

WW Production With a Jet Veto Made Simple(r)

Luke Arpino

December 6, 2017

University of Sussex

Jet Vetoes: What and Why?

Jet Vetoes: What?

Jet Vetoes: What?

Jet Veto

A cut on the maximum p_t of a jet: $p_t^{(J)} \leq p_{t,veto}$

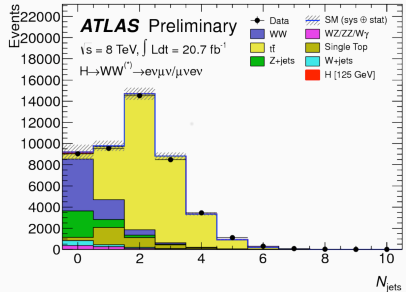
Jet Vetoes: Why?

Jet Vetoes: Why?

- $H \rightarrow WW \rightarrow l^+ \nu l^- \bar{\nu}$
important channel for
measurement of Higgs spin
and coupling

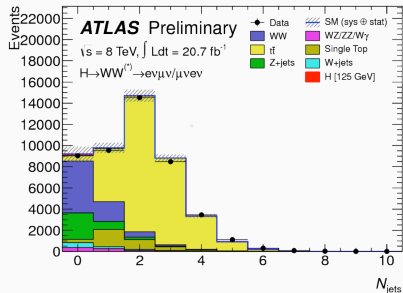
Jet Vetoes: Why?

- $H \rightarrow WW \rightarrow l^+ \nu l^- \bar{\nu}$
important channel for
measurement of Higgs spin
and coupling



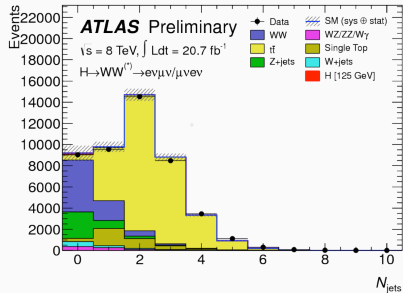
Jet Vetoes: Why?

- $H \rightarrow WW \rightarrow l^+ \nu l^- \bar{\nu}$
important channel for
measurement of Higgs spin
and coupling
- Massive $t\bar{t}$ background from
 $t \rightarrow bW$



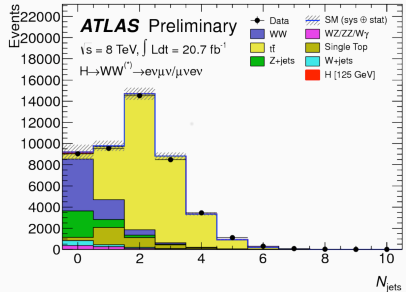
Jet Vetoes: Why?

- $H \rightarrow WW \rightarrow l^+ \nu l^- \bar{\nu}$
important channel for
measurement of Higgs spin
and coupling
- Massive $t\bar{t}$ background from
 $t \rightarrow bW$
- 0-jet bin least polluted by
tops



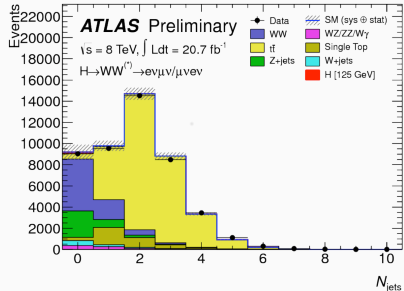
Jet Vetoes: Why?

- $H \rightarrow WW \rightarrow l^+ \nu l^- \bar{\nu}$
important channel for measurement of Higgs spin and coupling
- Massive $t\bar{t}$ background from $t \rightarrow bW$
- 0-jet bin least polluted by tops
- Veto on jets very efficient at reducing background



Jet Vetoes: Why?

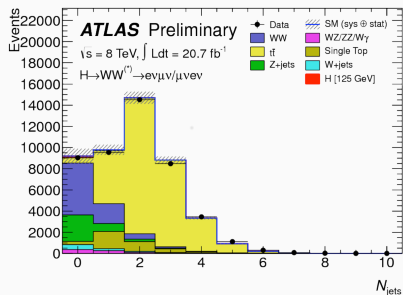
- $H \rightarrow WW \rightarrow l^+ \nu l^- \bar{\nu}$
important channel for measurement of Higgs spin and coupling
- Massive $t\bar{t}$ background from $t \rightarrow bW$
- 0-jet bin least polluted by tops
- Veto on jets very efficient at reducing background



Higgs + 0-jet cross section becomes the quantity of interest

Jet Vetoes: Why?

- $H \rightarrow WW \rightarrow l^+ \nu l^- \bar{\nu}$
important channel for measurement of Higgs spin and coupling
- Massive $t\bar{t}$ background from $t \rightarrow bW$
- 0-jet bin least polluted by tops
- Veto on jets very efficient at reducing background



Higgs + 0-jet cross section becomes the quantity of interest

Need to quantify how the jet veto reduces the Higgs cross section and the non $t\bar{t}$ cross section

This Talk

- Lots of work already done on the Higgs production cross section: $NNLL + N^3LO + LL_R$ known

- Lots of work already done on the Higgs production cross section: $NNLL + N^3LO + LL_R$ known
- WW is an interesting laboratory for new physics (even without the Higgs)
 - Top partners
 - Contact interactions

1. Jet veto resummation in a nutshell
2. Automation of jet veto resummation
3. A case study: $pp \rightarrow WW$

Jet Veto Resummation in a Nutshell

In the presence of tight jet vetoes resummation is required

Resummation

In the presence of tight jet vetoes resummation is required

Look at jets in the soft $\omega \ll M$

Resummation

In the presence of tight jet vetoes resummation is required

Look at jets in the soft $\omega \ll M$ and collinear $\theta \ll 1$ limit:

Resummation

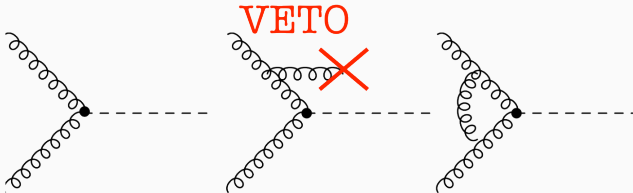
In the presence of tight jet vetoes resummation is required

Look at jets in the soft $\omega \ll M$ and collinear $\theta \ll 1$ limit:

