

Institute of High Energy Physics Chinese Academy of Sciences

Future Colliders

Lecture-2

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- 2-body phase space 10⁻² suppression;
- Gluon PDF large enhancement at mid-x region;
 - Collinear enhancement, pj~zpg.





Quiz







• The EWSB in the SM (tree-level).

$$\begin{split} V\left(\varphi\right) &= \mu^2 \left(\varphi^{\dagger}\varphi\right) + \lambda \left(\varphi^{\dagger}\varphi\right)^2 \\ \varphi &= \begin{pmatrix} G^+ \\ v + \frac{1}{\sqrt{2}} \left(H^0 + iG^0\right) \end{pmatrix} \quad \left\langle\varphi\right\rangle = \begin{pmatrix} 0 \\ v \end{pmatrix} \end{split}$$



$$\begin{split} D_{\mu}\varphi &= \\ \left(\begin{array}{c} \partial_{\mu}G^{+} + \frac{i}{2} \left(gW_{\mu}^{3} + g'B_{\mu} \right)G^{+} + \frac{ig}{2}W_{\mu}^{+} \left(\sqrt{2}v + H^{0} + iG^{0}\right) \\ \frac{1}{\sqrt{2}}\partial_{\mu}H^{0} + \frac{i}{\sqrt{2}}\partial_{\mu}G^{0} + \frac{ig}{\sqrt{2}}W_{\mu}^{-}G^{+} + \frac{i}{2\sqrt{2}} \left(-gW_{\mu}^{3} + g'B_{\mu} \right) \left(\sqrt{2}v + H^{0} + iG^{0}\right) \\ \end{array} \right) \\ &= \left(\begin{array}{c} \partial_{\mu}G^{+} + ieA_{\mu}G^{+} + \frac{ie(c_{W}^{2} - s_{W}^{2})}{2s_{W}c_{W}}Z_{\mu}G^{+} + \frac{ie}{2s_{W}}W_{\mu}^{+} \left(\sqrt{2}v + H^{0} + iG^{0}\right) \\ \frac{1}{\sqrt{2}}\partial_{\mu}H^{0} + \frac{i}{\sqrt{2}}\partial_{\mu}G^{0} + \frac{ie}{\sqrt{2}s_{W}}W_{\mu}^{-}G^{+} - \frac{ie}{2\sqrt{2}s_{W}c_{W}}Z_{\mu}^{0} \left(\sqrt{2}v + H^{0} + iG^{0}\right) \end{array} \right) \end{split}$$

$$\begin{pmatrix} A^0 \\ Z^0 \end{pmatrix} = \begin{pmatrix} \cos\theta_W & \sin\theta_W \\ -\sin\theta_W & \cos\theta_W \end{pmatrix} \begin{pmatrix} B \\ W^3 \end{pmatrix} \qquad \frac{e^2v^2}{2s_W^2} W^-_\mu W^{+\mu} + \frac{e^2v^2}{4s_W^2 c_W^2} Z^0_\mu Z^{0\mu}$$



