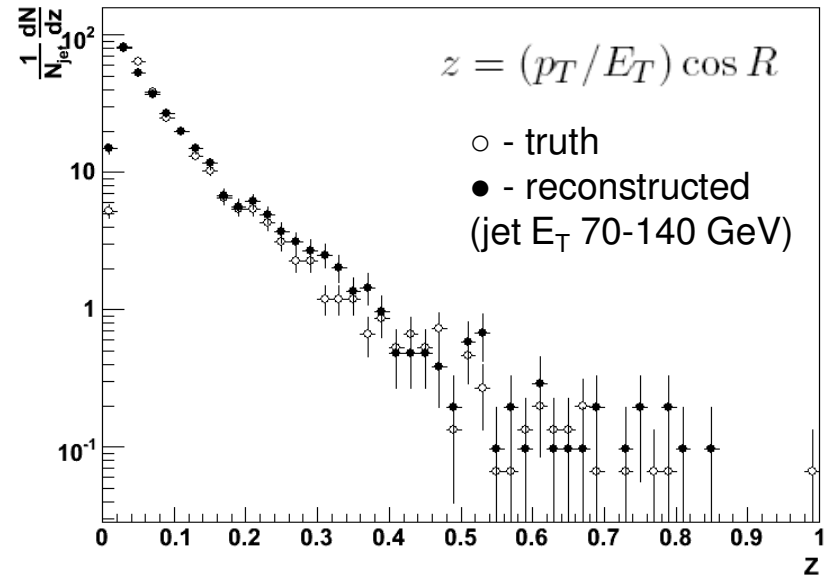
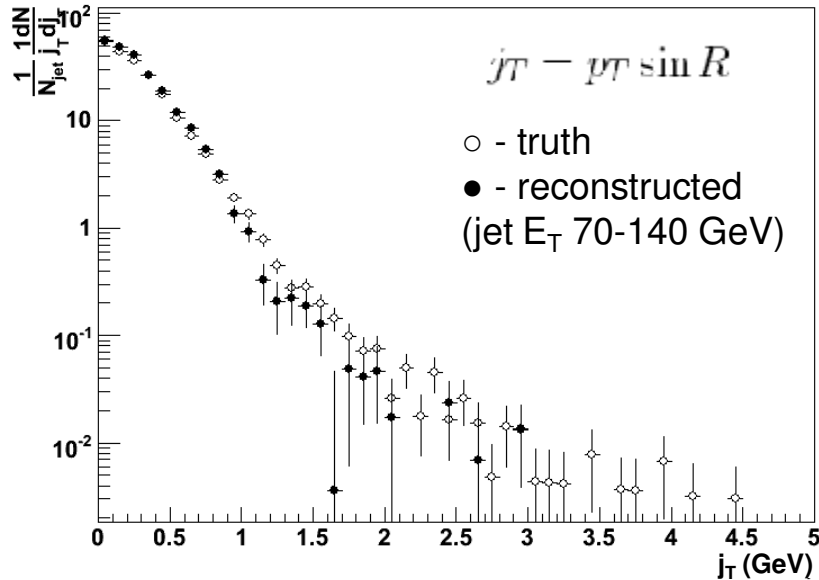


Area of jets delivered by different algorithms is very different in the noisy environment



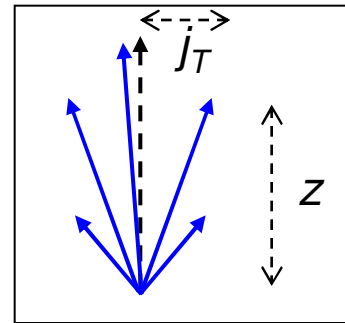


# Fragmentation function and $j_T$



## Reconstruction procedure:

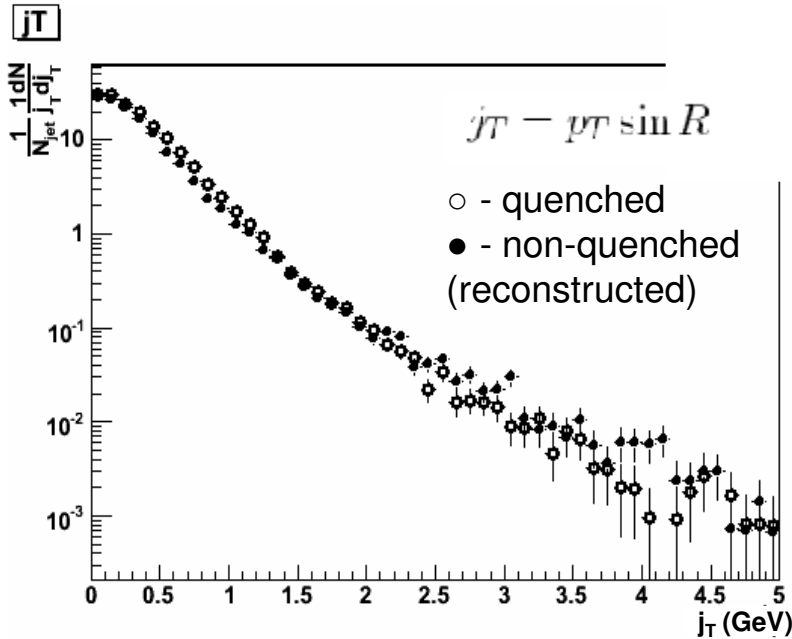
- tracks are matched to calorimeter towers of a jet
- $j_T$  and  $z$  for tracks above 2 GeV is computed
- background distributions of  $j_T$  and  $z$  are computed using tracks that match with HIJING particles, these distributions are subtracted, correction for the jet position resolution is applied



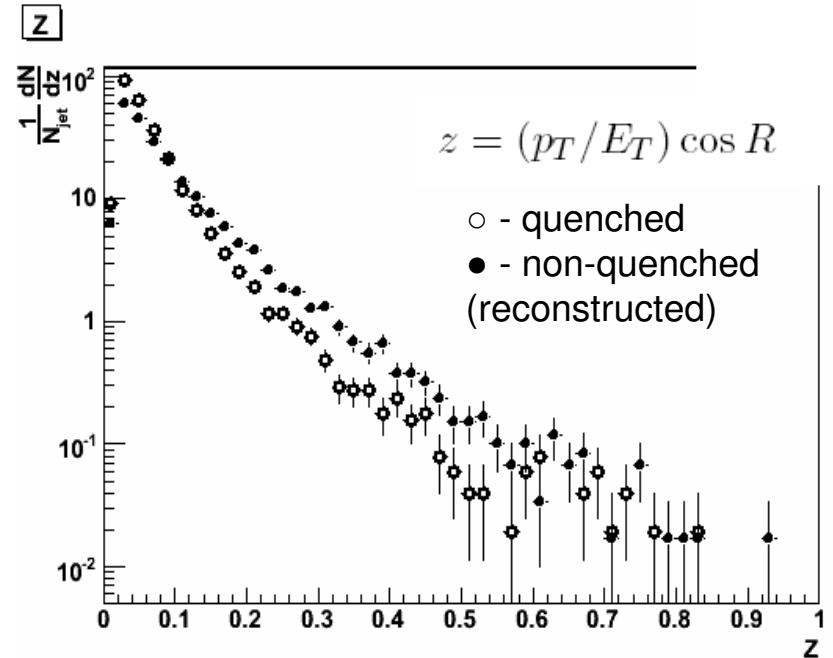
... we can well reproduce  $j_T$  distribution and fragmentation function



# Fragmentation function and $j_T$ from PYQUEN – reconstructed



Large  $j_T$  suppressed  $\Leftrightarrow$  gluons radiated from large angles



Low  $z$  enhanced, higher  $z$  suppressed  $\Leftrightarrow$  leading particle suppressed, redistribution of energy out of a jet core