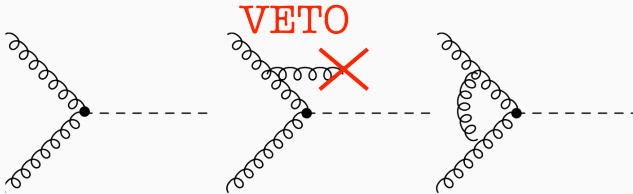
Resummation

In the presence of tight jet vetoes resummation is required

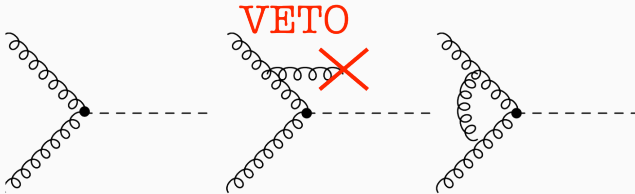
Look at jets in the soft $\omega \ll M$ and collinear $\theta \ll 1$ limit:



Resummation



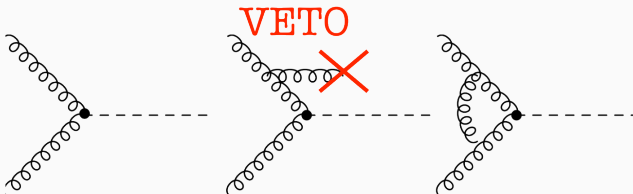
Resummation



$$\begin{aligned}\sigma_{0\text{-jet}}(p_{t,\text{veto}}) &\simeq \sigma_0 \left(1 + C \frac{\alpha_s}{\pi} \int \frac{d\omega}{\omega} \frac{d\theta^2}{\theta^2} (\Theta(p_{t,\text{veto}} - \omega\theta) - 1) \right) \\ &\simeq \sigma_0 \left(1 - 2C \frac{\alpha_s}{\pi} \log^2 \frac{M}{p_{t,\text{veto}}} + \dots \right)\end{aligned}$$

$C = C_A$ for gluons, $C = C_F$ for quarks

Resummation



$$\begin{aligned}\sigma_{0\text{-jet}}(p_{t,\text{veto}}) &\simeq \sigma_0 \left(1 + C \frac{\alpha_s}{\pi} \int \frac{d\omega}{\omega} \frac{d\theta^2}{\theta^2} (\Theta(p_{t,\text{veto}} - \omega\theta) - 1) \right) \\ &\simeq \sigma_0 \left(1 - 2C \frac{\alpha_s}{\pi} \log^2 \frac{M}{p_{t,\text{veto}}} + \dots \right)\end{aligned}$$

$C = C_A$ for gluons, $C = C_F$ for quarks

Remnant logs left after cancellation of real and virtual divergences

Resummation of Large Logs

Resummation of Large Logs

- NLL resummed calculations for jet observables automated in numerical code CAESAR, provided the observable satisfies some applicability conditions (i.e. rIRC safe, continuously global)

Resummation of Large Logs

- NLL resummed calculations for jet observables automated in numerical code CAESAR, provided the observable satisfies some applicability conditions (i.e. rIRC safe, continuously global)
- First observation is that the jet veto is within the scope of CAESAR, in fact at NLL the cross-section with a veto is just a Sudakov form factor

Resummation of Large Logs

- NLL resummed calculations for jet observables automated in numerical code CAESAR, provided the observable satisfies some applicability conditions (i.e. rIRC safe, continuously global)
- First observation is that the jet veto is within the scope of CAESAR, in fact at NLL the cross-section with a veto is just a Sudakov form factor
- NNLL resummation calculated by extending CAESAR prescription and using the known NNLL resummation of the Higgs p_t spectrum. This calculation is implemented in the numerical code JetVHeto

Resummation of Large Logs

- NLL resummed calculations for jet observables automated in numerical code CAESAR, provided the observable satisfies some applicability conditions (i.e. rIRC safe, continuously global)
- First observation is that the jet veto is within the scope of CAESAR, in fact at NLL the cross-section with a veto is just a Sudakov form factor
- NNLL resummation calculated by extending CAESAR prescription and using the known NNLL resummation of the Higgs p_t spectrum. This calculation is implemented in the numerical code JetVHeto
- Other methods based on SCET

The NNLL master formulae

The NNLL master formulae

$$\Sigma^{(J)}(p_{t,\text{veto}}) \sim \mathcal{L}(\mu_F = p_{t,\text{veto}}) \otimes \mathcal{M}_B^2 \otimes e^{-R(p_{t,\text{veto}})} (1 + \delta\mathcal{F}(p_{t,\text{veto}}))$$

The NNLL master formulae

$$\Sigma^{(J)}(p_{t,\text{veto}}) \sim \mathcal{L}(\mu_F = p_{t,\text{veto}}) \otimes \mathcal{M}_B^2 \otimes e^{-R(p_{t,\text{veto}})} (1 + \delta\mathcal{F}(p_{t,\text{veto}}))$$

- The usual QCD parton luminosities

The NNLL master formulae

$$\Sigma^{(J)}(p_{t,\text{veto}}) \sim \mathcal{L}(\mu_F = p_{t,\text{veto}}) \otimes \mathcal{M}_B^2 \otimes e^{-R(p_{t,\text{veto}})} (1 + \delta\mathcal{F}(p_{t,\text{veto}}))$$

- The usual QCD parton luminosities
- Sudakov form factor

NNLL Matching

This is not quite the full story, need to match the resummation against the fixed order

NLL Matching

This is not quite the full story, need to match the resummation against the fixed order

Schematically:

$$\Sigma_{matched}(p_{t,veto}) = \Sigma_{res}(p_{t,veto}) + \Sigma_{f.o.}(p_{t,veto}) - \Sigma_{res}(p_{t,veto}) \Big|_{\text{expanded in } \alpha_s}$$

NNLL Matching

This is not quite the full story, need to match the resummation against the fixed order

Schematically:

$$\Sigma_{matched}(p_{t,veto}) = \Sigma_{res}(p_{t,veto}) + \Sigma_{f.o.}(p_{t,veto}) - \Sigma_{res}(p_{t,veto}) \Big|_{\text{expanded in } \alpha_s}$$

- Σ_{res} valid for $p_{t,veto} \ll M$

NNLL Matching

This is not quite the full story, need to match the resummation against the fixed order

