

# Cooling Methods

AVA International School of Low Energy Antimatter Physics 2018  
Sebastian Gerber  
Cern, Politecnico di Milano

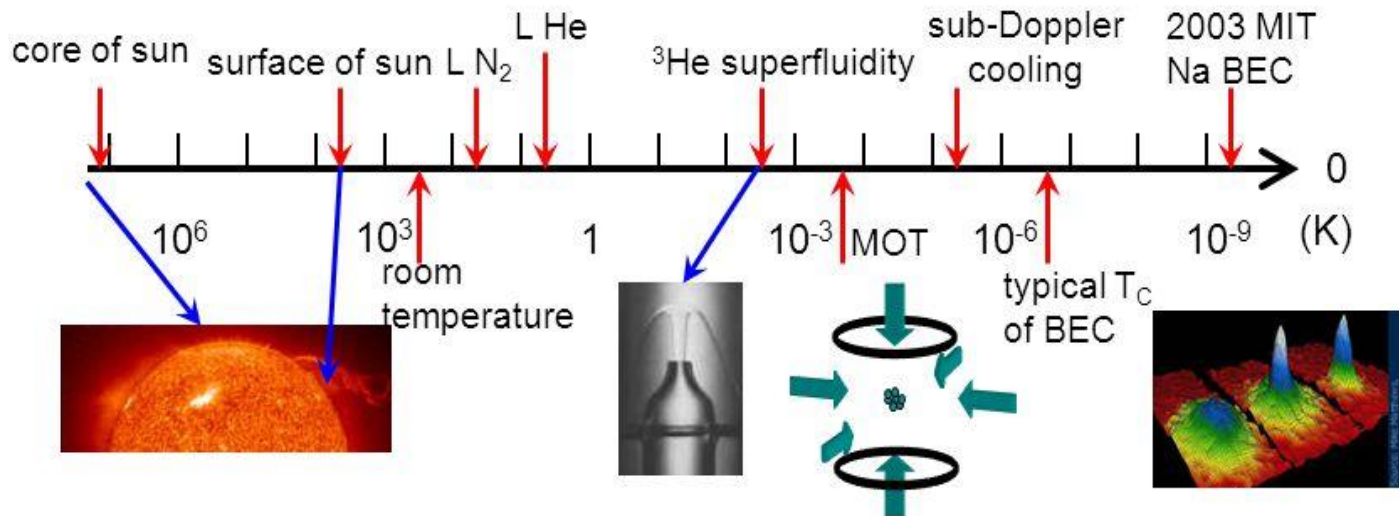


# Motivation for cooling $\bar{p}$



Precision experiments need well defined boundary conditions of specimens

- Isolated  $\bar{p}$  (Penning trap in vacuum  $<10^{-11}$  mbar)
- Small particle numbers require low temperature to gather statistics
- Low  $\bar{p}$  temperature (Doppler broadening of  $\bar{H}$  internal transitions / de Broglie for gravity exp.)

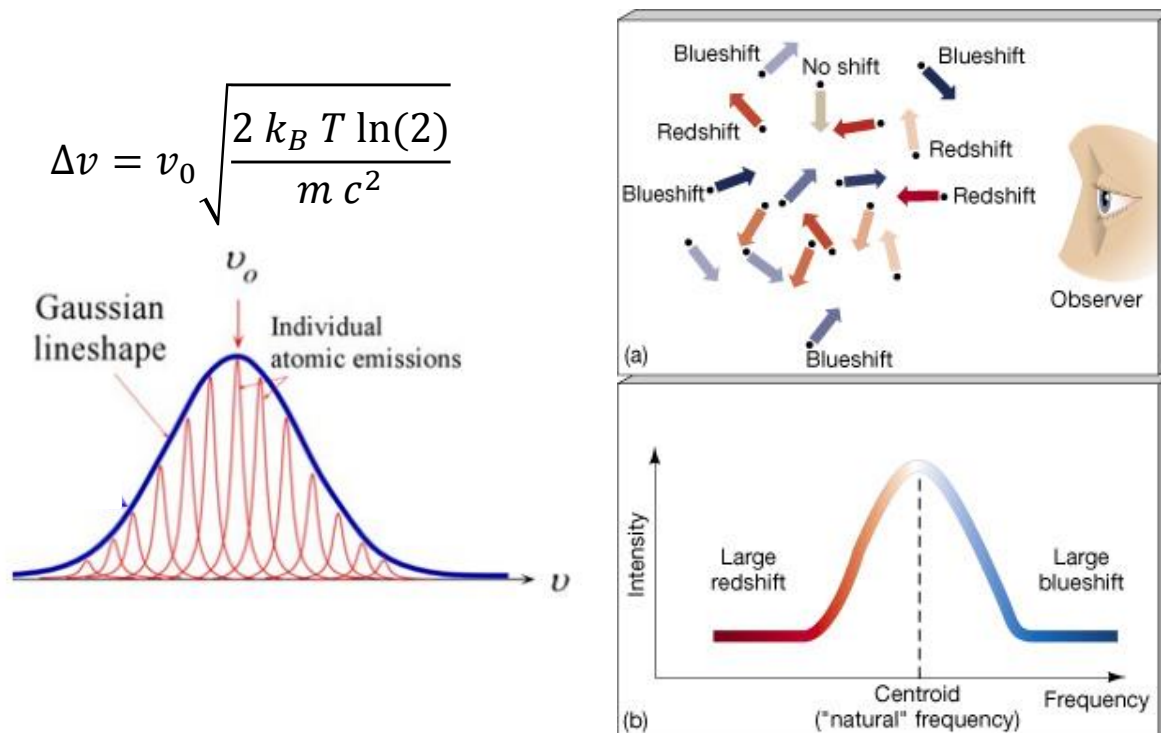


# Motivation for cooling $\bar{p}$



Precision experiments need well defined boundary conditions of specimens

- Isolated  $\bar{p}$  (Penning trap in vacuum  $<10^{-11}$  mbar)
- Low  $\bar{p}$  temperature (**Doppler broadening of  $\bar{H}$  internal transitions** / de Broglie for gravity exp.)



Spectroscopy on  $\bar{H}$  1s-2s at 0.54 K,  $\Delta v/v=10^{-13}$  M. Ahmadi et al., *Nature*, **541**, 506 (2017)

Spectroscopy on  $\bar{H}$  hyperfine at 0.54 K,  $\Delta v/v=10^{-13}$  M. Ahmadi et al., *Nature*, **548**, 66 (2017)

Spectroscopy on H 1s-2s  $\Delta v/v=10^{-15}$  A. Matveev et al., *Phys. Rev. Lett.*, **110**, 230801 (2013)

# Motivation for cooling $\bar{p}$

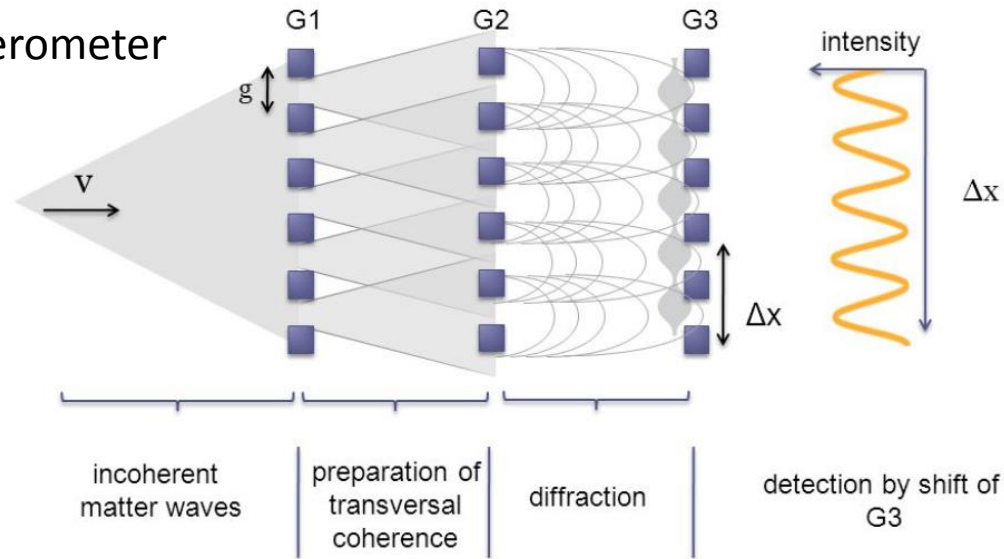


Precision experiments need well defined boundary conditions of specimens

- Isolated  $\bar{p}$  (Penning trap in vacuum  $<10^{-11}$  mbar)
- Low  $\bar{p}$  temperature (Doppler broadening of  $\bar{H}$  internal transitions / **de Broglie for gravity exp.**)

$$\lambda = \frac{h}{m v}$$

Talbot-Lau atom interferometer



Currently no experimental WEP test available for antimatter with high precision

normal matter  $\Delta g/g$ : H. Mueller et al., *Nature*, **463**, 926 (2010)

normal matter WEP: S. Baessler et al., *Phys. Rev. Lett.*, **83**, 3583 (1999)



- Cooling in the Antiproton decelerator
  - Stochastic cooling
  - Electron cooling
  
- Cooling in Penning traps
  - Electron cooling
  - Laser cooling



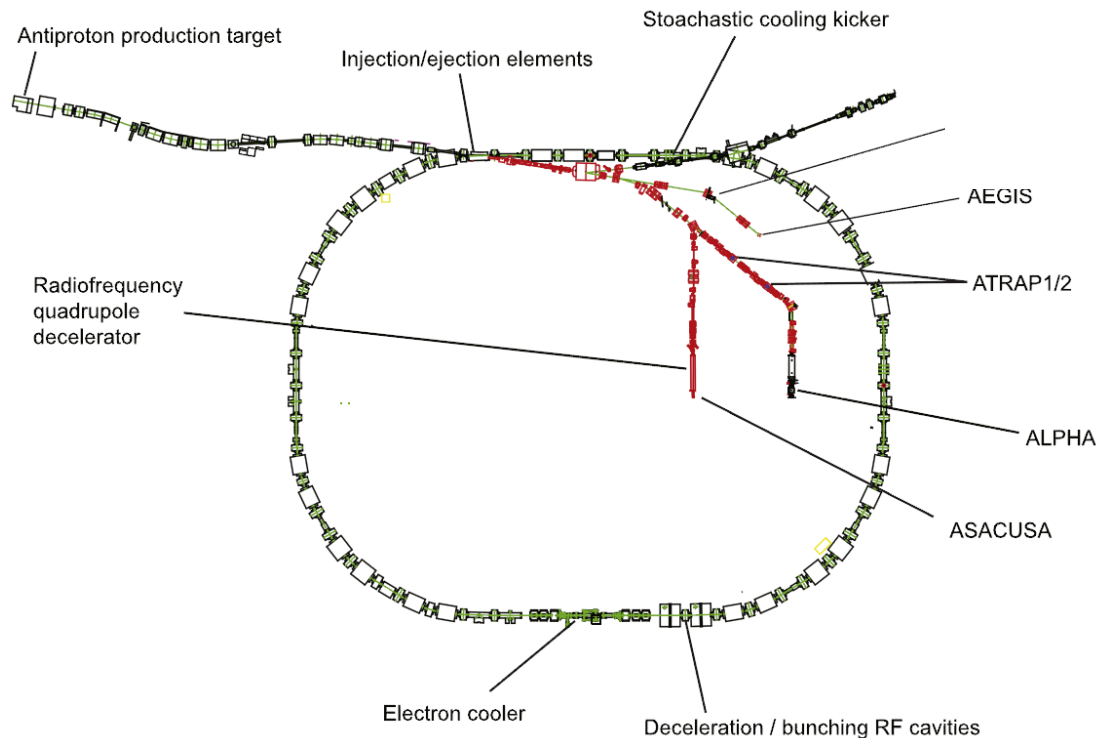
# Antiproton Decelerator (AD)



- Replacement of LEAR (shutdown 1996)
- AD started operation in 1999

## Specs:

- Produces  $\bar{p}$  of 5.3 MeV energy,
- Cycle time  $\sim 100$  s,
- Pulsed beams with  $>3 \times 10^7 \bar{p}$ ,
- Emittance  $\sim 0.3\pi$  mm mrad



# Antiproton Decelerator (AD)



- Replacement of LEAR (shutdown 1996)      Specs:
- A

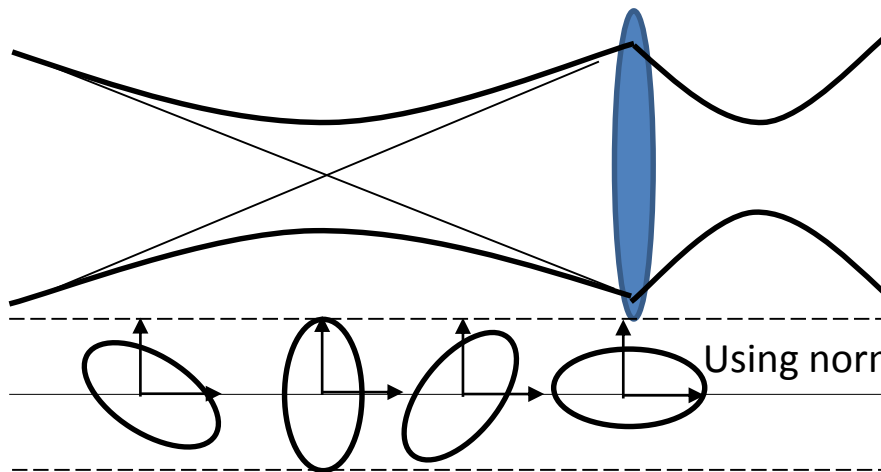
What is emittance?

Liouville's equation.

$$\frac{d\rho}{dt} = \frac{\partial \rho}{\partial t} + \sum_{i=0}^n \left( \frac{\partial \rho}{\partial q_i} \dot{q}_i + \frac{\partial \rho}{\partial p_i} \dot{p}_i \right) = 0 \quad \frac{d\rho}{dt} = -\{\rho, H\}$$

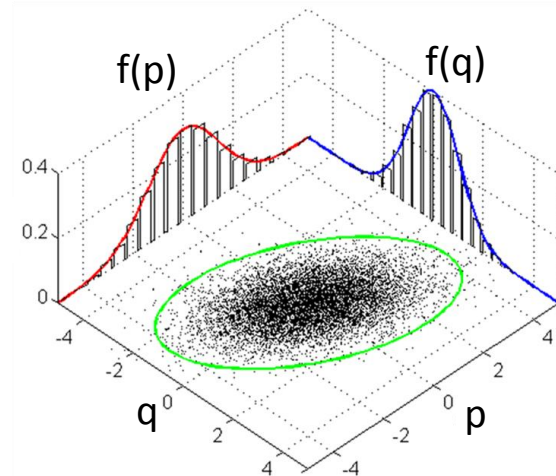
Theorem: Emittance can not be changes by lenses

*Emittance only changeable by:  
radiation damping, stochastic and electron cooling.*



Using normal electrostatic lens, area stays the same

coordinate  $q$  – momentum  $p$   
phase space  $\rho$  at  $t_0$



Emittance=Area

Hz



## 1. Antiproton production

$$p(26 \text{ GeV}/c) + p_{\text{Ir-target}} \rightarrow p + p + p + \bar{p} (3.6 \text{ GeV}/c)$$

$\bar{p}$  emittance  $200 \pi \text{ mm mrad}$   
 $\Delta p/p$  momentum spread  $\sim 6\%$

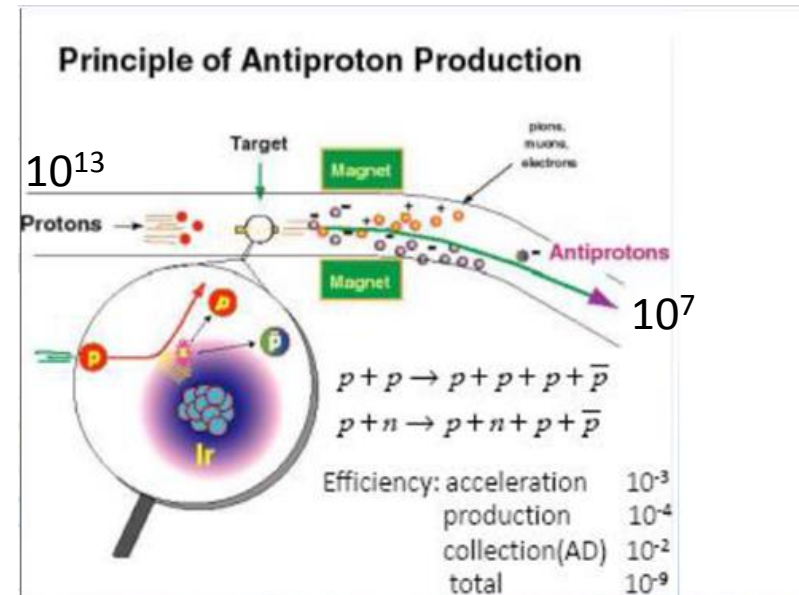
## 2. RF bunch rotation:

stretches  $\bar{p}$  bunch from 30 m to 190 m  
new  $\Delta p/p \sim 1.5\%$  ( $L \cdot \Delta p/p$  is conserved)

## 3. $\bar{p}$ cooling energy cascade of

RF cavity deceleration, **stochastic** and **electron** cooling:

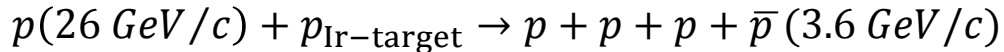
Target:  $3.5 \text{ GeV}/c \rightarrow$  AD cooling:  $100 \text{ MeV}/c \rightarrow 5.3 \text{ MeV}/c$







## 1. Antiproton production



$\bar{p}$  emittance  $200 \pi \text{ mm mrad}$

$\Delta p/p$  momentum spread  $\sim 6\%$

## 2. RF bunch rotation:

stretches  $\bar{p}$  bunch from 30 m to 190 m

new  $\Delta p/p \sim 1.5\%$  ( $L \cdot \Delta p/p$  is conserved)

## 3. $\bar{p}$ cooling energy cascade of

RF cavity deceleration, **stochastic** and **electron** cooling:

Target:  $3.5 \text{ GeV}/c \rightarrow$  AD cooling:  $100 \text{ MeV}/c \rightarrow 5.3 \text{ MeV}/c$

Why?