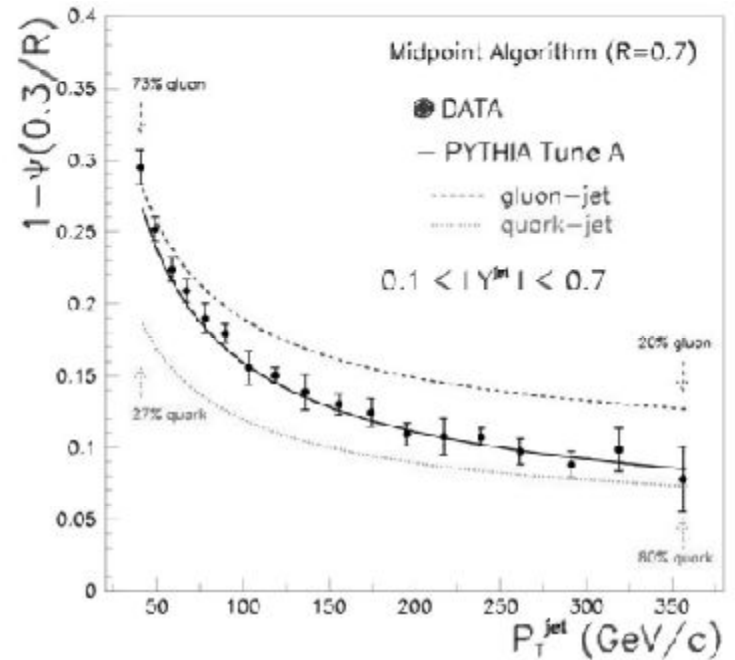
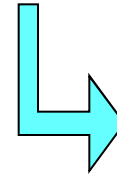
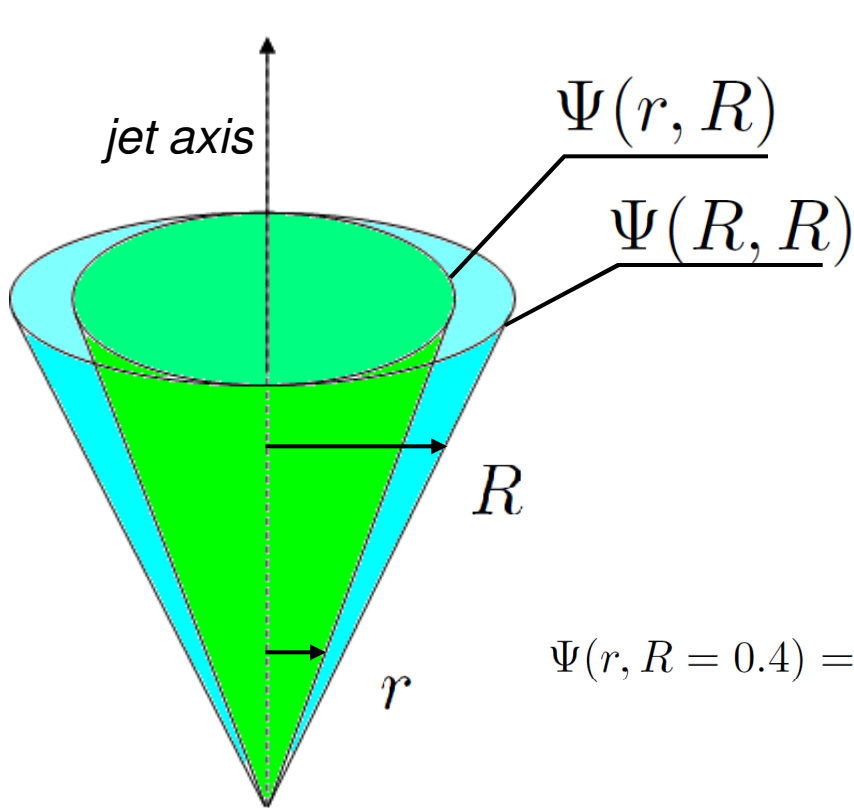
... if the quenching is of that order we should be able to measure it



# Jet shapes



- Measure the energy flow inside the jet at the calorimeter level
- Well defined in QCD, measured at Tevatron



$$\Psi(r, R = 0.4) = \frac{\int_0^r E_T(\rho) d\rho}{\int_0^R E_T(\rho) d\rho}$$

... integral jet shape

$$\psi(r, R = 0.4) = \frac{d\Psi(r, R)}{dr}$$

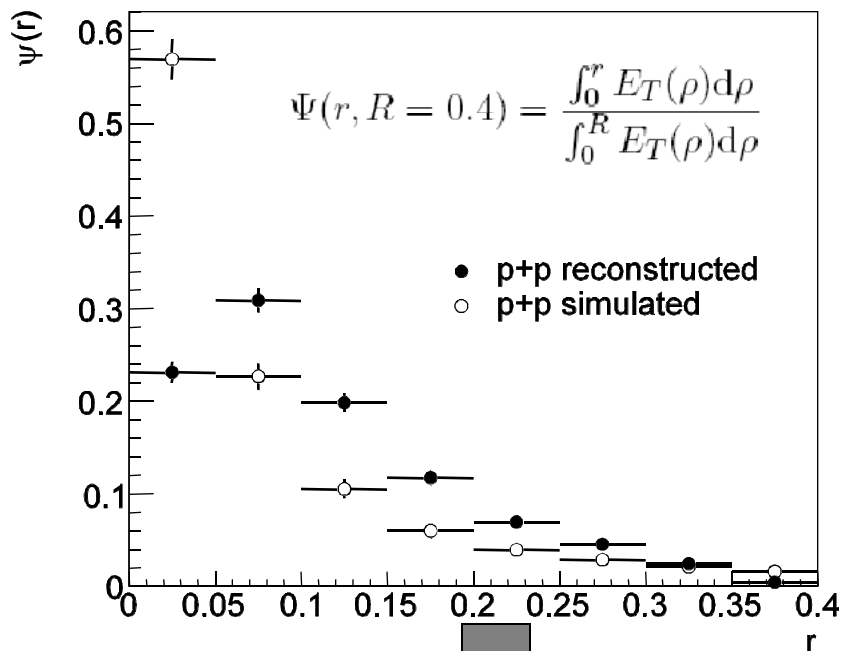
... differential jet shape



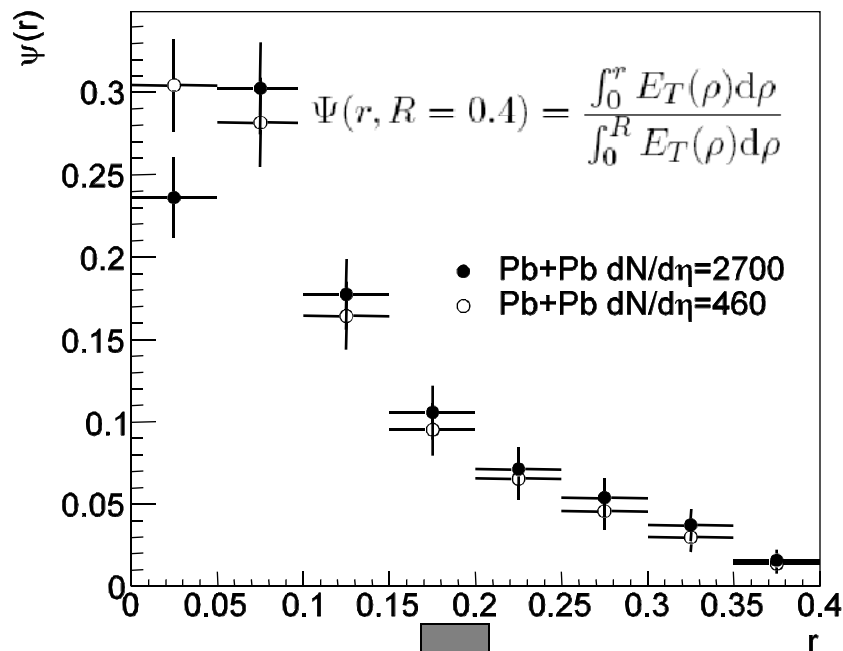
# Problems of jet shapes



## 1. Effect of the calorimeter



## 2. Effect of the background (only applying subtraction is not enough)



How to match “closed” to “open”?

Deconvolution

Better estimation of the jet position

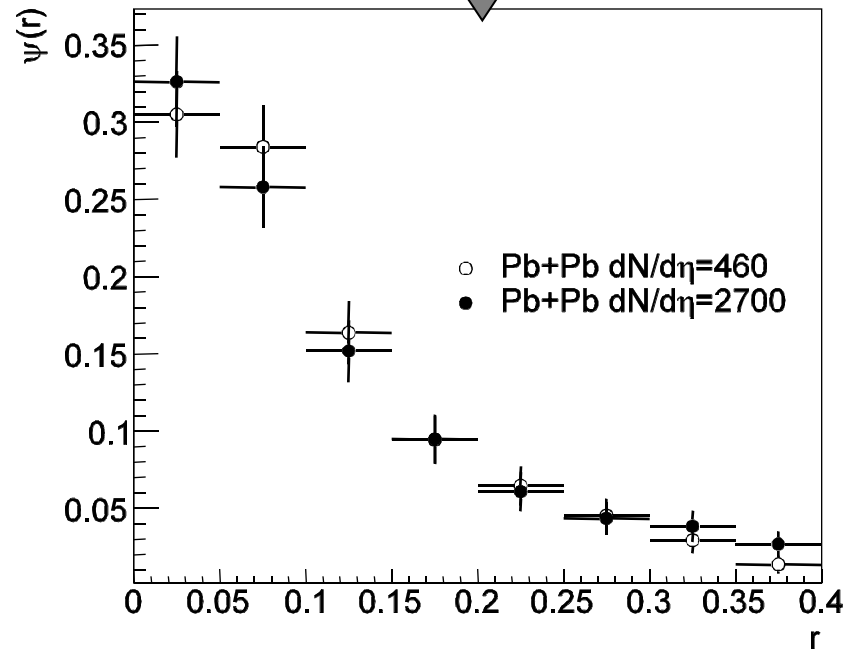
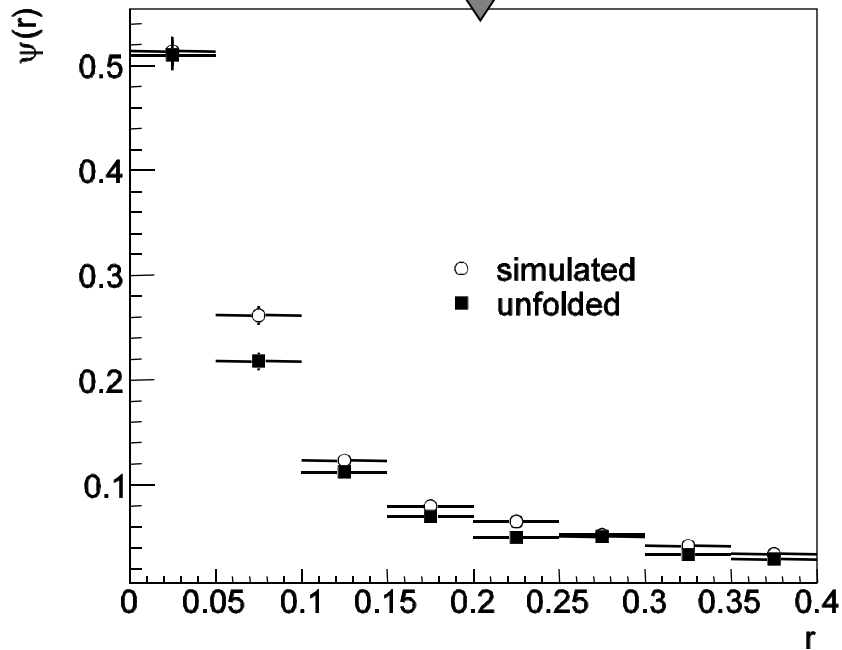


