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# B and K Physics in the scenario of one extra Fermion family

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# Plan of talk

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- CP violation and New Physics
- $B$  and  $K$  Physics : SM and BSM
- Analysis and Results : SM4
- Conclusions

# CP violation and NP

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- CP violation is crucial for the fact that we are here ... that matter dominates over antimatter in our universe
- New sources of CP violation must exist in nature  $\Rightarrow$  Strong *CP* problem, History of matter and antimatter, Hierarchy problem...
- New measurement of CP asymmetries in meson decays are sensitive to new sources of CP violation that come from physics at the few TeV

CP violation provides an excellent probe for NP

# $CP$ violation and $K$ Physics

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Direct  $\rightarrow CP$  phase appear in decay,

Indirect  $\rightarrow CP$  phase appear in mixing

- Indirect  $CP$  violation in  $K \rightarrow \pi\pi$  and  $K \rightarrow \pi l\nu$  is given by

$$|\epsilon_k| = (2.32 \pm 0.007) \times 10^{-3}$$

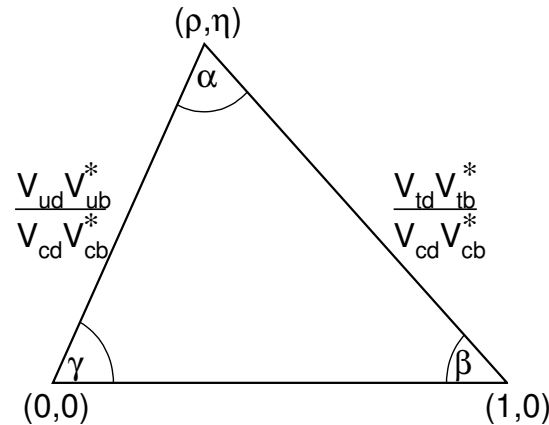
- Direct  $CP$  violation in  $K \rightarrow \pi\pi$  decays is given by

$$\epsilon'/\epsilon = (1.65 \pm 0.26) \times 10^{-3}$$

- $K_L \rightarrow \pi^0 \nu \bar{\nu}$  appears to be the best decay to measure

$$Im(\lambda_t)$$

# An important Goal of B Physics



The area of each triangle =  $|J|/2$ ,  $J \approx \lambda^6 A \eta$

- $|V_{ub}|^2 \propto \rho^2 + \eta^2 \rightarrow$  Charmless semileptonic decay rates
- $S_{\psi K_s} = \sin 2\beta$ ,  $CP$  asymmetry in  $B \rightarrow \psi K_s$ ,  
 $e^{i\beta} = 1 - \rho + i\eta$
- $e^{i\gamma} = \rho + i\eta$ ,  $\gamma$  is sensitive to various  $B \rightarrow DK$  decay rates
- $\alpha$  is sensitive to  $B \rightarrow \pi\pi, \rho\pi, \rho\rho$

# Motivation: $B$ - $CP$ anomalies

- $CP$  violation in Gold plated decay mode  $B \rightarrow \psi K_s \Rightarrow S_{\psi K_s} = 0.672 \pm 0.024$ 
  - SM predictions:  $\sin 2\beta_{no V_{ub}} = 0.87 \pm 0.09$
  - $\sin 2\beta_{fullfit} = 0.75 \pm 0.04$   
 $\Rightarrow 1.7$  to  $2.1 \sigma$  deviation.. Lunghi and Soni, 2008
- $\sin 2\beta$  extracted from  $B_d \rightarrow \phi K_s \Rightarrow 0.39 \pm 0.17$
- $\Delta A_{CP} \equiv A_{CP}(B^- \rightarrow K^- \pi^0) - A_{CP}(\bar{B}^0 \rightarrow K^- \pi^+) = (14.4 \pm 2.9)\%$ ,  $QCDF$  prediction  $\rightarrow (2.5 \pm 1.5)\%$ .
- $CP$  asymmetry in gold plated decay mode  $B_s \rightarrow \psi \phi$ ,  $S_{\psi \phi} = \sin \phi_s \Rightarrow [-0.90, -0.17]$ ,  $2\sigma$  range... CDF and D0, 2008,  
 $\phi_s = \arg[-\frac{M_{12}}{\Gamma_{12}}] \rightarrow CP$  violating weak phase..  $\Rightarrow S_{\psi \phi}^{SM} \approx 0$