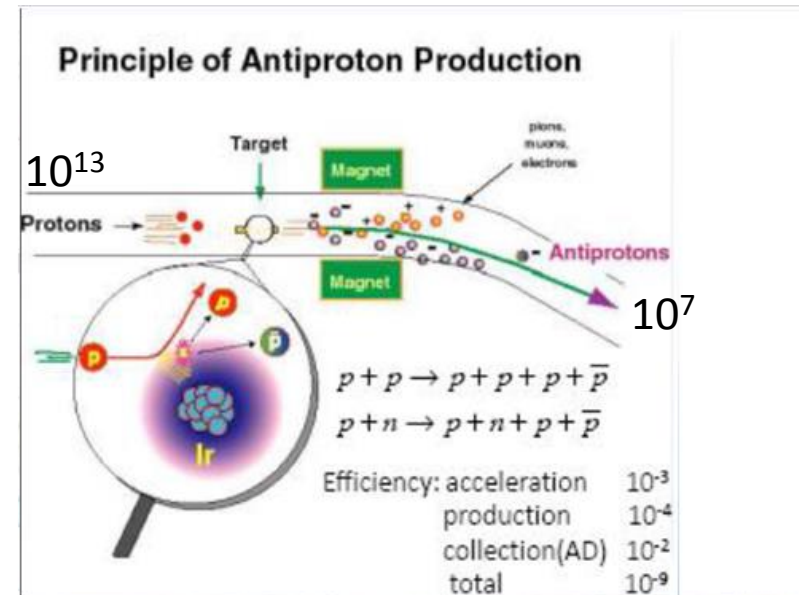
$E \sim 3.6 \text{ GeV}$  must be reduced by  $\sim 10^6$  before  $\bar{p}$  are trappable in Penning trap (10 kV electrodes)

$E \sim 3.6 \text{ GeV}$  must be reduced by  $\sim 10^{10}$  before  $\bar{H}$  production can start ( $\sim 300 \text{ K}$ )

$E \sim 3.6 \text{ GeV}$  must be reduced by  $\sim 10^{13}$  before produced  $\bar{H}$  is trappable in magnetic trap ( $\sim 0.5 \text{ K}$ )

$E \sim 3.6 \text{ GeV}$  must be reduced by  $\sim 10^{15}$  before a future  $\bar{H}$  atomic fountain (nK)

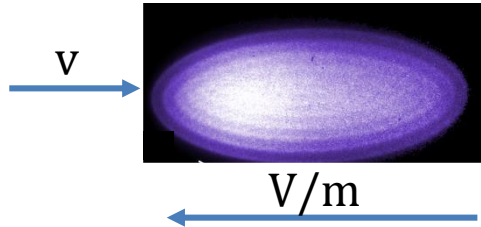
$\rightarrow$  Importance of  $\bar{p}$  cooling





Why is stochastic cooling needed?

A simple deceleration of a cloud  $N$  of  $\bar{p}$  would lead to an increase in phase-space density  $D$  [1]



$$D = \frac{N}{\sqrt{E_h E_v} L \Delta p/p}$$

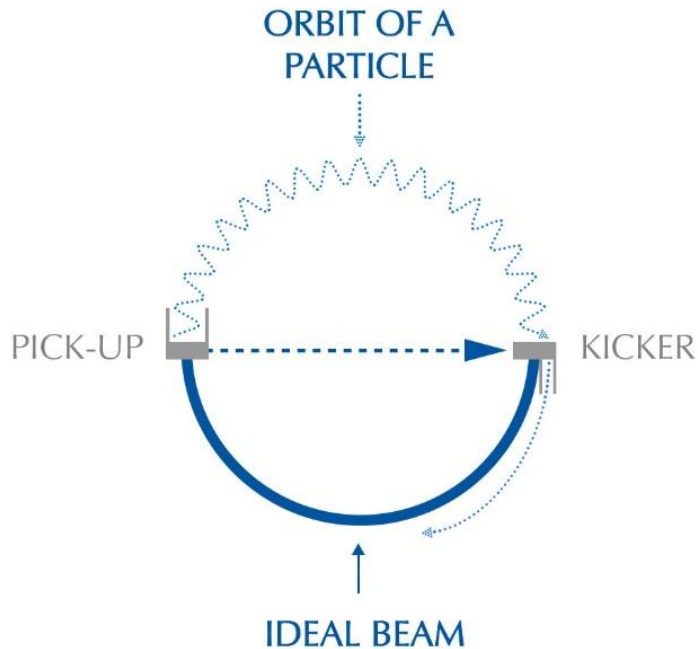
$E_h, E_v$  ...horizontal, vertical emittances  
 $L$  ...longitudinal spread  
 $N$ ... Number of particles  
 $\Delta p/p$  ...momentum spread

# Stochastic cooling



Initially wide spread of momentum and angle emission at  $\bar{p}$  production.

As particles travel around, a detector or "pick up" measures their motion and sends a signal across the ring to a corrector, the kicker, which adjusts their angles.



## Pick up electrode:

Detect  $\Delta p_i$  and  $\Delta x_i$  of  $\bar{p}$  subgroups  $i$ , relative to ideal orbiting  $\bar{p}$

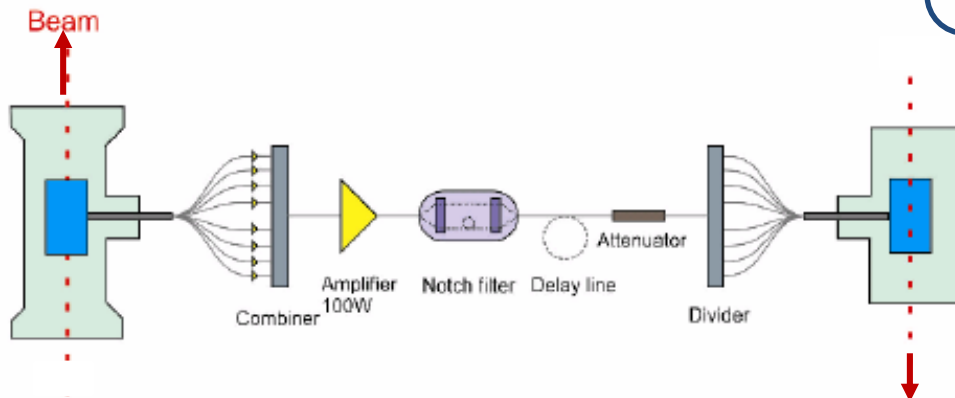
## Kicker:

Apply corrective pulses to electrodes

## Repeating steps:

$\Delta p/p \sim 0.07\%$  from initially  $\sim 1.5\%$

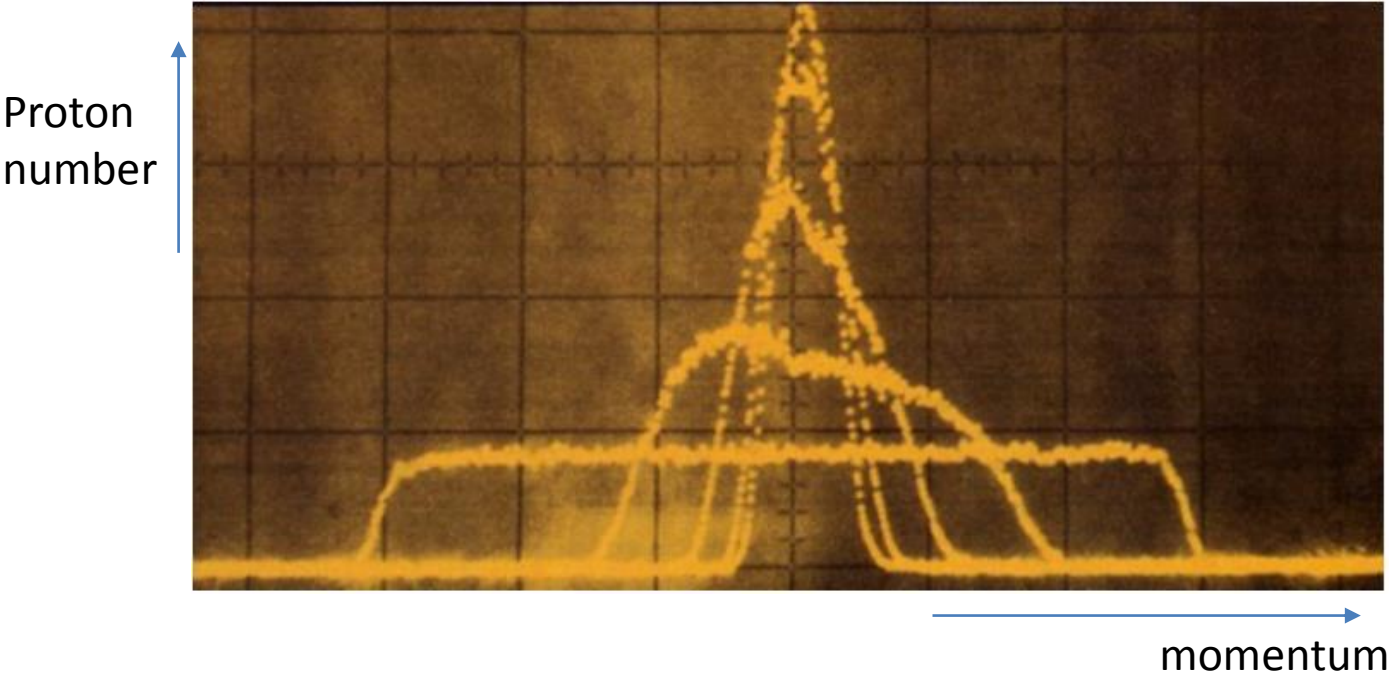
Emittance  $\sim 3\pi$  mm mrad from initially  $\sim 200\pi$



# Stochastic cooling



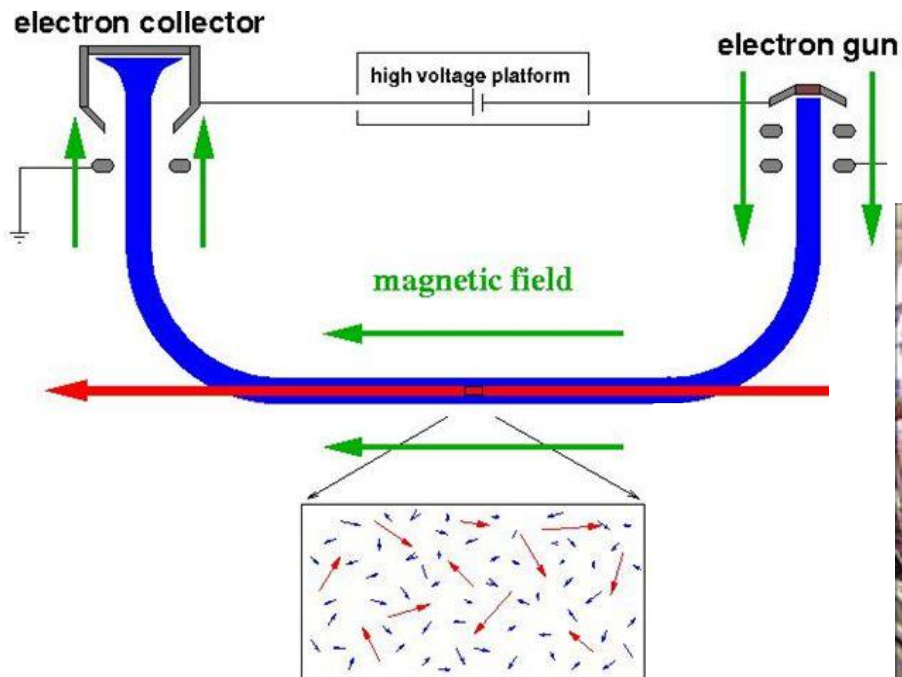
Tested first time 1977 together with electron cooling at CERN in the ICE (initial cooling experiment)



# Electron cooling

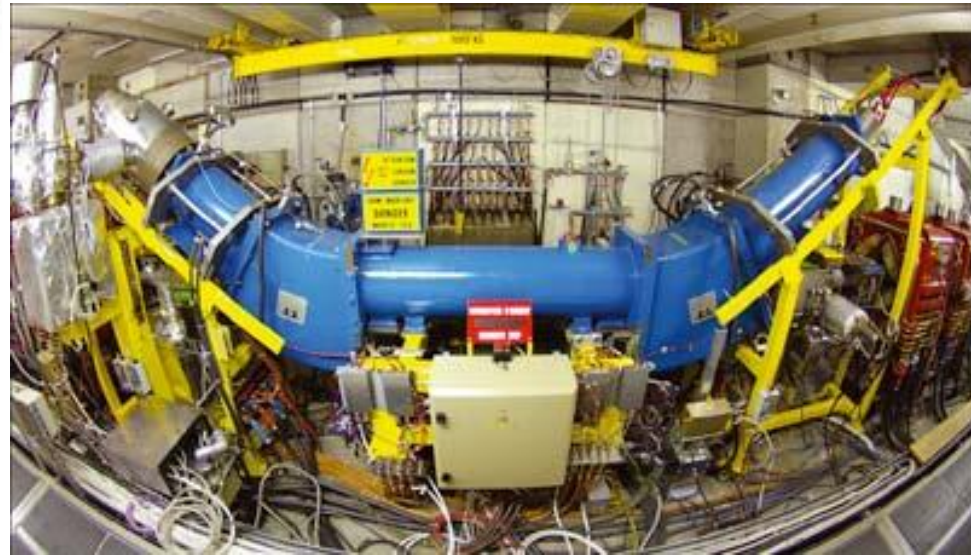


- **Superposition** of cold intense  $e^-$  beam with  $\bar{p}$  at **same velocity**
- Momentum transfer by Coulomb collisions
- Cooling results from energy loss in co-moving gas of free electrons



*In the AD:*

$\bar{p}$  beam merged with  $\varnothing$  20 mm  $e^-$  beam of  $\sim 3$  A collinearly over 2 m

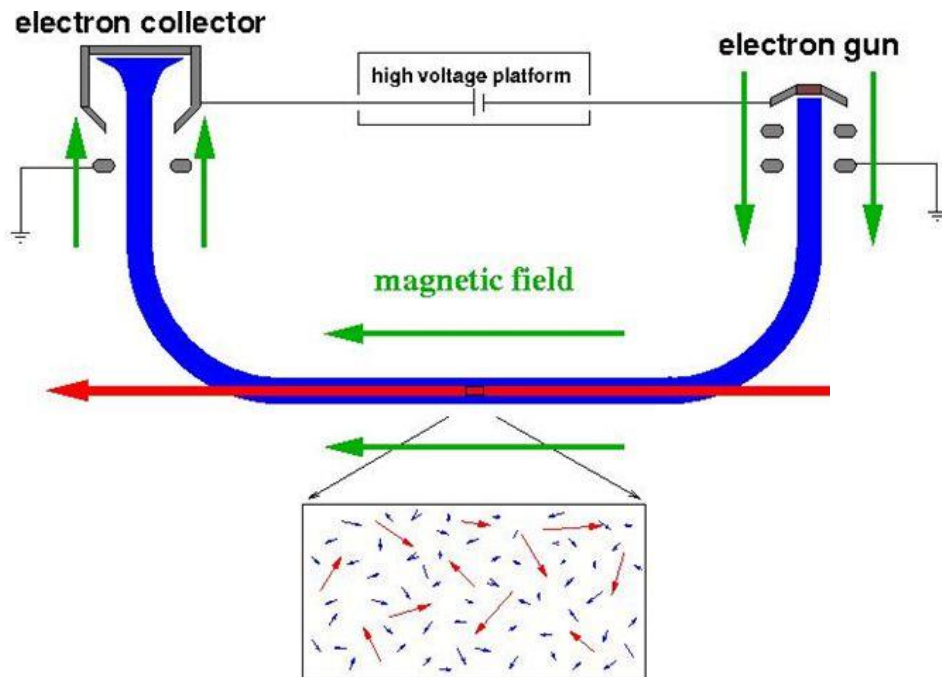




# Electron cooling



- **Superposition** of cold intense  $e^-$  beam with  $\bar{p}$  at **same velocity**
- Momentum transfer by Coulomb collisions
- Cooling results from energy loss in co-moving gas of free electrons



$$v_e = \beta_e c = \beta_p c = v_p$$

$$E_e = \frac{m_e}{m_p} E_p$$

$m_e$ ...electron mass

$m_p$ ... $\bar{p}$  mass

$E_e$ ...electron kinetic energy

$E_p$ ... $\bar{p}$  kinetic energy

e.g. 220 keV **electrons** cool 400 MeV  $\bar{p}$

In the beam frame:

**Cold electrons** interact with

**hot  $\bar{p}$**

# AD machine cycle

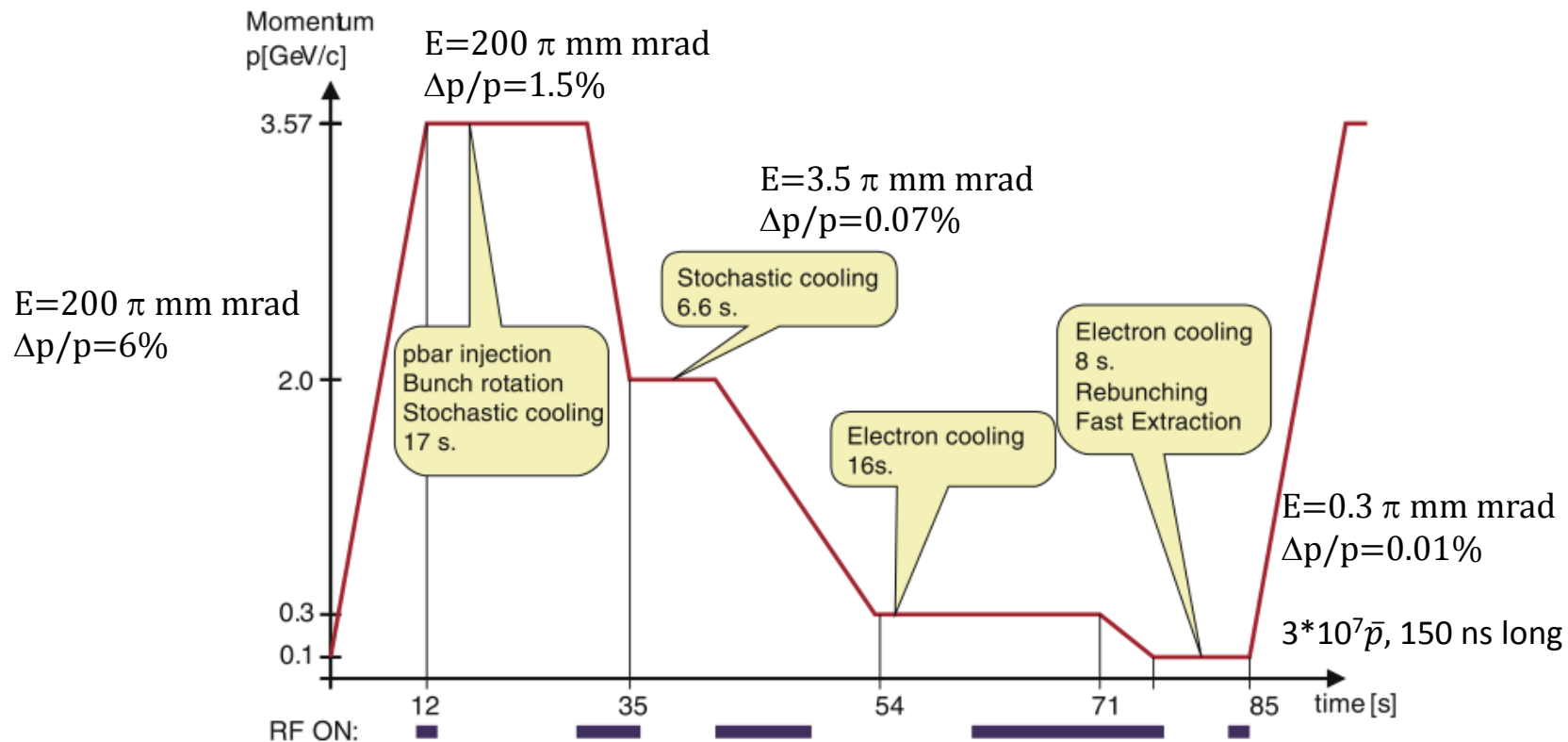


**Bunch rotation:** reduce  $\Delta p/p$

**Stochastic cooling:** reduce emittance, reduce  $\Delta p/p$

**Deceleration in radiofrequency cavities:** reduce p

**Electron cooling:** reduce emittance, increase D

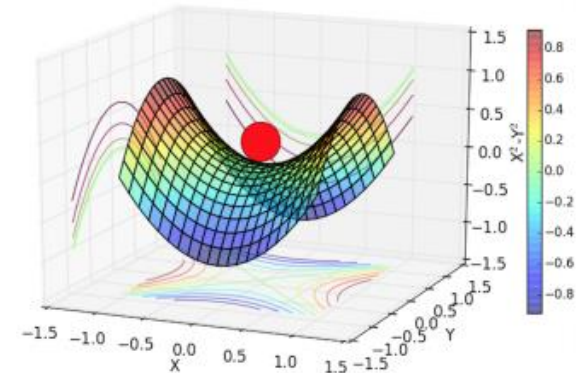
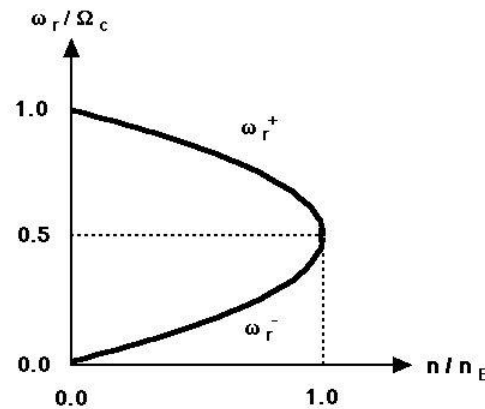
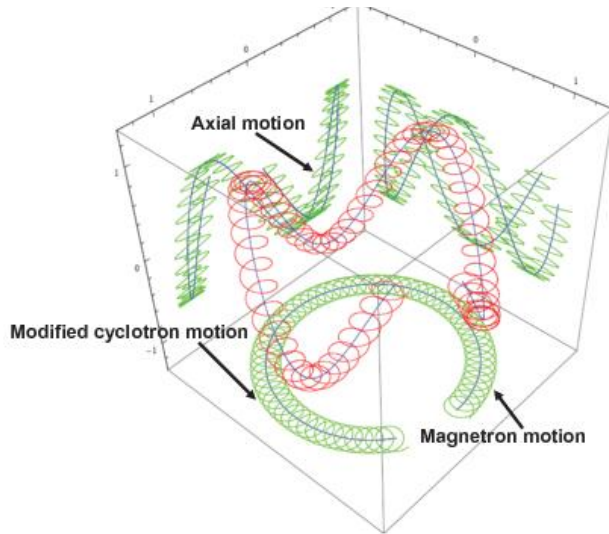
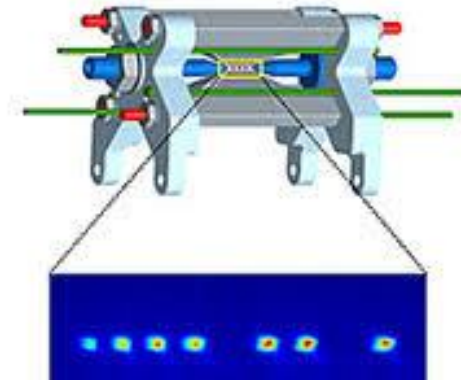
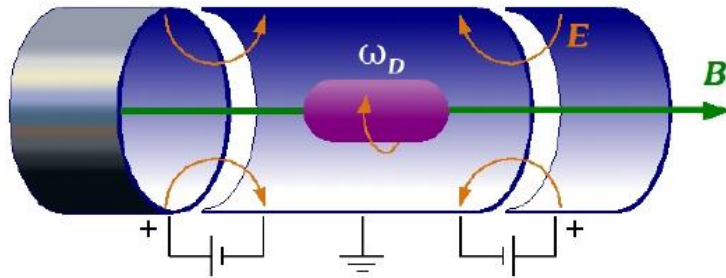


# Cooling in Particle Traps



Penning traps, uses E and B fields

Paul trap, uses rf-E fields



$$\omega_c = e B / m \quad \omega_c^2 = \omega_+^2 + \omega_-^2 + \omega_z^2 \quad \omega_c = \omega_+ + \omega_-$$



## $\bar{p}$ cooling mechanism in Penning traps:

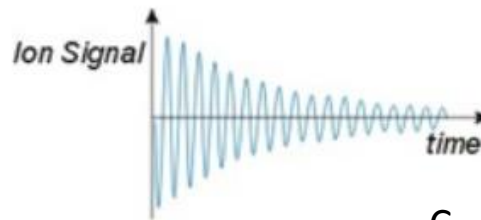
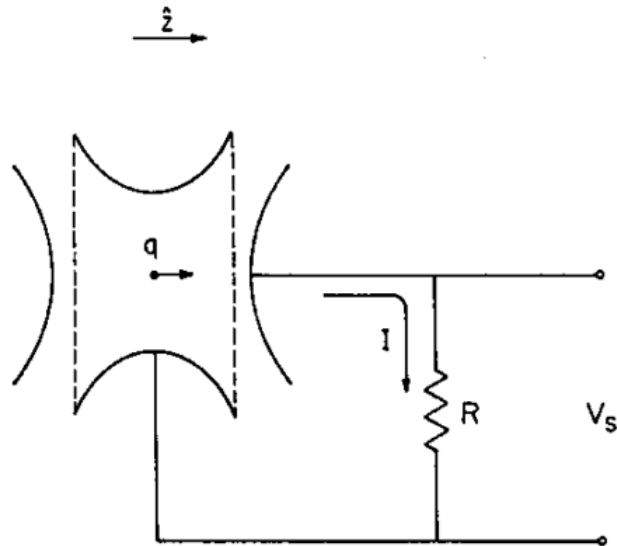
- Resistive cooling
- Stochastic cooling
- Sympathetic radiation electron cooling
- Evaporative / adiabatic cooling
- Sympathetic laser cooling using anion species:

**Os<sup>-</sup>** spectroscopy: *U. Warring et al., Phys. Rev. Lett.* **102** 043001 (2009)

**La<sup>-</sup>** spectroscopy: *E. Jordan et al., Phys. Rev. Lett.* **115** 113001 (2015)

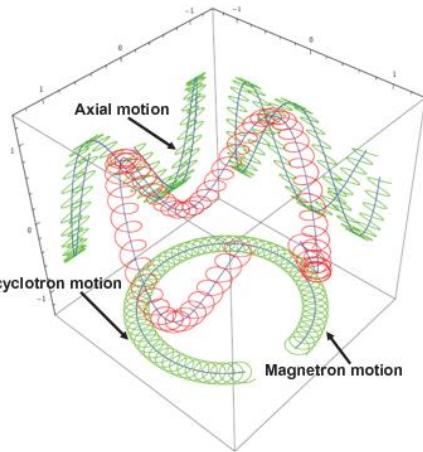
**C<sub>2</sub><sup>-</sup>** proposal: *P. Yzombard et al., Phys. Rev. Lett.* **114** 213001 (2015)

# Resistive cooling in Penning traps



$$\tau_R = \left( \frac{2z}{eC} \right)^2 \frac{m}{R}$$

C... constant of about unity on trap shape  
z...dimension of trap (~1 cm)  
R...resistance (100 kOhm gives 270 s for p)



- proven method for single particles in Penning traps
- difficult for many particles because of broad plasma modes  
i.e. weak coupling to high-Q LRC circuit

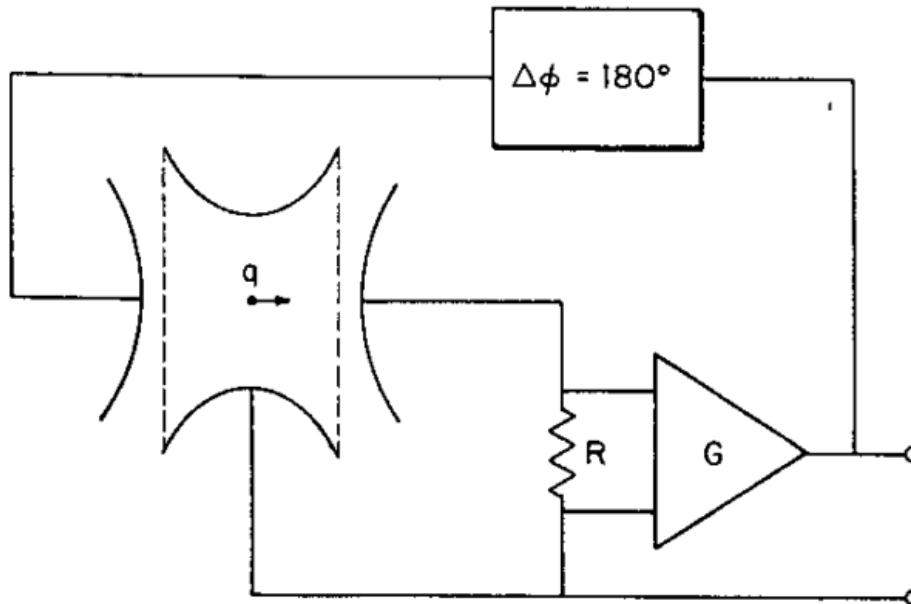


# Stochastic cooling in Penning traps



Stochastic cooling, negative feedback cooling

Signal from particle is amplified and fed back onto trap electrodes with 180° phase shift



$$\tau_{St} = \frac{1}{1 + G} \tau_R$$

G... amplifier gain

- proven method for single particles in Penning traps
- difficult for many particles because of broad plasma modes  
i.e. weak coupling to high-Q LRC circuit

# Electron cooling in Penning traps



Electron cooling using coupled rate equations:

$$\frac{d}{dt} T_p = -\frac{1}{\tau_{elcool}} (T_p - T_e)$$

$$\frac{d}{dt} T_e = \frac{N_p}{N_e} \frac{1}{\tau_{elcool}} (T_p - T_e) - \frac{1}{\tau_{rad}} (T_e - T_{trap})$$

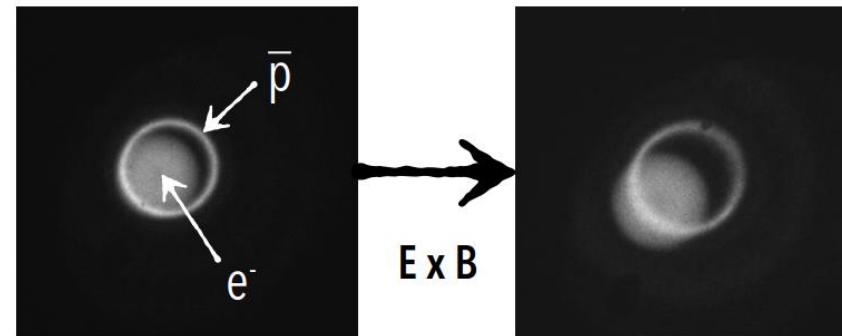
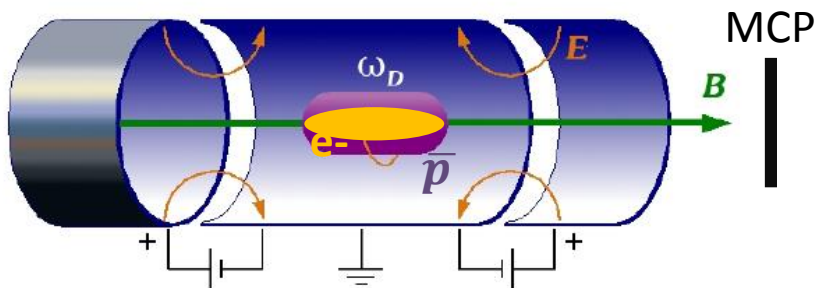
$T_p$ ...  $\bar{p}$  temperature

$T_e$ ...electron temperature

$T_{trap}$ ...temperature of trap electrodes

$N_p$ ...nr. of  $\bar{p}$

$N_e$ ...nr. of electrons



Centrifugal separation, causes  $e^-$  to be on axis

# Electron cooling in Penning traps



Electron cooling using coupled rate equations:

$$\frac{d}{dt} T_p = -\frac{1}{\tau_{elcool}} (T_p - T_e)$$

$T_p$ ...  $\bar{p}$  temperature

$T_e$ ...electron temperature

$$\frac{d}{dt} T_e = \frac{N_p}{N_e} \frac{1}{\tau_{elcool}} (T_p - T_e) - \frac{1}{\tau_{rad}} (T_e - T_{trap})$$

$T_{trap}$ ...temperature of trap electrodes

$N_p$ ...nr. of  $\bar{p}$

$N_e$ ...nr. of electrons

$$\tau_{rad} = \frac{3\pi m_e^3 c^3 \epsilon_0}{B^2 e^4}$$

Synchrotron radiation term

$B$ ... magnetic field of Penning trap

$$\tau_{elcool} = \frac{3m_p m_e c^3 (4\pi\epsilon_0)^4}{8\sqrt{2\pi} n_e e^4} \frac{1}{\ln\left(4000 n_e^{-1/2} T_e^{1/2} \left(T_e + \frac{T_p}{1836} + \frac{\sqrt{T_e T_p}}{21}\right)\right)} \left(\frac{kT_p}{m_p c^2} + \frac{kT_e}{m_e c^2}\right)^{3/2}$$

