

BSM physics with neutron oscillations

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1. An experimentalist perspective
2. A theory perspective
3. A new experiment

An experimentalist's perspective

Hypothesis: baryon number is weakly violated. How do we look for it ?

Need processes in which only BNV takes place.

Single nucleon decay searches, eg, $p \rightarrow \pi^0 + e^+$?

$$\Rightarrow |\Delta B| = 1, |\Delta L| = 1 !$$

Decays without leptons, eg, $p \rightarrow \pi + \pi$, impossible due to angular momentum conservation.

$|\Delta B| \neq 0, \Delta L = 0$ observables restricted by Nature.

$n \rightarrow \bar{n}, n'$ and dinucleon decay searches sensitive to BNV -only.

Free $n \rightarrow \bar{n}, n'$ searches \Rightarrow cleanest experimental and theoretical approach.

Searches for BNV

Decay mode Partial mean life ($\times 10^{30}$ yrs)

$N \rightarrow e^+ \pi^-$	> 2000 (n), > 8200 (p)
$N \rightarrow \mu^+ \pi^-$	> 1000 (n), > 6600 (p)
$N \rightarrow \nu \pi$	> 1100 (n), > 390 (p)
$p \rightarrow e^+ \eta$	> 4200
$p \rightarrow \mu^+ \eta$	> 1300
$n \rightarrow \nu \eta$	> 158
$N \rightarrow e^+ \rho^-$	> 217 (n), > 710 (p)
$N \rightarrow \mu^+ \rho^-$	> 228 (n), > 160 (p)
$N \rightarrow \nu \rho$	> 19 (n), > 162 (p)
$p \rightarrow e^+ \omega$	> 320
$p \rightarrow \mu^+ \omega$	> 780
$n \rightarrow \nu \omega$	> 108
$N \rightarrow e^+ K^-$	> 17 (n), > 1000 (p)
$N \rightarrow \mu^+ K^-$	> 26 (n), > 1600 (p)
$N \rightarrow \nu K$	> 86 (n), > 5900 (p)
$n \rightarrow \nu K_S^0$	> 260
$p \rightarrow e^+ K^*(892)^0$	> 84
$N \rightarrow \nu K^*(892)$	> 78 (n), > 51 (p)
$p \rightarrow e^+ \pi^+ \pi^-$	> 82
$p \rightarrow e^+ \pi^0 \pi^0$	> 147
$n \rightarrow e^+ \pi^- \pi^0$	> 52
$p \rightarrow \mu^+ \pi^+ \pi^-$	> 133
$p \rightarrow \mu^+ \pi^0 \pi^0$	> 101
$n \rightarrow \mu^+ \pi^- \pi^0$	> 74
$n \rightarrow e^+ K^0 \pi^-$	> 18
$n \rightarrow e^- \pi^+$	> 65
$n \rightarrow \mu^- \pi^0$	> 49
$n \rightarrow e^- \mu^+$	> 52
$n \rightarrow \mu^- \mu^+$	> 7
$n \rightarrow e^- K^+$	> 32
$n \rightarrow \mu^- K^+$	> 57
$p \rightarrow e^- \pi^+ \pi^+$	> 30
$n \rightarrow e^- \pi^+ \pi^0$	> 29
$p \rightarrow \mu^- \pi^+ \pi^+$	> 17
$n \rightarrow \mu^- \pi^+ \pi^0$	> 34
$p \rightarrow e^- \pi^+ K^+$	> 75
$p \rightarrow \mu^- \pi^+ K^+$	> 245

$p \rightarrow e^+ \gamma$	> 670
$p \rightarrow \mu^+ \gamma$	> 478
$n \rightarrow \nu \gamma$	> 28
$p \rightarrow e^+ \gamma \gamma$	> 100
$n \rightarrow \nu \gamma \gamma$	> 219
$p \rightarrow e^+ e^+ e^-$	> 793
$p \rightarrow e^+ \mu^+ \mu^-$	> 359
$p \rightarrow e^+ \mu \nu$	> 170
$n \rightarrow e^+ e^- \nu$	> 257
$n \rightarrow \mu^+ e^- \nu$	> 83
$n \rightarrow \mu^+ \mu^- \nu$	> 79
$p \rightarrow \mu^+ e^+ e^-$	> 529
$p \rightarrow \mu^+ \mu^+ \mu^-$	> 675
$p \rightarrow \mu^- \nu \nu$	> 220
$p \rightarrow e^+ \mu^+ \mu^-$	> 6
$n \rightarrow 3\nu$	> 0.0005
$N \rightarrow e^+ \text{anything}$	> 0.6 (n, p)
$N \rightarrow \mu^+ \text{anything}$	> 12 (n, p)
$N \rightarrow e^- \pi^0 \text{anything}$	> 0.6 (n, p)
$p\bar{p} \rightarrow \pi^+ \pi^-$	> 0.7
$p\bar{n} \rightarrow \pi^+ \pi^0$	> 2
$n\bar{n} \rightarrow \pi^+ \pi^-$	> 0.7
$n\bar{n} \rightarrow \pi^0 \pi^0$	> 3.4
$p\bar{n} \rightarrow K^+ K^+$	> 170
$p\bar{p} \rightarrow e^+ e^-$	> 5.8
$p\bar{p} \rightarrow e^+ \mu^\pm$	> 3.6
$p\bar{p} \rightarrow \mu^+ \mu^+$	> 1.7
$p\bar{n} \rightarrow e^+ \bar{\nu}$	> 2.8
$p\bar{n} \rightarrow \mu^+ \bar{\nu}$	> 1.6
$p\bar{n} \rightarrow \tau^+ \bar{\nu}_\tau$	> 1.0
$n\bar{n} \rightarrow \nu_e \bar{\nu}_e$	> 1.4
$n\bar{n} \rightarrow \nu_\mu \bar{\nu}_\mu$	> 1.4

Few searches for BNV only.

$n \rightarrow \bar{n} \Rightarrow$ nuclear instability.

Highly suppressed due to
non-degeneracy of n, \bar{n}

$$\tau_{n \rightarrow \bar{n}} : 2.5 \times 10^8 \text{ s}$$

(Super-K, model-dependent)

$$\Delta B \neq 0, \Delta L \neq 0$$

$$\Delta B \neq 0, \Delta L = 0$$

Last search with free neutrons (ILL) $\tau_{n \rightarrow \bar{n}} > 0.8 \times 10^8 \text{ s}$

New search proposed for European Spallation Source.

Baryon and lepton number violation

- BN, LN “accidental” SM symmetries at perturbative level
 - BNV, LNV in SM non-perturbatively (eg instantons)
 - $B-L$ is conserved, not B, L separately.
- BNV, LNV needed for baryogenesis and leptogenesis
- BNV, LNV generic features of SM extensions (eg SUSY, GUTs..)
- Need to explore the possible selection rules:

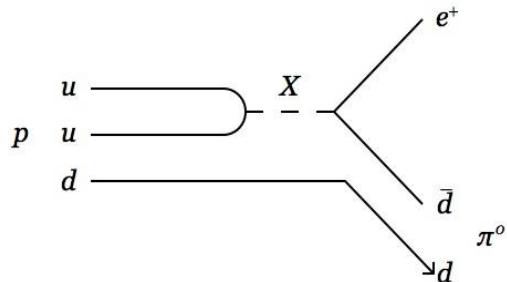
$$\Delta B \neq 0, \Delta L = 0, \Delta[B - L] \neq 0$$

$$\Delta B = 0, \Delta L \neq 0, \Delta[B - L] \neq 0$$

$$\Delta L \neq 0, \Delta B \neq 0, \Delta[B - L] = 0$$

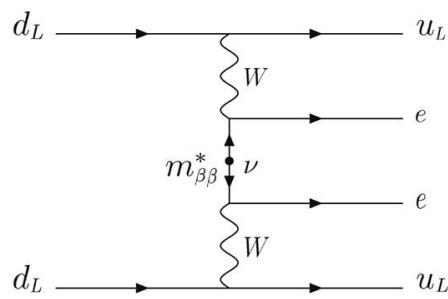
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Complementary B,L -violation observables



