

# WE-HERAEUS PHYSICS SCHOOL 2015

## NEW THEORETICAL RESULTS IN UPC

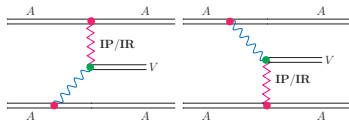
Mariola Klusek-Gawenda

Institute of Nuclear Physics PAS Kraków

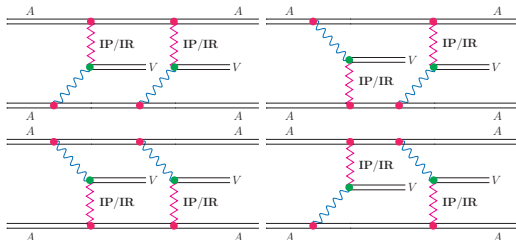


# MAIN GOAL - VECTOR MESONS PRODUCTION

- ▶ Single  $\rho^0$  &  $J/\psi$  production



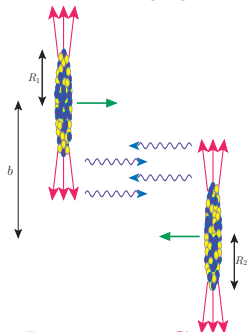
- ▶ Double-scattering mechanism



$$\text{Br}(\rho^0 \rho^0 \rightarrow \pi^+ \pi^- \pi^+ \pi^-) \simeq 100\%$$

STAR, CMS & ALICE data

## EQUIVALENT PHOTON APPROXIMATION

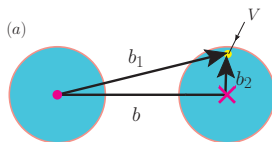


The strong electromagnetic field is a source of photons that can induce electromagnetic reactions in ion-ion collisions.

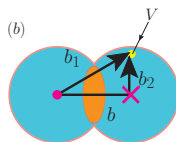
ULTRAPERIPHERAL COLLISIONS

SEMI-CENTRAL COLLISIONS

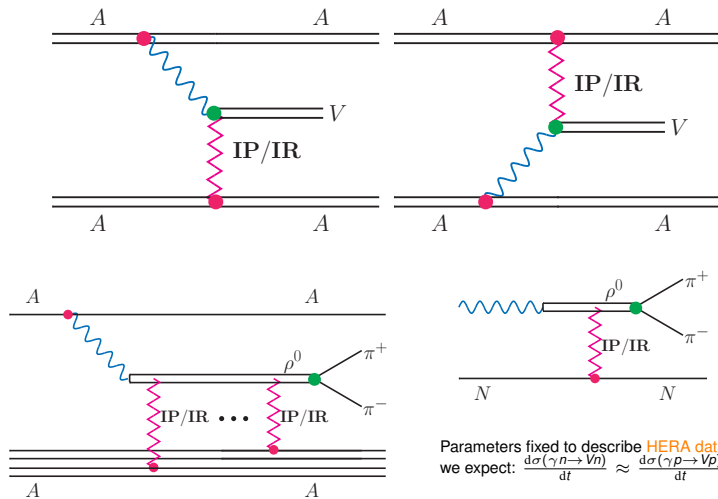
$$b > R_{min} = R_1 + R_2$$



$$b \leq R_{min}$$

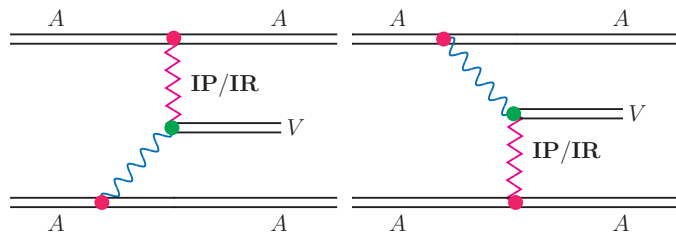


## PHOTOPRODUCTION OF VECTOR MESON



Parameters fixed to describe HERA data  
we expect:  $\frac{d\sigma(\gamma n \rightarrow Vn)}{dt} \approx \frac{d\sigma(\gamma p \rightarrow Vp)}{dt}$

