



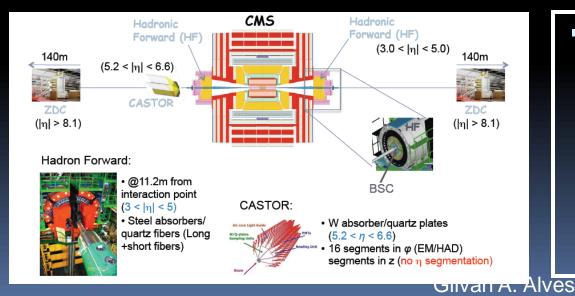
Review of diffractive and electromagnetic processes in CMS/LHC

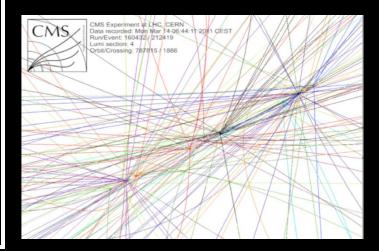
Gilvan A. Alves - Lafex/CBPF for the CMS collaboration

International Wilhelm and Else Heraeus Physics School Bad Honnef 2015

The setting:CMS@LHC

- High energy and high luminosity
 - Allows high statistics precision measurements, and sensitivity to "rare" processes (hard diffraction, exclusive production)
 - But high luminosity comes with high
 "pileup" average 2-8 in 2010/2011, 21 in
 2012
 - Low pileup needed for some analysis





- Good detector coverage
 - Tracking to $|\eta| < 2.4$
 - Hadronic calorimeter (HF) to |η| < 5
 - Forward calorimeters (cover -6.6< η< -5.2 (CASTOR) and |η| > 8.1 (ZDC)







- Studying the exclusive production at CMS
 - $\gamma \gamma \longrightarrow W^+ W^-$
- Measurement of diffraction dissociation
 - SD
 - DD
- Z and γ + jets production
- DY in association with jets
- Many other interesting results not covered here
- <u>https://twiki.cern.ch/twiki/bin/view/CMSPublic/PhysicsResultsFSQ</u>





Studying the exclusive production

CMS-PAS-FSQ-13-008

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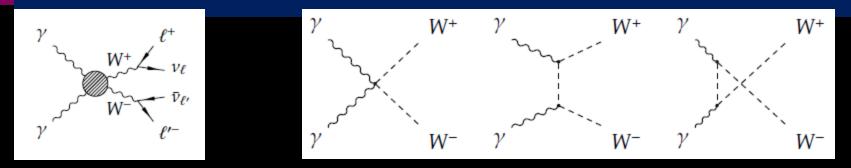
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Exclusive $\gamma\gamma \rightarrow WW$

CMS-PAS-FSQ-13-008





- Triple and quartic coupling in SM
 - Any deviation can signal new physics
- BSM contributions via effective Lagrangian (dimension 6+8 operators)

$$L_{6}^{0} = \frac{-e^{2}}{8} \frac{a_{0}^{W}}{\Lambda^{2}} F_{\mu\nu} F^{\mu\nu} W^{+\alpha} W_{\alpha}^{-} - \frac{e^{2}}{16 \cos^{2} \Theta_{W}} \frac{a_{0}^{Z}}{\Lambda^{2}} F_{\mu\nu} F^{\mu\nu} Z^{\alpha} Z_{\alpha},$$

$$L_{6}^{C} = \frac{-e^{2}}{16} \frac{a_{C}^{W}}{\Lambda^{2}} F_{\mu\alpha} F^{\mu\beta} (W^{+\alpha} W_{\beta}^{-} - W^{-\alpha} W_{\beta}^{+}) - \frac{e^{2}}{16 \cos^{2} \Theta_{W}} \frac{a_{C}^{Z}}{\Lambda^{2}} F_{\mu\alpha} F^{\mu\beta} Z^{\alpha} Z_{\beta},$$

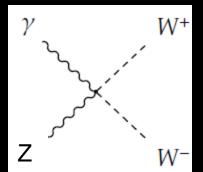
AQGC[†] parameters Λ : scale for New Physics

- Form factors introduced to preserve unitarity
 - W_{γγ}: γγ center of mass energy
 - Λ_{cutoff} : energy cutoff scale ($\Lambda_{\text{cutoff}} \rightarrow \infty$ = no form factor)

 $a_{0,C}^{W}
ightarrow a_{0,C}^{W} \left(W_{\gamma\gamma}^{2}
ight) = a_{0,C}^{W} \left(1 + \frac{W_{\gamma\gamma}^{2}}{\Lambda_{max}^{2}} \right)^{-2}$







- Anomalous quartic coupling include Z_{\(\nabla\)} WW vertex
 - Constrain this vertex (null) give relations between 6 -8 operators

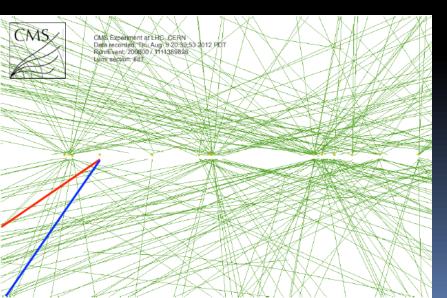
$$\frac{a_0^W}{\Lambda^2} = -\frac{4M_W^2}{g^2} \frac{f_{M,0}}{\Lambda^4} - \frac{8M_W^2}{g'^2} \frac{f_{M,2}}{\Lambda^4}$$
$$\frac{a_C^W}{\Lambda^2} = \frac{4M_W^2}{g^2} \frac{f_{M,1}}{\Lambda^4} + \frac{8M_W^2}{g'^2} \frac{f_{M,3}}{\Lambda^4}$$

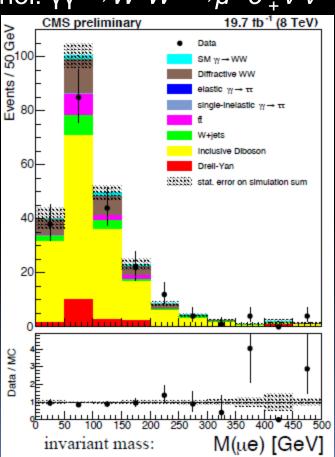
- $g = e/sin \Theta_w$
- g' = $e/\cos \Theta_w$

