

VIRTUAL CORRECTIONS

TO

MASSLESS QCD

PROCESSES WITH

D-DIMENSIONAL

UNITARITY

ACHILLEAS
LAZOPOULOS
RADCOR 2009
ETH ZURICH

Introduction

Introduction

- Massless quarks and gluons the necessary ingredients for any physical process.
- Four jets - landmark NLO calculation.
- $b\bar{b}b\bar{b}$ production with massless bottoms.

Introduction: the formula

$$\sigma_{nlo} = \int_n d\sigma_{tree} + \boxed{\int_n d\sigma_{virt}} + \int_{n+1} d\sigma_{real}$$

WITHIN D-DIMENSIONAL UNITARITY

Giele, Kunszt, Melnikov [0801.2237]

Ellis, Giele, Kunszt [0708.2398]

Ellis, Giele, Kunszt, Melnikov [0806.3467]

Zanderighi [...], Winter [...], Schulze [...]

AL[0812.2998]

Introduction: Color decomposition

$$A_{\text{gluons}}^{NLO} = \sum_{\sigma} CF_1(\sigma) \text{ [diagram 1]} + \frac{n_f}{N_c} \sum_{\sigma'} CF_2(\sigma') \text{ [diagram 2]}$$

The first diagram shows a gluon exchange between two vertices, each with two external gluon lines. The second diagram shows a fermion loop with two external gluon lines and two internal fermion lines.

$$A_{q\bar{q}+Ng}^{NLO} = \sum_{\sigma} CF_1(\sigma) \text{ [diagram 3]} + \frac{n_f}{N_c} \sum_{\sigma'} CF_2(\sigma') \text{ [diagram 4]} + \sum_{\sigma''} CF_3(\sigma'') \text{ [diagram 5]} \text{ [diagram 6]}$$

The third diagram shows a gluon exchange between two vertices, each with two external fermion lines. The fourth diagram shows a fermion loop with two external gluon lines and two internal fermion lines. The fifth diagram shows a gluon exchange between two vertices, each with two external gluon lines. The sixth diagram shows a gluon exchange between two vertices, each with two external gluon lines.

✓ No general formula for more than one fermion pair

Introduction: D-dimensional Unitarity

$$A_{\text{I}}^{D_s} = A_0 + D_s A_1 \quad A^{FDH} = 2A^{(6)} - A^{(8)}$$


$$A_{\text{II}}^{D_s} = 2^{D_s/2-1} A_0 \quad A^{FDH} = 8A^{(6)} = 16A^{(8)}$$


“Integrand level reduction ... partial fractioning of the amplitude over the standard base of master integrals ... OPP system”

Introduction: the OPP system in DDU

Finding $A^{(D_s)}$

$$\bar{e}_{I_5,s} = \mathcal{P}|_{l_s}^{I_5} = e_{I,0}$$

$$\bar{d}_{I_4,s} = \mathcal{P}|_{l_s}^{I_4} - \sum_{J_5/I_4} \frac{\bar{e}_{J_5/I_4}}{D_{k_1}(l_s)} = \sum_r d_{I_4,r} f_r(l_s)$$

$$\bar{c}_{I_3,s} = \mathcal{P}|_{l_s}^{I_3} - \sum_{J_5/I_3} \frac{\bar{e}_{J_5/I_3}}{D_{k_1}(l_s)D_{k_2}(l_s)} - \sum_{J_4/I_3} \frac{\bar{d}_{J_4/I_3}(l_s)}{D_{k_1}(l_s)} = \sum_r c_{I_3,r} g_r(l_s)$$

$$\bar{b}_{I_2,s} = \mathcal{P}|_{l_s}^{I_2} - \sum_{J_5/I_2} \frac{\bar{e}_{J_5/I_2}}{D_{k_1}(l_s)D_{k_2}(l_s)D_{k_3}(l_s)} - \sum_{J_4/I_2} \frac{\bar{d}_{J_4/I_2}(l_s)}{D_{k_1}(l_s)D_{k_2}(l_s)} - \sum_{J_3/I_2} \frac{\bar{c}_{J_3/I_2}(l_s)}{D_{k_1}(l_s)} = \sum_r b_{I_2,r} h_r(l_s)$$

b) Solve for those

a) Fix loop momentum to evaluate numerically those

Introduction: the OPP system in DDU

Finding $A^{(D_s)}$

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 \bar{e}_{I_5,s} &= \mathcal{P}|_{l_s}^{I_5} = e_{I,0} && l_0 \\
 \bar{d}_{I_4,s} &= \mathcal{P}|_{l_s}^{I_4} - \sum_{J_5/I_4} \frac{\bar{e}_{J_5/I_4}}{D_{k_1}(l_s)} = \sum_r d_{I_4,r} f_r(l_s) && l_0, l_1 \quad l_2, l_3, l_4 \\
 \bar{c}_{I_3,s} &= \mathcal{P}|_{l_s}^{I_3} - \sum_{J_5/I_3} \frac{\bar{e}_{J_5/I_3}}{D_{k_1}(l_s)D_{k_2}(l_s)} - \sum_{J_4/I_3} \frac{\bar{d}_{J_4/I_3}(l_s)}{D_{k_1}(l_s)} = \sum_r c_{I_3,r} g_r(l_s) && l_0 \dots 6 \quad l_7, l_8, l_9 \\
 \bar{b}_{I_2,s} &= \mathcal{P}|_{l_s}^{I_2} - \sum_{J_5/I_2} \frac{\bar{e}_{J_5/I_2}}{D_{k_1}(l_s)D_{k_2}(l_s)D_{k_3}(l_s)} - \sum_{J_4/I_2} \frac{\bar{d}_{J_4/I_2}(l_s)}{D_{k_1}(l_s)D_{k_2}(l_s)} - \sum_{J_3/I_2} \frac{\bar{c}_{J_3/I_2}(l_s)}{D_{k_1}(l_s)} = \sum_r b_{I_2,r} h_r(l_s) \\
 &&& l_0 \dots 8 \quad l_9
 \end{aligned}$$

Cuts in D_s

Note: The system is linear $2A^{(6)} - A^{(8)} \rightarrow 2\mathcal{P}^{(6)}|_l^I - \mathcal{P}^{(8)}|_l^I$

Introduction: the OPP system in DDU

Finding $A^{(D_s)}$

$$\begin{aligned}
 \bar{e}_{I_5,s} &= \mathcal{P}|_{l_s}^{I_5} = e_{I,0} && l_0 \\
 \bar{d}_{I_4,s} &= \mathcal{P}|_{l_s}^{I_4} - \sum_{J_5/I_4} \frac{\bar{e}_{J_5/I_4}}{D_{k_1}(l_s)} = \sum_r d_{I_4,r} f_r(l_s) && l_0, l_1 \quad l_2, l_3, l_4 \\
 \bar{c}_{I_3,s} &= \mathcal{P}|_{l_s}^{I_3} - \sum_{J_5/I_3} \frac{\bar{e}_{J_5/I_3}}{D_{k_1}(l_s)D_{k_2}(l_s)} - \sum_{J_4/I_3} \frac{\bar{d}_{J_4/I_3}(l_s)}{D_{k_1}(l_s)} = \sum_r c_{I_3,r} g_r(l_s) && l_0 \dots 6 \quad l_7, l_8, l_9 \\
 \bar{b}_{I_2,s} &= \mathcal{P}|_{l_s}^{I_2} - \sum_{J_5/I_2} \frac{\bar{e}_{J_5/I_2}}{D_{k_1}(l_s)D_{k_2}(l_s)D_{k_3}(l_s)} - \sum_{J_4/I_2} \frac{\bar{d}_{J_4/I_2}(l_s)}{D_{k_1}(l_s)D_{k_2}(l_s)} - \sum_{J_3/I_2} \frac{\bar{c}_{J_3/I_2}(l_s)}{D_{k_1}(l_s)} = \sum_r b_{I_2,r} h_r(l_s) \\
 &&& l_0 \dots 8 \quad l_9
 \end{aligned}$$