$$p \rightarrow e^+ + \pi^0$$

$$\Delta B = 1, \Delta L = 1$$

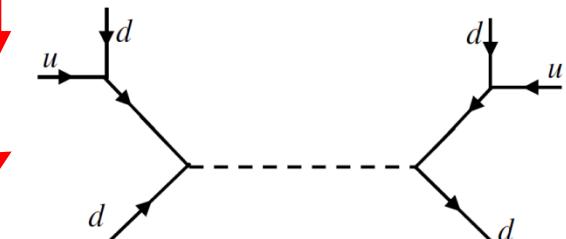


$$0\nu\beta\beta$$

$$\Delta B = 0, \Delta L = 2$$

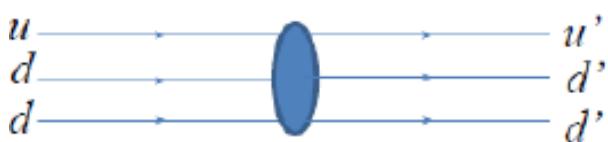
Symbiosis

Neutron
oscillation



$$n \rightarrow \bar{n}, NN \text{ decay.}$$

$$\Delta B = 2, \Delta L = 0$$



$$n \rightarrow n' \text{ (mirror).}$$

$$\Delta B = 1, \Delta L = 0$$

Neutron-antineutron oscillations

- R -parity violating supersymmetry, minimal flavour violation SUSY
- Unification models: $M \sim 10^{15}$ GeV
- Left-right symmetric models ($n\bar{n}$ and $0\nu 2\beta$)
- Extra dimensions models
- Post-sphaleron baryogenesis
- Bottom-up EFT - baryogenesis (C. Grojean, B. Shakya, J. Wells, Z. Zhang ([arxiv: 1806.00011 -hep-ph](#))

High precision $n \rightarrow \bar{n}$ search

⇒ Scan over wide range of phase space for generic BNV

+

⇒ model constraints.

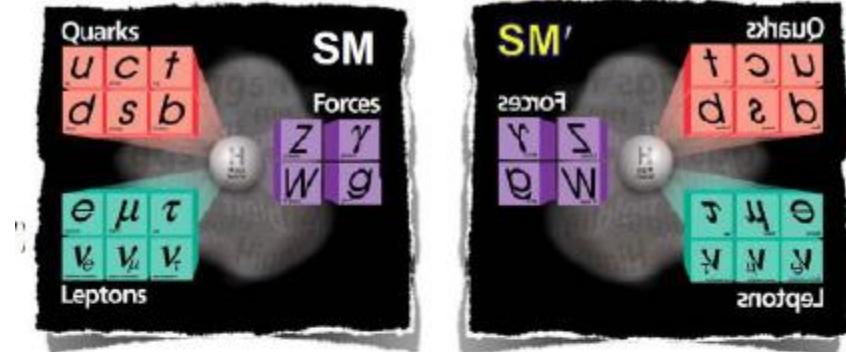
Mirror neutrons

Hidden/mirror sector

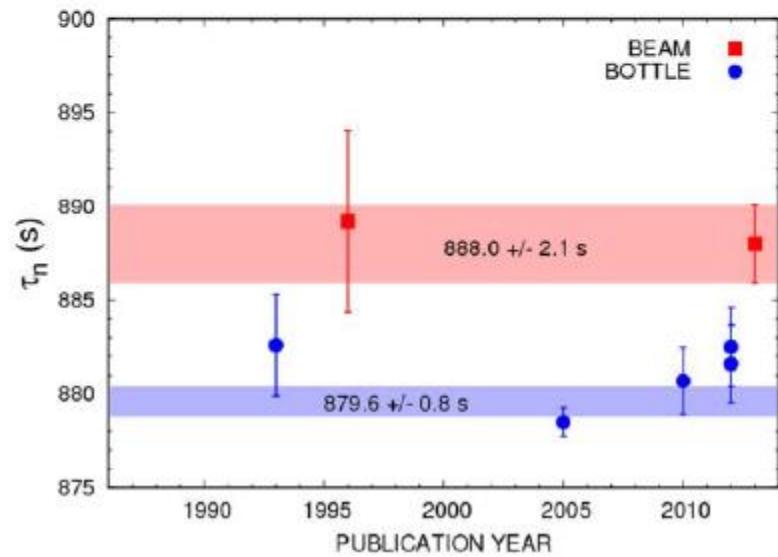
Restores parity symmetry

Mixing for $Q=0$ particles, eg $n \rightarrow n'$

Mirror matter: dark matter candidates



Explain neutron lifetime discrepancy
in bottle and beam experiments.



Operator analysis

Six quark operators O_i of dimension-9 :

$$\begin{aligned}
 (u_R d_R d_R)^2 &\equiv \epsilon_{abc} u_{R\dot{\alpha}}^a d_R^{\dot{\alpha}b} d_{R\dot{\gamma}}^c \epsilon_{def} u_{R\dot{\beta}}^d d_R^{\dot{\beta}e} d_R^{\dot{\gamma}f} \\
 (u_R d_R d_L)^2 &\equiv \epsilon_{abc} u_{R\dot{\alpha}}^a d_R^{\dot{\alpha}b} d_L^{\dot{\gamma}c} \epsilon_{def} u_{R\dot{\beta}}^d d_R^{\dot{\beta}e} d_{L\dot{\gamma}}^f \\
 (u_L d_L d_R)^2 &\equiv \epsilon_{abc} u_L^{\alpha a} d_{L\alpha}^b d_{R\dot{\gamma}}^c \epsilon_{def} u_L^{\beta d} d_{L\beta}^e d_R^{\dot{\gamma}f} \\
 (u_R d_R s_R)^2 &\equiv \epsilon_{abc} u_{R\dot{\alpha}}^a d_R^{\dot{\alpha}b} s_{R\dot{\gamma}}^c \epsilon_{def} u_{R\dot{\beta}}^d d_R^{\dot{\beta}e} s_R^{\dot{\gamma}f}.
 \end{aligned}$$

↑
 $n \rightarrow \bar{n}, NN \rightarrow \pi\pi$
↓
 $NN \rightarrow KK$

Eg $n \rightarrow \bar{n}$: $\langle \bar{n} | H_{eff} | n \rangle = \frac{1}{M_X^5} \sum_i \kappa_i \langle \bar{n} | O_i | n \rangle$

Short distance (RPV SUSY): κ_i

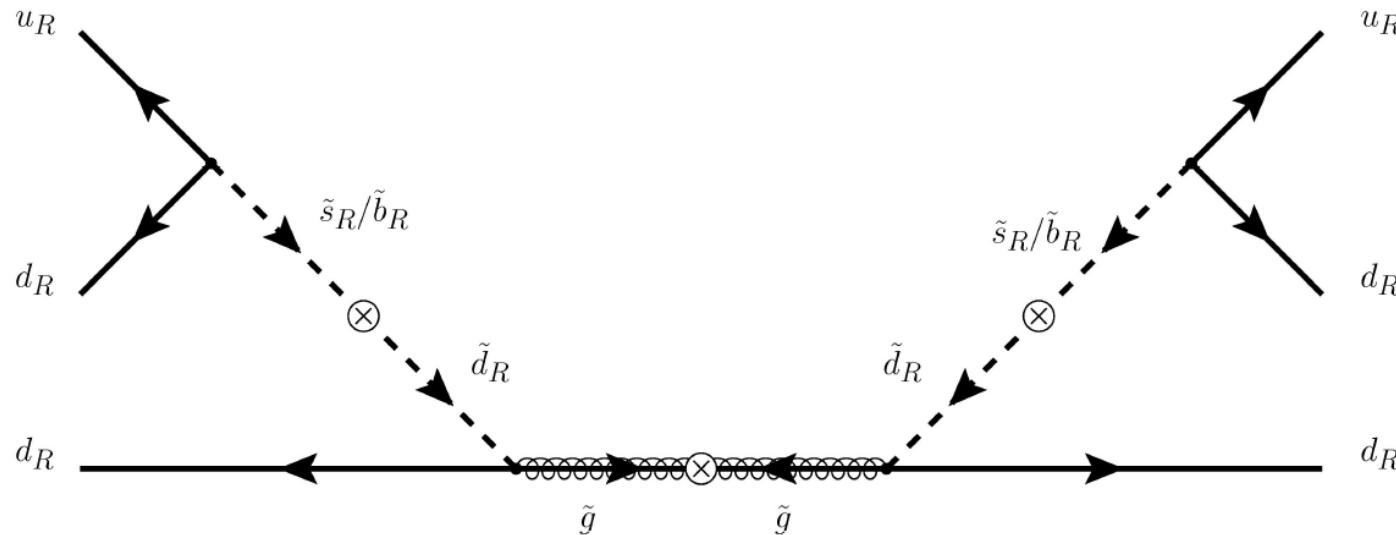
Long distance Hadronic ME: $\langle \bar{n} | O_i | n \rangle$: Λ_{QCD}^6

Oscillation time $\tau = \frac{1}{\langle \bar{n} | H_{eff} | n \rangle} \cdot \frac{M_X^5}{\kappa \Lambda_{QCD}^6}$

BNV in RPV-SUSY scenarios

$$W_{RPV} = \frac{1}{2} \lambda''_{ijk} \bar{U}_i \bar{D}_j \bar{D}_k ; \quad \lambda''_{ijk} = -\lambda''_{ikj} \Rightarrow \lambda''_{111} = 0$$

1st gen. quarks \Leftrightarrow 2nd/3rd gen. squarks $(\lambda''_{112}, \lambda''_{113})$.

