



Dynamic DM Production

Finite Temperature Effects and the Vev Flip Flop

Michael J. Baker

Blois, 06 June 2018

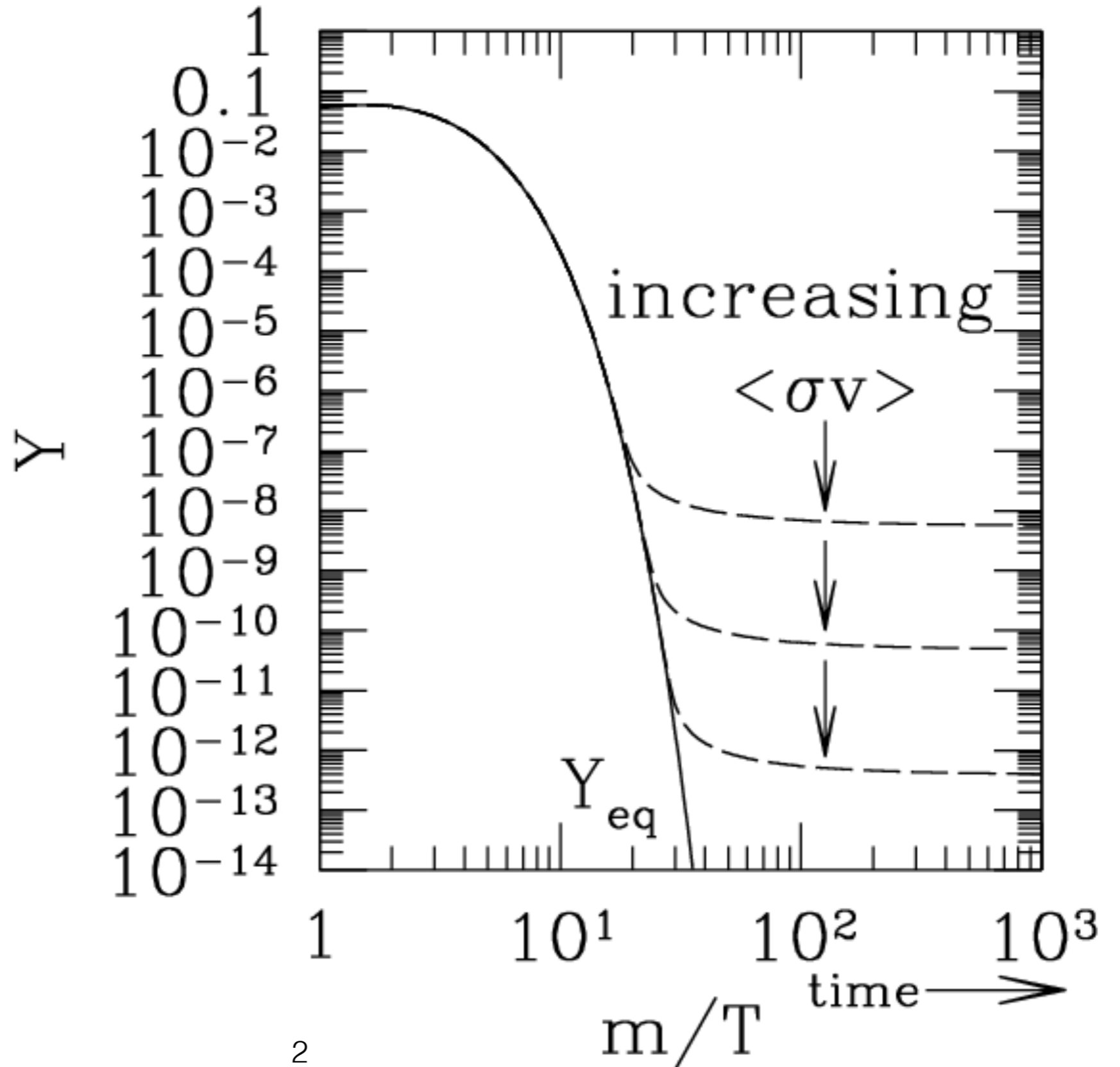
Based on [1608.07578](#) (PRL)
and [1712.03962](#) (JHEP)

MJB, M. Breitbach, L. Mittnacht, J. Kopp

- The thermal WIMP paradigm

1009.3690 — Gondolo & Gelmini

- 1) Assume a new, heavy particle is in thermal equilibrium in the early universe.

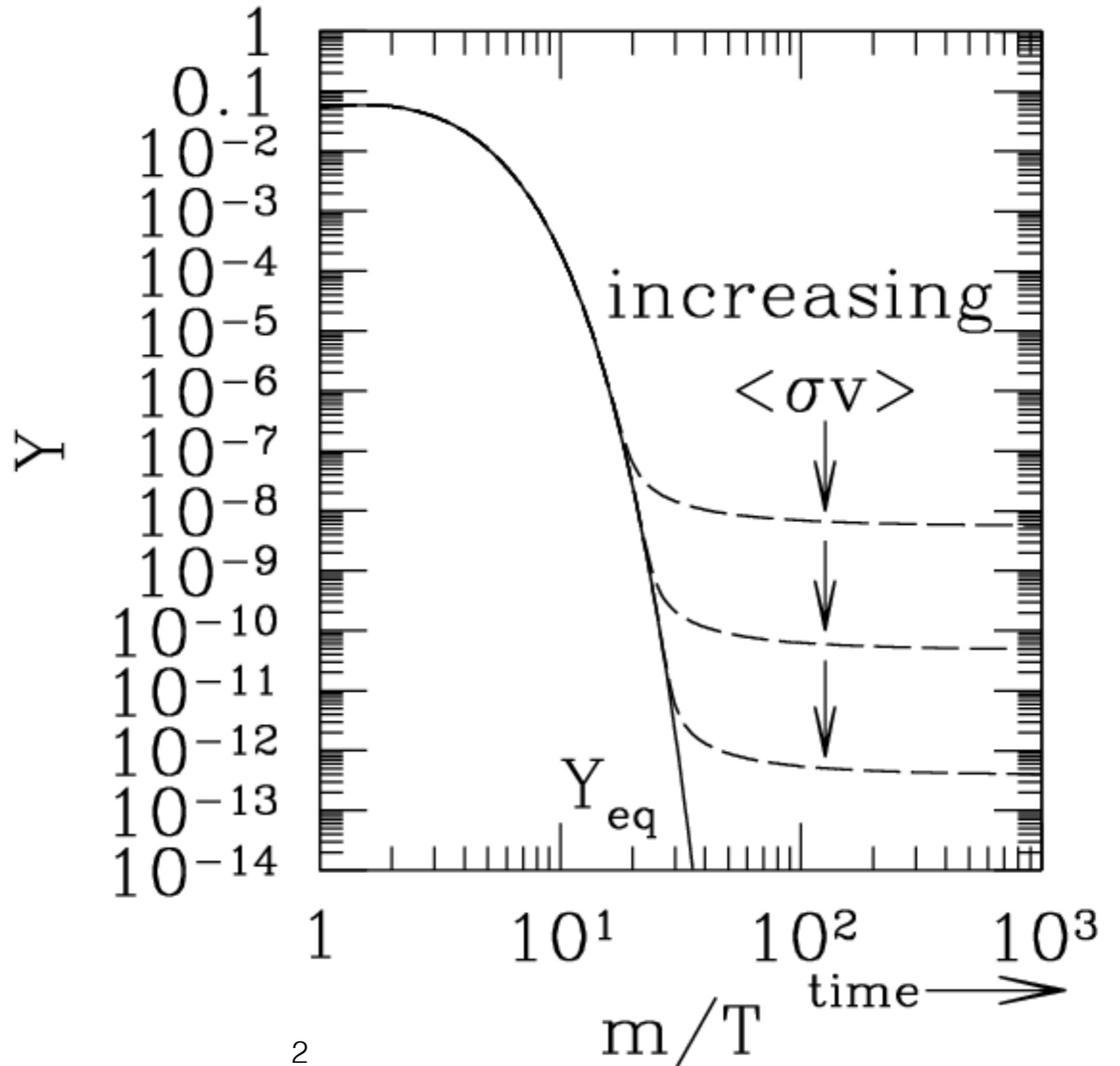


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- 2) Expansion rate greater than annihilation rate leads to freeze-out.

$$H \gtrsim \Gamma$$



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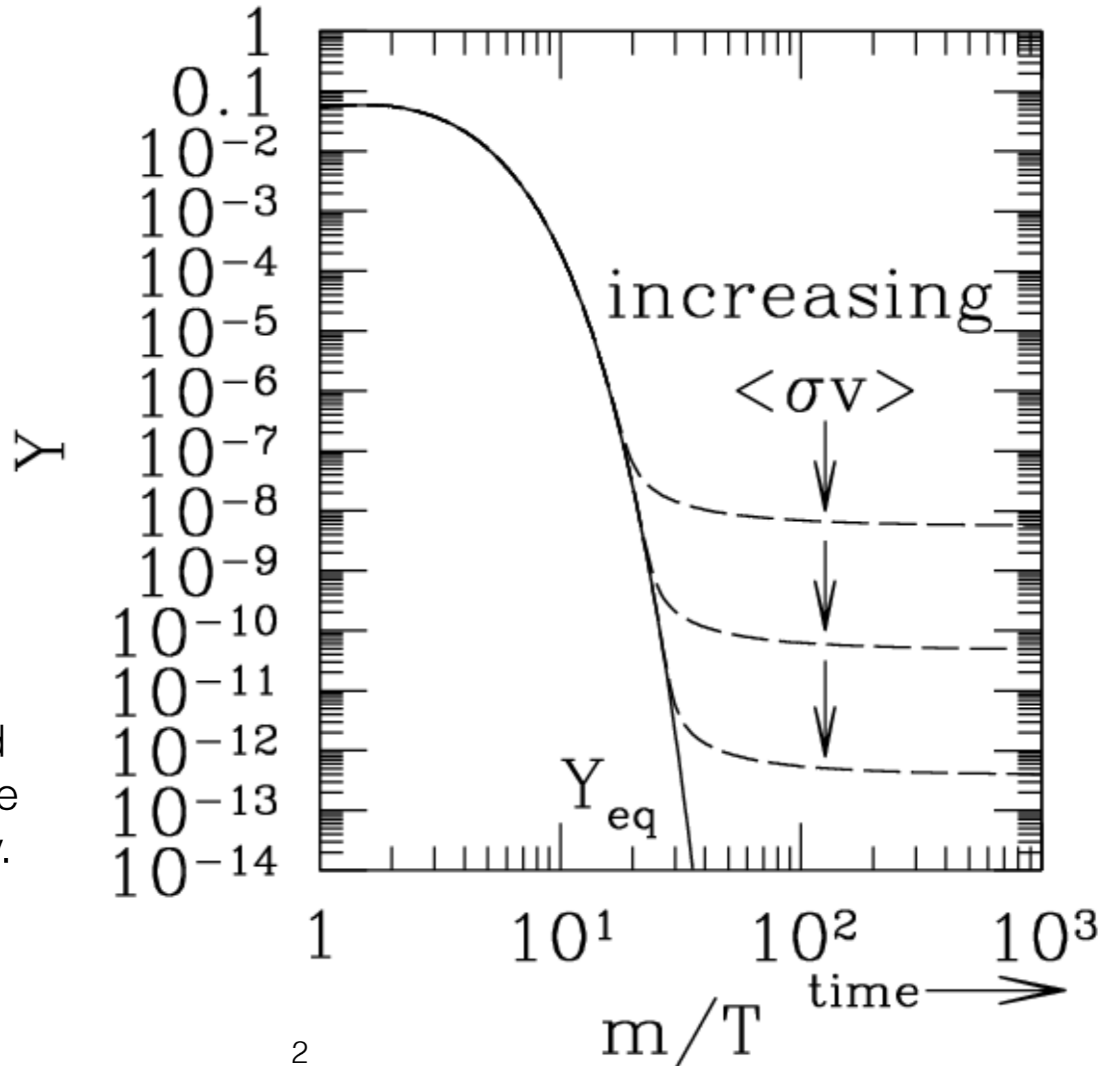
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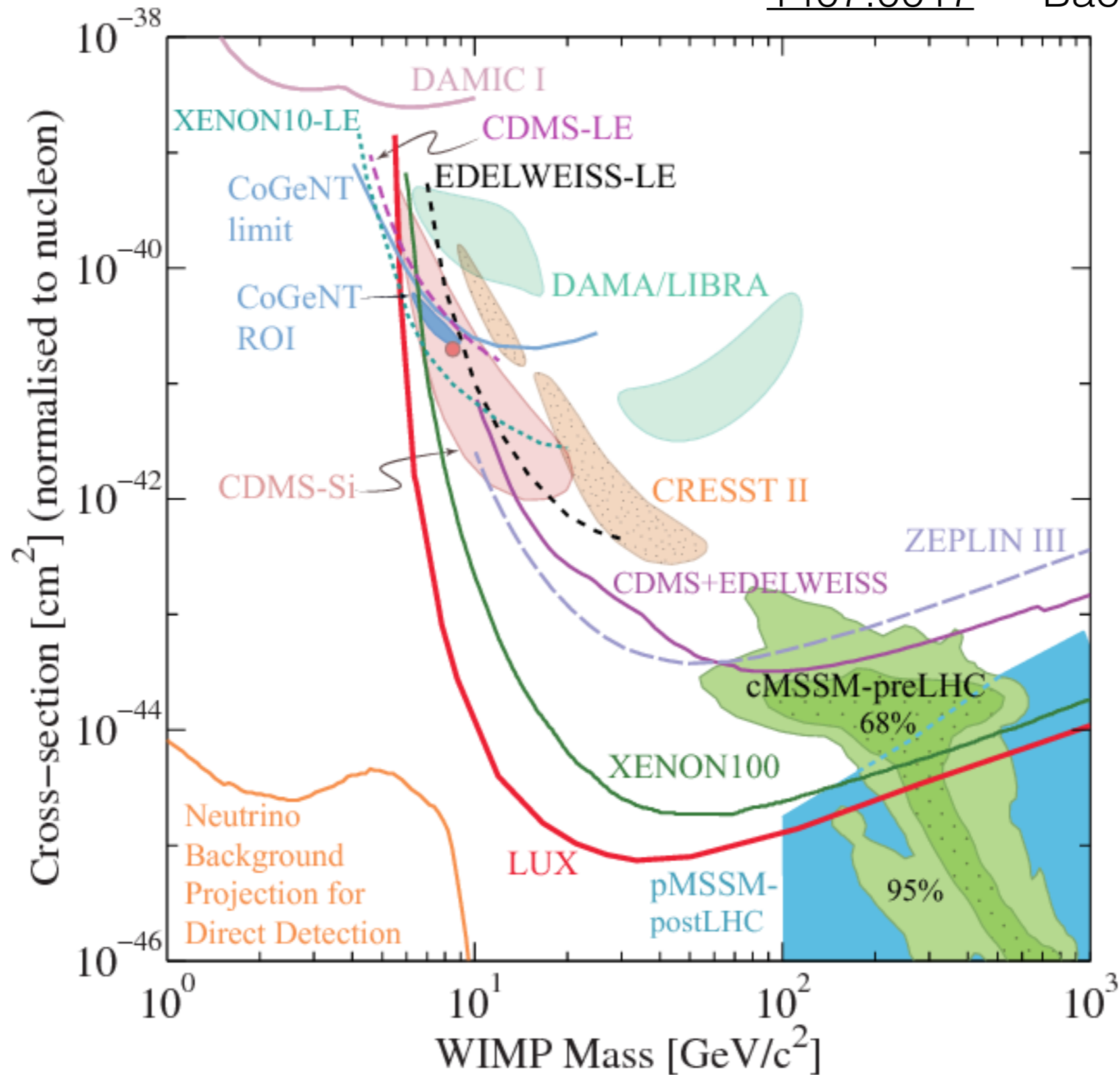
$$H \gtrsim \Gamma$$

3) Weak scale masses and order one couplings give ~ observed relic density.



- The thermal WIMP paradigm

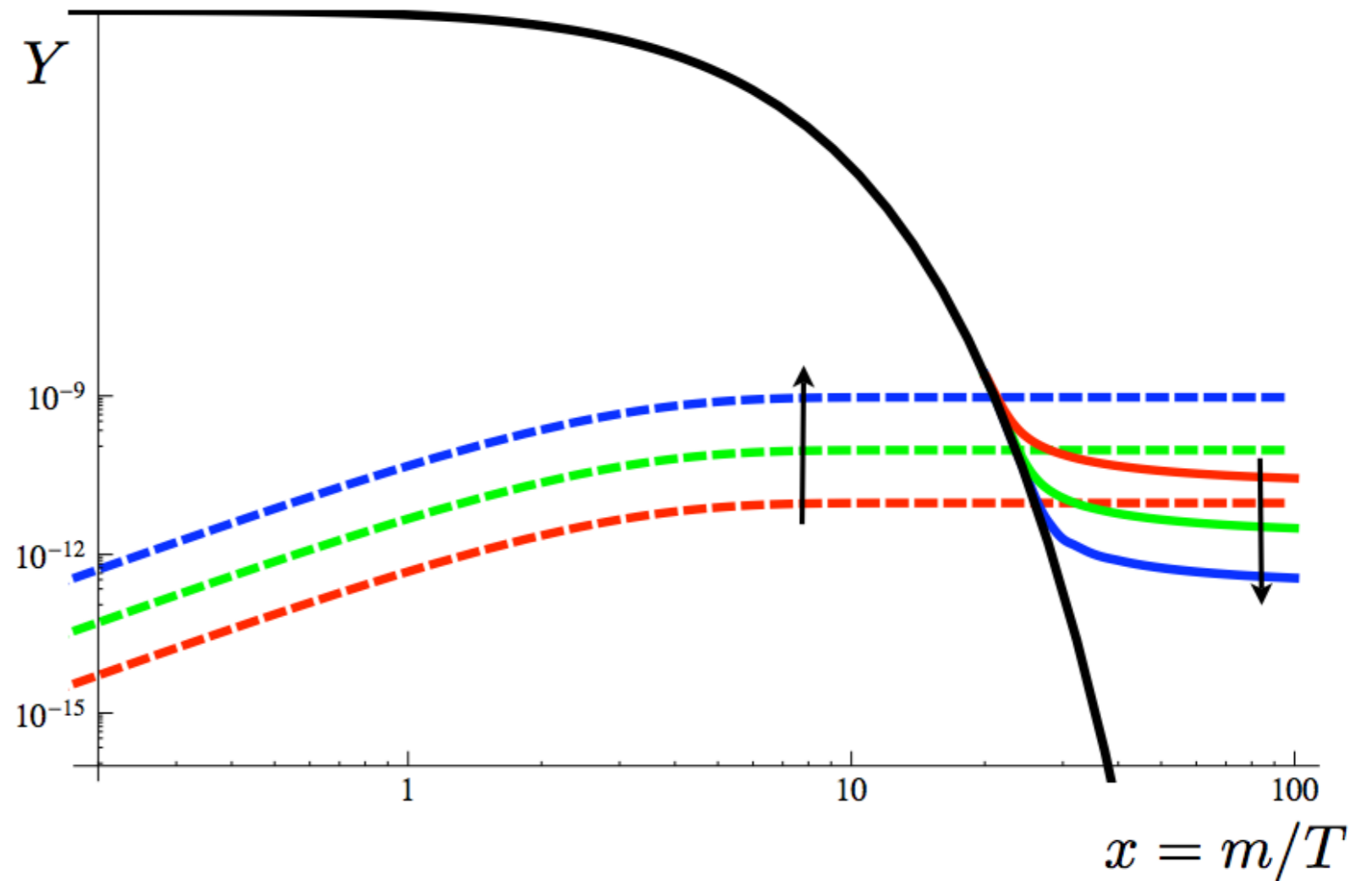
1407.0017 — Baer, Howard et al.