Sample selection $\gamma\gamma \rightarrow WW$



- Elastic and proton-dissociative contributions
 - $pp \rightarrow p W^+W^-p$
 - $pp \rightarrow p^{(*)} W^+W^-p^{(*)}$
- Unlike-flavor dilepton decay channel: $\gamma \gamma \rightarrow W^+ W^- \rightarrow \mu e_+^- v \overline{v}$
 - Avoid large backgrounds
- Data sample 19.7 fb⁻¹ @ 8 TeV
 - □ p_T(µ , e) > 20 GeV
 - No other track in vertex



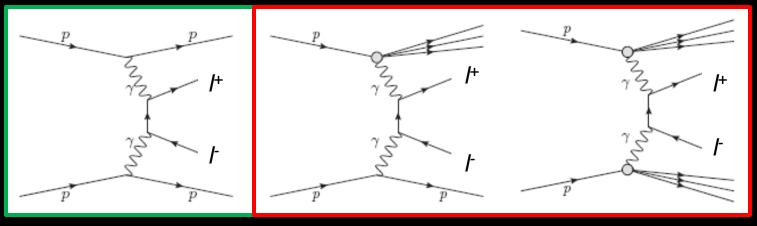


$\gamma \gamma \rightarrow \mu \mu$, ee control samples



Exclusive production

Proton dissociation

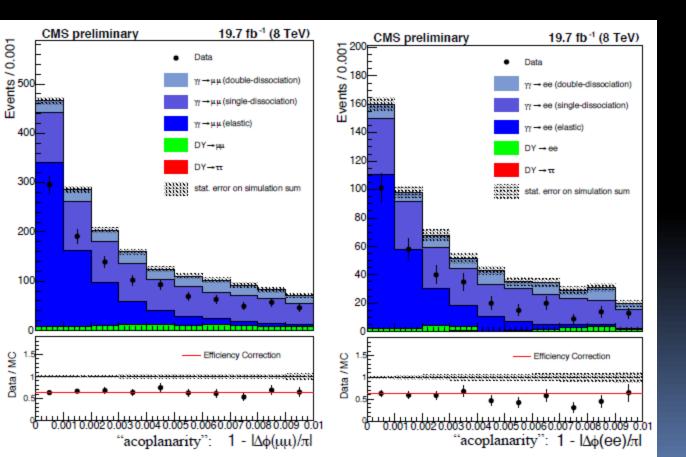


- Exclusive lepton production $pp \rightarrow p + p$
 - Well known QED like "Standard Candle"
 - Study exclusive selection efficiency
- Extract proton dissociation factor from $\gamma\gamma \rightarrow l^+l^-$ control sample
 - $\gamma\gamma \rightarrow WW$ simulation do not include proton dissociation
 - Apply dissociation factor from control sample

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$\gamma \gamma \rightarrow \mu \mu$, ee control samples

- Select high purity elastic events
 - Acoplanarity $1-|\Delta \phi(\ell+\ell-)/\pi| < 0.01$ (back to back leptons)
 - Invariant mass outside Z resonance -> m(l+l-)<70 GeV, m(l+l-)>106 GeV



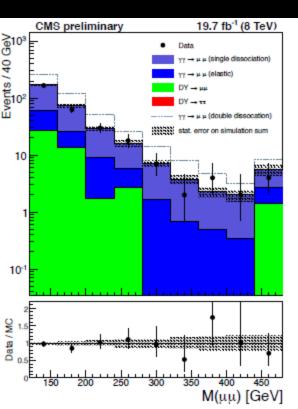
data-simulation
ratio: 0.63
apply as efficiency
correction

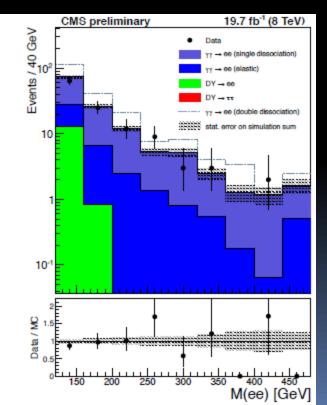
$\gamma \gamma \rightarrow \mu \mu$, ee proton dissociation



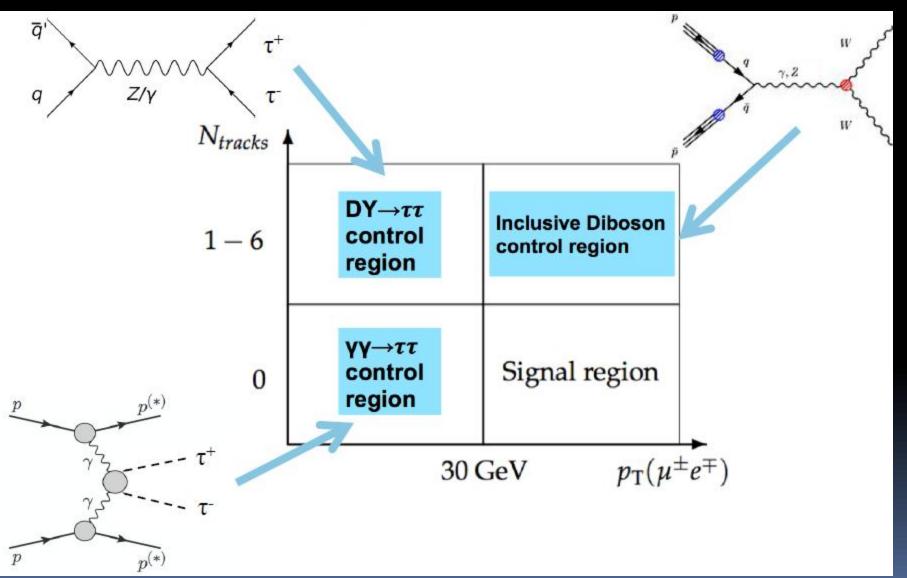
- → Use kinematic region close to $\gamma\gamma \rightarrow W+W-m(I^+I^-) > 160 \text{ GeV}$
- \rightarrow A correction factor is estimated from data / MC elastic
- → This factor applied to the predicted cross section for $\gamma\gamma \rightarrow W+W$ -

$$\begin{split} F &= \frac{N_{\mu\mu \ \rm data} - N_{\rm DY}}{N_{\rm elastic}} \bigg|_{m(\mu^+\mu^-) > 160 \, {\rm GeV}} \\ F &= 4.10 {\pm} 0.43 \end{split}$$