The EWSB in the SM (tree-level). • $V(\varphi) = \mu^2 \left(\varphi^{\dagger}\varphi\right) + \lambda \left(\varphi^{\dagger}\varphi\right)^2$ $\varphi = \begin{pmatrix} G^+ \\ v + \frac{1}{\sigma} (H^0 + iG^0) \end{pmatrix} \quad \langle \varphi \rangle = \begin{pmatrix} 0 \\ v \end{pmatrix}$ $D_{\mu}\varphi =$ $\begin{pmatrix} \partial_{\mu}G^{+} + \frac{i}{2} \left(gW_{\mu}^{3} + g'B_{\mu} \right) G^{+} + \frac{ig}{2}W_{\mu}^{+} \left(\sqrt{2}v + H^{0} + iG^{0} \right) \\ \frac{1}{\sqrt{2}}\partial_{\mu}H^{0} + \frac{i}{\sqrt{2}}\partial_{\mu}G^{0} + \frac{ig}{\sqrt{2}}W_{\mu}^{-}G^{+} + \frac{i}{2\sqrt{2}} \left(-gW_{\mu}^{3} + g'B_{\mu} \right) \left(\sqrt{2}v + H^{0} + iG^{0} \right) \end{pmatrix}$ $\partial_{\mu}G^{+} + ieA_{\mu}G^{+} + \frac{ie(c_{W}^{2} - s_{W}^{2})}{2surgen}Z_{\mu}G^{+} + \frac{ie}{2sur}W_{\mu}^{+}(\sqrt{2}v + H^{0} + iG^{0})$ **THINKING**: Try to give the result when the scalar belongs to the (2j+1, y) representation of the $SU(2) \times U(1)$ group. $\begin{pmatrix} A^0 \\ Z^0 \end{pmatrix} = \begin{pmatrix} \cos\theta_W & \sin\theta_W \\ -\sin\theta_W & \cos\theta_W \end{pmatrix} \begin{pmatrix} B \\ W^3 \end{pmatrix} \qquad \frac{e^2v^2}{2s_W^2} W^-_\mu W^{+\mu} + \frac{e^2v^2}{4s_W^2 c_W^2} Z^0_\mu Z^{0\mu}$



• The EWSB in the SM (tree-level).

$$\mathcal{L}_{\text{Yukawa}} = -Y_u^{ij} \epsilon_{ab} \bar{Q}_i^a \varphi^{b*} u_{Rj} - Y_d^{ij} \bar{Q}_i^a \varphi_a d_{Rj} - Y_e^{ij} \bar{L}_i^a \varphi_a e_{Rj} + \text{h.c.}$$
$$Y_u^{ij} v \bar{u}_{Li} u_{Rj} + Y_d^{ij} v \bar{d}_{Li} d_{Rj} + Y_e^{ij} v \bar{e}_{Li} e_{Rj} + \text{h.c.}$$

Polar decomposition theorem:

$$\begin{split} u_L &\to U_L^u u_L, \ u_R \to U_R^u u_R, \ d_L \to U_L^d d_L, \ d_R \to U_R^d d_R \\ \mathcal{L}_{\text{Yukawa}} &= -\frac{m_{ui}}{v} \epsilon_{ab} \bar{Q}_i^a \varphi^{b*} u_{Ri} - \frac{m_{di}}{v} \bar{Q}_i^a \varphi_a d_{Ri} - \frac{m_{ei}}{v} \bar{L}_i^a \varphi_a e_{Ri} + \text{h.c.} \\ V_{CKM} &= U_L^{u\dagger} U_L^d \end{split}$$



• The vertices with Higgs boson.

 $\mathcal{L}_{\text{Yukawa}} = -\frac{m_{ui}}{v} \epsilon_{ab} \bar{Q}_i^a \varphi^{b*} u_{Ri} - \frac{m_{di}}{v} \bar{Q}_i^a \varphi_a d_{Ri} - \frac{m_{ei}}{v} \bar{L}_i^a \varphi_a e_{Ri} + \text{h.c.}$ $\underbrace{\mathbf{H}^0}_{\mathbf{f}} - i \frac{m_f}{\sqrt{2}v} \propto m_f$

	u-quark	d-quark	s-quark	c-quark	b-quark	t-quark	е	μ	τ
mass (GeV)	0.0022	0.0047	0.096	1.28	4.18	160	0.000511	0.105	1.776
Yukawa	0.000009	0.00002	0.0004	0.005	0.017	0.65	0.000002	0.0004	0.0072



• The vertices with Higgs boson.

$$\begin{pmatrix} \partial_{\mu}G^{+} + ieA_{\mu}G^{+} + \frac{ie(c_{W}^{2} - s_{W}^{2})}{2s_{W}c_{W}}Z_{\mu}G^{+} + \frac{ie}{2s_{W}}W_{\mu}^{+}\left(\sqrt{2}v + H^{0} + iG^{0}\right) \\ \frac{1}{\sqrt{2}}\partial_{\mu}H^{0} + \frac{i}{\sqrt{2}}\partial_{\mu}G^{0} + \frac{ie}{\sqrt{2}s_{W}}W_{\mu}^{-}G^{+} - \frac{ie}{2\sqrt{2}s_{W}c_{W}}Z_{\mu}^{0}\left(\sqrt{2}v + H^{0} + iG^{0}\right) \end{pmatrix}$$





