## Summary: Highlights & Prospects

Amarjit Soni HET, BNL

"Second Workshop on Beyond 3 Generation Standard Model --- New Fermions at the Crossroads of Tevatron and LHC"

NTU, Taipei (Jan 14-16, 2010)

### Scope of the NTU Wrkshp

- Highlights of CERN Wrkshp: Unel
- DEWSB& barogenesis: Holdom, PQ, Gerhard, Kikukawa, Hashimoto, Bar-Shalom
- Possible cracks in CKM(SM): Punzi, Zieminska, Nandi, Lacker, Nakao, CYMa, Musy, FFLee
- EWPConstraints:Plehn, Ozcan, Antipin
- CKM constraints: Lenz, Lacker, CYMa, Nandi
- FCNC &CP: Arhrib, CHShen, Nebot, Bashiry
- Direct Searches (quarks):Whiteson,Bose,Plehn,Ozcan
- Direct Searches: Leptons: Antipin, Ozcan
- Dark Matter: Alwall
- Vector-like /isosinglet quarks: Branco, Nebot
- 4th Gen @SuperB/Bellell: Paoloni,PChang

# Sincere apologies for omissions, likely to be many despite considerable efforts

#### **Outline**

- The scope
- Prologue
- A Renaissance: may be joyous?
- Possible cracks in CKM?
- Prospects for settling the cracks.
- How much room is there?
- Indirect searches
- Direct searches
- DEWSB/Composite Higgs
- The Epilogue: Challenging Yogi

## UNEL

## Four Statements about the Fourth Generation

- 1)The fourth generation (SM4) is not excluded by EW precision data.
- 2)The SM4 addresses some of the currently open questions.
- 3)The SM4 can accommodate emerging possible hints of new physics.
- 4)LHC has the potential to discover or fully exclude the SM4.

arXiv:0904.4698 PMC Phys

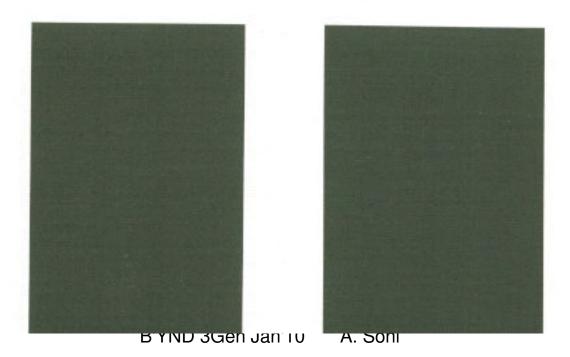
B. Holdom, W.S. Hou, T. Hurth, M. Mangano, S. Sultansoy, G. Ünel

- Once popular, an additional 4th generation was no longer "a la mode". Last workshop around late 1980s
  - ⇒misinformation disseminated by PDG contributed to this omission.

## The Fourth Family of Quarks and Leptons ~1987

Second International Symposium

Editors DAVID B. CLINE • AMARJIT SONI



#### Properties of a Composite Higgs Particle in a Dynamically Broken Formulation of the Standard Model

J. Carpenter, R. Norton, and S. Siegemund-Broka

Department of Physics, University of California, Los Angeles, California 90024

A. Soni

Brookhaven National Laboratory, Upton, New York 11973 (Received 1 September 1989) my 2 m

We suggest that dynamical breakdown of electroweak symmetry in an SU(3)×SU(2)×U(1) model of a single fermion family without fundamental scalars can reproduce the entire scalar sector of the standard model through fermion bound states and loop-induced couplings. The physical Higgs boson is essentially a  $t\bar{t}$  bound state. The top-quark and Higgs-boson masses are approximately 300 and 400 GeV, respectively. The strengths of the induced couplings among the scalars do not indicate a strongly interacting scalar sector.

 $(300)^{2}$   $(m_{t}-m_{b})$   $(m_{t}-m_{t})$  B'YND 3Gen Jan'10 A. Son

me ~500 Gev

## Origin of EWSB:Possibly a glorious renaissance

II. Strong dynamics {WEXD, TC, WTC....} Quarks/Leptons with O(500 GeV) masses make strong dynamics unavoidable.... The natural mass scale for Higgs like object becomes  $>> m_Z \sim$  susy higgs mass scale

### Top ~10 reasons

#1: udcstbgz

- 89->
- Instead of ~121 can do with 7
- Another 4 reasons:1234
- Another 3 reasons: e mu tau

OSCILL8

 $\sim 80 - 89$ 

### Possible cracks in CKM?

Related talks here by:

Punzi, Zieminska, Nandi, Lacker, Nakao, CYMa, Bashiry, Musy

Based also on Lunghi+AS

#### Summary of B-CP Anomalies (~'07-'09)

- Fitted ("SM-predicted") value of sin  $2\beta(\phi_1)$  vs directly measured, a) via golden tree decays
  - b) via loop decays
  - Dir CP in K+ $\pi$  vs K+ $\pi$ 0
  - Bs->ψφ (esp. significant since 1. Its theoretically very clean(Gold plated) II. It essentially follows from others...Consequently very important that Fermilab follows it up & clarifies it with very high priority).
- •\Each ~2 to 3.5 σ

also Bona Fal

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### Should also bear in mind

- Br (B -> T), 1,56 Weakest
- AFB (B ->K^\* I I)
- 2.56 Just BELLE +BABAR

• (g-2)

 $\sim 36$ 

#### **Lunghi + AS, arXiv:0903.5059**

$\epsilon_K$ , $\Delta M_{B_q}$ , $\epsilon_K$ , $\Delta M_{B_q}$ , $\epsilon_K$ , $\Delta M_{B_q}$ ,		in i	0.885±0.082 0.846±0.069 0.747±0.029
t→ces	tree		0.672±0.024
φK <sup>0</sup>			0.44+0.17
η'K <sup>0</sup>	penguin (elean)	1 - 1	$0.59\pm0.07$
(ø,7')K			0.57±0.065
$K_SK_SK_S$			0.74±0.17
π <sup>0</sup> Ε <sup>0</sup>			0.57+0.17
$\rho^0 K_\delta$		1 4	$0.63^{+0.17}_{-0.21}$
$\omega \mathbb{K}_S$	penguin (other) +		0.45±0.24
$f_0K_S$	,		0.62+0.11
$\pi^0\pi^0K_S$	-		-0.52±0.41
$\phi \pi^0 K_B$			0.97 + 0.03
K+K-K0		1-0-1	0.82±0.07
-1.5	-1.0 -0.5 0.0	0.5 1.	0 1.5

mode	$w/out V_{ub}$	with $V_{ub}$
$S_{\psi K_S}$	$2.4 \sigma$	$2.0 \ \sigma$
$S_{\phi K_S}$	$2.2 \sigma$	$1.8~\sigma$
$S_{\eta'K_S}$	$2.6 \sigma$	$2.1 \sigma$
$S_{(\phi-\eta')K_S}$	$2.9 \sigma$	$2.5 \ \sigma$

#### Lunghi+AS,arXiv.0707.0212

(Sin 2  $\beta$  = 0.78+-.04)

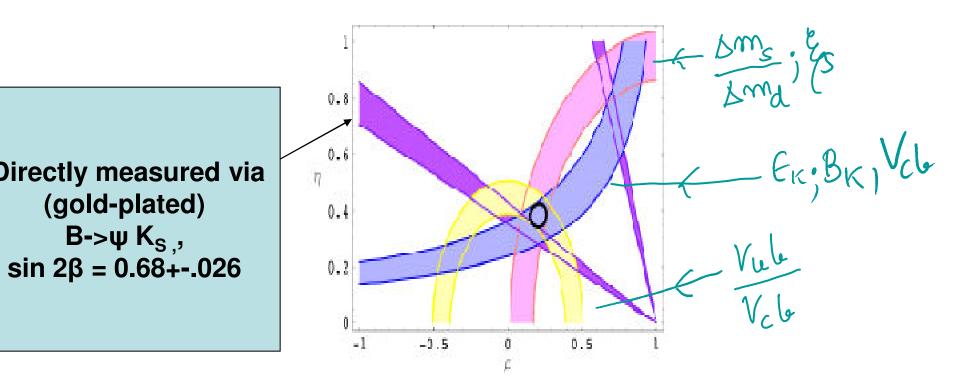
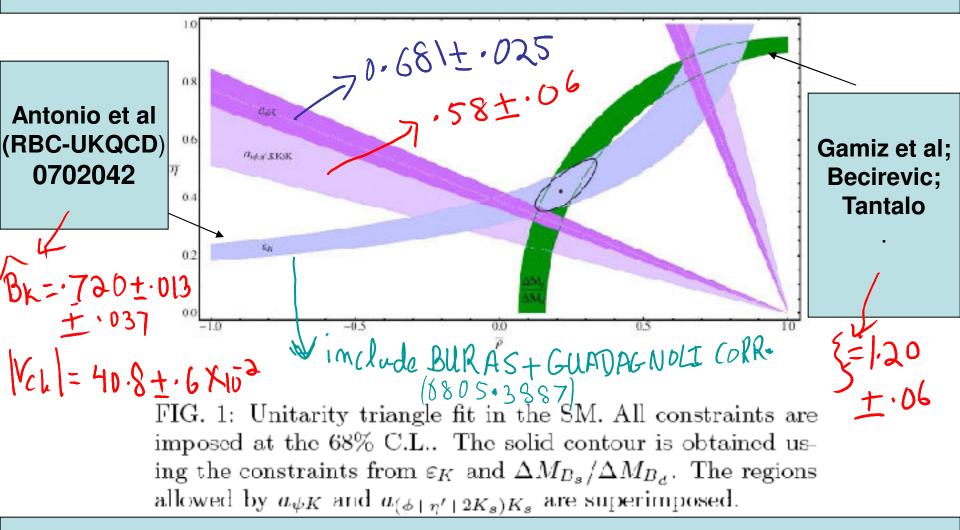


Figure 1: Unitarity triangle fit in the SM. The constraints from  $|V_{ub}/V_{cb}|$ ,  $\varepsilon_K$ ,  $\Delta M_{B_s}/\Delta M_{B_d}$  are included in the fit; the region allowed by  $a_{\psi K}$  is superimposed.