# Table of Inputs

A. Soni, A. K. Alok, A. Giri, R. Mohanta and SN, arxiv:0807.1971 [hep-ph]

$B_K = 0.72 \pm 0.05$ , Antonio <i>et al.</i> , 2008 $f_{bs}\sqrt{B_{bs}} = 0.281 \pm 0.021$ GeV $\Delta M_s = (17.77 \pm 0.12)ps^{-1}$ $\Delta M_d = (0.507 \pm 0.005)ps^{-1}$ $\xi_s = 1.2 \pm 0.06$ , Gamiz <i>et al.</i> , 2007 $R_{bb} = 0.216 \pm 0.001$ $ V_{cb}  = (40.8 \pm 0.6) \times 10^{-3}$	$ \epsilon_k  \times 10^3 = 2.32 \pm 0.007$ $BR(K^+ \rightarrow \pi^+ \nu\nu) = (0.147_{-0.089}^{+0.130}) \times 10^{-9}$ $BR(B \rightarrow X_s \gamma) = (3.55 \pm 0.25) \times 10^{-4}$ $BR(B \rightarrow X_s \ell^+ \ell^-) = (0.44 \pm 0.12) \times 10^{-6}$ $\gamma = (75.0 \pm 22.0)^\circ$ $ V_{ub}  = (37.2 \pm 2.7) \times 10^{-4}$ $T_4 = 0.11 \pm 0.14$
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Constrain used  $\Rightarrow$  **CKM unitarity + Oblique corrections +**  
 $Br(B \rightarrow X_s \gamma) + \epsilon_k + \Delta M_d + \frac{\Delta M_d}{\Delta M_s} + Br(B \rightarrow X_s \ell^+ \ell^-) +$   
 $BR(K^+ \rightarrow \pi^+ \nu\nu)$

Allowed ranges for the parameters,  $\lambda_{t'}^s = \mathbf{V}_{t'b} \mathbf{V}_{t's}^*$ , phase of  $\mathbf{V}_{t's} = \phi_{t'}^s$ , and upper limit for  $\lambda_{t'}^d = \mathbf{V}_{t'b} \mathbf{V}_{t'd}^*$

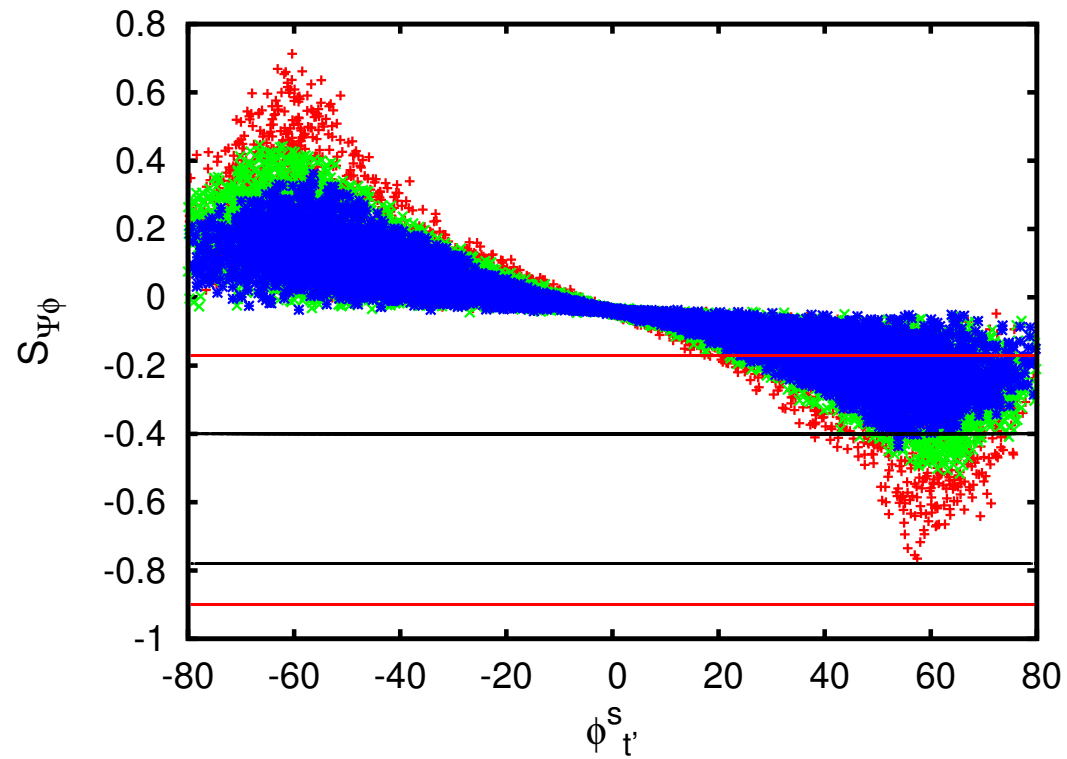
$m_{t'} (GeV)$	400	600
$\lambda_{t'}^s \times 10^2$	(0.08 - 1.4)	(0.05 - 0.6)
$\phi_{t'}^s$	-80 $\rightarrow$ 80	-80 $\rightarrow$ 80
$\lambda_{t'}^d \times 10^3$	< 2.0	< 1.6

Upper bound on the SM4 coupling reducing with the mass of fourth family...