Cuts in D_s



Note: when $l_s \in 4D \rightarrow A^{(6)} = A^{(8)} = A^{(4)}$

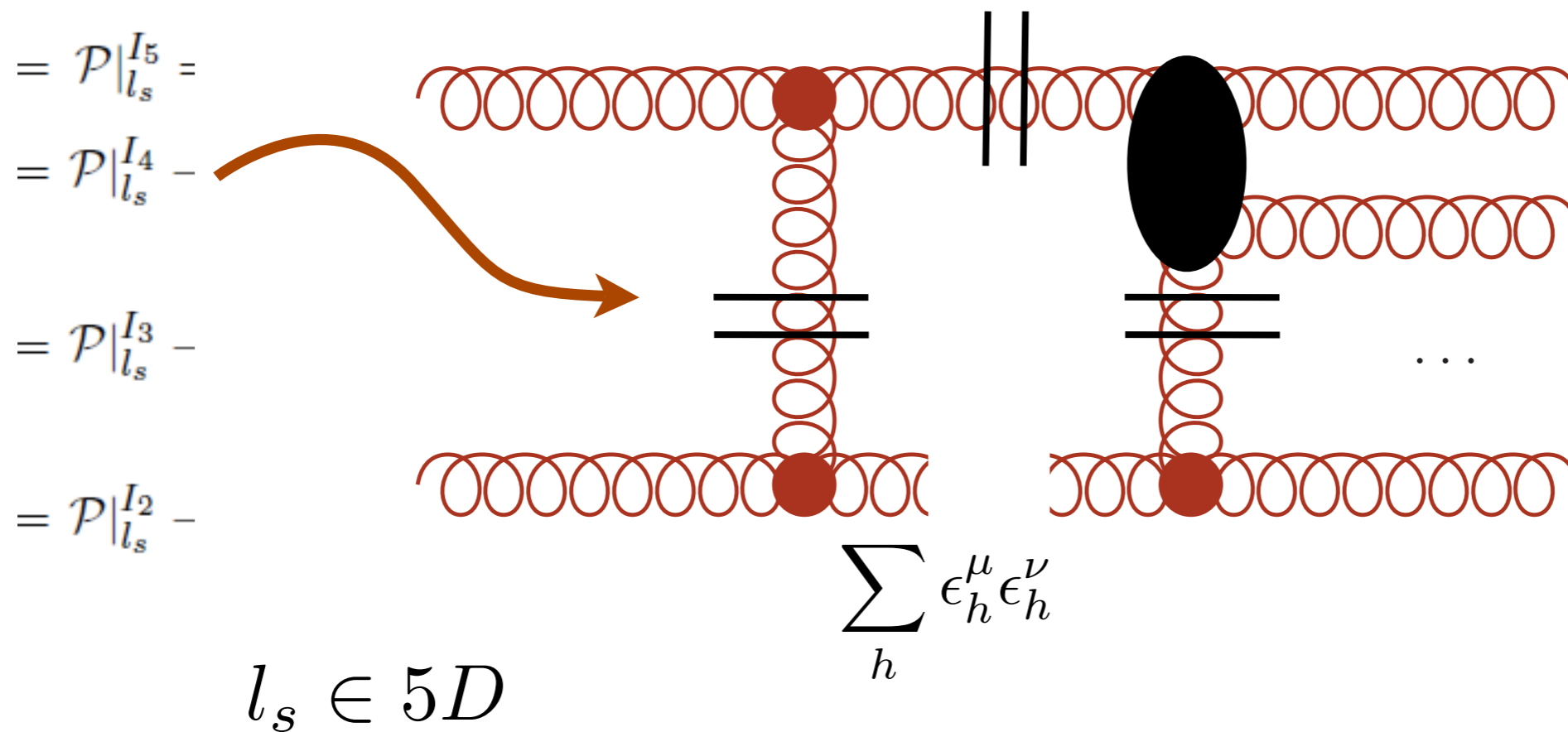
Cuts in 4D



NOT FERMION LOOP

Introduction: the OPP system in DDU

Finding $A^{(D_s)}$

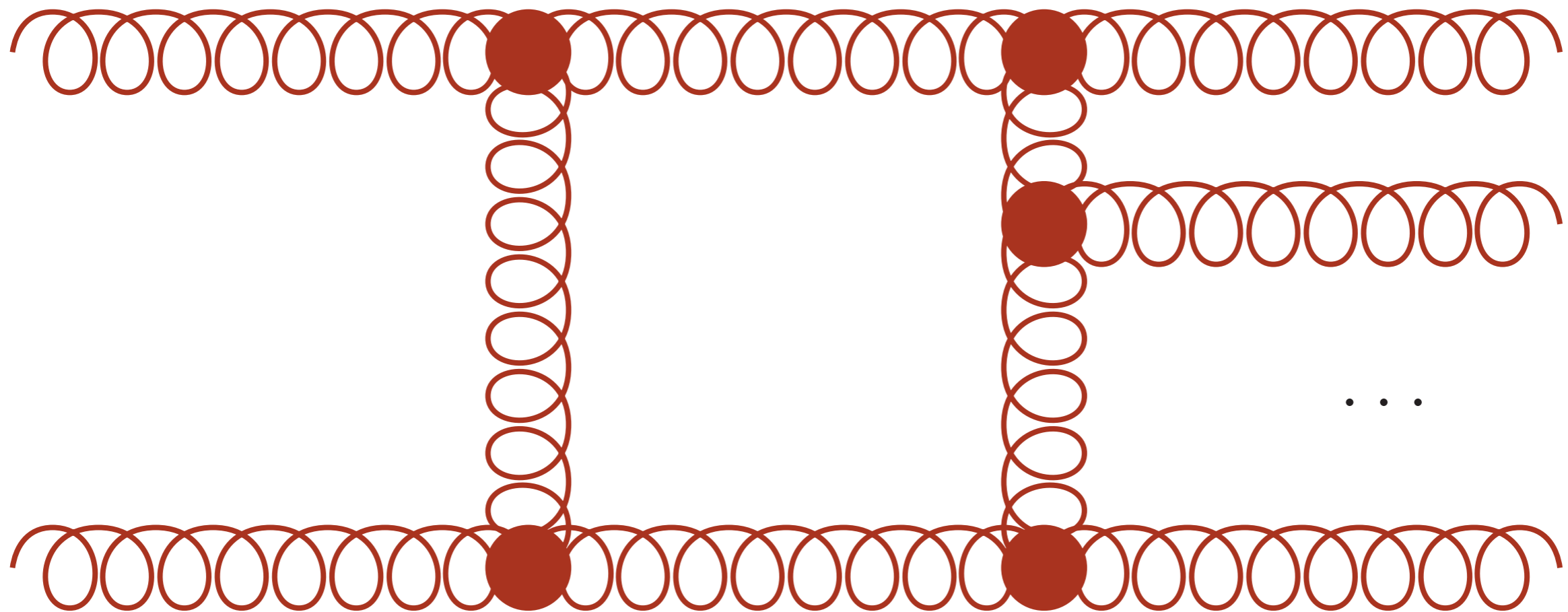


$$\mathcal{P}|_{l_s}^{I_x} = 2\mathcal{P}^{(6)}|_{l_s}^{I_x} - \mathcal{P}^{(8)}|_{l_s}^{I_x} = \mathcal{P}^{(6)}|_{l_s}^{I_x} - \mathcal{P}^{(8^*)}|_{l_s}^{I_x}$$