Important to Examine only DeltaF=2 observables:Leave out Vub sin 2  $\beta$  = 0.87+-.09{Lunghi+AS,hep-ph/08034340} (became possible only due significantly reduced error in  $B_{\kappa}$ )



2.1-2.7 σ- deviation from the directly measured values of sin 2 β requires careful follow-up

#### $B_K$ - Preliminary Results

#### C.Kelly for RBC-UKQCD @ LATTICE'09 Beijing, 7/26/09

▶ 24³ paper value (*Phys.Rev.Lett.*100:032001,2008): 
$$_{R}$$
  $_{R}$   $_{$ 

New value:

$$B_K^{\overline{\mathrm{MS}}}(2\,\mathrm{GeV}) = 0.537(19)$$
 NEW (PRELIMINARY)

Source	24 <sup>3</sup> Mag	New Mag.	%	Comment	
stat	0.010	0.006	1.1	Stat err. from sim. fit	
ChPT	0.010	0.013	2.4	analytic expansion $\sim NLO^2$ est.	
FV	0.005	??	??	NLO ChPT FV corrections	
NPR	0.013	0.013	2.4	Comb. stat + sys err.	
				(mult. ex + non-ex MOM schemes)	
Scaling	0.021	0	0	Scaling inc. in fits	
Unphys. $m_h$	0.005	0	0	Reweighting	
Total	0.030	0.019	3.5	Preliminary	

Aubin, Laiho, Vande Water arXiv:0905.3947

$$B_K^{\overline{\rm MS},{\rm NDR}}(2{\rm GeV}) = 0.527(6)(20)$$





### Brief (~25 years) History of B<sub>K</sub>

, ~'83 DGH use K+ lifetime + LOChPT + SU(3)->  $B_K \sim 0.33...$  no error estimate, no scale dependence....

UNCONTROLLABLE APPROXIMATION =>

~'84 Lattice method for WME born...many attempts & improvements for B<sub>K</sub> evaluations

~'98 JLQCDstaggered B<sub>K</sub> (2GeV)= 0.628(42)quenched(~110).

- ~'97 1st B<sub>K</sub> with DWQ(T.Blum&A.S),0.628(47) quenched.
- ~'01 RBC B<sub>K</sub> with DWQ, quenched=0.532(11) quenched
- $\sim$ '05 RBC, nf=2, dyn. DWQ, B<sub>K</sub> =0.563(21)(39)(30)
- $\sim$  '06 Gimnez et al (HPQCD; stagg.) 2+1,  $B_K = 0.618(18)(19)(30)(130)$
- ~07,RBC-UKQCD DWQ 2+1 .....0.524(10)(28)

DWQ lower  $B_K$  -> requiring larger CKM-phase

- ~'08 Target 2+1 dyn. DWQ, B<sub>K</sub> with total error 5%
- ~'09 : error below ~4%

**NO LONGER THE LIMITING ERROR!!** 

BBG-69N PLB'88 70+.07

.720(13)(37)

116

time has constrained KM parameters well enough that a knowledge of the B parameter can be used to put a nontrivial bound on the mass of the top quark. Thus, for example, with E = .33 the standard model with three generations requires  $n_{\rm top} \gtrsim 40$  GeV for it to successfully account for the experimentally observed value of  $\varepsilon$ .

Left-right symmetric theories constitute a popular means for extending the standard nodel. The  $K^0 \overline{K}^0$  mean difference can be used to learn about the lower bound on the mass scale of such theories. These bounds depend linearly on  $N_{LR} \equiv \langle K^0 | \overline{s}^\gamma_\mu (1-\gamma_5) d | \overline{s}^\gamma_\nu (1+\gamma_5) d | \overline{K}^0 \rangle$ , and the current bound (1.6 TeV for the mass of the right handed gauge bosons) was deduced by assuming vacuum saturation. The vill therefore be useful to calculate  $M_{LR}$  on the lattice.

The matrix elements of some penguin operators control in the standard model another CF violation parameter, namely  $\epsilon'/\epsilon$ .  $^{6,8)}$  Indeed efforts are now undervay for an improved measurement of this important parameter. In the absence of a reliable calculation for these parameters, the experimental measurements, often achieved at tremendous effort, cannot be used effectively for constraining the theory. It is therefore clearly important to see how far one can go with MC techniques in alleviating this old but very difficult problem of non-leptonic weak decays.



#### **Courtesy: Tom Browder**

Critical Role of the B factories in the verification of the KM hypothesis was recognized and cited by the Nobel Foundation

A single irreducible phase in the weak interaction matrix accounts for most of the CPV observed in kaons and B's.

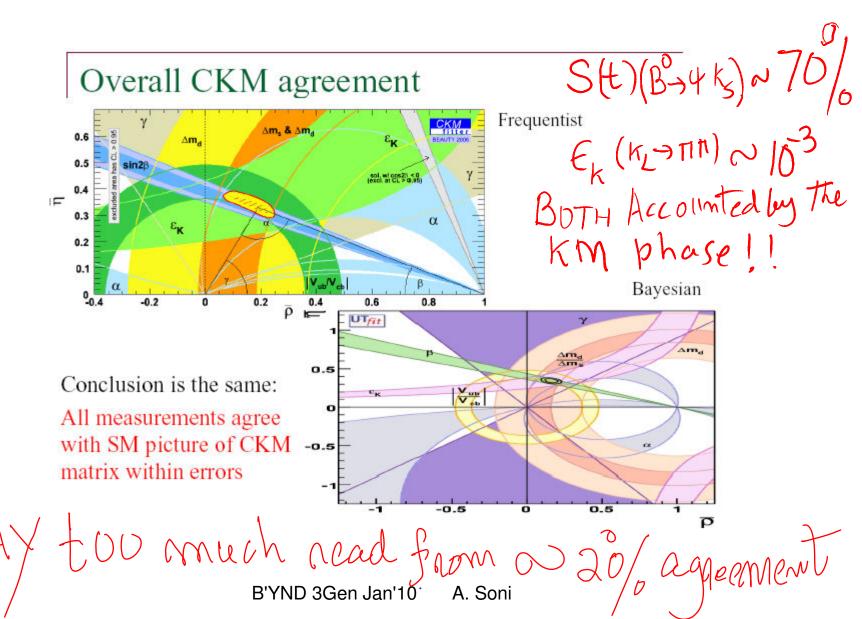


CP violating effects in the B sector are O(1) rather than  $O(10^{-3})$  as in the kaon system.

#### **Lunghi + AS, arXiv:0903.5059**

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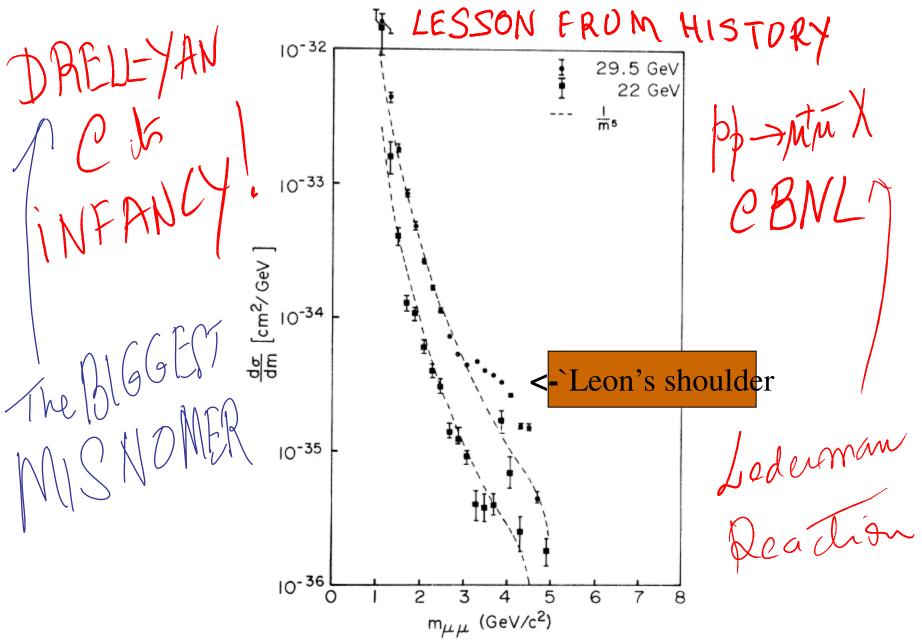


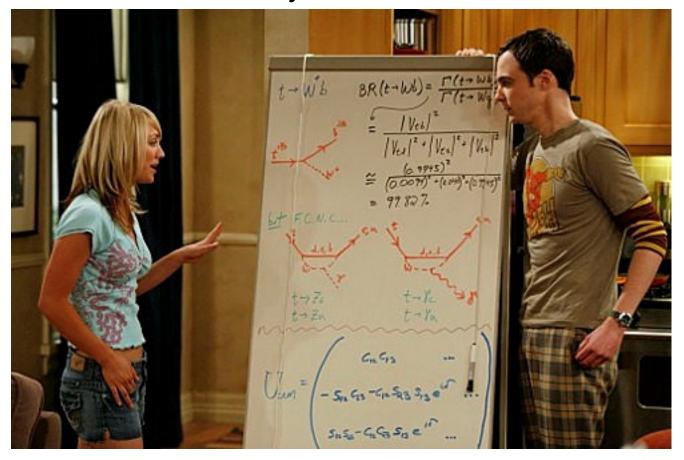
FIG. 15. Experimental cross sections at two energies compared with a simple  $1/m^5$  continuum.