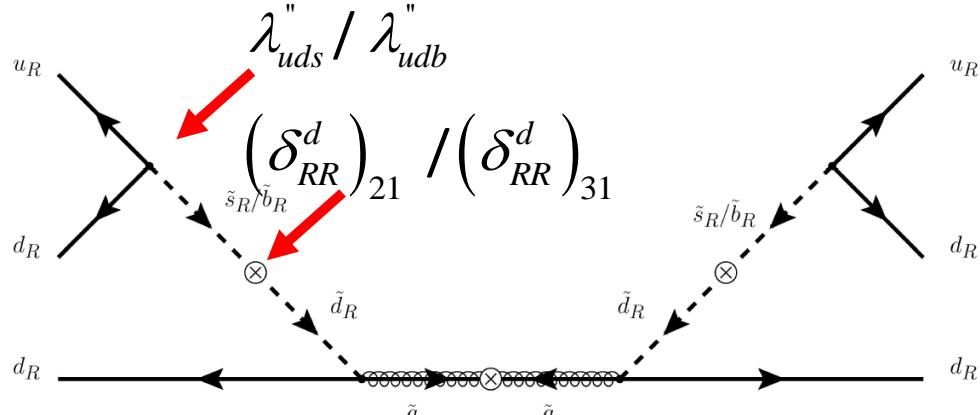


Simplified models: sparticle masses, Yukawa coupling and mixing terms.

JHEP 1605 (2016) 144 (arXiv:1602.04821)

L. Calibbi, G. Ferretti, D. Milstead, C. Petersson, R. Pöttgen

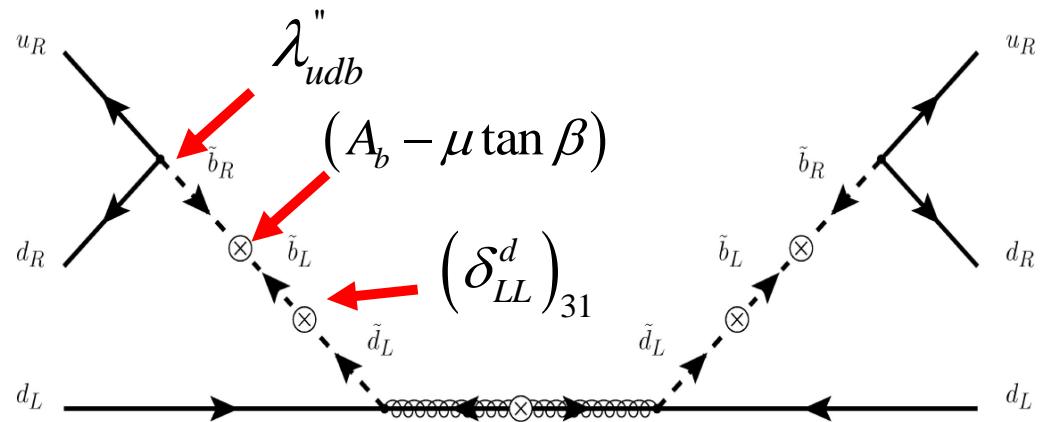
Scenarios – strong



Zwirner

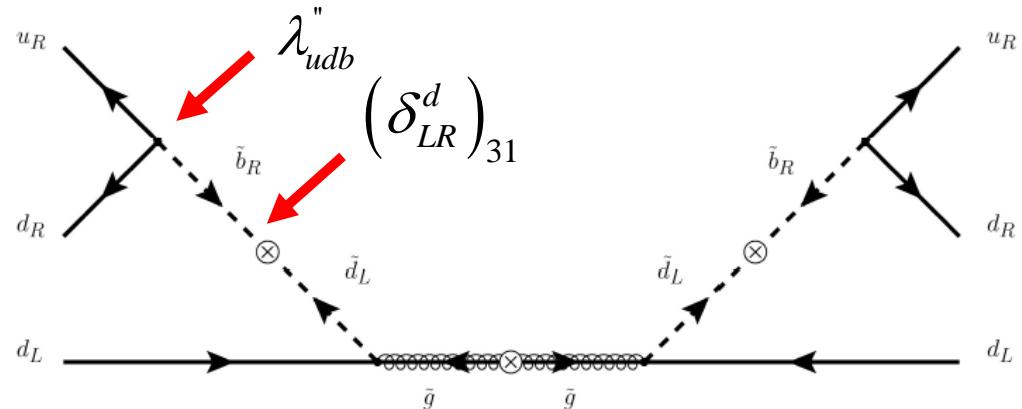
$$Z_1 : \lambda''_{uds}, (\delta^d_{RR})_{21}$$

$$Z_2 : \lambda''_{udb}, (\delta^d_{RR})_{31}$$



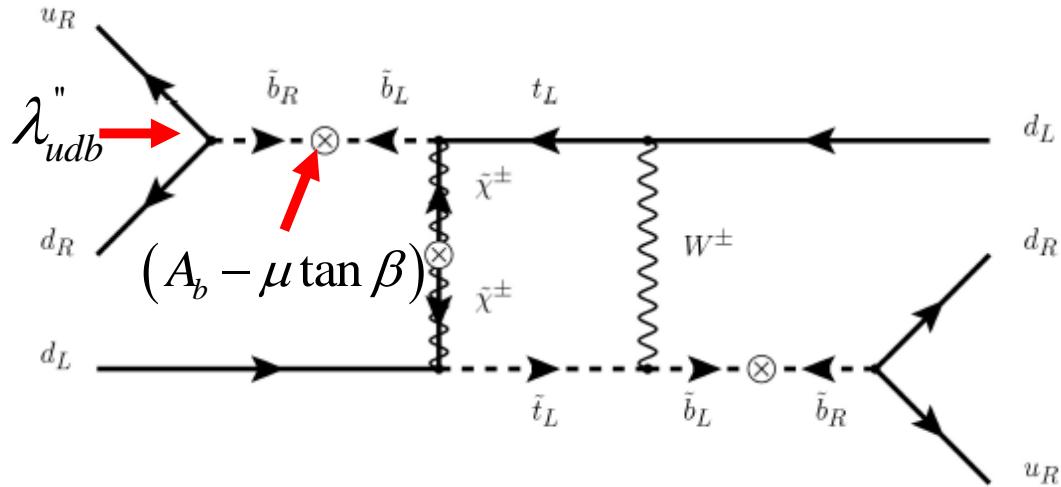
Barbieri and Masiero

$$BM_1 : \lambda''_{udb}, (\delta^d_{LL})_{31}, (A_b - \mu \tan \beta)$$



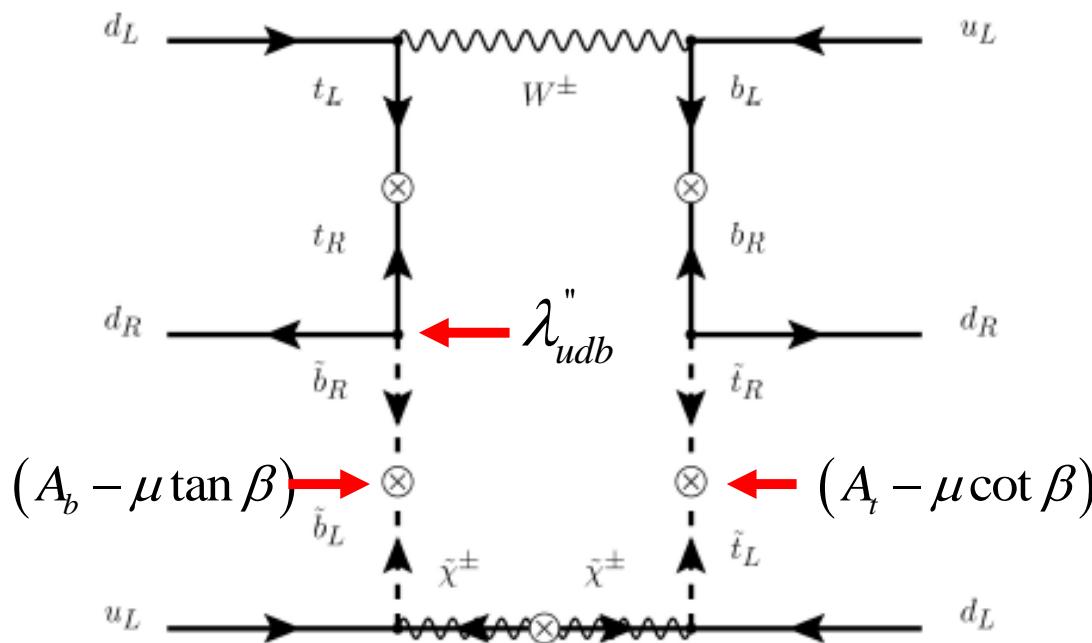
$$BM_2 : \lambda''_{udb}, (\delta^d_{LR})_{31}$$