$n_e$ ...density of electrons

$m_e$ ...electron mass

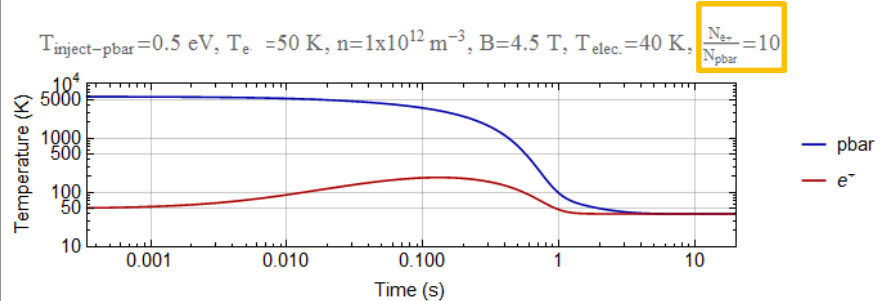
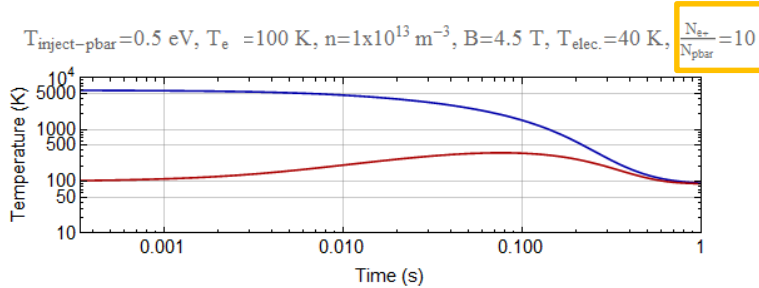
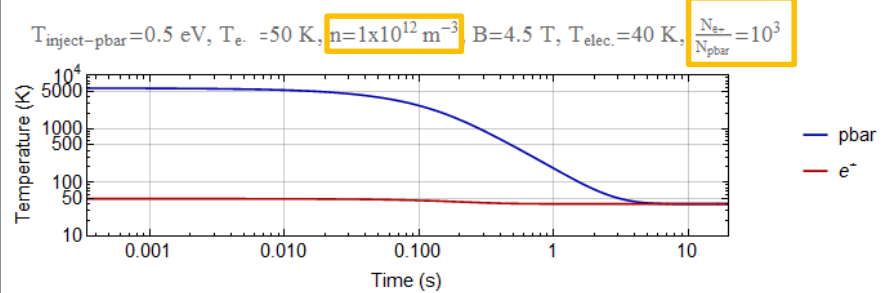
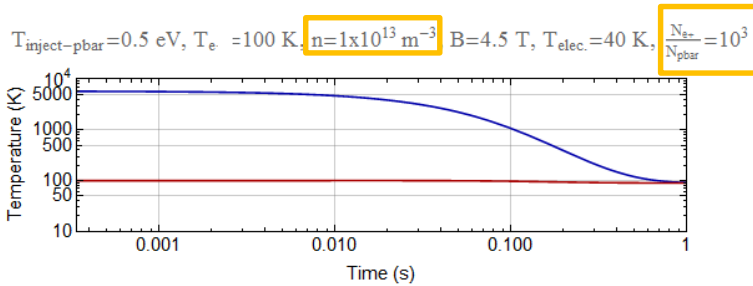
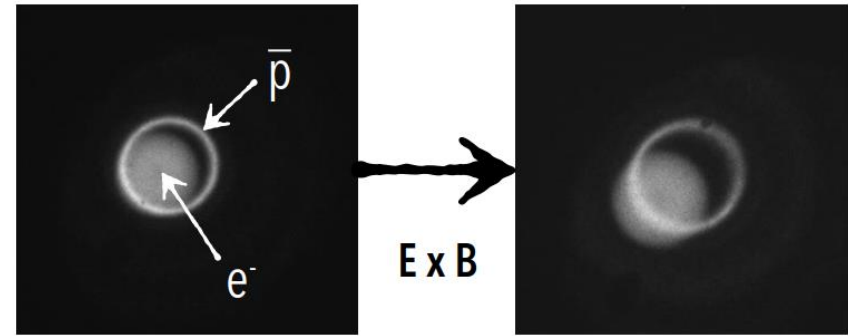
$m_p$ ... $\bar{p}$  mass

# Electron cooling in Penning traps



$$\frac{d}{dt} T_p = -\frac{1}{\tau_{elcool}} (T_p - T_e)$$

$$\frac{d}{dt} T_e = \frac{N_p}{N_e} \frac{1}{\tau_{elcool}} (T_p - T_e) - \frac{1}{\tau_{rad}} (T_e - T_{trap})$$



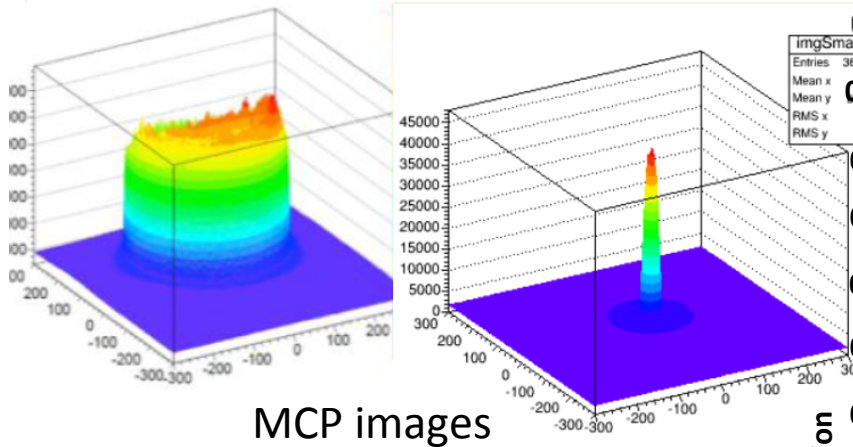
→ Great method for many particles, limited by B field and environmental temperature

# Electron cooling in Penning traps

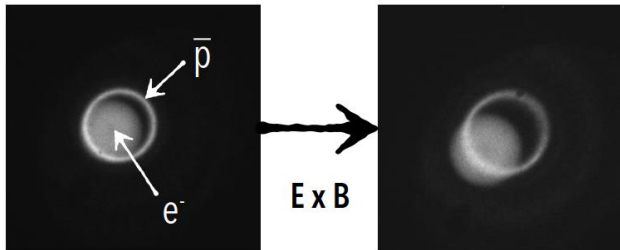


Example of a measurement:

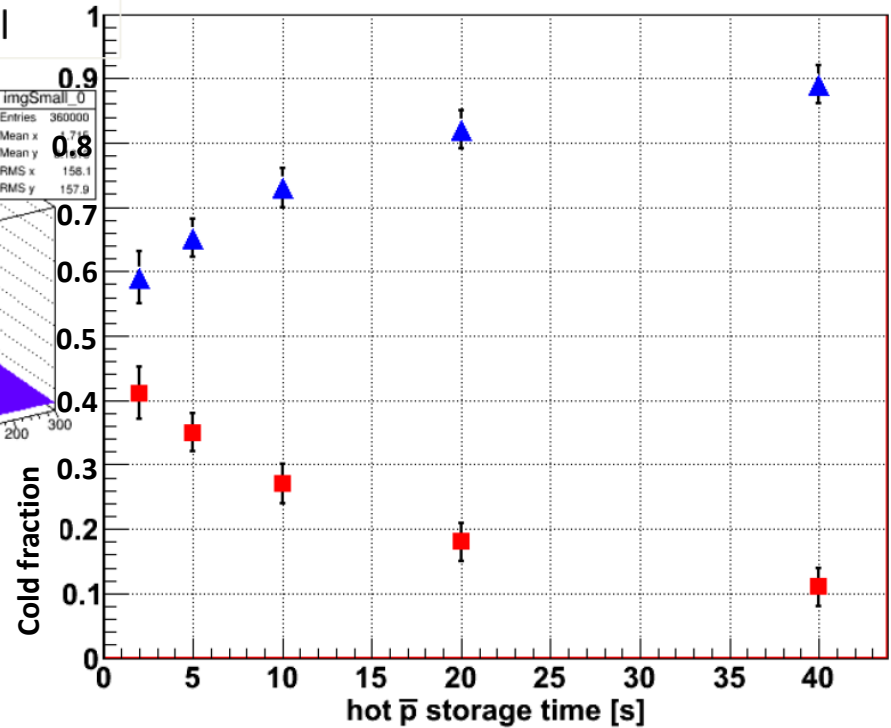
Radial compression of the Trapped plasma with RF field: Rotating Wall



MCP images



Cold and hot antiproton fractions vs time of cooling



$$\tau_{rad} \propto \frac{m^3}{B^2}$$

$$\begin{array}{l} e^-, e^+ \quad \tau_{rad} \cong 0.1 \text{ sec} @ 5T \\ \bar{p} \quad \tau_{rad} \cong 10^9 \text{ sec} @ 5T \end{array}$$

Cyclotron radiation + Coulomb collisions = thermal equilibrium for e- and pbar

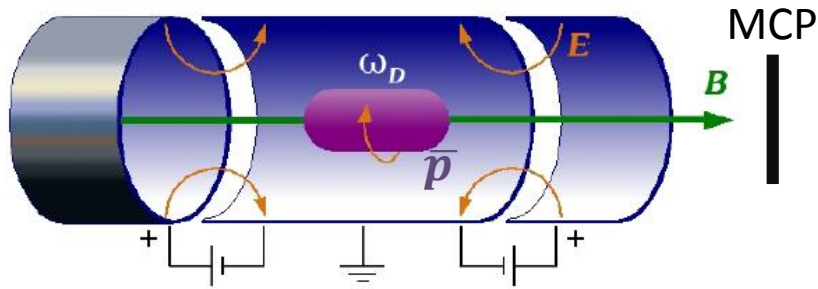
Final energy estimation: about 100 K



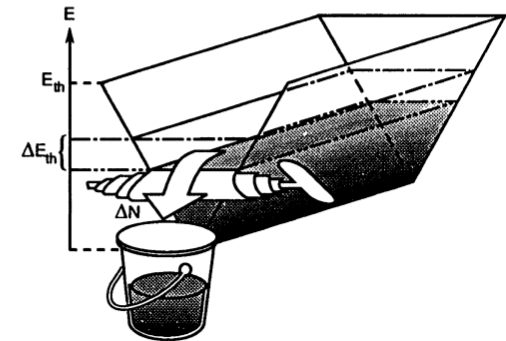
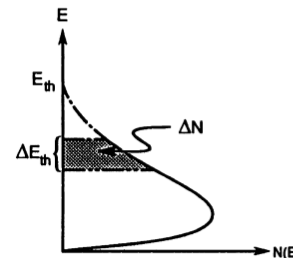
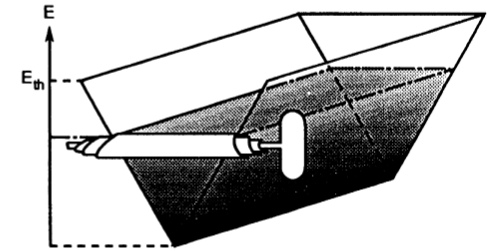
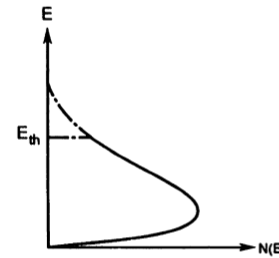
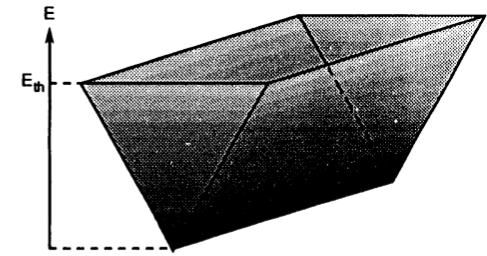
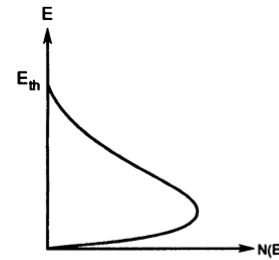
# Evaporative cooling in Penning traps



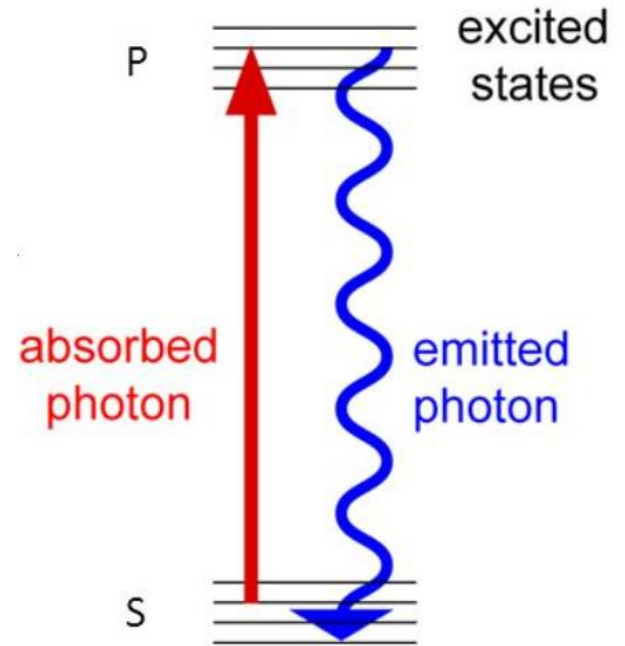
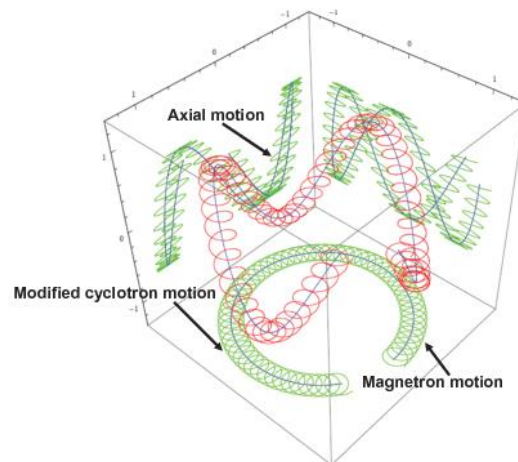
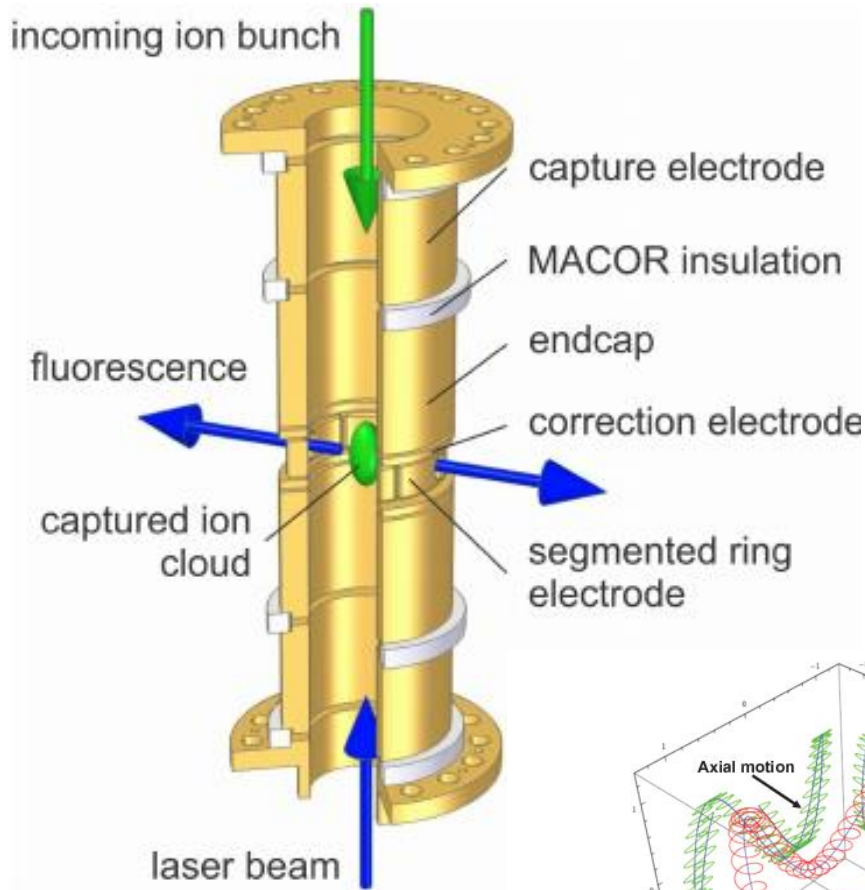
Lower axial confinement voltage to open trap



- Often used in AD experiments
- Limited by  $\bar{p}$  numbers and axial confinement
- Reach typically  $\sim 100$  K



# Laser cooling in Penning traps



$$\hbar\omega_{\text{abs}} < \hbar\omega_{\text{em}}$$

**=> COOLING !**

(Need a 2 level system)

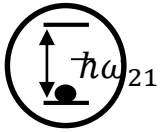


**Three possible laser cooling methods for anions  
(to sympathetically cool  $\bar{p}$ ) :**

- 1)...Doppler laser cooling**
- 2)...photodetachment cooling**
- 3)...dipole force Sisyphus cooling**