# Problems of jet shapes



Deconvolution:  
Bayesian unfolding

Better estimation of the jet  
position: using the method  
of smaller cones



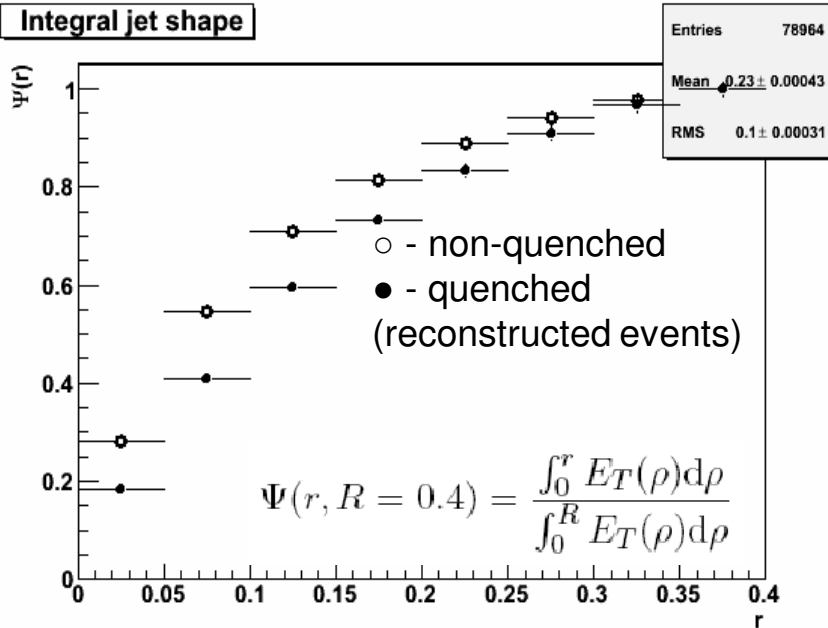




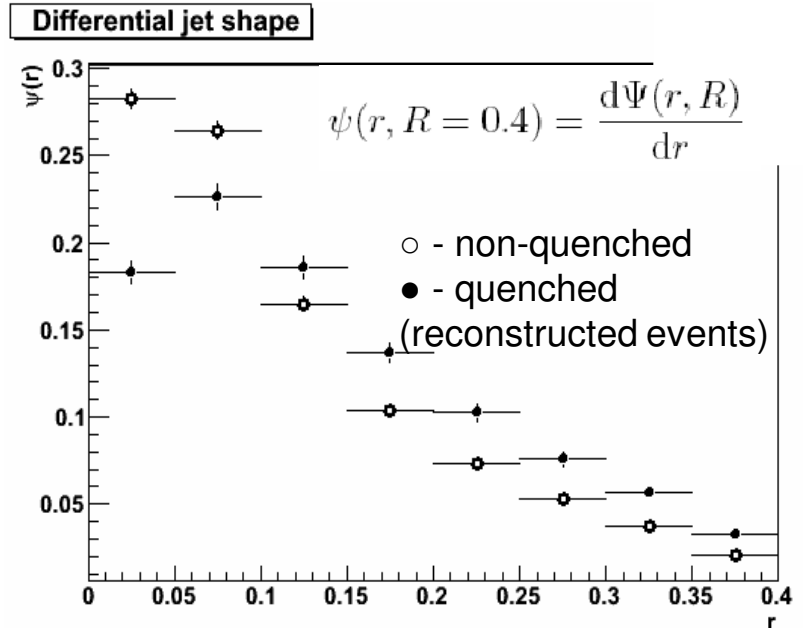
# Jet shapes from PYQUEN – reconstructed



Integral jet shape



Differential jet shape



- Result at the **after the reconstruction**
- Jet quenching effect still well visible

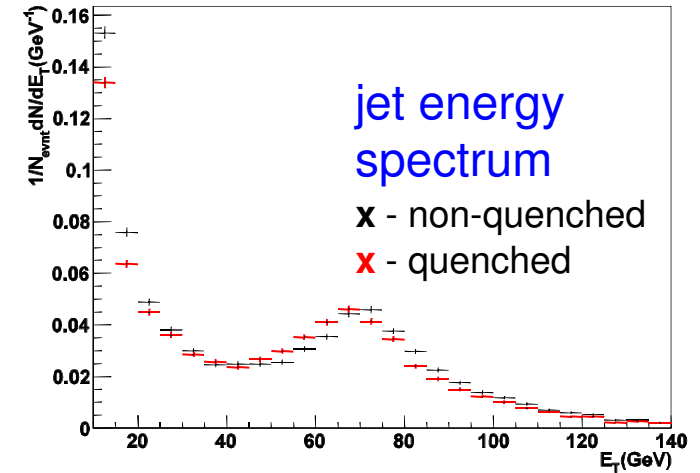
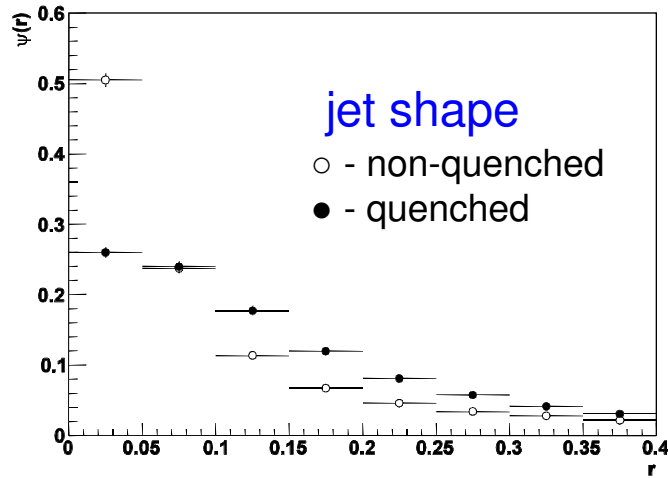
... if the quenching is of that order we should be able to measure it



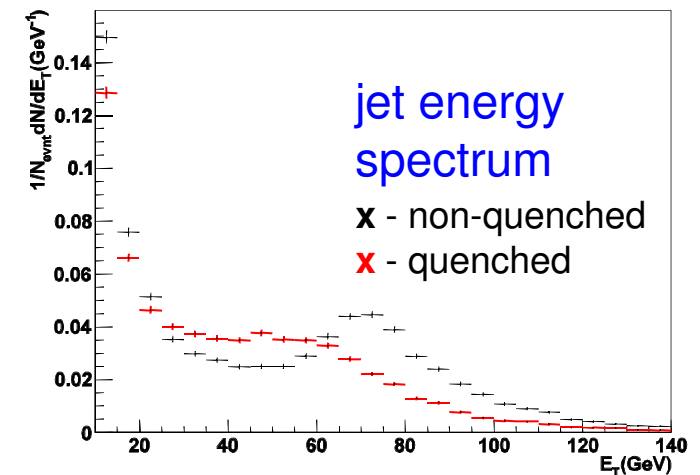
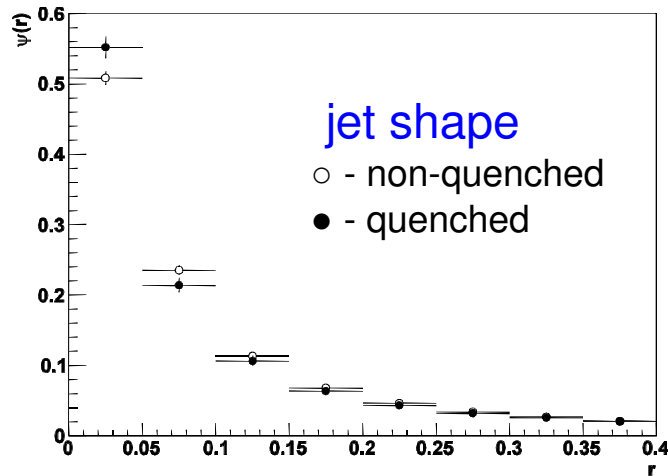
# Complementary quantities to measure different energy loss mechanisms



PYQUEN  
with only  
radiative  
energy loss



PYQUEN  
with only  
collisional  
energy loss





# Summary



- We didn't talk about:
  - di-jet azimuthal correlations
  - $\gamma$ -jet correlations
  - $p_{OUT}$ , radial moments, ...  
... all these quantities can be used to distinguish different energy loss scenarios.
- ATLAS is very good detector for the jet heavy ion physics (it is also good detector for soft heavy ion physics, quarkonia, and other important measurements).
- We have in our hands tools for reconstruction and analysis of heavy ion data.
- Bearing in mind that p+p with full luminosity is expected to be similar to peripheral Pb+Pb collisions, we can say that these tools could be useful also for p+p measurements.