Primitives

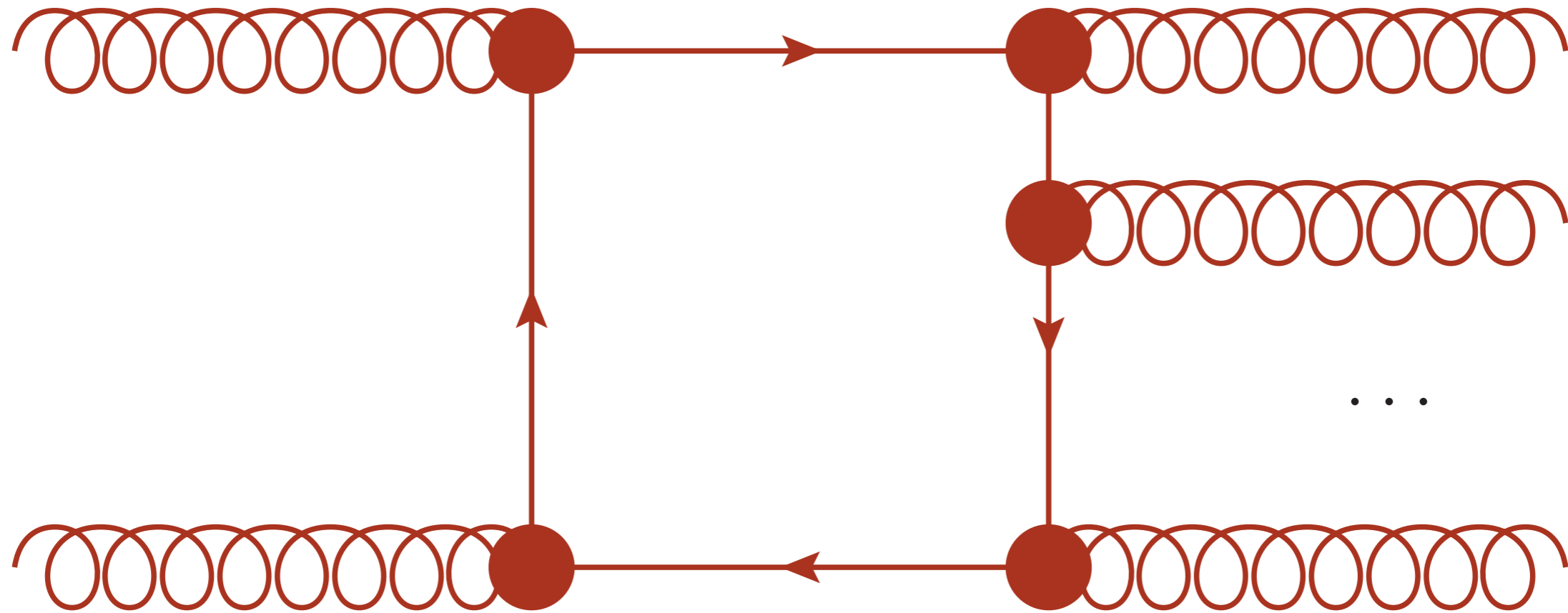
Gluons

Primitives: purely gluonic



- ✓ The most dangerous numerically
- ✓ The fastest for fixed N

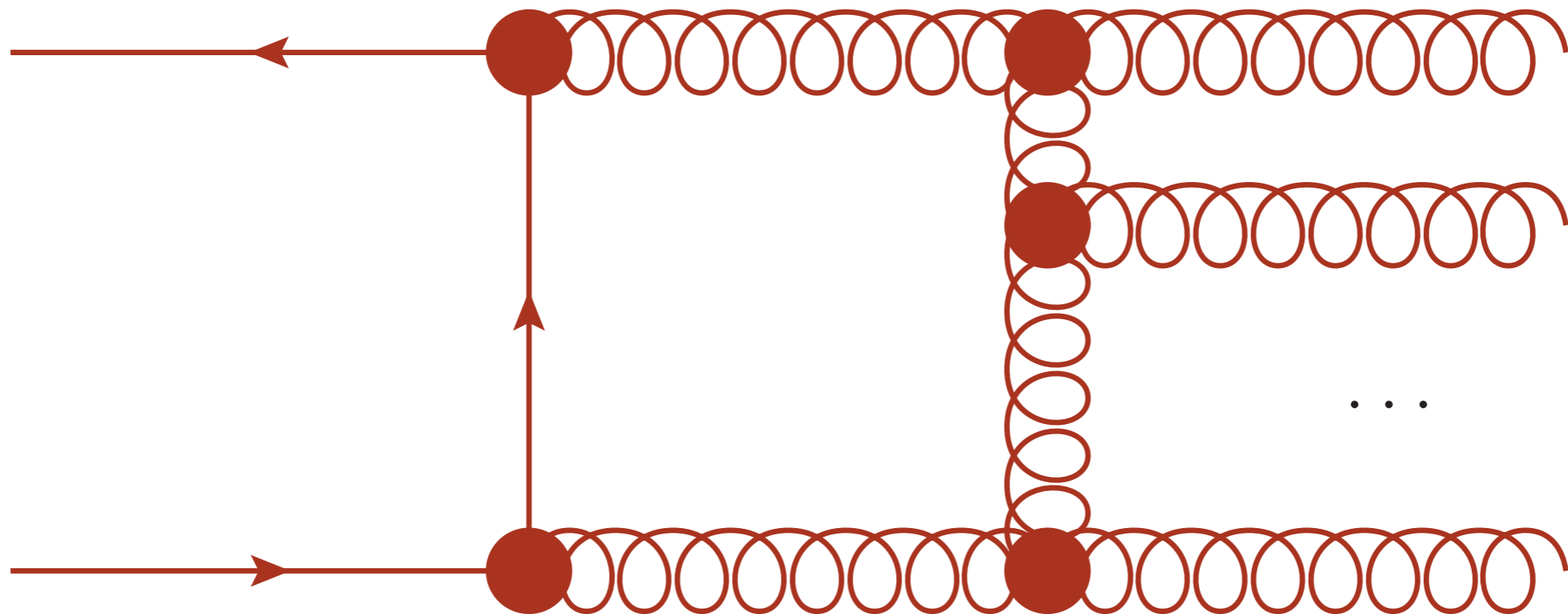
Primitives: gluonic - fermion loop



- ✓ Much more stable
- ✓ Proportional to n_f/N_c

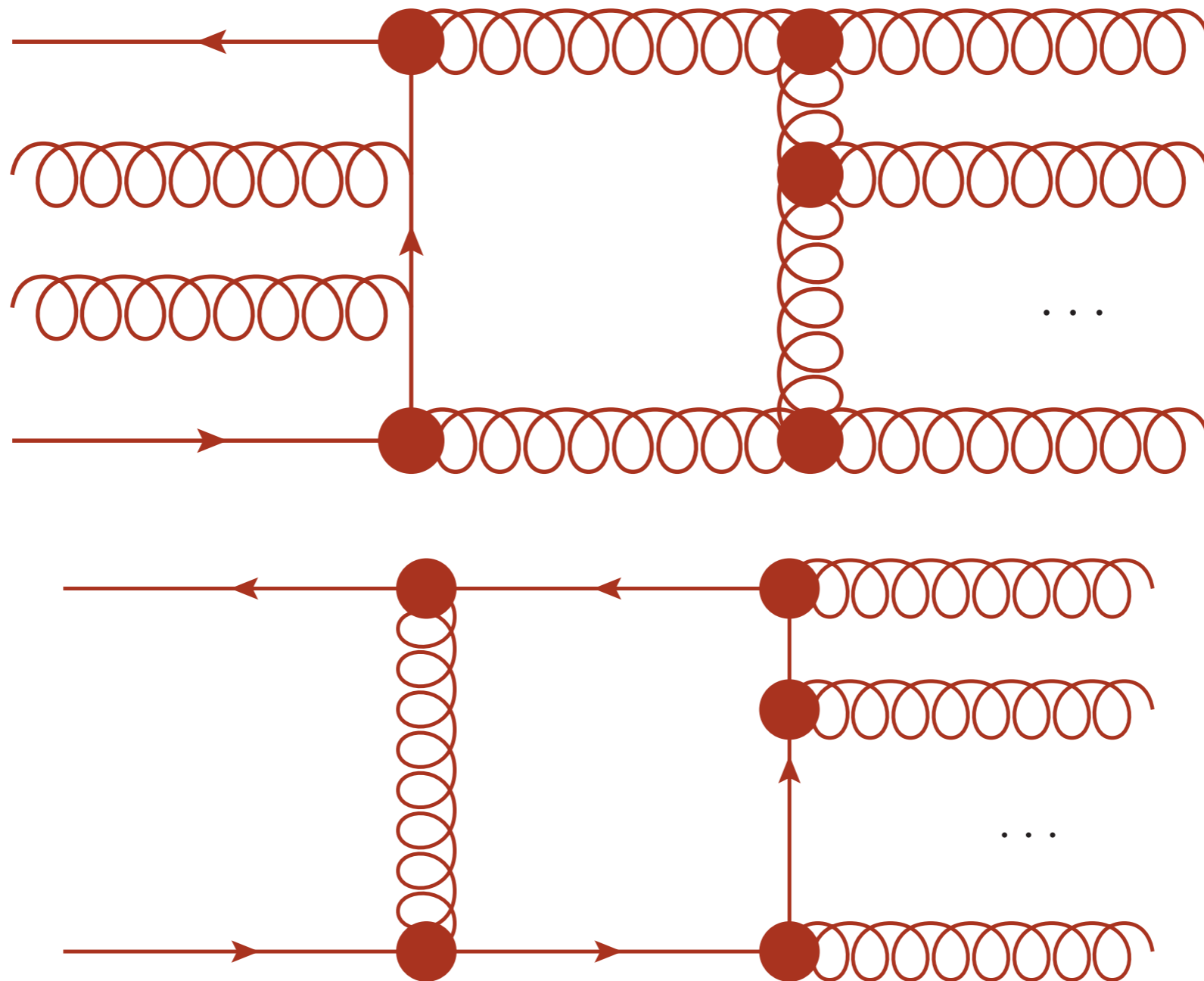
Fermions

Primitives: single fermion line



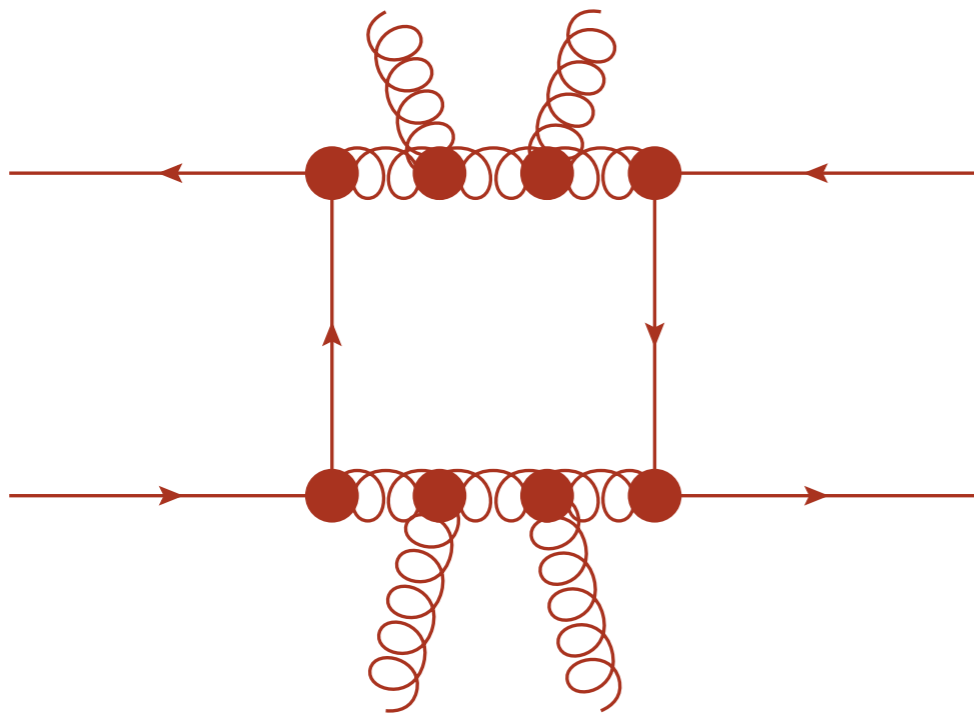
✓ Leading color

Primitives: single fermion line

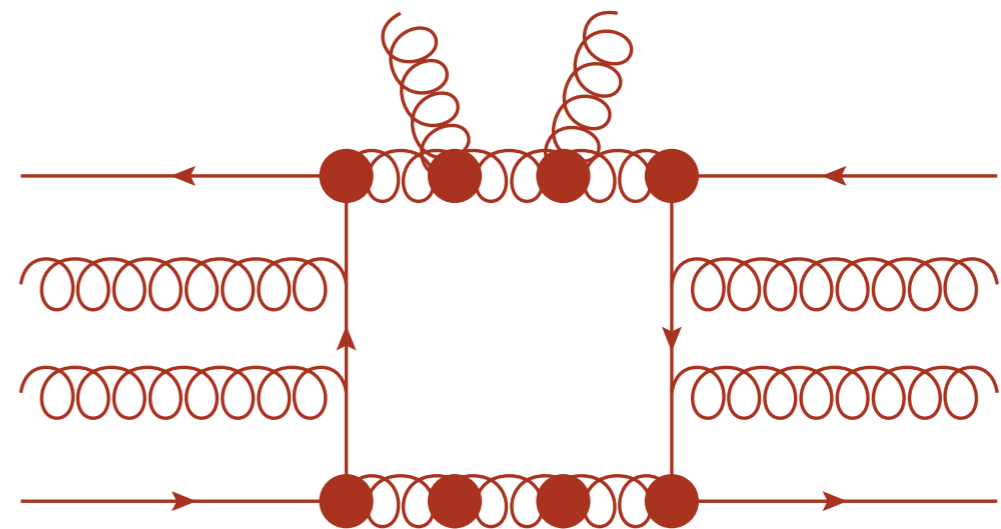


✓ $1/N_c$ suppressed

Primitives: two fermion lines

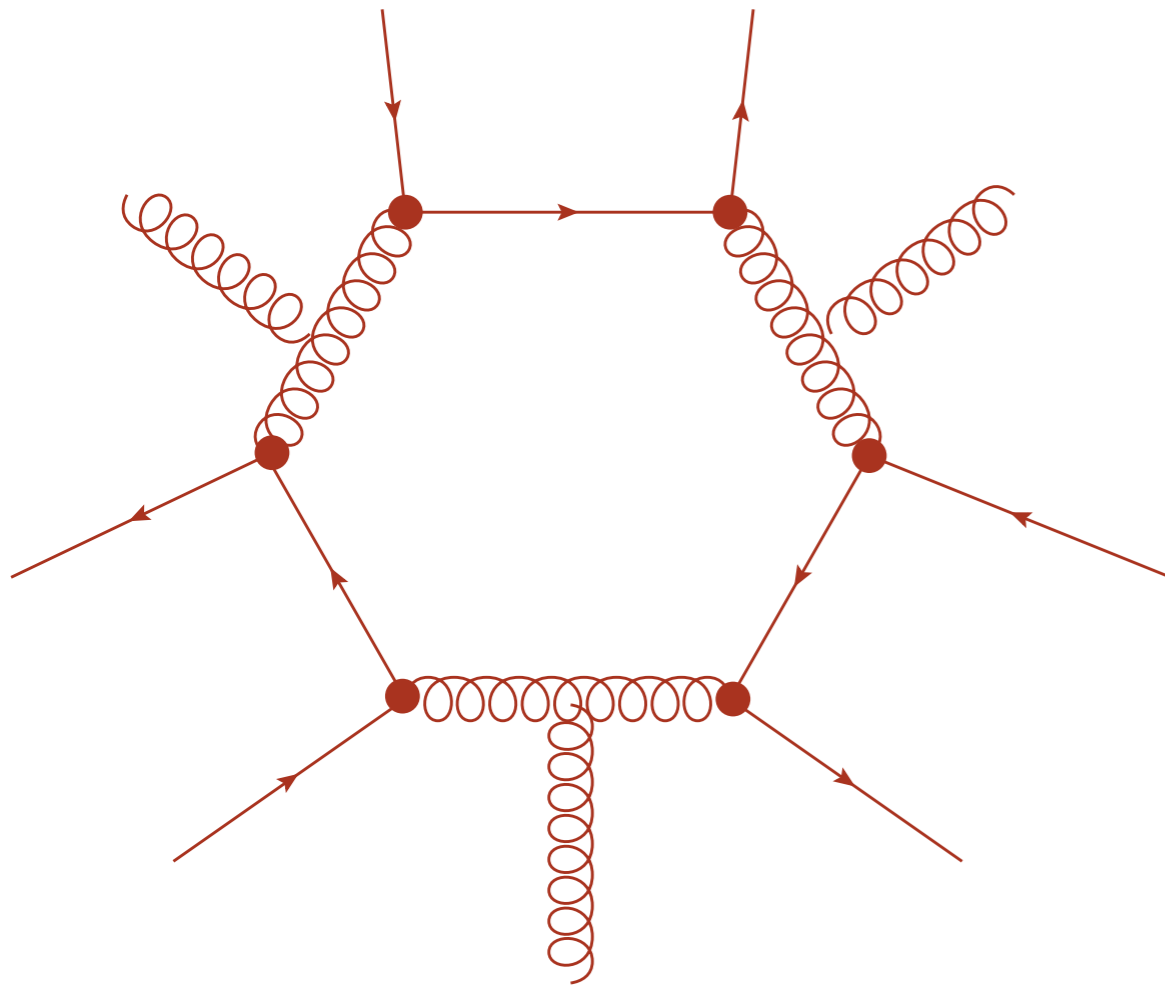


- ✓ Leading color
- ✓ Fermion line direction

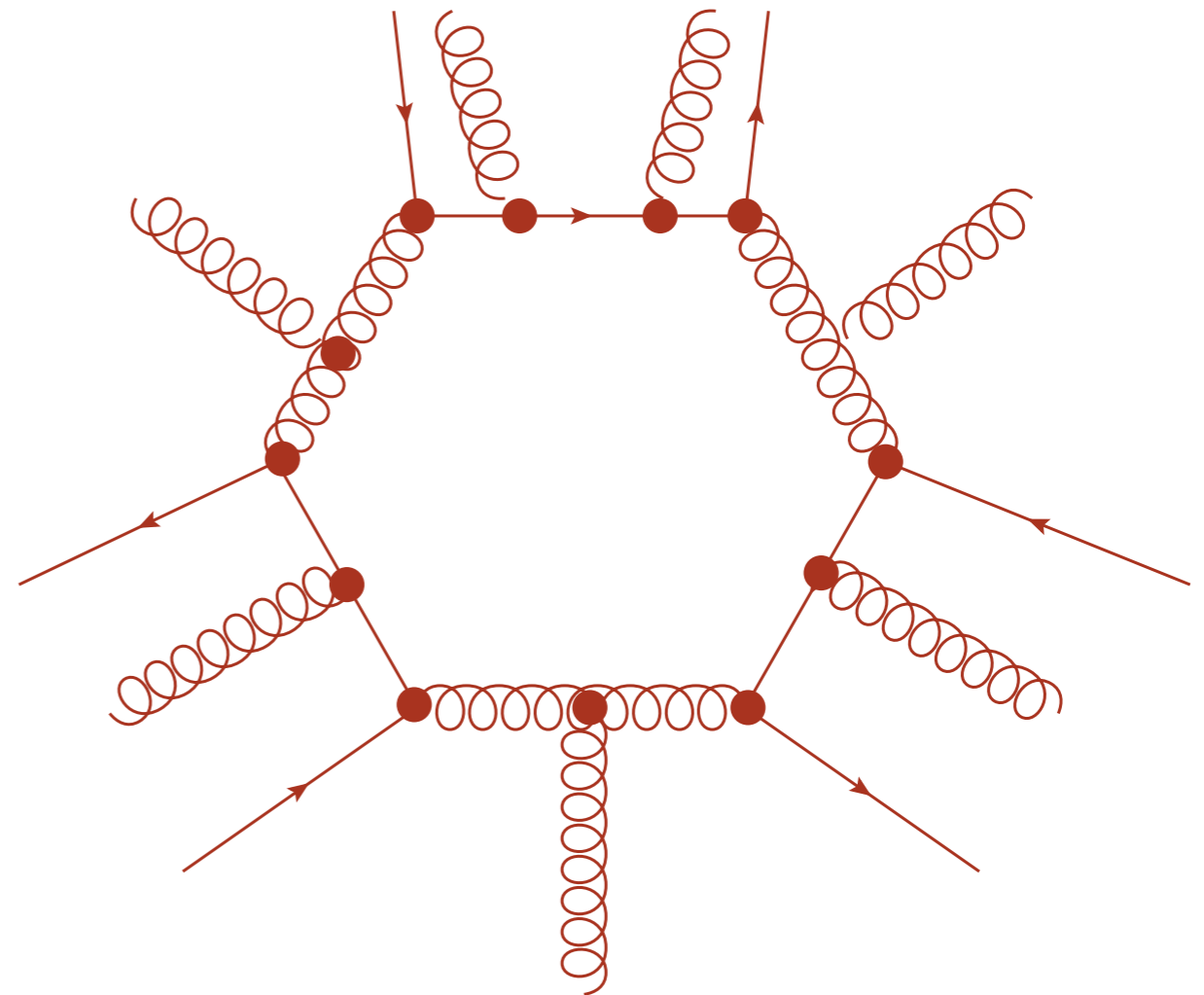


- ✓ $1/N_c$ suppressed

Primitives: three fermion lines



- ✓ Leading Color
- ✓ Fermion line direction



- ✓ $1/N_c$ suppressed

Primitives

“All necessary ingredients for a virtual amplitude with massless QCD partons are in place, **tested**, ready for production mode”.

Numerical Stability

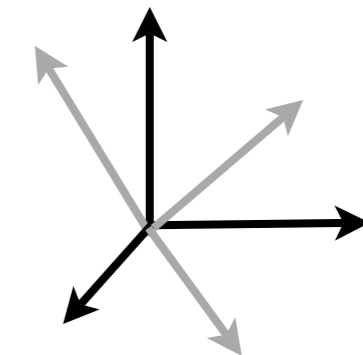
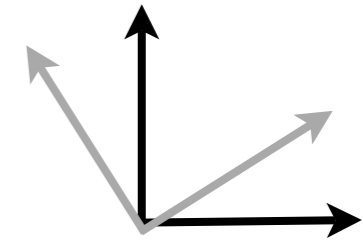
Numerical Stability: accidental instabilities

TRIPLE CUT

$$l^\mu = V^\mu + a_1 n_1^\mu + a_2 n_2^\mu + a_5 n_5^\mu$$

DOUBLE CUT

$$l^\mu = V^\mu + a_1 n_1^\mu + a_2 n_2^\mu + a_3 n_3^\mu + a_5 n_5^\mu$$



Diagnosed by **redundant** OPP equation
Solved by picking another set of n-vectors

