

RG Fixed Point with a 4th Generation

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- Implications

Common objection to non-SUSY SM with fundamental scalars:

The need for fundamental scalar fields in the theory of weak and electromagnetic forces is a serious flaw. They require unnatural adjustments of bare constants (Wilson, Susskind,...)

• Recall:
$$V = \mu_0^2 \phi^{\dagger} \phi + \lambda (\phi^{\dagger} \phi)^2$$

 $\lambda > 0 \text{ and } \mu_0^2 < 0 \Rightarrow \text{SSB of SM:}$
 $\langle \phi^0 \rangle = \sqrt{\frac{-\mu_0^2}{2\lambda}} \approx 174 \, GeV.$

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• Quantum corrections to μ_0^2 can drastically alter that value.

• Flaw: Corrections to the scalar mass e.g. from a fermion loop: $\mu^2 = \mu_0^2 - \frac{g_f^2}{16\pi^2}\Lambda_{max}^2 + \dots$ $(\Lambda_{max}: UV Cutoff)$ ("Quadratic divergence")

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- Why is it a flaw? If the physical cut-off $\Lambda_{max} \sim M_{Pl} \sim 10^{19} \, GeV$, one expects $\mu^2 \sim O(M_{Pl}^2)$ (assuming $g_f = O(1)$) unless...

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• One fine-tunes the cancellation to at least 34 decimal places to make $\mu \sim \Lambda_{EW} \sim O(100 \, GeV) \Rightarrow$ Naturalness and Hierarchy (between M_{Pl} and Λ_{EW}) problems.

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- Each momentum slice represents a distinct set of physics, e.g. $\lambda \phi^4 \Rightarrow$ Same equation of motion but with different couplings for different slices.
- Divergence arises when $n \to \infty$: The physics contribution from each momentum slice is finite but there is an infinite number of such slices.

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- UV cutoff Λ_{max} large (e.g. M_{Pl}): Need to cancel the "quadratic divergence". Very attractive symmetry: SUSY.
- Exact SUSY: Cancellation between bosons and fermions to all orders \Rightarrow No "quadratic divergence". However, the absence so far of SUSY partners of SM particles \Rightarrow Broken SUSY. To avoid reintroducing the "quadratic divergence" \Rightarrow Softly Broken SUSY \Rightarrow $\delta\mu^2 \propto m_{soft}^2$. The requirement $m_{soft} \leq \Lambda_{EW} \Rightarrow$ SUSY partners < O(TeV).

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- Bright Side: A possible "solution" to the hierarchy problem; Dark Matter candidates; Lots of particles to search for;...
- Flip Side: Lots of parameters: More than one hundred; FCNC problems; How SUSY is broken is still unclear; Has its own fine-tuning problem: the μ-problem,...

$\Lambda_{max} \sim O(TeV)$: Composite Higgs

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- Many interesting Top-quark-like Condensate Models discussed in the workshop

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- Even in a massless theory, scale invariance can be broken at the quantum level.
- Divergence of dilatation current = $\theta^{\mu}_{\mu} \propto \beta(g)$
- If $\beta(g) = 0$ and there are no explicit masses, scale invariance is valid \Rightarrow Absence of quadratic and quartic divergences.

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- However, could some approximation such as $\beta_{2-loop}(g^*) = 0$ hint at the energy scale where scale invariance might be restored?
- In particular, under what conditions would that energy scale be $\Lambda_{max} \approx O(TeV)$ with minimal enlargement of the 3-generation SM?
- Would a heavy 4th generation do it and what are the implications? Work done with Chi Xiong.

What can a heavy 4th generation do?

● A heavy Higgs-4th Yukawa system can give rise to condensates and bound states of 4th generation fermions ⇒ Implications on the vacuum structure of the SM and the number of Higgs doublets (a mixture of fundamental and composites)

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- A heavy Higgs-4th Yukawa system can give rise to condensates and bound states of 4th generation fermions ⇒ Implications on the vacuum structure of the SM and the number of Higgs doublets (a mixture of fundamental and composites)
- The appearance of quasi-fixed points in the 2-loop approximation at a scale $\Lambda_{FP} \sim O(TeV)$ hints at a possible scale-invariant theory above Λ_{FP} .

Higgs-Yukawa system with four generations

• Run the 2-loop RG equations: $16\pi^2 \frac{dY}{dt} = \beta_Y$ from Λ_{EW} on, with $Y = \lambda, g_t^2, g_q^2, g_l^2$ (quartic, top, 4th quark, 4th lepton couplings).

