

Generalised Unitarity For Massive One-Loop Amplitudes

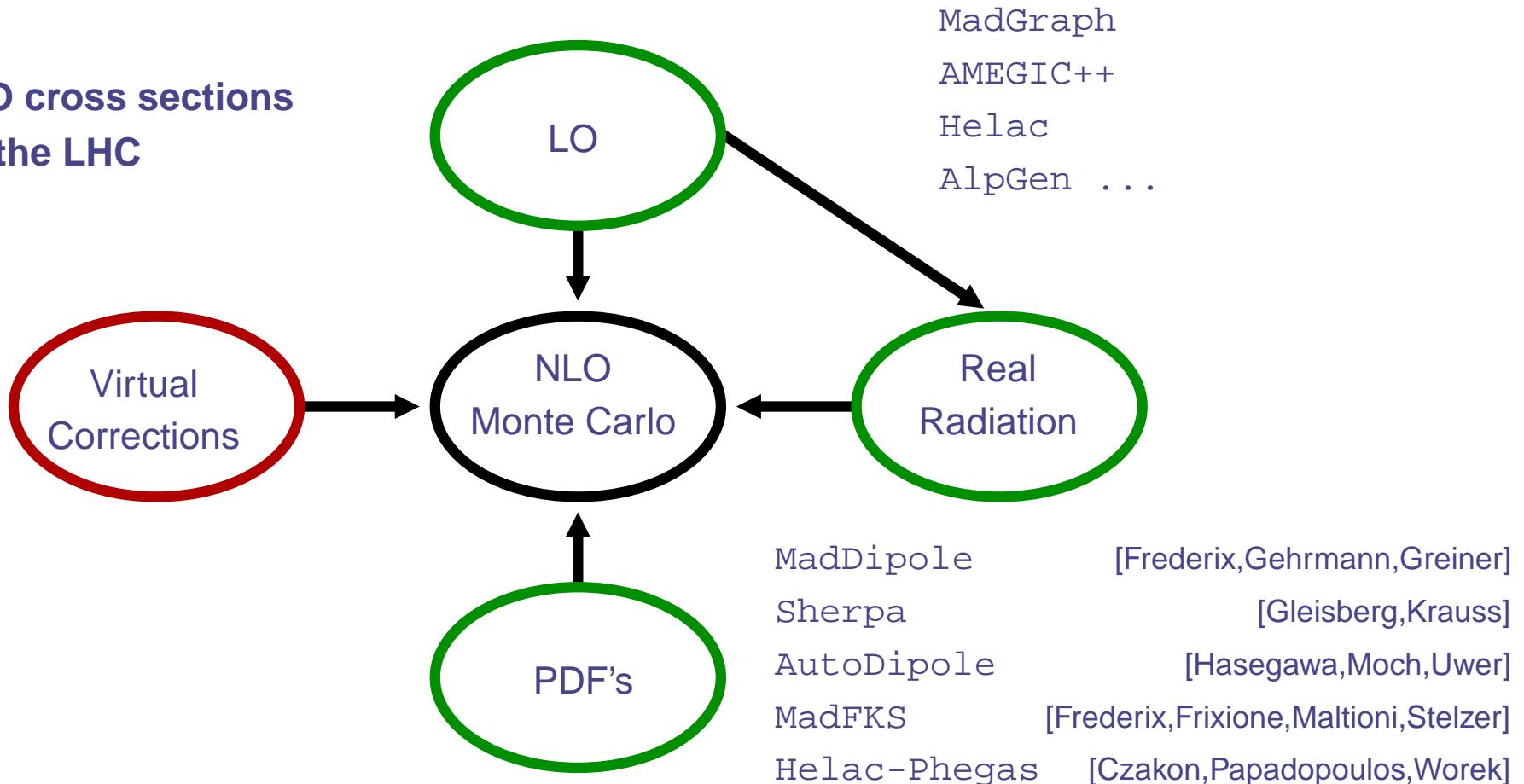
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27th October 2009