$$\begin{aligned} \bar{e}_{I_5, s} &= \mathcal{P}|_{l_s}^{I_5} = e_{I, 0} \\ \bar{d}_{I_4, s} &= \mathcal{P}|_{l_s}^{I_4} - \sum_{J_5/I_4} \frac{\bar{e}_{J_5/I_4}}{D_{k_1}(l_s)} = \sum_r d_{I_4, r} f_r(l_s) \\ \bar{c}_{I_3, s} &= \mathcal{P}|_{l_s}^{I_3} - \sum_{J_5/I_3} \frac{\bar{e}_{J_5/I_3}}{D_{k_1}(l_s) D_{k_2}(l_s)} - \sum_{J_4/I_3} \frac{\bar{d}_{J_4/I_3}(l_s)}{D_{k_1}(l_s)} = \sum_r c_{I_3, r} g_r(l_s) \\ \bar{b}_{I_2, s} &= \mathcal{P}|_{l_s}^{I_2} - \sum_{J_5/I_2} \frac{\bar{e}_{J_5/I_2}}{D_{k_1}(l_s) D_{k_2}(l_s) D_{k_3}(l_s)} - \sum_{J_4/I_2} \frac{\bar{d}_{J_4/I_2}(l_s)}{D_{k_1}(l_s) D_{k_2}(l_s)} - \sum_{J_3/I_2} \frac{\bar{c}_{J_3/I_2}(l_s)}{D_{k_1}(l_s)} = \sum_r b_{I_2, r} h_r(l_s) \end{aligned}$$

$\xrightarrow{\hspace{10em}}$

$$\frac{\bar{d}(l_s)}{(l_s + q)^2}$$

There is a special
check for this
implemented

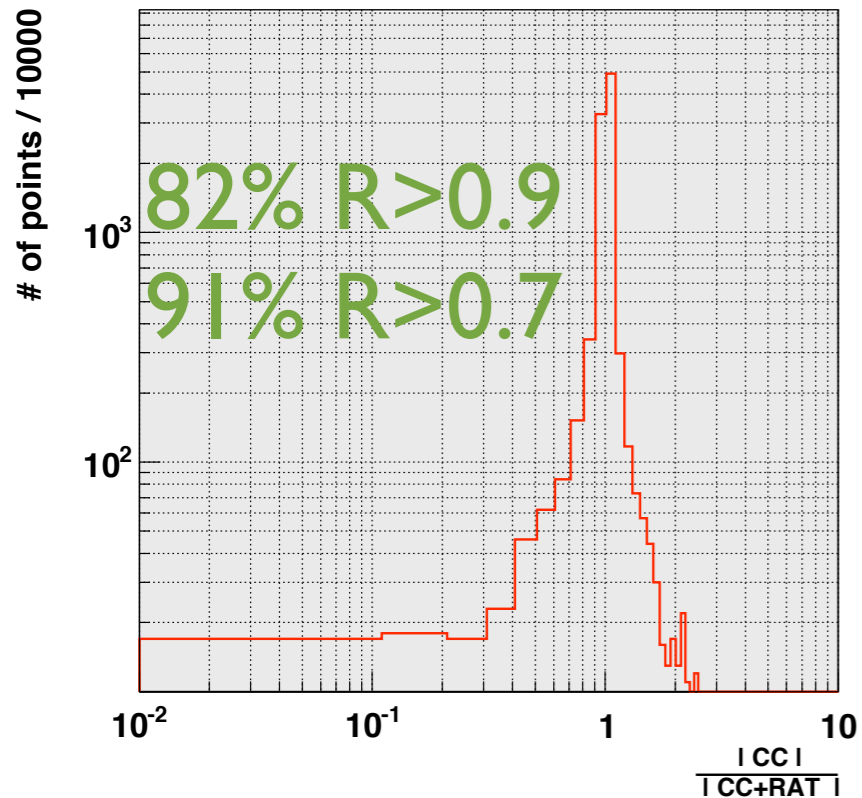
Numerical Stability: pentagon contamination

$$\begin{aligned}
 \bar{e}_{I_{5,s}} &= \mathcal{P}|_{l_s}^{I_5} = e_{I,0} \\
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 \bar{c}_{I_{3,s}} &= \mathcal{P}|_{l_s}^{I_3} - \sum_{J_5/I_3} \frac{\bar{e}_{J_5/I_3}}{D_{k_1}(l_s) D_{k_2}(l_s)} - \sum_{J_4/I_3} \frac{\bar{d}_{J_4/I_3}(l_s)}{D_{k_1}(l_s)} = \sum_r c_{I_{3,r}} g_r(l_s) \\
 \bar{b}_{I_{2,s}} &= \mathcal{P}|_{l_s}^{I_2} - \sum_{J_5/I_2} \frac{\bar{e}_{J_5/I_2}}{D_{k_1}(l_s) D_{k_2}(l_s) D_{k_3}(l_s)} - \sum_{J_4/I_2} \frac{\bar{d}_{J_4/I_2}(l_s)}{D_{k_1}(l_s) D_{k_2}(l_s)} - \sum_{J_3/I_2} \frac{\bar{c}_{J_3/I_2}(l_s)}{D_{k_1}(l_s)} = \sum_r b_{I_{2,r}} h_r(l_s)
 \end{aligned}$$