- Thermal Freeze-in

0911.1120 — Hall, Jedamzik, March-Russell & West

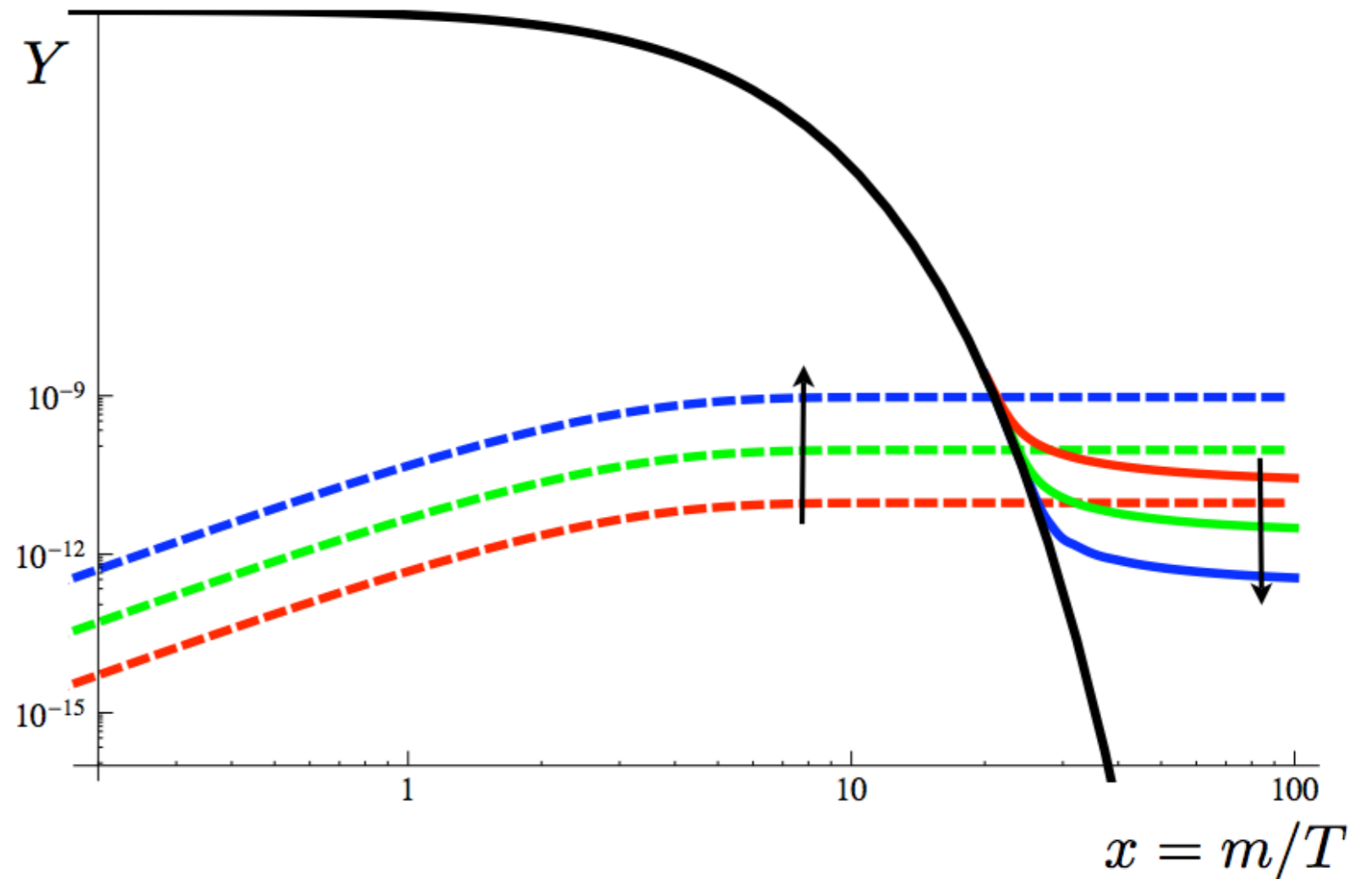
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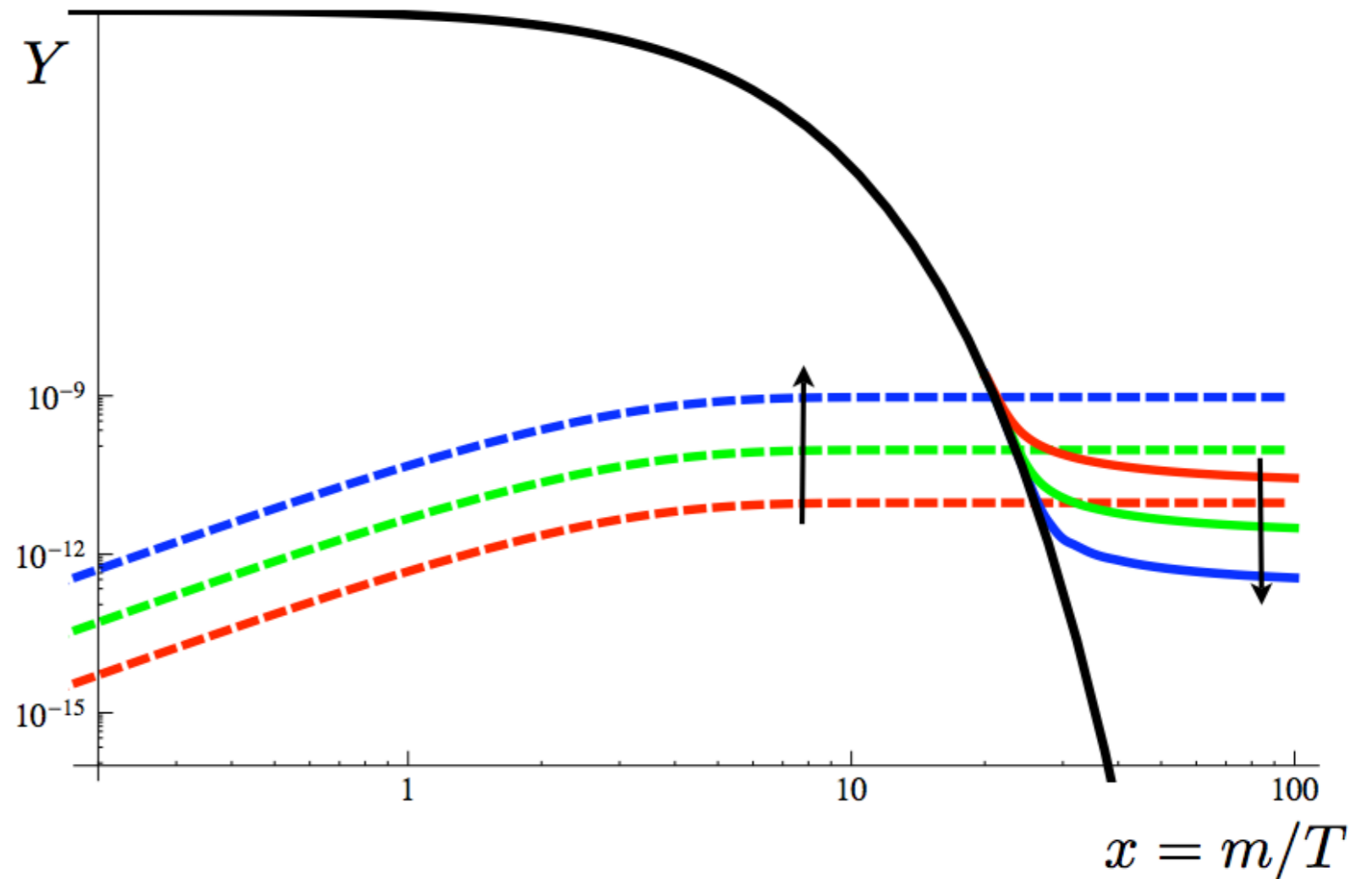
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• Thermal Freeze-in

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- 1) Assume the new, heavy particle initially has negligible abundance.
- 2) If it interacts very weakly with the thermal bath, it will come towards equilibrium, but it won't reach it.
- 3) Whereas freeze-out is determined by physics at $x \sim 20 - 30$, freeze-in is determined at $x \sim 2 - 5$.



- Overview
-

In this talk we will discuss three BSM scenarios:

- Kinematically induced freeze-in
- Vev induced freeze-in
- Decaying dark matter

- Overview
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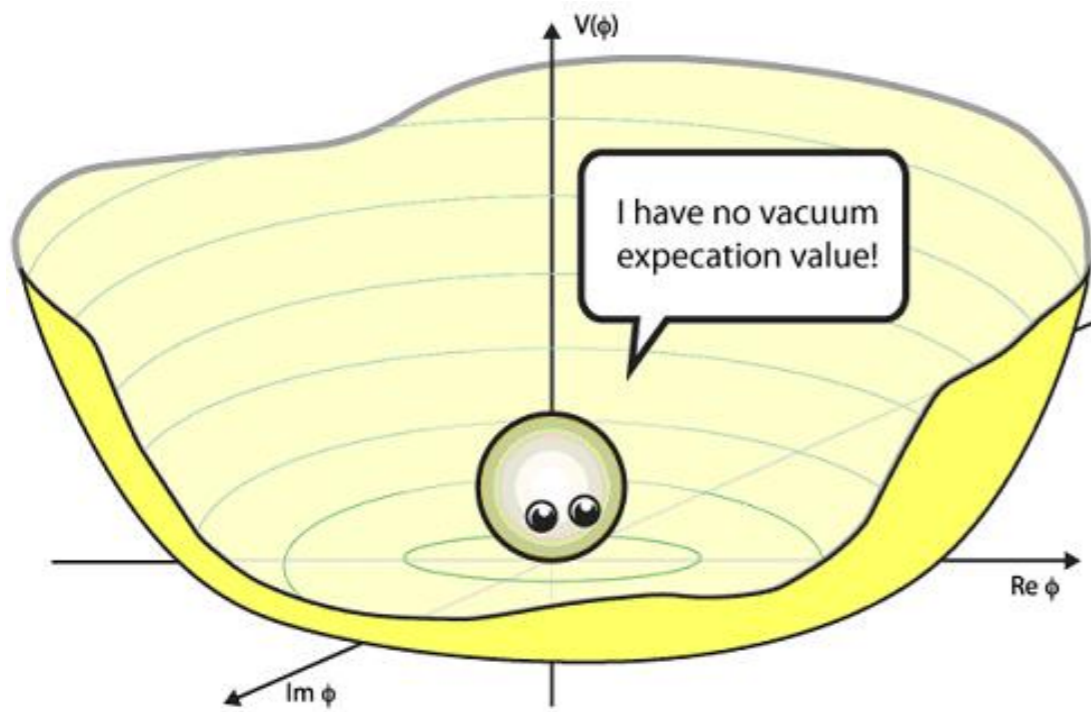
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Our main tools will be:

- The effective potential at finite temperature
- Boltzmann equations

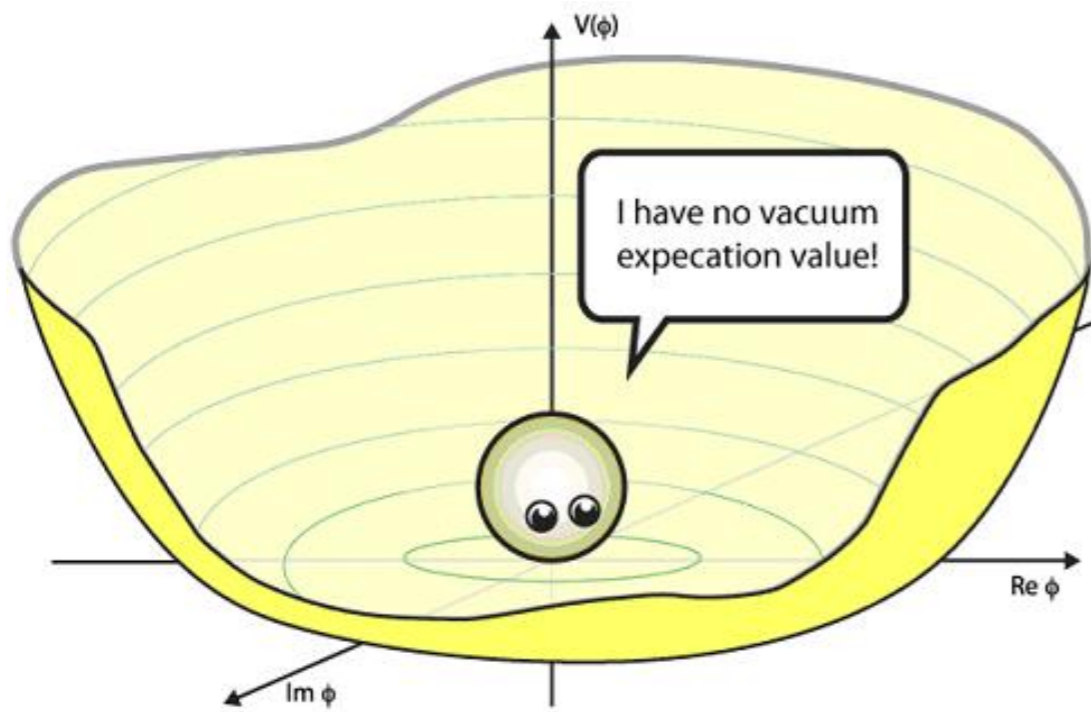
- The Standard Model Effective Potential

$$T \gg T_{EW}$$

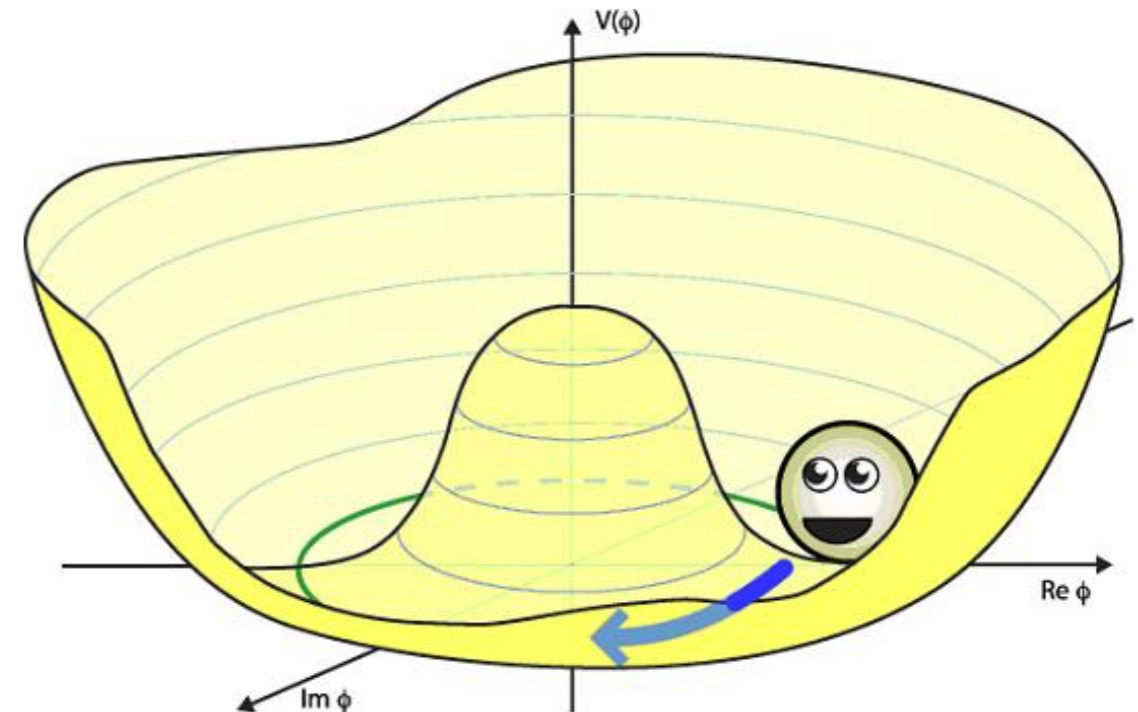


- The Standard Model Effective Potential

$$T \gg T_{EW}$$



$$T \sim 0$$



Kinematically Induced Freeze-in

- Kinematically Induced Freeze-in

Field	Spin	\mathbb{Z}_2	mass scale
S	0	+1	$m_S \sim 5 \text{ GeV}$
χ	$\frac{1}{2}$	-1	$m_\chi \sim 50 \text{ GeV}$
ψ	$\frac{1}{2}$	-1	$m_\psi \sim 50 \text{ GeV}$

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 & + \frac{\lambda_{S3}}{3!} \mu_S S^3 + \lambda_{p3} \mu_S S (H^\dagger H) + \frac{\lambda_{p4}}{2} S^2 (H^\dagger H)
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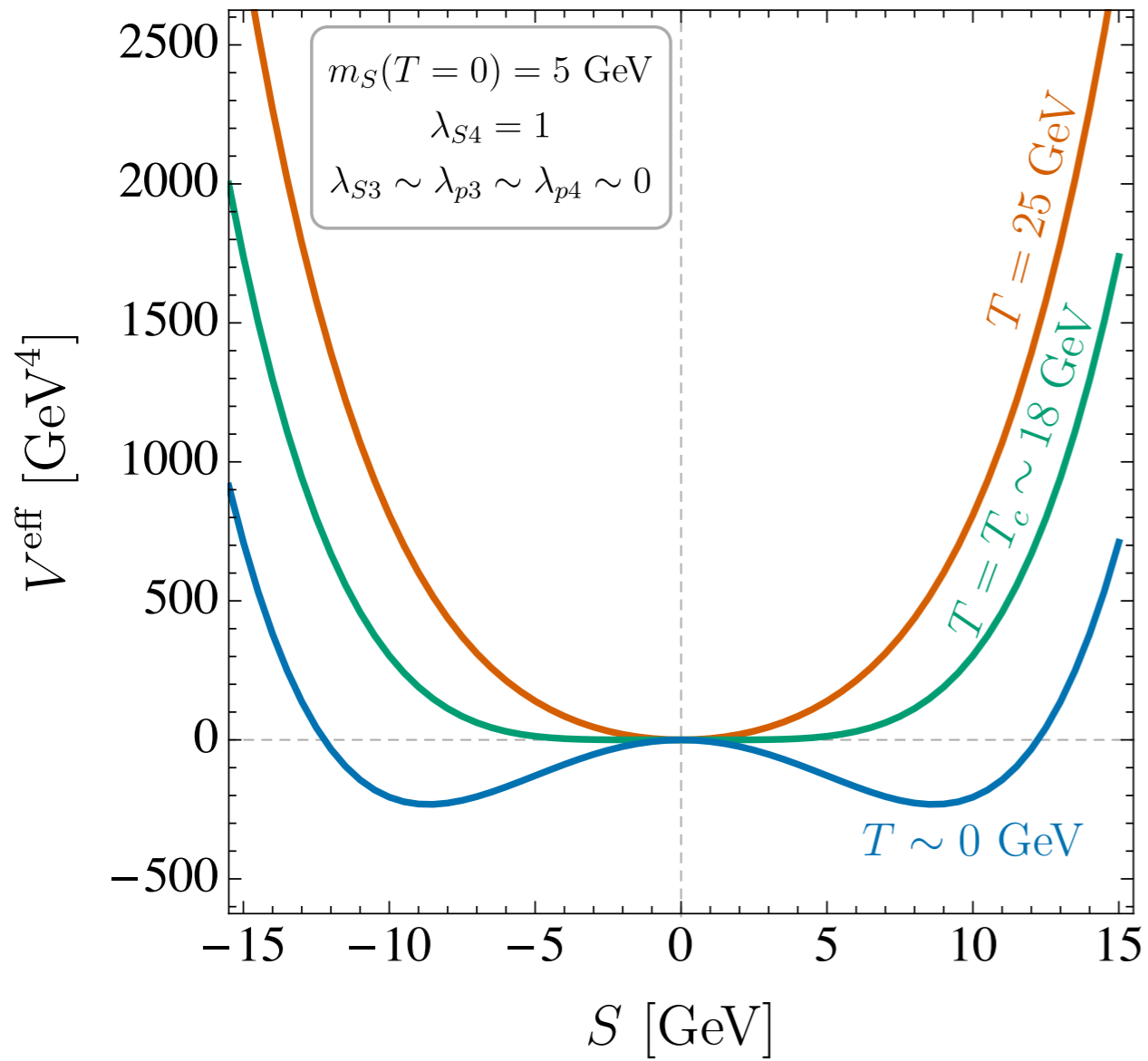
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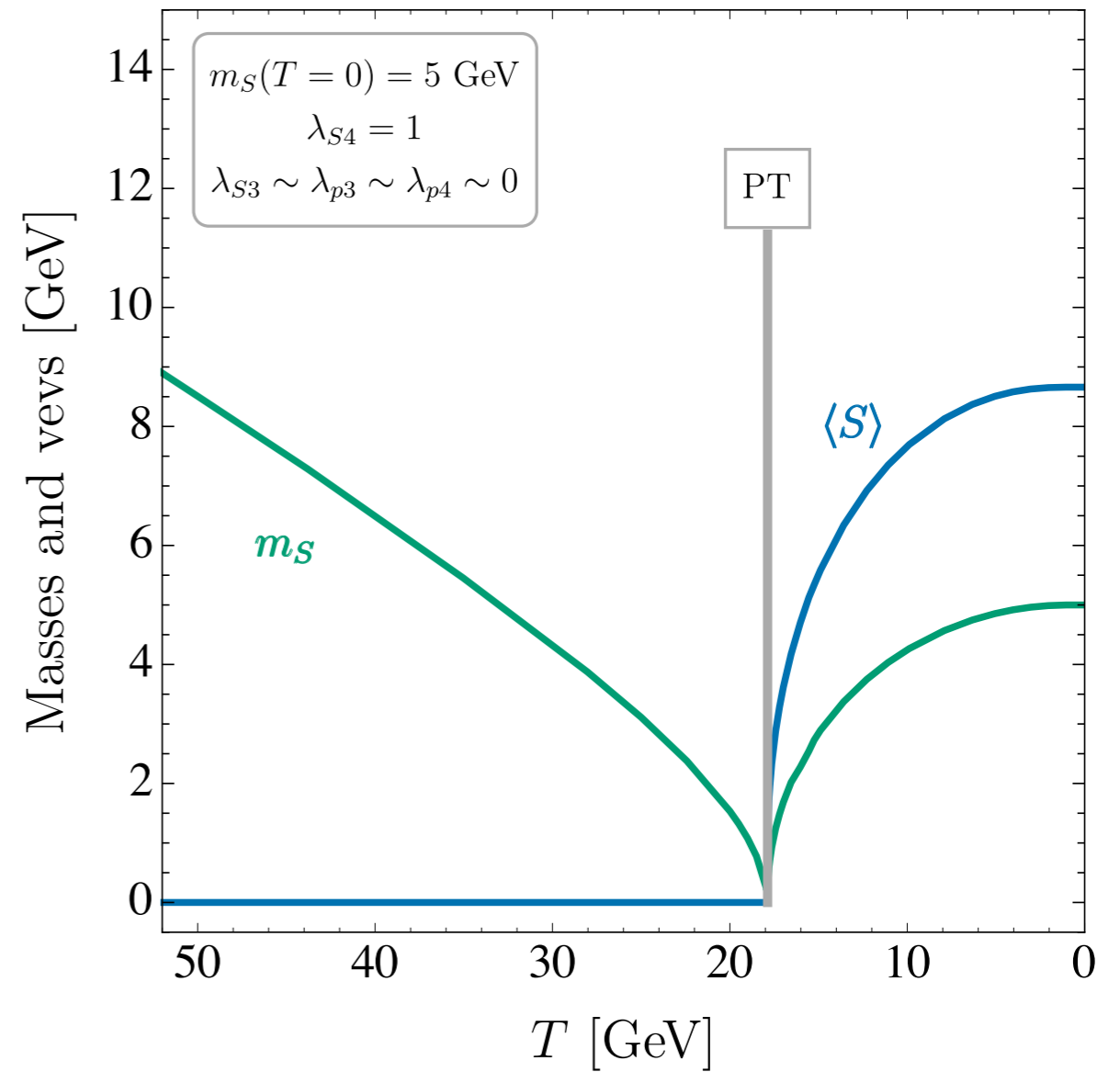
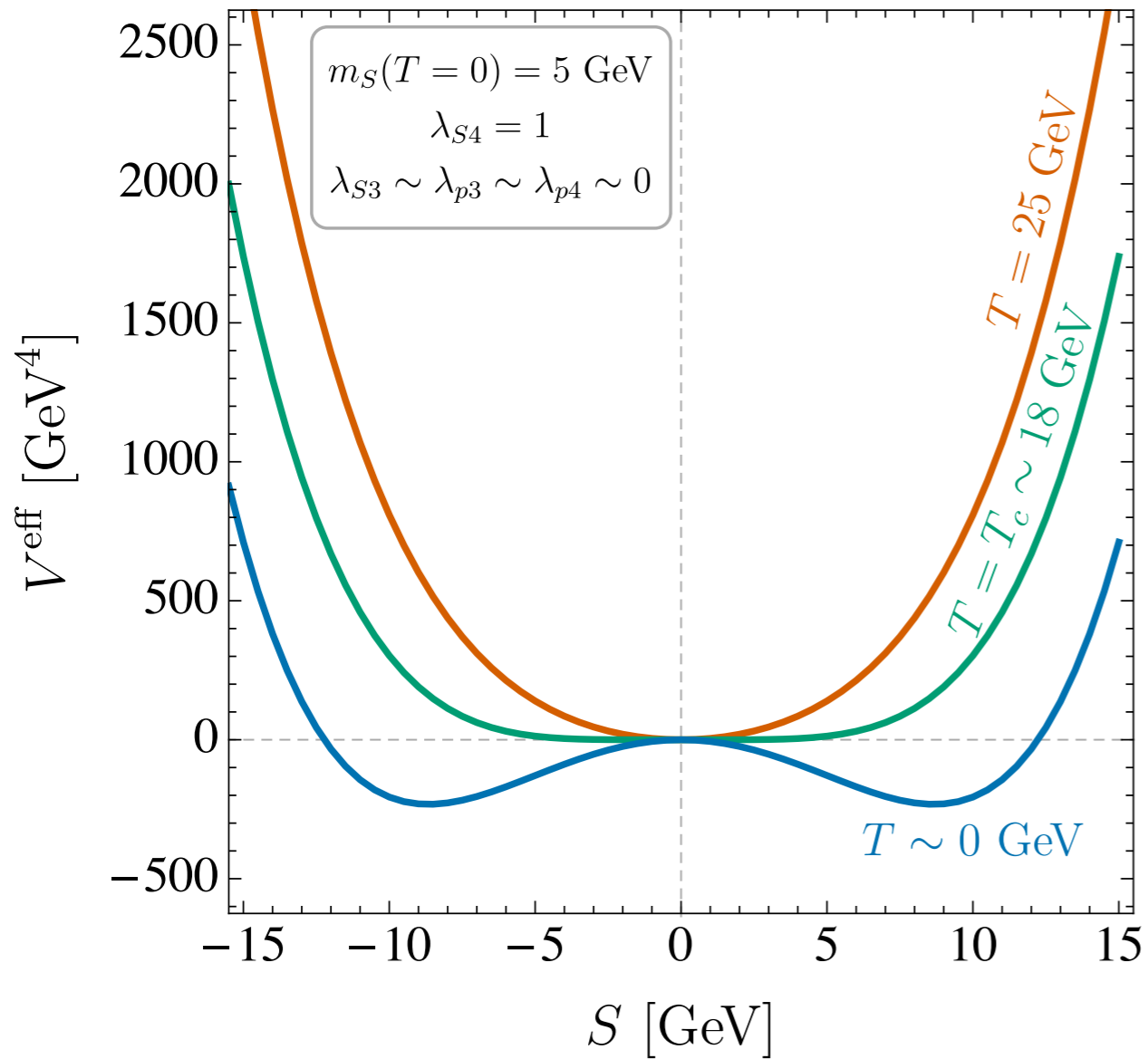
$$\lambda_{S3} \sim 0 \quad \lambda_{S4} \sim 1$$