• The vertices with Higgs boson.



$$-6\sqrt{2}i\lambda v = -\frac{3iem_h^2}{2m_W s_W} -6i\lambda v = -6i\lambda v = -6i\lambda v$$
$$\sim -1.1vi\lambda v = -6i\lambda v = -6i\lambda v$$

$$-6i\lambda = -\frac{3ie^2m_h^2}{4m_W^2s_W^2}$$
$$\sim -0.77i$$



• Low energy theorem:

$$\lim_{p_h \to 0} \mathcal{M}(A \to B + h) = \frac{1}{\sqrt{2}v} \left(\sum_f m_f \frac{\partial}{\partial m_f} + \sum_{V=W,Z} m_V \frac{\partial}{\partial m_V} \right) \mathcal{M}(A \to B)$$

• The vertices with Higgs boson (loop induced).

$$\lim_{p_h \to 0} \mathcal{M}(g \to g + h) \approx \frac{m_t}{\sqrt{2}v} \frac{\partial}{\partial m_t} \mathcal{M}(g \to g)$$

$$= \frac{2\alpha_S}{\pi} \delta^{ab} T_F(p^2 g^{\mu\nu} - p^{\mu} p^{\nu}) \frac{m_t}{\sqrt{2}v} \frac{1}{3m_t}$$

$$= \frac{\alpha_S}{3\sqrt{2}\pi v} \delta^{ab} (p^2 g^{\mu\nu} - p^{\mu} p^{\nu})$$

$$\Rightarrow \mathcal{L}_{\text{eff}} = \frac{\alpha_S}{12\sqrt{2}\pi v} h G^a_{\mu\nu} G^{a,\mu\nu}$$



• Low energy theorem:

$$\lim_{p_h \to 0} \mathcal{M}(A \to B + h) = \frac{1}{\sqrt{2}v} \left(\sum_f m_f \frac{\partial}{\partial m_f} + \sum_{V=W,Z} m_V \frac{\partial}{\partial m_V} \right) \mathcal{M}(A \to B)$$

• The vertices with Higgs boson (loop induced).











- Thresholds:
 - WW: 2×80=160GeV;
 - ZZ: 2×91=181GeV;
 - Zh: 91+125=216GeV;
 - tt: 2×173=346GeV.





s-channel processes: s⁻¹ suppression in the large s limit from the denominator in the propagator.



• The helicity configuration (in SM, J=1):





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$$\begin{array}{c} \mathbf{Z} \\ \mathbf{P} \\ \mathbf{$$

• The longitudinal mode of Z is the Goldstone boson.

$$\begin{pmatrix} \partial_{\mu}G^{+} + ieA_{\mu}G^{+} + \frac{ie(c_{W}^{2} - s_{W}^{2})}{2s_{W}c_{W}}Z_{\mu}G^{+} + \frac{ie}{2s_{W}}W_{\mu}^{+}\left(\sqrt{2}v + H^{0} + iG^{0}\right) \\ \frac{1}{\sqrt{2}}\partial_{\mu}H^{0} + \underbrace{\frac{i}{\sqrt{2}}\partial_{\mu}G^{0}}_{\sqrt{2}s_{W}} + \frac{ie}{\sqrt{2}s_{W}}W_{\mu}^{-}G^{+} - \underbrace{\frac{ie}{2\sqrt{2}s_{W}c_{W}}}_{2\sqrt{2}s_{W}c_{W}}Z_{\mu}^{0}\left(\sqrt{2}v + H^{0} + iG^{0}\right) \end{pmatrix}$$



- The helicity configuration (in SM, J=1).
- The (massless) electron and positron can only couple to the transverse Z*.
- The Higgs boson couples to both transverse and longitude Z.

$$\langle Z(\lambda)H|Z^*(m)\rangle = \frac{gm_Z}{c_W}d^1_{m,\lambda}(\theta)\Gamma_{\lambda 0}$$

• The longitudinal mode of Z is the Goldstone boson.