## SINGLE VECTOR MESON PRODUCTION



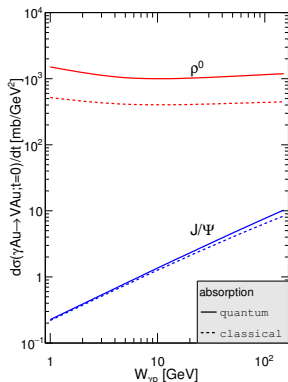
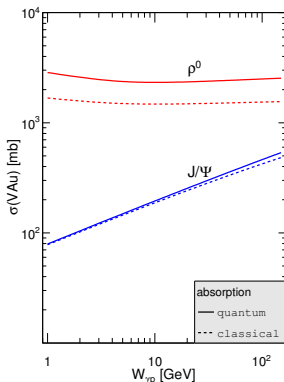
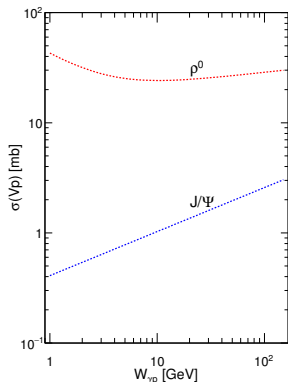
$$\frac{d\sigma_{A_1 A_2 \rightarrow A_1 A_2 V}}{d^2 b dy} = \frac{dP_{\gamma \mathbf{P}}(b, y)}{dy} + \frac{dP_{\mathbf{P} \gamma}(b, y)}{dy} \quad (1)$$

$$\frac{dP_{1/2}(b, y)}{dy} = \omega_{1/2} N(\omega_{1/2}, b) \sigma_{\gamma A_{2/1} \rightarrow V A_{2/1}}(W_{\gamma A_{2/1}}) \quad (2)$$

- ▶ point-like  $F(q) = 1$
- ▶ realistic  $F(q) = \frac{4\pi}{q} \int \rho(r) \sin(qr) r dr$
- ▶ monopole  $F(q) = \frac{\Lambda^2}{\Lambda^2 + q^2}$ ,  $\Lambda = \sqrt{\frac{6}{\langle r^2 \rangle}}$

$$\frac{d\sigma_{\gamma p \rightarrow Vp}(t=0)}{dt}$$

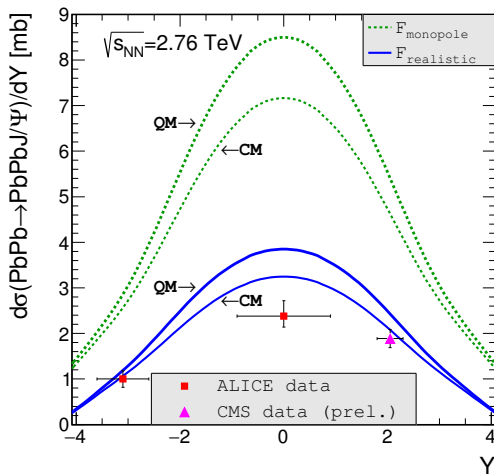
← HERA data



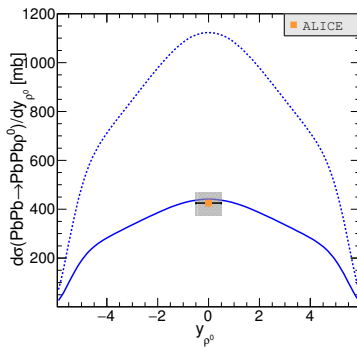
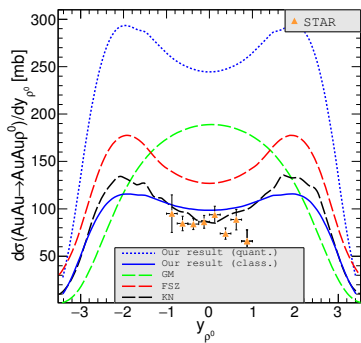
- ▶  $\sigma_{tot}^{CM}(VA) = \int d^2\mathbf{r} (1 - \exp(-\sigma_{tot}(Vp)T_A(\mathbf{r})))$
- ▶  $\sigma_{tot}^{QM}(VA) = 2 \int d^2\mathbf{r} (1 - \exp(-\frac{1}{2}\sigma_{tot}(Vp)T_A(\mathbf{r})))$