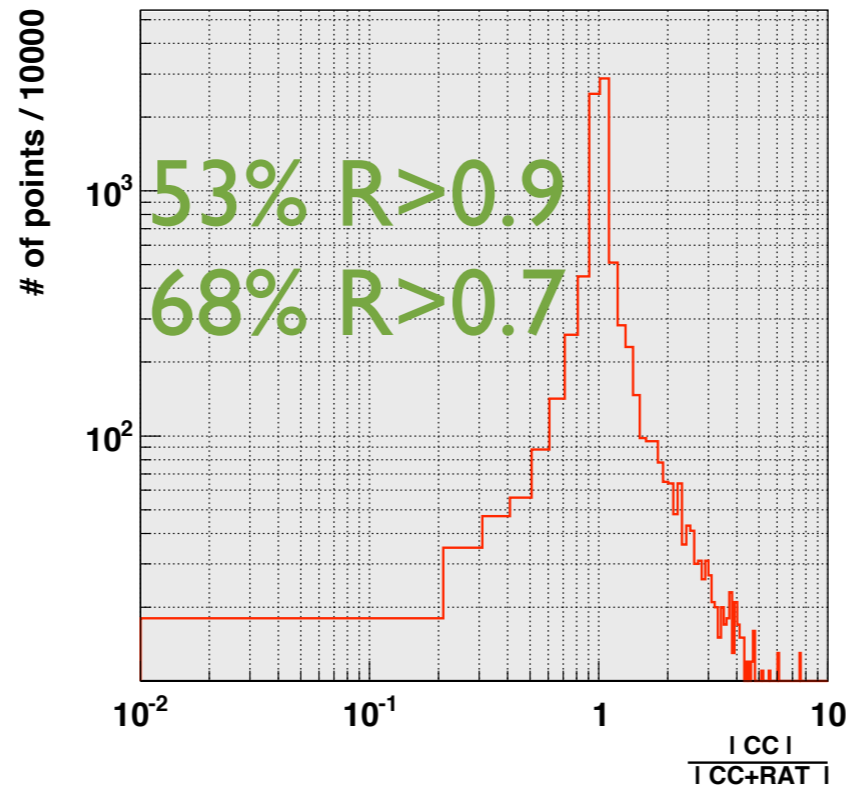
- ✓ Large cancellations in cases of small gram dets.
- ✓ Particularly for the pure gluonic case.
- ✓ Even true for the Cut Constructible part even though it can also be evaluated in 4d, i.e. without pentagons.
- ✓ Would be desirable to remove the pentagon contamination from the CC part.
- ✓ And to control its effect better in Rational.

Numerical Stability: pentagon contamination

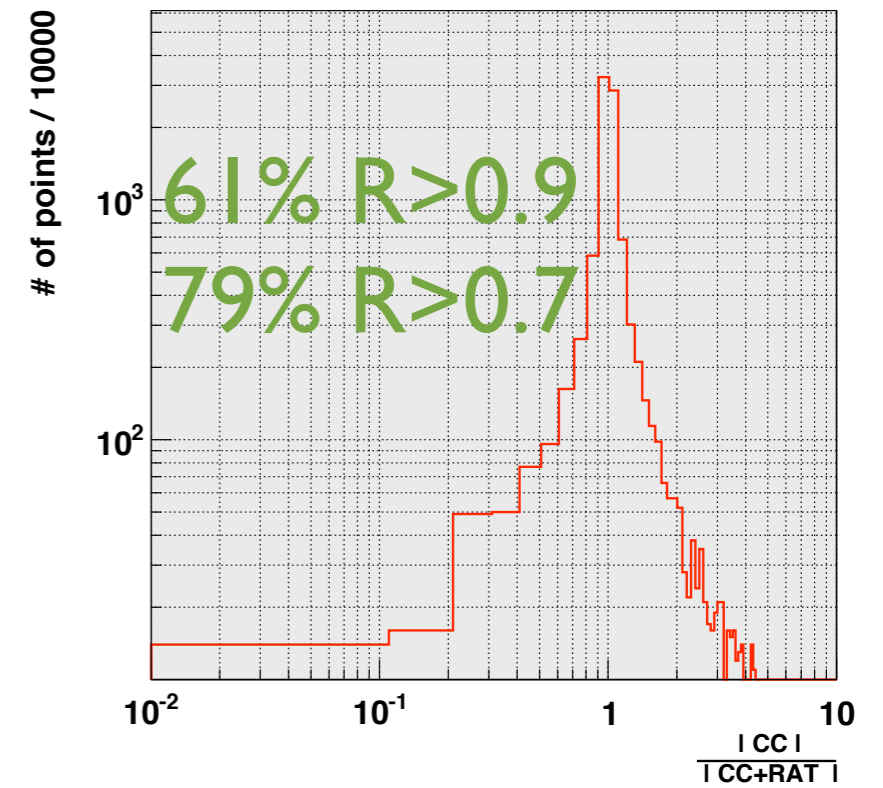
gg->gggg (+++---) cc / finite



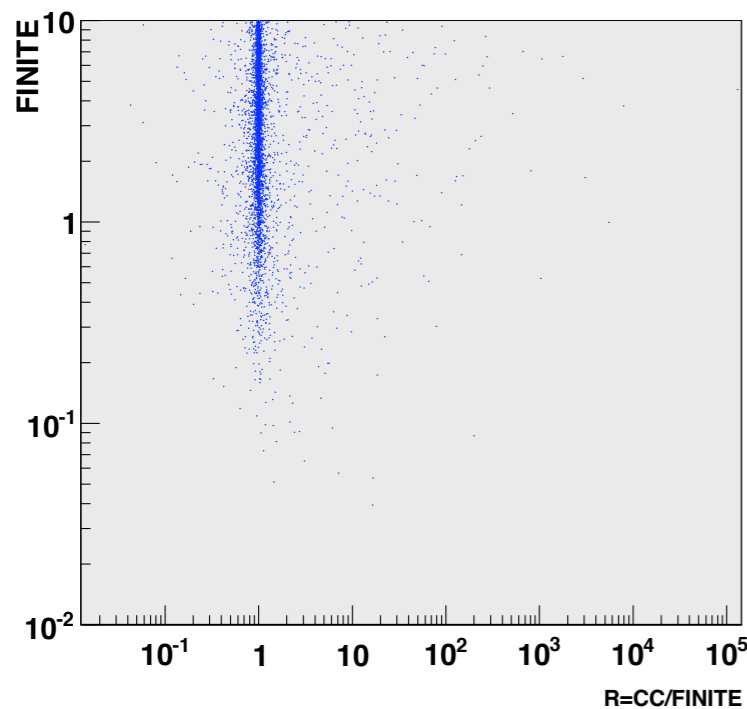
gg->gggg (++++) cc / finite



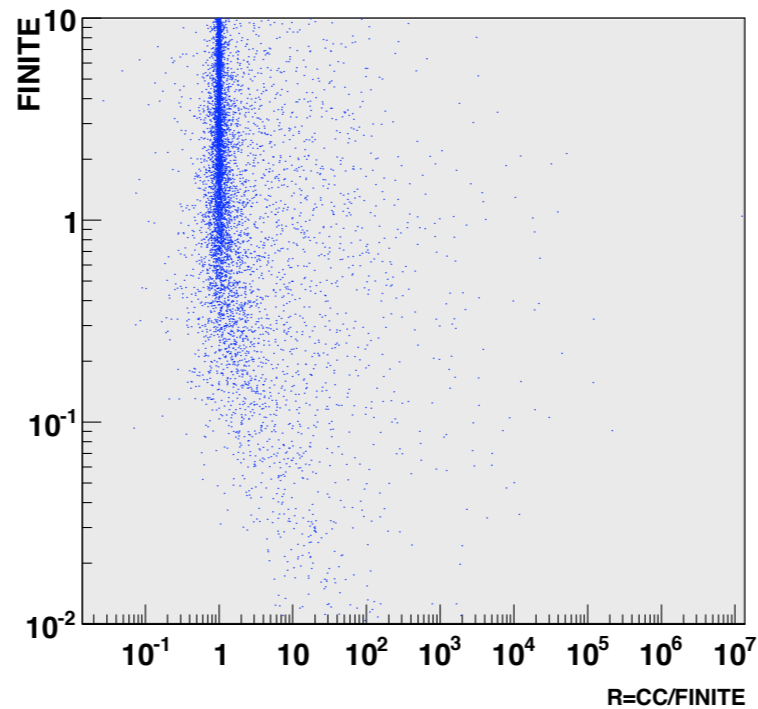
q \bar{q} ->gggg (++---++) cc / finite



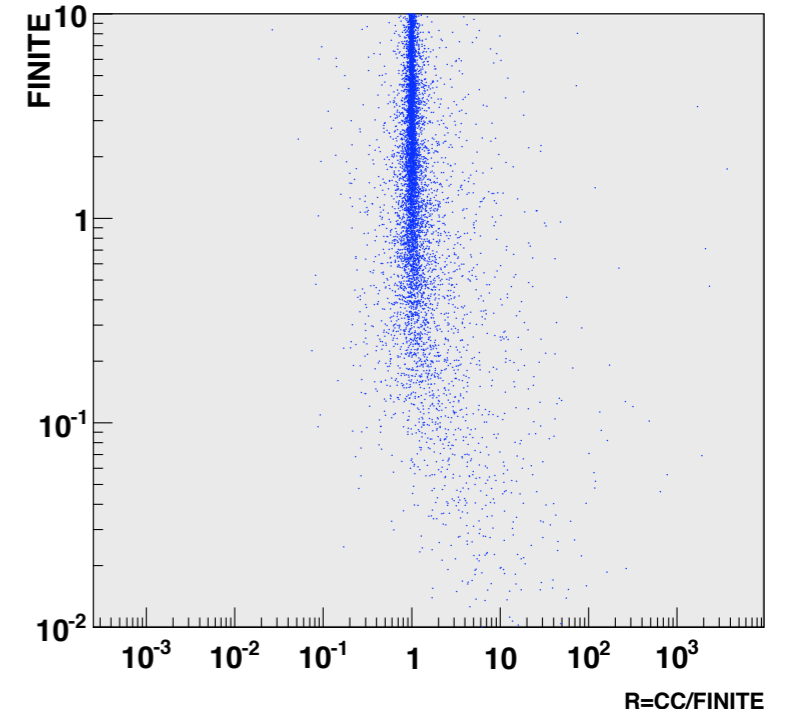
gg->gggg (++++)



gg->gggg (+++---)



q \bar{q} ->gggg (++---++)