Radcor, Ascona, Switzerland

QCD for the LHC

NLO cross sections
for the LHC



POWHEG

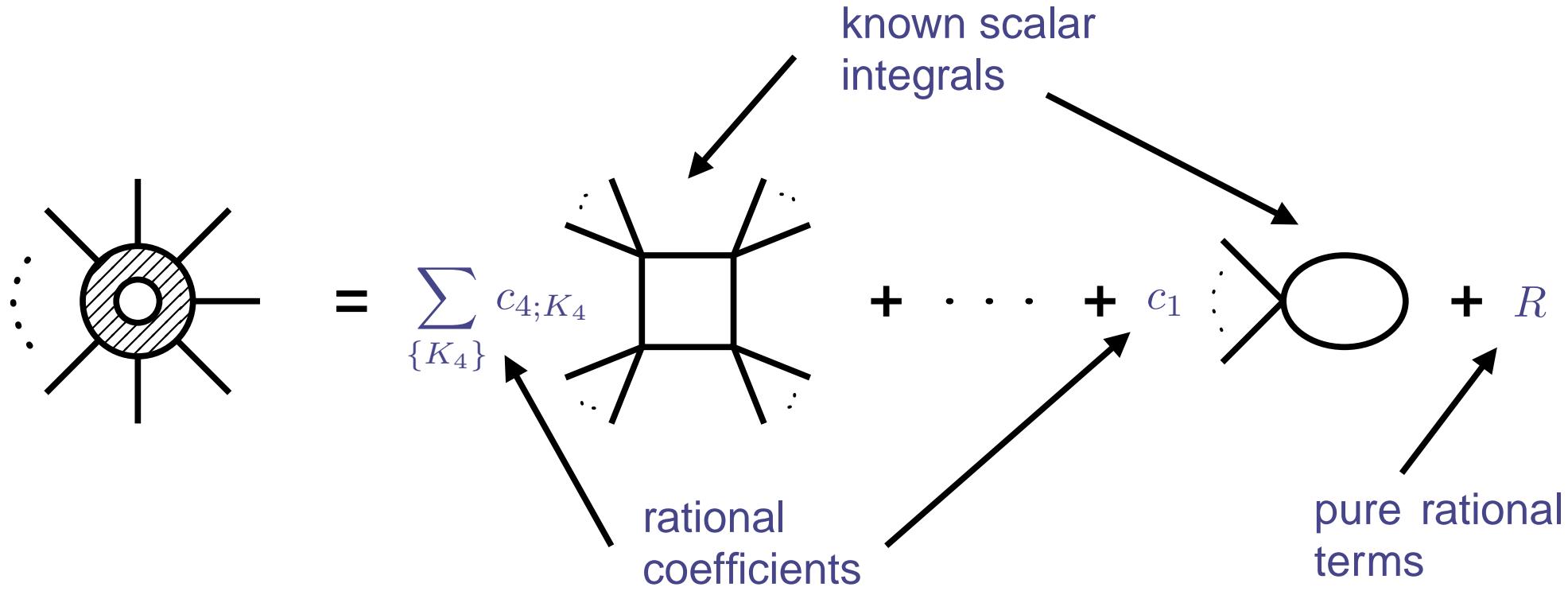
[Nason et al]

MC@NLO [Frixione, Webber]

Computations of Virtual Corrections

- A lot of recent progress in computational methods for virtual corrections:
Bern,Dixon,Dunbar,Kosower,Britto,Cachazo,Feng,Mastrolia,
Ossola,Papadopoulos,Pittau,Ellis,Giele,Kunszt,Melnikov,Forde,...
- Automated numerical approaches:
[BlackHat, Rocket, CutTools/Helac-1loop, GOLEM, Denner et al., ...]
- Efficiency:
 - Numerical stability
 - Fast numerical evaluation
 - On-shell simplifications
- Rest of this talk: Concentrate on analytic approach to generalised unitarity