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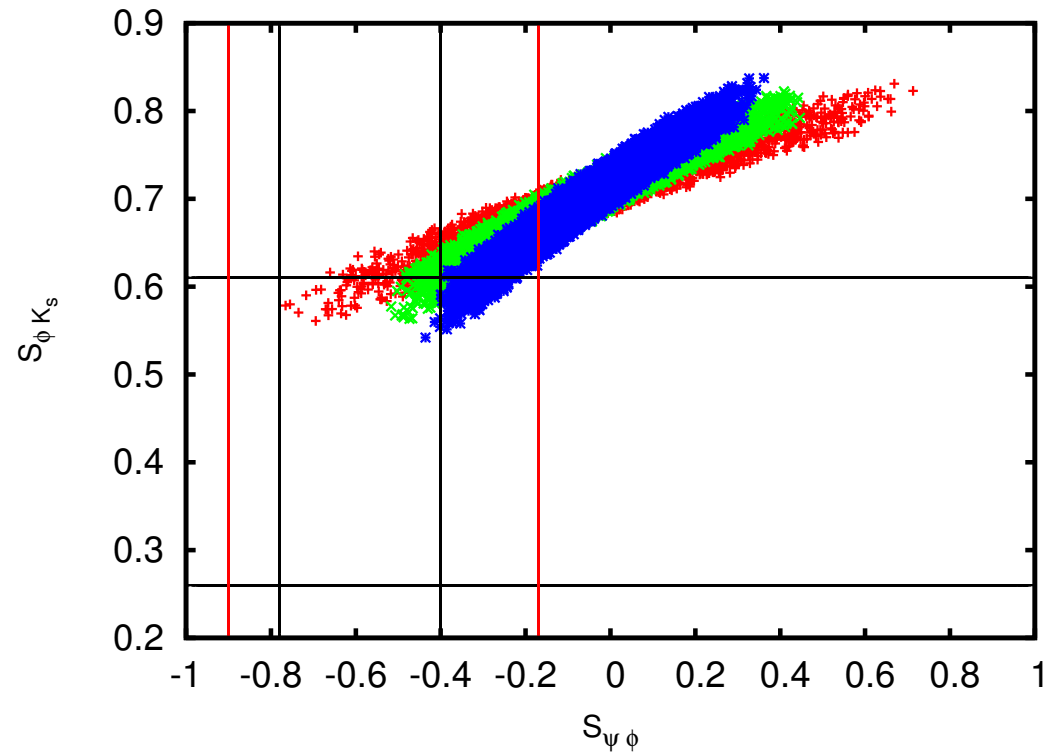
Variation of  $S_{\psi\phi}$  with the new phase  $\phi_{t'}^s$



Present experimental bound on  $S_{\psi\phi}$  allows positive range for  $\phi_{t'}^s$

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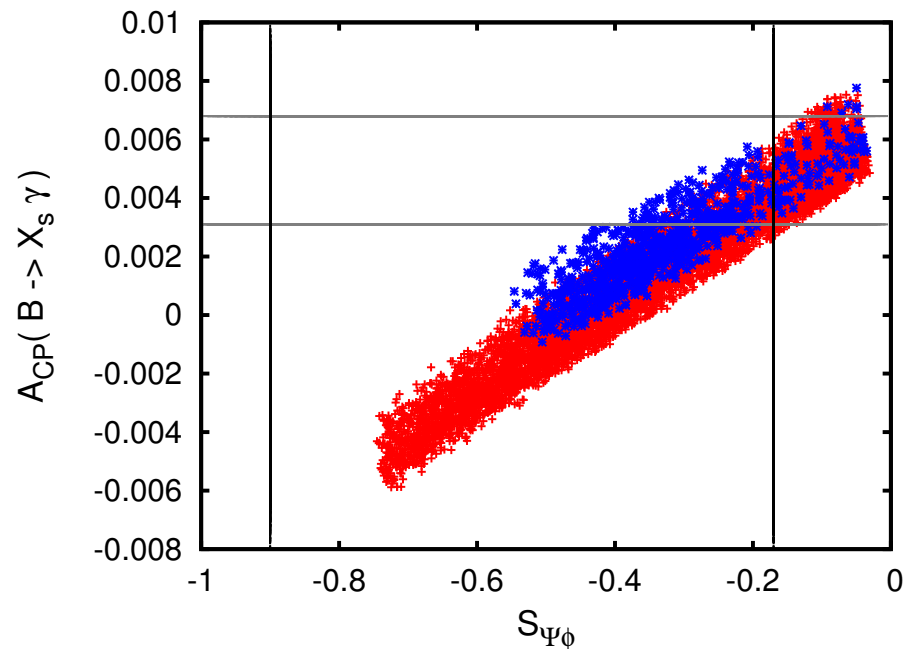
Correlation between  $S_{\psi\phi}$  with  $S_{\psi K_s}$



correlation suggests that  $m_{t'} \approx (400 - 600) \text{ GeV}$  is more favoured

Correlation between  $S_{\psi\phi}$  and  $CP$  asymmetry in  $B \rightarrow X_s \gamma$