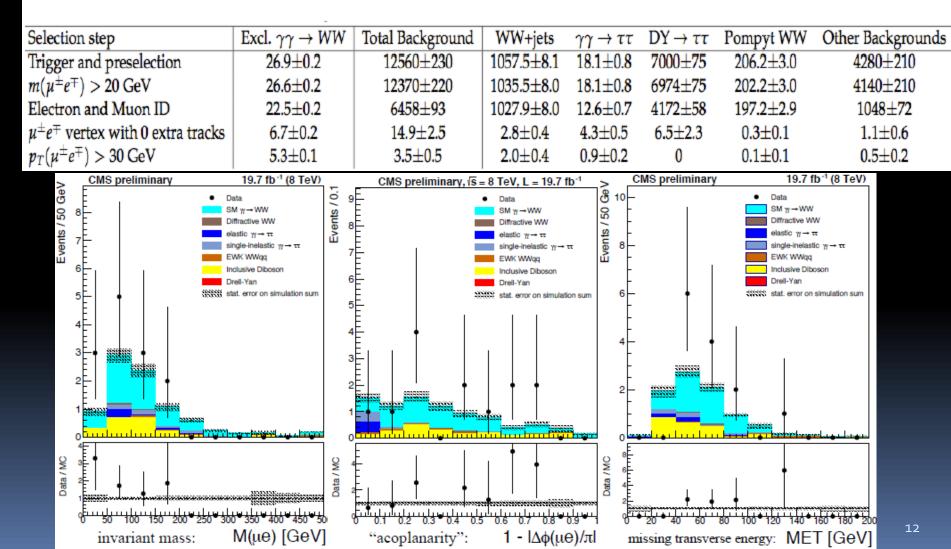


Signal and background $\gamma\gamma \rightarrow WW$



Signal and background $\gamma\gamma \rightarrow WW$

Number of expected signal and background events in simulation passing each selection step, normalized to an integrated luminosity of 19.7 fb⁻¹ :



Results $\gamma\gamma \rightarrow WW$

- 13 events observed
 - Expected 5.3 \pm 0.1 signal plus 3.5 \pm 0.5 background events
 - Data is 3.6 σ above background only hypothesis

measured exclusive W⁺W⁻ production cross section times branching ratio:

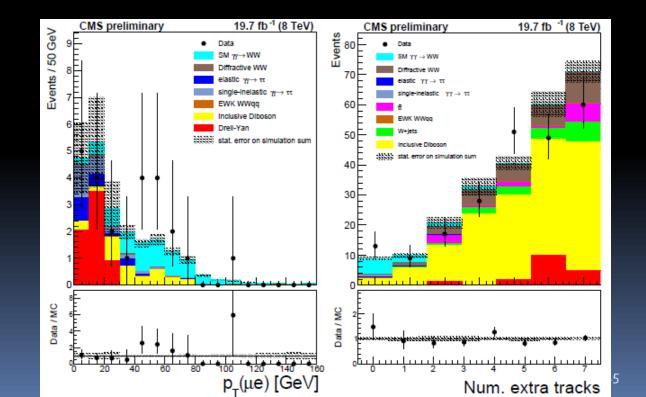
$$\sigma (pp \rightarrow p^{(*)}W^+W^- p^{(*)} \rightarrow p^{(*)}\mu^{\pm}e^{\mp} p^{(*)}) = 12.3^{+5.5}_{-4.4} \text{ fb}$$

■ SM – 6.9 ± 0.6 fb

Search for $AQGC\gamma\gamma \rightarrow WW$

R.

- Two bins are used looking for excess over SM
 - 30 GeV < $p_T(\mu e_+)$ < 130 GeV (SM dominated)
 - $p_T(\mu e_+) > 130 \text{ GeV}$ (Anomalous dominated)



95% CL intervals $\gamma\gamma \rightarrow WW$

Dimension 6 operators

$$-1.1 \times 10^{-4} < a_0^W / \Lambda^2 < 1.0 \times 10^{-4} \text{ GeV}^{-2} \ (a_C^W / \Lambda^2 = 0, \Lambda_{\text{cutoff}} = 500 \text{ GeV}), -4.2 \times 10^{-4} < a_C^W / \Lambda^2 < 3.4 \times 10^{-4} \text{ GeV}^{-2} \ (a_0^W / \Lambda^2 = 0, \Lambda_{\text{cutoff}} = 500 \text{ GeV}).$$

Compare to:

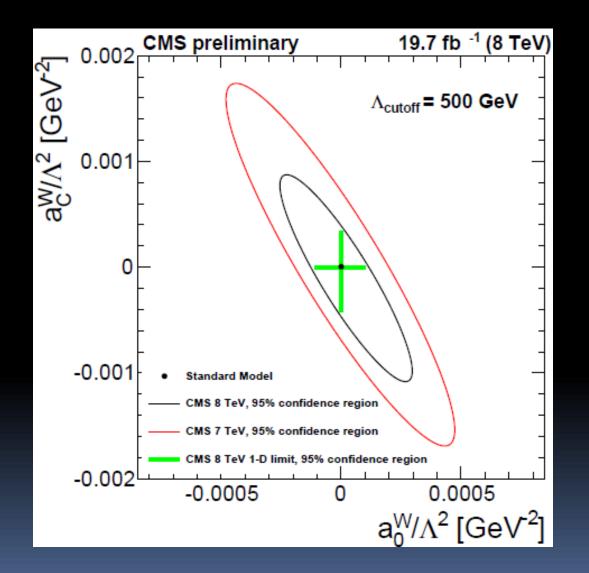
		OPAL (2004)	DØ (2013)	CMS (2013)
a_0^W/Λ^2 a_c^W/Λ^2	no form factor $\Lambda_{cutoff} = 500 \text{ GeV}$ no form factor $\Lambda_{cutoff} = 500 \text{ GeV}$	$\pm 2 \times 10^{-2}$ $^{+3.7}_{-5.2} \times 10^{-2}$	$\pm 4.3 \times 10^{-4}$ $\pm 2.5 \times 10^{-3}$ $\pm 1.5 \times 10^{-3}$ $\pm 9.2 \times 10^{-3}$	$\pm 4.0 \times 10^{-6}$ $\pm 1.5 \times 10^{-4}$ $\pm 1.5 \times 10^{-5}$ $\pm 5.0 \times 10^{-4}$