 $\Gamma_{00} = -E_Z/m_Z, \ \Gamma_{10} = -1$



- The helicity configuration (in SM, J=1).
- The (massless) electron and positron can only couple to the transverse Z*.
- The Higgs boson couples to both transverse and longitude Z.

$$\begin{array}{c} & \mathbf{Z} \\ \hline & \mathbf{C} \\ \hline & \mathbf{C} \\ H \end{array} \begin{pmatrix} \mathbf{Z} \\ \mathbf{C} \\ \mathbf{W} \end{pmatrix} = \frac{g m_Z}{c_W} d^1_{m,\lambda}(\theta) \Gamma_{\lambda 0} \\ \end{array}$$

• The longitudinal mode of Z is the Goldstone boson.

$$\frac{d\sigma}{d\cos\theta} = \frac{G_F^2 m_Z^6 (v_e^2 + a_e^2)\beta}{32\pi [(s - m_Z^2)^2 + m_Z^2 \Gamma_Z^2]} \left[|\Gamma_{00}^2| \sin^2\theta + |\Gamma_{10}|^2 (1 + \cos^2\theta) \right]
= \frac{G_F^2 m_Z^6 (v_e^2 + a_e^2)}{16\pi [(s - m_Z^2)^2 + m_Z^2 \Gamma_Z^2]} \left(1 + \frac{p_Z^2}{2m_Z^2} \sin^2\theta \right) \left[1 - \frac{(m_H + m_Z)^2}{s} \right]^{1/2} \left[1 - \frac{(m_H - m_Z)^2}{s} \right]^{1/2}$$



- The Electron and positron can be polarized.
- The initial state prepared by the collider should be described by a density matrix.
- The initial state electron and initial state positron is fully disentangled, so the density matrix can be decomposed as a tensor product of the electron density matrix and the positron density matrix.

$$\operatorname{Tr} \rho = 1, \quad \operatorname{Tr} \rho^2 \leq 1,$$

$$\rho = \rho_{++} |+\rangle \langle +|+\rho_{+-}|+\rangle \langle -|+\rho_{-+}|-\rangle \langle +|+\rho_{--}|-\rangle \langle -|$$

The transition probability is

$$|\mathcal{M}|^{2} = \operatorname{Tr}(\rho_{e^{+}} \otimes \rho_{e^{-}} \mathcal{M}^{\dagger} \mathcal{M})$$
$$= \sum \rho_{e^{+}}(\lambda_{e^{+}} \lambda_{e^{+}}') \rho_{e^{-}}(\lambda_{e^{-}} \lambda_{e^{-}}') \mathcal{M}^{\dagger}(\lambda_{e^{+}}' \lambda_{e^{-}}') \mathcal{M}(\lambda_{e^{+}} \lambda_{e^{-}})$$



• Unpolarized beam, pure state, and general case:

$$\rho = \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}, \quad \rho = |\Psi\rangle \langle \Psi|, \quad \rho = \frac{1}{2} \begin{pmatrix} 1 + P_L & P_1 - iP_2 \\ P_1 + iP_2 & 1 - P_L \end{pmatrix}$$

$$\begin{aligned} |\mathcal{M}|^2 &= \frac{1}{4} \left\{ (1 - P_{e^-})(1 + P_{e^+}) |F_{LR}|^2 + (1 + P_{e^-})(1 - P_{e^+}) |F_{RL}|^2 \\ &+ (1 - P_{e^-})(1 - P_{e^+}) |F_{LL}|^2 + (1 + P_{e^-})(1 + P_{e^+}) |F_{RR}|^2 \\ &- 2P_{e^-}^T P_{e^+}^T \{ [\cos(\phi_- - \phi_+) \operatorname{Re} (F_{RR} F_{LL}^*) + \cos(\phi_- + \phi_+ - 2\phi) \operatorname{Re} (F_{LR} F_{RL}^*)] + \dots \end{aligned} \end{aligned}$$