Schematically:

$$\Sigma_{\text{matched}}(p_{t,\text{veto}}) = \Sigma_{\text{res}}(p_{t,\text{veto}}) + \Sigma_{\text{f.o.}}(p_{t,\text{veto}}) - \Sigma_{\text{res}}(p_{t,\text{veto}}) \Big|_{\text{expanded in } \alpha_s}$$

- Σ_{res} valid for $p_{t,\text{veto}} \ll M$
- $\Sigma_{\text{f.o.}}$ valid for $p_{t,\text{veto}} \sim M$

NNLL Matching

This is not quite the full story, need to match the resummation against the fixed order

Schematically:

$$\Sigma_{matched}(p_{t,veto}) = \Sigma_{res}(p_{t,veto}) + \Sigma_{f.o.}(p_{t,veto}) - \Sigma_{res}(p_{t,veto}) \Big|_{\text{expanded in } \alpha_s}$$

- Σ_{res} valid for $p_{t,veto} \ll M$
- $\Sigma_{f.o.}$ valid for $p_{t,veto} \sim M$
- subtract $\Sigma_{res} \Big|_{\text{expanded in } \alpha_s}$ to remove double counting.

NNLL Matching

This is not quite the full story, need to match the resummation against the fixed order

Schematically:

$$\Sigma_{matched}(p_{t,veto}) = \Sigma_{res}(p_{t,veto}) + \Sigma_{f.o.}(p_{t,veto}) - \Sigma_{res}(p_{t,veto}) \Big|_{\text{expanded in } \alpha_s}$$

- Σ_{res} valid for $p_{t,veto} \ll M$
- $\Sigma_{f.o.}$ valid for $p_{t,veto} \sim M$
- subtract $\Sigma_{res} \Big|_{\text{expanded in } \alpha_s}$ to remove double counting.

Automated Resummation of Jet Veto

Current Status of Automation in Resummation

Current Status of Automation in Resummation

The Good

- Lots of resummation codes for many observables
- Cross checks between groups \implies correct resummation + code!

Current Status of Automation in Resummation

The Good

- Lots of resummation codes for many observables
- Cross checks between groups \implies correct resummation + code!

The Bad

- Most code is process specific
- Lots of reinventing the wheel
- Not exactly user friendly

Current Status of Automation in Resummation

The Good

- Lots of resummation codes for many observables
- Cross checks between groups \implies correct resummation + code!

The Bad

- Most code is process specific
- Lots of reinventing the wheel
- Not exactly user friendly

The Ugly

- Specialised resummation codes can become very large
- Implementation error prone (hardly unique to resummation)

Current Status of Automation in Resummation

The Good

- Lots of resummation codes for many observables
- Cross checks between groups \implies correct resummation + code!

The Bad

- Most code is process specific
- Lots of reinventing the wheel
- Not exactly user friendly

The Ugly

- Specialised resummation codes can become very large
- Implementation error prone (hardly unique to resummation)

There must be a better way!

Utilise existing Monte Carlo code

MCFM + Resummation

Monte Carlo for FeMtobarn processes

- (N)NLO Monte Carlo integrator

- (N)NLO Monte Carlo integrator
- Human readable implementation of matrix elements

- (N)NLO Monte Carlo integrator
- Human readable implementation of matrix elements
- NLO is implemented through a local subtraction scheme:
Catani-Seymour subtraction dipoles

- (N)NLO Monte Carlo integrator
- Human readable implementation of matrix elements
- NLO is implemented through a local subtraction scheme:
Catani-Seymour subtraction dipoles
- Local subtraction terms helpful

Calculation of cross sections at NLO

QCD NLO Cross Section

$$\delta\sigma^{NLO} = \int_m d\sigma^V + \int_{m+1} d\sigma^R$$

Calculation of cross sections at NLO

QCD NLO Cross Section

$$\delta\sigma^{NLO} = \int_m d\sigma^V + \int_{m+1} d\sigma^R$$

Problems

QCD NLO Cross Section

$$\delta\sigma^{NLO} = \int_m d\sigma^V + \int_{m+1} d\sigma^R$$

Problems

- virtual and real contributions are separately divergent

QCD NLO Cross Section

$$\delta\sigma^{NLO} = \int_m d\sigma^V + \int_{m+1} d\sigma^R$$

Problems

- virtual and real contributions are separately divergent
- infinities only cancel when combined

Calculation of cross sections at NLO

QCD NLO Cross Section

$$\delta\sigma^{NLO} = \int_m d\sigma^V + \int_{m+1} d\sigma^R$$

Problems

- virtual and real contributions are separately divergent
- infinities only cancel when combined
- integrated over different phase spaces before combination

Calculation of cross sections at NLO

QCD NLO Cross Section

$$\delta\sigma^{NLO} = \int_m d\sigma^V + \int_{m+1} d\sigma^R$$

Problems

- virtual and real contributions are separately divergent
- infinities only cancel when combined
- integrated over different phase spaces before combination

Completely unsuitable for numerical integration!

The Subtraction Method

The Subtraction Method

Subtraction Term

Cleverly add zero by introducing a subtraction term $d\sigma^A$:

$$\delta\sigma^{NLO} = \int_m d\sigma^V + \int_{m+1} [d\sigma^R - d\sigma^A] + \int_{m+1} d\sigma^A$$

The Subtraction Method

Subtraction Term

Cleverly add zero by introducing a subtraction term $d\sigma^A$:

$$\delta\sigma^{NLO} = \int_m d\sigma^V + \int_{m+1} [d\sigma^R - d\sigma^A] + \int_{m+1} d\sigma^A$$

- $d\sigma^A$ exactly matches the singular behaviour of $d\sigma^R$

The Subtraction Method

Subtraction Term

Cleverly add zero by introducing a subtraction term $d\sigma^A$:

$$\delta\sigma^{NLO} = \int_m d\sigma^V + \int_{m+1} [d\sigma^R - d\sigma^A] + \int_{m+1} d\sigma^A$$

- $d\sigma^A$ exactly matches the singular behaviour of $d\sigma^R$
- $d\sigma^A$ must be analytically integrable over the one-parton phase space leading to the divergences

The Subtraction Method

Subtraction Term

Cleverly add zero by introducing a subtraction term $d\sigma^A$:

$$\delta\sigma^{NLO} = \int_m \left[d\sigma^V + \int_1 d\sigma^A \right]_{\epsilon=0} + \int_{m+1} \left[(d\sigma^R)_{\epsilon=0} - (d\sigma^A)_{\epsilon=0} \right]$$