Dimension 8 operators (new from 8 TeV data)

$$\begin{split} -4.2 \times 10^{-10} &< f_{M,0} / \Lambda^4 < 3.8 \times 10^{-10} \,\text{GeV}^{-4} \,\left(\Lambda_{\text{cutoff}} = 500 \,\text{GeV}\right), \\ -16 \times 10^{-10} &< f_{M,1} / \Lambda^4 < 13 \times 10^{-10} \,\text{GeV}^{-4} \,\left(\Lambda_{\text{cutoff}} = 500 \,\text{GeV}\right), \\ -2.1 \times 10^{-10} &< f_{M,2} / \Lambda^4 < 1.9 \times 10^{-10} \,\text{GeV}^{-4} \,\left(\Lambda_{\text{cutoff}} = 500 \,\text{GeV}\right), \\ -8.0 \times 10^{-10} &< f_{M,3} / \Lambda^4 < 6.4 \times 10^{-10} \,\text{GeV}^{-4} \,\left(\Lambda_{\text{cutoff}} = 500 \,\text{GeV}\right). \end{split}$$

95% 2D limits $\gamma\gamma \rightarrow WW$









Diffractive Dissociation @ 7 TeV

CMS PAS FSQ 12-005 Phys. Rev. D 92, 012003 (2015)

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23/08/2015

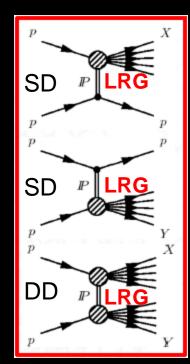
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Measurement of diffractive dissociation



- ≈ 25% of inelastic cross section can be attributed to diffractive processes → LRG
- Description based on Regge type models
- Diffractive measurements important for improving models
- Data sample 16.5 μ b⁻¹ low pileup (μ =0.14) @ 7 TeV
 - Minimum bias trigger
 - Hit in both BPTX and either BSC
- Offline selection
 - Large Rapidity Gap (LRG) tagging
 - At least 2 PF objects in BSC acceptance
 - No vertex requirement (low mass)
- MC simulation
 - PYTHIA8-MBR Minimum Bias Rockefeller model
 - PYTHIA8-4C for systematic studies

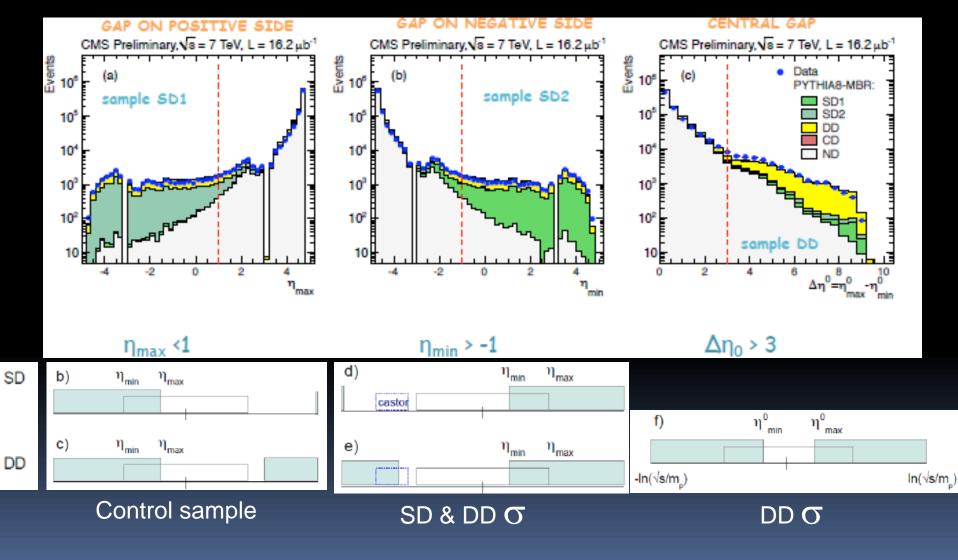




🗠 Experimental Topologies



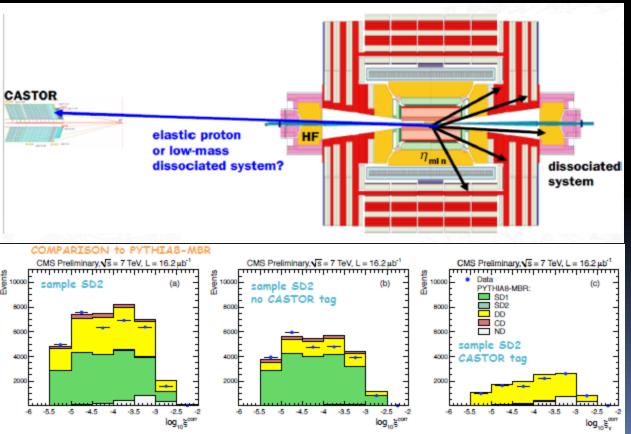
Based on the LRG position



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💬 Detector level distributions

- Proton fractional momentum loss $\xi = M^2 x/s$
 - M²x Mass of the dissociated system
- At detector level it is reconstructed as $\xi = \Sigma E^{i} p_{z}^{i}/\sqrt{s}$
 - Sum over all PF objects
 - ξ corrected (MC) for undetected particles (low E, low η)