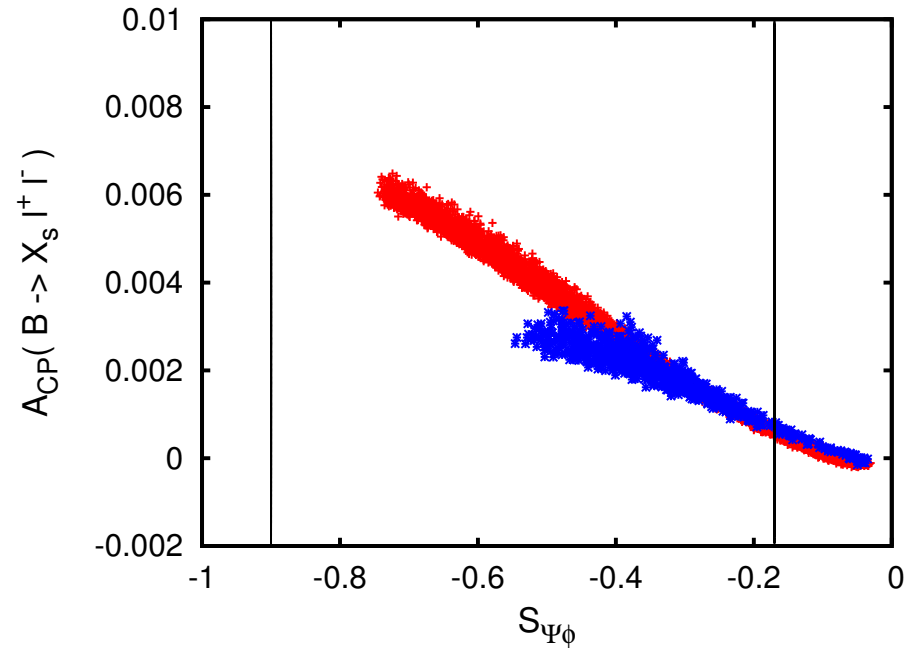
A. Soni, A. K. Alok, A. Giri, R. Mohanta and SN, in preparation



$$A_{CP}^{B \rightarrow X_s \gamma} \simeq \frac{10^{-2}}{|C_7^{\text{tot}}(m_b)|^2} \left\{ -1.82 \text{Im} [C_7^{\text{new}}] + 1.72 \text{Im} [C_8^{\text{new}}] - 4.46 \text{Im} [C_8^{\text{new}} C_7^{\text{new}*}] \right. \\ \left. + 3.21 \text{Im} [\epsilon_s (1 - 2.18 C_7^{\text{new}*} - 0.26 C_8^{\text{new}*})] \right\}$$

$A_{CP}$  is reducing with SM4 coupling.....

Correlation between  $S_{\psi\phi}$  and  $CP$  asymmetry in  $B \rightarrow X_s l^+ l^-$



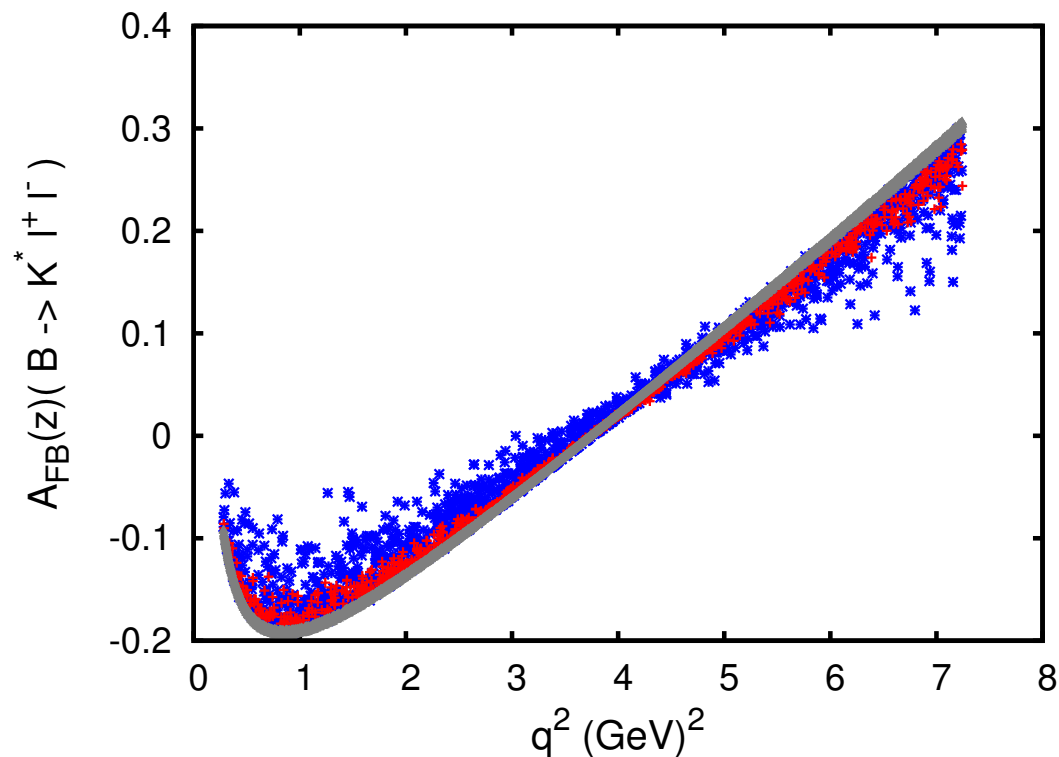
$$A_{CP} \propto \left[ \text{Im}(\lambda_{tu}^s) \left\{ 2(1+2z)\text{Im}(\xi_1\xi_2^*) - 12C_7\text{Im}(\xi_2) \right\} 2 \left( 1 + \frac{2t^2}{z} \right) + X_{im} \left\{ (1+2z)C_9^{t'} + 6C_7^{t'} \right\} \right]$$

$$\text{SM3} \Rightarrow \text{Im}(\xi_2) \propto \text{Im} \left( g(\hat{m}_u, z) - g(\hat{m}_c, z) \right) \approx \text{zero}$$

$$\text{SM4} \Rightarrow \text{Im}(\xi_2) \text{ unaffected but } \xi_1 \text{ increases}$$