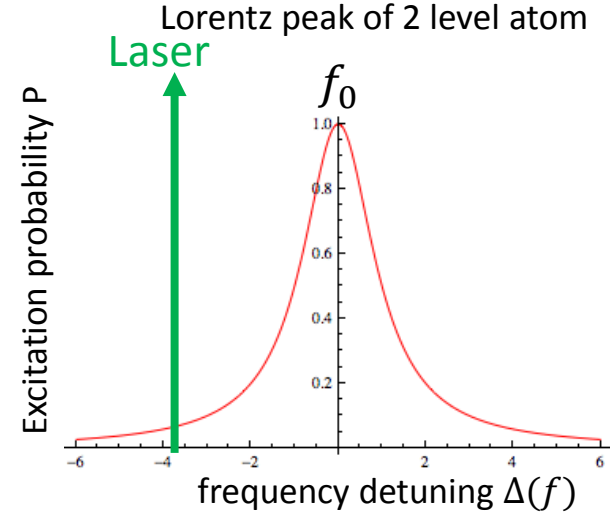
## Doppler laser cooling



Atom stands still, **no excitation**

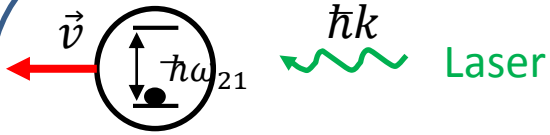
Atom sees laser without Doppler shift

$$f_D = \frac{2f_0}{c} \vec{v} = 0$$





## Doppler laser cooling

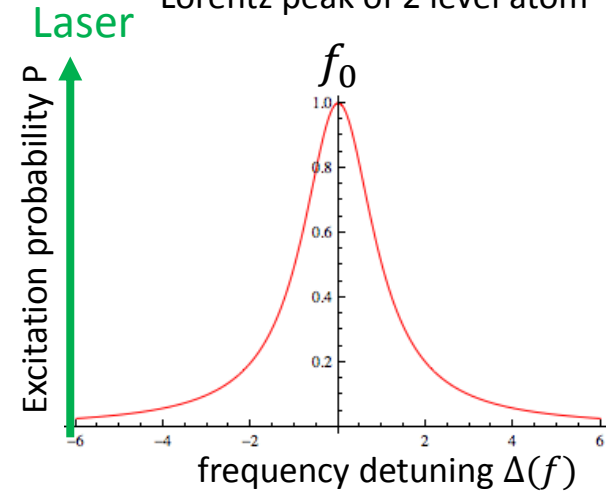


Atom moves away from the laser,  
**no excitation**

Atom sees laser red Doppler shifted

$$f_D = \frac{2f_0}{c} \vec{v}$$

Lorentz peak of 2 level atom

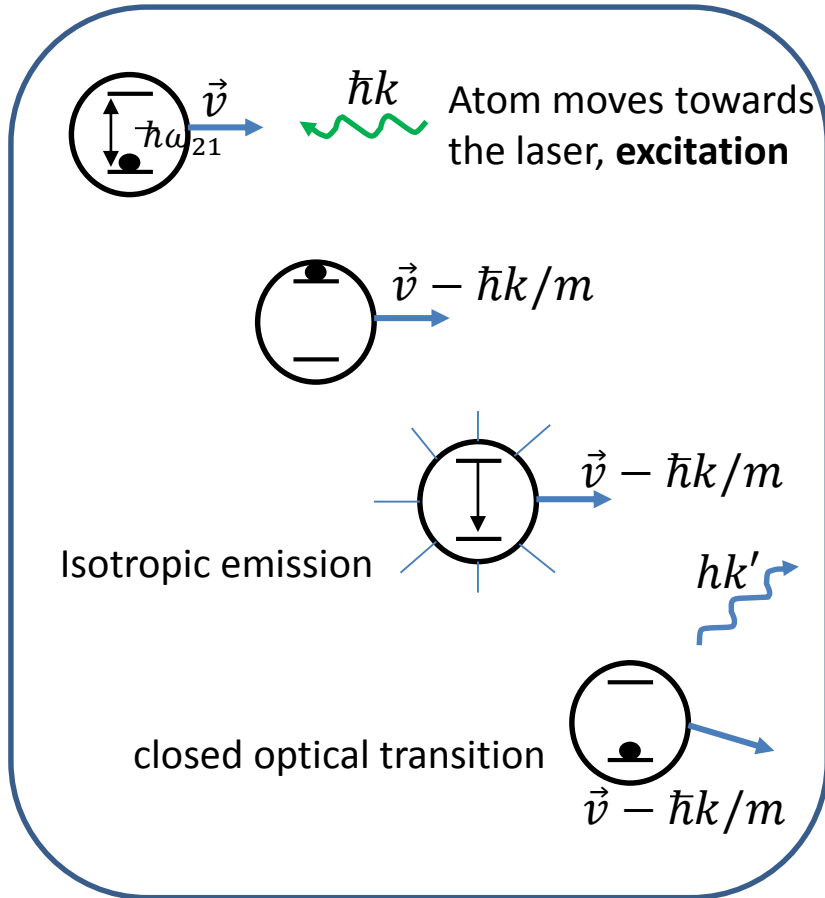




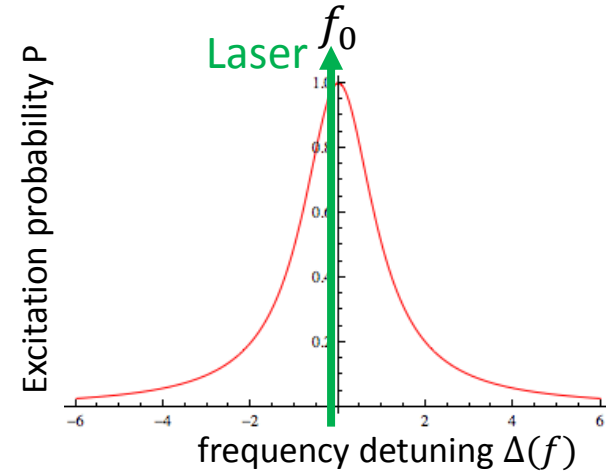
# Laser cooling in Penning traps



## Doppler laser cooling



Lorentz peak of 2 level atom



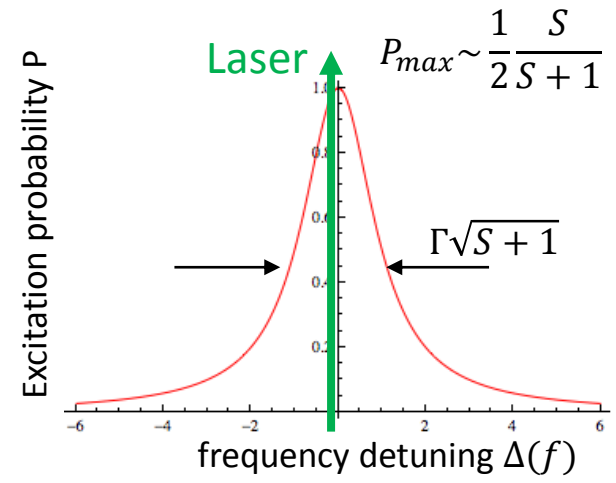
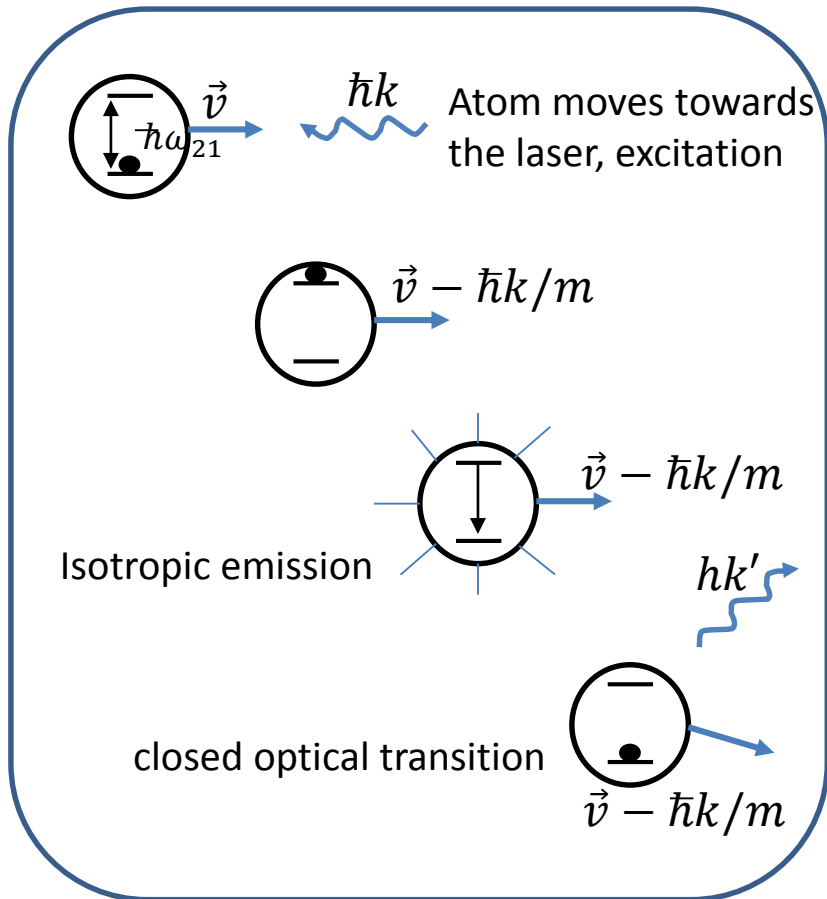
Atom sees laser blue Doppler shifted

$$f_D = \frac{2f_0}{c} \vec{v}$$

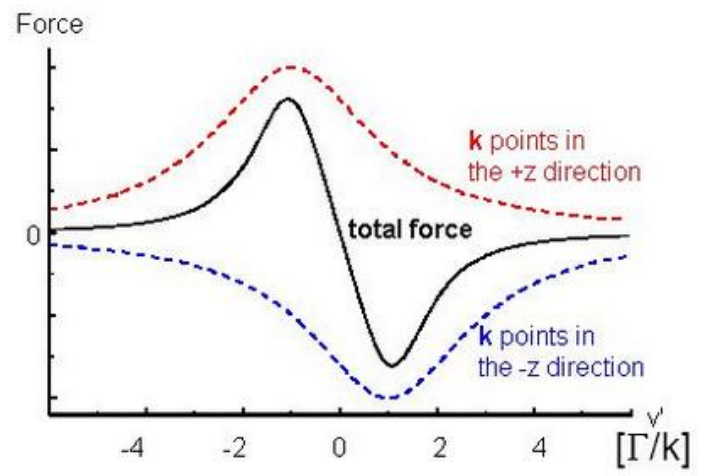
# Laser cooling in Penning traps



## Doppler laser cooling

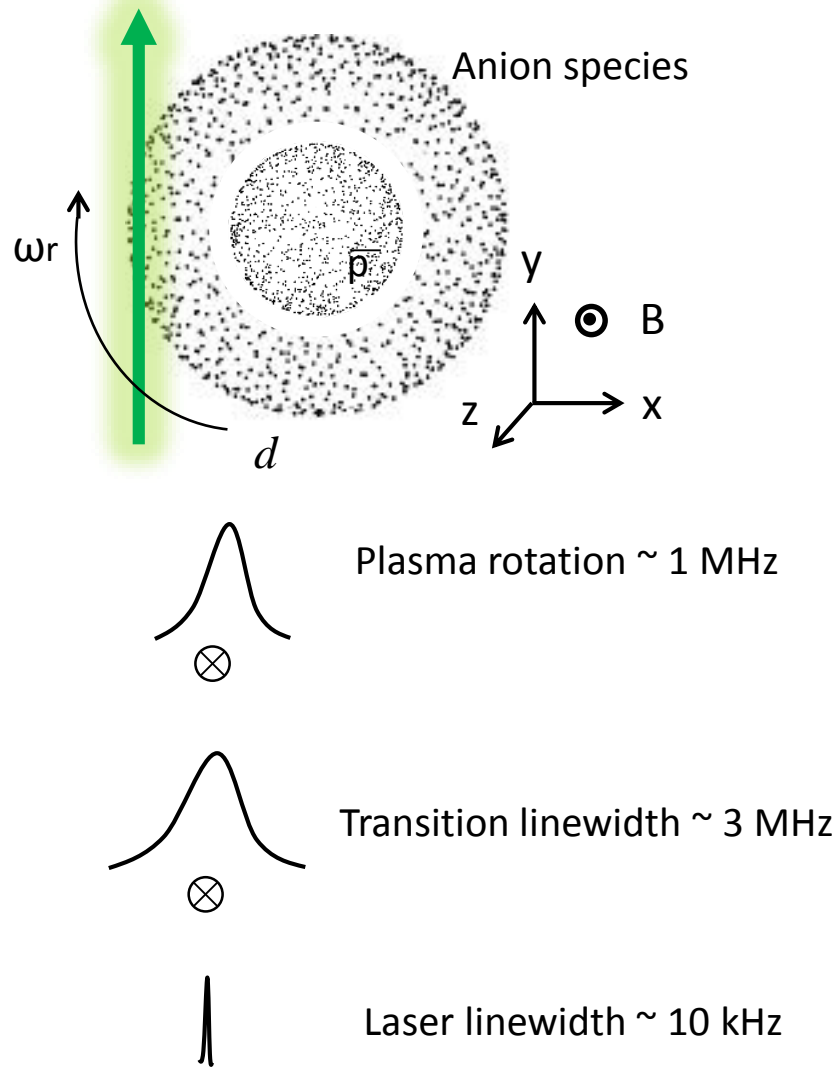


Cooling force 
$$F = \pm \frac{\hbar \vec{k}}{2} S \Gamma \frac{(\Gamma/2)^2}{(\omega - \omega_{21} - \vec{v} \vec{k}) + (\Gamma/2)^2 (1 + S)}$$



Minimal Doppler temperature 
$$T_D = \frac{\hbar \Gamma}{2 k_B}$$

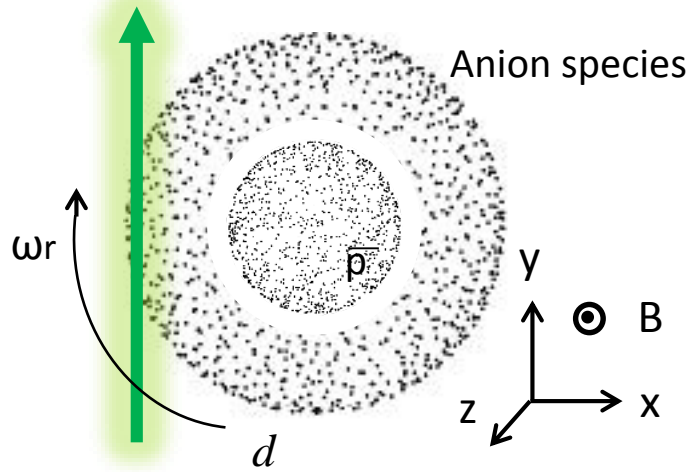
# Doppler cooling in Penning traps



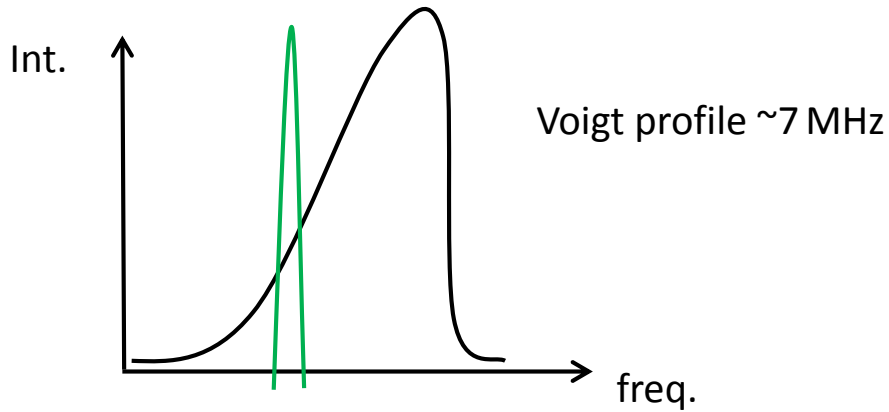
## Laser interacts both with 'magnetron' and 'cyclotron'

- If  $\omega_L < \omega_0$  and laser put to receding position  
→ Plasma size reduced
- If  $\omega_L < \omega_0$  and laser put to the approaching side  
→ Plasma size increased
- Cyclotron is laser cooled and plasma forms nearly circular magnetron shape.
- Transition blue Doppler shifted prop. to the distance from trap centre

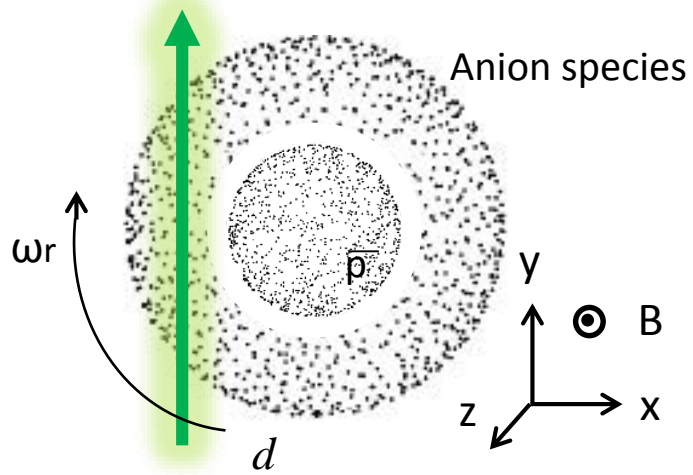
# Doppler cooling in Penning traps



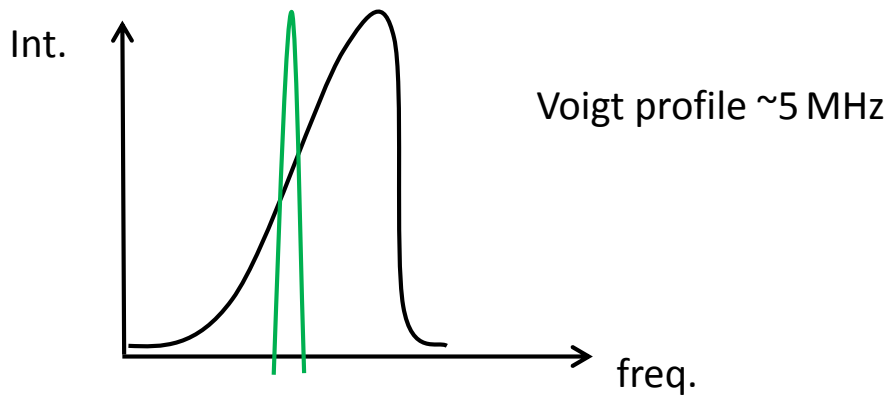
Laser on cooling transition:



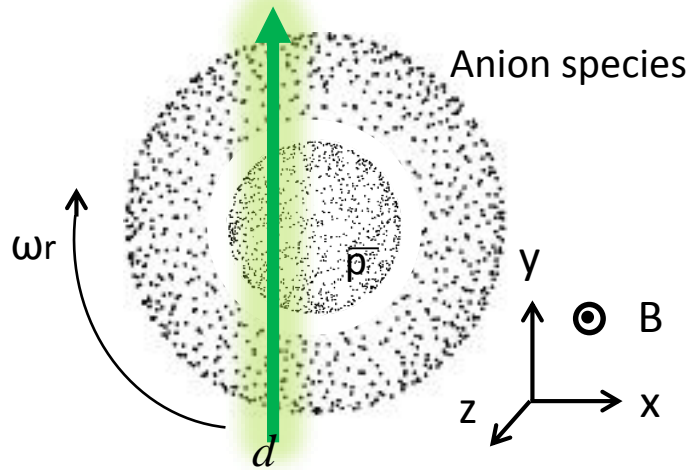
# Doppler cooling in Penning traps



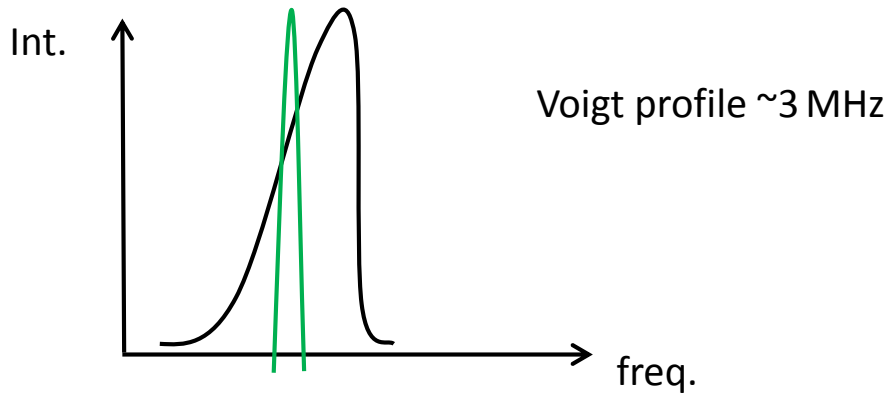
Laser on cooling transition:



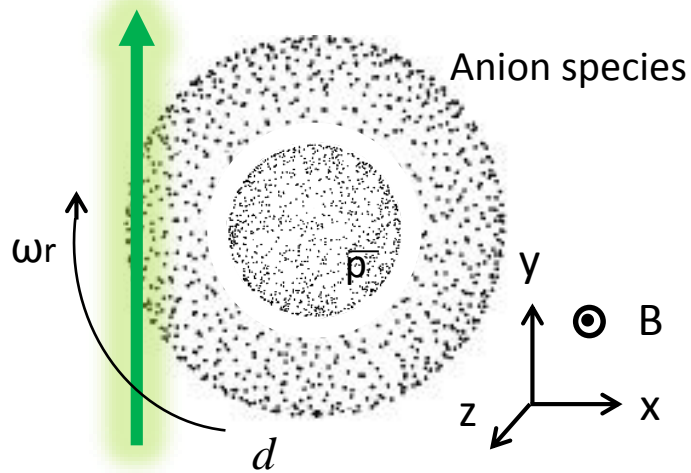
# Doppler cooling in Penning traps



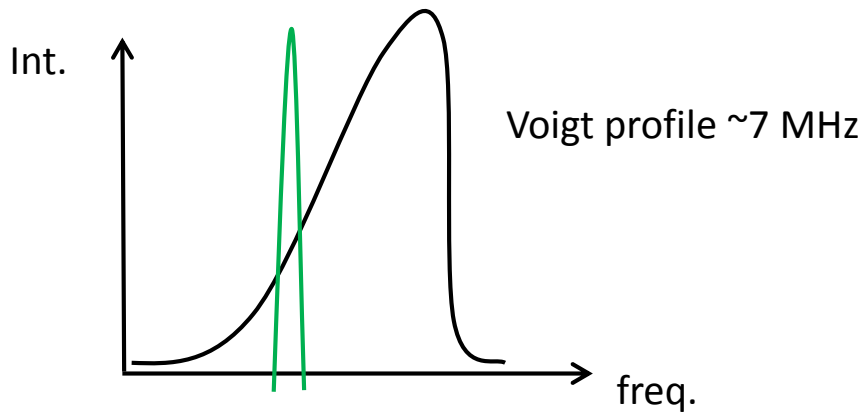
Laser on cooling transition:



# Doppler cooling in Penning traps

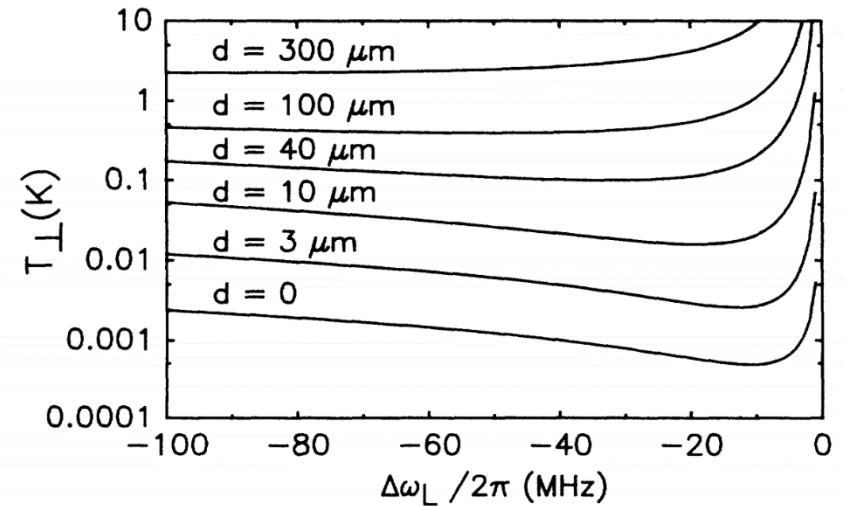


Laser on cooling transition:



Doppler cooling Penning trap: its limit

Be<sup>+</sup> plasma:  $\Delta = -20$  MHz at 313 nm  
→  $\mu$ K reached



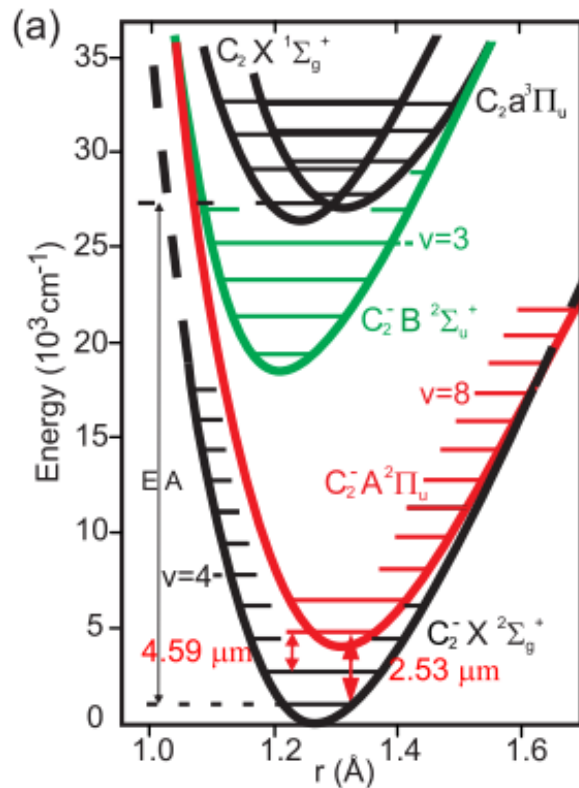
W. Itano et al., *Phys. Rev. A.* **38**, 5698 (1985)



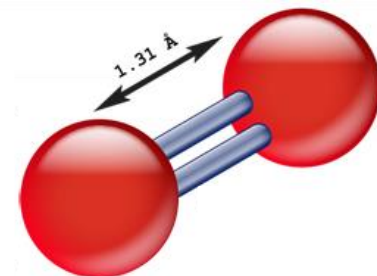
# Anion example: Molecule $C_2^-$



Ro-Vib level structure and molecular potential energy of  $C_2^-$



$C_2$  Homonuclear molecule



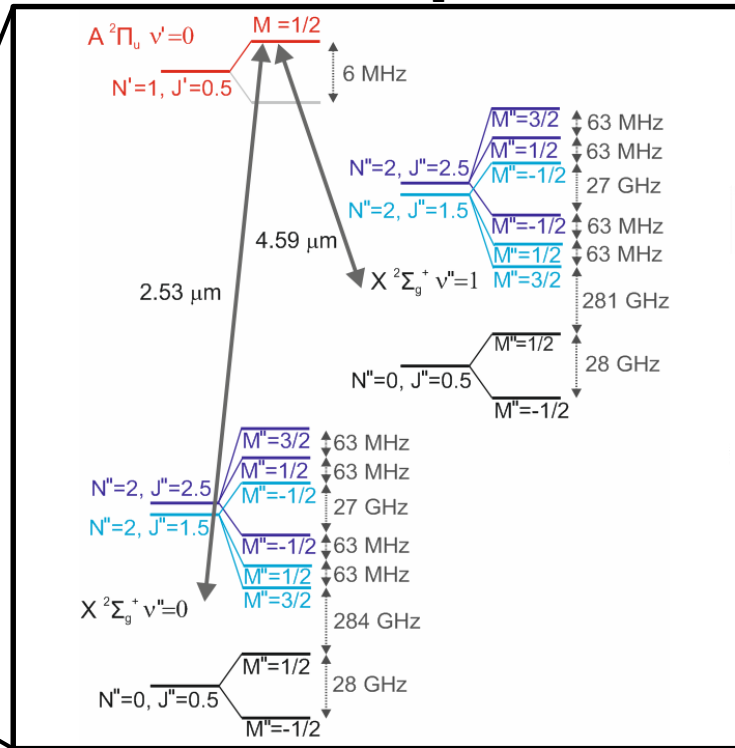
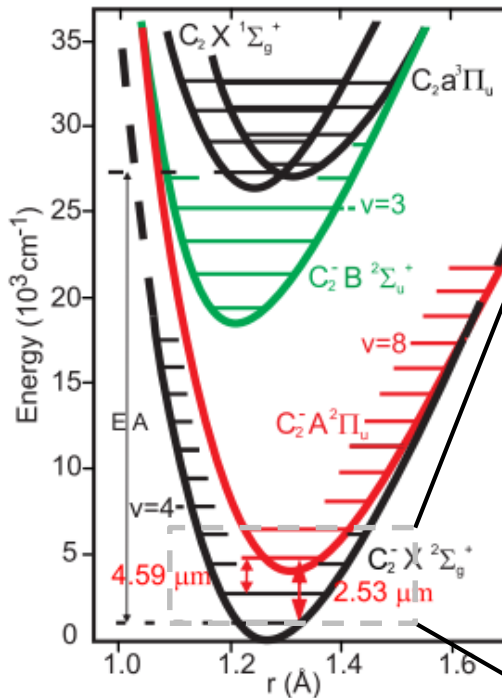
## $C_2^-$ benefits:

- Known level scheme to  $\nu/d\nu \sim 100$  MHz M. Tulej et al., *J. Raman Spectrosc.* **41**, 853 (2010)
- 2.53  $\mu\text{m}$  and 4.59  $\mu\text{m}$  dipole transition accessible with DFB diode lasers
- Production in supersonic gas expansion

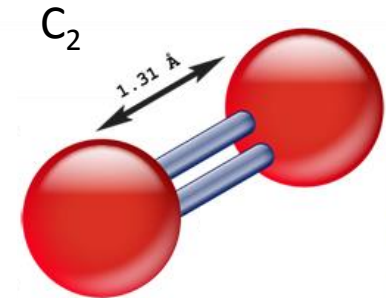
# Anion example: Molecule $C_2^-$



Ro-Vib level structure and molecular potential energy of  $C_2^-$  in B=1 T



Homonuclear molecule



## $C_2^-$ benefits:

- Known level scheme to  $v/dv \sim 100$  MHz M. Tulej et al., *J. Raman Spectrosc.* **41**, 853 (2010)
- 2.53  $\mu\text{m}$  and 4.59  $\mu\text{m}$  dipole transition accessible with DFB diode lasers
- Production in supersonic gas expansion

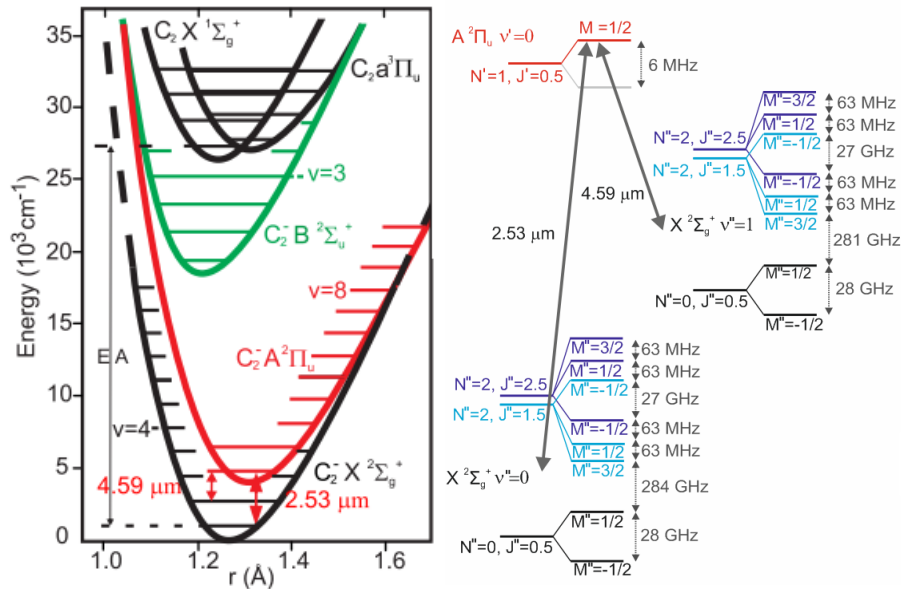
## $C_2^-$ challenges:

- 16 transitions for closed laser cycle with Zeeman splitting
- 20 kHz linewidth

# Doppler cooling simulation $C_2^- / \bar{p}$



2 x Doppler 2.53  $\mu\text{m}$  laser  
6 x Repumper at 2.53  $\mu\text{m}$ , 4.59  $\mu\text{m}$



→ Crystallization effects at low T  
(trap and plasma geometry, B,  $\Gamma_{coupling}$ )

H. Totsuji et al., *AIP Conf. Proc.*, **498**, 77 (1999)

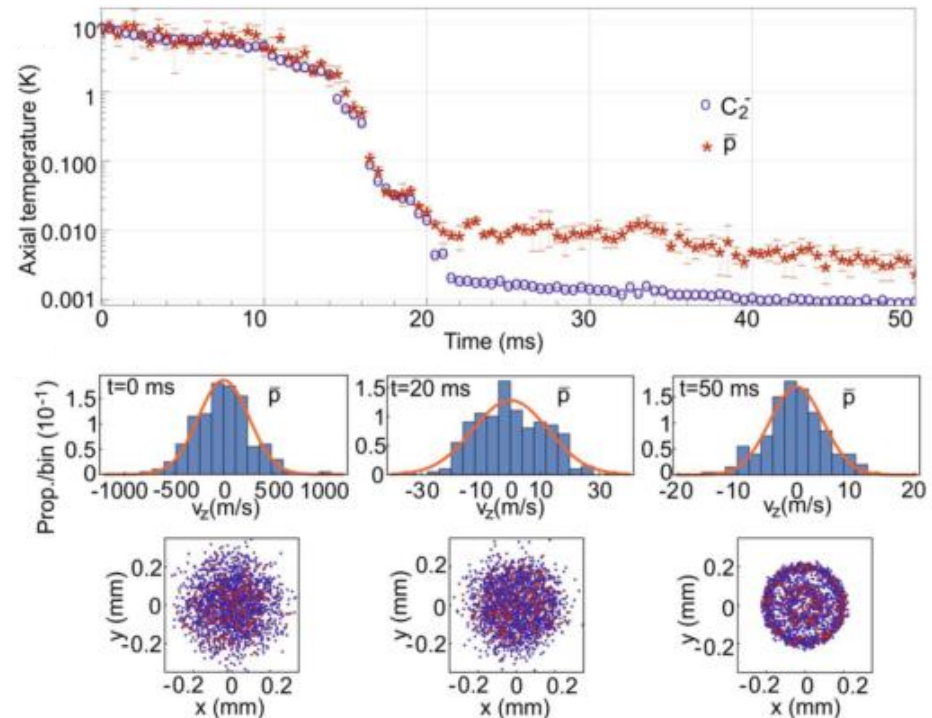
T. Mitchell et al., *Science*, **282**, 1290 (1998)

→ ~10 mK temperature within minutes

## simulation

$n_{C_2^-} / \bar{p} = 10^7 \text{ cm}^{-3}$ , B=1 T, initialized at 11 K

GPU calculation of 1000  $C_2^- / 100 \bar{p}$  particle-particle interactions  
Coulomb factor increased by  $10^4$



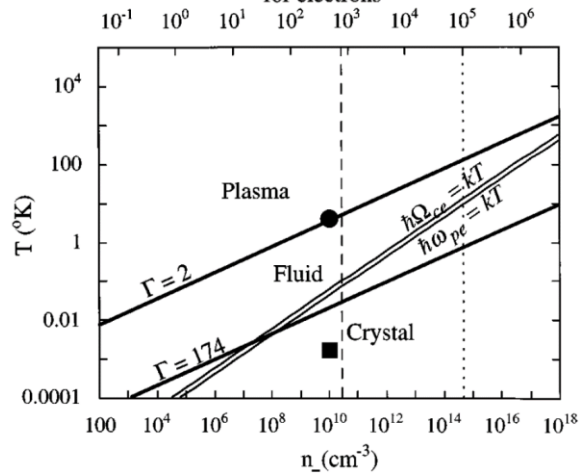
Detuning adjusted in 4 steps from  $\Delta\nu = -39 \text{ MHz}$  to  $\Delta\nu = -3.4 \text{ MHz}$ .  
Laser linewidth  $\delta\nu = 1 \text{ MHz}$

# Doppler cooling simulation $C_2^- / \bar{p}$



## Phase diagram B - T - n:

minimum B(Gauss)  
for electrons



Malmberg and O'Neil (1977)

FIGURE 1. Relaxation of energies. From left to right, parallel component of antiprotons, parallel, and perpendicular components of electrons.