$$\sigma_{\gamma A \rightarrow VA}(W_{\gamma A_{2/1}}) = \frac{d\sigma_{\gamma A \rightarrow VA}(t=0)}{dt} \int_{-\infty}^{t_{max}} dt \left| F_A(t) \right|^2$$

# $J/\psi$ MESON PRODUCTION



# $\rho^0$ MESON PRODUCTION

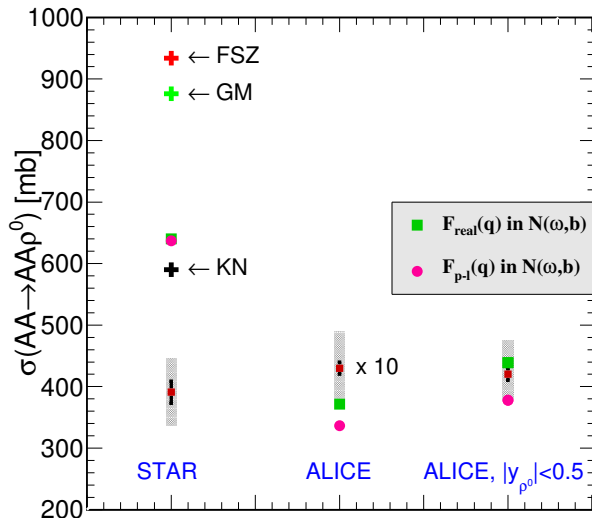


**GM** - V.P. Gonçalves and M.V.T. Machado, "The QCD pomeron in ultraperipheral heavy ion collisions. IV. Photonuclear production of vector mesons", Eur. Phys. J. **C40** (2005) 519,

**FSZ** - L. Frankfurt, M. Strikman and M. Zhalov, "Signals for black body limit in coherent ultraperipheral heavy ion collisions", Phys. Lett. **B537** (2002) 51,

**KN** - S. Klein and J. Nystrand, "Exclusive vector meson production in relativistic heavy ion collisions", Phys. Rev. **C60** (1999) 014903



SINGLE  $\rho^0$  MESON PRODUCTION

# SMEARING OF $\rho^0$ MASS

Drell-Söding contribution:

$$A(m) = \mathcal{A}_{BW} \frac{\sqrt{mm_{\rho^0}\Gamma_{\rho^0}(m)}}{m^2 - m_{\rho^0}^2 + im_{\rho^0}\Gamma_{\rho^0}(m)} + \mathcal{B}_{\pi\pi}$$

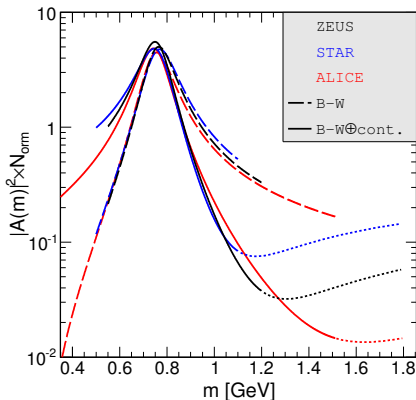
$$\Gamma_{\rho^0}(m) = \Gamma_{\rho^0} \frac{m_{\rho^0}}{m} \left( \frac{m^2 - 4m_{\pi}^2}{m_{\rho^0}^2 - 4m_{\pi}^2} \right)^{3/2}$$