$$\lambda_{p3} \sim 0 \quad \lambda_{p4} \sim 10^{-3}$$

- The Effective Potential - Second Order PT



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$$\mathcal{L} \supset y_\chi \bar{\chi} S \chi + y_\psi \bar{\psi} S \psi + [y_{\psi\chi} \bar{\psi} S \chi + h.c.]$$

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$$y_{\psi\chi} \sim 10^{-12}$$

$$y_\chi \sim 0$$

$$y_\psi \sim 0.01$$

- Kinematically Induced Freeze-in

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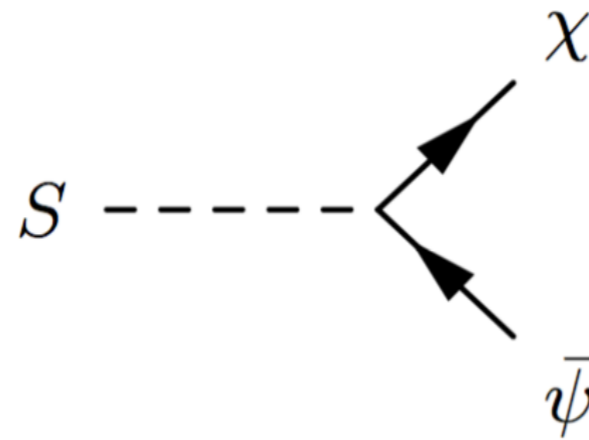
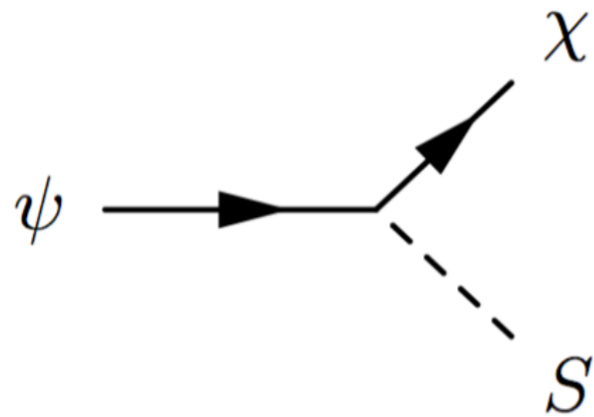
we'll assume ψ stays in thermal equilibrium

- Kinematically Induced Freeze-in
-

Assume χ begins with negligible abundance

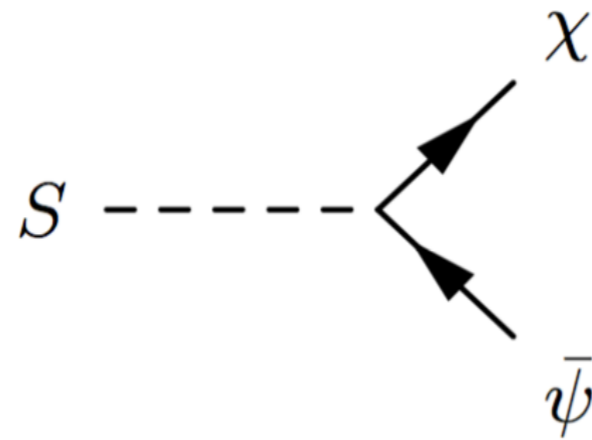
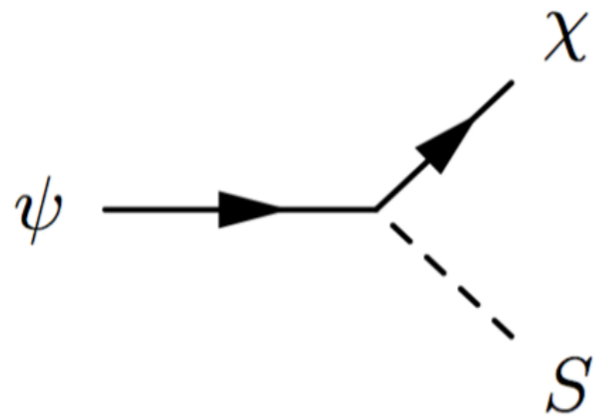
- Kinematically Induced Freeze-in

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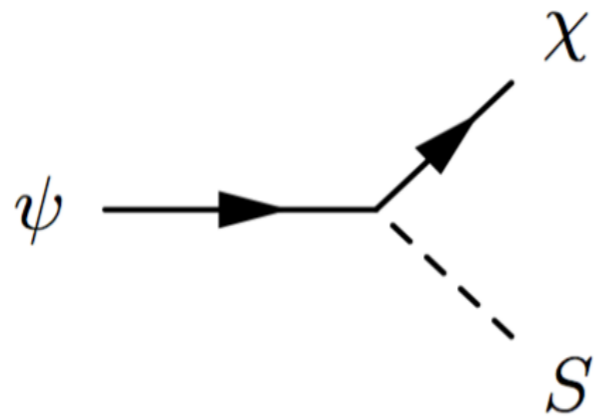
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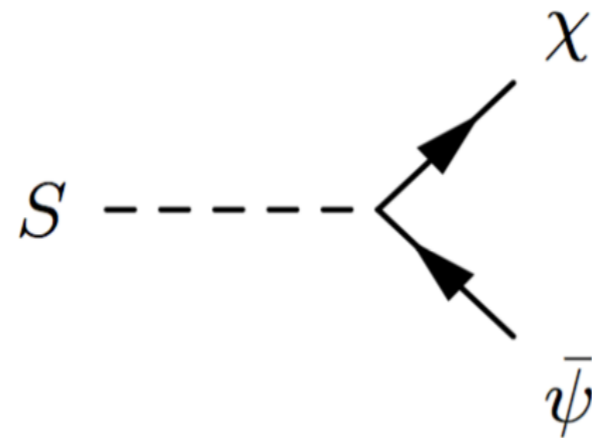
$$m_\psi > m_\chi + m_S(T)$$

- Kinematically Induced Freeze-in

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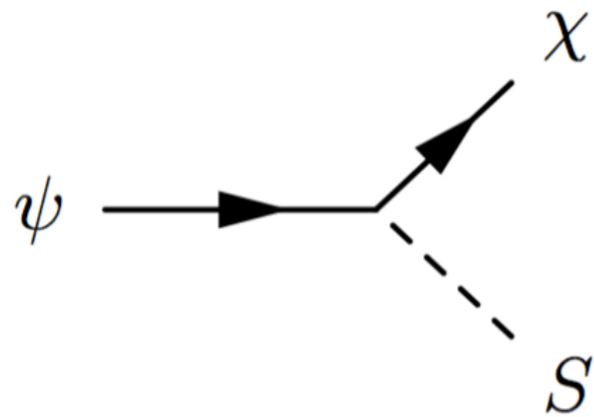
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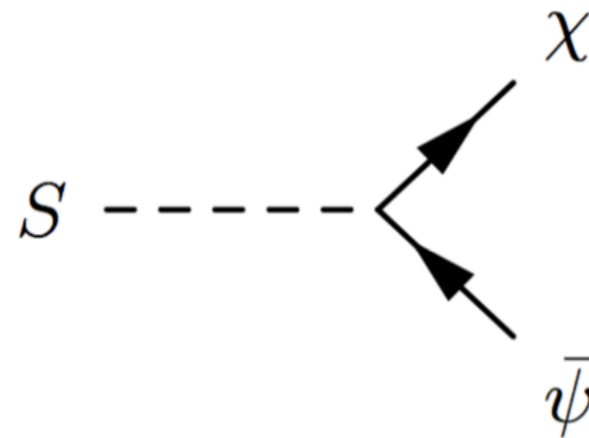
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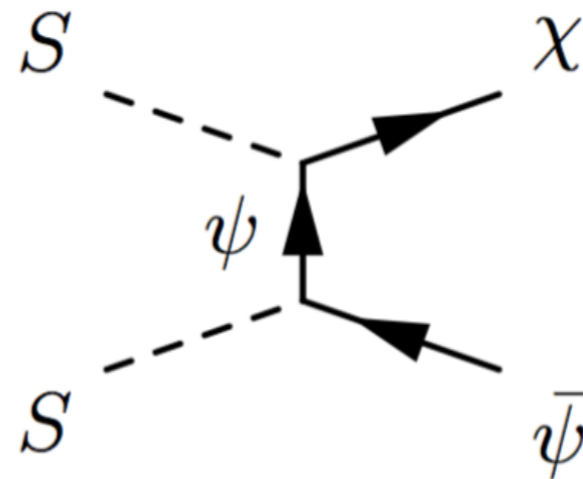
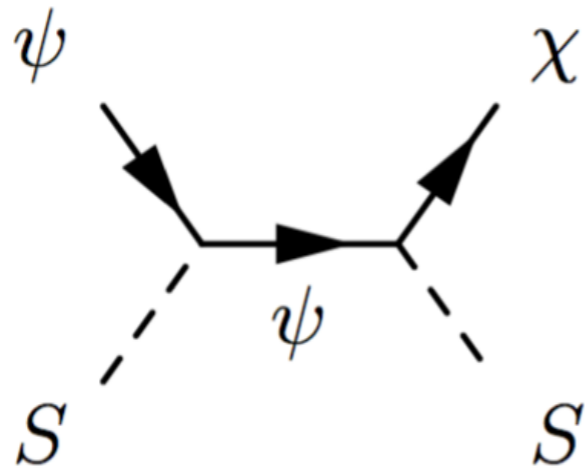
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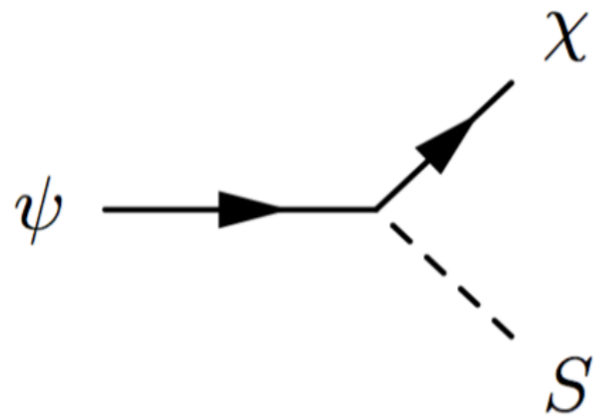


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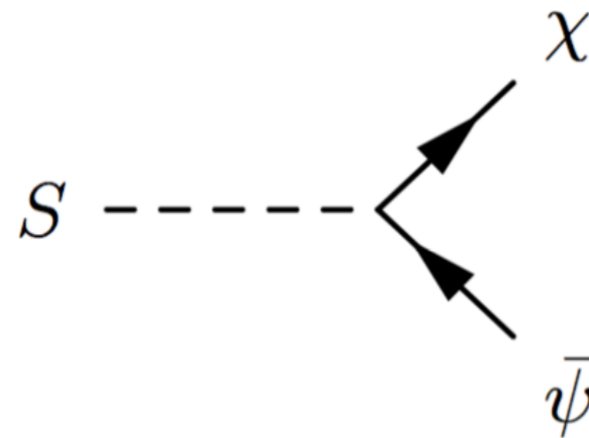


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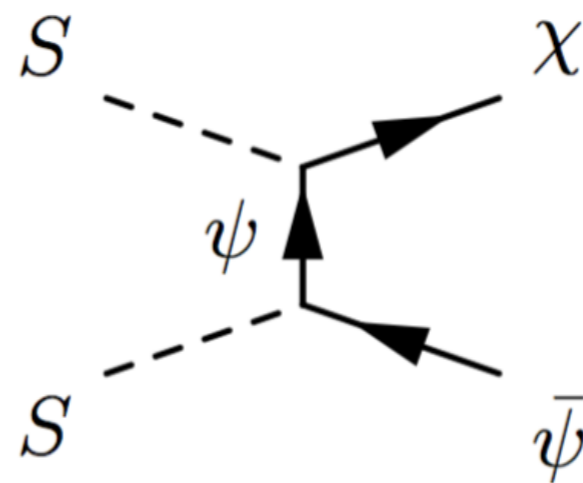
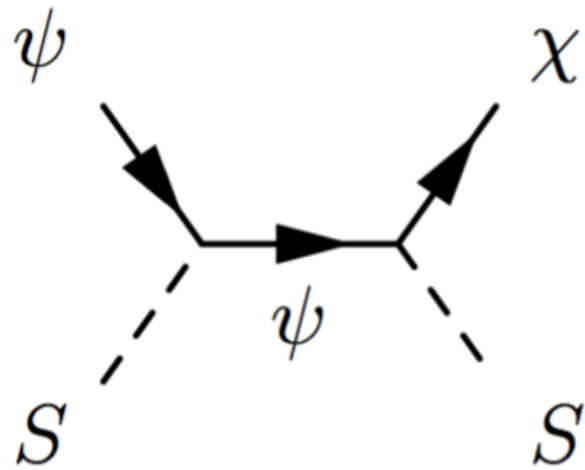
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$$m_S(T) > m_\chi + m_\psi$$



$$T \gtrsim m_\chi$$

- Freeze-in - Boltzmann Equations

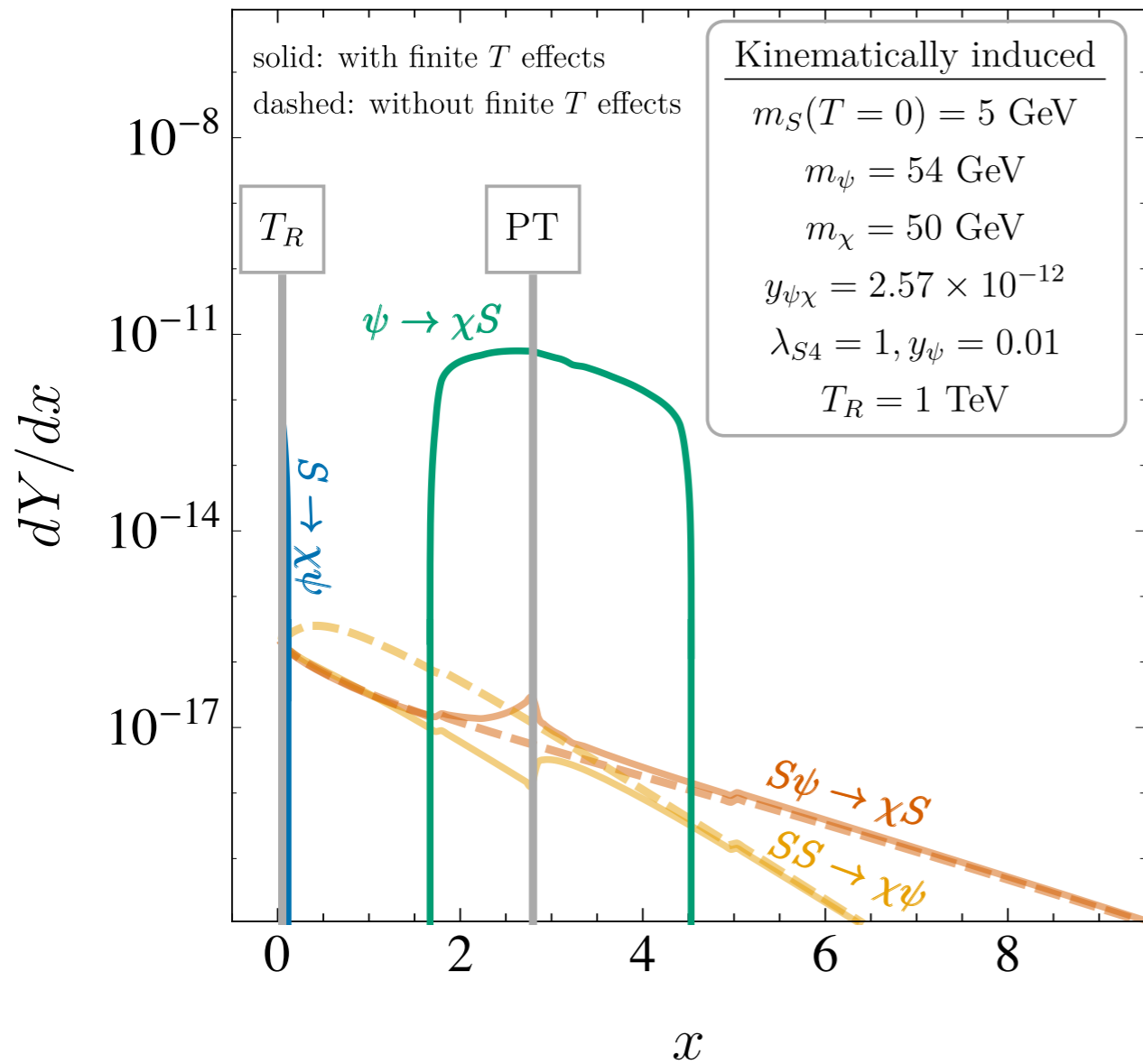
$$\frac{dY_\chi}{dx} = \frac{g_{B_1} m_{B_1}^2}{2\pi^2} \frac{m_\chi}{H(x) s(x) x^2} \Gamma(B_1 \rightarrow \chi B_2) K_1\left(\frac{m_{B_1} x}{m_\chi}\right)$$