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Forward Backward asymmetry as a function of  $q^2$  (low  $q^2$  region), grey region corresponds to the SM prediction

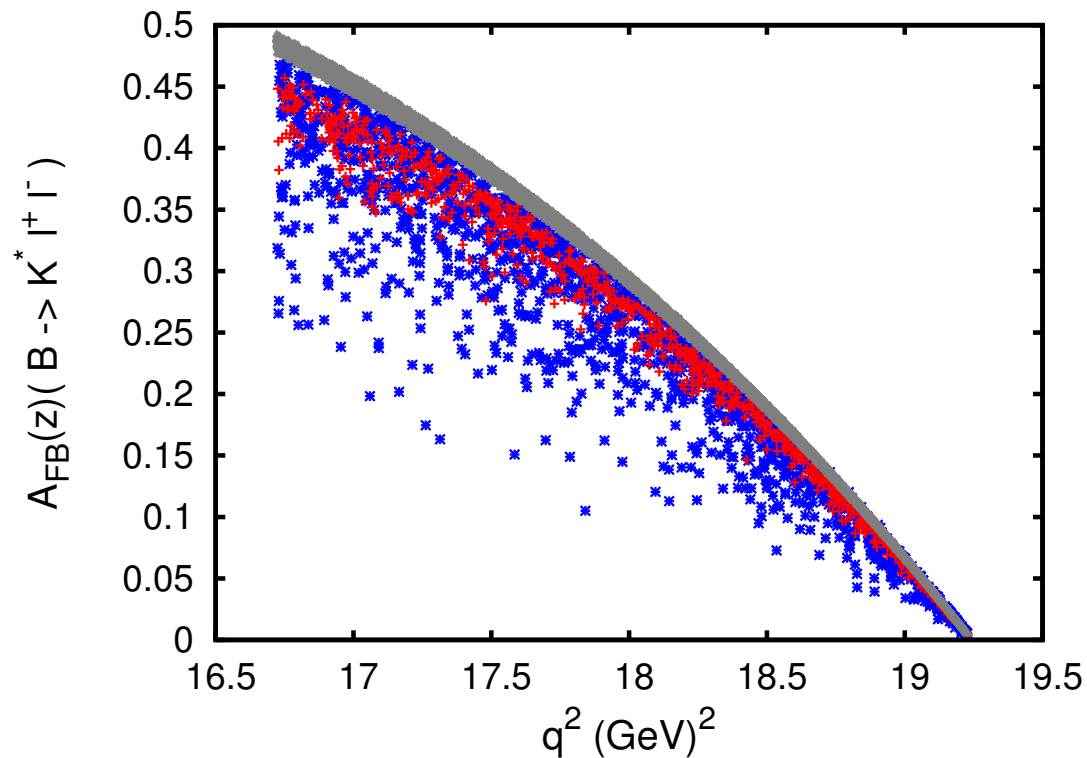


Small deviation from SM3 can be observed for  $q^2$  in between  $0.6 \rightarrow 1.0$ ....more prominent for

$m_{t'} = 600 \text{ GeV}$

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Forward Backward asymmetry as a function of  $q^2$  (high  $q^2$  region)