Parameter	ZEUS	STAR	ALICE
$m_{\rho^0}$ [GeV]	$0.77 \pm 0.002$	$0.775 \pm 0.003$	$0.761 \pm 0.0023$
$\Gamma_{\rho^0}$ [GeV]	$0.146 \pm 0.003$	$0.162 \pm 0.007$	$0.1502 \pm 5.5$
$\left  \frac{\mathcal{B}_{\pi\pi}}{\mathcal{A}_{BW}} \right $ [GeV $^{-1/2}$ ]	0.669	$0.89 \pm 0.08$	$0.5 \pm 0.04$
$m$ [GeV]	(0.55 – 1.2)	(0.5 – 1.1)	(0.28 – 1.512)

SMEARING OF  $\rho^0$  MASS

Drell-Söding contribution:

A



$$- + \mathcal{B}_{\pi\pi}$$

$$3/2$$

Parameter

 $m_{\rho^0}$  [GeV] $\Gamma_{\rho^0}$  [GeV] $\left| \frac{\mathcal{B}_{\pi\pi}}{\mathcal{A}_{B\mathcal{W}}} \right|$  [%] $m$  [GeV]

(0.55 – 1.2)

(0.5 – 1.1)

ALICE

3 0.761 ± 0.0023

7 0.1502 ± 5.5

0.5 ± 0.04

(0.28 – 1.512)

INTRODUCTION

EPA

THEORY

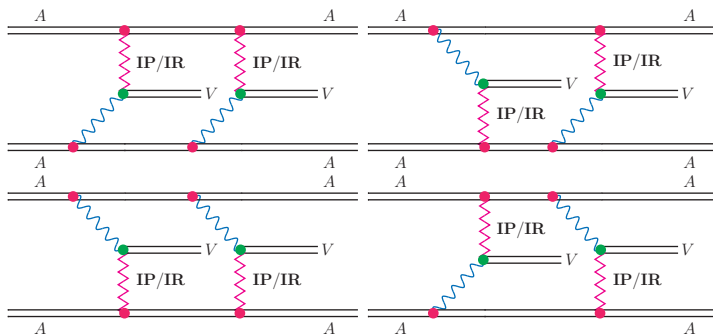
PHOTOPRODUCTION  
OF VECTOR  
MESONS

SINGLE MESON PRODUCTION

 $J/\psi$  MESON PRODUCTION $\rho^0$  MESON PRODUCTIONSMEARING OF  $\rho^0$  MASSDOUBLE-SCATTERING  
MECHANISM

CONCLUSION

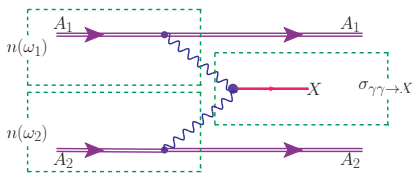
# DOUBLE-SCATTERING MECHANISM



$$\frac{d\sigma_{A_1 A_2 \rightarrow A_1 A_2 \rho^0 \rho^0}}{dy_1 dy_2} = \frac{1}{2} \int \left( \frac{dP_{\gamma \mathbf{P}}(b, y_1)}{dy_1} + \frac{dP_{\mathbf{P} \gamma}(b, y_1)}{dy_1} \right) \times \left( \frac{dP_{\gamma \mathbf{P}}(b, y_2)}{dy_2} + \frac{dP_{\mathbf{P} \gamma}(b, y_2)}{dy_2} \right) d^2 b$$

( $\rho^0$ 's have negligibly small transverse momenta)