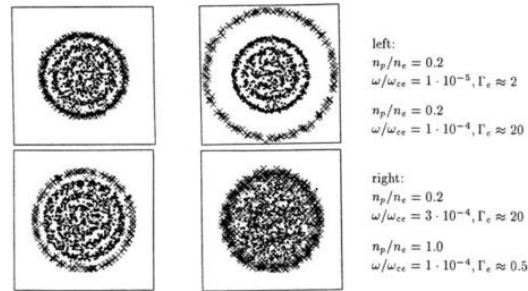


FIGURE 2. Examples of equilibrium distribution of electrons (dots) and antiprotons (crosses) observed along the magnetic field.

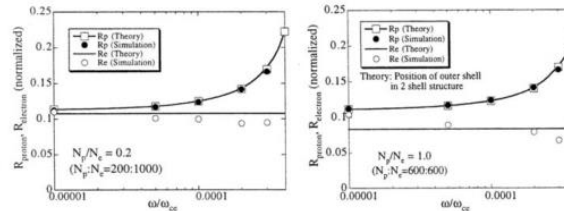
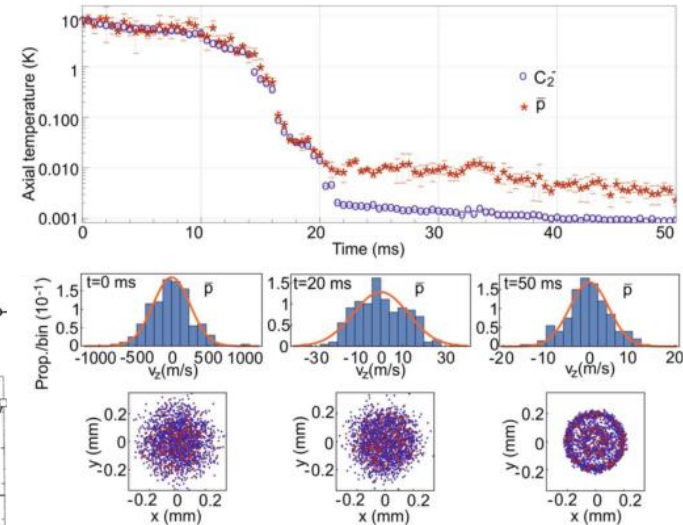


FIGURE 3. Position of antiproton shell and the radius of electron distribution. Simulation (symbols) and theory (lines).

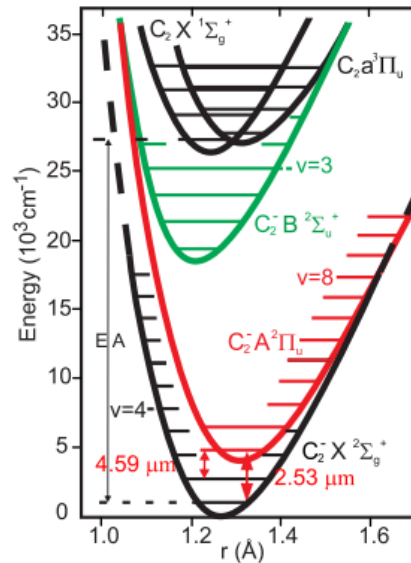
H. Totsuji et al., *AIP Conf. Proc.* **498**, 77 (1999)



# Doppler cooling simulation $C_2^- / \bar{p}$



2 x Doppler 2.5  $\mu\text{m}$  laser  
6 x Repumper at 2.5  $\mu\text{m}$ , 4.5  $\mu\text{m}$



→ Crystallization effects at low T  
(trap and plasma geometry, B,  $\Gamma_{coupling}$ )

H. Totsuji et al., *AIP Conf. Proc.*, **498**, 77 (1999)

T. Mitchell et al., *Science*, **282**, 1290 (1998)

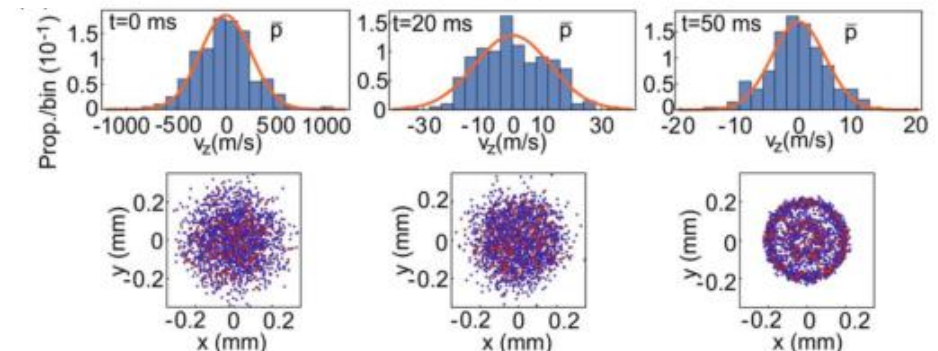
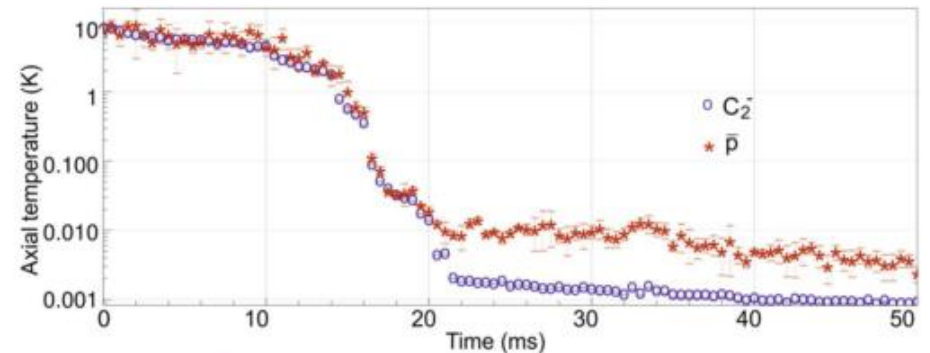
→ ~10 mK temperature within minutes

→ Possible, but challenging in Penning trap  
(scatter  $\sim 10^6$  photons with 8 lasers plus sidebands to <MHz stabilized)

## simulation

$n_{C_2^-} / \bar{p} = 10^7 \text{ cm}^{-3}$ , B=1 T, initialized at 11 K

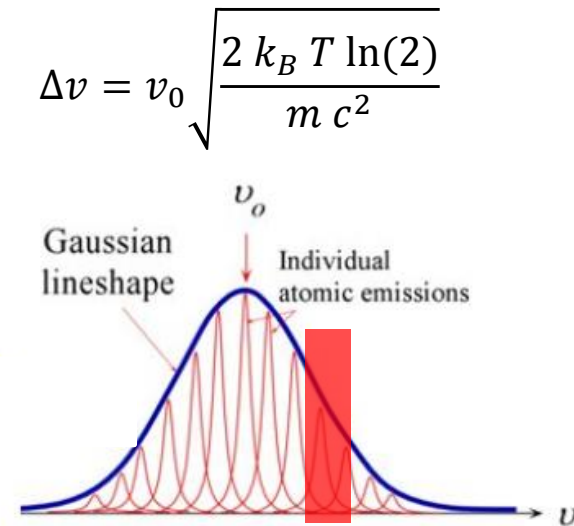
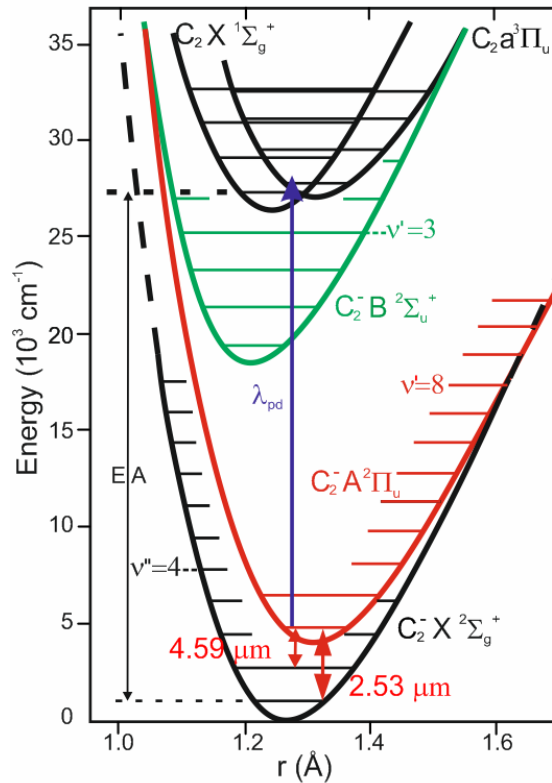
GPU calculation of 1000  $C_2^- / 100 \bar{p}$  particle-particle interactions  
Coulomb factor increased by  $10^4$



Detuning adjusted in 4 steps from  $\Delta\nu = -39 \text{ MHz}$  to  $\Delta\nu = -3.4 \text{ MHz}$ .  
Laser linewidth  $\delta\nu = 1 \text{ MHz}$



Ro-Vib level structure and molecular potential energy of  $C_2^-$



$$\Delta v = v_0 \sqrt{\frac{2 k_B T \ln(2)}{m c^2}}$$

2.53  $\mu\text{m}$  laser addresses  $\nu + \Delta\nu$

## Photodetachment cooling

- 2.53  $\mu\text{m}$  laser addresses high velocity fraction of molecules
- UV laser 399 nm photodetaches this fraction

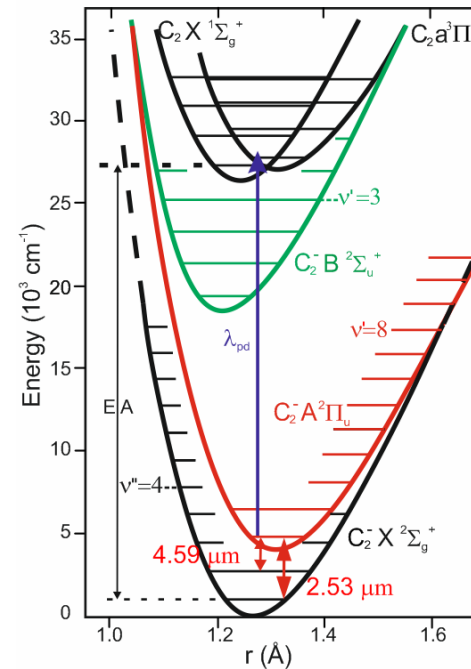
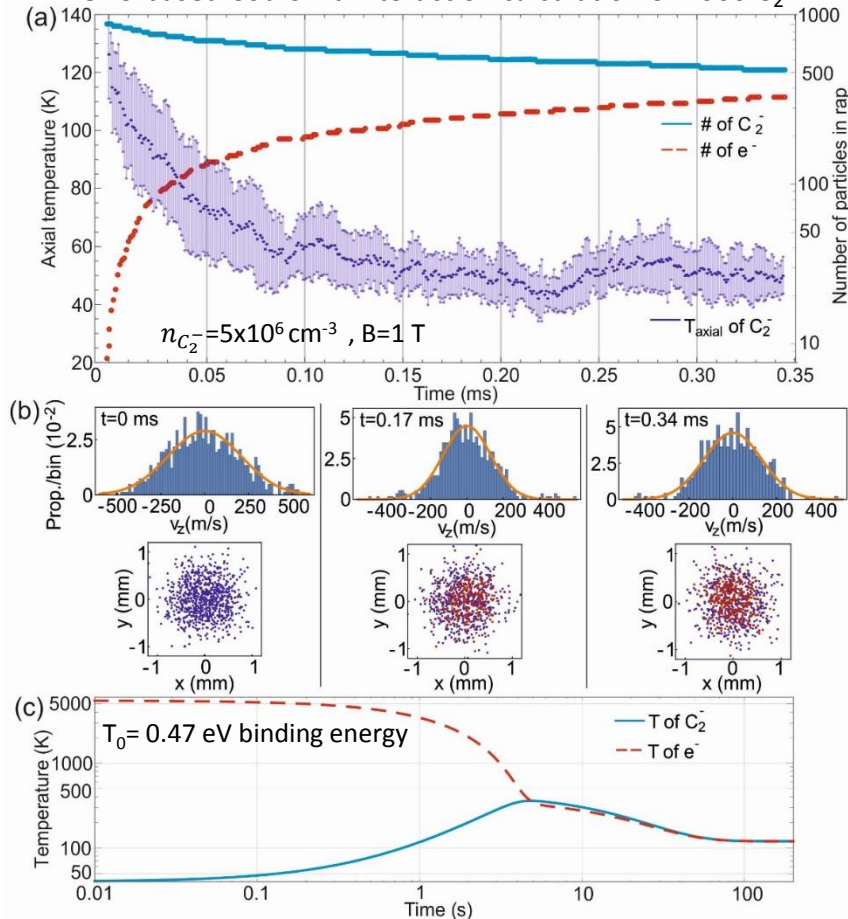




## simulation

1 x 2.5  $\mu\text{m}$  laser, 1 x photodetachment laser 399 nm (100 mW enhanced in  $F \sim 1000$  cavity)

GPU based Coulomb interaction calculation of 1000  $C_2^-$ :



Photodetachment rate:

$$\Gamma = \frac{I}{h\nu} \sigma \sim 65 \text{ kHz}$$

Photodetachment cross section:

$$\sigma / \text{cm}^2 \sim 1 \times 10^{-17}$$

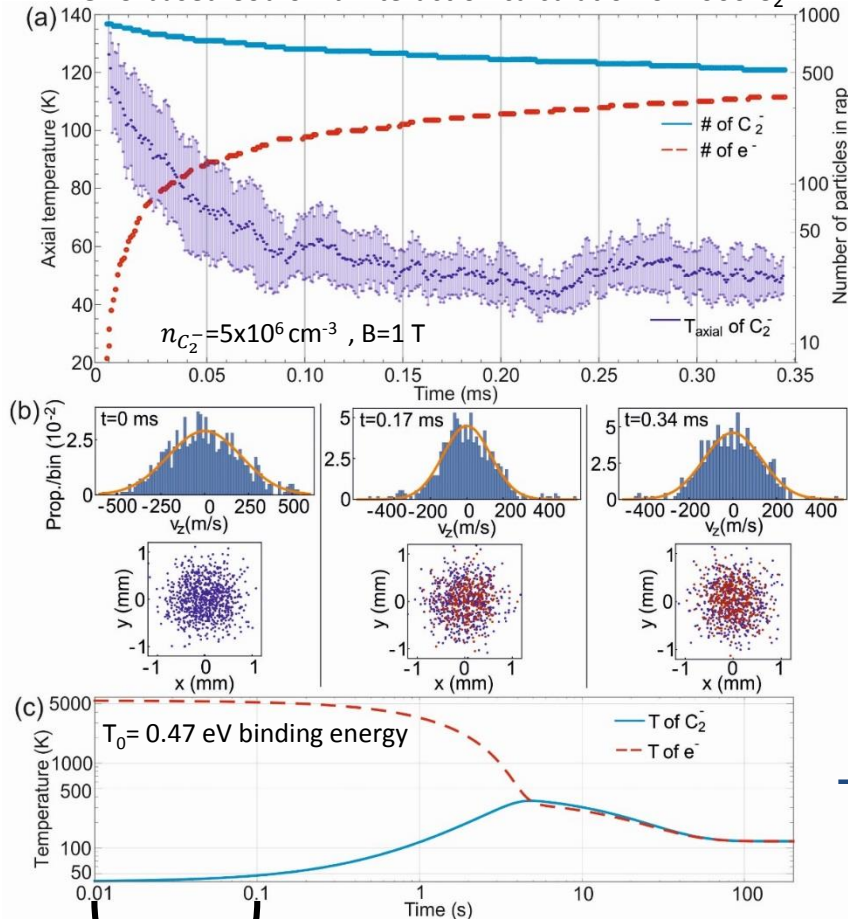




## simulation

1 x 2.5  $\mu\text{m}$  laser, 1 x photodetachment laser 399 nm (100 mW enhanced in  $F \sim 1000$  cavity)

GPU based Coulomb interaction calculation of 1000  $C_2^-$ :



- Factor  $\sim 2.5$  temperature reduction after 0.3 ms;  $e^-$  released with binding energy 0.47 eV as:  
 $C_2^- + h\nu \rightarrow C_2 + e^-$
- $e^-$  equilibrate with  $C_2^-$  in  $\sim 5 \text{ s}$  then thermalizing with 120 K environment after 100 s in 1 T.