Structure of One-Loop Amplitudes

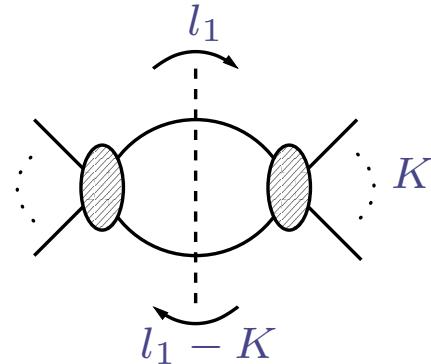


- General gauge theory amplitudes reduced to box topologies or simpler
 - [Passarino,Veltman;Melrose]
- Isolate logarithms with cuts → exploit on-shell simplifications
- General cutting principle:
 - apply δ -functions to L and R sides
 - generate and solve the linear system for the coefficients

Unitarity

- Cutkosky Rules to compute $\text{Disc}_p(A^{\text{1-loop}})$

$$\frac{1}{p^2 - m^2 + iO^+} \rightarrow i\delta^{(+)}(p^2 - m^2)$$



- Construction of amplitudes from double cuts
- Improved fitting with triple cuts
- Generalised cuts with complex momenta
- D -dimensional cuts can also access rational contributions

[Bern,Dixon,Dunbar,Kosower (1994)]

[Bern,Dixon,Kosower (1997)]

[Britto,Cachazo,Feng (2004)]

[Bern,Morgan (1995)]

[Bern,Dixon,Dunbar,Kosower (1997)]

[Giele,Kunszt,Melnikov;Anastasiou,Britto,Feng,Mastrolia,Kunszt;SB]

Massive Integral Basis

$$C_4 = \prod_{i=1}^4 A_i(l) \quad C_3 = \lim_{t \rightarrow \infty} \prod_{i=1}^3 A_i(l(t))|_{t^0} \quad C_2 \sim \lim_{y/t \rightarrow \infty} \prod_{i=1}^2 A_i(l(y, t))$$

$A_n^{(1)} =$

+

$\log(m^2)$

+

products of trees
→ integral coefficients

Rational

+

+

+

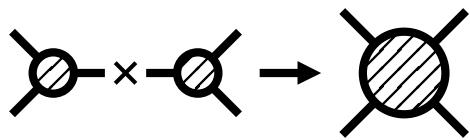
$$C_4^{[4]} = \lim_{\mu^2 \rightarrow \infty} \prod_{i=1}^4 A_i(l)|_{\mu^4} \quad C_3^{[2]} = \lim_{\mu^2, t \rightarrow \infty} \prod_{i=1}^3 A_i(l(t))|_{t^0, \mu^2} \quad C_2^{[2]} \sim \lim_{\mu^2, y/t \rightarrow \infty} \prod_{i=1}^2 A_i(l(y, t))$$

Tree-Level Amplitudes

- BCFW recursion generates compact analytic expressions

[Britto,Cachazo,Feng,Witten (2005)]

- Application of Cauchy's theorem and factorisation



$$0 = \oint A(z)/z = \sum_{\text{residues}} A(z)/z = \sum A_L \frac{1}{P^2} A_R + A(0)$$

- Simple generalisation to include internal masses [SB,Glover,Khoze,Svrček (2005)]
- Helicity basis for massive fermions for compact analytic formulae:

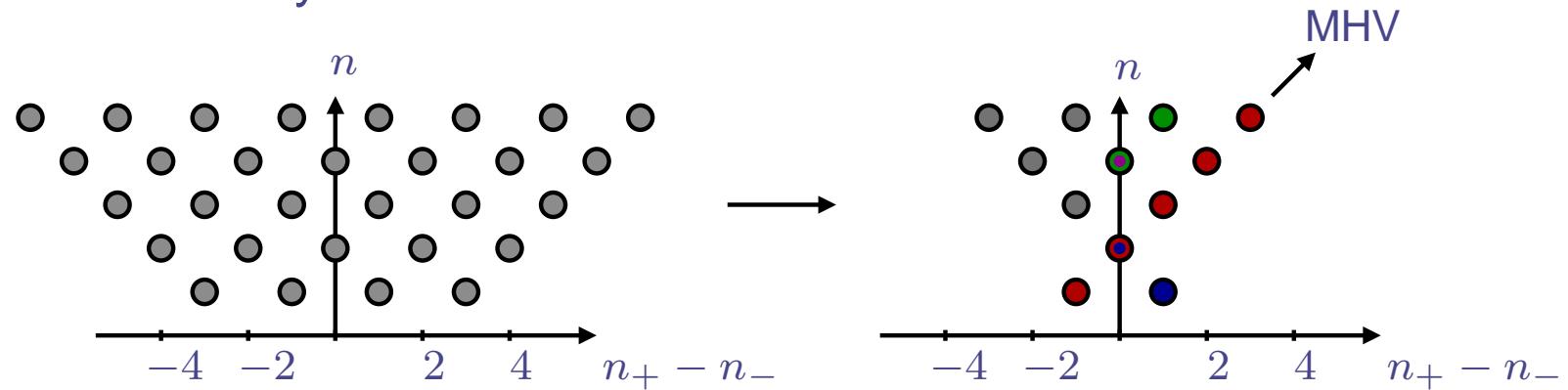
$$u_{\pm}(p, m; p^b, \eta) = \frac{(\not{p} + m)|\eta\mp\rangle}{\langle p^b \pm |\eta\mp\rangle}$$

[Kleiss,Stirling]

Gauge theories defined from basis of three-point on-shell vertices

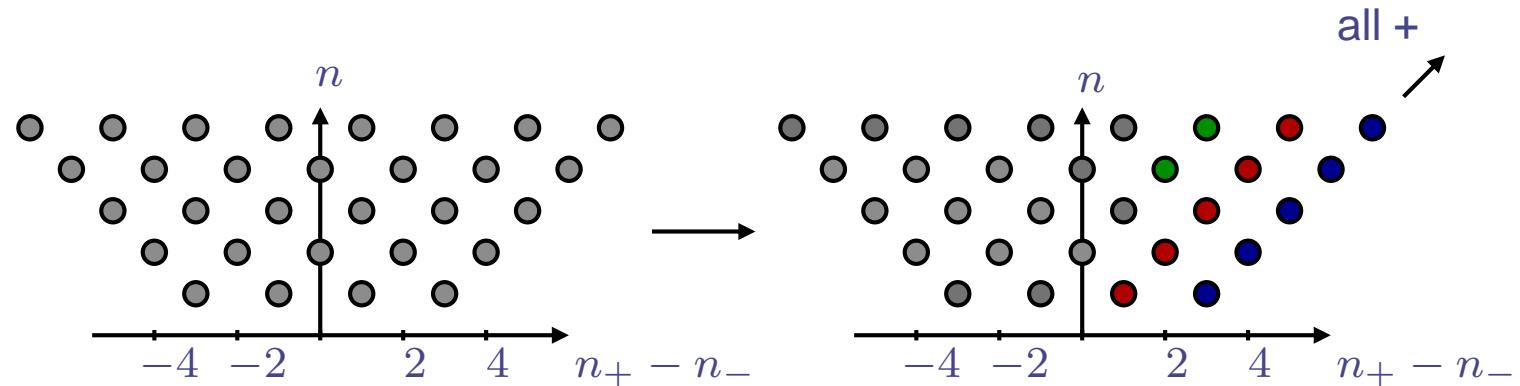
Massive Helicity Amplitudes

Massless QCD helicity structure:



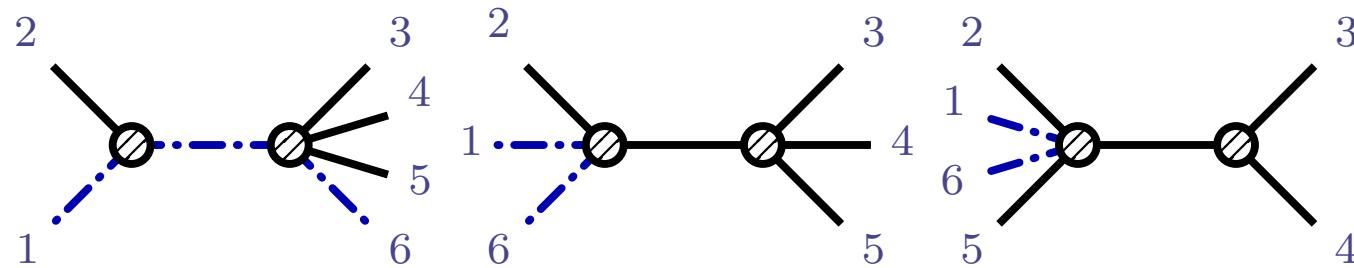
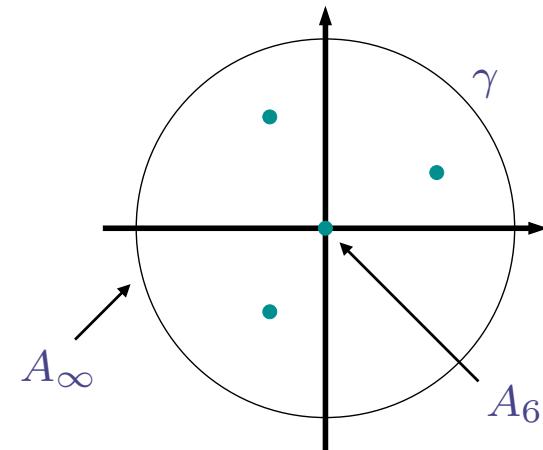
Symmetry relates massive quark helicity states:

$$u_-(p, m; p^\flat, \eta) = \frac{\langle p^\flat \eta \rangle}{m} u_+(p, m; \eta, p^\flat)$$



Tree-level Amplitudes: $t\bar{t} + 4g$

- Complex continuation
 $(p_2, p_3) \rightarrow (p_2(z), p_3(z))$
- Sum over residues and internal helicities



$$\langle ab \rangle \sim \sqrt{a.b}$$

$$A_6(1_t^+, 2^+, 3^+, 4^+, 5^+, 6_{\bar{t}}^+) = \frac{i m^3 \langle \eta_6 \eta_1 \rangle [5|6P_{45}P_{23}1|2]}{\langle 1^b \eta_1 \rangle \langle 6^b \eta_6 \rangle \langle 23 \rangle \langle 34 \rangle \langle 45 \rangle \langle 2|1|2] \langle 5|6|5] (s_{123} - m^2)}$$

Tree-level Amplitudes: $t\bar{t} + n(g)$

For this "all-plus" configuration we can look at higher point functions:

$$A_7(1_t^+, 2^+, 3^+, 4^+, 5^+, 6^+, 7_{\bar{t}}^+) = \frac{i m^3 \langle \eta_7 \eta_1 \rangle ((s_{1,3} - m^2) [6|7P_{56}P_{2,4}1|2] - m^2 [6|7P_{56}43|2])}{\langle 1^\flat \eta_1 \rangle \langle 7^\flat \eta_7 \rangle \prod_{\alpha=2}^4 \langle \alpha \alpha + 1 \rangle (s_{1,\alpha} - m^2)}$$

$$A_8(1_t^+, 2^+, 3^+, 4^+, 5^+, 6^+, 7^+, 8_{\bar{t}}^+) = \frac{i m^3 \langle \eta_8 \eta_1 \rangle}{\langle 1^\flat \eta_1 \rangle \langle 8^\flat \eta_8 \rangle \prod_{\alpha=2}^5 \langle \alpha \alpha + 1 \rangle (s_{1,\alpha} - m^2)} \times \\ \left((s_{1,4} - m^2) ((s_{1,3} - m^2) [7|8P_{67}P_{2,5}1|2] + m^2 [7|8P_{67}P_{45}3|2]) - m^2 [7|8P_{67}54P_{23}1|2] \right)$$