- $d\sigma^A$ exactly matches the singular behaviour of $d\sigma^R$
- $d\sigma^A$ must be analytically integrable over the one-parton phase space leading to the divergences

A Glimpse at the Method

A Glimpse at the Method

$$\Sigma^{(J)}(p_{t,\text{veto}}) \sim \mathcal{L}(p_{t,\text{veto}}) \otimes \mathcal{M}_B^2 \otimes e^{-R(p_{t,\text{veto}})} (1 + \delta\mathcal{F}(p_{t,\text{veto}}))$$

A Glimpse at the Method

$$\Sigma^{(J)}(p_{t,\text{veto}}) \sim \mathcal{L}(p_{t,\text{veto}}) \otimes \mathcal{M}_B^2 \otimes e^{-R(p_{t,\text{veto}})} (1 + \delta\mathcal{F}(p_{t,\text{veto}}))$$

- In a generic resummation the real radiation exponentiates and enters through the Sudakov form factor

A Glimpse at the Method

$$\Sigma^{(J)}(p_{t,\text{veto}}) \sim \mathcal{L}(p_{t,\text{veto}}) \otimes \mathcal{M}_B^2 \otimes e^{-R(p_{t,\text{veto}})} (1 + \delta\mathcal{F}(p_{t,\text{veto}}))$$

- In a generic resummation the real radiation exponentiates and enters through the Sudakov form factor
- Luminosity $\mathcal{L}(p_{t,\text{veto}}) \otimes \mathcal{M}_B^2$ lives in the Born phase space

A Glimpse at the Method

$$\Sigma^{(J)}(p_{t,veto}) \sim \mathcal{L}(p_{t,veto}) \otimes \mathcal{M}_B^2 \otimes e^{-R(p_{t,veto})} (1 + \delta\mathcal{F}(p_{t,veto}))$$

- In a generic resummation the real radiation exponentiates and enters through the Sudakov form factor
- Luminosity $\mathcal{L}(p_{t,veto}) \otimes \mathcal{M}_B^2$ lives in the Born phase space
- Modify the integrated subtraction terms to get the correct luminosity for the resummation

A Glimpse at the Method

$$\Sigma^{(J)}(p_{t,veto}) \sim \mathcal{L}(p_{t,veto}) \otimes \mathcal{M}_B^2 \otimes e^{-R(p_{t,veto})} (1 + \delta\mathcal{F}(p_{t,veto}))$$

- In a generic resummation the real radiation exponentiates and enters through the Sudakov form factor
- Luminosity $\mathcal{L}(p_{t,veto}) \otimes \mathcal{M}_B^2$ lives in the Born phase space
- Modify the integrated subtraction terms to get the correct luminosity for the resummation
- Sudakov form factor calculated and provided through an interface between JetVHeto and MCFM

A Glimpse at the Method

A Glimpse at the Method

$$\int_1 d\sigma^A \propto \left(\frac{\alpha_s}{2\pi}\right) c_\Gamma \left[\mathcal{V}^{end} \delta(1-z) + \mathcal{V}^{plus} + \mathcal{V}^{regular} \right]$$

A Glimpse at the Method

$$\int_1 d\sigma^A \propto \left(\frac{\alpha_s}{2\pi}\right) c_\Gamma \left[\mathcal{V}^{end} \delta(1-z) + \mathcal{V}^{plus} + \mathcal{V}^{regular} \right]$$

- Modify \mathcal{V} terms to include the resummation through the subtraction dipoles

A Glimpse at the Method

$$\int_1 d\sigma^A \propto \left(\frac{\alpha_s}{2\pi}\right) c_\Gamma \left[\mathcal{V}^{end} \delta(1-z) + \mathcal{V}^{plus} + \mathcal{V}^{regular} \right]$$

- Modify \mathcal{V} terms to include the resummation through the subtraction dipoles
- \mathcal{V}^{end} encodes the virtual corrections, not changed by the resummation

A Glimpse at the Method

$$\int_1 d\sigma^A \propto \left(\frac{\alpha_s}{2\pi}\right) c_\Gamma \left[\mathcal{V}^{end} \delta(1-z) + \mathcal{V}^{plus} + \mathcal{V}^{regular} \right]$$

- Modify \mathcal{V} terms to include the resummation through the subtraction dipoles
- \mathcal{V}^{end} encodes the virtual corrections, not changed by the resummation
- Calculate \mathcal{V}^{plus} and $\mathcal{V}^{regular}$ once and for all

Comments on the Implementation

Comments on the Implementation

- Can shuffle terms around: subtraction terms \Leftrightarrow Sudakov

Comments on the Implementation

- Can shuffle terms around: subtraction terms \Leftrightarrow Sudakov
- Some of the usual simplifications not applicable: real and virtual emissions have different characteristic scales
 - *virtual* $\sim M$
 - *real* $\sim p_{t,veto}$

pick up additional π^2 terms not in MCFM

Comments on the Implementation

- Can shuffle terms around: subtraction terms \Leftrightarrow Sudakov
- Some of the usual simplifications not applicable: real and virtual emissions have different characteristic scales

- *virtual* $\sim M$
- *real* $\sim p_{t,\text{veto}}$

pick up additional π^2 terms not in MCFM

- Regularisation schemes
 - MCFM uses FDH scheme for most calculations
 - Need to match back to \overline{MS}

- Can resum jet vetoes for any colour singlet at NNLL

- Can resum jet vetoes for any colour singlet at NNLL
- Keep everything included in MCFM: spin correlations, experimental cuts, interference effects, exclusive decay products, . . .

- Can resum jet vetoes for any colour singlet at NNLL
- Keep everything included in MCFM: spin correlations, experimental cuts, interference effects, exclusive decay products, . . .
- Better validation, already know most of the setup is correct

Automated Resummation

- Can resum jet vetoes for any colour singlet at NNLL
- Keep everything included in MCFM: spin correlations, experimental cuts, interference effects, exclusive decay products, . . .
- Better validation, already know most of the setup is correct
- Less work, more time to spend working on new resummations

Automated Resummation

- Can resum jet vetoes for any colour singlet at NNLL
- Keep everything included in MCFM: spin correlations, experimental cuts, interference effects, exclusive decay products, . . .
- Better validation, already know most of the setup is correct
- Less work, more time to spend working on new resummations
- Can be brought to other Monte Carlos:
POWHEG, MG5_aMC@NLO, . . .