- Freeze-in - Boltzmann Equations

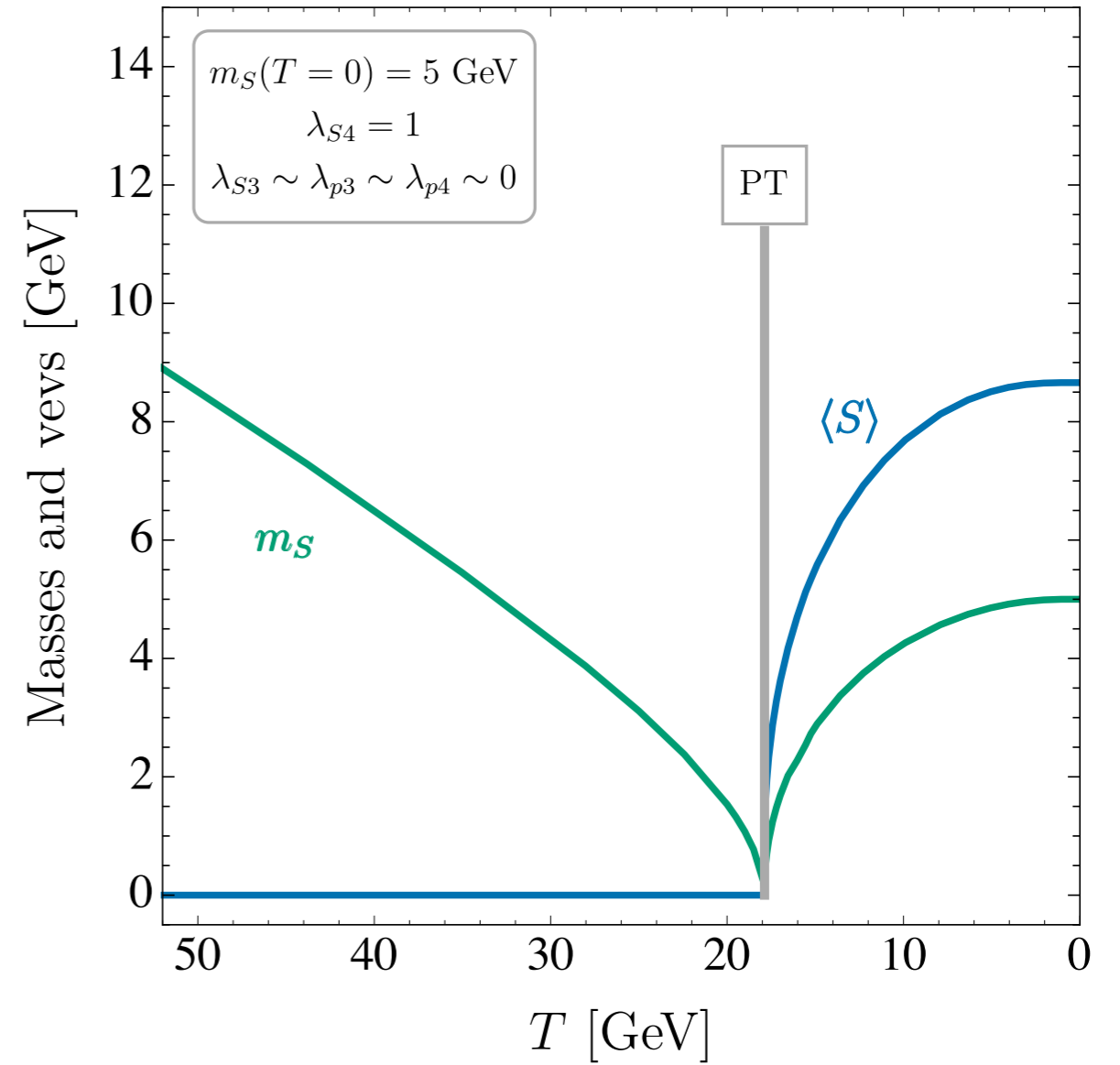
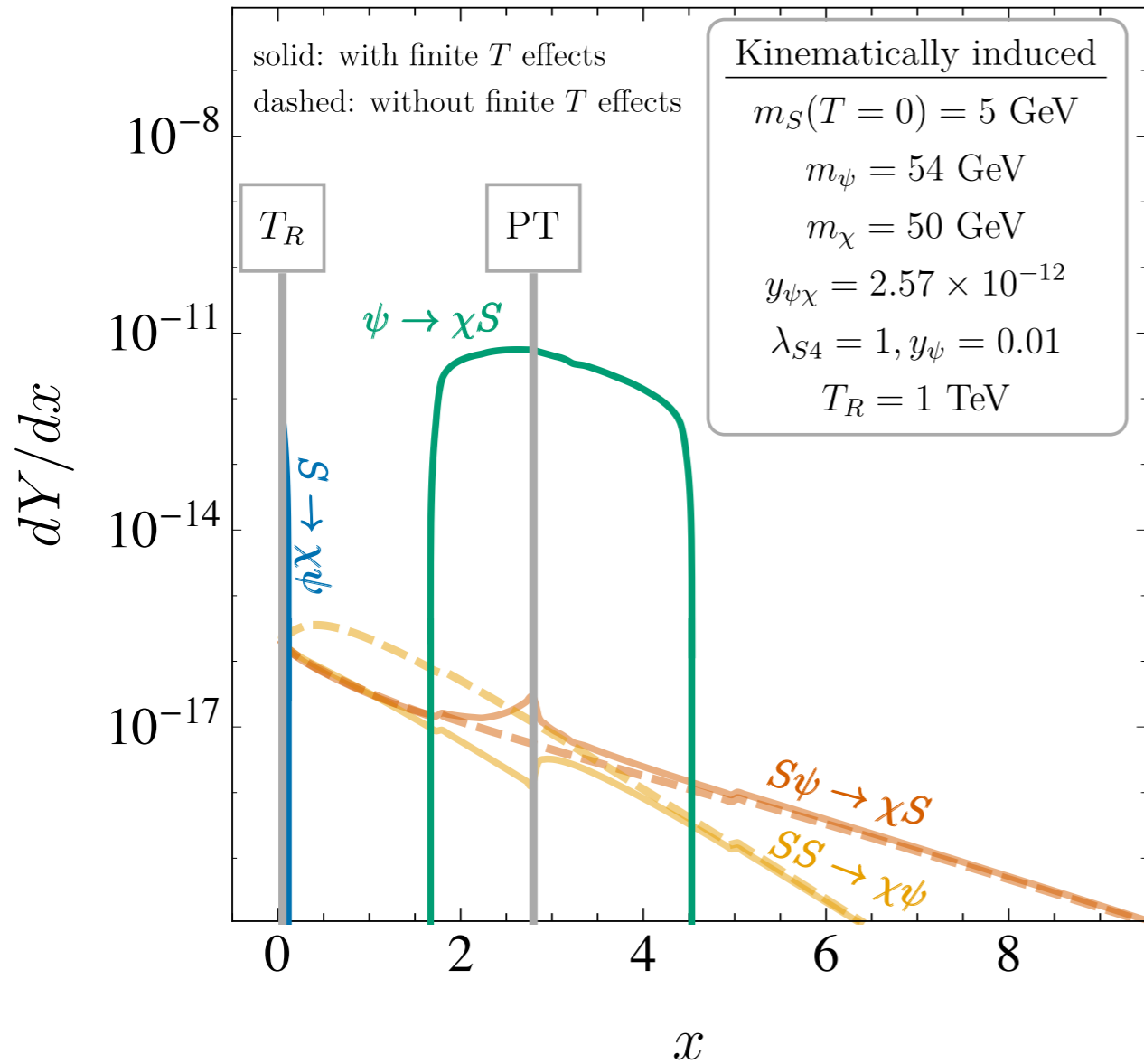
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$$\frac{dY_\chi}{dx} = \frac{g_{B_1} g_{B_2}}{32\pi^4} \frac{m_\chi}{H(x)s(x)x^2} \int_{(m_{B_1} + m_{B_2})^2}^{\infty} ds 4p_{B_1 B_2}^2 \sigma(B_1 B_2 \rightarrow \chi B_3) K_1\left(\frac{x\sqrt{s}}{m_\chi}\right)$$

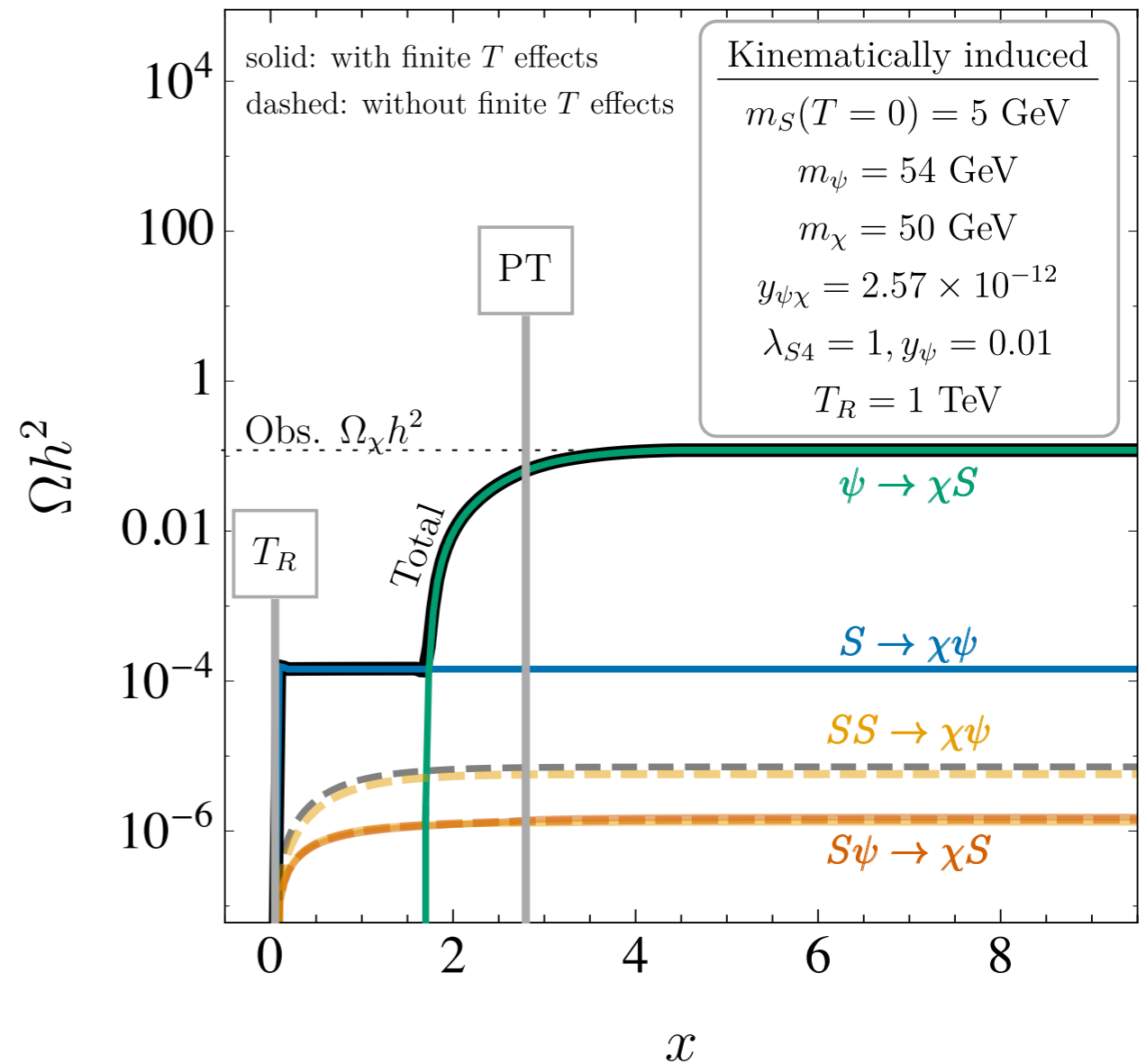
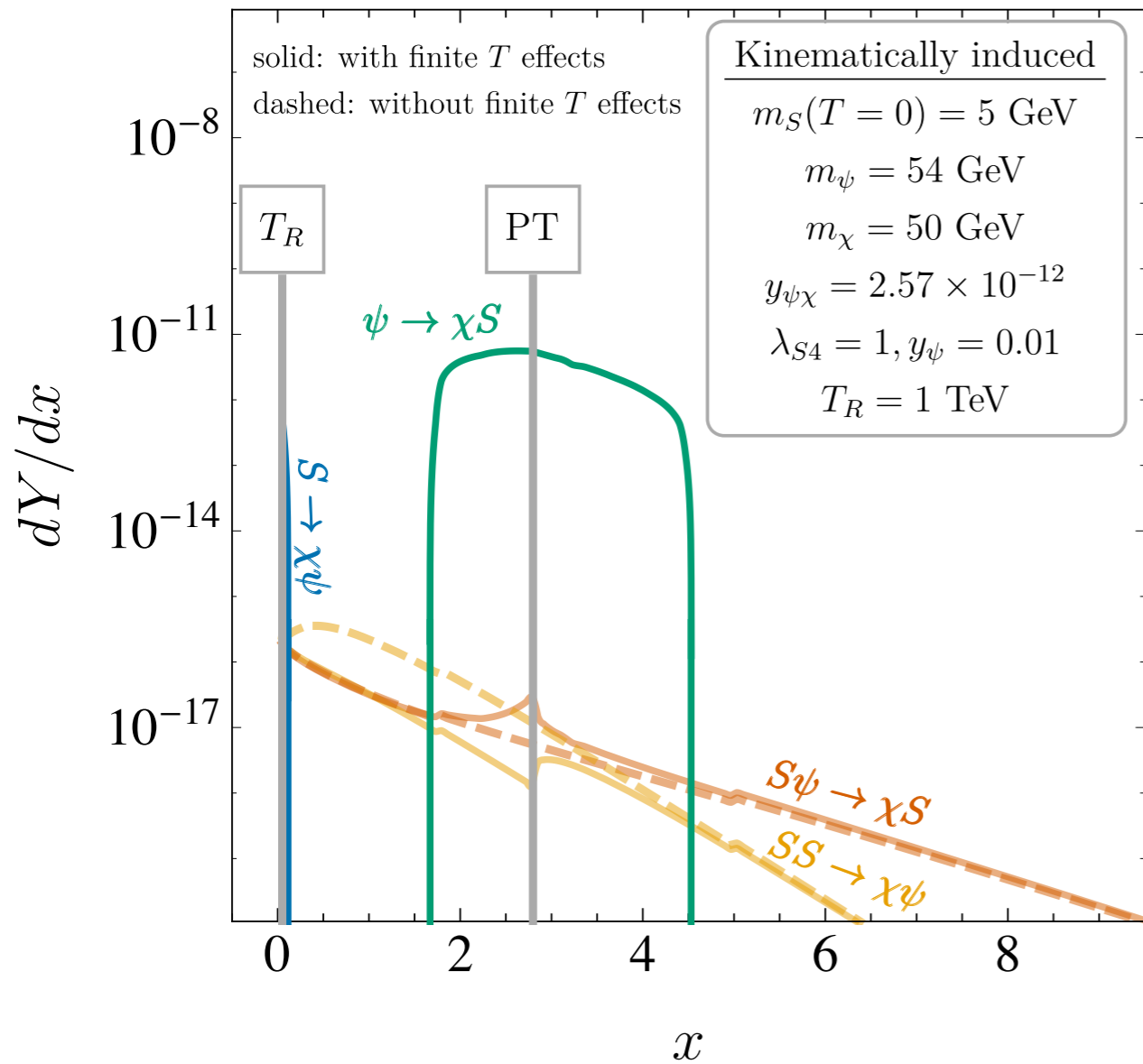
- Kinematically Induced Freeze-in - Results



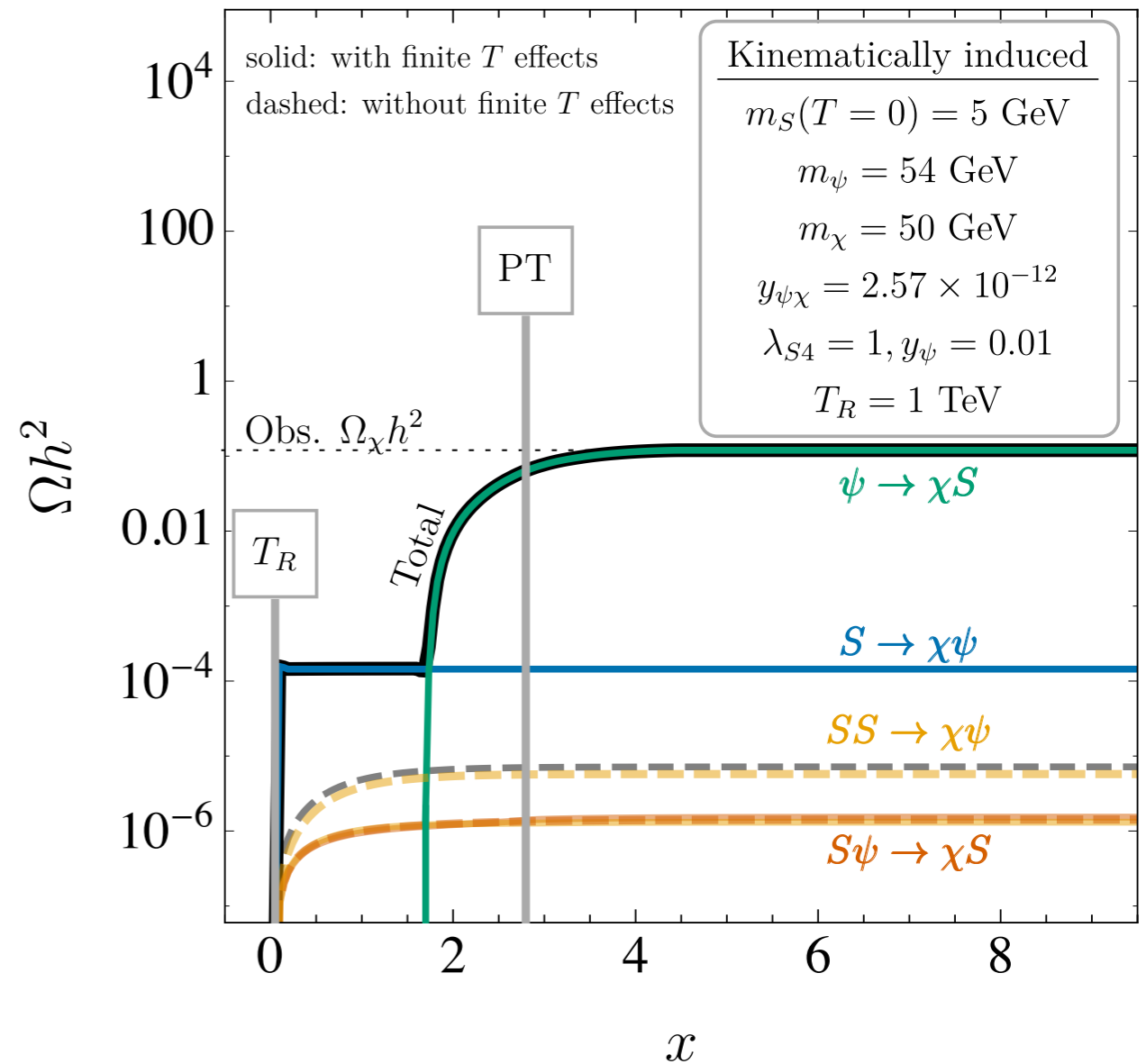
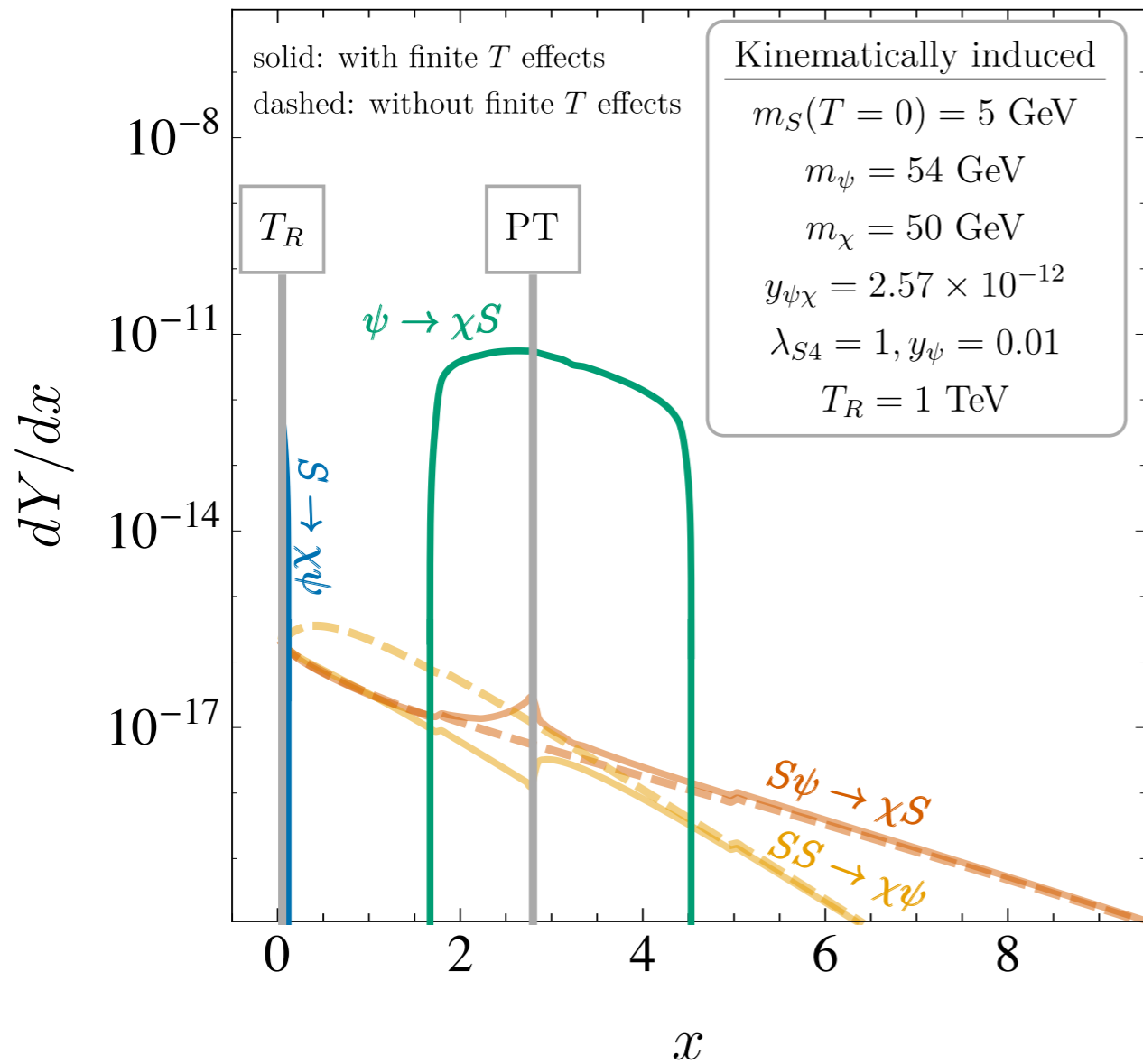
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Thermal effects increase relic abundance by orders of magnitude