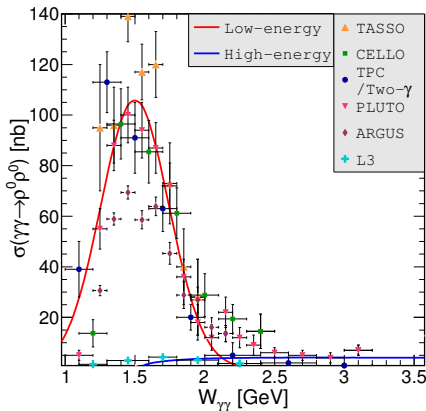
# $AA \rightarrow AA\rho^0\rho^0 - \gamma\gamma$ FUSION



$$n(\omega) = \int_{R_{min}}^{\infty} 2\pi b db N(\omega, b)$$

## Reference:

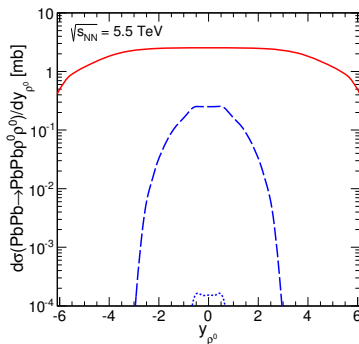
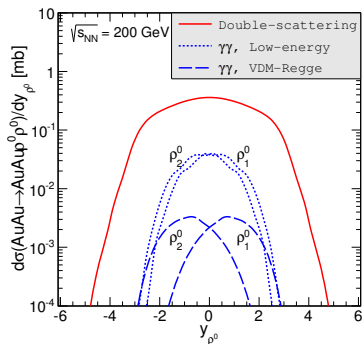
M. Klusek, W. Schäfer and A. Szczurek,  
 Phys.Lett. **B674** (2009) 92,  
 "Exclusive production of  $\rho^0\rho^0$  pairs in  $\gamma\gamma$   
 collisions at RHIC"  
 + back-up slide



# DOUBLE-SCATTERING MECHANISM




VS

$\gamma\gamma$  FUSION



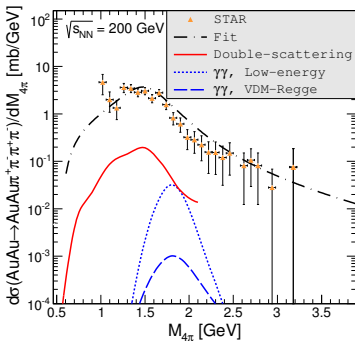
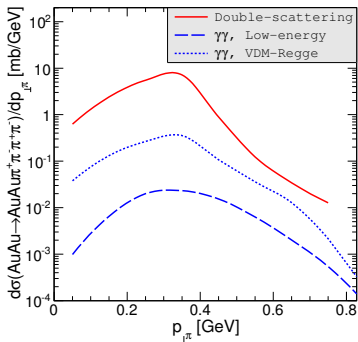
## Reference:

M. Klusek-Gawenda and A. Szczurek, Phys. Rev. **C89** (2014) 024912

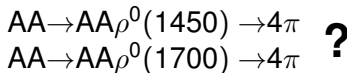
"Double-scattering mechanism in the exclusive  $AA \rightarrow AA p^0$  reaction in ultrarelativistic collisions",   

# $\pi^+\pi^-\pi^+\pi^-$ PRODUCTION @ RHIC

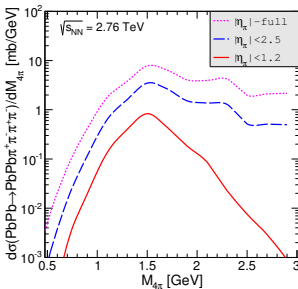
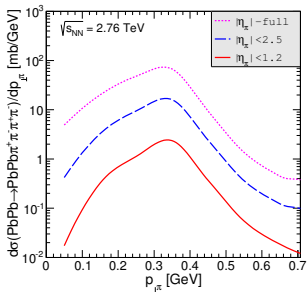
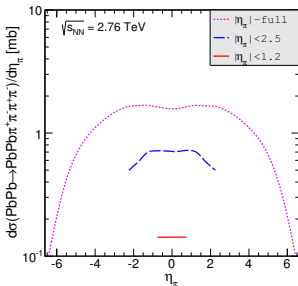
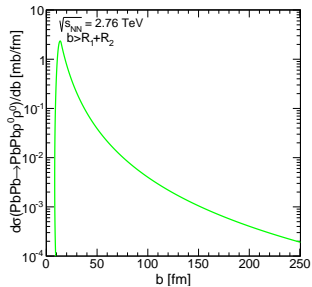
$$|\eta_\pi| < 1$$



missing  
mechanisms:

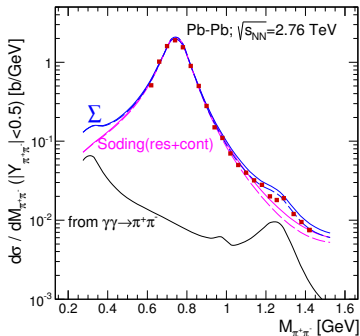
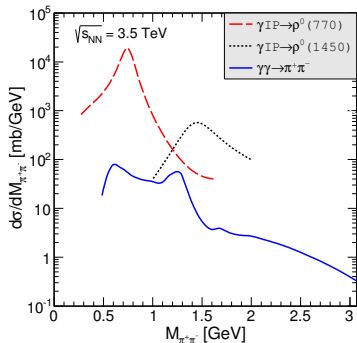


# $\pi^+\pi^-\pi^+\pi^-$ PRODUCTION @ LHC





# TWO-PION PRODUCTION



## Reference:

M. Klusek-Gawenda and A. Szczurek,  
Phys. Rev. **C87** (2013) 054908

" $\pi^+\pi^-$  and  $\pi^0\pi^0$  pair production in photon-photon  
and in ultraperipheral ultrarelativistic heavy ion  
collisions",

**Drell-Söding+ $f_2(1270)$**

colored solid lines -  $\Gamma_{\rho^0} = 150.2$  MeV

colored dashed lines -  $\Gamma_{\rho^0} = 140$  MeV

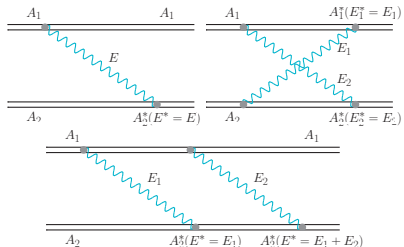
ALICE data arXiv:1503.09177

# CONCLUSIONS

- ▶ Impact parameter space approach
  - ▶ Good description of
    - ▶ STAR and ALICE data for  $\rho^0(770)$  production
    - ▶ CMS and ALICE data for  $J/\psi$  production
  - ▶ Smearing of  $\rho^0$  meson
  - ▶ Comparison of four-pion production via  $\rho^0\rho^0$  production
    - ▶  $\gamma\gamma$  fusion
    - ▶ nuclear double-photoproduction (**very large**)
- with STAR data

# CONCLUSIONS

## ▶ Multiple Coulomb excitations

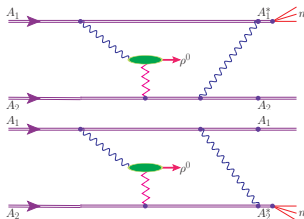
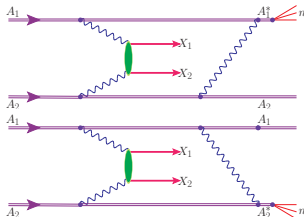
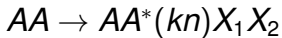


### Reference:

M. Kłusek-Gawenda, M. Ciemala, W. Schäfer and A. Szczurek,  
Phys. Rev. **C89** (2014) 054907,

"Electromagnetic excitation of nuclei and neutron evaporation  
in ultrarelativistic ultraperipheral heavy ion collisions"

associated with  $\rho^0\rho^0$  or  $\rho^0$   
production may cause additional  
excitation of one or both nuclei to  
the giant resonance region