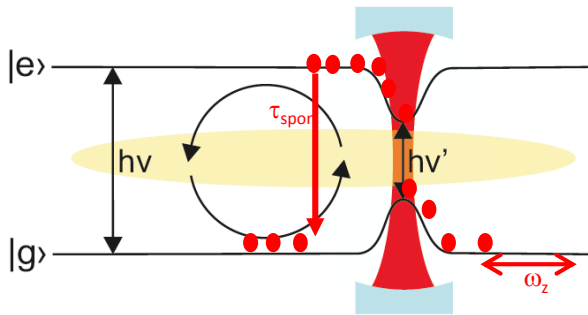
→ Simple method, but only for small time window

$\sim 100 \text{ ms}$  experimental time window at lower temperature

# Dipole force Sisyphus cooling in Penning trap

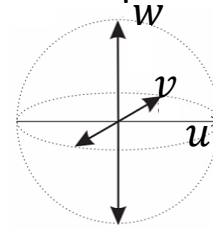


- 1 x 2.5  $\mu\text{m}$   $\sim$  1.5 GHz detuned dipole laser
- Spontaneous decay lifetime  $\tau_{\text{spont}} = 50 \mu\text{s} >$  axial trap motion
- $T_{\text{min}}$  given by potential depth  $U_{\text{dipole}}$  to  $\sim$  mK



Force of light-matter interaction for a 2 level-system:

Bloch sphere



$$|\psi\rangle\langle\psi| = \begin{pmatrix} \rho_{11} & \rho_{12} \\ \rho_{21} & \rho_{22} \end{pmatrix}$$

$$u = \rho_{12} + \rho_{21} \quad \text{in-phase}$$

$$v = -i(\rho_{12} - \rho_{21}) \quad \text{quadrature}$$

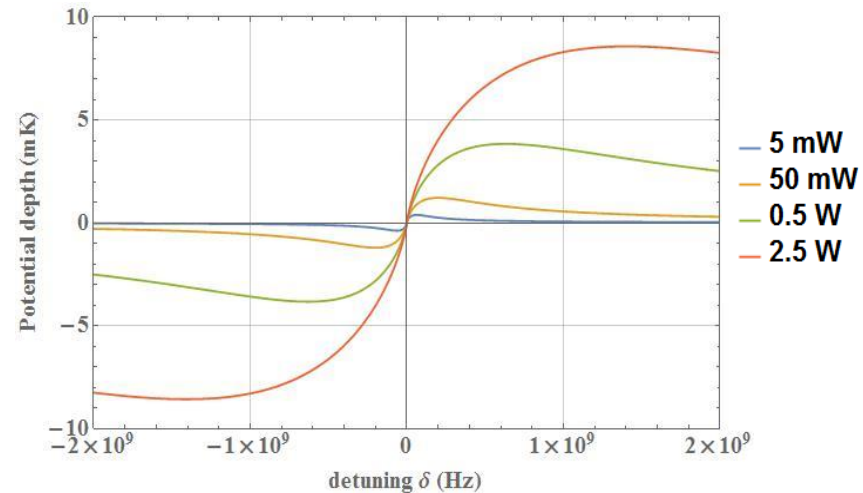
$$w = \rho_{11} - \rho_{22} \quad , \quad |R|^2 = u^2 + v^2 + w^2 = 1$$

$$F = \frac{-q \chi_{12}}{2} \left[ u \frac{\partial E}{\partial z} + v E k \right]$$

$$F = F_{\text{dipole}} + F_{\text{scatter}}$$

$$F_{\text{dipole}} = \nabla U_{\text{dipole}} \quad , \quad U_{\text{dipole}} = \hbar \delta \ln \left[ 1 + \frac{2\Omega^2}{4\delta^2 + \Gamma^2} \right]$$

Dipole potential depth vs laser detuning, waist = 0.2 mm



Sisyphus scheme: J. Dalibard and C. Cohen-Tannoudji, *J. Opt. Soc. Am. B*, **2**, 11 (1985)

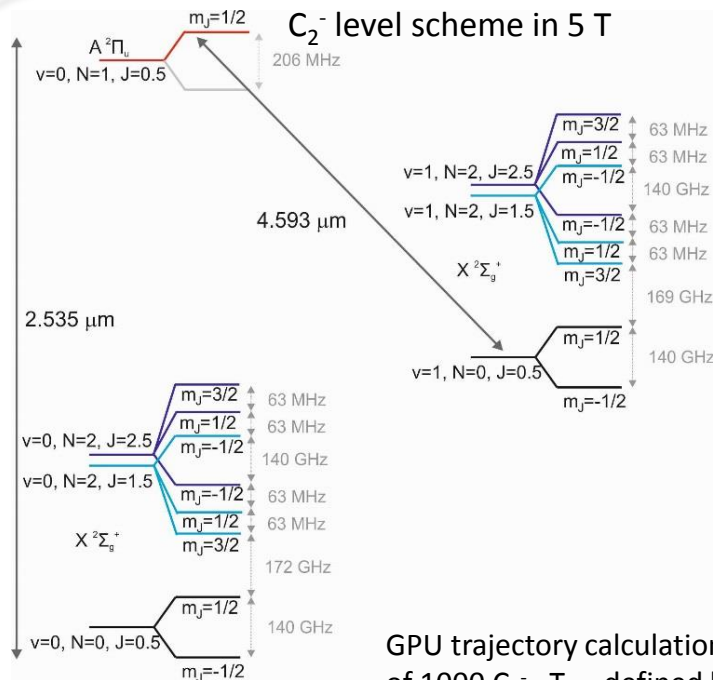
Sisyphus cooling in magnetic traps: S. Wu, W. Phillips, J. V. Porto et al., *Phys. Rev. Lett.* **106**, 213001 (2011)

# Dipole force Sisyphus cooling in Penning trap

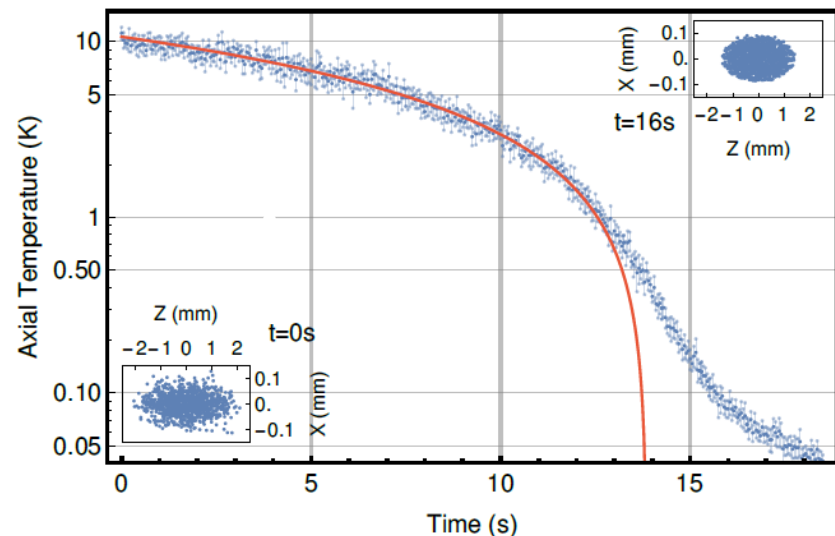


## simulation

- 1 x 2.5  $\mu\text{m}$   $\sim$  1.5 GHz detuned dipole laser ( $U_{\text{dipole}}=7$  mK)
- Spontaneous decay  $\tau_{\text{spont}}=50$   $\mu\text{s}$   $>$  axial trap motion
- 1 x pump at 2.5  $\mu\text{m}$ , 4 x repumper at 2.5  $\mu\text{m}$  and 4.5  $\mu\text{m}$



$n_{C_2^-}=10^7$   $\text{cm}^{-3}$ ,  $B=5$  T, Dipole=2.5 W, Pump=10 mW,  $w=0.2$  mm,  $F=1000$



GPU trajectory calculation including Monte Carlo for optical transitions of 1000  $C_2^-$ .  $T_{\text{min}}$  defined by potential depth  $U_{\text{dipole}}$  to  $\sim 1$  mK

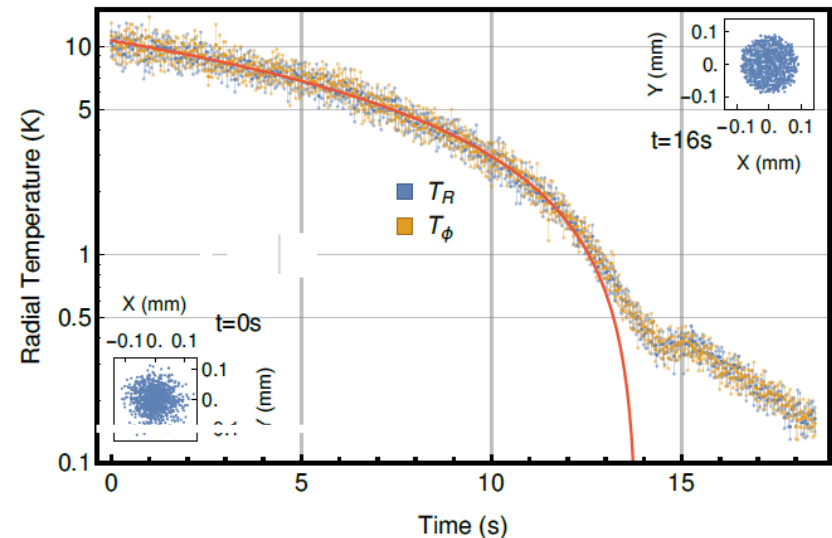
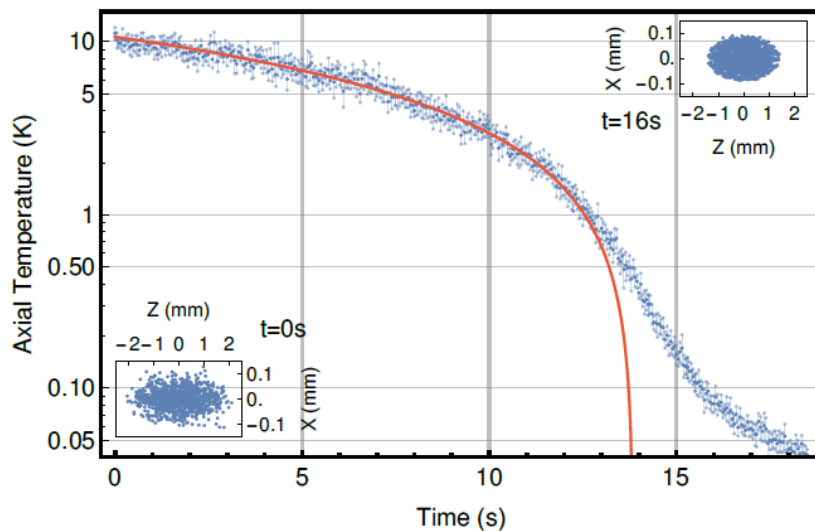
# Dipole force Sisyphus cooling in Penning trap



## simulation

- 1 x 2.5  $\mu\text{m}$   $\sim$ 1.5 GHz detuned dipole laser ( $U_{\text{dipole}}=7$  mK)
- Spontaneous decay  $\tau_{\text{spont}}=50$   $\mu\text{s}$   $>$  axial trap motion
- 1 x pump at 2.5  $\mu\text{m}$ , 4 x repumper at 2.5  $\mu\text{m}$  and 4.5  $\mu\text{m}$

$n_{\text{C}_2^-}=10^7$   $\text{cm}^{-3}$ ,  $B=5$  T, Dipole=2.5 W, Pump=10 mW,  $w=0.2$  mm,  $F=1000$

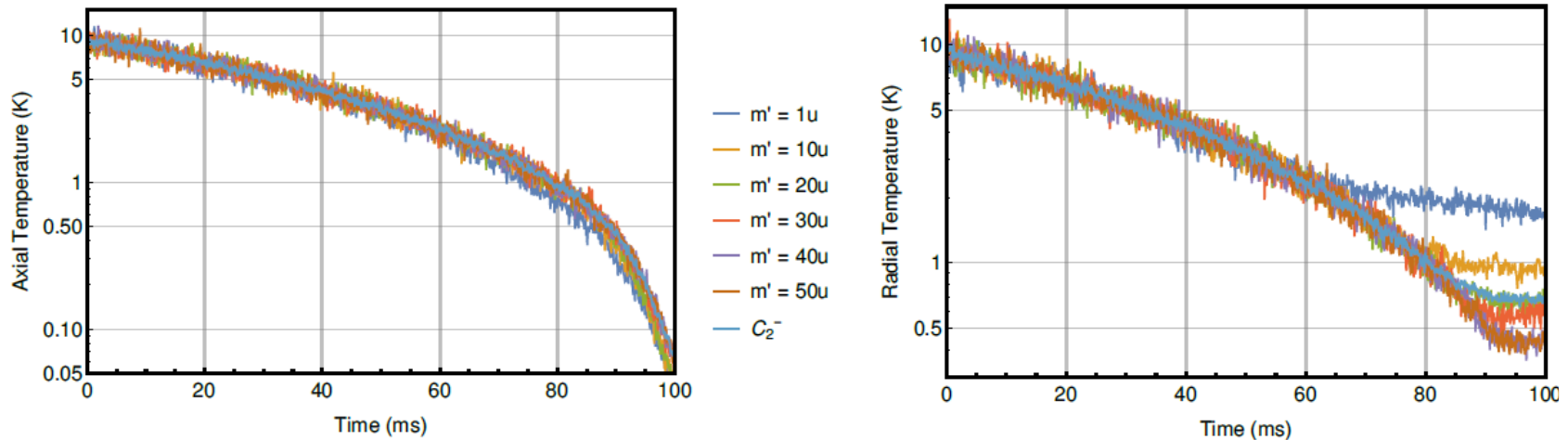


GPU trajectory calculation including Monte Carlo for optical transitions of 1000  $\text{C}_2^-$ .  $T_{\text{min}}$  defined by potential depth  $U_{\text{dipole}}$  to  $\sim$ 1 mK



## simulation

$n_{C_2^-} = 10^7 \text{ cm}^{-3}$ ,  $B = 5 \text{ T}$ , Dipole = 2.5 W, Pump = 10 mW,  $w = 0.2 \text{ mm}$ ,  $F = 1000$



GPU trajectory calculation including Monte Carlo for optical transitions of  $1000 C_2^- / 100 m'$ .  $T_{\min}$  defined by potential depth  $U_{\text{dipole}}$  to  $\sim 1 \text{ mK}$

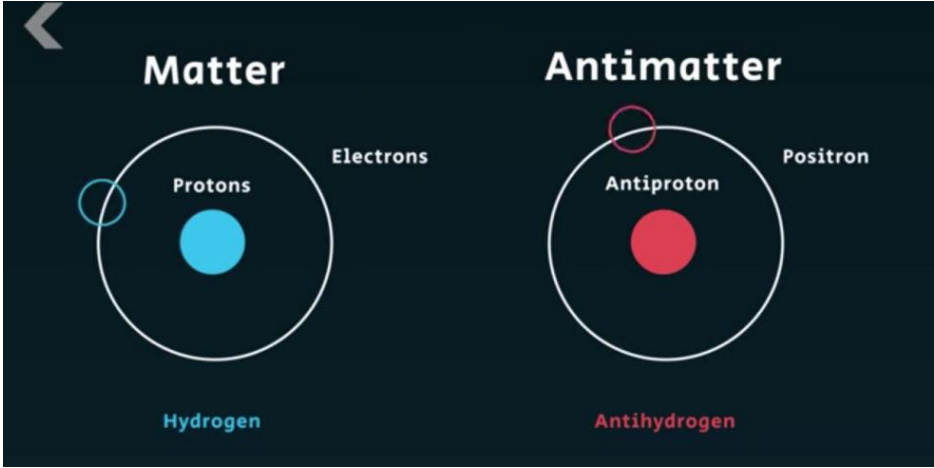
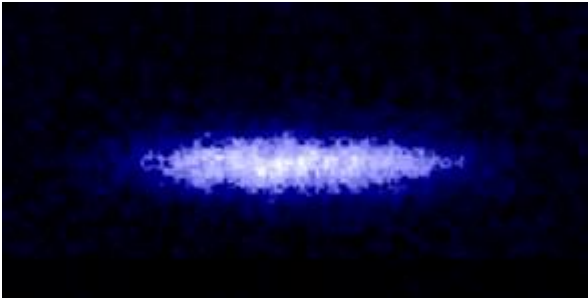
→ mK  $\bar{p}$  feasible (scatter  $\sim 10^4$  photons, 6x lasers drift stabilized to cell/wavemeter)

→ sub-K  $\bar{H}$  production becomes feasible with current Ps technology ( $e^+$  are 1822x lighter)

Goal for  $\bar{H}$ :  $T_{\bar{p}} \sim 10 \text{ mK}$  limited by 100 K Ps momentum transfer to 100 mK  $\bar{H}$



# Stay tuned for ultracold antiproton and antihydrogen physics



Thank you very much for your attention!