Vev Induced Freeze-in

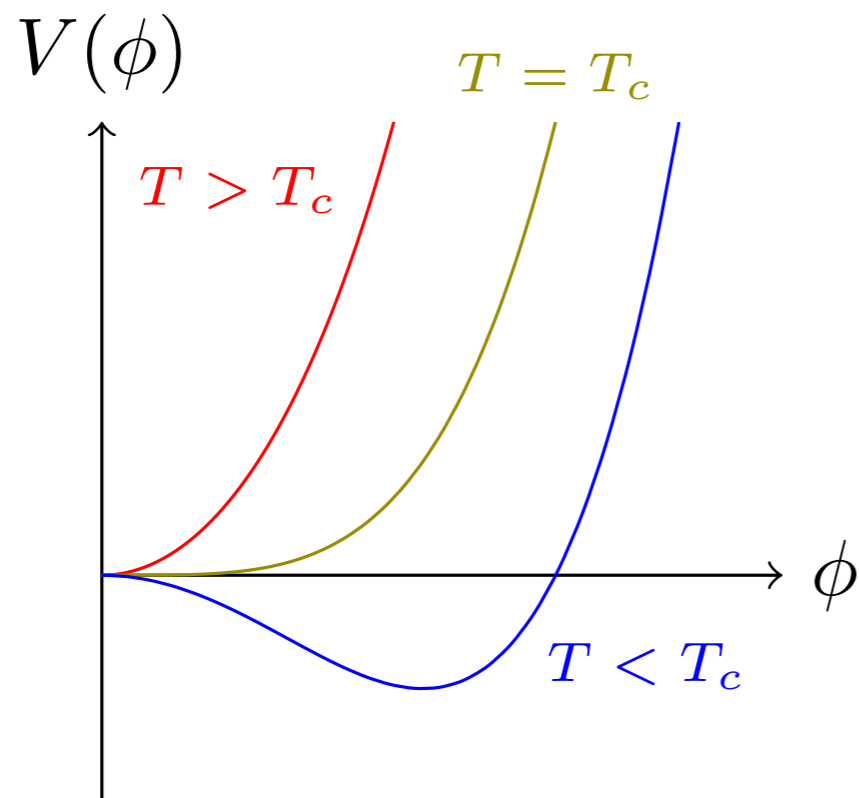
- Vev Flip-Flop: Phase Transitions
-

Brief (1 slide) aside on phase transitions...

- Vev Flip-Flop: Phase Transitions

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Second Order Phase Transitions



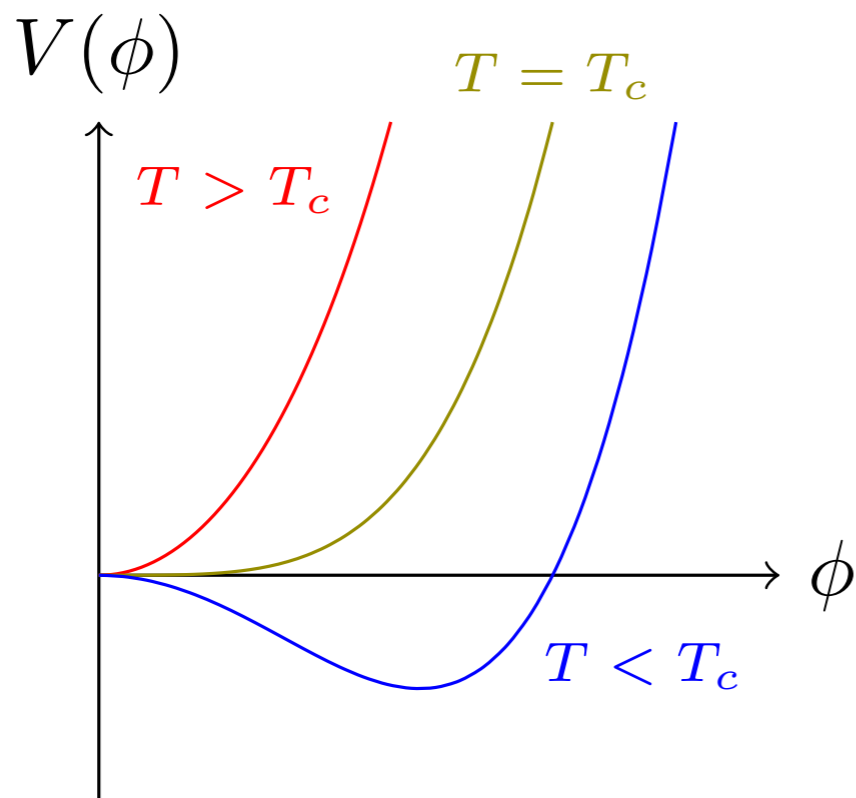
$$T_{\text{PT}} = T_c$$

- Vev Flip-Flop: Phase Transitions

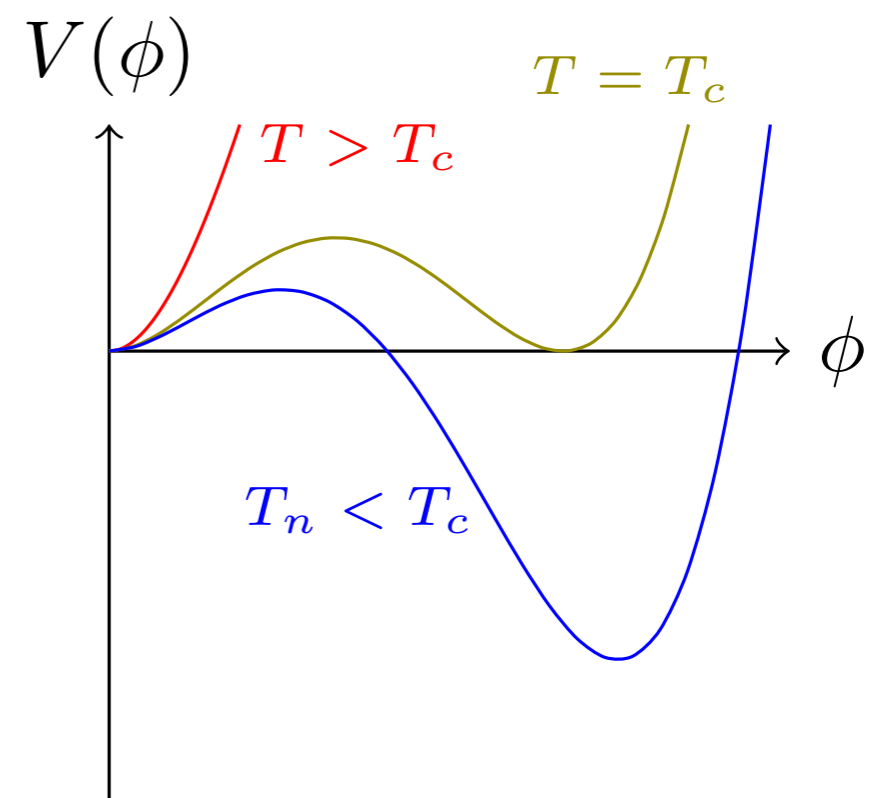
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Second Order Phase Transitions

First Order Phase Transitions



$$T_{\text{PT}} = T_c$$



$$T_n < T_c$$

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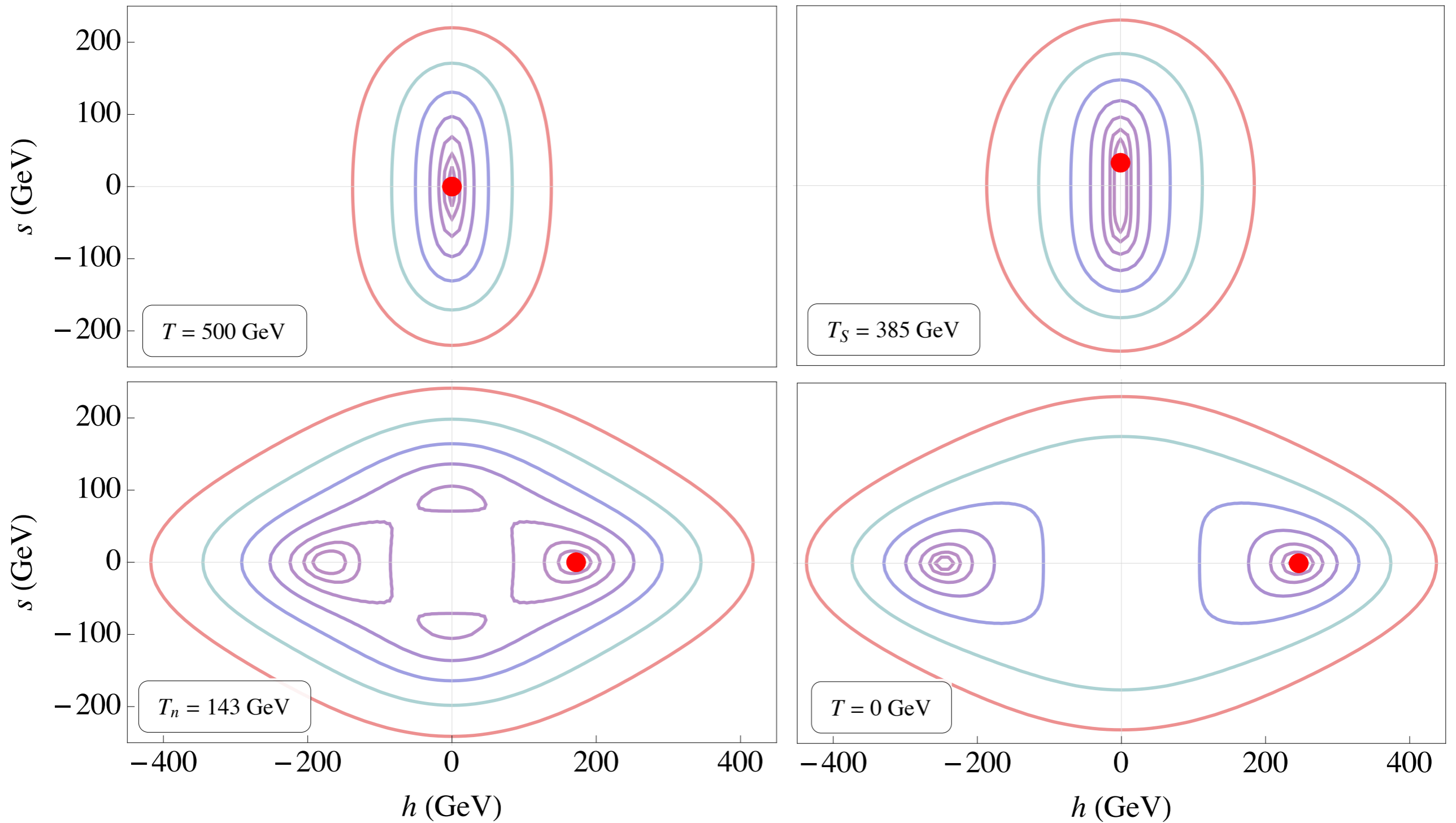
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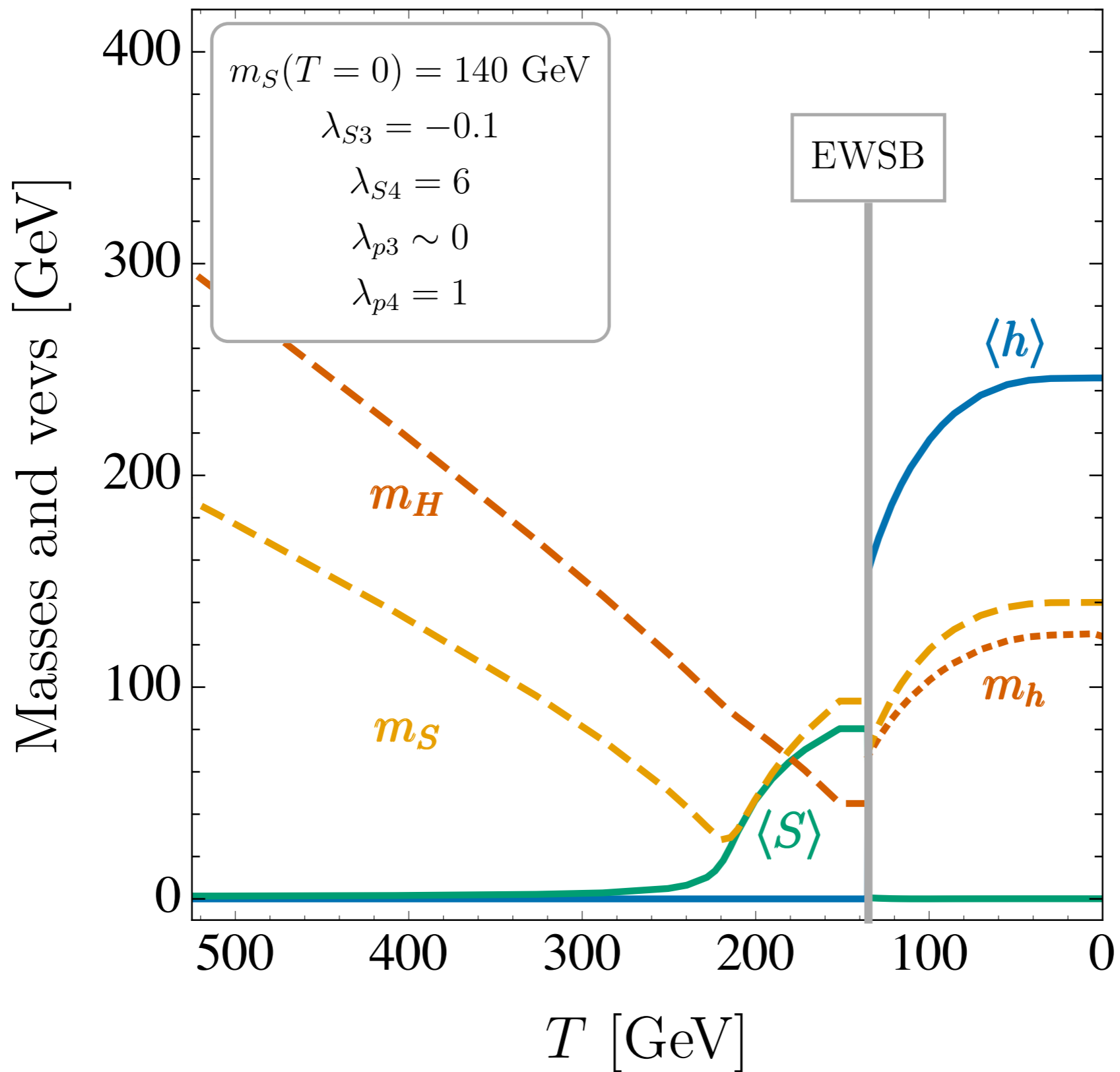
$$\lambda_{S3} \ll 1 \quad \lambda_{S4} \sim 1$$

$$\lambda_{p3} \sim 0 \quad \lambda_{p4} \sim 1$$

- Vev Flip-Flop: Effective Potential



- Vev Induced Freeze-in

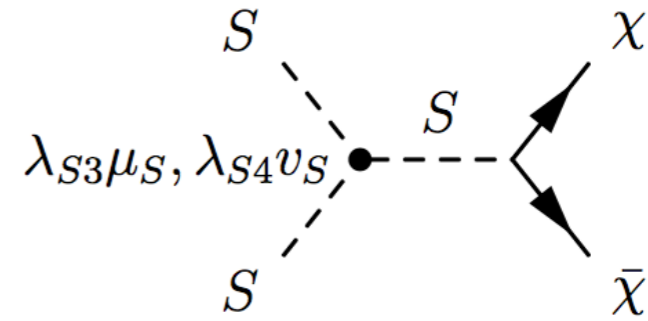
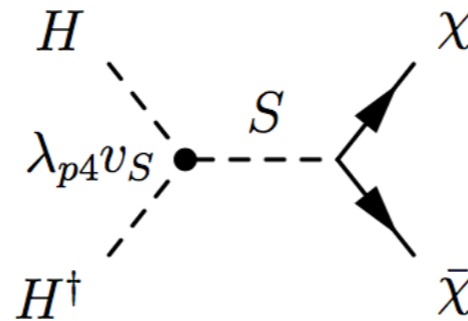
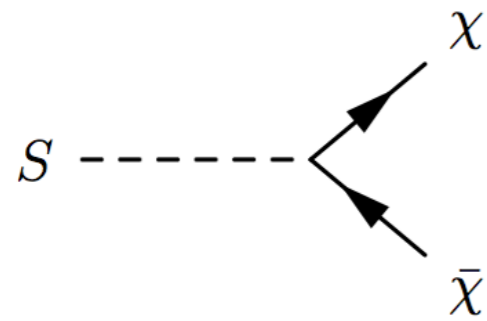