Leon Lederman

Deserves a 2<sup>nd</sup> NP for inventing the reaction: junk + junk -> gems +X which has led to the discoveries of J, Upsilon, W, Z, top,...and remains the most powerful exploratory tool in our arsenal!!

### LHC/Super B factory synergy discussion on US TV comedy



D. Saltzberg,ScienceAdvisor

CBS, "Big Bang Theory" averages 9 million viewers per episode.

Courtesy Tom Browder

#### A lesson from history (I)

"A special search at Dubna was carried out by E. Okonov and his group. They did not find a single  $K_1 \rightarrow \pi^+ \pi^-$  event among 600 decays into charged particles [12] (Anikira et al., JETP 1962). At that stage the search was terminated by the administration of the Lab. The group was unlucky."

-Lev Okun, "The Vacuum as Seen from Moscow"

1964: BF=  $2 \times 10^{-3}$  Character Senson, FITCH A failure of imagination ? Lack of patience ?  $4 \times 10^{-4}$ 

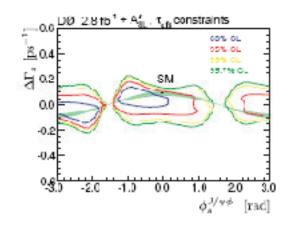
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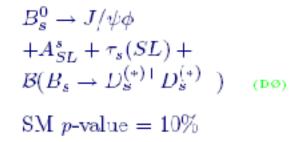
#### ZIEMINSKA

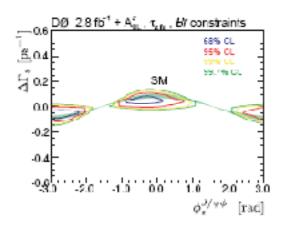


#### $\Delta\Gamma_s$ vs $\phi_s^{J/\psi\phi}$ with Additional Constraints

$$B_s^0 o J/\psi \phi$$
 (DØ) 
$$+A_{SL}^s + \tau_s(SL) \quad \text{(World Average)}$$
 SM  $p\text{-value} = 12\%$ 

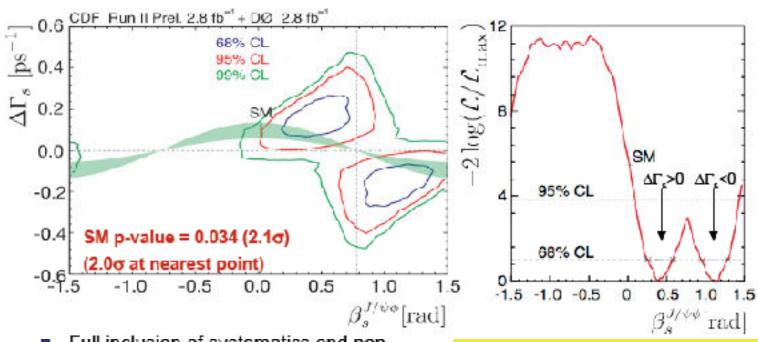






#### Combined Tevatron result

[http://tevbwg fnal gov]



- Full inclusion of systematics and non-Gaussian effects. No external constraints.
- Compared to HFAC 2008: Larger CDF sample + Better accounting for tails → same level of SM agreement.

β<sub>s</sub>J/**v**Φ **range:** [0.27,0.59] U [0.97,1.30] @68% [0.10,1.42] @95% 1-D p-value for SM= 0 020 (2 3σ)

## Settling the fate of the cracks

# Highest priority item in the particle physics community for now(at least for flavor physics)

### Other items needing priority

- From the lattice B<sub>K</sub> with other (than DWQ) methods such as overlap and/or other groups
- $\xi_s$  SU3 breaking ratio for  $\Delta m_s/\Delta m_d$  with **DWQ**/overlap
- f<sub>B</sub> with DWQ/overlap
   MORE ACCURATE A<sub>FB</sub> from Expt
- · More accurate Br(B-> > ∫ nom expt.

## WHODUNIT?

A. Soni

### Honest answer &

- Don't really know (too many possibilities...)
- But theoretically the most interesting possibility is that we may be witnessing RANDALL-SUNDRUM 99 Dawning of the age of

"Warped Flavordynamics"

(Gives a simultaneous resoltution to Planck-EW hierarchy AND the flavor puzzle [through fermion "geography"] ).

Tree level FCNC are severely suppressed through "RS-GIM" mechanism (Agashe, Perez & AS'04)

Thus remarkably RS-leads to lowering of  $\Lambda_{flavor}$  from ~1000 TeV to < 20 TeV (possibly just a few TeV if you allow small amount of tuning) NEUBERTOUR 108
BURAS et d GROSSMAN-NEUBERT, 99 GERGHETTA-POMAROL

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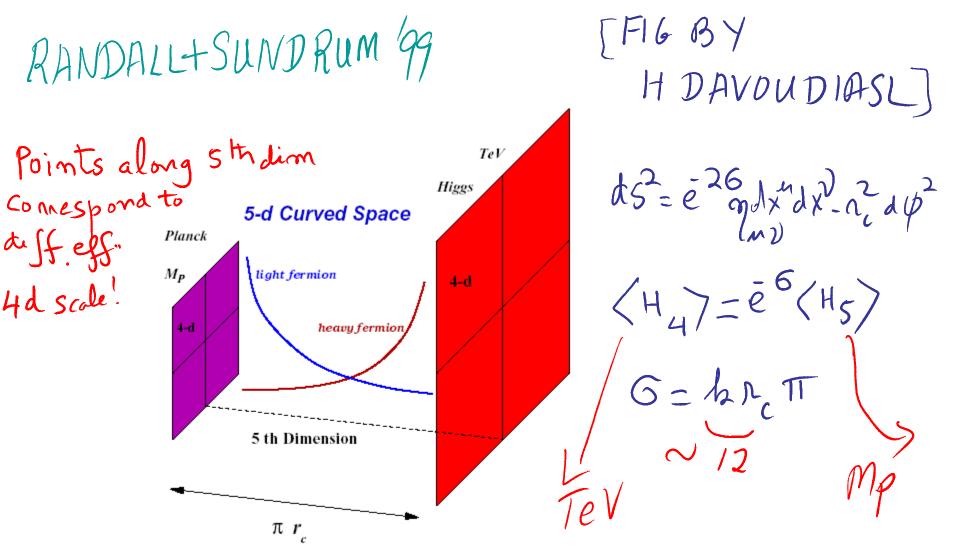


Figure 1: Warped geometry with flavor from fermion localization. The Higgs field resides on the TeV-brane. The size of the extra dimension is  $\pi r_c \sim M_P^{-1}$ .