- All-order formula possible

[Forde,Kosower]

[Schwinn,Weinzierl]

[Ferrario,Rodrigo,Talavera]

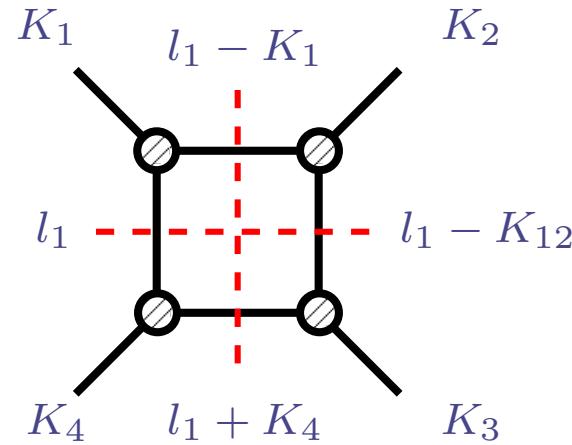
Analytic Solutions for Quadruple Cuts

- Quadruple cut freezes loop momentum in 4 dimensions

[BCF]

- Two complex solutions to on-shell constraints

$$\{(l_1 - K_{1,i-1})^2 = m_i^2\}$$



- Parametrise loop momentum with two projected massless vectors:

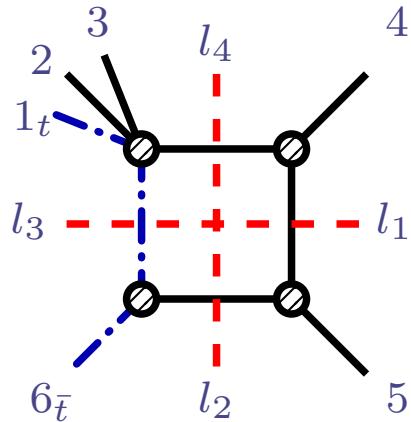
[Forde;Ossola,Papaodopoulos,Pittau]

$$l^\nu = a(K_1^\flat)^\nu + b(K_2^\flat)^\nu + c \langle K_1^\flat | \gamma^\nu | K_2^\flat \rangle + \left(d - \frac{m_1^2}{2K_1^\flat \cdot K_2^\flat} \right) \langle K_2^\flat | \gamma^\nu | K_1^\flat \rangle$$

- Sum over solutions to obtain integral coefficient

$$C_{4;K_1|K_2|K_3|K_4} = \frac{1}{2} \sum_{\sigma} A_1 A_2 A_2 A_4(l^\sigma)$$

Application to $t\bar{t} + 4g$



$$\{l_1^2 = 0, (l_1 - 5)^2 = 0, (l_1 - 5 - 6)^2 = m^2, (l + 4)^2 = 0\}$$

$$l_1^{(1)} = \frac{\langle 5|6|5]}{2\langle 4|6|5]}\langle 4|\gamma^\mu|5], \quad l_1^{(2)} = \frac{\langle 5|6|5]}{2\langle 5|6|4]}\langle 5|\gamma^\mu|4]$$

$$C_{4;4|5|6|123}^{[L]}(1_t^+, 2^+, 3^+, 4^+, 5^+, 6_{\bar{t}}^+) =$$

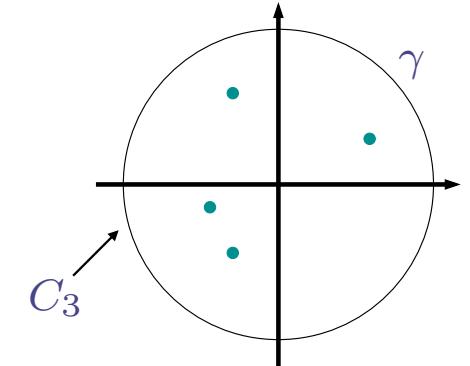
$$\frac{im^3\langle\eta_6\eta_1\rangle[54]}{\langle\eta_66^\flat\rangle\langle 1^\flat\eta_1\rangle\langle 23\rangle\langle 3P_{45}65\rangle\langle 2|1|2]} \left(\frac{s_{45}[4|P_{23}1|2]m^2}{(s_{123} - m^2)} + \frac{\langle 5|6|5]\langle 3|P_{45}P_{16}1|2]}{2\langle 34\rangle} \right)$$