- Vev Induced Freeze-in

Again, assume χ begins with negligible abundance

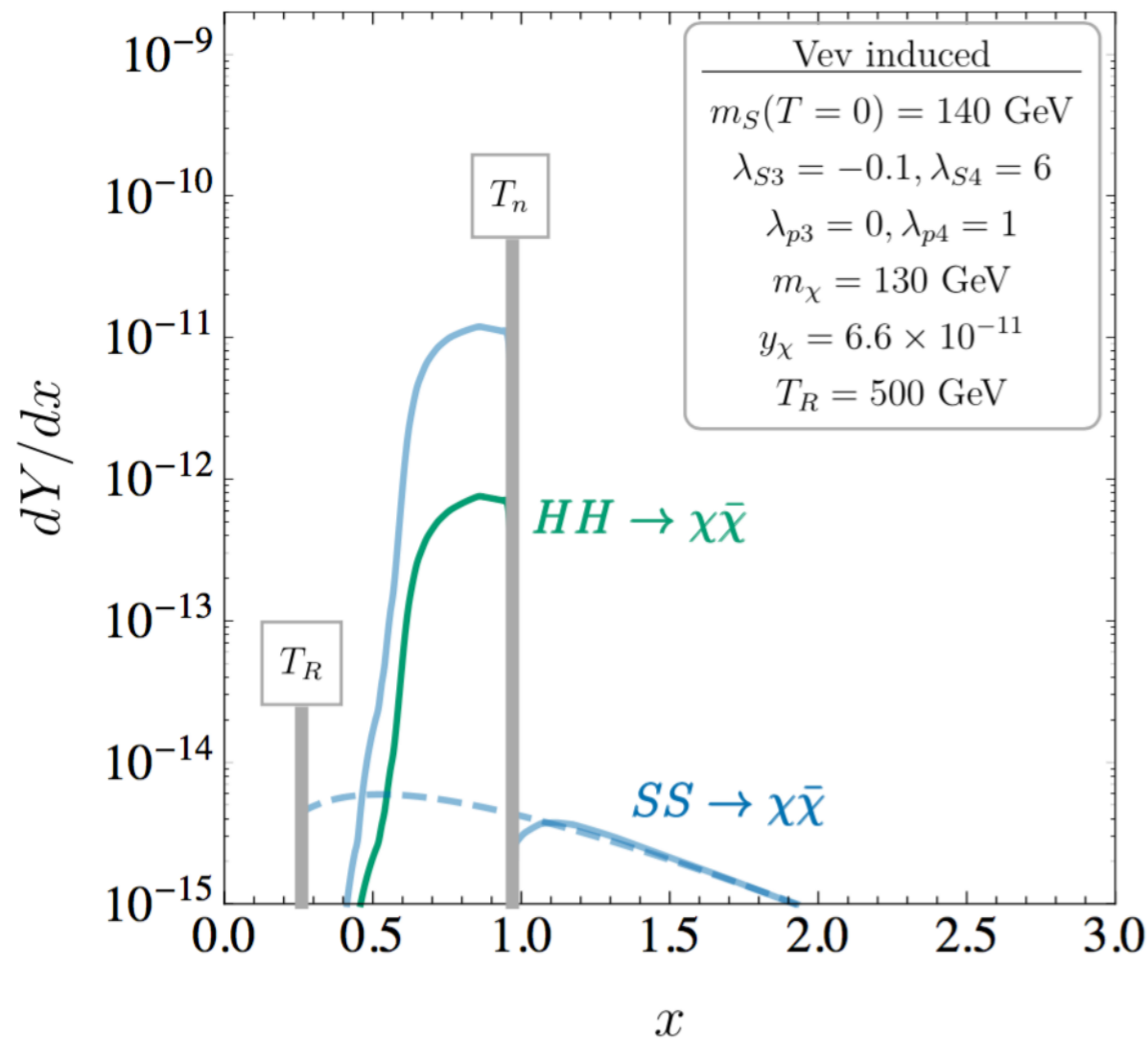
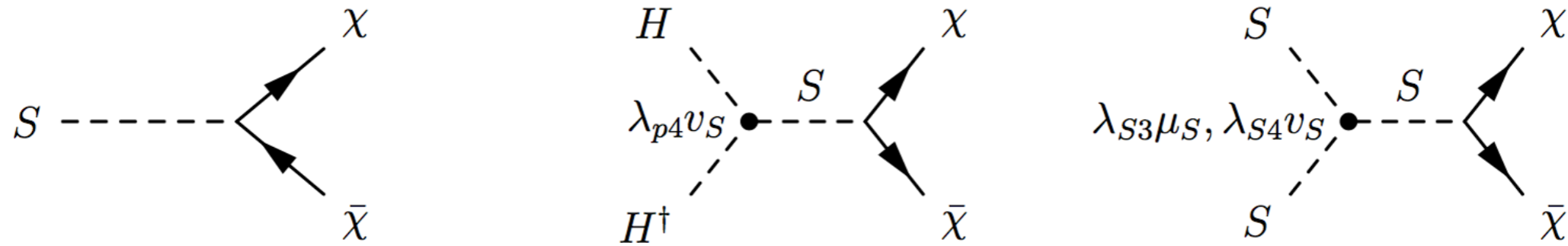
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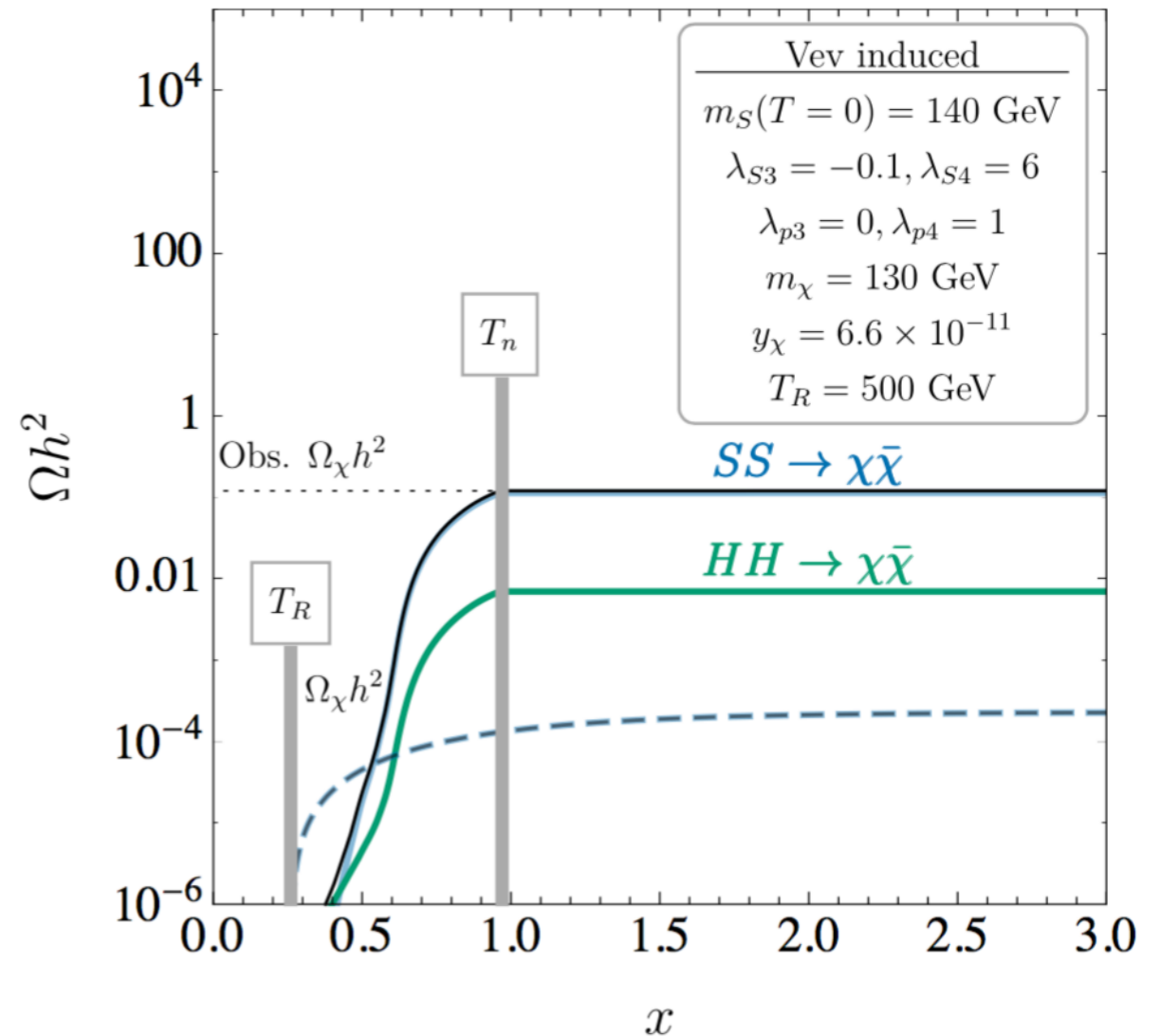
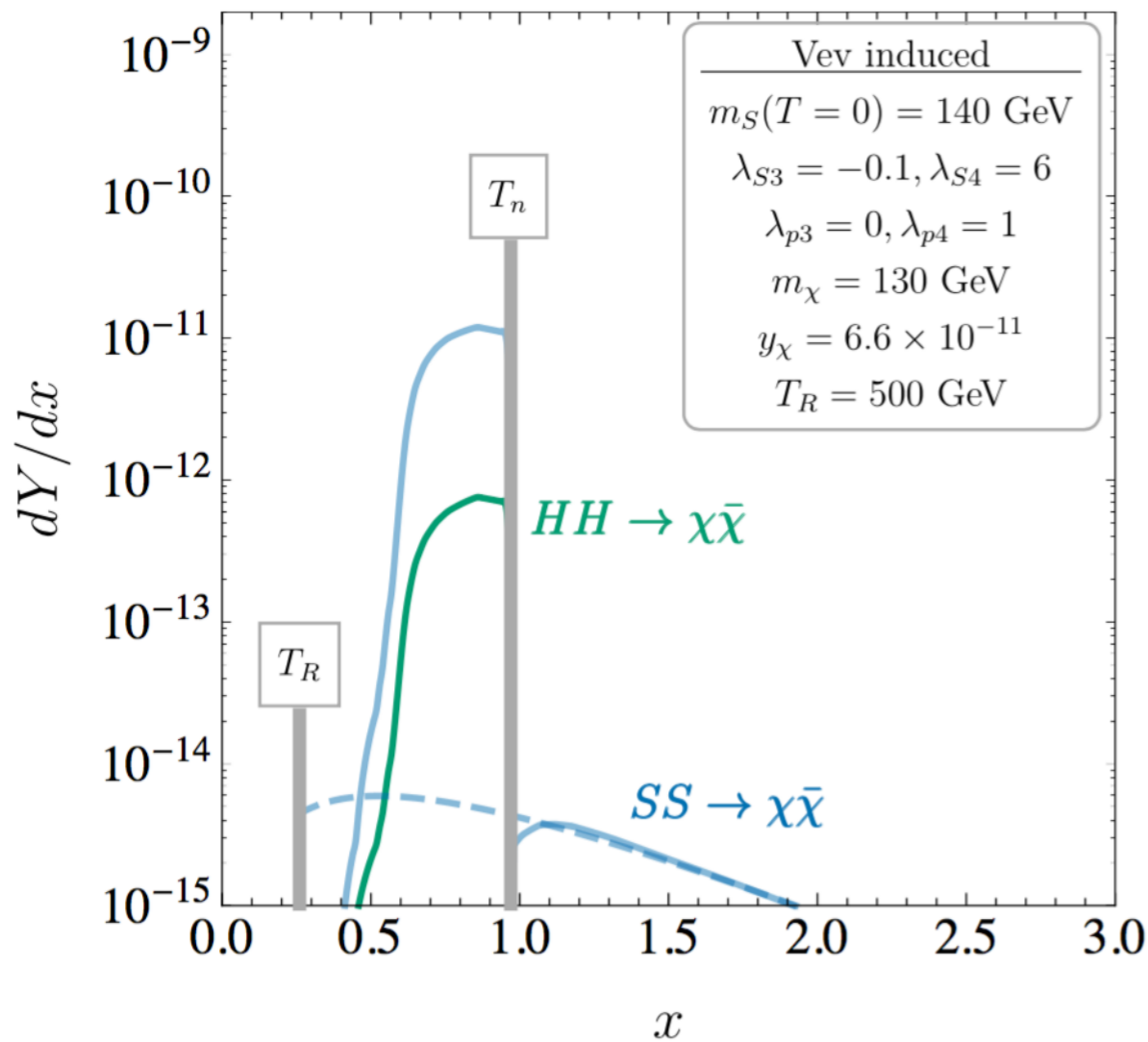
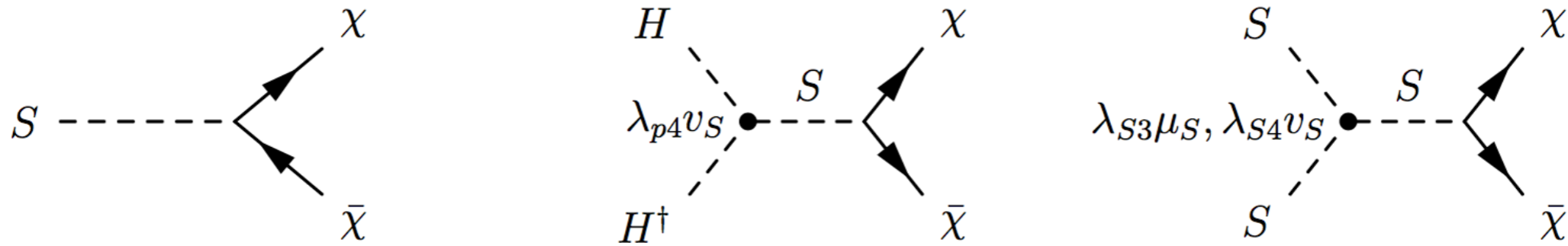
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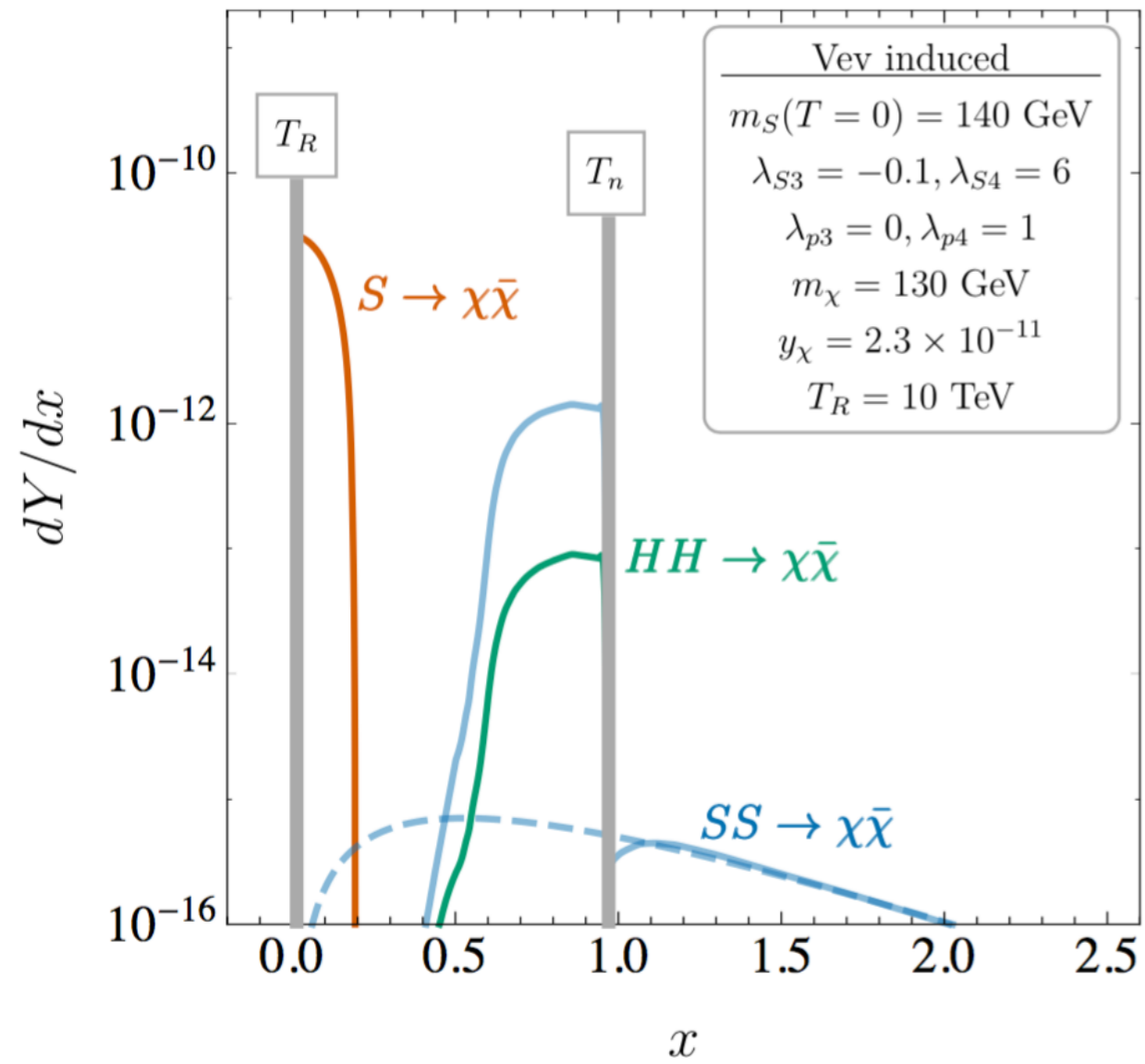