Scenarios – electroweak



Goity and Sher :

$$\lambda_{udb}'', (A_b - \mu \tan \beta)$$



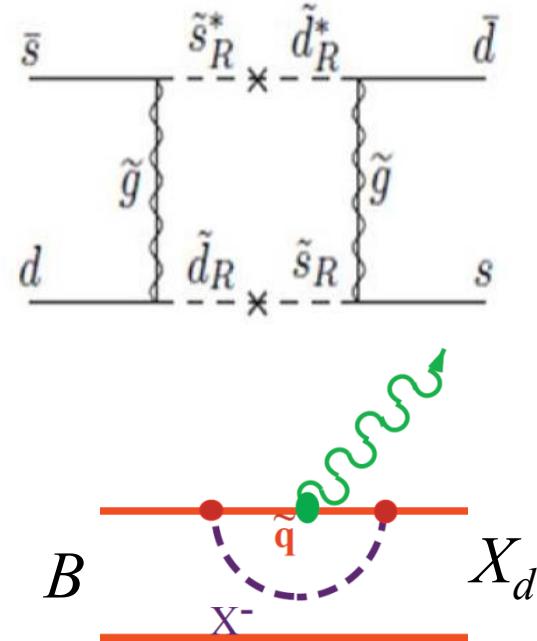
Chang and Keung :

$$\lambda_{tdb}'', (A_b - \mu \tan \beta), (A_t - \mu \cot \beta)$$

Experimental constraints

Flavour/mixing

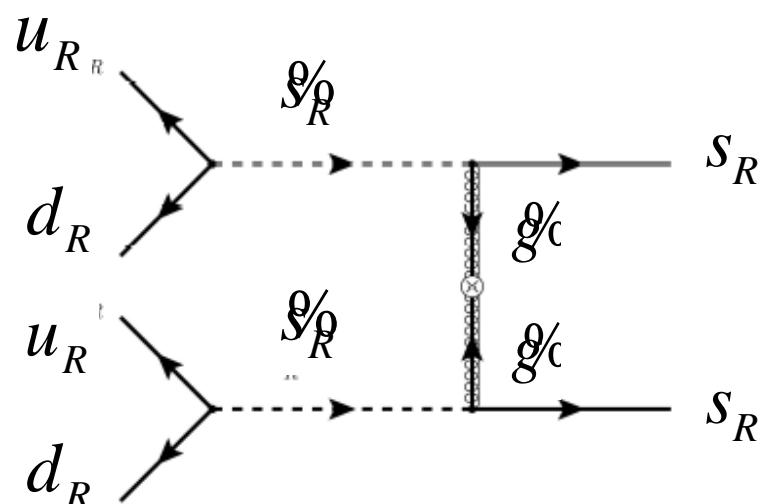
Observable	Parameter
Kaon mixing	$(\delta_{RR}^d)_{21}$
B-mixing	$(\delta_{RR}^d)_{31}$
$b \rightarrow d + \gamma$	$\mu \tan \beta, (\delta_{RR}^d)_{31}$



Low energy BNV

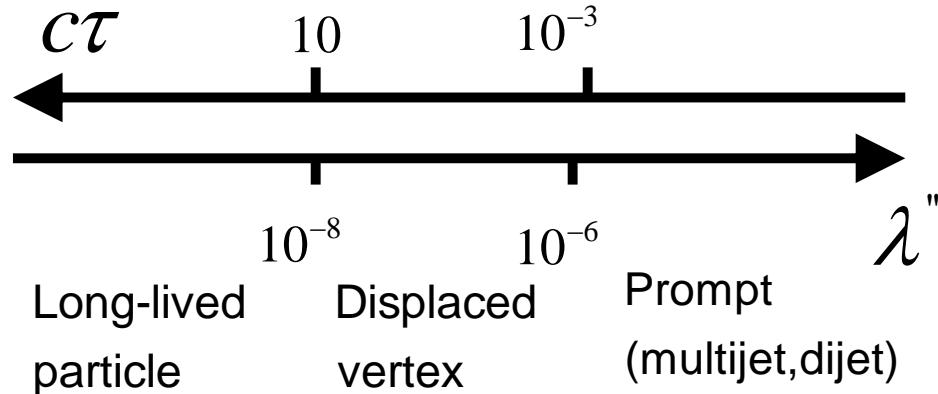
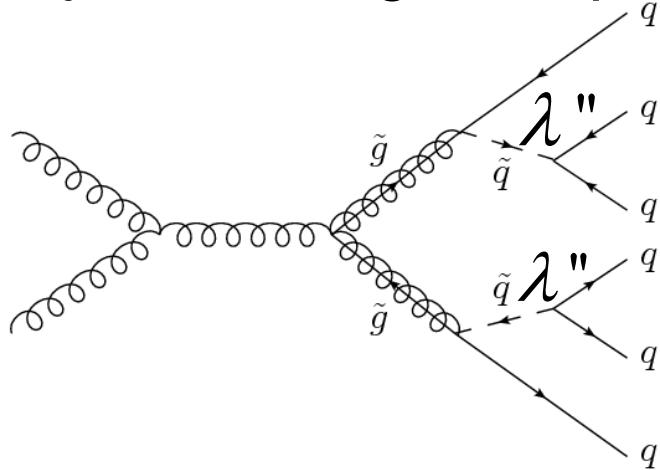
Observable	Parameter
$n \rightarrow \bar{n}$	$\lambda''_{112}, \lambda''_{113}$
$NN \rightarrow \text{mesons}$	