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# Back-up slides

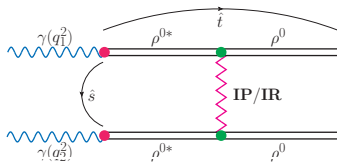
ELEMENTARY CROSS SECTION  $\gamma\gamma \rightarrow \rho^0\rho^0$ 

$$\sigma_{\gamma\gamma \rightarrow \rho^0\rho^0}^{\text{high-energy}} = \int_{\hat{t}_{\min}(\hat{s})}^{\hat{t}_{\max}(\hat{s})} \frac{d\sigma_{\gamma\gamma \rightarrow \rho^0\rho^0}^{\text{high-energy}}}{d\hat{t}} d\hat{t} \quad (3)$$

$$\frac{d\sigma_{\gamma\gamma \rightarrow \rho^0\rho^0}^{\text{high-energy}}}{d\hat{t}} = \frac{1}{16\pi\hat{s}} \left| \mathcal{M}_{\gamma\gamma \rightarrow \rho^0\rho^0}(\hat{s}, \hat{t}; q_1, q_2) \right|^2 \quad (4)$$

$$\mathcal{M}_{\gamma\gamma \rightarrow \rho^0\rho^0}(\hat{s}, \hat{t}; q_1, q_2) = C_{\gamma \rightarrow \rho^0} C_{\gamma \rightarrow \rho^0} \mathcal{M}_{\rho^{0*}\rho^{0*} \rightarrow \rho^0\rho^0}(\hat{s}, \hat{t}; q_1, q_2) \quad (5)$$

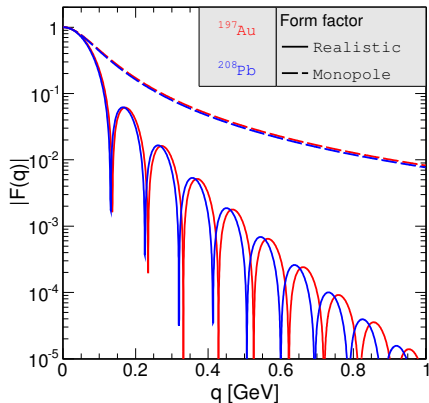
$$\begin{aligned} \mathcal{M}_{\rho^{0*}\rho^{0*} \rightarrow \rho^0\rho^0}(\hat{s}, \hat{t}; q_1, q_2) &= \left( \eta_{\mathbf{P}}(\hat{s}, \hat{t}) C_{\mathbf{P}} \left( \frac{\hat{s}}{s_0} \right)^{\alpha_{\mathbf{P}}(\hat{t})-1} + \eta_{\mathbf{R}}(\hat{s}, \hat{t}) C_{\mathbf{R}} \left( \frac{\hat{s}}{s_0} \right)^{\alpha_{\mathbf{R}}(\hat{t})-1} \right) \\ &\times \hat{s} F(\hat{t}; q_1^2 \approx 0) F(\hat{t}; q_2^2 \approx 0) \end{aligned} \quad (6)$$



## FORM FACTOR

REALISTIC  $F_{em}$ 

$$F(q) = \frac{4\pi}{q} \int \rho(r) \sin(qr) r dr$$

MONOPOLE  $F_{em}$ 

$$F(q^2) = \frac{\Lambda^2}{\Lambda^2 + q^2}$$

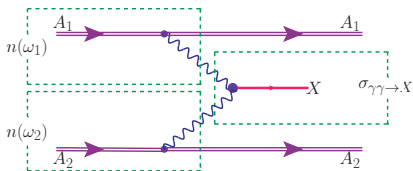
$$\Lambda = \sqrt{\frac{6}{\langle r^2 \rangle}}$$

- ▶  $^{197}\text{Au} \Rightarrow \sqrt{\langle r^2 \rangle} = 5.3 \text{ fm}$ ,  
 $\Lambda = 91 \text{ MeV}$ ,
- ▶  $^{208}\text{Pb} \Rightarrow \sqrt{\langle r^2 \rangle} = 5.5 \text{ fm}$ ,  
 $\Lambda = 88 \text{ MeV}$ .