# Backup slides



# Backup slides



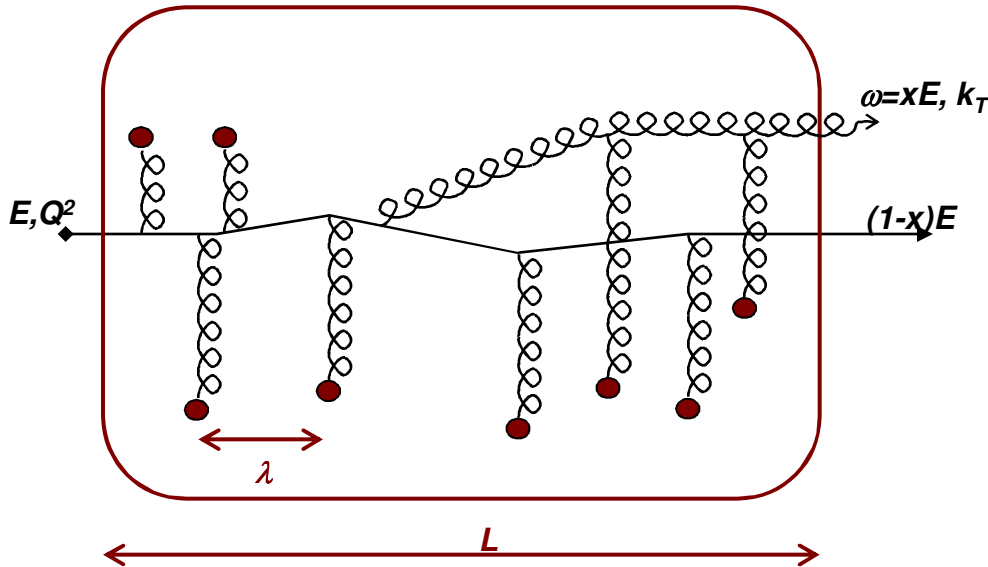
## Path integral in opacity (ASW)

$$\epsilon = \frac{1}{E_q} \int_0^\infty \omega \frac{dI_{med}}{d\omega} d\omega$$



$$D_{q \rightarrow h}^{med}(x, Q^2) = \int_0^1 dc P(c) \frac{1}{1-\epsilon} D_{q \rightarrow h} \left( \frac{x}{1-\epsilon}, Q^2 \right)$$

$$\begin{aligned} \omega \frac{dI_{med}}{d\omega dk} &= \frac{\alpha_s C_F}{(2\pi)^2 \omega^2} 2\text{Re} \int_0^\infty dy_l \int_{y_l}^\infty d\bar{y}_l \int d^2 u \\ &\times \epsilon^{-i\mathbf{k}_l \cdot \mathbf{u}} \epsilon^{\frac{1}{2} \int_{\bar{y}_l}^\infty d\xi n(\xi) \sigma(\mathbf{u})} \frac{\partial}{\partial \mathbf{y}} \cdot \frac{\partial}{\partial \mathbf{u}} \\ &\times \int_{\mathbf{y}=\mathbf{r}(y_l)}^{\mathbf{u}=\mathbf{r}(\bar{y}_l)} D\mathbf{r} \exp \left[ i \int_{y_l}^{\bar{y}_l} d\xi \frac{\omega}{2} \left( \dot{\mathbf{r}}^2 - \frac{n(\xi) \sigma(\mathbf{r})}{i2\omega} \right) \right] \end{aligned}$$



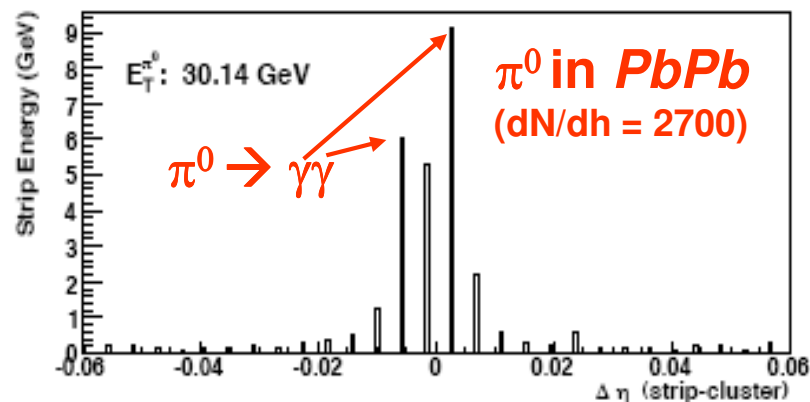
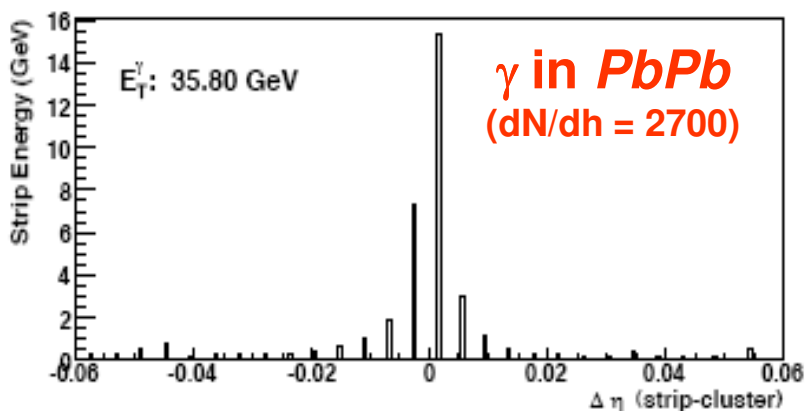
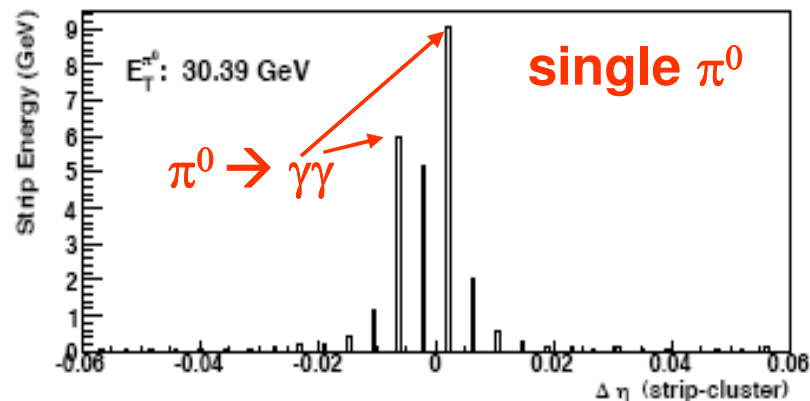
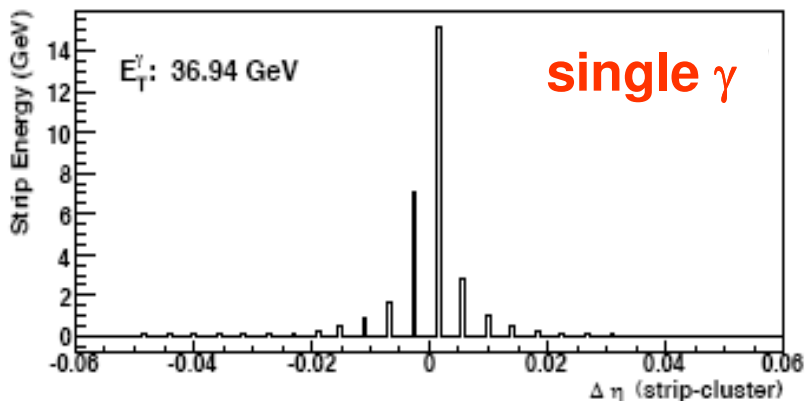
**Gluon energy spectrum can be computed in 2 limits:**

- few hard scatterings (GLV)
- many soft scatterings (BDMPS)