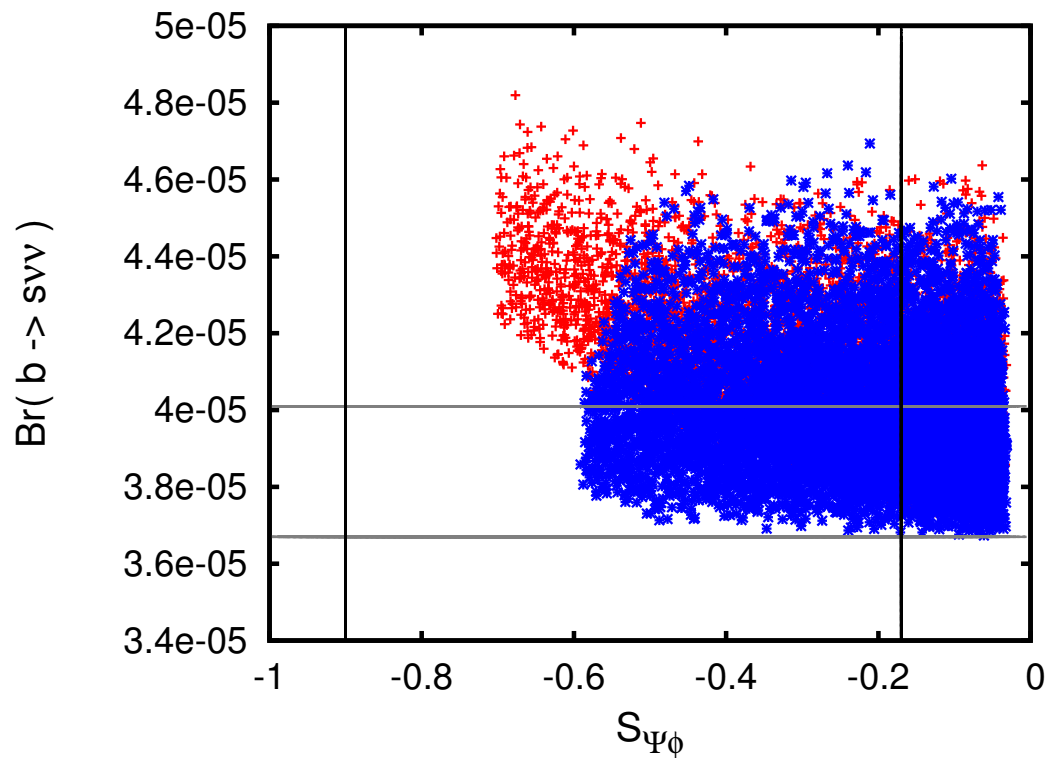


For  $q^2$  in between **16.0**  $\rightarrow$  **18.0**, 10% to 30% deviation from SM3 prediction could be observed for  $m_{t'} = 400$  and  $600$  GeV respectively

$q^2 (\text{GeV}^2/c^2)$	$A_{FB}$			
	exp	SM	$m'_t = 400 \text{ GeV}$	$m'_t = 600 \text{ GeV}$
0.6 – 1.0	$0.47^{+0.26}_{-0.33}$	$-0.18 \rightarrow -0.19$	$-0.13 \rightarrow -0.19$	$-0.08 \rightarrow -0.19$
1.0 – 6.0	$0.26^{+0.28}_{-0.31}$	$-0.2 \rightarrow 0.2$	$-0.2 \rightarrow 0.2$	$-0.2 \rightarrow 0.2$
6.0 – 8.0	$0.45^{+0.21}_{-0.26}$	$0.19 \rightarrow 0.30$	$0.17 \rightarrow 0.28$	$0.11 \rightarrow 0.30$
16.5 – 18.0	$0.66^{+0.12}_{-0.16}$	$0.28 \rightarrow 0.49$	$0.25 \rightarrow 0.45$	$0.15 \rightarrow 0.47$
18.0 – 19.5	For $q^2 > 16$	$0.003 \rightarrow 0.30$	$0.003 \rightarrow 0.27$	$0.003 \rightarrow 0.28$

SM4 predicted ranges are consistent with the present experimental bound

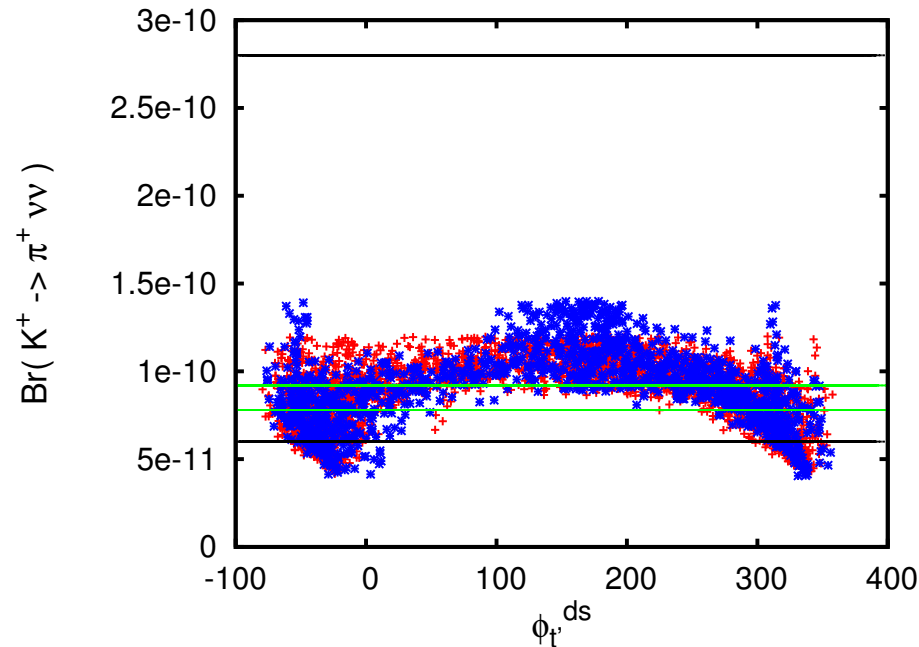
Correlation between  $S_{\psi\phi}$  and  $Br(\mathbf{b} \rightarrow s\nu\bar{\nu})$ , grey horizontal lines represent the SM limit



$$Br(B \rightarrow X_s \nu \bar{\nu}) \propto \frac{\alpha^2}{2\pi^2 \sin^4 \Theta_W} \left| V_{tb}^* V_{ts} X_0(x_t) \right|^2 \left| 1 + \frac{V_{t'b}^* V_{t's}}{V_{tb}^* V_{ts}} \frac{X_0(x_{t'})}{X_0(x_t)} \right|^2$$



Variation of  $\mathbf{Br}(K^+ \rightarrow \pi^+ \nu \bar{\nu})$  with  $\phi_{t'}^{\mathbf{ds}} = \phi_{t'}^{\mathbf{d}} - \phi_{t'}^{\mathbf{s}}$ , green horizontal lines correspond to SM limit