Numerical Stability: the cure for pentagon contamination

$$\bar{d}_{I_4,s} = \bar{d}_{I_4,s}^{(1)} + \bar{d}_{I_4,s}^{(2)}$$

$$\bar{d}_{I_4,s}^{(1)} = \mathcal{P}|_{l_s}^{I_4} = \sum_r d_{I_4,r}^{(1)} f_r(l_s)$$

$$\bar{d}_{I_4,s}^{(2)} = - \sum_{J_5/I_4} \frac{\bar{e}_{J_5/I_4}}{D_{k_1}(l_s)} = \sum_r d_{I_4,r}^{(2)} f_r(l_s)$$

$$\bar{c}_{I_3,s} = \bar{c}_{I_3,s}^{(1)} + \bar{c}_{I_3,s}^{(2)}$$

$$\bar{c}_{I_3,s}^{(1)} = \mathcal{P}|_{l_s}^{I_3} - \sum_{J_4/I_3} \frac{\bar{d}_{J_4/I_3}^{(1)}(l_s)}{D_{k_1}(l_s)} = \sum_r c_{I_3,r}^{(1)} g_r(l_s)$$

$$\bar{c}_{I_3,s}^{(2)} = - \sum_{J_5/I_3} \frac{\bar{e}_{J_5/I_3}}{D_{k_1}(l_s) D_{k_2}(l_s)} - \sum_{J_4/I_3} \frac{\bar{d}_{J_4/I_3}^{(2)}(l_s)}{D_{k_1}(l_s)} = \sum_r c_{I_3,r}^{(2)} g_r(l_s)$$

$$\bar{b}_{I_2,s} = \bar{b}_{I_2,s}^{(1)} + \bar{b}_{I_2,s}^{(2)}$$

$$\bar{b}_{I_2,s}^{(1)} = \mathcal{P}|_{l_s}^{I_2} - \sum_{J_4/I_2} \frac{\bar{d}_{J_4/I_2}^{(1)}(l_s)}{D_{k_1}(l_s) D_{k_2}(l_s)} - \sum_{J_3/I_2} \frac{\bar{c}_{J_3/I_2}^{(1)}(l_s)}{D_{k_1}(l_s)} = \sum_r b_{I_2,r}^{(1)} h_r(l_s)$$

$$\bar{b}_{I_2,s}^{(2)} = - \sum_{J_5/I_2} \frac{\bar{e}_{J_5/I_2}}{D_{k_1}(l_s) D_{k_2}(l_s) D_{k_3}(l_s)} - \sum_{J_4/I_2} \frac{\bar{d}_{J_4/I_2}^{(2)}(l_s)}{D_{k_1}(l_s) D_{k_2}(l_s)} - \sum_{J_3/I_2} \frac{\bar{c}_{J_3/I_2}^{(2)}(l_s)}{D_{k_1}(l_s)} = \sum_r b_{I_2,r}^{(2)} h_r(l_s)$$

Numerical Stability: the cure for pentagon contamination



$$\bar{d}_{I_4,s} = \bar{d}_{I_4,s}^{(1)} + \bar{d}_{I_4,s}^{(2)}$$

$$\bar{d}_{I_4,s}^{(1)} = \mathcal{P}|_{l_s}^{I_4} = \sum_r d_{I_4,r}^{(1)} f_r(l_s)$$

$$\bar{d}_{I_4,s}^{(2)} = - \sum_{J_5/I_4} \frac{\bar{e}_{J_5/I_4}}{D_{k_1}(l_s)} = \sum_r d_{I_4,r}^{(2)} f_r(l_s) \quad \times \quad (d_0^{(2)} \text{ cancels against pentagon reduced to boxes})$$

$$\bar{c}_{I_3,s} = \bar{c}_{I_3,s}^{(1)} + \bar{c}_{I_3,s}^{(2)}$$

$$\bar{c}_{I_3,s}^{(1)} = \mathcal{P}|_{l_s}^{I_3} - \sum_{J_4/I_3} \frac{\bar{d}_{J_4/I_3}^{(1)}(l_s)}{D_{k_1}(l_s)} = \sum_r c_{I_3,r}^{(1)} g_r(l_s)$$

$$\bar{c}_{I_3,s}^{(2)} = - \sum_{J_5/I_3} \frac{\bar{e}_{J_5/I_3}}{D_{k_1}(l_s) D_{k_2}(l_s)} - \sum_{J_4/I_3} \frac{\bar{d}_{J_4/I_3}^{(2)}(l_s)}{D_{k_1}(l_s)} = \sum_r c_{I_3,r}^{(2)} g_r(l_s) \quad \times$$

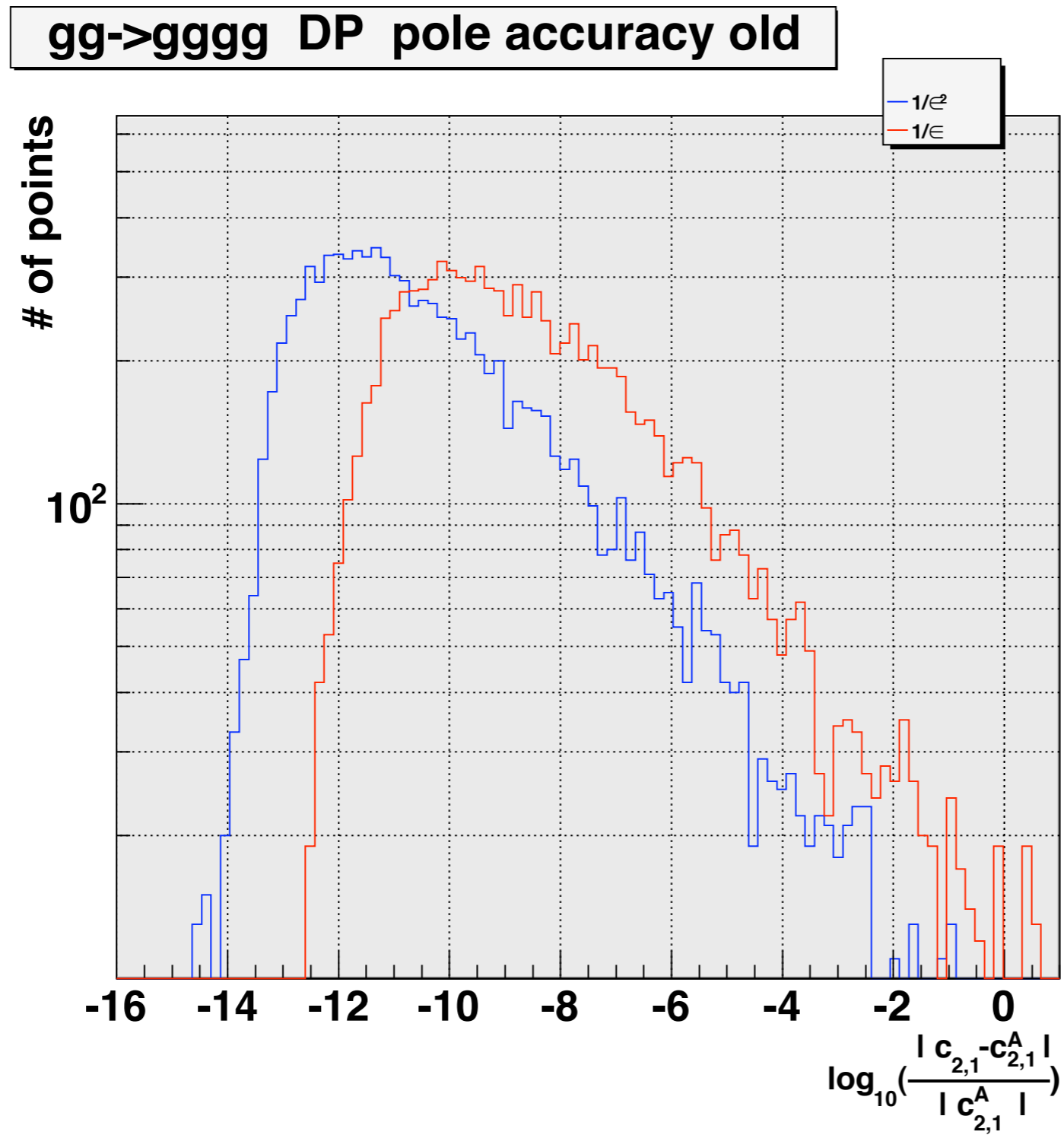
$$\bar{b}_{I_2,s} = \bar{b}_{I_2,s}^{(1)} + \bar{b}_{I_2,s}^{(2)}$$

$$\bar{b}_{I_2,s}^{(1)} = \mathcal{P}|_{l_s}^{I_2} - \sum_{J_4/I_2} \frac{\bar{d}_{J_4/I_2}^{(1)}(l_s)}{D_{k_1}(l_s) D_{k_2}(l_s)} - \sum_{J_3/I_2} \frac{\bar{c}_{J_3/I_2}^{(1)}(l_s)}{D_{k_1}(l_s)} = \sum_r b_{I_2,r}^{(1)} h_r(l_s)$$

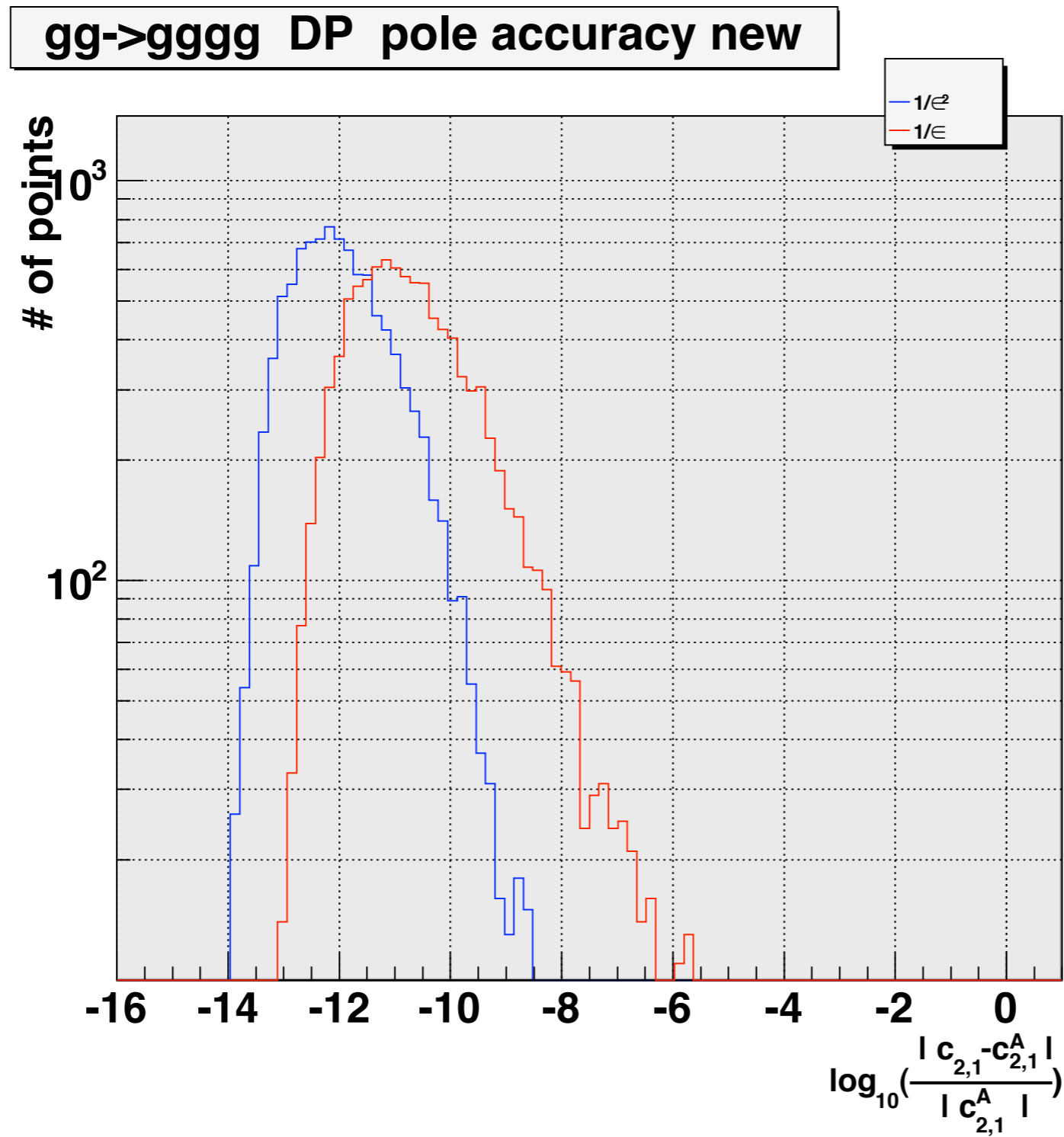
$$\bar{b}_{I_2,s}^{(2)} = - \sum_{J_5/I_2} \frac{\bar{e}_{J_5/I_2}}{D_{k_1}(l_s) D_{k_2}(l_s) D_{k_3}(l_s)} - \sum_{J_4/I_2} \frac{\bar{d}_{J_4/I_2}^{(2)}(l_s)}{D_{k_1}(l_s) D_{k_2}(l_s)} - \sum_{J_3/I_2} \frac{\bar{c}_{J_3/I_2}^{(2)}(l_s)}{D_{k_1}(l_s)} = \sum_r b_{I_2,r}^{(2)} h_r(l_s) \quad \times$$

✓ Pentagon decoupling from 4-D!