**Other approaches: Higher Twists (direct computation of fragmentation function – formally similar to DGLAP), Finite Temperature approach (AMY), techniques which use AdS/CFT correspondence**

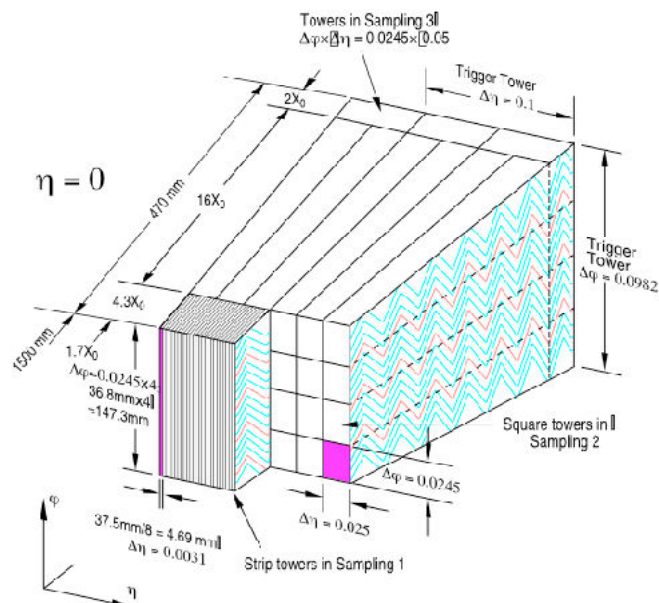


# Backup slides: $\gamma$ -isolation





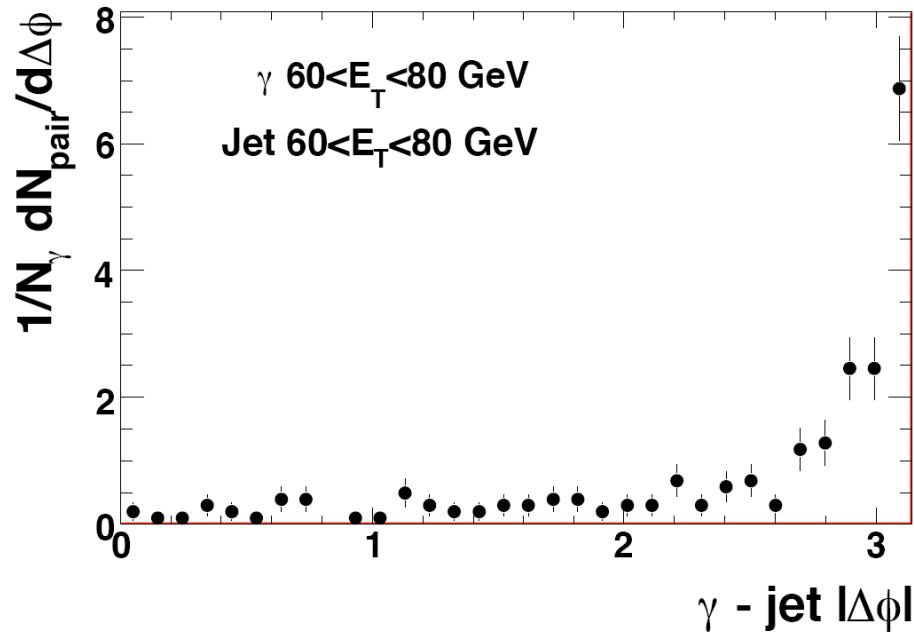
# Backup slides: $\gamma$ -isolation



- can benefit from excellent longitudinal segmentation (0.003 for the first sampling of EMCAL)
- a set of cuts to distinguish a direct photons from neutral hadrons – based on the shower shape
- double peaked or wide showers rejected
- the most important cuts:
  - the fraction of energy in 8 strips around the core of 6 strips
  - energy of second peak minus the minimum energy between the two peaks



# Backup slides: $\gamma$ -jet correlations



- $\gamma$ -jet correlation measurements can help jet analysis at low  $E_T$ , can be used for the fake rejection
- important for in-medium fragmentation studies