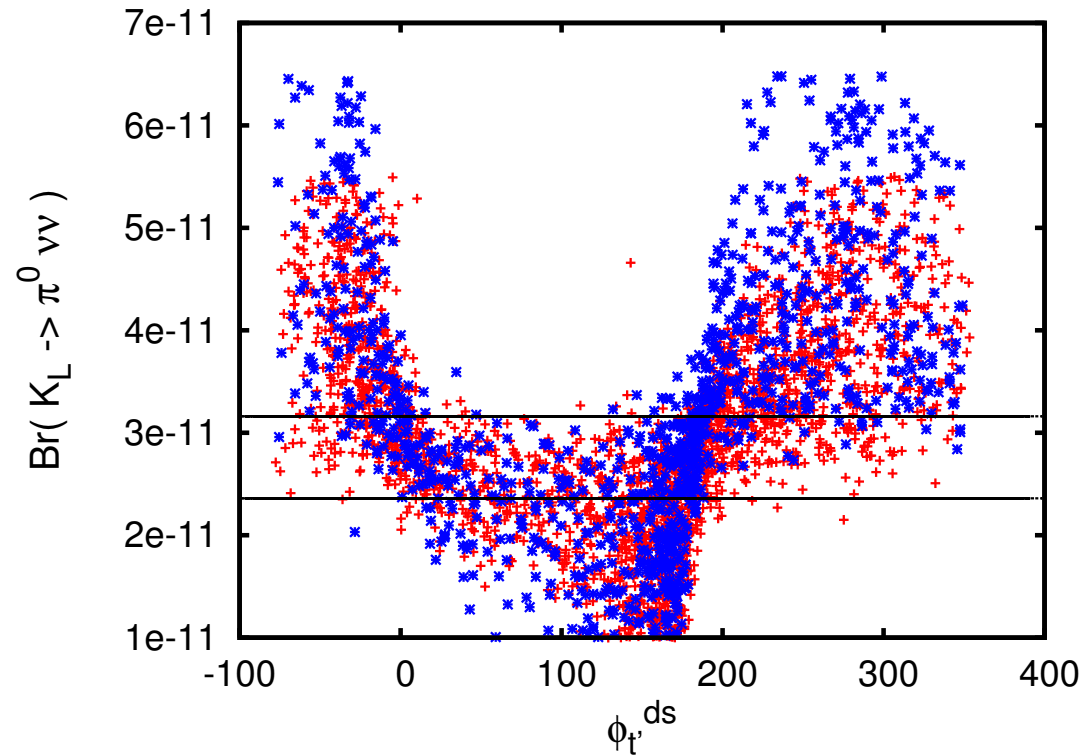


$$Br(K^+ \rightarrow \pi^+ \nu \bar{\nu}) \propto$$

$$\left[ \left( \frac{\text{Im}\lambda_t}{\lambda^5} X(x_t) + \frac{\text{Im}\lambda_{t'}}{\lambda^5} X(x_{t'}) \right)^2 + \left( \frac{\text{Re}\lambda_c}{\lambda} P_0(X) + \frac{\text{Re}\lambda_t}{\lambda^5} X(x_t) + \frac{\text{Re}\lambda_{t'}}{\lambda^5} X(x_{t'}) \right)^2 \right]$$

$$Br(K^+ \rightarrow \pi^+ \nu \bar{\nu}) = (4.0 \rightarrow 12.0) \times 10^{-11} \Rightarrow 1\sigma \text{ range in SM4}$$

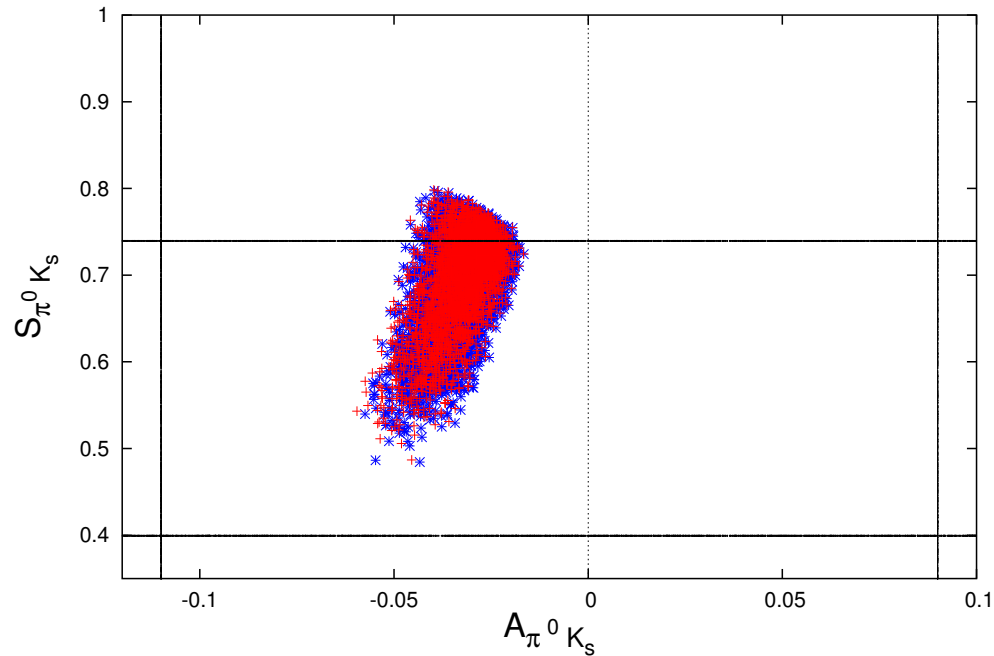
Variation of  $\text{Br}(K_L \rightarrow \pi^0 \nu \bar{\nu})$  with  $\phi_{t'}^{\text{ds}} = \phi_{t'}^{\text{d}} - \phi_{t'}^{\text{s}}$ , black horizontal lines correspond to SM limit