#### Simultaneous resolution to hierarchy and flavor puzzles

#### Fermion "geography" (localization) naturally explains:

Grossman&Neubert; Gherghetta&Pomarol; Davoudiasl, Hewett & Rizzo

- Why they are light (or heavy)
- FCNC for light quarks are severely suppressed
- RS-GIM MECHANISM (Agashe, Perez, AS'04)
  flavor changing transitions though at the tree
  level (resulting from rotation from interaction
  to mass basis)are suppressed roughly to the
  same level as the loop in SM
- Most flavor violations are driven by the top

## Contrasting B-Factory Signals from WEXD with those from SM

Agashe,Perez &AS, PRL'04
(Then for simplicity assumed Bd-mixing is SM)
O(1) uncertainties stressed. NOTE these are genuine PREDICTIONS

	$\Delta m_{B_s}$	$S_{B_s  o \psi \phi}$	$S_{B_d  o \phi K_s}$	$Br[b \rightarrow sl^+l^-]$	$S_{B_{d,s} \to K^*, \phi \gamma}$	$S_{B_{d,s}  o  ho,K^*\gamma}$	
RS1	$\Delta m_{B_s}^{\rm SM}[1+O(1)]$	0(1)	$\sin 2\beta \pm O(.2)$	$Br^{\text{SM}}[1+O(1)]$	\ \ /	O(1)	73
SM	$\Delta m_{B_s}^{\rm SM}$	$\lambda_c^2$	$\sin 2\beta$	$Br^{\rm SM}$	$\frac{m_s}{m_b} \left( \sin 2\beta, \lambda_c^2 \right)$	$\frac{m_d}{m_b} \left( \lambda_c^2, \sin 2\beta \right)$	

Recently many very nice extensions (Buras, Falkowski, Perez, Weiler, Neubert) et al

## EXTREMELY INTERESTING SUBTELTY of warped models

- Maldacena conjecture
- "Warped Flavordynamics" is DUAL to strong dynamics->

SIMPLEST 4d Explanation of B,Bs CP anomalies is SM4: SM with 4 generations

# Bundman & Rold axiv: 0710.0623

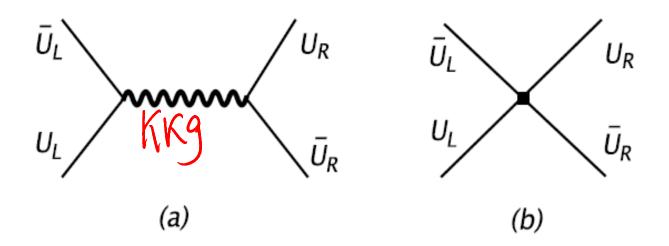


Figure 1: Two contributions to four-fermion interactions of the up-type fourth-generation quark: (a) from the interactions with a KK gluon; (b) from the four-fermion interactions induced by the bulk operators of (2.11).

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### HEAVY fourth family cannot co-exist with light Higgs

### Higgs description loses meaning

- $m_{t',b'} \approx 600 \text{ GeV}$  close to the unitarity upper bound
- Goldstone bosons of electroweak symmetry breaking couple strongly to t', b', so strong interactions unitarize WW scattering

### Higgs is no longer needed

- strong interactions can also be responsible for the Goldstone bosons
- a fermion condensate replaces the Higgs vev
- so where do quark and lepton masses come from?

what causes 
$$\langle \overline{\Psi}\Psi \rangle = \Lambda_{ew}^3$$
?

- new strong UNBROKEN gauge interaction—technicolor or
- new strong BROKEN gauge interaction—similar to original NJL model
- for the latter we can identify Ψ with the fourth family

$$\frac{1}{\Lambda_{fl}^2} \overline{\Psi} \Psi \overline{\psi} \psi \qquad \qquad \frac{1}{{\Lambda_{fl}'}^2} \overline{\Psi} \Psi \overline{\Psi} \Psi$$

$$\Lambda_{fl} \gg \Lambda'_{fl} \approx \Lambda_{ew} \approx 1 \text{ TeV}$$

BDB HOLDSM

• flavor gauge interactions, partially broken at  $\Lambda_{fl}$ , with a remnant broken at  $\Lambda'_{fl}$ 

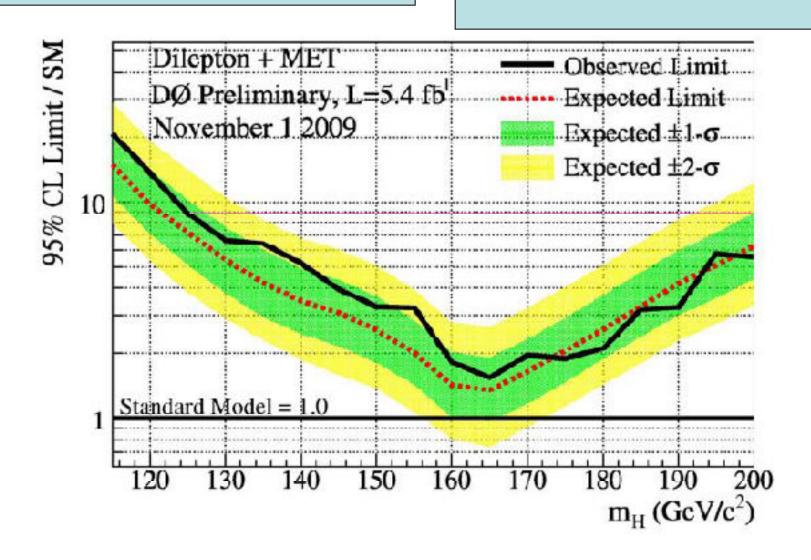
# Heavy 4th Generation's scenario

### What can a heavy 4th generation do?

- A heavy Higgs-4th Yukawa system can give rise to condensates and bound states of 4th generation fermions ⇒ Implications on the vacuum structure of the SM and the number of Higgs doublets (a mixture of fundamental and composites)
- The appearance of quasi-fixed points in the 2-loop approximation at a scale  $\Lambda_{FP} \sim O(TeV)$  hints at a possible scale-invariant theory above  $\Lambda_{FP}$ .

#### Tulika Bose, D0

#### In SM4 mH > 200 GeV



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# The "3-prong composite solution" to the SM4

(figure taken from: Chanowitz, Phys.Rev.D79:113008,2009)

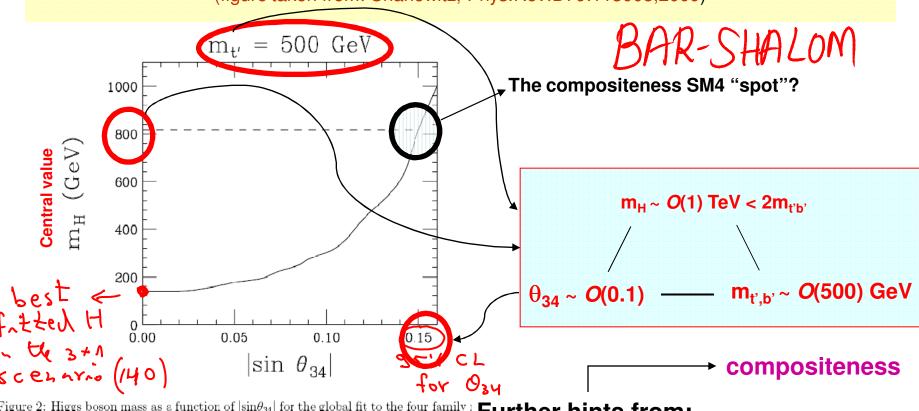
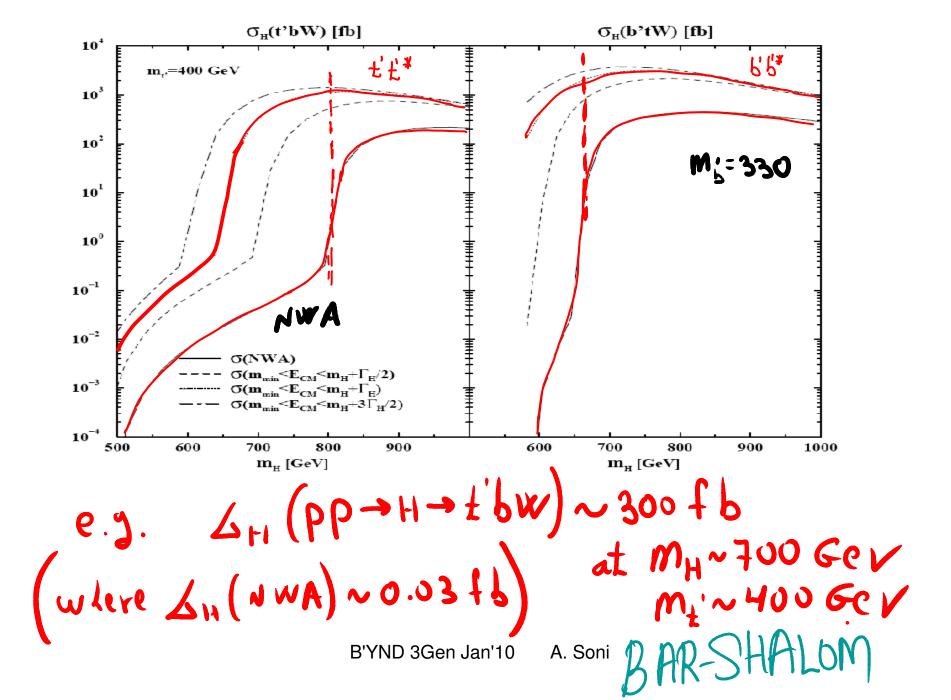


Figure 2: Higgs boson mass as a function of  $|\sin\theta_{34}|$  for the global fit to the four family: Further hints from: with  $m_{t'} = 500$  GeV. The horizontal line indicates the 95% confidence interval for  $|\sin\theta_{34}|$ 

$$m_{t'} - m_{b'} \approx \left(1 + \frac{\ln \frac{m_H}{115 \text{ GeV}}}{5}\right) \times 50 \text{ GeV}$$
 Direct searches

Flavor data

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LENZ, LACKER, CYMA, NANDI...