Castor tag selects low mass systems $Mx \approx 3.2 \text{ GeV}$

Separate SD & DD

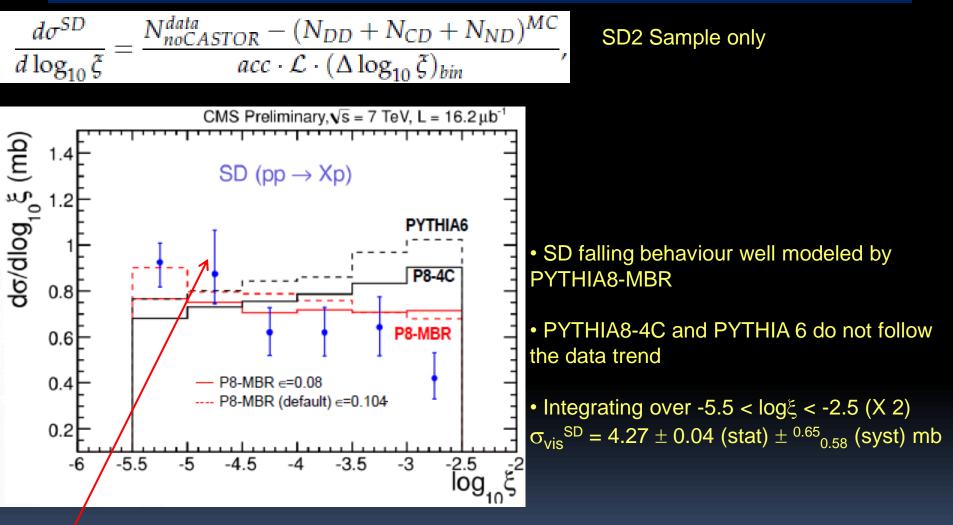






Measurements: SD cross section





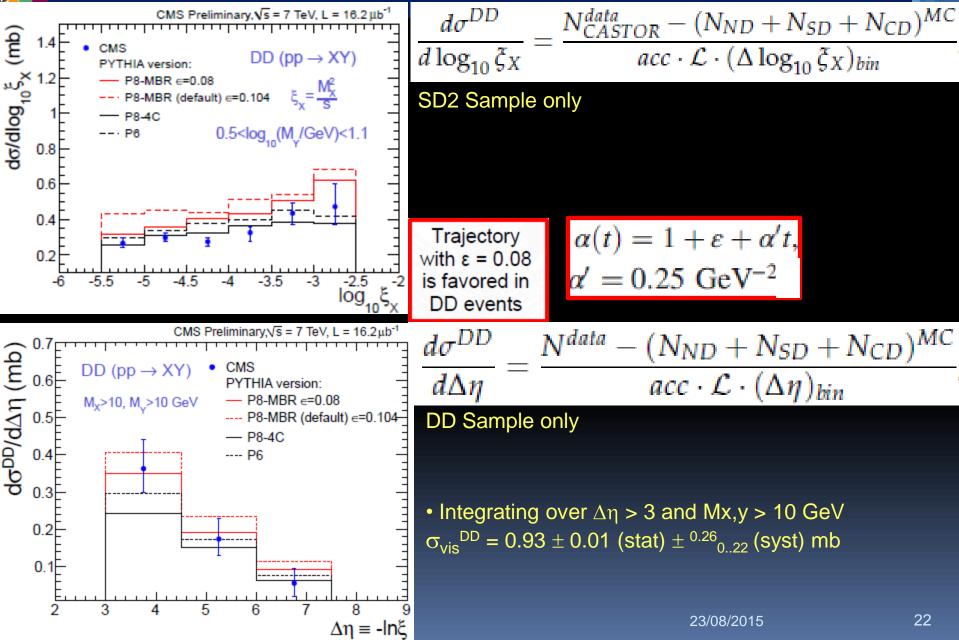
 Systematic dominated by energy scale and background subtraction

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CM

Measurements: DD cross section

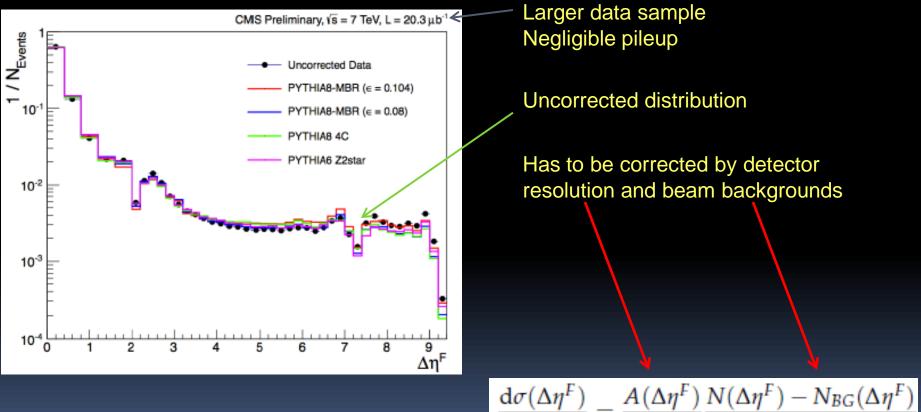




Measurements: LRG cross section



- Difficult to measure the whole $Mx \rightarrow$ measure size of LRG
 - Inclusive measure the largest forward gap $\Delta \eta_F = \max(4.7 \eta_{max}, 4.7 + \eta_{min})$
 - largest gap between each edge of the detector and the position in η of the first particle moving away from the edge