Limits from Super-K



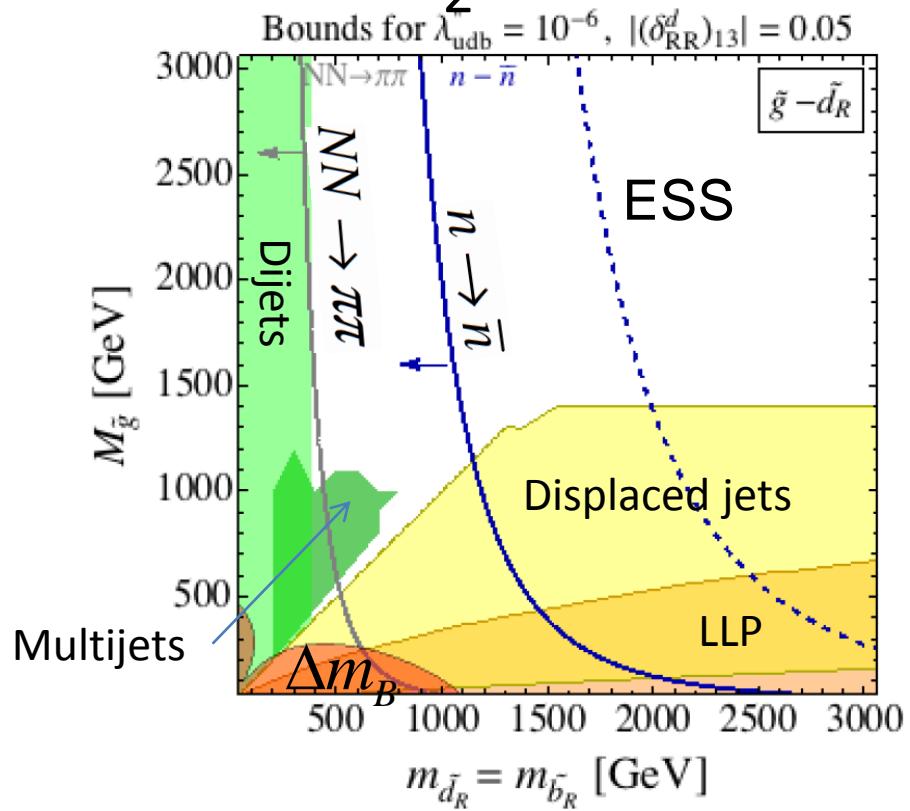
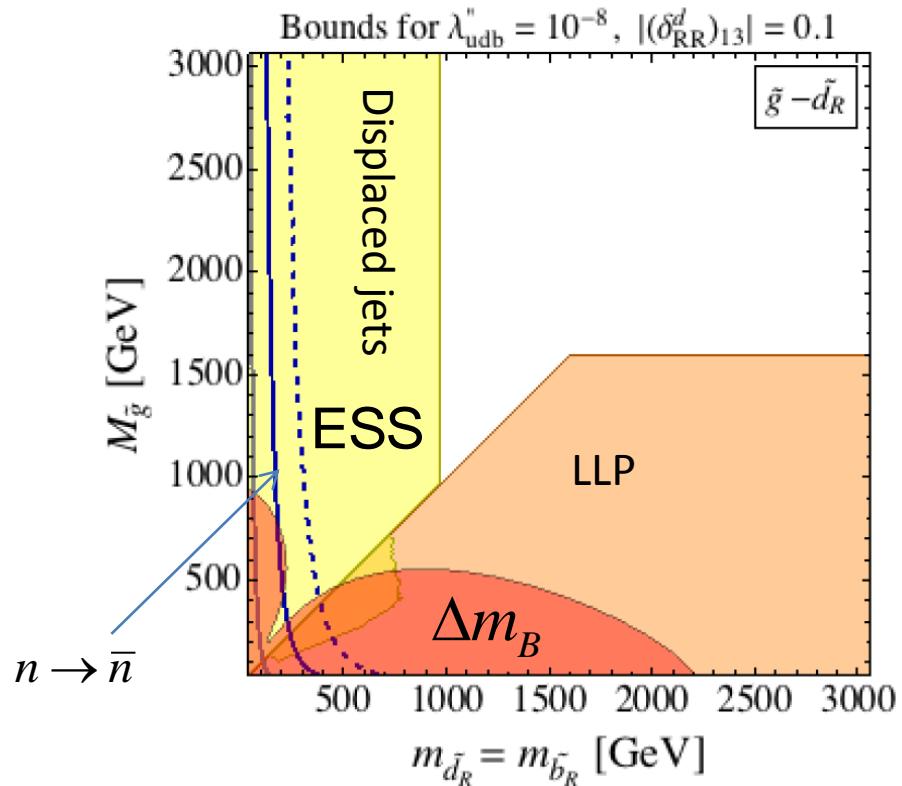
Experimental constraints - LHC

Multijet and long-lived particle signatures



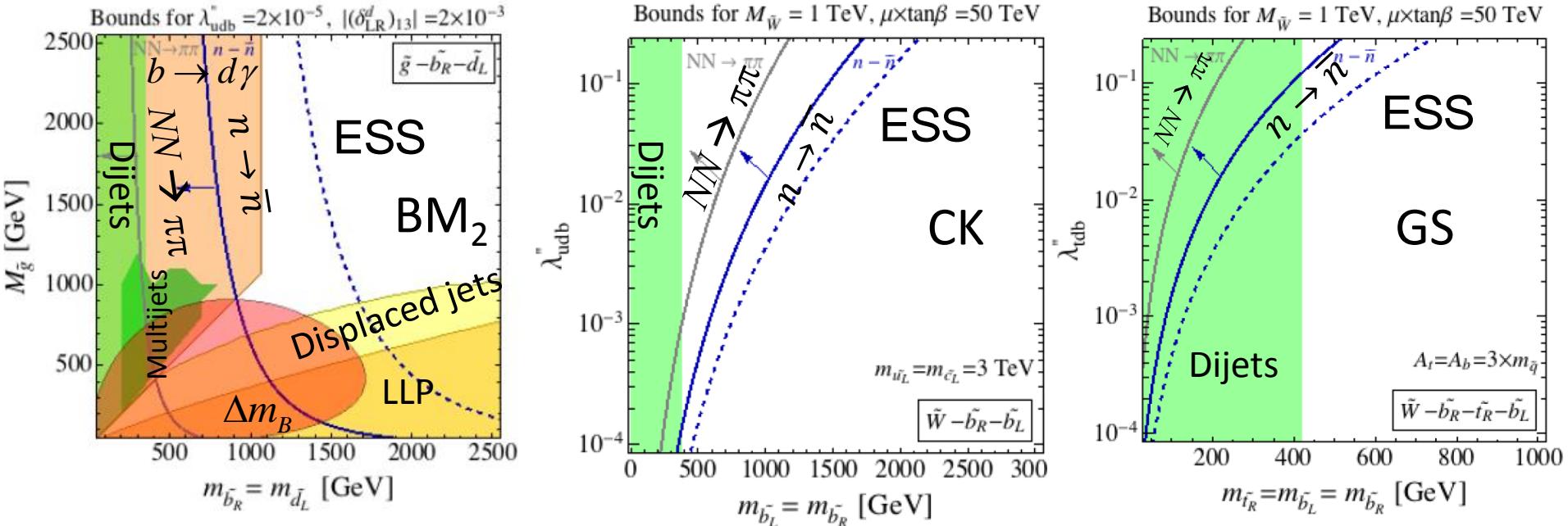
Observable	Parameter
ATLAS Multijets (Arxiv:1602.04821 hep-ex) recast with Madgraph+Pythia+Delphes	$\lambda''_{112}, \lambda''_{113}$
CMS Dijets Arxiv:1412.7706	
ATLAS/CMS Displaced vertex+ long-lived particle recast (arxiv:1503.05923, 1505.00784 hep-ph, CMS-PAS-EXO-15-010)	