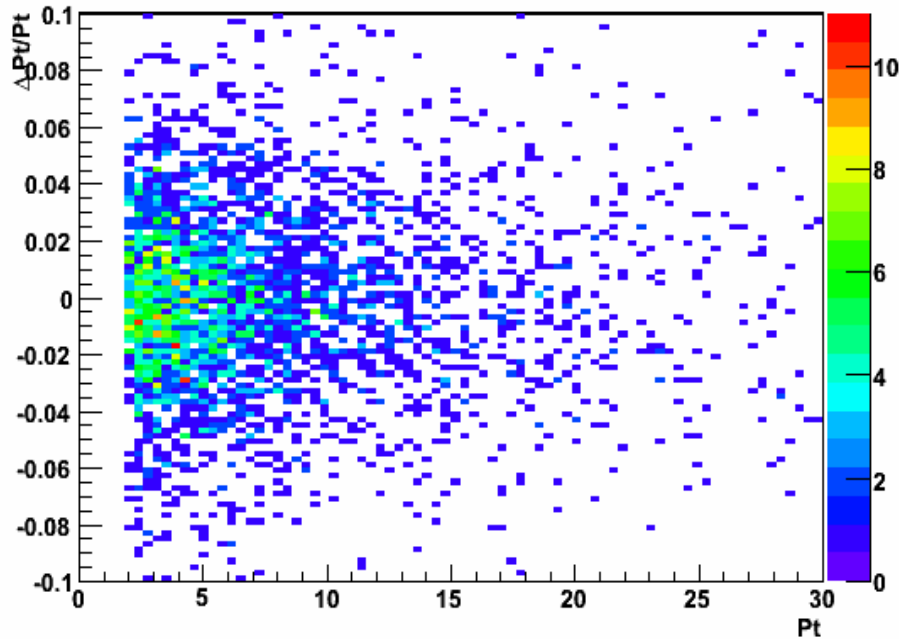


# Backup slides: tracking performance



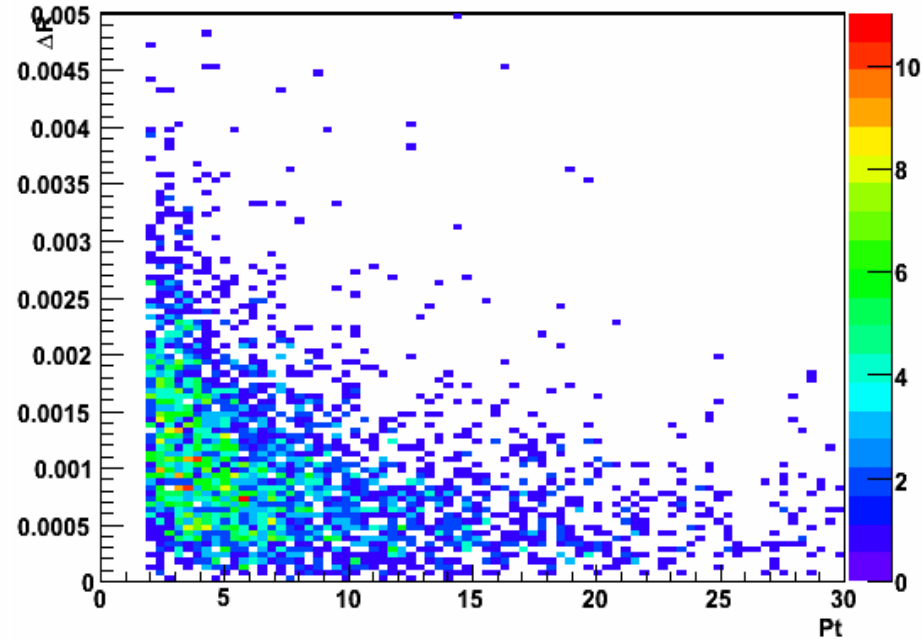
Pt vs.  $\Delta Pt/Pt$

Integral 3268



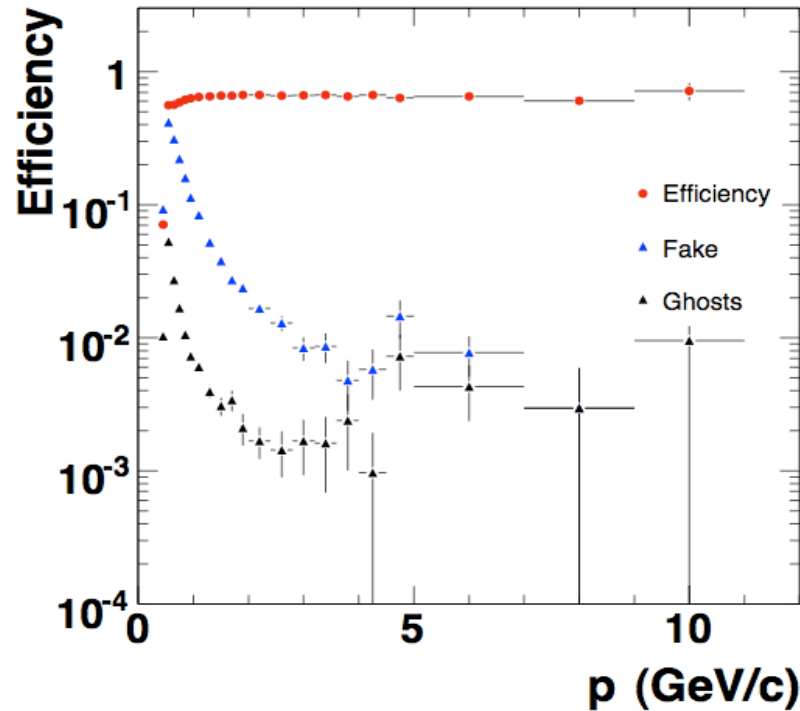
Pt vs.  $\Delta R$

Integral 3190





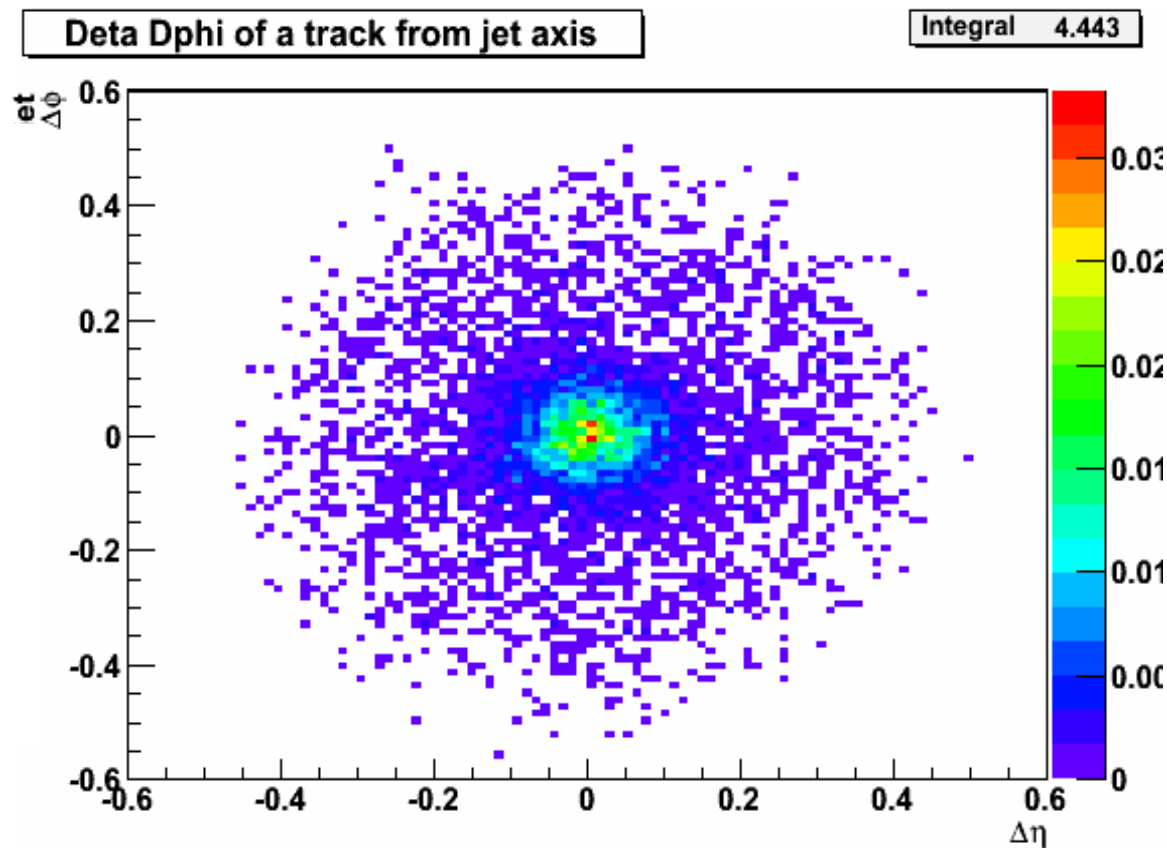
# Backup slides: tracking efficiency



Tracking efficiency  $\sim 70\%$  for the most central collisions



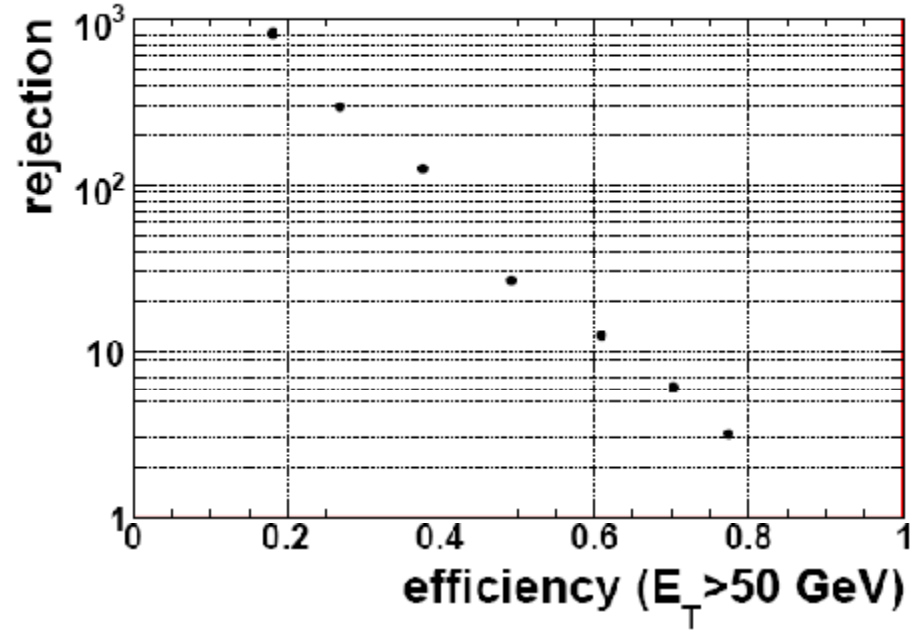
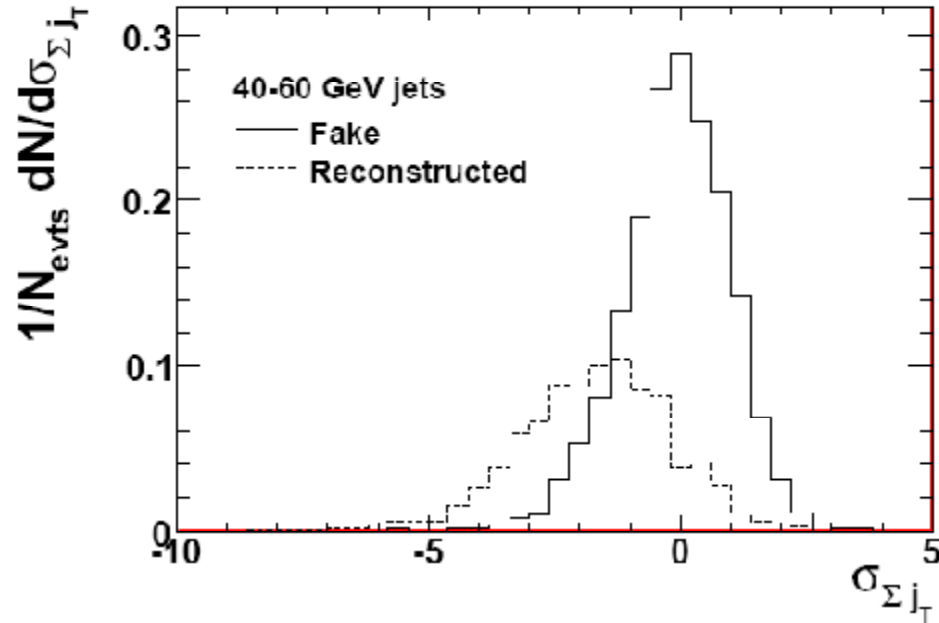
# Backup slides: jet's tracks