$$\text{Br}(K_L \rightarrow \pi^0 \nu \bar{\nu}) \propto \left[ \left( \frac{\text{Im} \lambda_t}{\lambda^5} X(x_t) + \frac{\text{Im} \lambda_{t'}}{\lambda^5} X(x_{t'}) \right)^2 \right]$$

$$\text{Br}(K_L \rightarrow \pi^0 \nu \bar{\nu}) = (1.0 \rightarrow 5.2) \times 10^{-11} \Rightarrow 1 \sigma \text{ range in SM4}$$

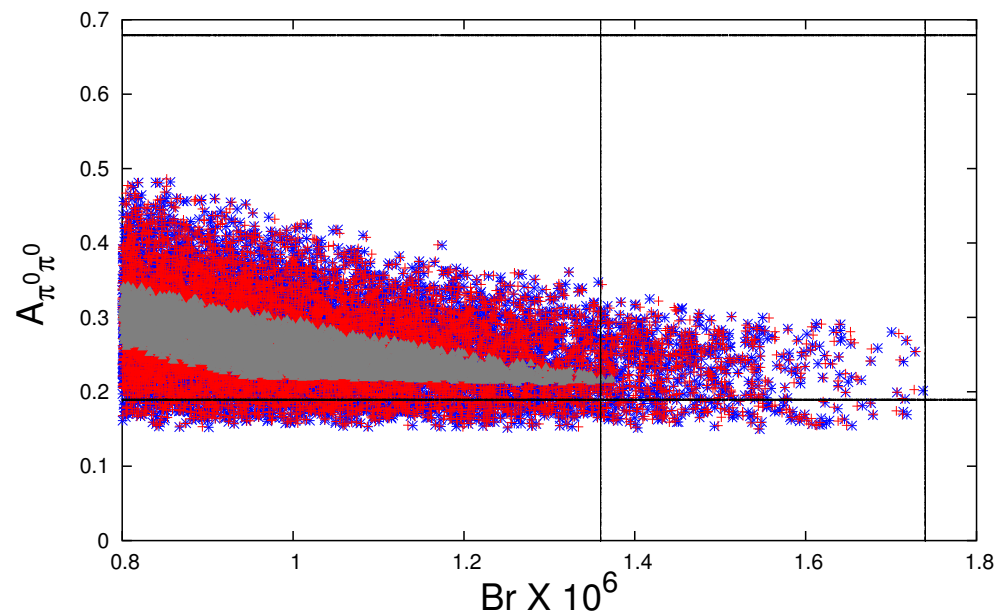
Correlation Plots between the mixing induced  $CP$  asymmetry  $S_{\pi^0 K_S}$  and the direct  $CP$  asymmetry  $A_{\pi^0 K_S}$



SM3 contribution  $\Rightarrow S_{\pi^0 K_S} \approx (0.7 \rightarrow 0.8)$  and  $A_{\pi^0 K_S} \approx (0 \rightarrow -0.05) \Rightarrow$  For  $S_{\pi^0 K_S}$  SM4 prediction have few  $\sigma$  deviations from that of SM3

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correlation plot between the direct *CP* asymmetry and the *CP* – averaged branching ratio for the  $B^0 \rightarrow \pi^0 \pi^0$  process



Large deviations for  $A_{\pi^0 \pi^0}^{0,0} \dots$

# Conclusions

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- Shortly after the discovery of third generation, a fourth generation was an obvious extension  $\Rightarrow$  Consistent with Electroweak precision test
- Heavier quarks with masses  $(400 - 600) \text{ GeV}$  can play a crucial role in dynamical electroweak-symmetry breaking...
- SM4 offers a simple explanation to many anomalies seen in  $B$  and  $B_s$  system..
- It is not clear whether we are able to reproduce the full size of the observed discrepancy in  $\Delta A_{CP}$   $\Rightarrow$  Could be a limitation of QCDF model..
- Several processes wherein SM4 causes large deviations from the expectations of SM3; for example  $\mathbf{b} \rightarrow s\nu\bar{\nu}$ ,  $\mathbf{K}_L \rightarrow \pi^0\nu\bar{\nu}$ ,  $\mathbf{A}_{\pi^0\pi^0}$ ,  $\mathbf{S}_{\psi\phi}$  *e.t.c*  $\Rightarrow$  may provide further indirect evidence for an additional family of quark

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Thank you !