(Chanocottz)

# HOW much room is there?

### UT angles: "γ"

Problem within 4<sup>th</sup> generation scenario:

$$\gamma = Arg \left( -\frac{V_{ud}V_{ub}^*}{V_{ud}V_{ub}^*} \right)$$
 but the following CKM element combinations appear in weak decay or

also K-K mixing

1) GLW

a) 
$$D^0 \rightarrow K^+ K^ -Arg \left( \frac{V_{us}V_{ob}^*}{V_{cs}V_{ub}^*} \frac{V_{cs}V_{us}^*}{V_{us}V_{cs}^*} \right) = \gamma - \tau \bigcirc Arg \left( \frac{V_{ud}V_{us}^*}{V_{cd}V_{us}^*} \frac{V_{ud}V_{us}^*}{V_{cd}V_{us}^*} \right)$$
b)  $D^0 \rightarrow \pi^+ \pi$   $-Arg \left( \frac{V_{us}V_{ob}^*}{V_{cs}V_{ub}^*} \frac{V_{cd}V_{ud}^*}{V_{ud}V_{cd}^*} \right) = \gamma - \tau \bigcirc Arg \left( \frac{V_{ud}V_{us}^*}{V_{cd}V_{us}^*} \frac{SM: \pi \text{ to very good approximation}}{V_{cd}V_{us}^*} \right)$ 
c)  $D^0 \rightarrow K_S \pi^0 \text{ etc. } -Arg \left( \frac{V_{us}V_{ob}^*}{V_{cs}V_{ub}^*} \frac{V_{cs}V_{us}^*}{V_{us}V_{cs}^*} \right) = \gamma - \pi \bigcirc Arg \left( \frac{V_{ud}V_{us}^*}{V_{cd}V_{us}^*} \frac{SM: \pi \text{ to very good approximation}}{V_{cd}V_{us}^*} \right)$ 

b) 
$$D^0 \rightarrow \pi^+ \pi$$

$$Arg\left(\frac{V_{us}V_{cb}^{*}}{V_{cs}V_{ub}^{*}}\frac{V_{cd}V_{ud}^{*}}{V_{ud}V_{cd}^{*}}\right) = \gamma - r \bigoplus Arg$$

$$rg\left(\frac{V_{\omega}V_{\omega}^{*}}{V_{\omega}V_{\omega}^{*}}\frac{V_{\omega}V_{\omega}^{*}}{V_{\omega}V_{\omega}^{*}}\right) = \gamma - \pi \bigcirc Arg$$

WORK IN progress

2) ADS: 
$$D^0 \rightarrow K^+ \pi - Arg \left( \frac{V_{ux} V_{ch}^* V_{ud} V_{ux}^*}{V_{ud} V_{ux}^* V_{ud} V_{cs}^*} \right) = \gamma - \pi$$

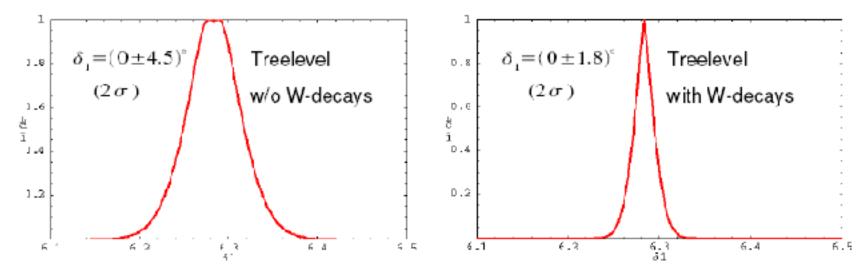
Dalitz

Analogous possibilities depending on final state in Dalitz plane.

### UT angles: "γ"

Define 
$$\delta_1 = Arg \left( \frac{V_{ud} V_{us}^*}{V_{cd} V_{cs}^*} \right) + \pi$$
 => Fit for deviation from  $2\pi = 0$  within  $4^{th}$  generation scenario

(Already considered in Kurimoto & Tomita, Prog.Theor.Phys. 98 (1997) 967; sign error in eq. (18))



=> γ-extraction in ADS and GLW within 4th generation scenario

like in 3-generation scenario to a very good approximation (1° systematics @ 1<sub>o</sub>) Should be analogous in Dalitz plot analysis but not thought through in detail yet

#### **In Progress**



#### Mixing with the 4th Family XIV

#### Result:

As expected: most points belong to SM3 like parameters

Unexpected: large mixing not yet excluded

 $V_{td}, V_{ts}, V_{tb}$  can differ considerably from SM3-fit values

LENZ

#### Ultraconservative allowed ranges

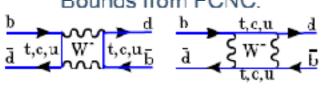
$$\theta_{14} \leq 0.04 \approx 1.27\lambda^{2}$$
 $\theta_{24} \leq 0.25 \approx 0.9\lambda^{1}$ 
 $\theta_{34} < 0.8 \approx 0.8\lambda^{0}$ 
 $\delta_{14}, \delta_{24} \leq 2\pi$ 

No "nice" Wolfenstein expansion possible



### Mixing with the 4th Family VII





$$\Delta_X := \frac{X_{SM4}}{X_{SM3}}$$

■ K-Mixing:

$$Re(\Delta_K) = 1 \pm 0.5 \, (0.25)$$

$$Re(\Delta_K) = 1 \pm 0.5 (0.25)$$
  $Im(\Delta_K) = 0 \pm 0.3 (0.15)$ 

 $\blacksquare$   $B_d$ -Mixing:

$$|\Delta_{B_d}| = 1 \pm 0.3 (0.1)$$

$$|\Delta_{B_d}| = 1 \pm 0.3 \, (0.1) \qquad Arg(\Delta_{B_d}) = 0 \pm 10^\circ \, (5^\circ)$$

■ B<sub>s</sub>-Mixing:

$$|\Delta_{B_s}| = 1 \pm 0.3 \, (0.1)$$
  $Arg(\Delta_{B_s}) = free$ 

$$Arg(\Delta_{B_s}) = free$$

 $\blacksquare b \rightarrow s\gamma$ 

$$\Delta_{b \to s \gamma} = 1 \pm 0.15 \, (0.07)$$

classic examples of evasion ob decoupling

# Early (~87-88) studies on 4th gen.

- Hou, Willey and AS, PRL (88)..b->s I I...
- Hou, AS, Steger, PRL 87.....b-> s g
- Hou, AS, Steger, PLB 87
   4X4 mixing matrix and b -> s gamma

mportance of B-decays for searching 4<sup>th</sup> gen. due to non-decou emphasized long ago

### CYMA (+ HOU

# Use parametrization matrix [Hou Soni and Steger PLB 192 (1987)

1st Row

$$\begin{split} V_{ud} &= -c_{13}\,s_{14}\,s_{12}\,s_{24}\,e^{i\,\Phi_{a5}-i\,\Phi_{d}} - s_{13}\,c_{24}\,s_{14}\,s_{34}\,e^{-i\,\Phi_{d}-i\,\phi_{7}} + c_{14}\,c_{13}\,c_{12} \\ V_{us} &= c_{13}\,c_{24}\,s_{12} - s_{13}\,s_{34}\,s_{24}\,e^{-i\,\Phi_{ab}-i\,\phi_{3}} \\ V_{ub} &= s_{13}\,c_{34}\,e^{-i\,\phi_{5}} \\ V_{ub'} &= c_{14}\,c_{13}\,s_{12}\,s_{24}\,e^{i\,\Psi_{ab}} + c_{13}\,c_{12}\,s_{14}\,e^{i\,\Psi_{d}} + s_{13}\,c_{14}\,c_{24}\,e^{-i\,\phi_{3}}\,s_{34} \end{split}$$

$$egin{aligned} V_{t'd} &= -c_{34}\,c_{24}\,s_{14}\,e^{-i\,\Phi_d} \ V_{t's} &= -c_{34}\,s_{24}\,e^{-i\,\Phi_{ab}} \ V_{t'b} &= -s_{34} \ V_{t'b'} &= c_{14}\,c_{34}\,c_{24} \end{aligned}$$

#### 2nd Row

```
\begin{split} V_{2d} &= s_{14} \, s_{13} \, s_{12} \, s_{24} \, s_{23} \, e^{i \, \Phi_{ab} - i \, \Phi_{a} + i \, \phi_{b}} - c_{12} \, c_{23} \, s_{14} \, s_{24} \, e^{i \, \Phi_{ab} - i \, \Phi_{d}} - c_{13} \, c_{24} \, s_{14} \, s_{34} \, s_{25} \, e^{-i \, \Phi_{d}} - c_{14} \, c_{12} \, e^{i \, \phi_{b}} \, s_{13} \, s_{23} - c_{14} \, c_{23} \, s_{12} \, c_{24} \, c_{25} \\ V_{5a} &= -c_{13} \, s_{34} \, s_{24} \, s_{23} \, e^{-i \, \Phi_{ab}} - c_{24} \, e^{i \, \phi_{3}} \, s_{13} \, s_{12} \, s_{23} + c_{12} \, c_{24} \, c_{23} \end{split}
```

 $\begin{aligned} V_{cb} &= c_{13}\,c_{34}\,s_{23} \\ V_{cb'} &= -c_{14}\,s_{13}\,s_{12}\,s_{24}\,s_{23}\,e^{i\,\Phi_{cb}+i\,\phi_3} + c_{14}\,c_{12}\,c_{23}\,s_{24}\,e^{i\,\Phi_{cb}} - c_{12}\,s_{14}\,s_{13}\,s_{23}\,e^{i\,\Phi_{c}+i\,\phi_3} - c_{23}\,s_{14}\,s_{12}\,e^{i\,\Phi_{d}} + c_{14}\,c_{13}\,c_{24}\,s_{34}\,s_{23} \end{aligned}$ 