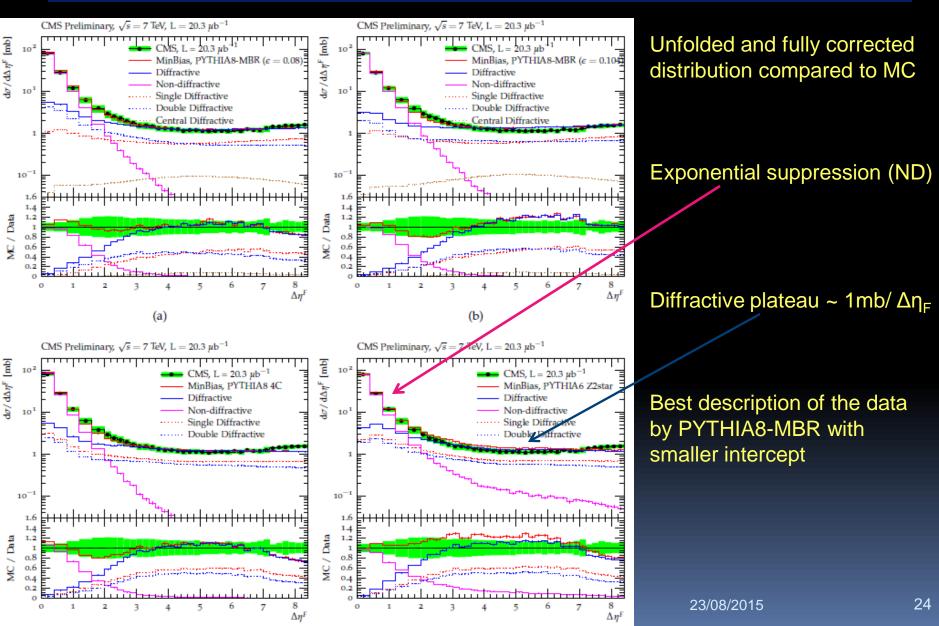


 $\Delta \eta_{\rm bin}$



Measurements: LRG cross section

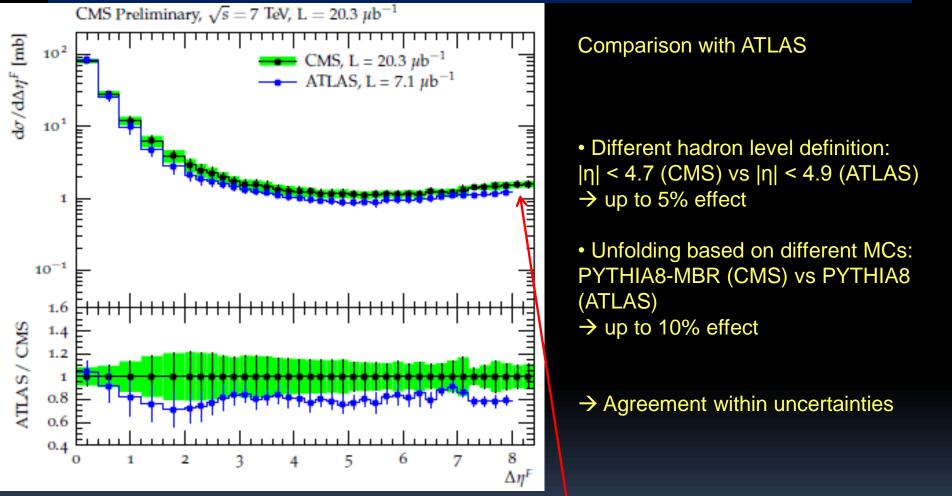






Measurements: LRG cross section





□ → CMS result extends ATLAS measurement by 0.4 unit of gap size





Z+*jets* and γ + *jets* production @ 8 TeV

CMS PAS SMP-14-005

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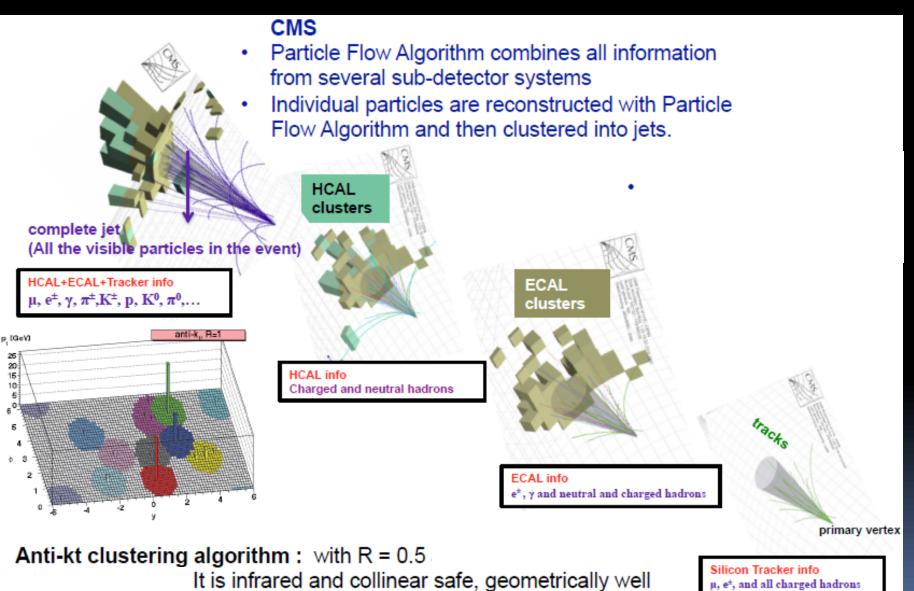


\mathcal{C} Z+jets $|\gamma + j$ ets Motivation

- Important test of the SM as well BSM searches
 - NLO calculation available (BLACKHAT)
 - Cross section ratio useful for estimate higher order effects
 - Can help reduce uncertainties on Z -> invisible
 - Used in BSM searches
 - Measure cross sections and ratio as function of p_T
- Data sample & selection
 - 19.7 fb⁻¹ collected @ 8 TeV
 - Trigger on high p_T isolated leptons (Z -> I⁺I⁻) and photons (~17 GeV)
 - Offline
 - m(l⁺l⁻) compatible with a Z boson (71 GeV< m(l⁺l⁻) <111 GeV)
 - photons in barrel ($|\eta| < 1.4$)
 - PF-jets, anti-k_T clustering R=0.5, p_T>30 GeV
 - Z+jets and γ+jets analyzed separately
 - Comparison with MADGRAPH+PHYTHIA 6(LO+PS+K_{FEWZ}) and BLACKHAT
 23/08/2015



💬 Jet Reconstruction @ CMS



defined, and tends to cluster around the hard energy deposits.

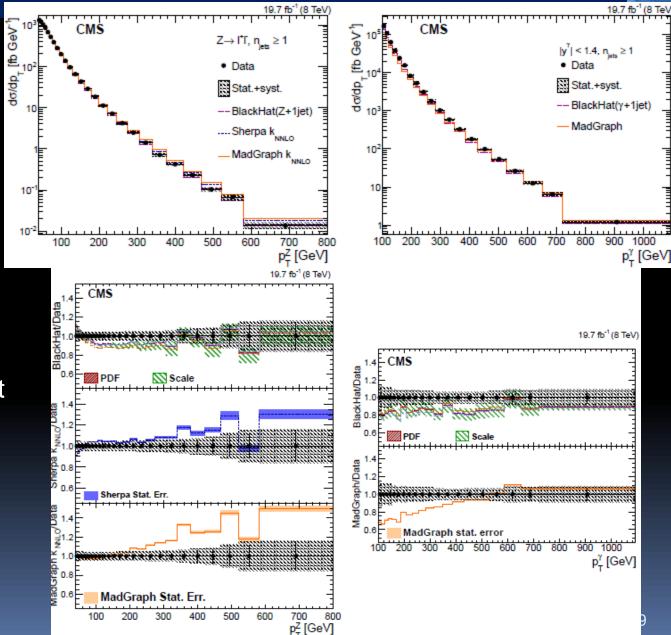
\bigcirc Z+jets (γ + jets Measurements



• Z/γ + \geq 1 Jet

BLACKHAT

- Mostly flat
- Underestimate at low p_T
- MADGRAPH
 - Increase w/ p_T
 - Overestimate at high p_T