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- Generically, $\beta/16\pi^2 = g(\beta_0(\alpha/4\pi) + \beta_1(\alpha/4\pi)^2 + ..)$ (see e.g.Abbott 1980)

Higgs-top Yukawa system

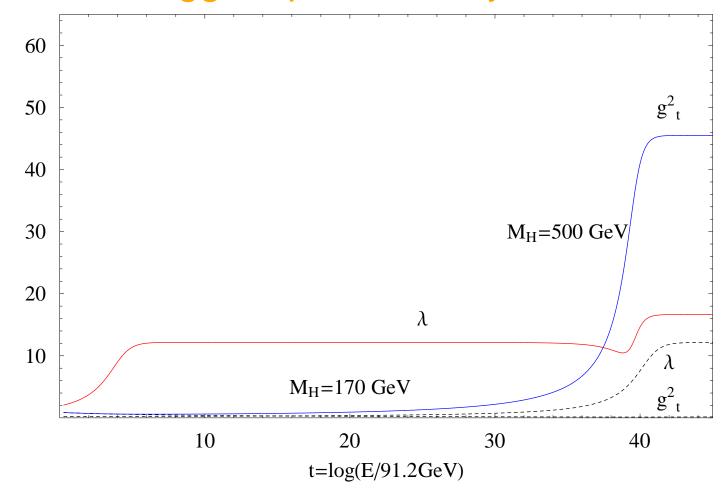
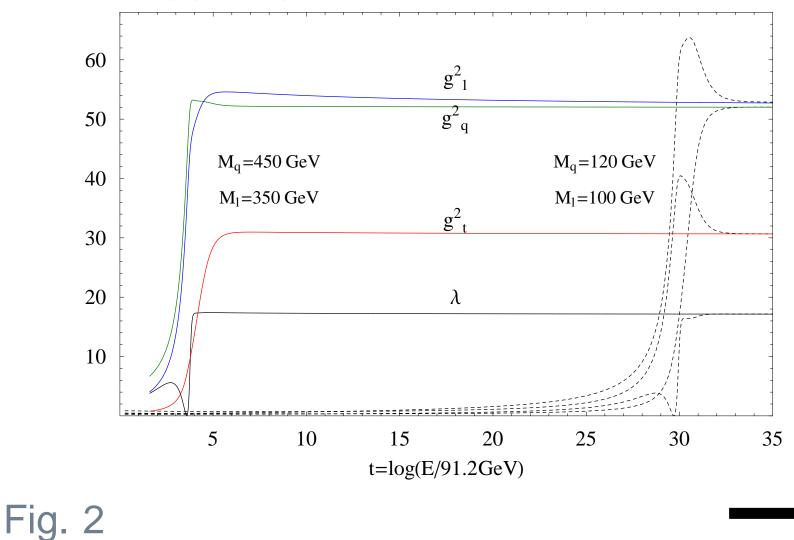


Fig.

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Higgs-4 generation Yukawa system



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- In Fig. 2, the light 4th generation (120 GeV and 100 GeV) case is shown only for comparison. It is excluded experimentally.
- The masses shown in Fig. 2 will be modified when extra dynamical Higgs doublets are included (discussed below)
- Notice the values of the quasi fixed point expansion parameters: $\alpha_{q,l}^*/4\pi \approx 0.3$, $\alpha_t^*/4\pi \approx 0.2$, $\lambda^*/16\pi^2 \approx 0.1$. More on this below.

Higgs-top Yukawa system

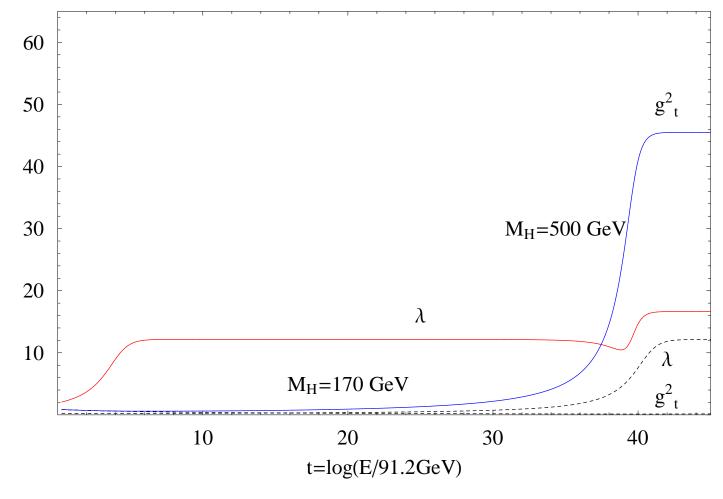


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Comments on Higgs-top Yukawa system: Fig. 1

In Fig. 1, light Higgs (170 GeV) ⇒ No quasi-fixed point below the Planck scale $(t_{pl} \approx 39)$.