#### 3rd Row

$$\begin{split} V_{td} &= c_{23}\,s_{14}\,s_{13}\,s_{12}\,s_{24}\,e^{i\Phi_{ab}-i\,\Phi_{d}+i\,\phi_{a}} + c_{12}\,s_{14}\,s_{24}\,s_{23}\,e^{i\,\Phi_{ab}-i\,\Phi_{d}} - c_{13}\,c_{24}\,c_{23}\,s_{14}\,s_{34}\,e^{-i\,\Phi_{d}} + c_{14}\,s_{12}\,s_{23} - c_{14}\,c_{12}\,c_{23}\,e^{i\,\phi_{a}}\,s_{13} \\ V_{ts} &= -c_{13}\,c_{23}\,s_{34}\,s_{24}\,e^{-i\,\Phi_{ab}} - c_{12}\,c_{24}\,s_{23} - c_{24}\,c_{23}\,e^{i\,\phi_{a}}\,s_{13}\,s_{12} \\ V_{tb} &= c_{13}\,c_{34}\,c_{23} \\ V_{tb'} &= -c_{14}\,c_{23}\,s_{13}\,s_{12}\,s_{24}\,e^{i\,\Phi_{ab}+i\,\phi_{2}} - c_{14}\,c_{12}\,s_{24}\,s_{23}\,e^{i\,\Phi_{ab}} - c_{12}\,c_{23}\,s_{14}\,s_{13}\,e^{i\,\Phi_{ab}+i\,\phi_{2}} + s_{14}\,s_{12}\,s_{23}\,e^{i\,\Phi_{d}} - c_{14}\,c_{13}\,c_{24}\,c_{23}\,s_{34} \end{split}$$

# CYMA (+ HOW)

### Conclusion

For 
$$f_{Bs}\sqrt{B_{Bs}} = 266(18)$$
 MEV,  $m_{t} = 500$  GEV

- I. For  $V_{ts}^* V_{th}^{-1} = 0.006 e^{i75^\circ}$ ,  $|V_{tb}^{-1}| < 0.13 \text{ on } Z \to bb$ ,  $|V_{tb}^{-1}| > 0.06 \text{ on } D\text{-}D \text{ mixing , best global fit at } |V_{tb}^{-1}| = 0.08$
- II.  $\sin 2\Phi_{Bs} \sim -0.33 \in [0.00, -0.40],$ Larger  $m_{\mu}$  implies larger  $\sin 2\Phi_{Bs}$
- III. Future measurement  $\mathcal{B}(K_L \to \pi^0 \nu \overline{\nu})$  and  $\sin 2\Phi_{Bd}$  can determine  $V_{rd}$ .

# NANDI

Allowed ranges for the parameters,  $\lambda_{\mathbf{t'}}^s = \mathbf{V_{t'b}V_{t's}^*}$ , phase of  $\mathbf{V_{t's}} = \phi_{\mathbf{t'}}^s$  and upper limit for  $\lambda_{\mathbf{t'}}^{\mathbf{d}} = \mathbf{V_{t'b}V_{t'd}^*}$ 

$m_{t'}\left(GeV ight)$	400	600	
$\lambda_{t'}^s  imes 10^2$	(0.08 - 1.4)	(0.05 - 0.6)	
$\phi^s_{t'}$	-80 → 80	-80 → 80	
$\lambda_{t'}^d  imes 10^3$	< 2.0	< 1.6	

Upper bound on the SM4 coupling reducing with the mass of fourth family...

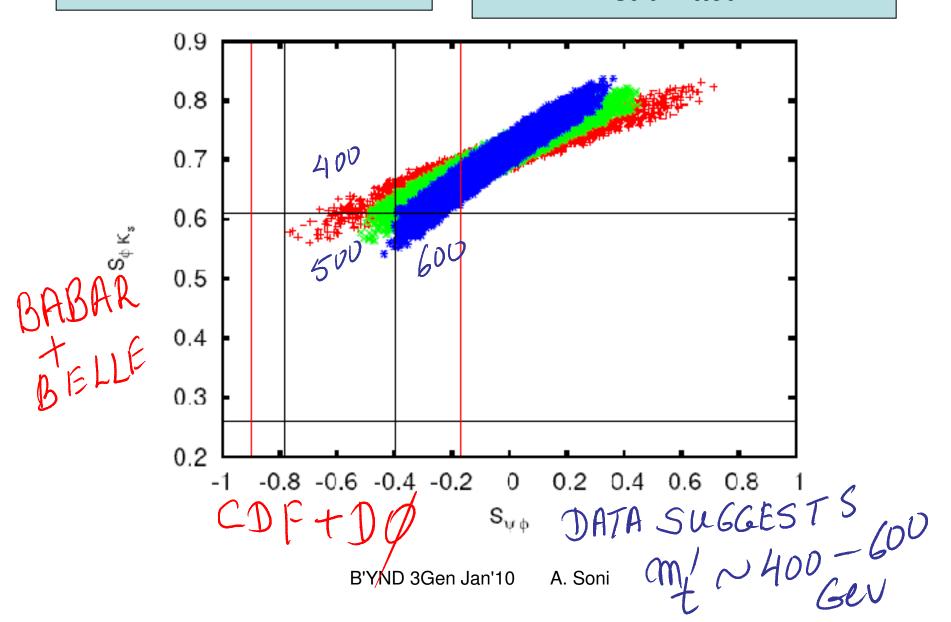
Constrain used ⇒ CKM unitarity + Oblique corrections +

$$Br(B \to X_s \gamma) + \epsilon_k + \Delta M_d + \frac{\Delta M_d}{\Delta M_s} + Br(B \to X_s \ell^+ \ell^-) +$$

- 
$$\mathcal{BR}(K^+ \to \pi^+ \nu \nu)$$



# arXiv:0807.1971, and to be submitted



d(S)

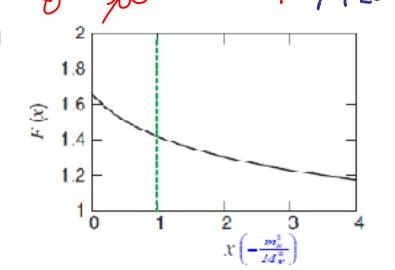
# **THUS**

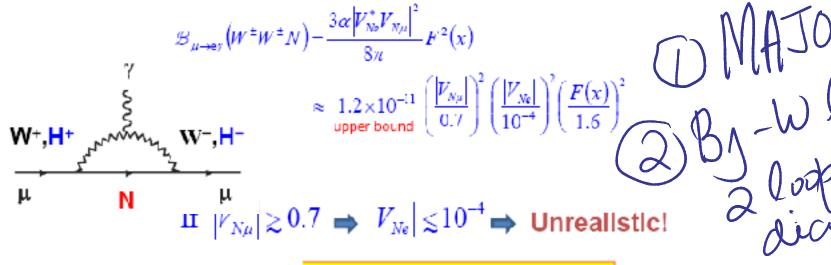
ly constantion

- The CKM-paradigm of CP violation accounts for the observed CP patterns to an accuracy of about 15%!
- SM3-CKM predicted value of sin2β tends to be high compared to direct (ψ K) measurements by about 15-20%...t is dominant
- Hierarchical structure of SM4 mixing matrix NATURALLY lets t' be subdominant here but due to its large mass (and decoupling theorem) not negligible Leads to Small ~ 15.1 deviations
- Dynamics of EW gauge interactions (evasion of decoupling theorem) by EWpenguins and the large mt' plays an important role in the large "isospin" violating  $\Delta A_{CP}$  (K π)
- SM3 says B<sub>s</sub> mixing has negligible CP-odd phase therein t' plays a dominant role (& t is subdominant)

#### HOU, MA & LEE, PRD'09

- F(x), an Inami-Lim loopfunction, is well-behaved and bounded
- F(1) = 17/12 : not singular
   Incorrectly rendered in W. Huo and T. F. Feng. arXiv:0301153
   and K. R. Lynch arXiv:0108081
- Since F(x) ~ 1.4, |V<sub>Nµ</sub>| needs to be ≥ 0.7 to reach within 2o of △a<sub>µ</sub>
- On the other hand,