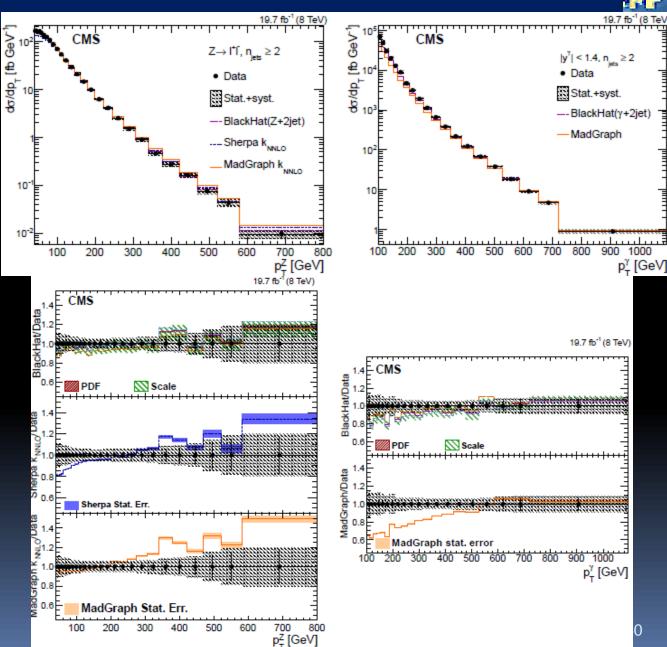


\bigcirc Z+jets $|\gamma + j$ ets Measurements





Same behavior

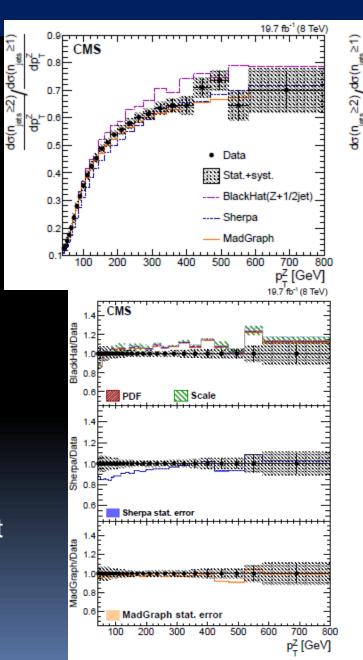


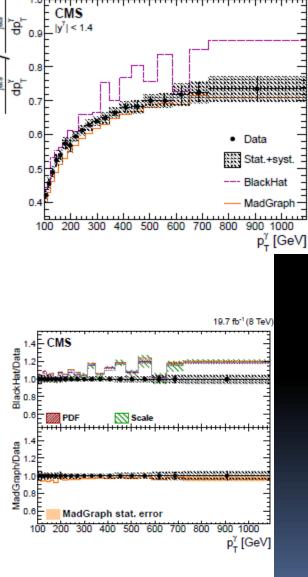
\bigcirc Z+jets $|\gamma + j$ ets Ratio Measurements



Z/γ + ≥2 Jets /
 Z/γ + ≥1 Jets

- BLACKHAT
 - Overestimate at high p_T
- MADGRAPHGood agreement





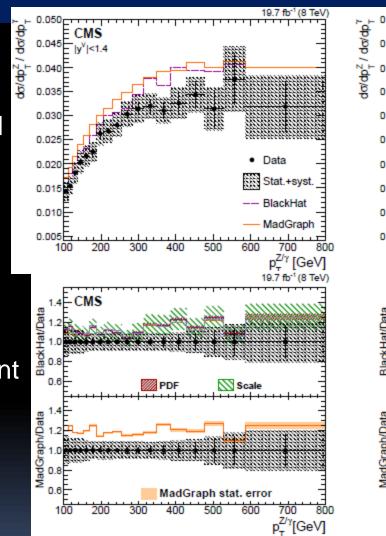
\bigcirc Z+jets $|\gamma + j$ ets Ratio of cross sections

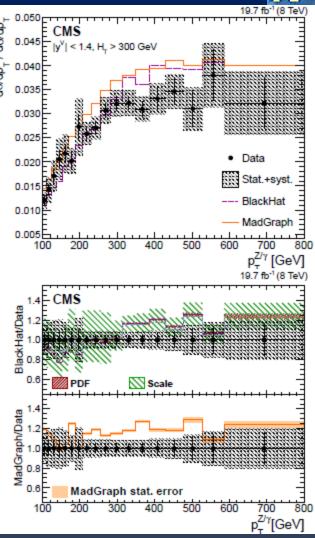


- H_T scalar p_T sum of all jets
- H_T > 300 GeV used in BSM analysis

- Both MC
 - Overestimate
 - Shape agreement

Ratio @ plateau





3/08/2015

 $R_{\rm dilep} = \frac{\sigma_{\rm Z \to \ell^+ \ell^-}(p_{\rm T}^{\rm Z} > 314\,{\rm GeV})}{\sigma_{\gamma}(p_{\rm T}^{\gamma} > 314\,{\rm GeV})} = 0.0322 \pm 0.0008\,{\rm (stat)} \pm 0.0020\,{\rm (syst)}.$





Drell - Yan in association with jets

CMS PAS FSQ-13-003

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CM Drell-Yan in association with jets

- DY lepton pair production
 - Well defined in SM (NNLO in pQCD)
 - At small p_T requires soft gluon resummation in all orders
 - DY in association w/ jets shift the pT spectrum
 - Allows resummation effects calculation in perturbative regime
- DY + jets at high mass (~125 GeV)
 - Important for Higgs production comparison
- Data sample and selection
 - 4.9 fb⁻¹ collected @ 7 TeV
 - Trigger on 2 high p_T isolated muons (p_T^{lead} > 20 GeV, p_T^{sublead} >10 GeV)
 - Offline
 - B 30 GeV < m(µ⁺µ⁻) < 1500 GeV</p>
 - 2 muons w/ |η| < 2.1; Jets w/ |η| < 4.5
 - PF-jets, anti-k_T clustering R=0.5, p_T>30 GeV
 - Comparison with MADGRAPH+PHYTHIA 6(LO+PS) 23/08/2015