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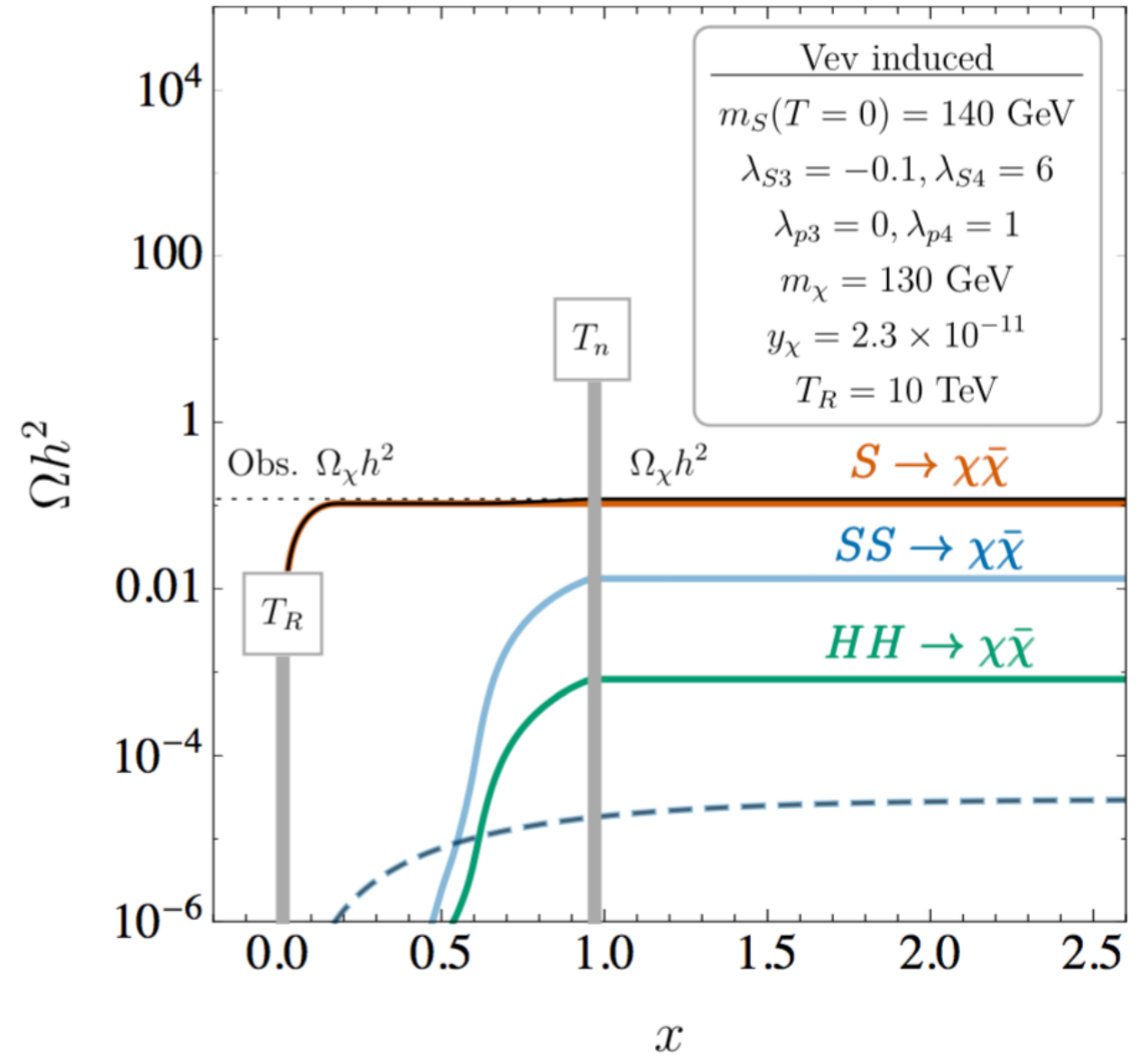
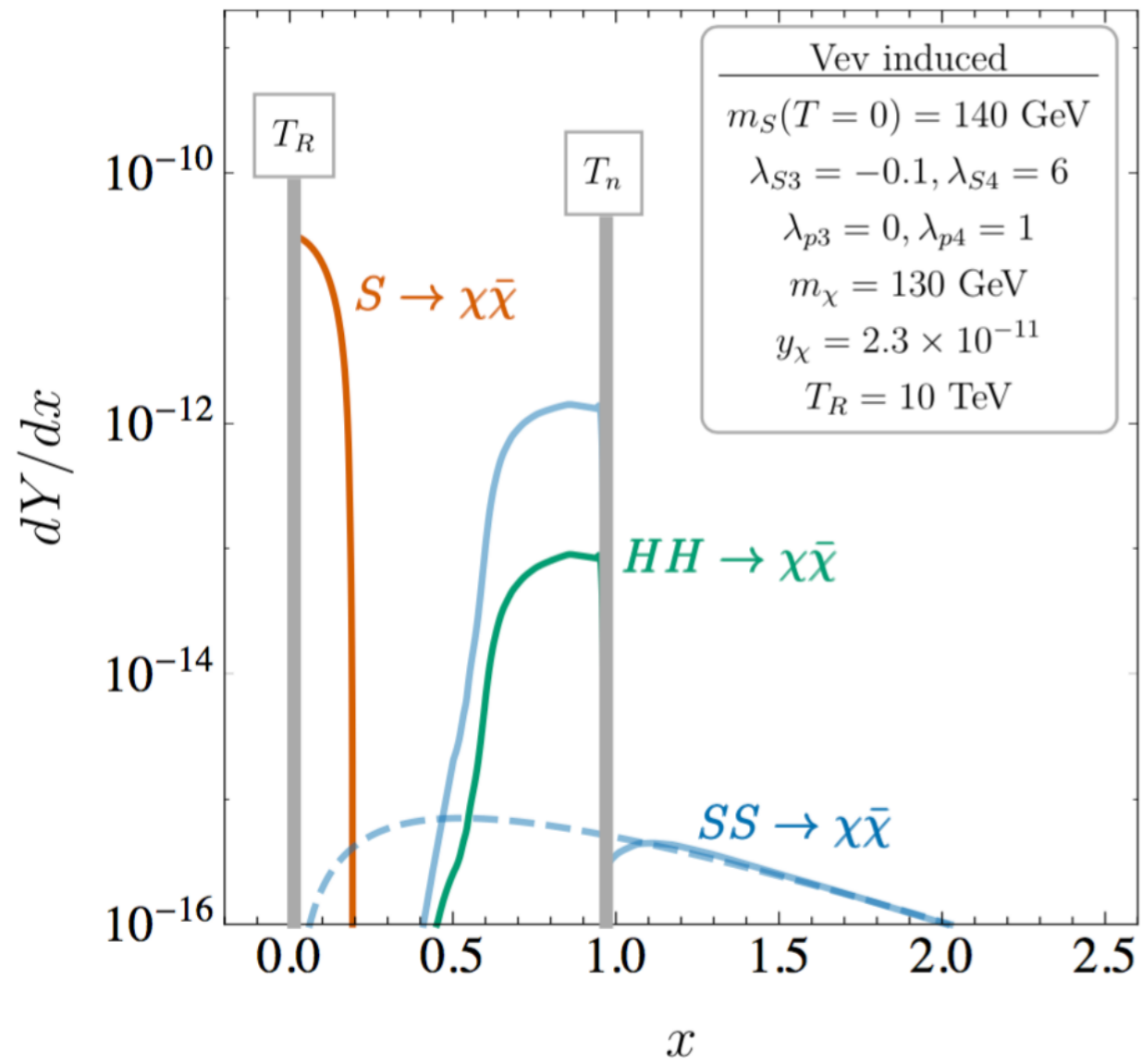
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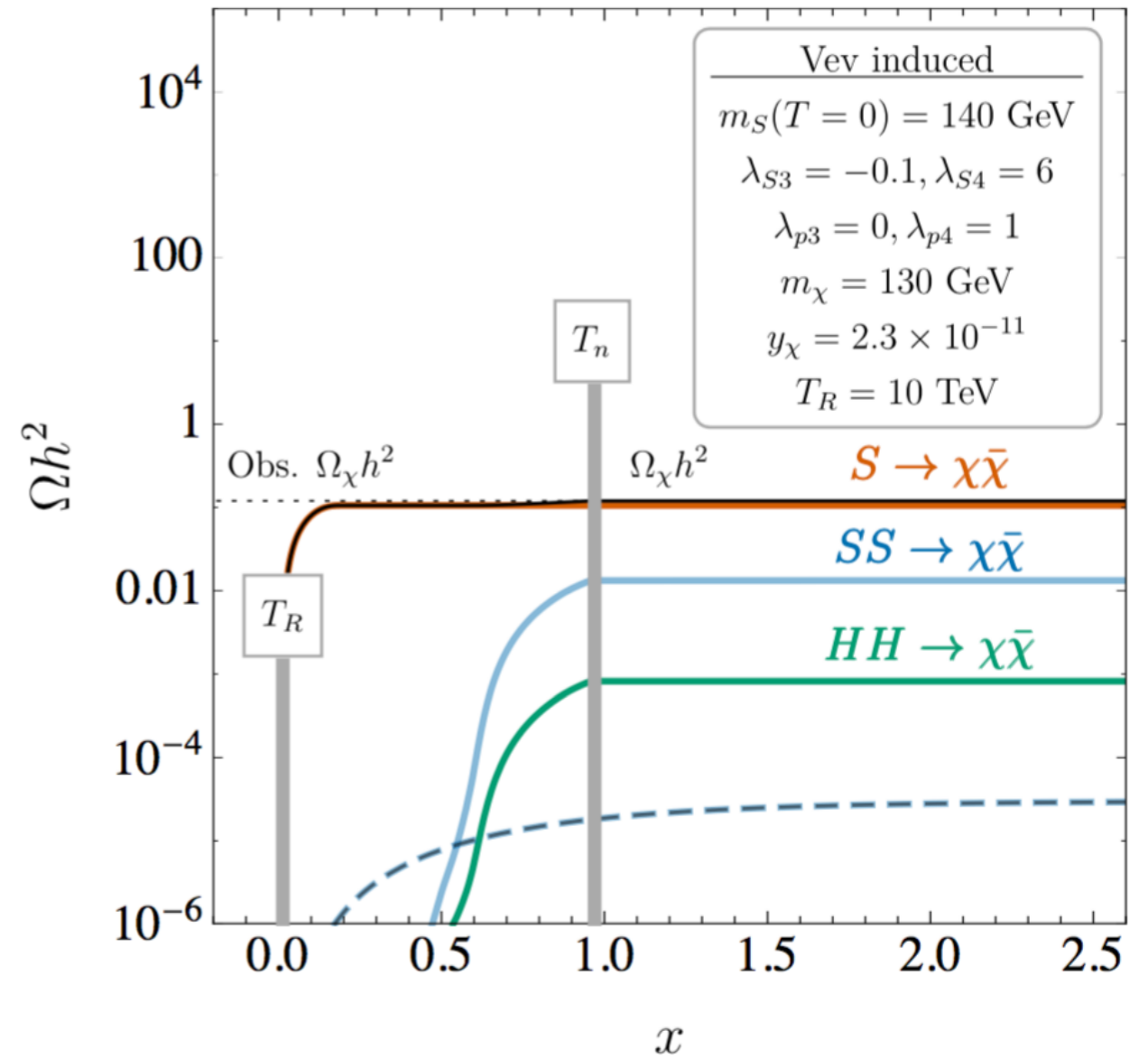
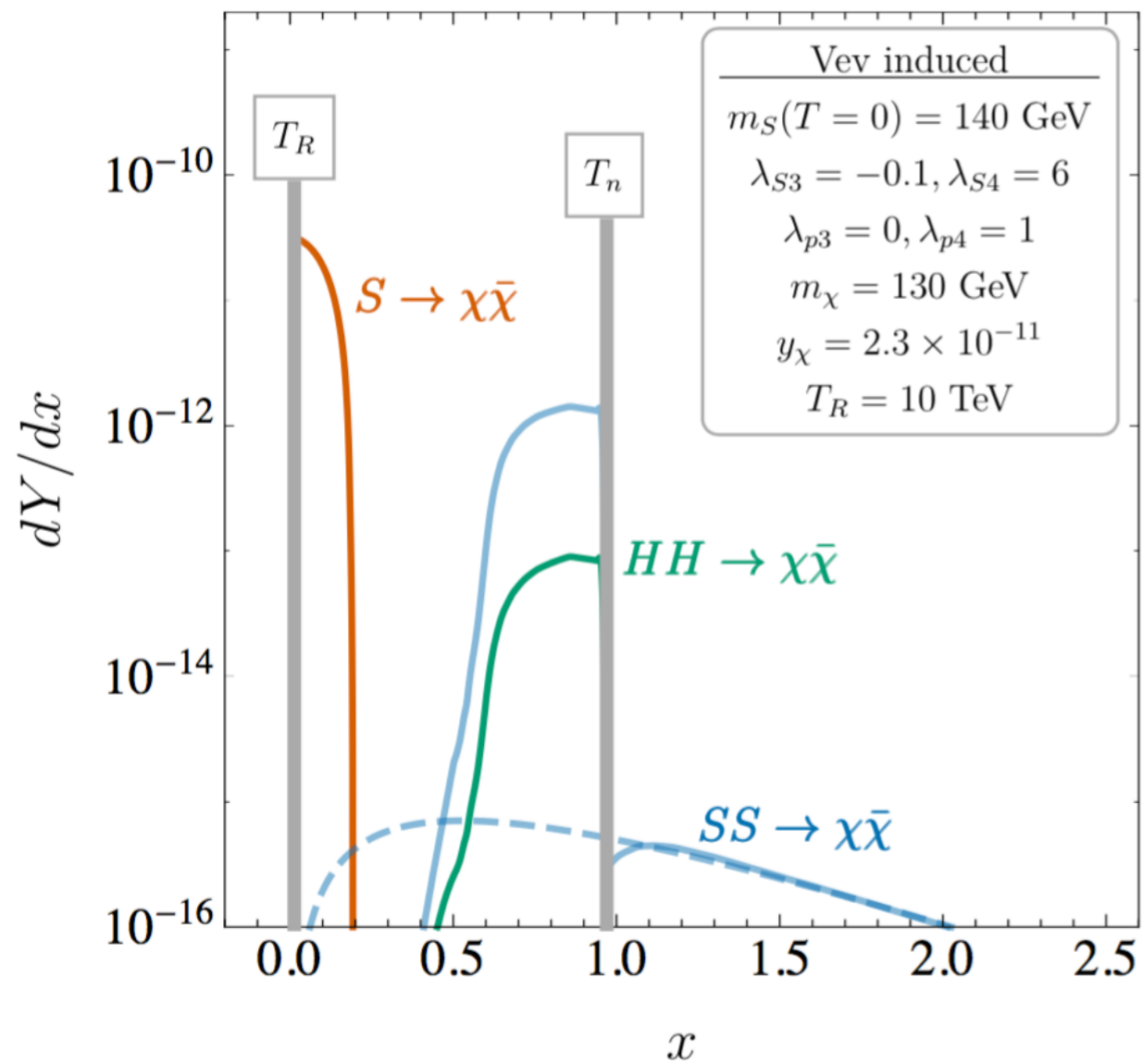
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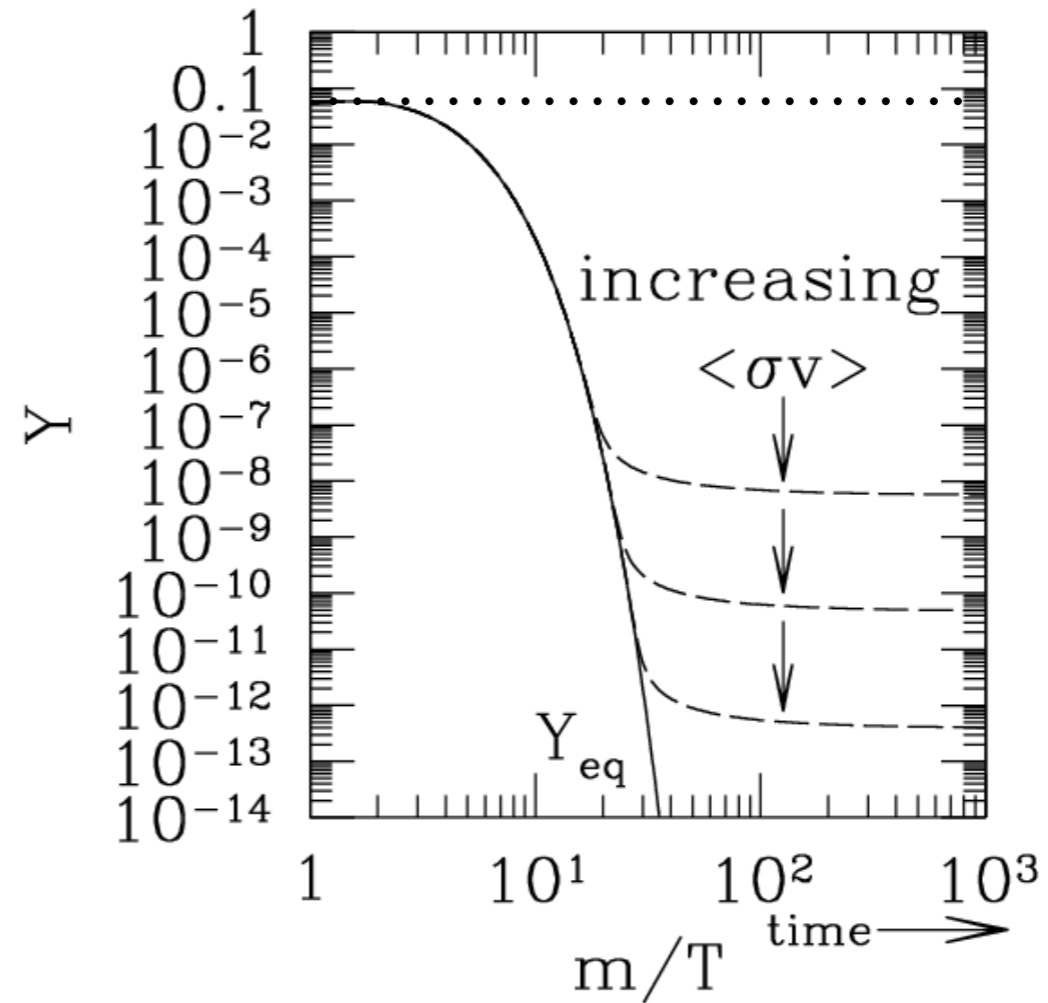
Decaying Dark Matter

- Vev Flip-Flop: Model

Field	Spin	SM	\mathbb{Z}_3	mass scale
χ	$\frac{1}{2}$	(1, 1, 0)	$\chi \rightarrow e^{2\pi i/3} \chi$	TeV
S	0	(1, 3, 0)	$S \rightarrow e^{2\pi i/3} S$	100 GeV
Ψ	$\frac{1}{2}$	(1, 3, 0)	$\Psi \rightarrow e^{-2\pi i/3} \Psi$	TeV
Ψ'	$\frac{1}{2}$	(1, 3, 0)	$\Psi' \rightarrow e^{-2\pi i/3} \Psi'$	TeV

- Vev Flip-Flop: Dark Matter Decay

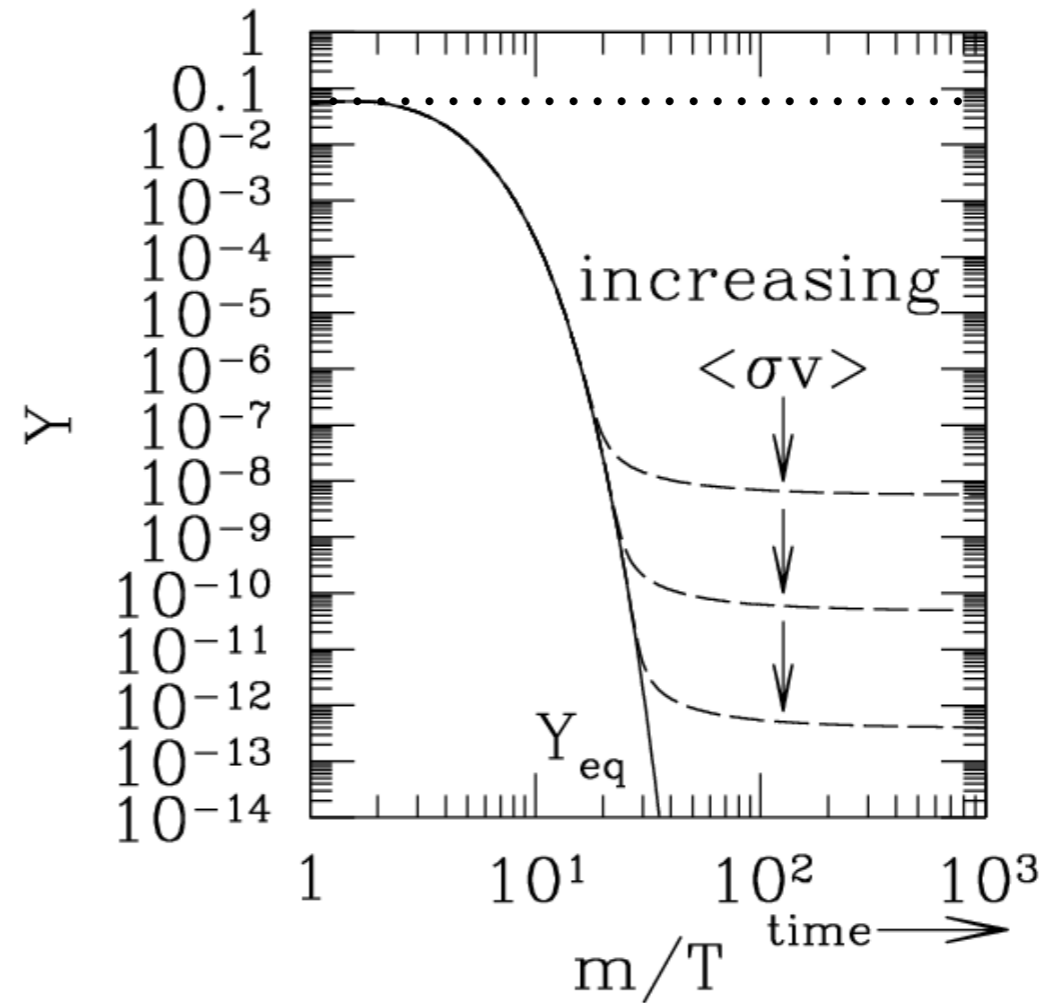
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$$\mathcal{L} \supset y_\chi S \bar{\chi} \Psi + y'_\chi S \bar{\chi} \Psi'$$

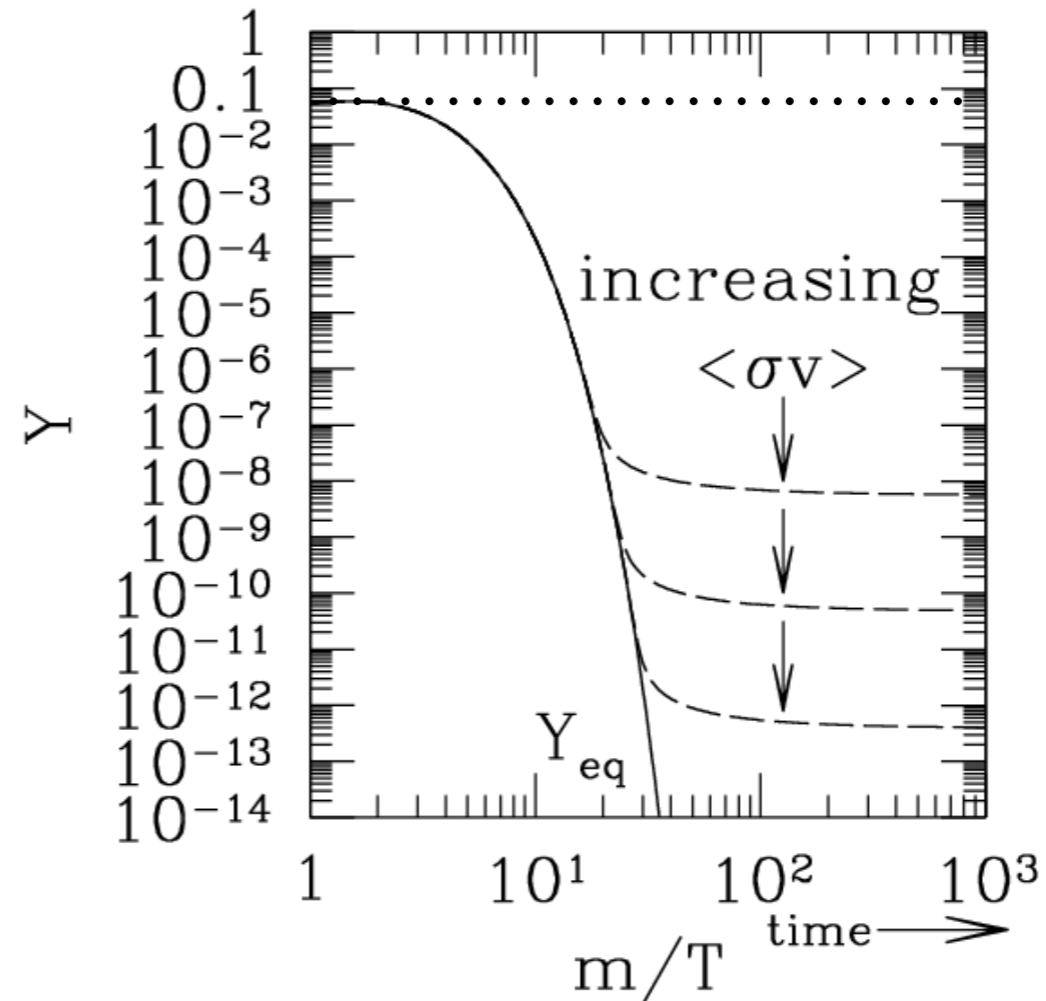


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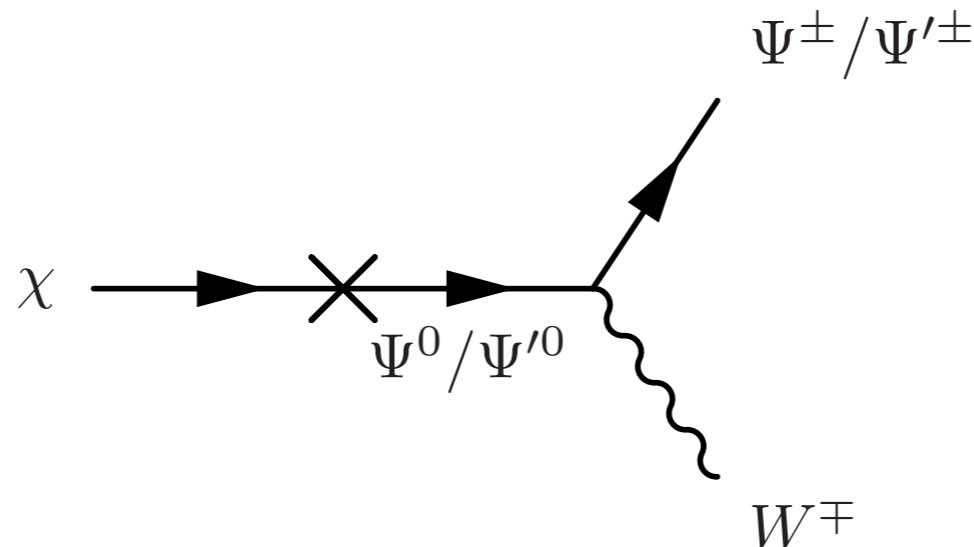
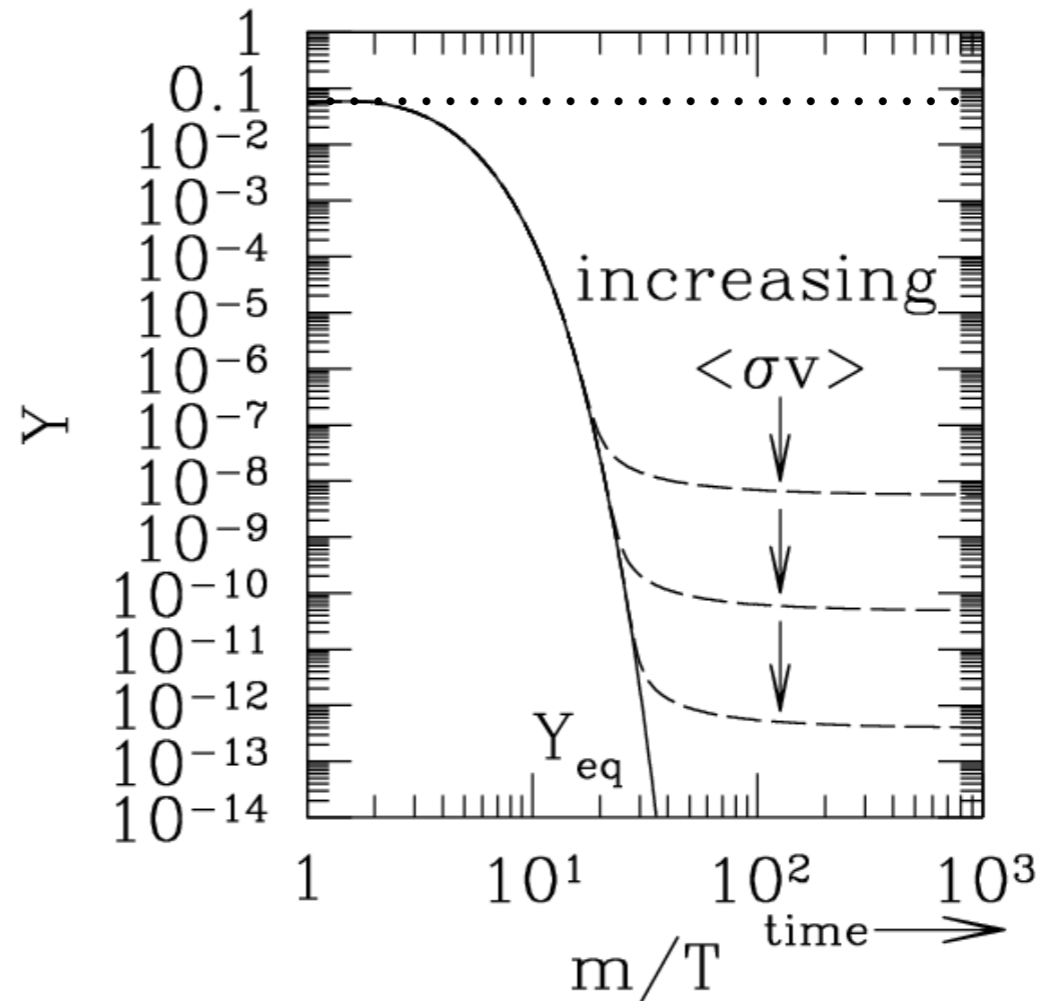