Automated Resummation

- Can resum jet vetoes for any colour singlet at NNLL
- Keep everything included in MCFM: spin correlations, experimental cuts, interference effects, exclusive decay products, ...
- Better validation, already know most of the setup is correct
- Less work, more time to spend working on new resummations
- Can be brought to other Monte Carlos:
POWHEG, MG5_aMC@NLO, ...

Completely general!

A case study: $pp \rightarrow WW$

The Setup

The Setup

Consider the SMEFT inspired Lagrangian:

$$\mathcal{L} \supseteq -\kappa_t \frac{m_t}{v} H (\bar{t}_R t_L + h.c.) + \kappa_g \frac{\alpha_s}{12\pi} \frac{H}{v} G_{\mu\nu}^a G^{\mu\nu,a}$$

The Setup

Consider the SMEFT inspired Lagrangian:

$$\mathcal{L} \supseteq -\kappa_t \frac{m_t}{v} H (\bar{t}_R t_L + h.c.) + \kappa_g \frac{\alpha_s}{12\pi} \frac{H}{v} G_{\mu\nu}^a G^{\mu\nu,a}$$

For this to be phenomenologically relevant:

The Setup

Consider the SMEFT inspired Lagrangian:

$$\mathcal{L} \supseteq -\kappa_t \frac{m_t}{v} H (\bar{t}_R t_L + h.c.) + \kappa_g \frac{\alpha_s}{12\pi} \frac{H}{v} G_{\mu\nu}^a G^{\mu\nu,a}$$

For this to be phenomenologically relevant:

- Higgs total cross section must be kept constant: strongly constrained by experimental measurements

The Setup

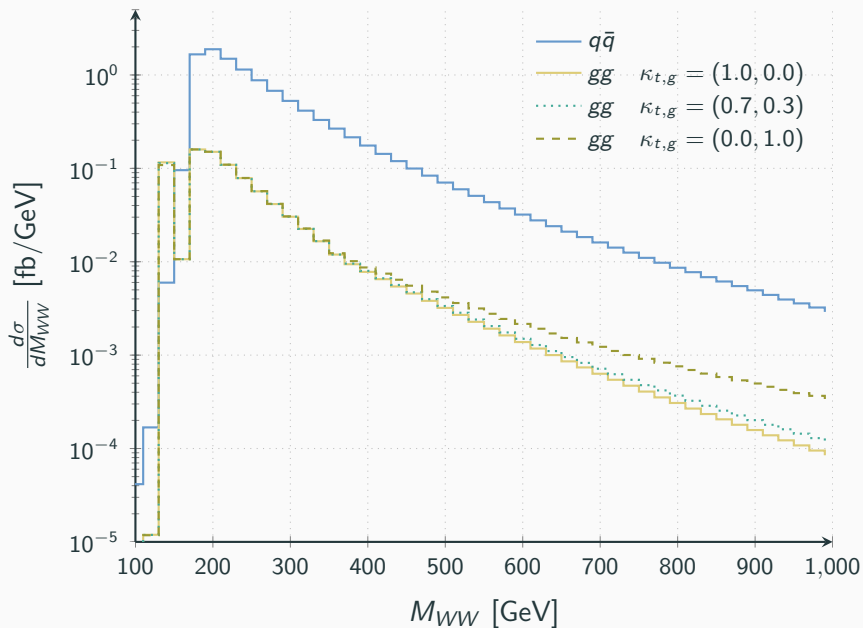
Consider the SMEFT inspired Lagrangian:

$$\mathcal{L} \supseteq -\kappa_t \frac{m_t}{v} H (\bar{t}_R t_L + h.c.) + \kappa_g \frac{\alpha_s}{12\pi} \frac{H}{v} G_{\mu\nu}^a G^{\mu\nu,a}$$

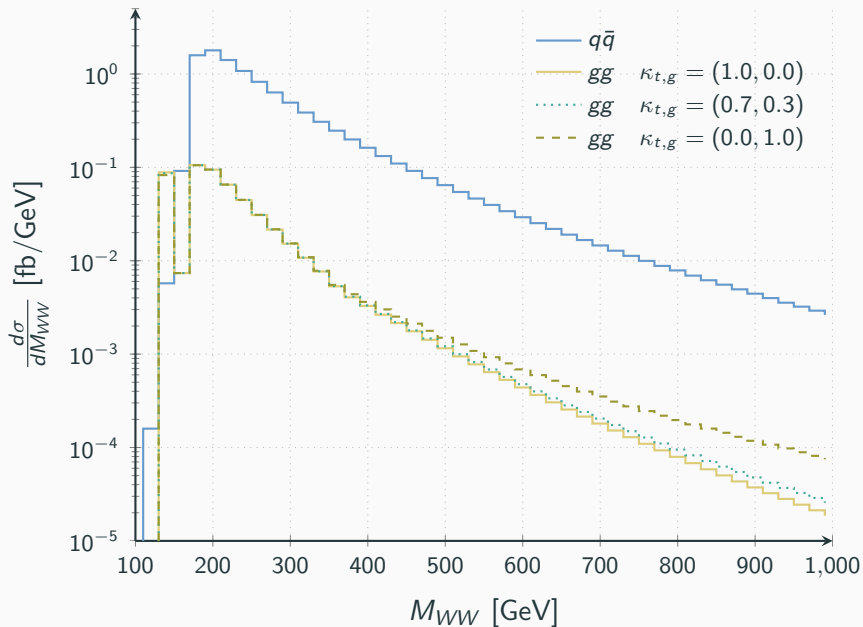
For this to be phenomenologically relevant:

- Higgs total cross section must be kept constant: strongly constrained by experimental measurements
- Vary ggH and $t\bar{t}H$ couplings together
+ fix $\kappa_t + \kappa_g = 1$