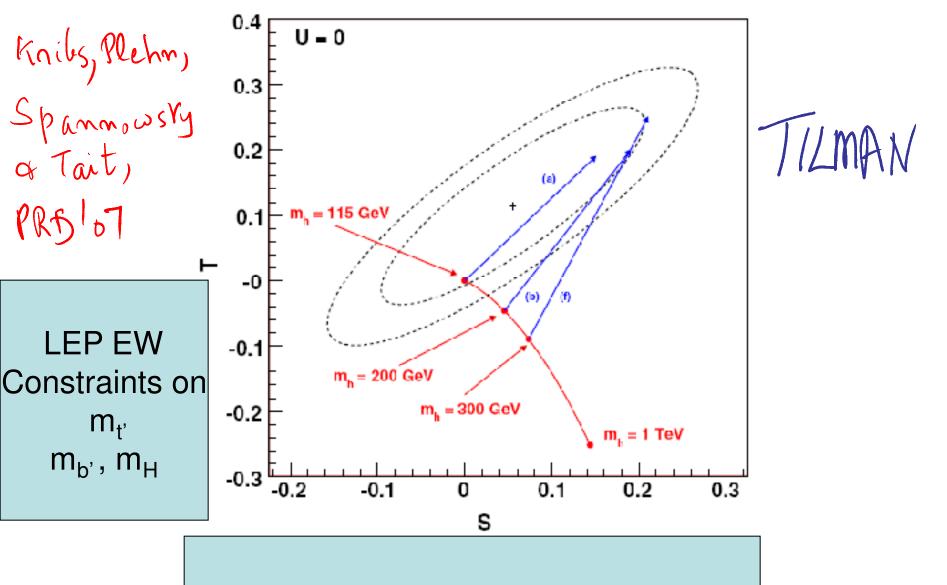




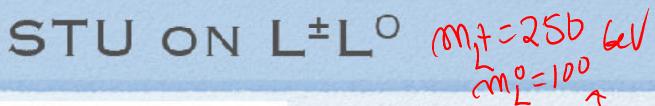
△a, unlikely from W+W-N

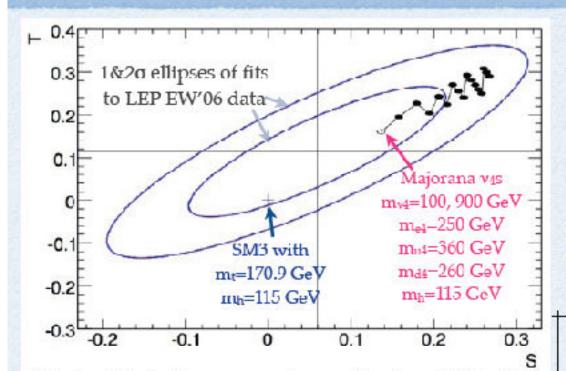
B'YND 3Gen Jan'10 A. Soni

because of  $\mu \rightarrow \epsilon \gamma$ 



4<sup>th</sup> family is not inconsistent with LEP EWPC (despite perhaps wishful thinking of some Influential PDG authors)





Black solid circles represent magoing from 150 to 900. GeV in steps of 50 GeV, while the best value of mu4 goes slowly up to 390 CeV.

Implemented exact one-loop calculations from:

- B.A.Kniehl & H.G.Khors, PRD48(1993)225.
- H.J.He, N.Polonsky & S.F.Su PRD64(2001)053004.
- => Referance scenario can easily be accomodated for light or heavy Higgs alike.

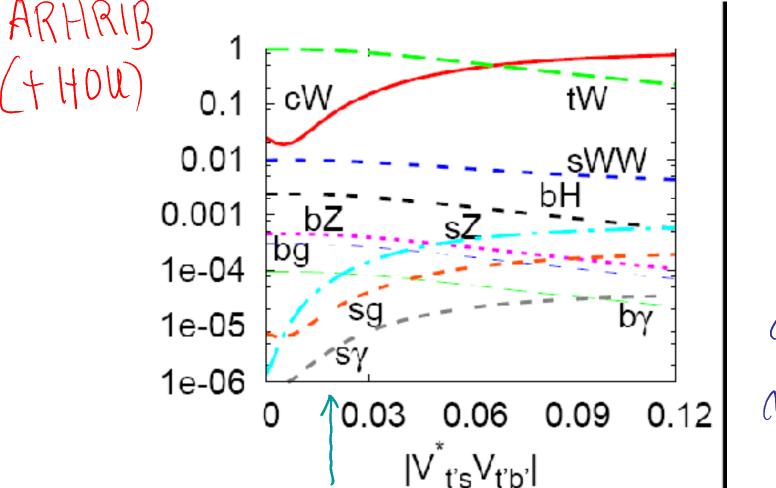
After all cuts, 1 fb-1 @ 14 TeV:

- 7 signal events
- Rudimentary background estimate = 0.22 events

Further details in: V. E. Özcan, S. Sultansoy, G. Unel, A Possible Discovery Channel for New Charged Leptons at the LHC, J. Phys. G 36 (2009) 095002.

C ALSO ANTIPIN

# Indirect effects: Ahrib,CHShen,Bashiry,CYMa, Nandi,(Eilam M T)...



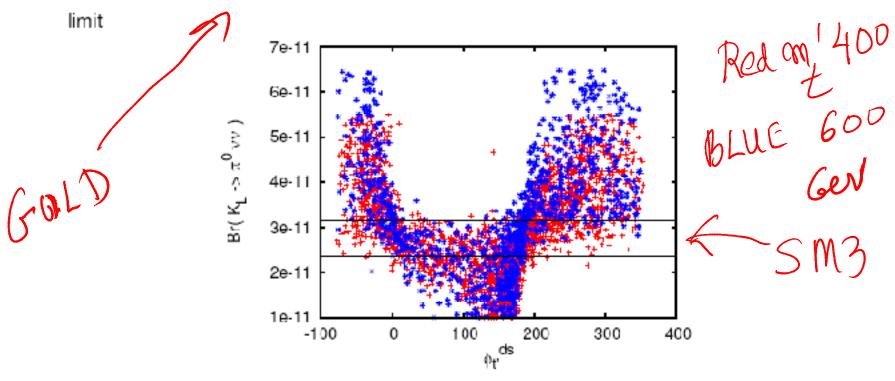
m/ ~340

Figure 6:  $m_{b'} \iff 340 \text{ GeV}, m_{t'} = m_{b'} = 50 \text{ GeV},$  $V_{ts}V_{tb'} = 0.01 e^{-i \cdot 10^{\circ}} \text{ and } \arg(V_{t's}^*V_{t'b'}) = 70^{\circ}$ 

CALSO EILAM, MELIC, TRAMPET/C
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# NAVDI

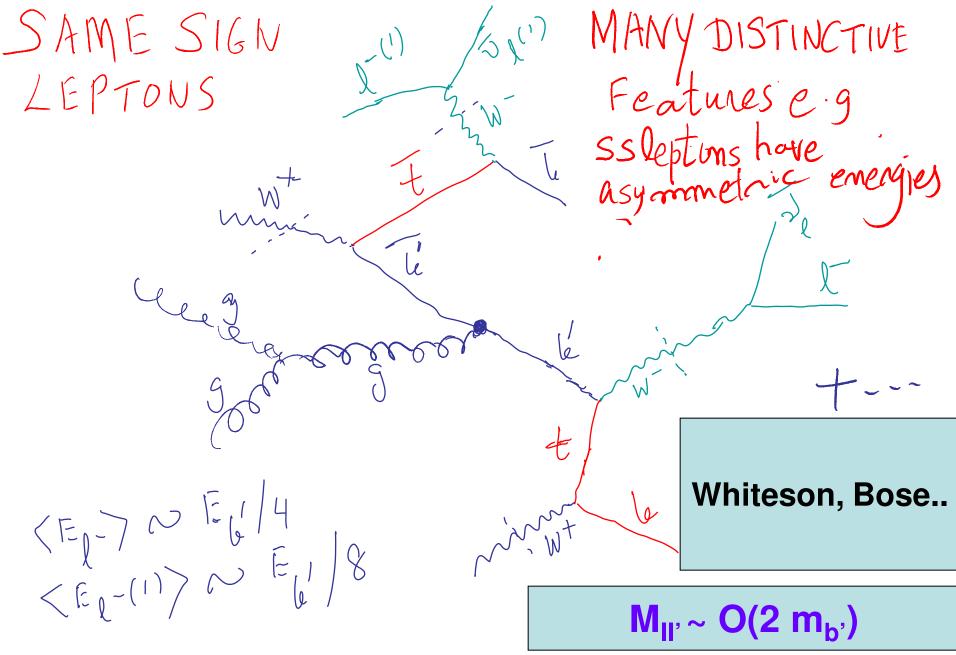
Variation of  $\mathbf{Br}(\mathbf{K_L} \to \pi^0 \nu \bar{\nu})$  with  $\phi^{\mathbf{ds}}_{\mathbf{t'}} = \phi^{\mathbf{d}}_{\mathbf{t'}} - \phi^{\mathbf{s}}_{\mathbf{t'}}$ , black horizontal lines correspond to SM



$$Br(K_L \to \pi^0 \nu \bar{\nu}) \propto \left[ \left( \frac{Im\lambda_t}{\lambda^5} X(x_t) + \frac{Im\lambda_{t'}}{\lambda^5} X(x_{t'}) \right)^2 \right]$$
  
 $Br(K_L \to \pi^0 \nu \bar{\nu}) = (1.0 \to 5.2) \times 10^{-11} \Rightarrow 1 \sigma \text{ range in SM4}$ 

## DIRECT SEARCHES

Whiteson, Bose, Ozcan, Antipin...