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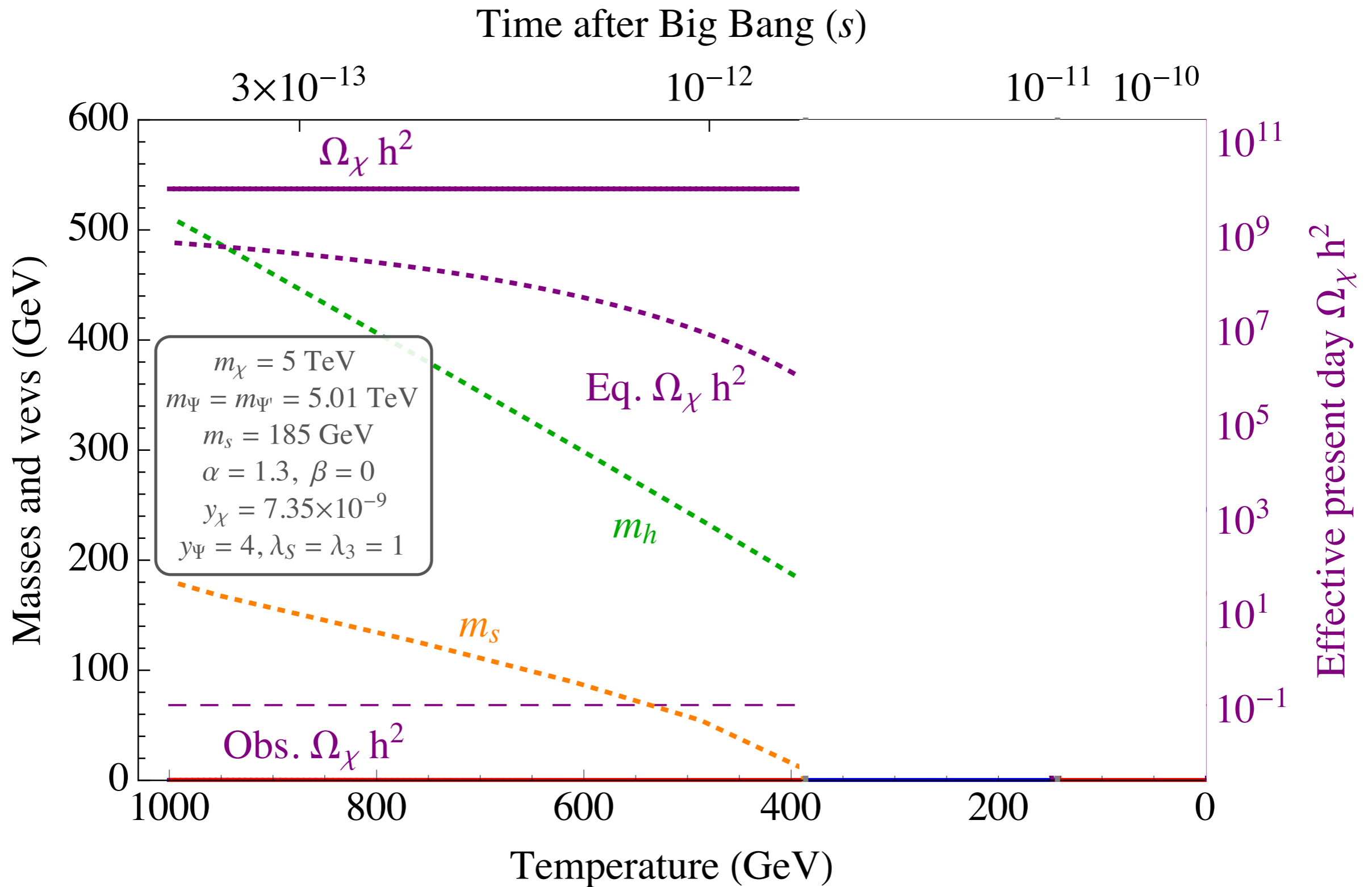
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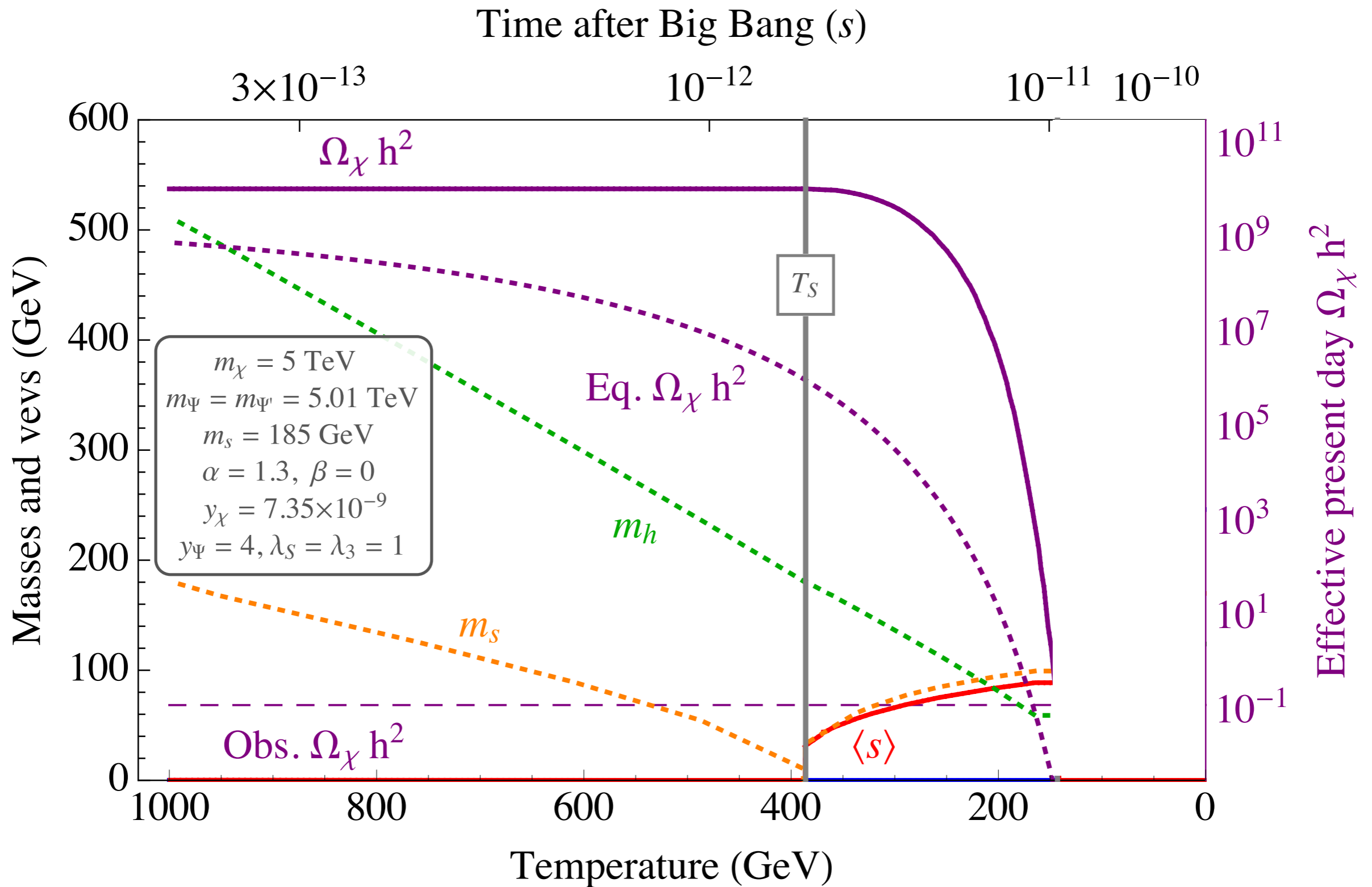
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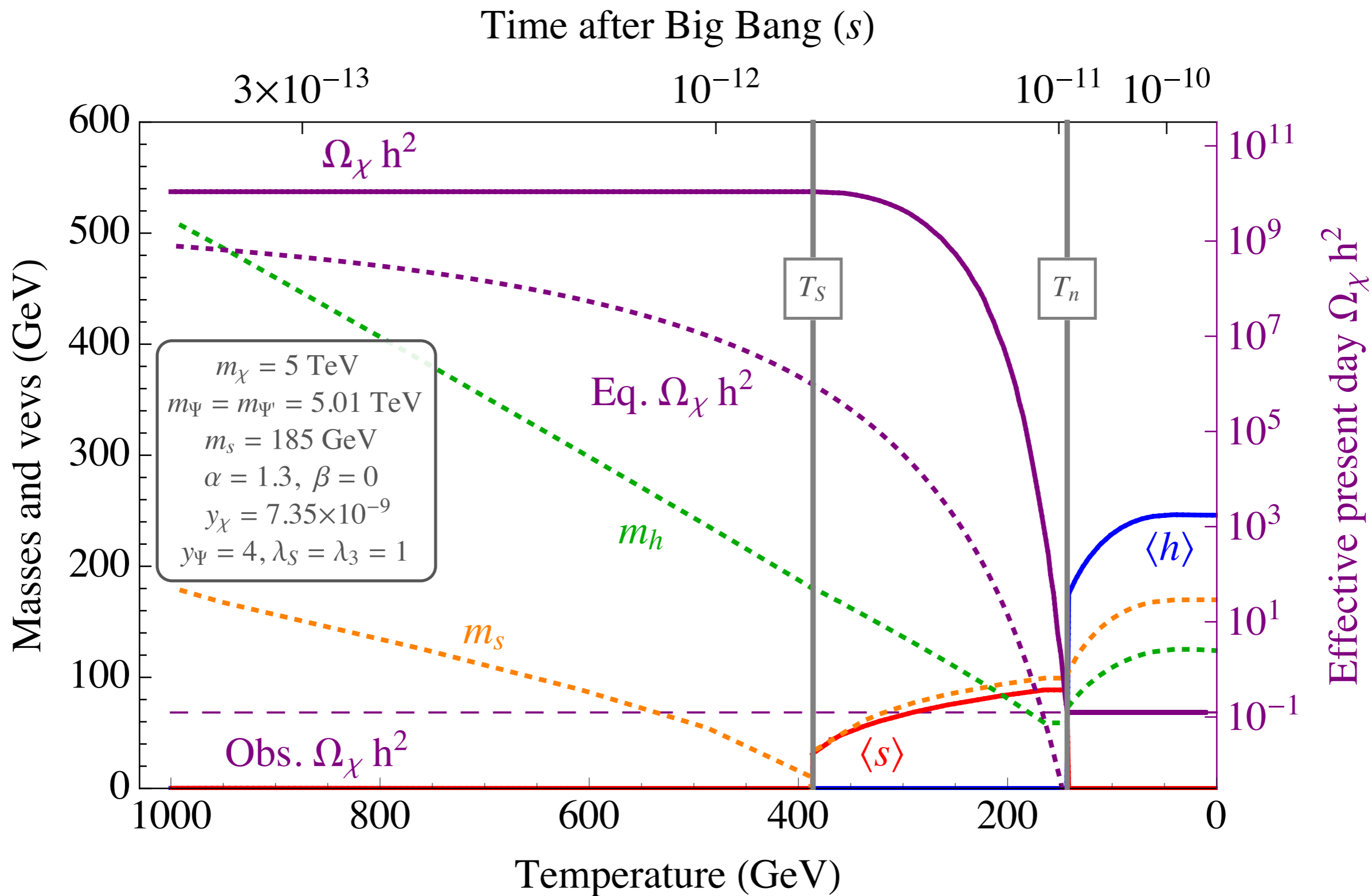
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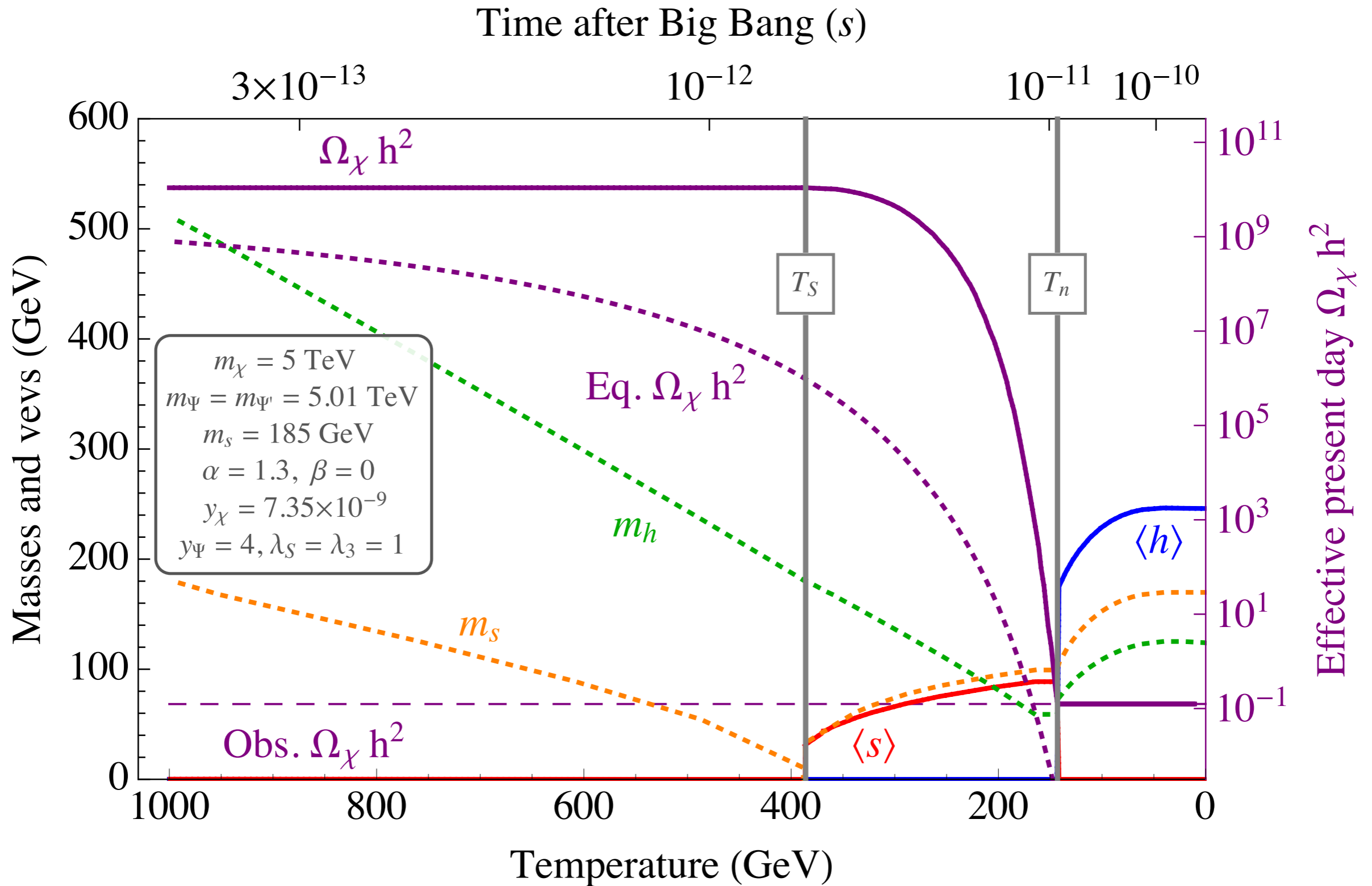
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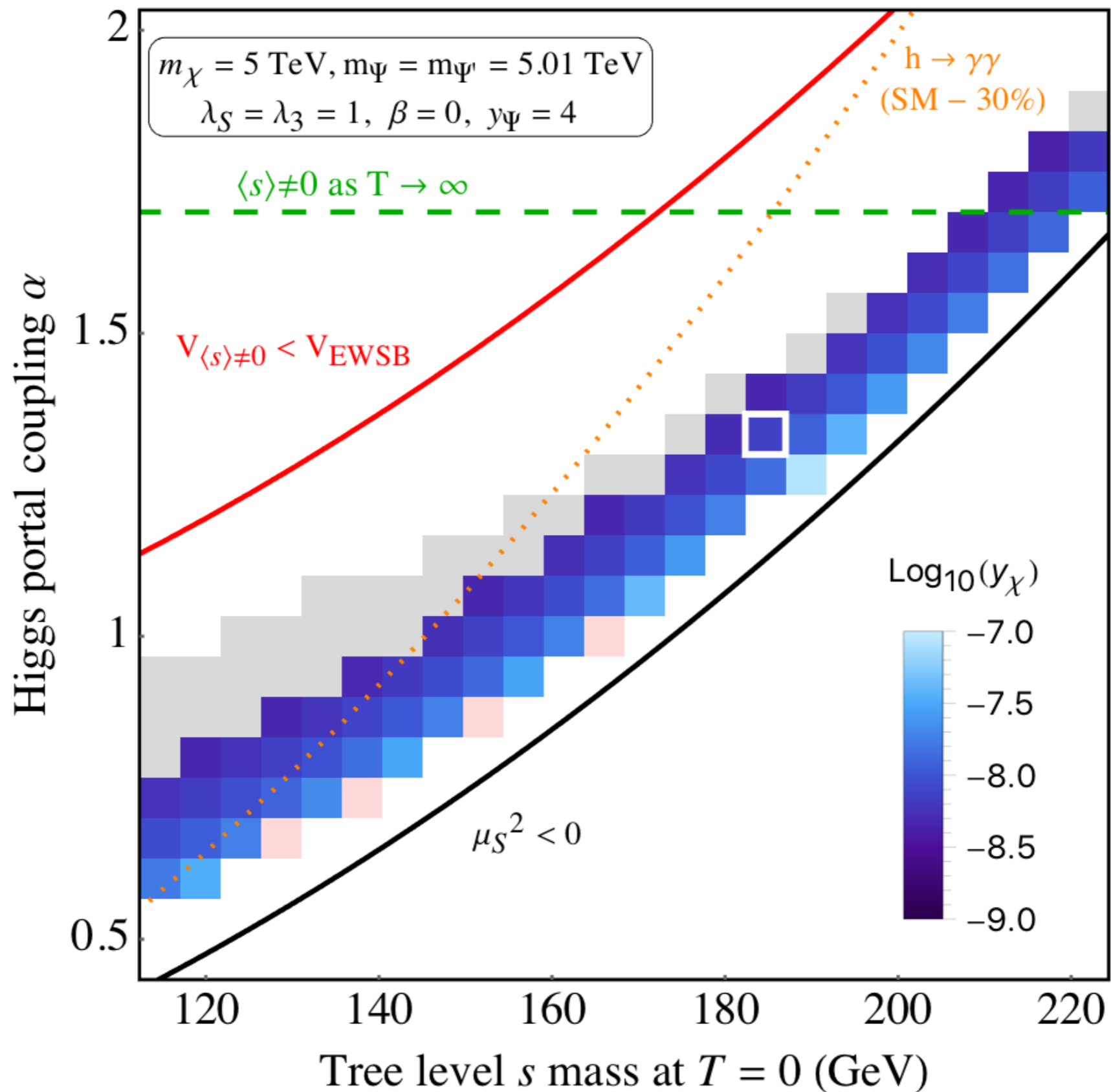


- Vev Flip-Flop: Dark Matter Decay



New production mechanism, determined by finite temperature effects

- Vev Flip-Flop: Parameter Space



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- Outlook:
 - More realisations,
 - LHC phenomenology,
 - gravitational wave signals,
 - baryogenesis,...

- Vev Flip-Flop: Possible tests
-

Direct and Indirect Detection

- Vev Flip-Flop: Possible tests
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Detection of χ hindered by small couplings

Best prospect from subdominant population of s, Ψ^0, Ψ'^0

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Small mass splitting implies soft decay products

Mono- χ is several orders of magnitude too weak

Possibilities in disappearing charged tracks