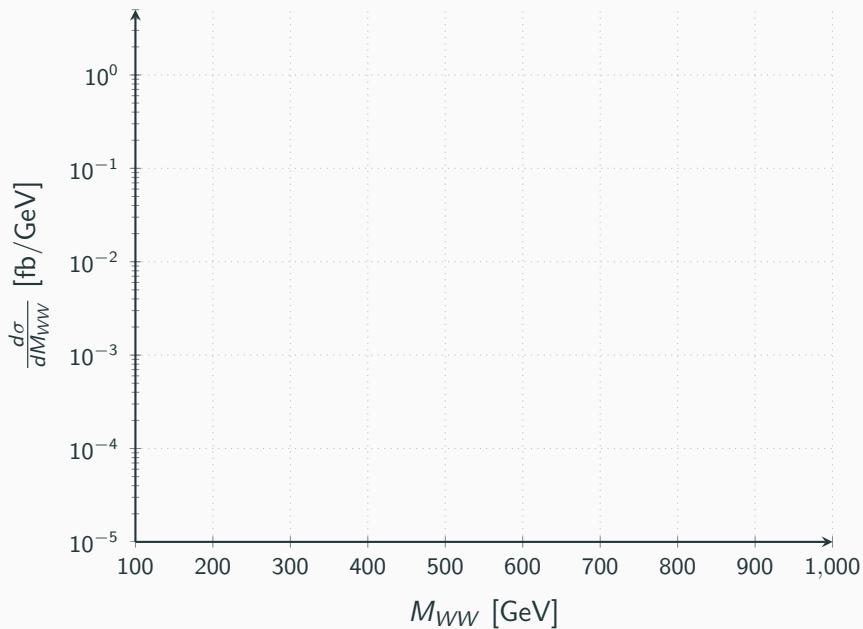
$q\bar{q}(gg) \rightarrow WW \rightarrow l^+\nu l^-\bar{\nu}$ (fixed-order)



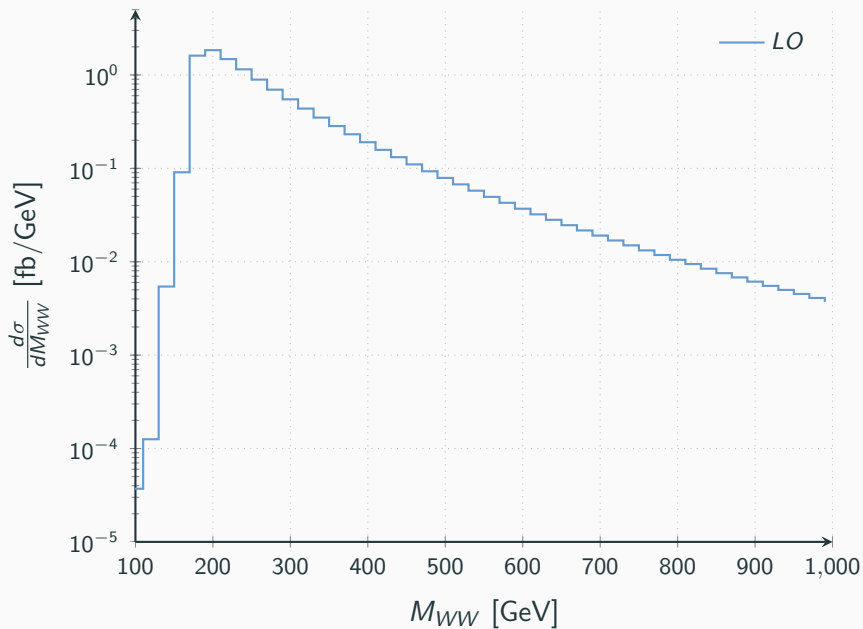
$q\bar{q}(gg) \rightarrow WW \rightarrow l^+\nu l^-\bar{\nu}$ (resummed)



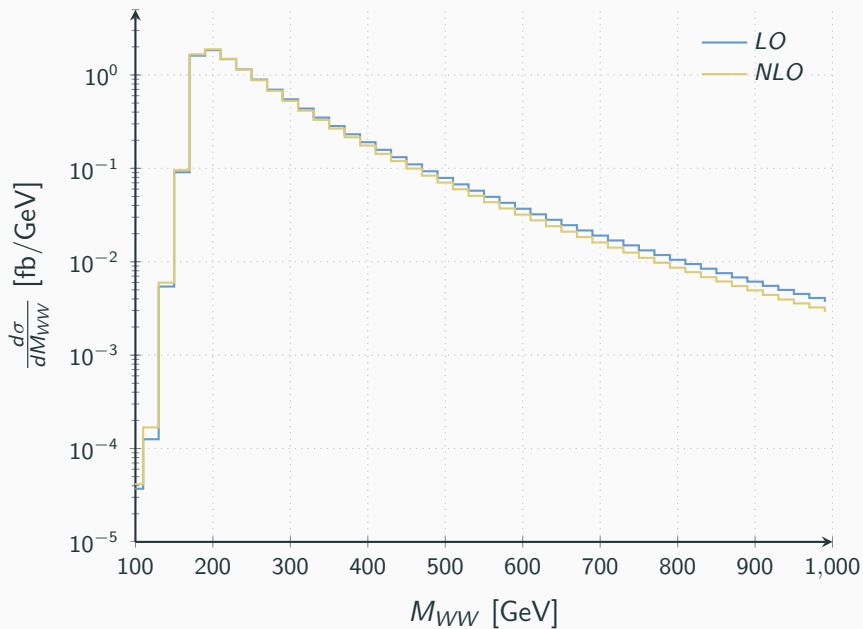
$q\bar{q} \rightarrow WW \rightarrow l^+\nu l^-\bar{\nu}$ (fixed-order vs. resummed)



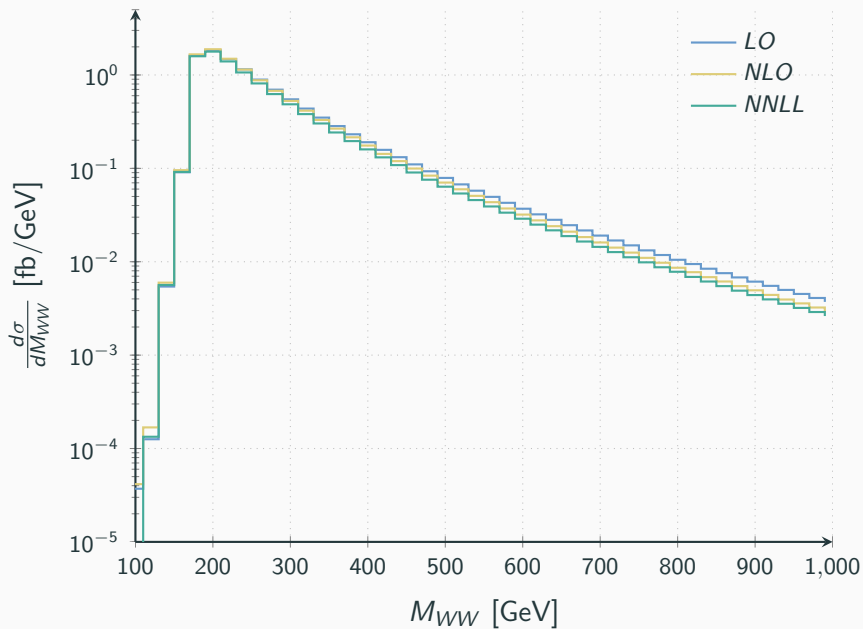
$q\bar{q} \rightarrow WW \rightarrow l^+\nu l^-\bar{\nu}$ (fixed-order vs. resummed)