Shape of the jet from tracking



# Backup slides: fake jet rejection

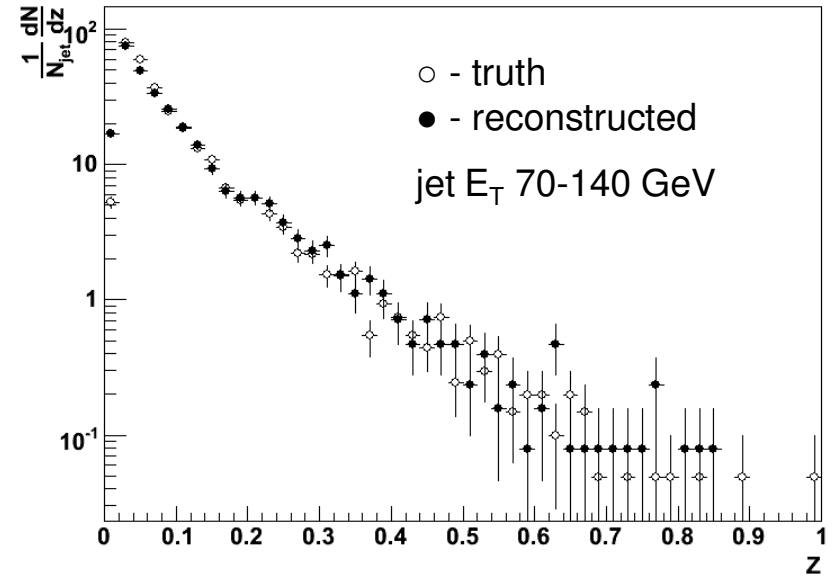
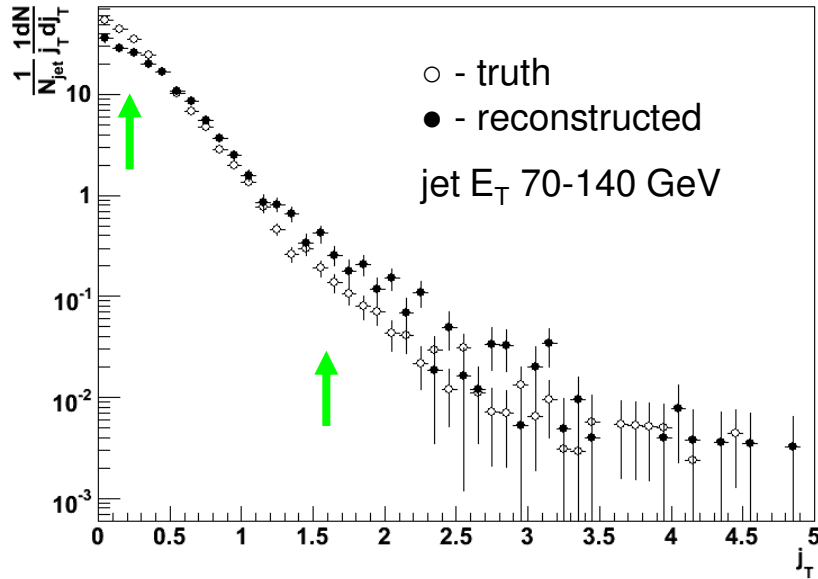


$$\lambda j_T = \sum_{\text{cell} \in \text{jet}} p_{j_T, \text{cell}} \sin R_{\text{cell}}$$

$$\sigma_{\Sigma j_T} = \frac{\Sigma j_T - \langle \Sigma j_T \rangle (E_T)}{\sigma(E_T)}$$



# Backup slides: Sensitivity to the jet position



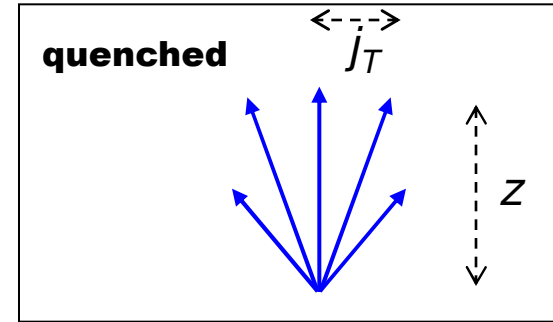
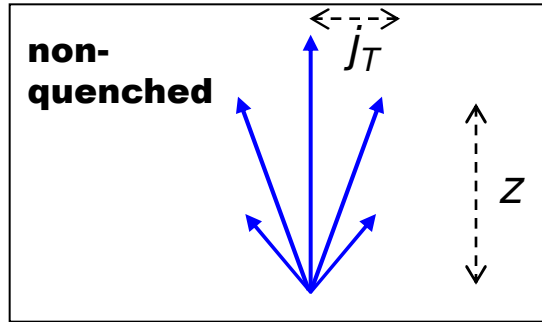
Background distributions subtracted but the correction on the jet position resolution not applied => visible underestimation at small  $j_T$  ( $z$  is not affected much by the jet position resolution). Discrepancy can be removed using jet position determined by the cone reconstruction with smaller  $R$ .



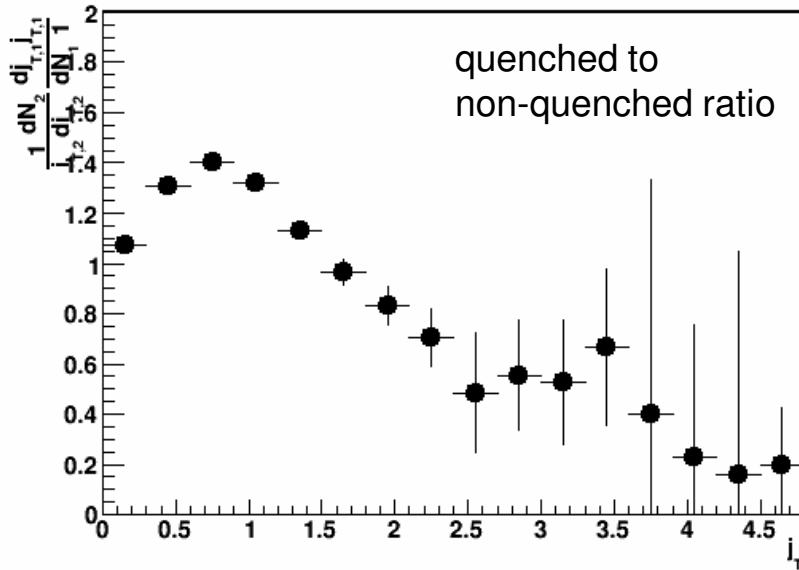
# Backup slides: PYQUEN simulations - ratios



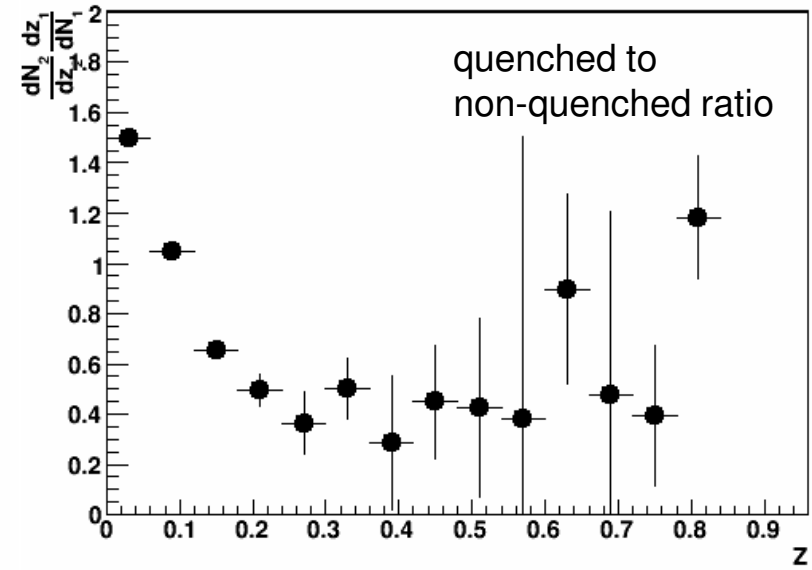
illustration ...



**jT**

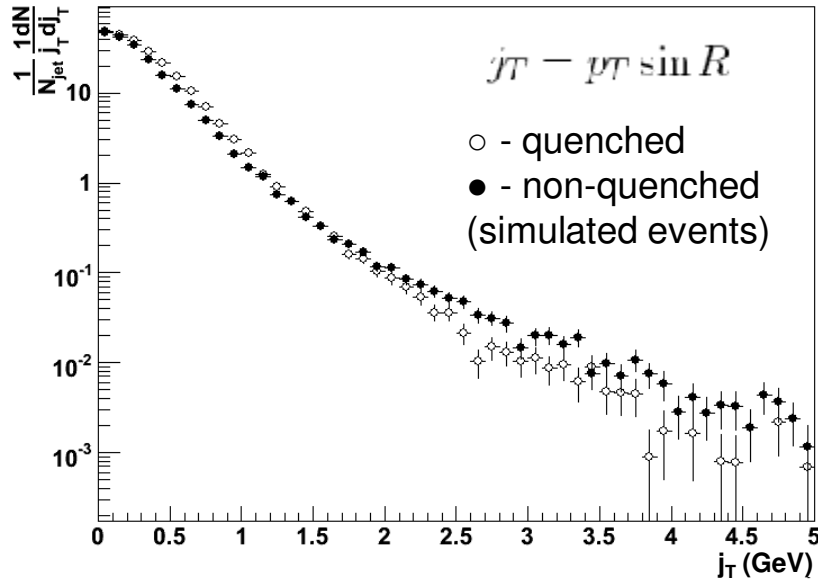


**Z**

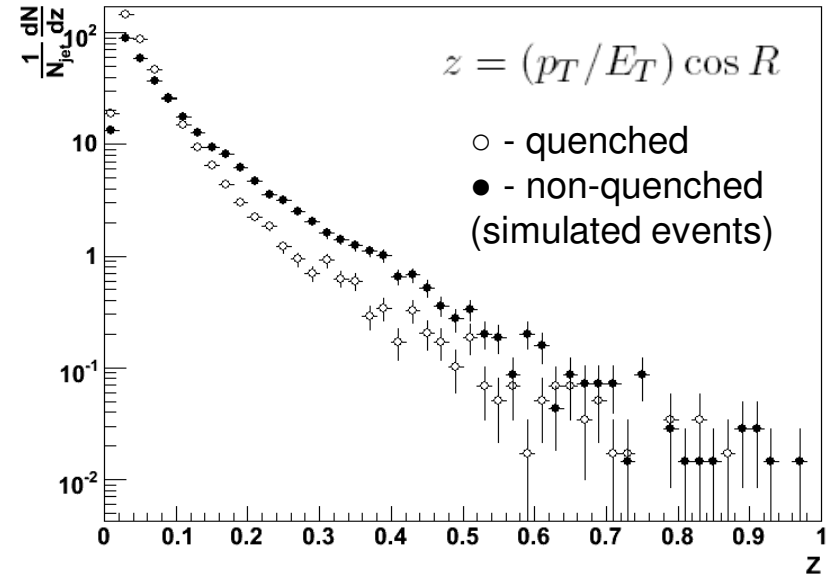




# Fragmentation function and $j_T$ from PYQUEN – generator level



Large  $j_T$  suppressed  $\Leftrightarrow$  gluons radiated from large angles



Low  $z$  enhanced, higher  $z$  suppressed  $\Leftrightarrow$  leading particle suppressed, redistribution of energy out of a jet core