Numerical Stability: improved CC plot



Numerical Stability: improved CC plot



Numerical Stability: the cure for pentagon contamination

5D

$$\bar{d}_{I_4,s} = \bar{d}_{I_4,s}^{(1)} + \bar{d}_{I_4,s}^{(2)}$$

$$\bar{d}_{I_4,s}^{(1)} = \mathcal{P}|_{l_s}^{I_4} = \sum_r d_{I_4,r}^{(1)} f_r(l_s)$$

$$\bar{d}_{I_4,s}^{(2)} = - \sum_{J_5/I_4} \frac{\bar{e}_{J_5/I_4}}{D_{k_1}(l_s)} = \sum_r d_{I_4,r}^{(2)} f_r(l_s)$$

! can be handled with care

$$\bar{c}_{I_3,s} = \bar{c}_{I_3,s}^{(1)} + \bar{c}_{I_3,s}^{(2)}$$

$$\bar{c}_{I_3,s}^{(1)} = \mathcal{P}|_{l_s}^{I_3} - \sum_{J_4/I_3} \frac{\bar{d}_{J_4/I_3}^{(1)}(l_s)}{D_{k_1}(l_s)} = \sum_r c_{I_3,r}^{(1)} g_r(l_s)$$

$$\bar{c}_{I_3,s}^{(2)} = - \sum_{J_5/I_3} \frac{\bar{e}_{J_5/I_3}}{D_{k_1}(l_s) D_{k_2}(l_s)} - \sum_{J_4/I_3} \frac{\bar{d}_{J_4/I_3}^{(2)}(l_s)}{D_{k_1}(l_s)} = \sum_r c_{I_3,r}^{(2)} g_r(l_s)$$

! can be handled with care

$$\bar{b}_{I_2,s} = \bar{b}_{I_2,s}^{(1)} + \bar{b}_{I_2,s}^{(2)}$$

$$\bar{b}_{I_2,s}^{(1)} = \mathcal{P}|_{l_s}^{I_2} - \sum_{J_4/I_2} \frac{\bar{d}_{J_4/I_2}^{(1)}(l_s)}{D_{k_1}(l_s) D_{k_2}(l_s)} - \sum_{J_3/I_2} \frac{\bar{c}_{J_3/I_2}^{(1)}(l_s)}{D_{k_1}(l_s)} = \sum_r b_{I_2,r}^{(1)} h_r(l_s)$$

$$\bar{b}_{I_2,s}^{(2)} = - \sum_{J_5/I_2} \frac{\bar{e}_{J_5/I_2}}{D_{k_1}(l_s) D_{k_2}(l_s) D_{k_3}(l_s)} - \sum_{J_4/I_2} \frac{\bar{d}_{J_4/I_2}^{(2)}(l_s)}{D_{k_1}(l_s) D_{k_2}(l_s)} - \sum_{J_3/I_2} \frac{\bar{c}_{J_3/I_2}^{(2)}(l_s)}{D_{k_1}(l_s)} = \sum_r b_{I_2,r}^{(2)} h_r(l_s)$$

!

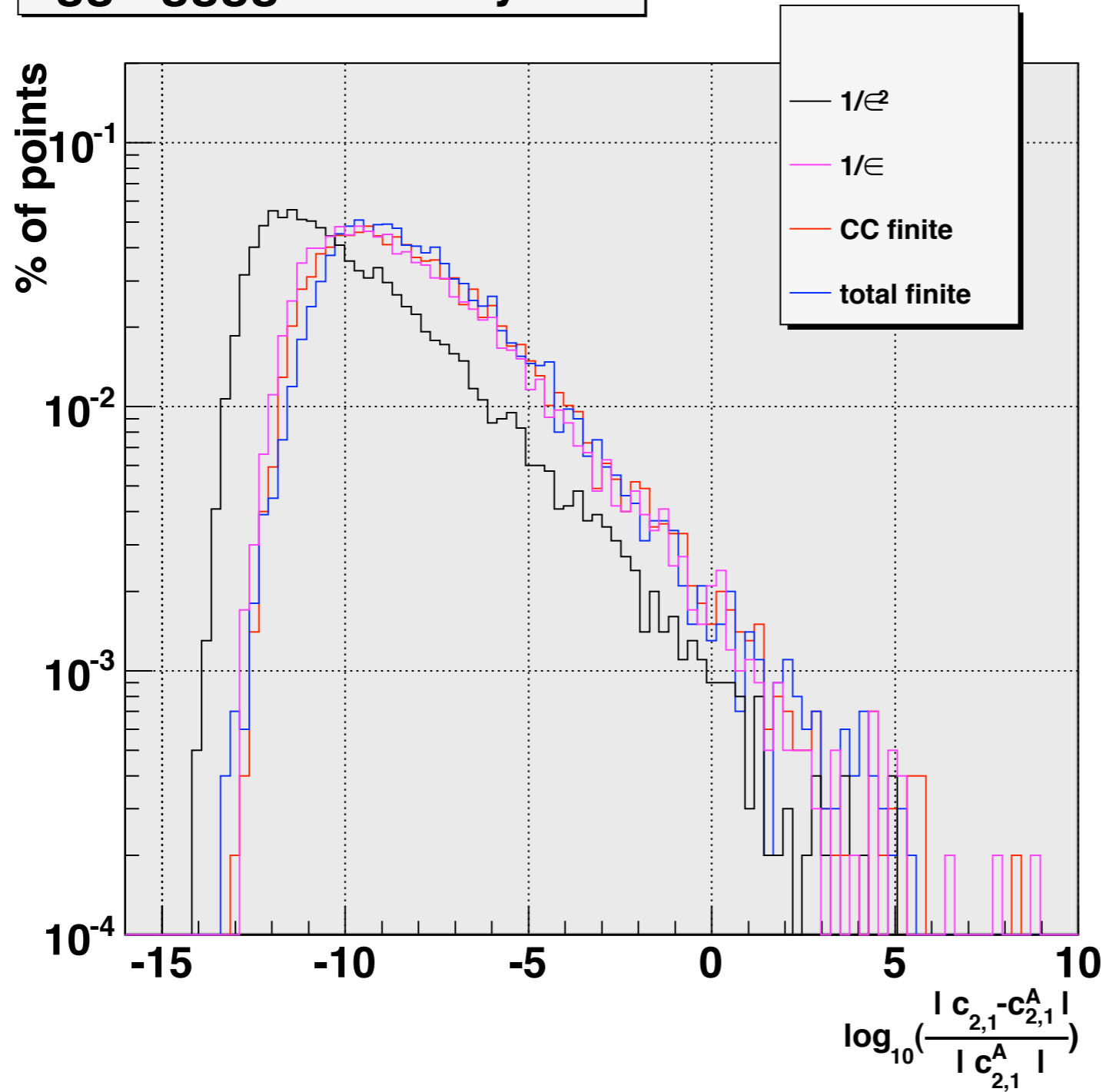
✓ Both subsystems contribute to the rational part

Numerical Stability: the cure for pentagon contamination

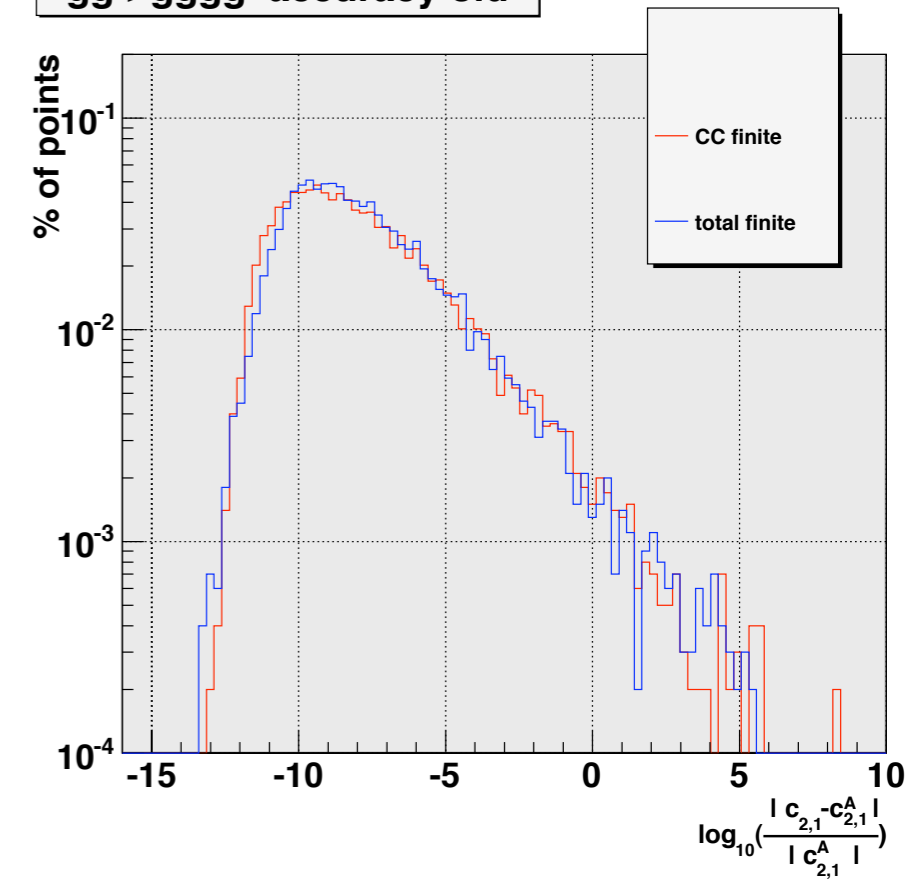
“The pentagon coefficient should always be factored out of any subtractions”