In the literature:

$$\Lambda = (80 - 90) \text{ GeV}$$

## NUCLEAR CROSS SECTION



$$n(\omega) = \int_{R_{min}}^{\infty} 2\pi b db N(\omega, b) \quad (7)$$

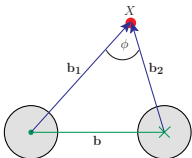
$$\sigma_{A_1 A_2 \rightarrow A_1 A_2 X} = \int d\omega_1 d\omega_2 n(\omega_1) n(\omega_2) \sigma_{\gamma\gamma \rightarrow X}(\omega_1, \omega_2)$$

$$= \dots$$

$$= \int N(\omega_1, \mathbf{b}_1) N(\omega_2, \mathbf{b}_2) S_{abs}^2(\mathbf{b})$$

$$\times \sigma_{\gamma\gamma \rightarrow X}(\sqrt{S_{A_1 A_2}})$$

$$\times 2\pi b db d\bar{b}_x d\bar{b}_y \frac{W_{\gamma\gamma}}{2} dW_{\gamma\gamma} dY_X \quad (8)$$







# MUTUAL ELECTROMAGNETIC DISSOCIATION

WE-HERAEUS  
PHYSICS SCHOOL

MARIOLA K-G

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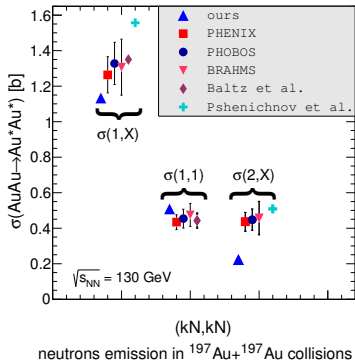
CONCLUSION

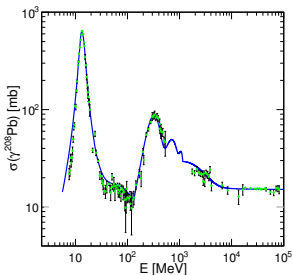
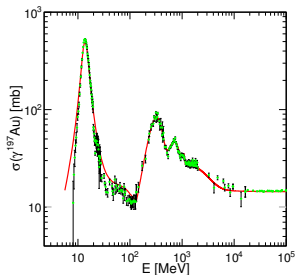
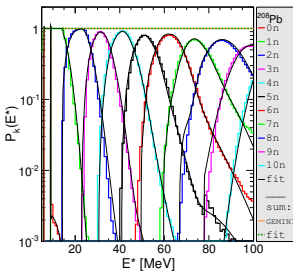
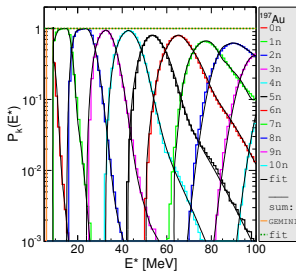
$$\sigma_{tot}(A_1 A_2 \rightarrow A_1^* A_2^*) = \int d^2\mathbf{b} P_{surv}(\mathbf{b}) P_{A_2}^{exc}(\mathbf{b}) P_{A_1}^{exc}(\mathbf{b})$$

$$P_{surv}(\mathbf{b}) \approx \theta(|\mathbf{b}| - (R_1 - R_2))$$

$$P_{A_2}^{exc}(\mathbf{b}) \approx \bar{n}_{A_2}(\mathbf{b}) \exp[-\bar{n}_{A_2}(\mathbf{b})]$$

$$\bar{n}_{A_2}(\mathbf{b}) = \int_{E_{min}}^{\infty} dE N_{A_1}(E, \mathbf{b}) \sigma_{tot}(\gamma A_2; E)$$



Photoabsorption  $\sigma(\gamma A \rightarrow A) \downarrow$ Probability of neutron multiplicity  $\downarrow$ 

Excitation function  $\gamma^A X \rightarrow kn^{A-1} X \downarrow$ 