Other configurations straightforward and relatively compact

Automation : BCFW to box expressions (`FORM`, `maple`)

Triple and Double Cuts

- Lower point coefficients have unfixed integrations
- Analytic extraction via Laurent Expansion
 - Complex analysis again: Residue theorem



$$C_3 = \sum_{\sigma} \text{Inf}_t [A_1 A_2 A_3(l^{\sigma}(t))] \Big|_{t^0}$$

- Bubble coefficients also reduced to algebraic procedure

[Forde]

$$C_2 \sim \text{Inf}_{y,t} [A_1 A_2(l(y,t))] + \sum_{y_{\pm}} \text{Inf}_t [A_1 A_2 A_3(l(y_{\pm}, t))]$$

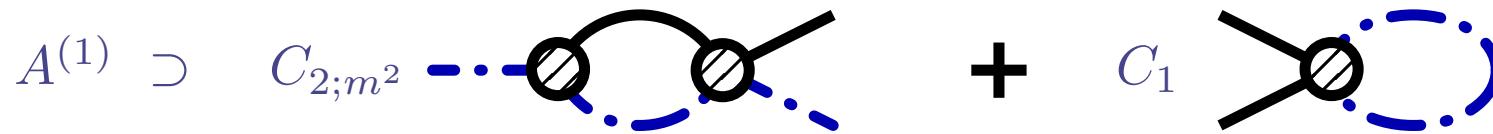
- Analysis of double cut: Stokes' Theorem
- Straightforward generalisation to massive case

[Mastrolia]

[Kilgore; SB]

Automation : closed analytic expressions (FORM, maple)

Wave-function Renormalisation and Tadpoles



- More remaining logarithms - $\log(m^2)$

$$A^{(1)} \supset \frac{C_{1/\epsilon}}{\epsilon} \left(\frac{\mu_R^2}{m^2} \right)^\epsilon + \mathcal{O}(\epsilon^0)$$

- Double cut contains mass renormalisation divergence [Ellis,Giele,Kunszt,Melnikov]
- $C_{1/\epsilon}$ fixed from universal IR behaviour, e.g. [Catani,Dittmaier,Trócsányi] [Mitov,Moch]

$$C_{1/\epsilon}^{[L]}(t\bar{t} + n(g)) = \frac{1}{4} A_n^{(0)} - \sum_{\{t,\bar{t}\} \in K_2} C_{2;K_2}^{[L]}$$

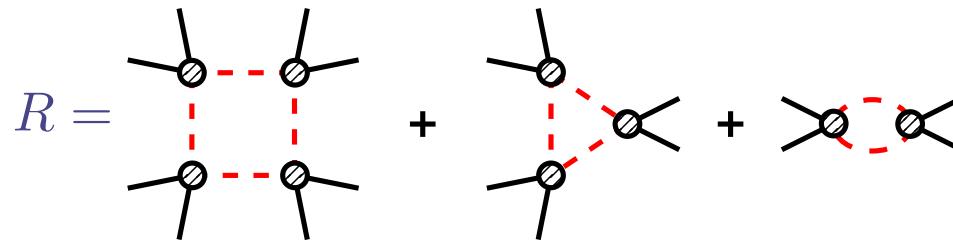
- Direct unitarity solution for tadpole (rational only) [Kilgore;Britto,Feng]

Rational Terms

- Simpler example: six-gluon amplitudes
- Analytic D-dimensional cuts

• $4 - 2\epsilon$ -dim. cuts \rightarrow massive cuts $\int d^{4-2\epsilon} l = \int d^{-2\epsilon} \mu \int d^4 \bar{l}$

[Giele,Kunszt,Melnikov]