Model exclusion – Z_2



LHC is the dominant player at $\sim \text{TeV}$

Model exclusion: BM_2 , CK, GS



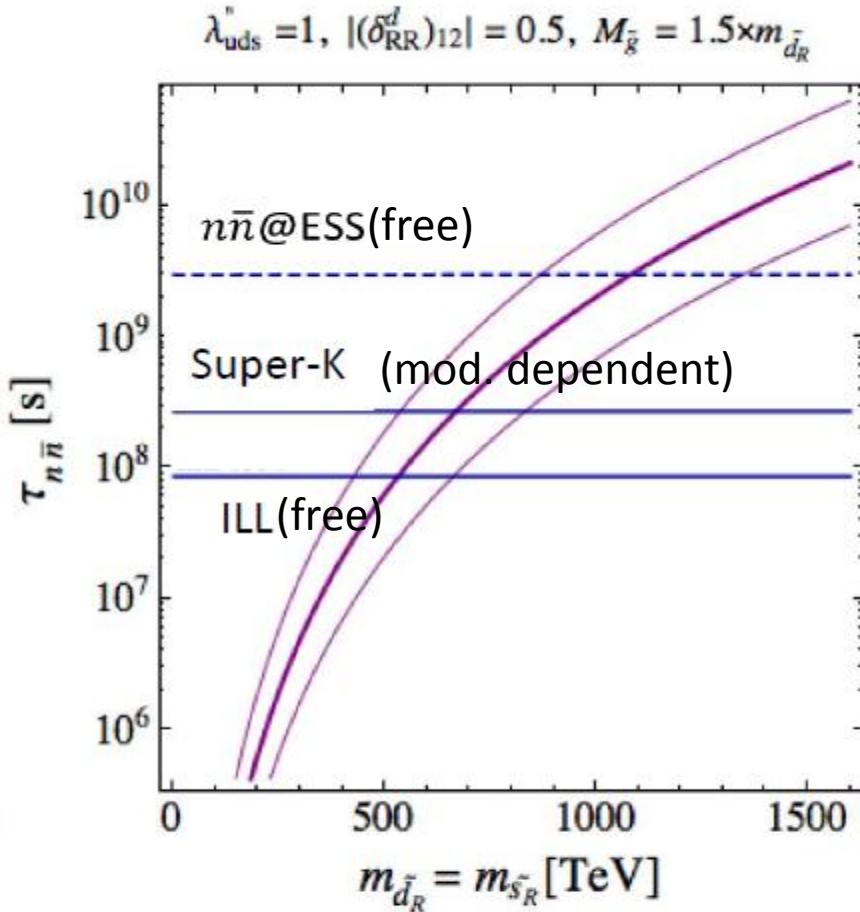
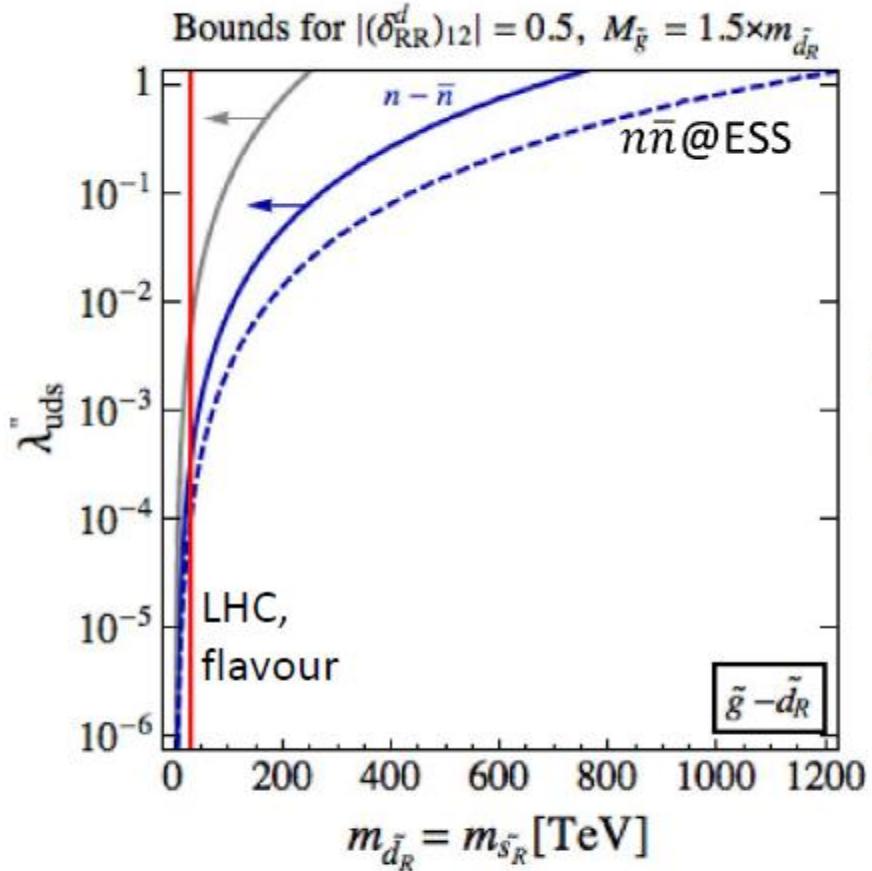
Consistent picture:

Dedicated BNV expts. sensitive to higher mass scales
than LHC and flavour experiments.

Dependent on the coupling and mixing values.

Searches are complementary.

Beyond the TeV scale



Constraints vanish for \gg TeV masses

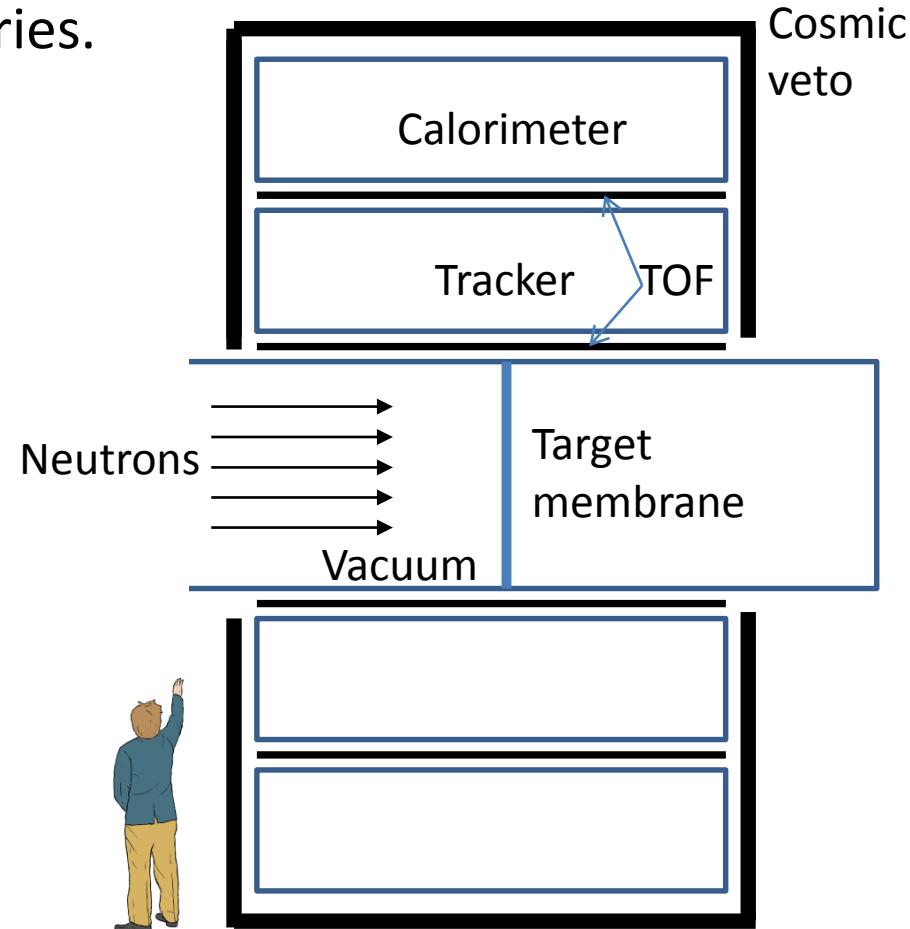
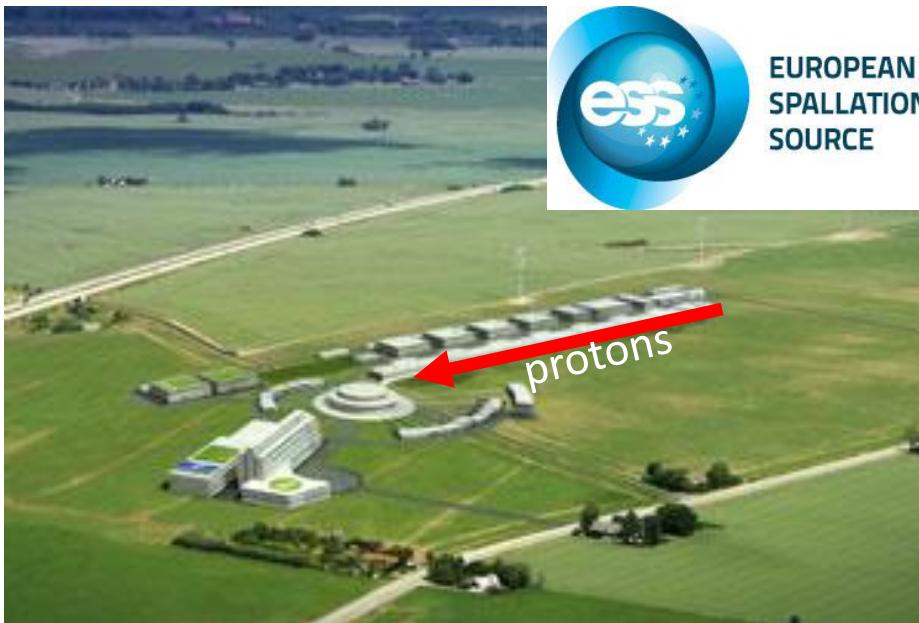
nnbar@ESS: extends mass range by up to ~ 400 TeV cf Super-K
: pushes into the PeV scale
: Reach beyond the LHC

A search for $n \rightarrow \bar{n}$ at the European Spallation Source

nnbar@ESS Collaboration: Co-Spokespersons G. Brooijmans, D. Milstead

Expression of Interest submitted to ESS.

Signatories from 26 institutes , 8 countries.