Numerical Stability: improved Finite plot

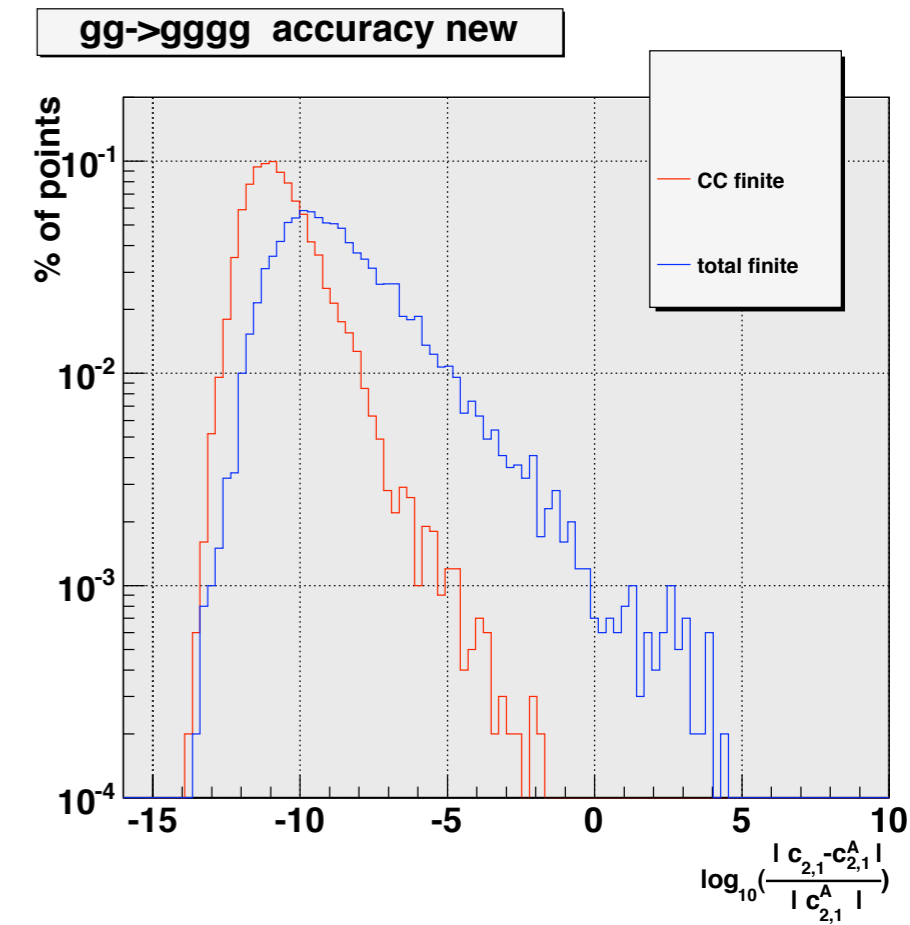
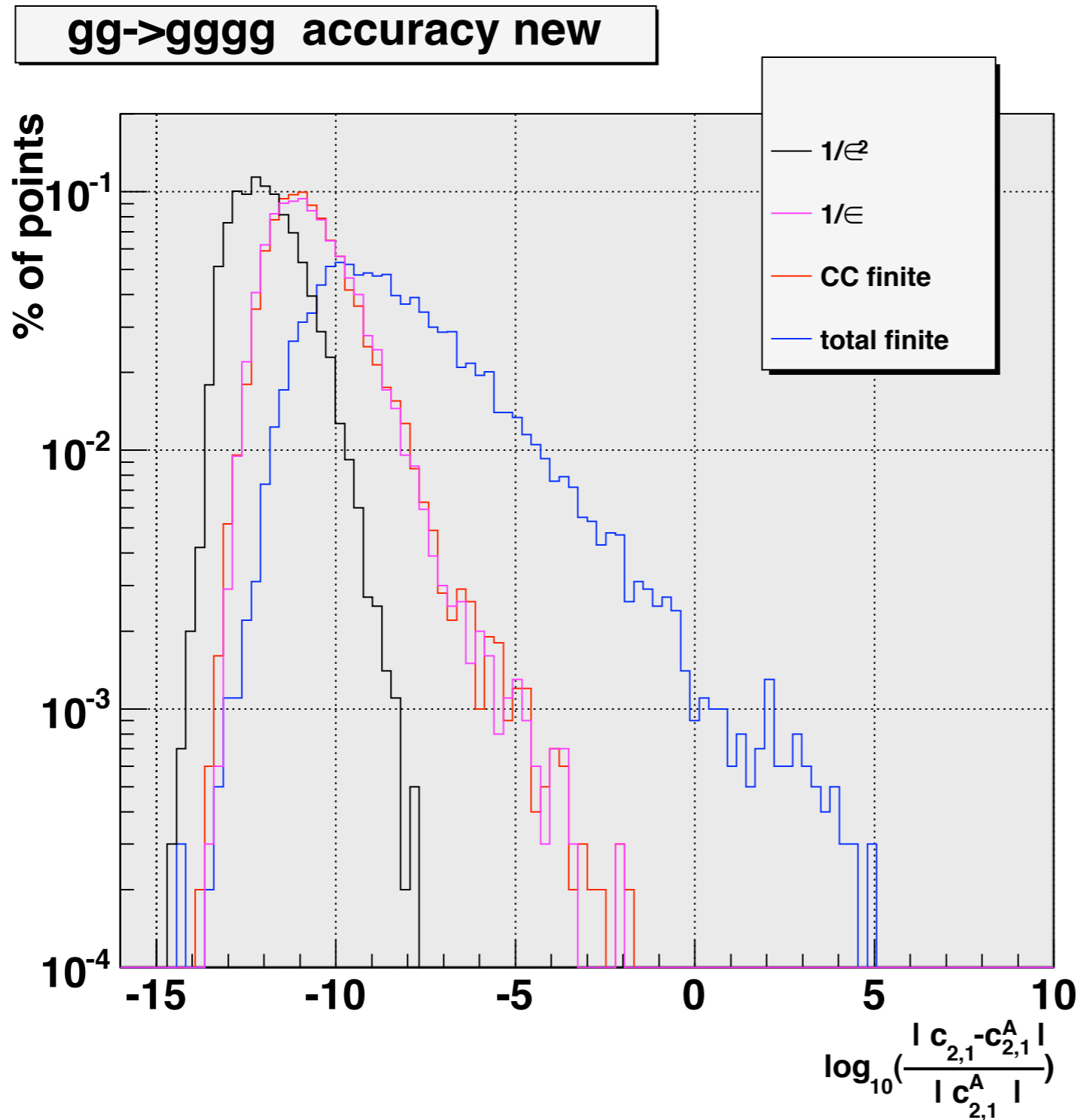
gg->gggg accuracy old



gg->gggg accuracy old

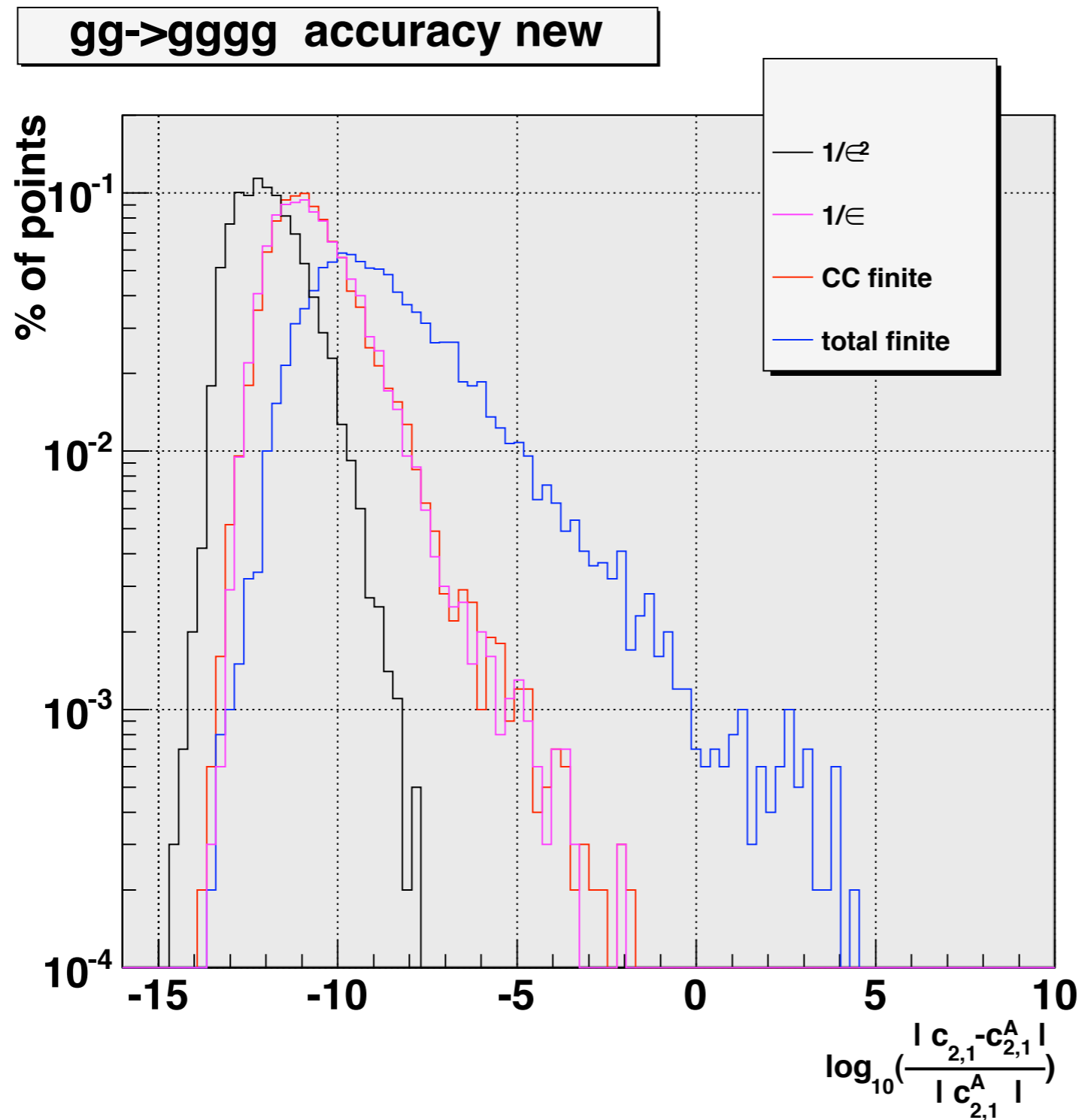


Numerical Stability: improved Finite plot



Splitting in two subsystems.

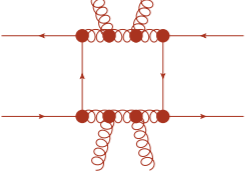
Numerical Stability: improved Finite plot



Improving the way the box coefficients are treated in subsystem 2

Performance

Performance table

	N=6	N=7	N=8
	50ms	141ms	350ms
	57ms	153ms	380ms
	60ms	155ms	373ms
	59ms	157ms	373ms
	60ms	152ms	369ms

QUAD
PENALTY
x100

“Thanks to improved accuracy, quadruple precision is only called rarely, which decreases drastically the realistic cpu time per PSP”.

Summary

Summary

- We can now do all primitives necessary for massless QCD partonic processes, including four- and six- fermion subprocesses.
- Numerical instability issues due to pentagon contamination are removed from the CC part and controlled much better at the RAT part.
- Ready for production mode.