• Extraction of rational terms from Taylor expansion in μ^2

[SB]

$$\begin{aligned} R = & -\frac{1}{6} \sum_{\{K_4\}} \text{Inf}_{\mu^2} \left[\prod_{i=1}^4 A_i \right] |_{\mu^4} - \frac{1}{2} \sum_{\{K_3\}} \text{Inf}_{\mu^2; t} \left[\prod_{i=3}^4 A_i \right] |_{\mu^2, t^0} \\ & - \sum_{\{K_2\}} \frac{K_2 - 3(m_1^2 + m_2^2)}{6} \left(\text{Inf}_{\mu^2; y; t} \left[\prod_{i=2}^4 A_i \right] |_{\mu^2, t^0, Y_i} + \text{Inf}_{\mu^2; t} \left[\prod_{i=3}^4 A_i(y(t)) \right] |_{\mu^2, T_i, Y_i} \right) \end{aligned}$$

Rational Terms

Example six-gluon all-plus

$$R_6^g(1^+, 2^+, 3^+, 4^+, 5^+, 6^+) = \sum_{\sigma \in S_n} \left(\frac{2i(s_{45}\langle 6|1+2|3][51][64]^2 - s_{46}\langle 5|1+2|3][54]^2[61]) [56]}{\langle 12\rangle\langle 23\rangle\text{tr}_5(5,4,6,1)\text{tr}_5(5,4,6,3)} \right. \\ + \frac{2i\langle 5|1+2|6]\langle 6|1+2|5][12][43][65]^2}{\langle 12\rangle\langle 34\rangle\text{tr}_5(5,2,6,1)\text{tr}_5(5,4,6,3)} \\ \left. + \frac{2i(\langle 3|1+2|3]\langle 6|1+2|6] - s_{36}s_{12}) [12][54][63]^2}{\langle 12\rangle\langle 45\rangle\text{tr}_5(2,3,6,1)\text{tr}_5(5,3,6,4)} \right)$$

- Closed form expressions suitable numerical evaluation
- All helicity configurations available
- Compact? Compare with on-shell recursion

[Berger,Bern,Dixon,Forde,Kosower]

Conclusions and Outlook

- Efficient constructions of one-loop amplitudes with generalised unitarity
- Simple extension to massive case [Britto,Feng,Mastrolia;Kilgore]
- Computation of scalar integral coefficients purely algebraic
- Complexity scales well with additional external legs
- Automated numerical tools
 - BlackHat , Rocket , CutTools/Helac-1loop ,
Lazopoulos , Giele+Winter , Melnikov+Schulze,...
- Analytic expressions: Speed, Stability,... needs further investigation
- Massive amplitudes contain more scales! Need to employ efficient tools
 - SM backgrounds to BSM signal @ LHC

A Bit of Numerical Analysis

- Test case: evaluation of $H + 2j$ virtual amplitudes: [Ciaran Williams' Talk]
 - Computation in $m_t \rightarrow \infty$ limit: Effective Higgs-gluon interaction
 - [739 Feynman diagrams: 12 boxes, 24 triangles, 8 bubbles]
 - [Existing semi-numerical evaluation: Ellis,Giele,Zanderighi (2006)]
 - Compact expressions for complete set of helicity amplitudes
 - $gggg$ channels: [Berger,Del-Duca,Dixon;SB,Glover,Risager,Mastrolia,Williams]
 - $q\bar{q}gg$ channels: [Dixon,Sofianatos;SB,Campbell,Ellis,Williams]
 - $q\bar{q}Q\bar{Q}$ channels: [Dixon,Sofianatos;Ellis,Giele,Zanderighi]
- Basic C++ implementation of spinor products and scalar integrals

amplitude	time
— — —	$60\mu\text{s}$
— — ++	$170\mu\text{s}$
— + — +	$180\mu\text{s}$
+ — —	$200\mu\text{s}$
$\sum_{\text{col.,hel.}} \text{Re}(A^1 A^0)$	10 ms

[CPU: 64 bit 3GHz Intel Core Duo]