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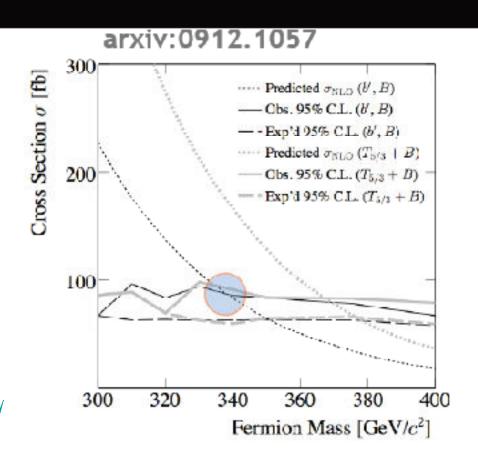
A. Soni

# Limits

WHITESON

### Limit

mb > 338 GeV mt / 7311 TEVATRON / 11 ~ 400 Ge)



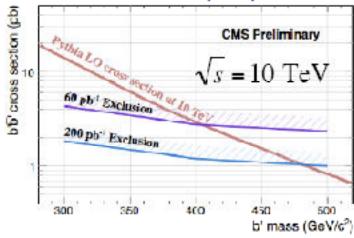
BOSE

### Heavy b' limits

m(b <sup>1</sup> ) @ 200pb <sup>-1</sup>	300 GeV	400 GeV	500 GeV
Cross-section	13.6 pb	2.8 pb	0.78 pb
Expected Yields	34.08	10.58	3.52
Background	1.08	1.08	1.08
Significance	9.0 σ	3.7 σ	1.4 σ

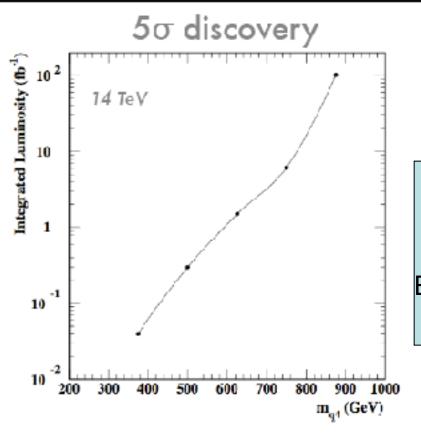
CM3

Exclude b' masses less than 485 (405) GeV with 200 (60) pb<sup>-1</sup>



Stringent limits can be set at the LHC with early data

# Sensitivity



WHITESON

LHC should be able to cover the entire Expected mass range

From ATLAS SN-ATLAS-2008-069

 $Z^* o \chi \chi$  and  $W^* o \ell^{\pm} \chi$  production channels.

# LHC Phenomenology II

#### Production and decay of new heavy leptons:

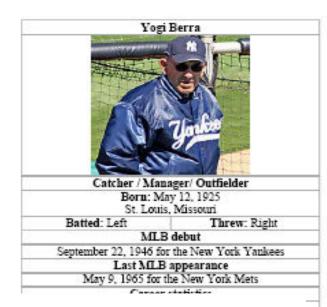
Table 1: Signal cross-sections  $\tau$  in (b) with the corresponding leading SM background for three scenarios described in the text. Pre-selection and selection criteria are also described in the text.

$\mu^{\perp}\mu^{\perp} X (M_{\chi_1} = 100 \text{ GeV})$	Pre-selected $\sigma(tb)$	Selected $\sigma(ib)$	Events/10 $tb^{-1}$	S/4/B
$\mu^{\pm}\chi_{1}: \mu^{\pm}\mu^{\pm} + 2 \text{ jets}$	1.11	0.6	6	7.25
$\chi_1 \chi_1 : \mu^{\pm} \mu^{\pm} + 4 \text{ jets}$	2.5	1.32	13.2	5
SM background	10.25	0.7	7	2
$\mu^{\pm}\mu^{\pm} X (M_{\chi_1} - 90 \text{ GeV})$	Pre selected $\sigma(fb)$	Selected $\sigma(fb)$	Events/10 fb <sup>-1</sup>	$S/\sqrt{B}$
$\mu^{+}\chi_{1}:\mu^{+}\mu^{+}+2$ jets	0.21	0.113	1.1	1.89
$\chi_1 \chi_1 : \mu^+ \mu^+ + 4 \text{ jets}$	0.73	0.39	3.9	
SM background	13.25	0.7	7	
$\mu^{\pm}\mu^{\pm} X (M_{\chi_2} = 135 \text{ GeV})$	Pre-selected $\sigma(\text{fb})$	Selected $\sigma(fb)$	Events/10 fb 1	$S/\sqrt{B}$
$\mu^{\perp}\chi_{1}:\mu^{\perp}\mu^{\perp}+2$ jets	2.1	1.1	11	6.95
$\chi_1 \chi_1 : \mu^{\perp} \mu^{\perp} + 4 \text{ jets}$	1.4	0.74	7.4	
SM background	10.25	0.7	7	
$\ell^+\ell^+\ell^\mp \times (M_{\gamma_1}=100 \text{ GeV})$	Pre-selected σ(fb)	Selected $\sigma(\mathrm{fb})$	Events/10 fb <sup>-1</sup>	$3/\sqrt{B}$

In the most optimistic cases early discoveries at the LHC might be expected.

# The Epilogue: Challenging Yogi

# The Legendary American philosopher



### The Future

 Yogi Berra: "Its difficult to make predictions, especially about the future"

- New York Yankees (1964, 1984-1985)
- New York Mets (1972-1975)

#### Career highlights and awards

- 15× All-Star selection (1948, 1949, 1950, 1951, 1952, 1953, 1954, 1955, 1956, 1957, 1958, 1959, 1960, 1961, 1962)
- 13× World Series champion (1947, 1949, 1950, 1951, 1952, 1953, 1956, 1958, 1961, 1962, 1969, 1977, 1978)
- 3× AL MVP (1951, 1954, 1955)
- New York Yankees #8 retired
- Major League Baseball All-Century Team

Member of the National

Baseball Hall of Fame A A Induction 1972

Vote 85 61% (second ballor)

\* Lifted from Stone Sheldon Stone

# Prospects: (why not?) Glorious

 By late 2010-mid'11: S(Bs->ψφ) nonvanishing to ~4 σ through combined efforts of CDF/D0/possibly LHCb!

- '11-12: Gold-plated H->ZZ!! (heavy, possibly broad)
- 12-13: Excess of same sign dileptons signaling mb'~450-500GeV!!!

### **BORING REPETITION?**

- If the mt' is heavy  $\sim$ (400-600) GeV, then for sure it will have serious role to play in EWSB .(NOTE CDF+D0 latest bound m<sub>t'</sub> > 350 GeV).
- It will clearly have significant impact on CP violation phenomena, given that now we will have 2 additional CP-odd phases
- It may play an interesting role in baryogenesis (W.-S. Hou, 0803.1234)
- CANNOT BE A CONVENTIONAL 4th Gen..mv4>mZ/2
- Possible DMC (if no mixing with lighter 3 nu's)..see e.g. Volovik'03
- It may open up possibilty of unification (PQ Hung,'98)
- Can be observed or ruled out at LHC

# Cons: "Cancellations"

- Extra contributions to EWP observables due mt',mb' need to be cancelled by the heavier "higgs"
- Similarly, |mt'-mb'| < ~ 60 GeV for mt' O(500 GeV)

  GeV)

  O(10).) Degeneracy
- So how much of a concern should one give to these cons?
- Let's just remember Δ(mn-mp)<O(0.1%)</li>
   We understand this now as due ISOSPIN

# Summary & Conclusions (I)

- While for now no compelling evidence against CKM-picture, several fairly sizeable effects ( $\sim 2$   $\sim 3.5$   $\sigma$ ) in B,Bs CP asymmetries are difficult to understand in SM3.
- Being careful, "conservative" in such instances means hunting down seriously the underlying cause.....
- CDF/D0 finding of Bs->psi phi "anomaly" is especially interesting since, 1) it is essentially predicted by the ones from BFs, 2) unlike the other effects, this is theoretically very clean i.e. GOLD-PLATED
  - Consequently, it is extremely important that FERMILAB follows it up & clarifies it with a very high priority.

# Summary & Conclusions (II)

- SM4 opens up important new avenues for baryogenesis, DMC, unification, and most likely also crucial for EWSB...thereby it may well lead to a possible resolution to the hierarchy puzzle
- Underlying nature of "4th gen." has to be significantly diff
- If the effects stand further scrutiny, SM4 with mt', mb' (400-600 GeV)
  provides a v. simple explanation.
- Direct verification at LHC should be relatively easy.
- Suggestive also of a HEAVY (therefore most likely composite), likely broad HIGGS
- Dynamical EW symm breaking, strong dynamics, the lattice & high intensity facilities (e.g. SBF) may well have a crucial role to play even in the era of the LHC

# THANK YOU Very much.

Local Organizing Committee: Paoti Chang (National Taiwan University) Yuan-Hann Chang (National Central University) Chuan-Hung Chen (National Cheng Kung University) Kai-Feng Chen (National Taiwan University) Kingman Cheung (National Tsing Hua University) Cheng-Wei Chiang (National Central University) Xiao-Gang He (National Taiwan University) George Hou (National Taiwan University, Chair) Song-Ming Wang (Academia Sinica) Kwei-Chou Yang (Chung Yuan Christian University) Tzu-Chiang Yuan (Academia Sinica) For a fanta stic meeting.

B'YND 3Gen Jan'10 A. Soni