First search for free neutron oscillation since 1991

Sensitivity increase $\times 10^3$ for $P(n \rightarrow \bar{n})$

The proposed program

Stage 1

HIBEAM - high intensity baryon extraction and measurement

Early to late 2020s

- Match or improve sensitivity to $P(n \rightarrow \bar{n})$ wrt previous search at ILL
- Search for mirror neutrons (regeneration)
- R&D for full experiment (*NNBAR*)

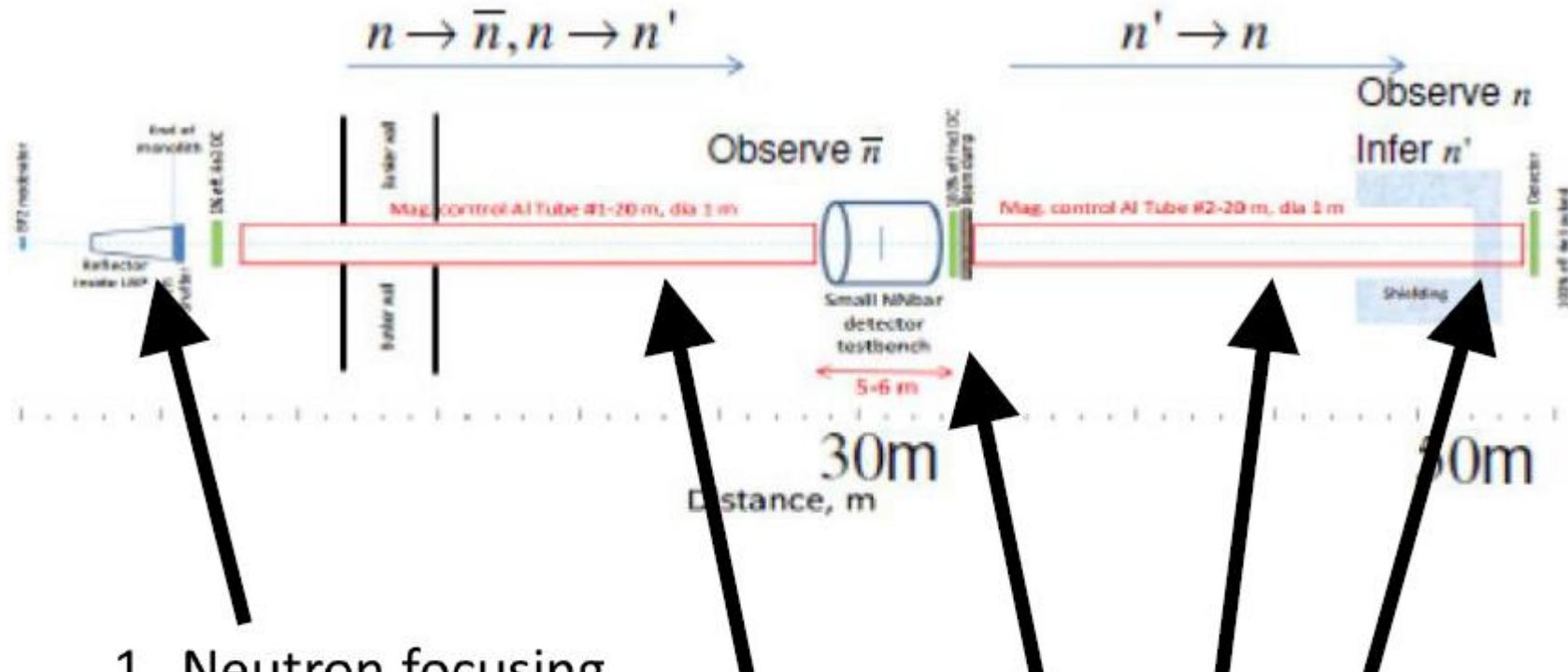
Stage 2

NNBAR experiment

Late 2020's + 5 years

- Improve sensitivity to $P(n \rightarrow \bar{n})$ by $\sim 10^3$
- Further mirror neutron searches

HIBEAM and nnbar



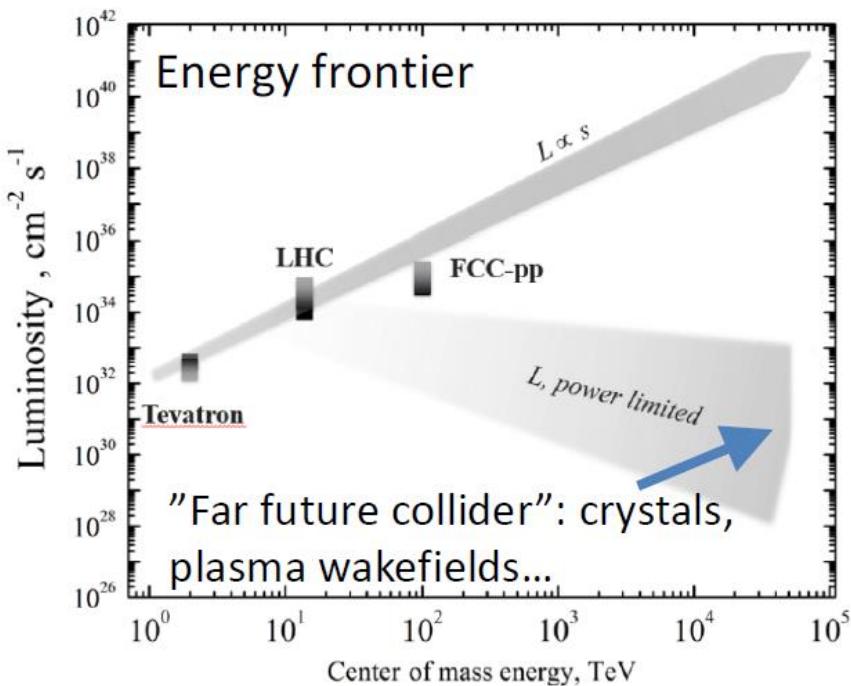
1. Neutron focusing
2. Passage through shielded tube
3. If $n \rightarrow \bar{n} \rightarrow$ annihilation of \bar{n} in C target
4. If $n \rightarrow n' \rightarrow$ passage through absorber
5. $n' \rightarrow n \rightarrow$ measure n in neutron counter

Nnbar experiment – extend to ~300m

Summary

- BNV-only searches are an important part of program to test symmetries, search for BSM physics and understand baryogenesis.
- If Nature chose BNV-only She hid it but we know how to look.
- RPV SUSY provides a framework to study sensitivity.
- A future dedicated $n \rightarrow \bar{n}, n'$ search can massively extend sensitivity.

We live in interesting times



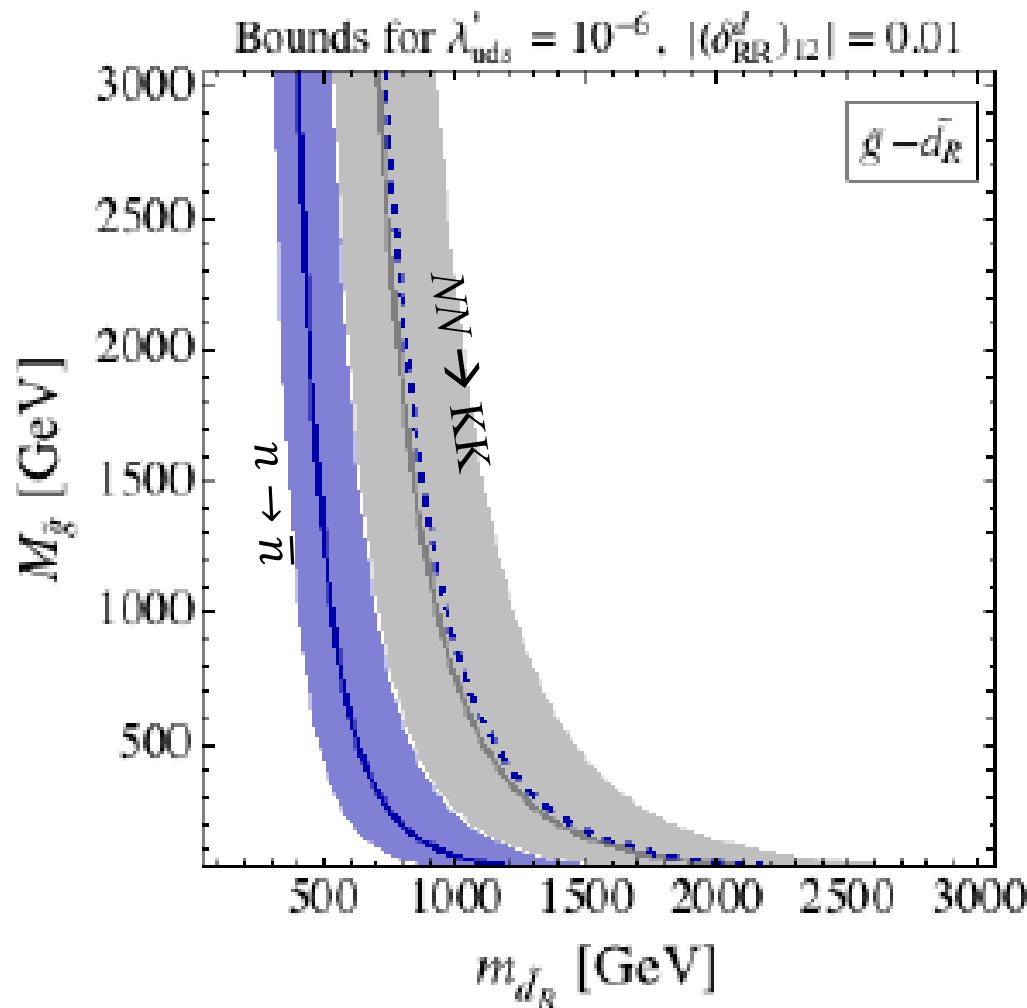
Future discoveries or walking a few km in a desert ?