 If the electron and positron are longitudinal polarized, because the left-handed current and the right-handed current interaction are not the same, the total cross section will be changed.

$$\begin{split} \sigma_{P_e - P_{e^+}} &= \frac{1}{4} \left\{ (1 + P_{e^-})(1 + P_{e^+})\sigma_{\mathrm{RR}} + (1 - P_{e^-})(1 - P_{e^+})\sigma_{\mathrm{LL}} \right. \\ &+ (1 + P_{e^-})(1 - P_{e^+})\sigma_{\mathrm{RL}} + (1 - P_{e^-})(1 + P_{e^+})\sigma_{\mathrm{LR}} \right\} \end{split}$$



• Electron positron collider with c.m. E ~240-250GeV: Higgs factory.





 The theoretical prediction of the total cross section has significant uncertainty.





The theoretical prediction of the total cross section has significant uncertainty.

$\sqrt{\epsilon}$ (CeV)		LO (fb)	NLO Weak (fb)		NNLO mixed electroweak-QCD (fb)			
Vs (Gev)		$\sigma^{(0)}$	$\sigma^{(\alpha)}$	$\sigma^{(0)} + \sigma^{(\alpha)}$	$\sigma_Z^{(\alpha\alpha_a)}$	$\sigma_{\gamma}^{(lpha lpha_{s})}$	$\sigma^{(\alpha \alpha_s)}$	$\sigma^{(0)} + \sigma^{(\alpha)} + \sigma^{(\alpha\alpha_s)}$
240	Total	223.14	6.64	229.78	2.42	0.008	2.43	232.21
	L	88.67	3.18	91.86	0.96	0.003	0.97	92.82
	Т	134.46	3.46	137.92	1.46	0.005	1.46	139.39
250	Total	223.12	6.08	229.20	2.42	0.009	2.42	231.63
	L	94.30	3.31	97.61	1.02	0.004	1.02	98.64
	Т	128.82	2.77	131.59	1.40	0.005	1.40	132.99

\sqrt{s}	schemes	$\sigma_{ m LO}~({ m fb})$	$\sigma_{ m NLO}~({ m fb})$	$\sigma_{\rm NNLO}$ (fb)
	$\alpha(0)$	223.14 ± 0.47	229.78 ± 0.77	$232.21_{-0.75-0.21}^{+0.75+0.10}$
240	$\alpha(M_Z)$	252.03 ± 0.60	$228.36\substack{+0.82\\-0.81}$	$231.28^{+0.80+0.12}_{-0.79-0.25}$
	G_{μ}	239.64 ± 0.06	$232.46^{+0.07}_{-0.07}$	$233.29^{+0.07+0.03}_{-0.06-0.07}$
	$\alpha(0)$	223.12 ± 0.47	229.20 ± 0.77	$231.63^{+0.75+0.12}_{-0.75-0.21}$
250	$\alpha(M_Z)$	252.01 ± 0.60	$227.67^{+0.82}_{-0.81}$	$230.58^{+0.80+0.14}_{-0.79-0.25}$
	G_{μ}	239.62 ± 0.06	231.82 ± 0.07	$232.65^{+0.07+0.04}_{-0.07-0.07}$





	Z	W	t	Н
Width (GeV)	2.4952	2.085	1.41	0.00407
Μ/ Γ	36.545	38.55	123	3.07×10 ⁴













Channel	Br (%)	# of Events	Var	Rel err (%)
bb	58.24	582400	763	0.1
CC	2.891	28910	170	0.6
au au	6.272	62720	250	0.4
μμ	0.02176	218	15	6.9
WW	21.37	213700	462	0.2
ZZ	2.61	26100	162	0.6
gg	8.187	81870	286	0.3
γγ	0.227	2270	48	2.1
Ζγ	0.153	1530	39	2.5



Next Lecture

- Higgs phenomenology at future lepton colliders.
- A short introduction of the top-quark mass and its measurement at future lepton colliders.