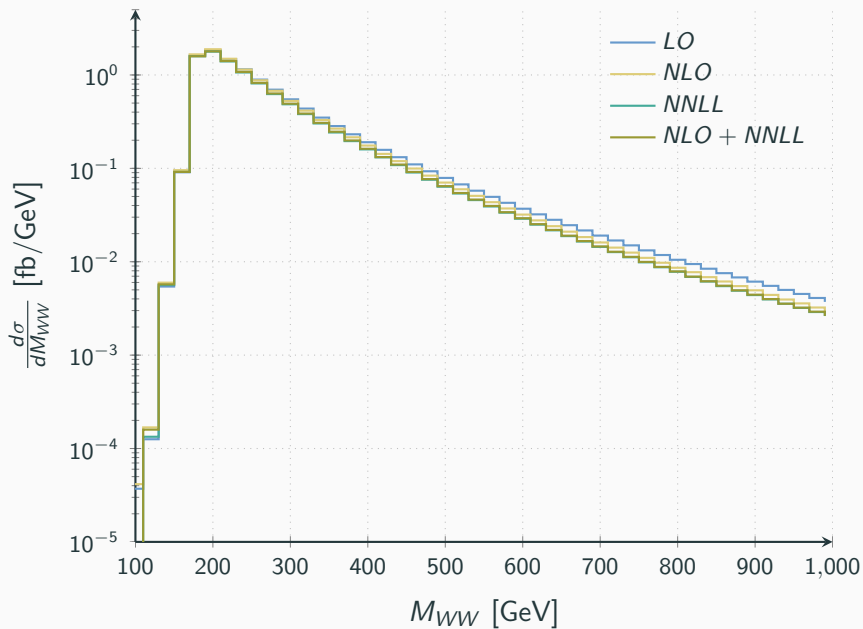
$q\bar{q} \rightarrow WW \rightarrow l^+\nu l^-\bar{\nu}$ (fixed-order vs. resummed)



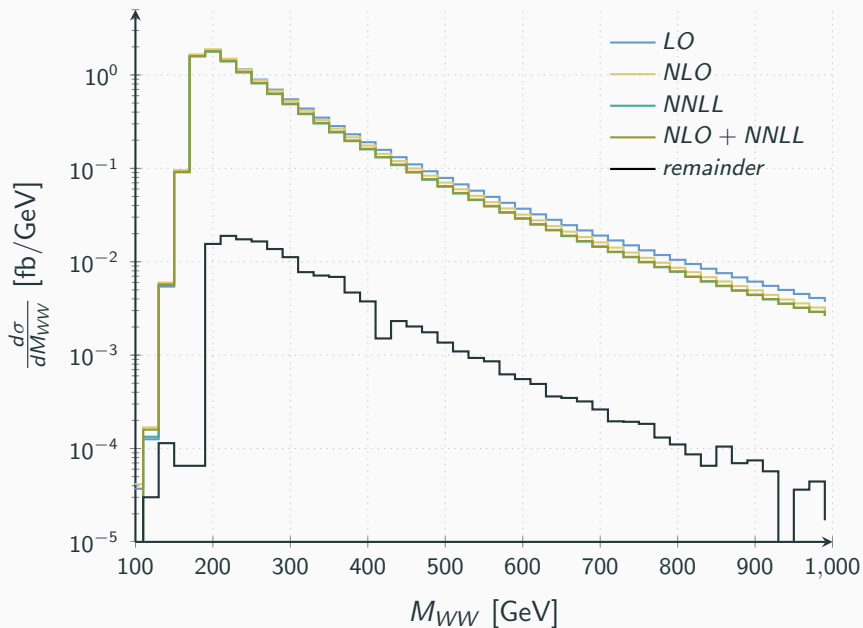
$q\bar{q} \rightarrow WW \rightarrow l^+\nu l^-\bar{\nu}$ (fixed-order vs. resummed)



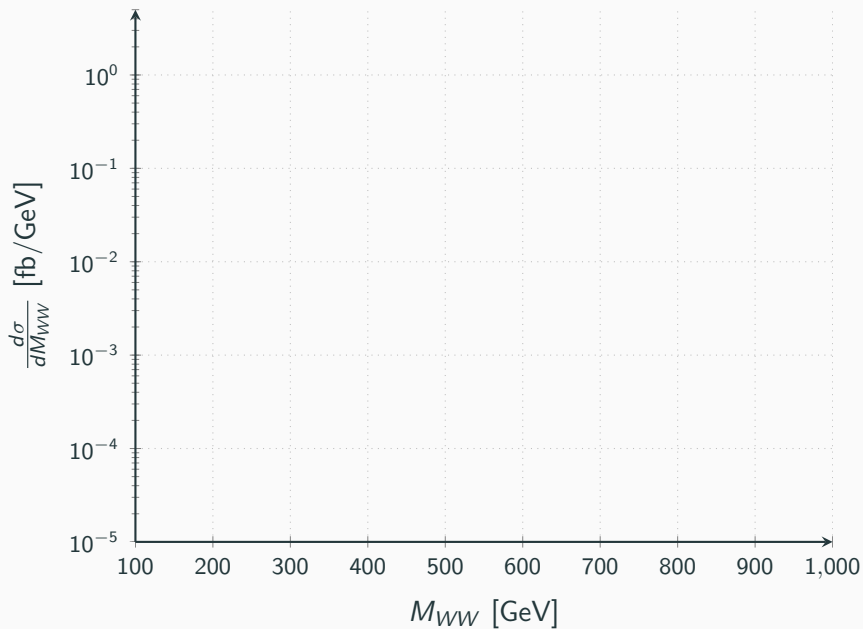
$q\bar{q} \rightarrow WW \rightarrow l^+\nu l^-\bar{\nu}$ (fixed-order vs. resummed)