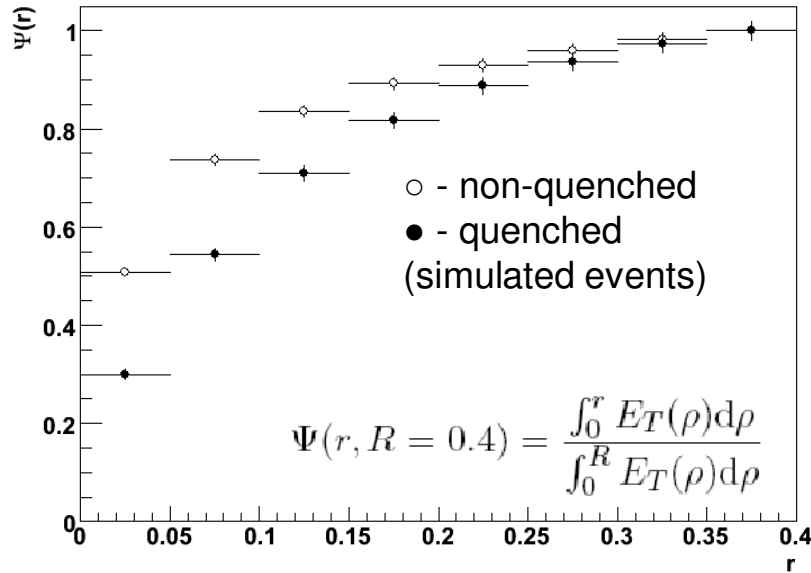
- Result at the **generator level**
- PYQUEN settings: default setting for quenching,  $b=0$ ,  $p_{T,min}=70$  GeV,  $PbPb$ , LHC energy



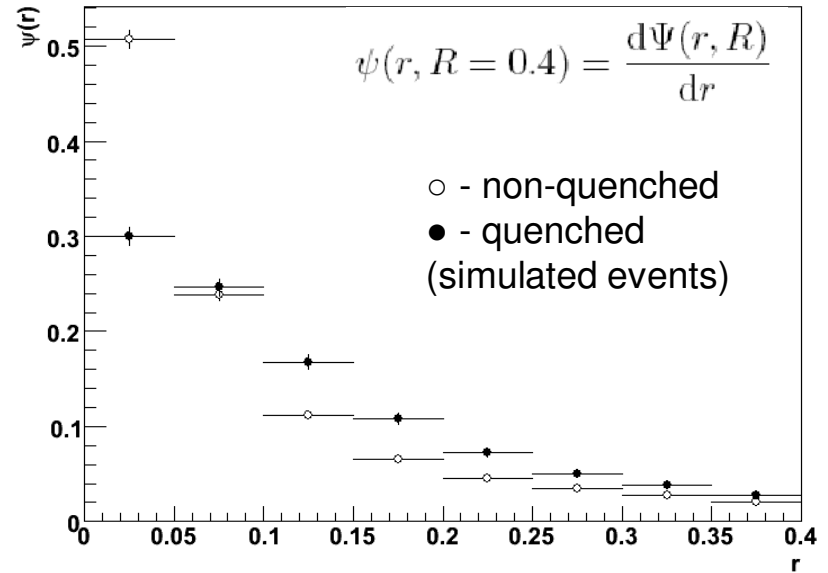
# Jet shapes from PYQUEN – generator level



Integral jet shape



Differential jet shape



- Result at the **generator level**
- Almost factor of two in the jet core
- Differential jet shape can show better the flow of the energy – energy is redistributed out of center of the jet

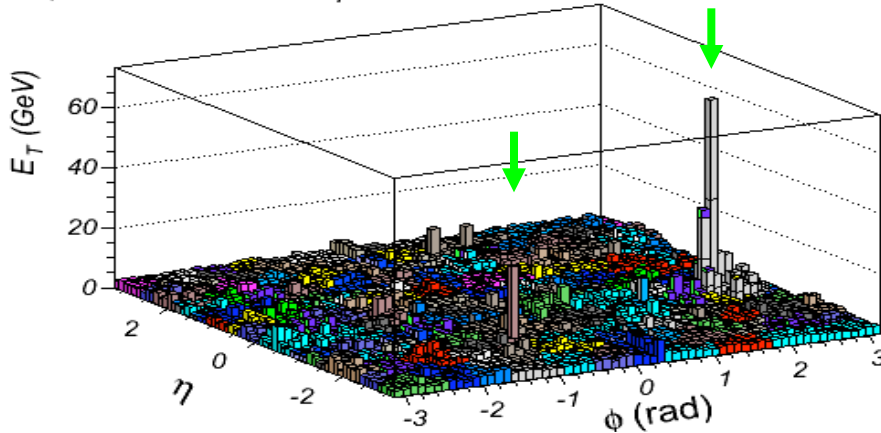




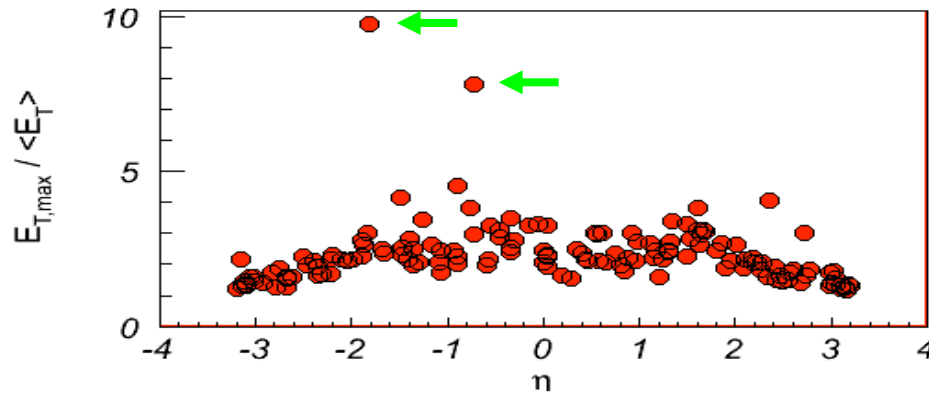
# Fast $k_T$ jet reconstruction strategy



Pythia+HIJING,  $K_T$  Finder  $R=0.4$ , EMB/EC



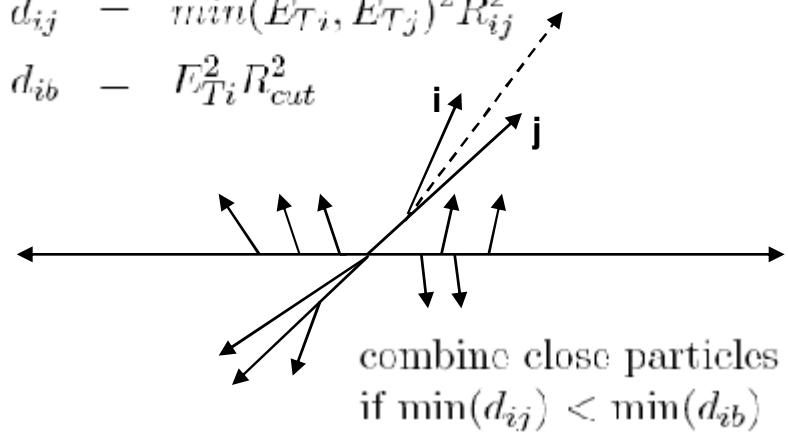
Pythia+HIJING,  $K_T$  Finder  $R=0.4$ , EMB/EC



## $k_T$ algorithm

$$d_{ij} = \min(E_{T_i}, E_{T_j})^2 R_{ij}^2$$

$$d_{ib} = E_{T_i}^2 R_{cut}^2$$



... reconstructs jets backwards along fragmentation chain

## $k_T$ jet reconstruction:

- run fast  $k_T$  algorithm
- separate jet from the background
- subtract background from jets
- calibrate jet energy



# Backup slides