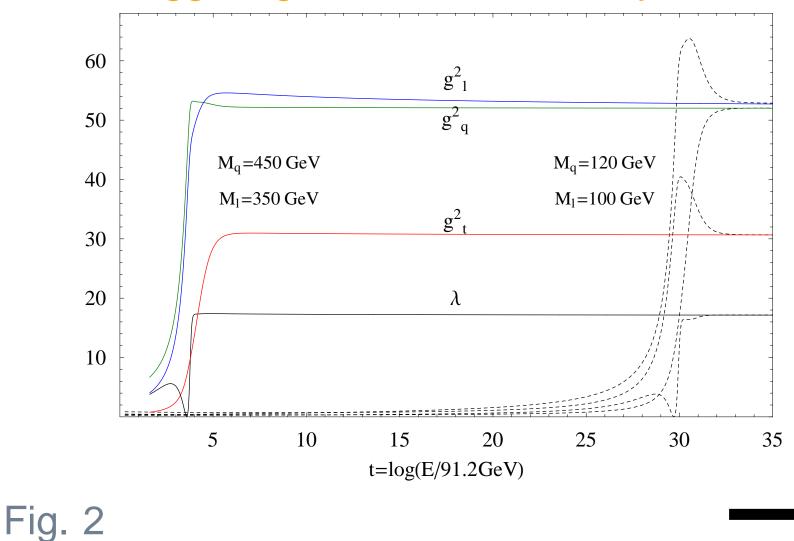
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- For heavy Higgs (500 GeV) \Rightarrow Quasi fixed point just above t_{pl}
- Notice at that quasi fixed point, $\alpha_t/4\pi \approx 0.2$

Higgs-4 generation Yukawa system



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Comments on Higgs-4 generation Yukawa system: Fig. 2

• The light $(M_q = 120 \, GeV, M_l = 100 \, GeV)$ and the heavy $(M_q = 450 \, GeV, M_l = 350 \, GeV)$ reach the same quasi-fixed point but at different momentum scales \Rightarrow Hint of some real fixed point so $\beta = 0 \Rightarrow$ Hint of scale invariance.

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- Vacuum stability requires $\lambda > 0 \Rightarrow$ Fig. 2 shows the minimum values of the initial λ which satisfy the stability criterion.

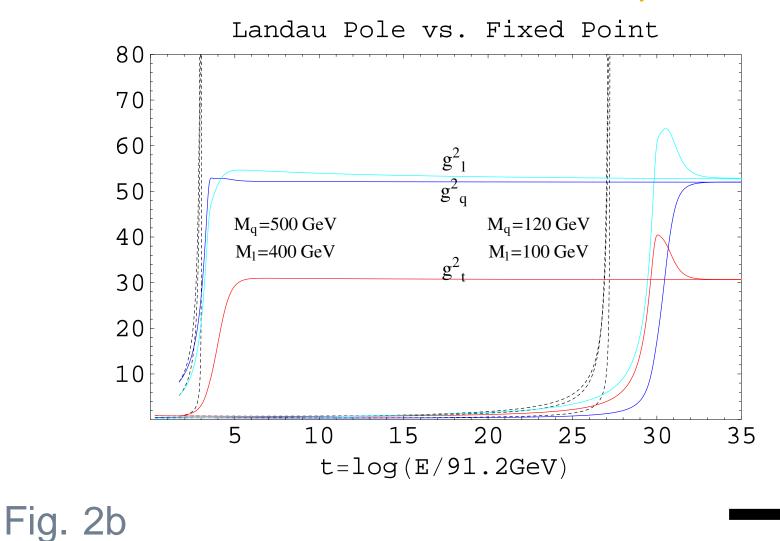
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- At the dip where $\lambda = 0$, the Yukawa coupling is large enough to have a "strong Coulomb-like" condensate formation of 4th generation fermions.

Landau Pole vs Quasi-fixed point



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- The Landau pole (one loop) and the Quasi-fixed point occur at around the same place Λ_{FP} .
- Hint at something occuring at Λ_{FP} for a heavy 4th generation: perhaps a scale-invariant theory?