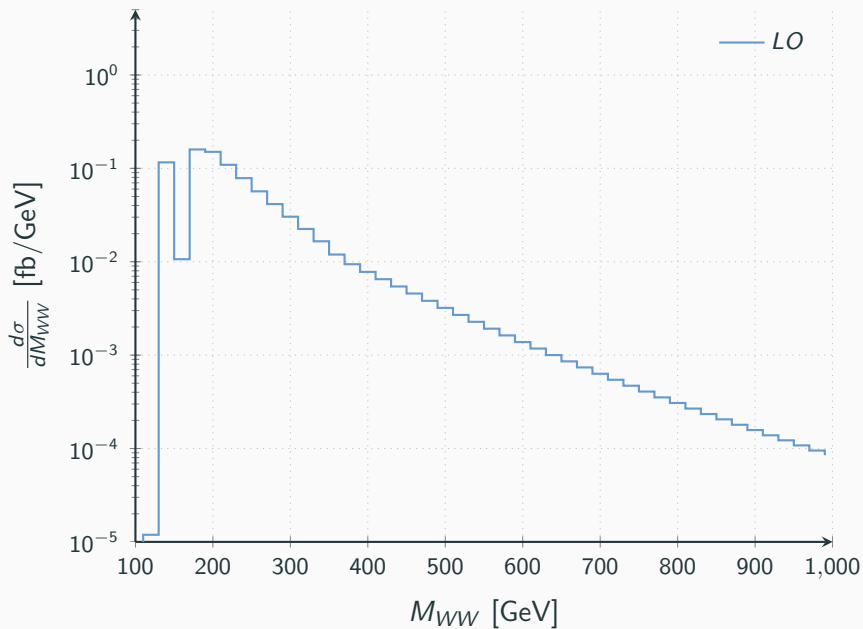
$q\bar{q} \rightarrow WW \rightarrow l^+\nu l^-\bar{\nu}$ (fixed-order vs. resummed)



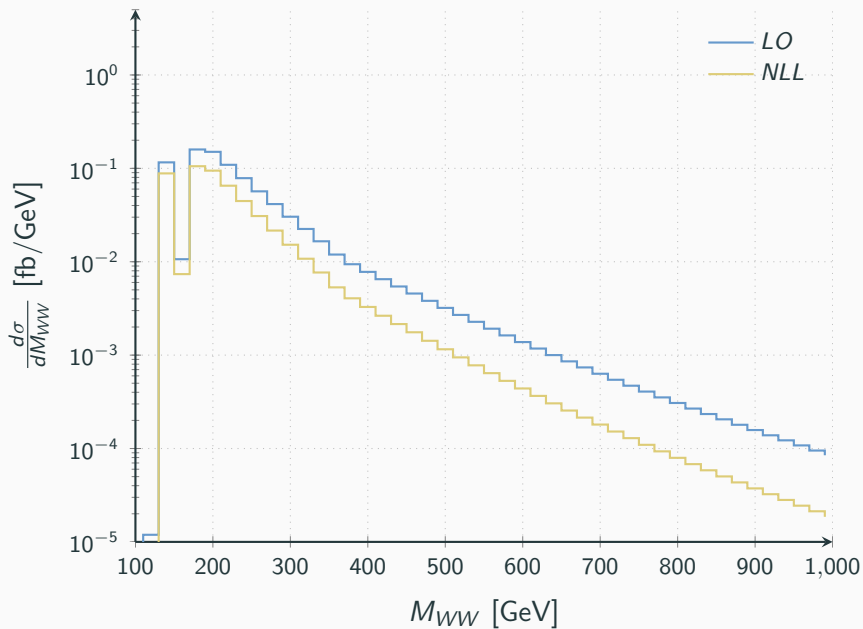
$gg \rightarrow WW \rightarrow l^+ \nu l^- \bar{\nu}$ (fixed-order vs. resummed)



$gg \rightarrow WW \rightarrow l^+ \nu l^- \bar{\nu}$ (fixed-order vs. resummed)



$gg \rightarrow WW \rightarrow l^+ \nu l^- \bar{\nu}$ (fixed-order vs. resummed)

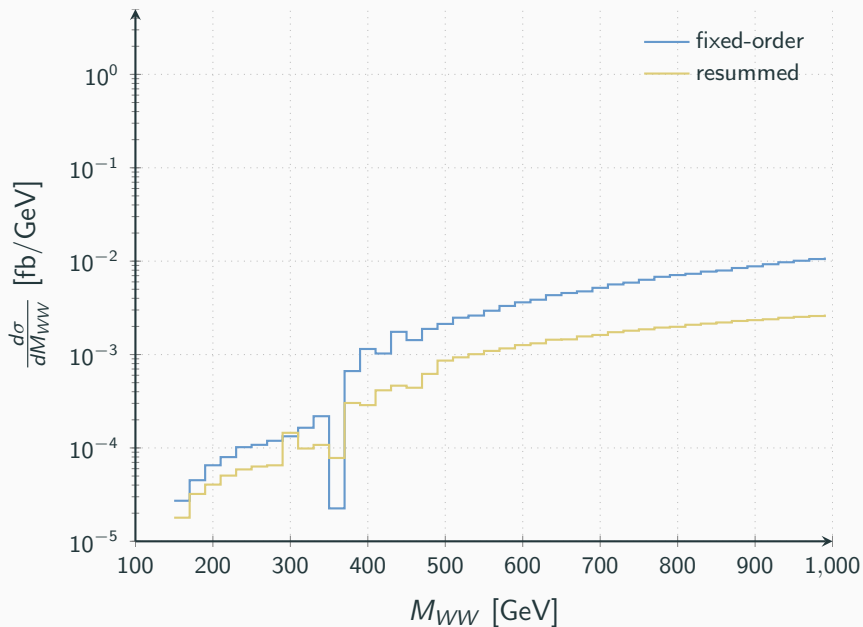


Comments and Analysis

- Resummation suppresses the gg initiated contributions much more than the $q\bar{q}$ initiated

- Resummation suppresses the gg initiated contributions much more than the $q\bar{q}$ initiated
- Parameterise the discrepancy with respect to the Standard Model:

$$\delta = \frac{\sigma_{gg}^{BSM} - \sigma_{gg}^{SM}}{\sigma^{SM}}$$

δ (fixed-order vs. resummation)

Conclusion

- Automation in resummation is the way to go for complicated studies: experimental cuts, spin correlated final states, . . .

- Automation in resummation is the way to go for complicated studies: experimental cuts, spin correlated final states, . . .
- WIP comparison with Parton Shower Monte Carlos