For the first time in 50 years, going to higher collision energies no longer offers a clear path to discoveries or fundamental insights.

Need a complementary set of collider +non-collider experiments with unique physics potentials and reach of energy scale.

Uncertainties from hadronic matrix elements



Operators and lifetime

Six quark operators O_i of dimension-9 :

$$(u_R d_R d_R)^2 \equiv \epsilon_{abc} u_{R\dot{\alpha}}^a d_R^{\dot{\alpha}b} d_R^c \epsilon_{def} u_{R\dot{\beta}}^d d_R^{\dot{\beta}e} d_R^{\dot{\gamma}f}$$

$$(u_R d_R d_L)^2 \equiv \epsilon_{abc} u_{R\dot{\alpha}}^a d_R^{\dot{\alpha}b} d_L^{\dot{\gamma}c} \epsilon_{def} u_{R\dot{\beta}}^d d_R^{\dot{\beta}e} d_L^f$$

$$(u_L d_L d_R)^2 \equiv \epsilon_{abc} u_L^{\alpha a} d_{L\alpha}^b d_R^c \epsilon_{def} u_L^{\beta d} d_{L\beta}^e d_R^{\dot{\gamma}f}$$

$$(u_R d_R s_R)^2 \equiv \epsilon_{abc} u_{R\dot{\alpha}}^a d_R^{\dot{\alpha}b} s_{R\dot{\gamma}}^c \epsilon_{def} u_{R\dot{\beta}}^d d_R^{\dot{\beta}e} s_R^{\dot{\gamma}f}.$$

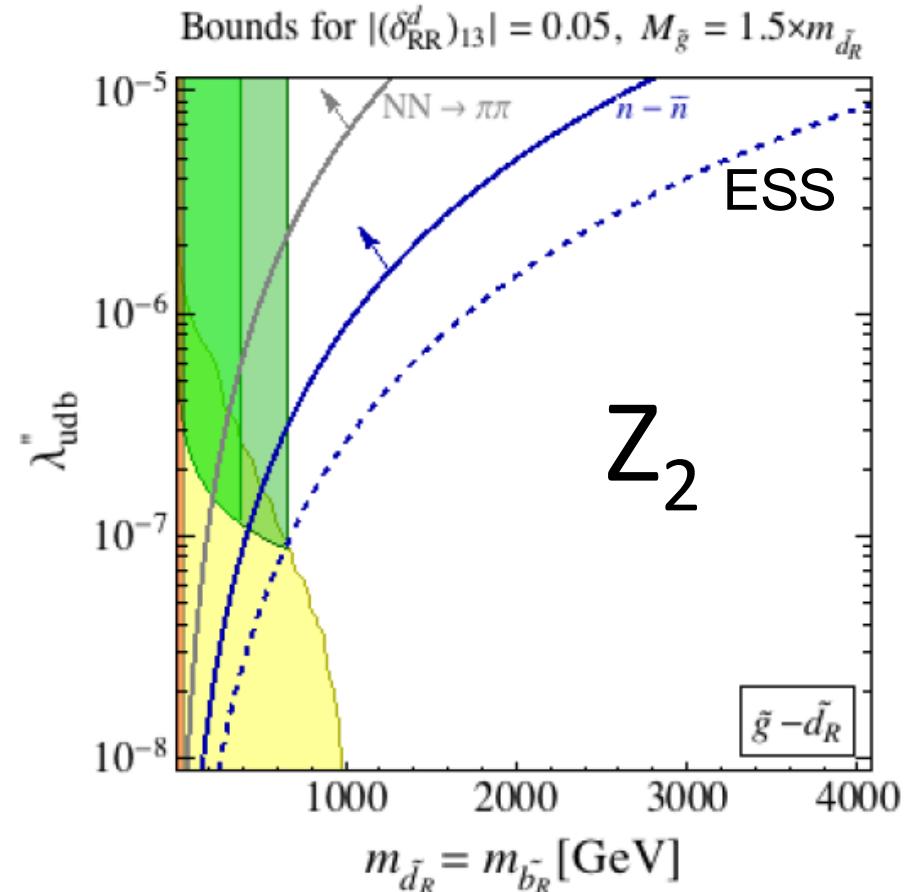
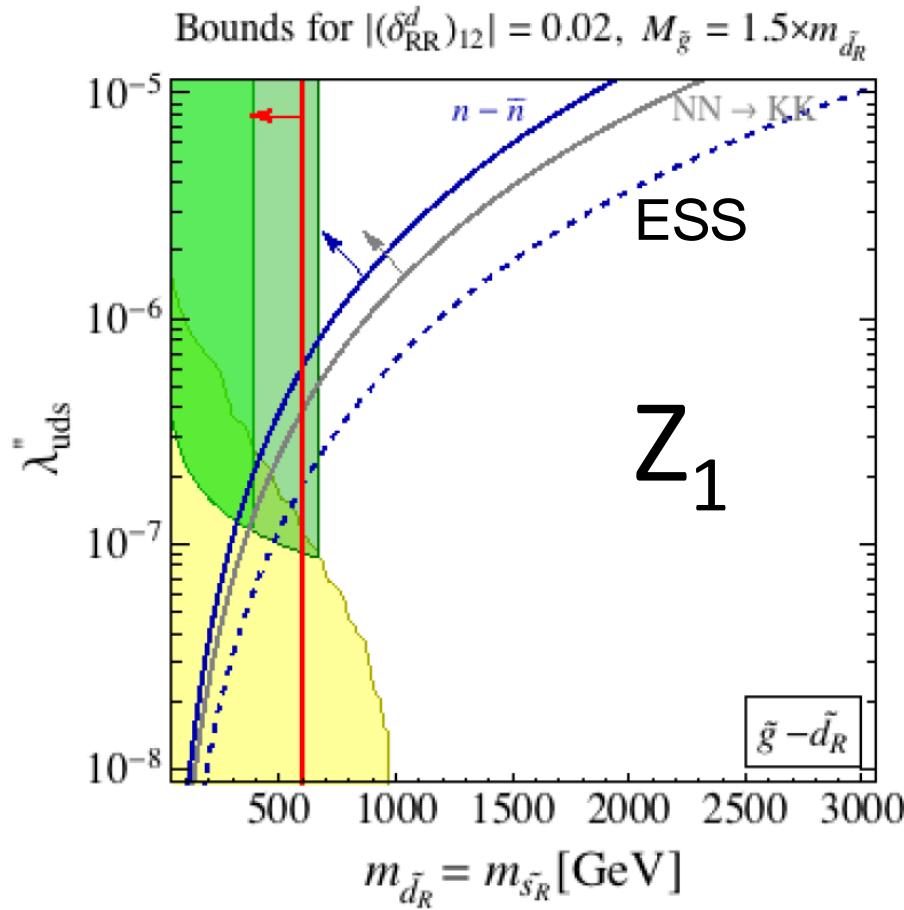
$n \rightarrow \bar{n}, NN \rightarrow \pi\pi$

$NN \rightarrow KK$

Eg $n \rightarrow \bar{n}$, Zwirner:

$$\tau = (2.5 \times 10^8 \text{ s}) \times \frac{(250 \text{ MeV})^6}{\langle \bar{n} | (u_R d_R d_R)^2 | n \rangle} \times \frac{m_{g_0}}{1.2 \text{ TeV}} \left(\frac{\bar{m}_D}{500 \text{ GeV}} \right)^4 \left(\frac{10^{-6}}{\lambda_{uk}''} \right)^2$$

Model exclusion – Z_1, Z_2



Dedicated BNV searches give sensitivity beyond the LHC.