Bound States and Condensates in Higgs-4 generation Yukawa system

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- Bound states of 4th generation quarks and leptons can get formed by the exchange of the Higgs scalar.
- For example, a non-relativistic short-range Yukawa potential for a $\overline{F}F$ bound state can be written as $V(r) = -\alpha_Y(r) \frac{e^{-m_H(r)r}}{r}$ with

 $\alpha_Y = \frac{m^2}{4\pi v^2}$

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• Numerical solution to the Schrödinger equation yields a constraint (negative bound state energy): $K_f = \frac{g_f^3}{16\pi\sqrt{\lambda}} > 1.68$.

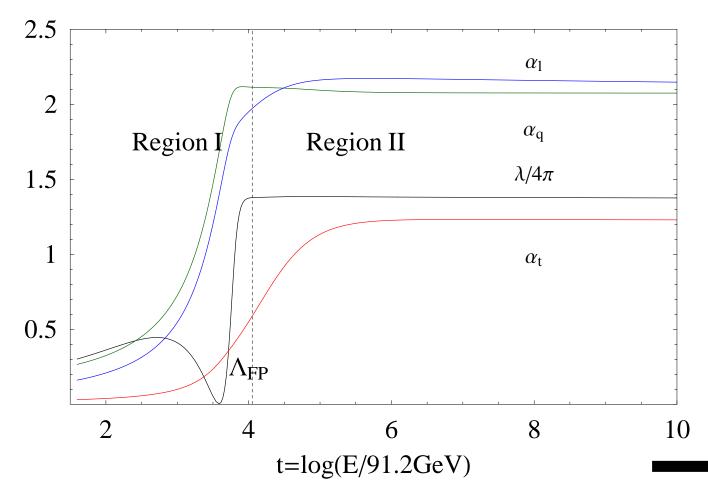
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- Around the quasi-fixed point, $K_q = 1.82$, $K_l = 1.92$, $K_t = 0.82 \Rightarrow \text{No } t\bar{t}$ bound state there.

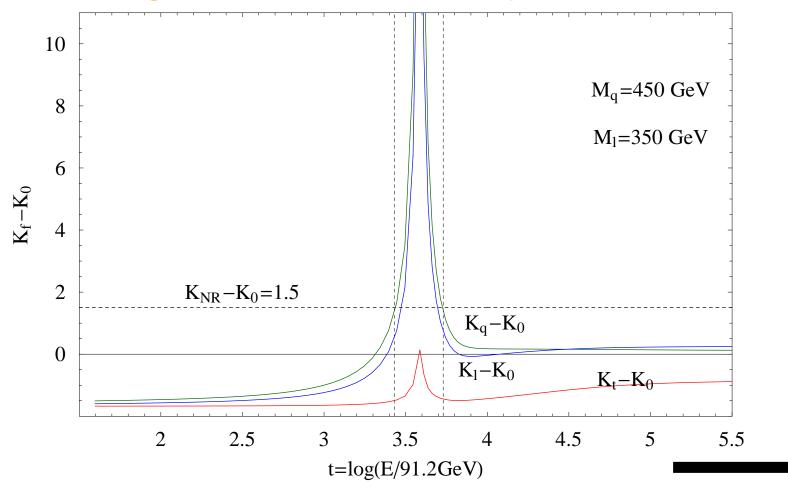
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- Binding energies: $E_q \approx -4.9$ GeV, $E_q \approx -16$ GeV ⇒ Very loose bound states in Region II of the following graph.

Bound States and Condensates in Higgs-4 generation Yukawa system



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- Region II shows the quasi fixed point ⇒ Loose bound states of the 4th generation fermions. This corresponds to the right hand side of Fig. 4.
- Region I shows the "dip" where $\lambda \to 0$. As one approaches the dip, the correlation length $\xi_H \sim 1/m_H$ goes from a short-range correlation (small ξ) to an infinite-range correlation ($\xi = \infty$) with a "strong Coulomb-like potential" $V_{dip}(r) = -\frac{\alpha_Y}{r} \Rightarrow$ Possibility of condensate formation.

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- A richer vacuum structure due to this possibility of condensate formation of 4th generation fermions.
- There could be an interesting phenomenology with these new composite Higgses.

• A heavy (~ 400 - 500 GeV) 4th generation can lead to a quasi fixed point at around $\Lambda_{FP} \sim O(TeV) \Rightarrow$ Hint of scale invariance in that region \Rightarrow Hint of a possible solution to the hierarchy problem

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- The vacuum structure (and its phenomenological implications) of the SM with a 4th generation is richer and more interesting.
- Much more work needs to be done!