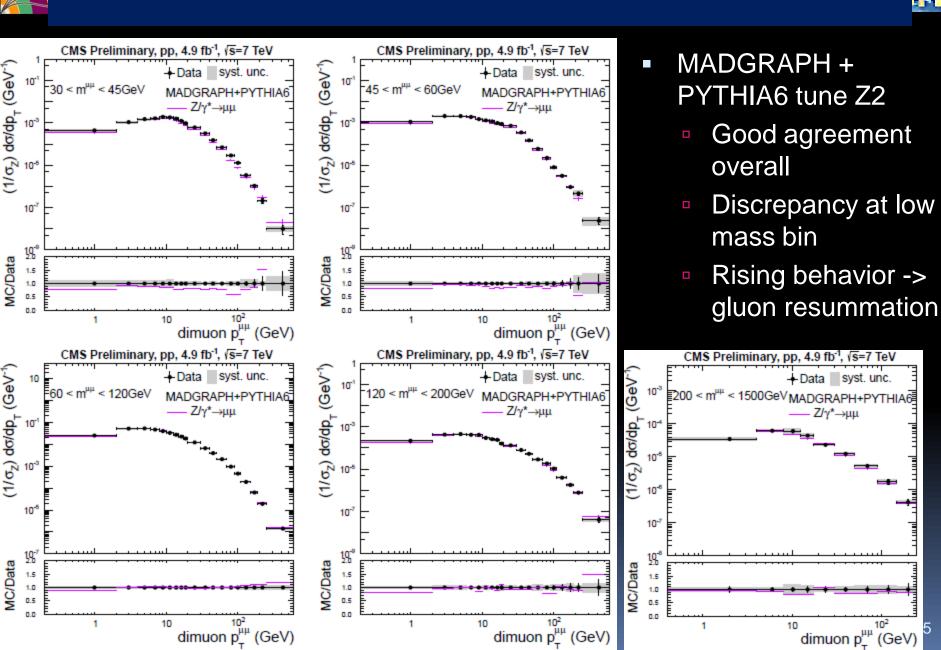
^{CM} Inclusive Drell-Yan -> µ+µ-



Z/γ*→μμ

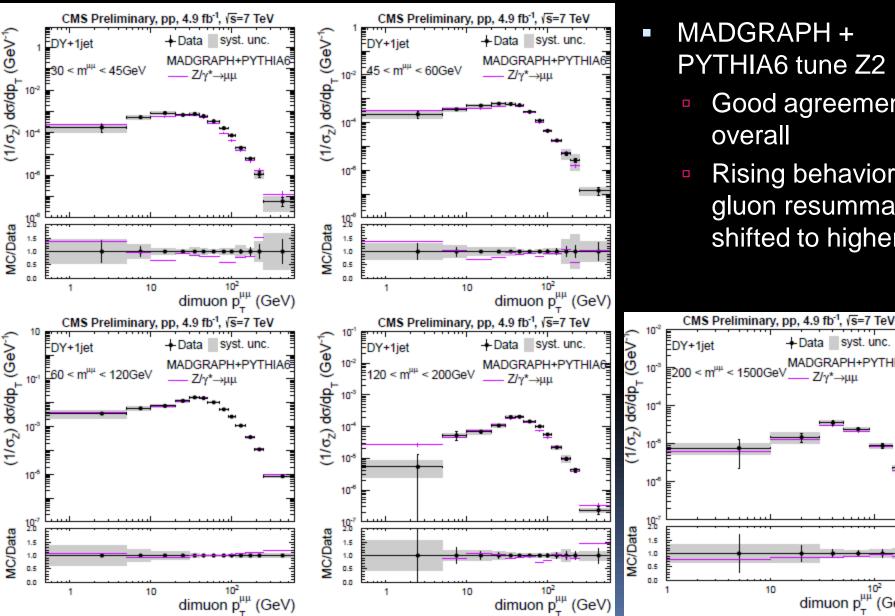
10²

dimuon p₊^{µµ} (GeV)



^{CM} Drell-Yan -> µ⁺µ⁻+ 1 Jet





- MADGRAPH + PYTHIA6 tune Z2
 - Good agreement overall
 - Rising behavior -> gluon resummation shifted to higher p_T

+ Data syst. unc.

Z/γ*→uu

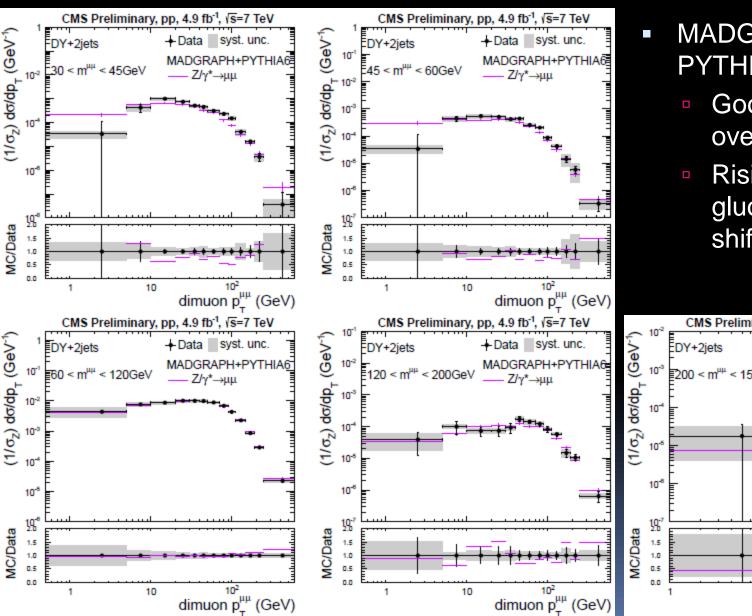
MADGRAPH+PYTHIA6

 10^{2}

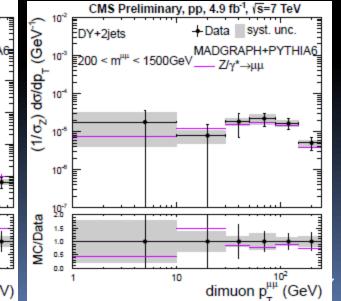
dimuon p_T^{µµ} (GeV) 6

\sim Drell-Yan -> $\mu^+\mu^-$ + 2 Jets





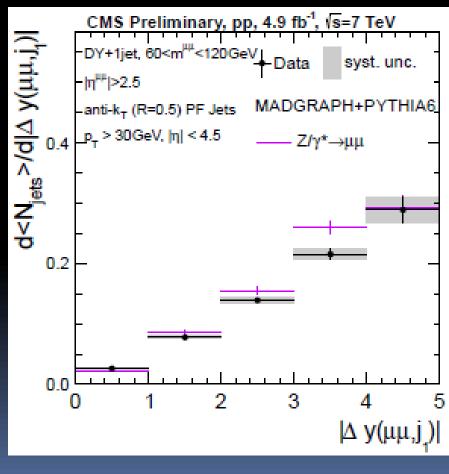
- MADGRAPH + PYTHIA6 tune Z2
 - Good agreement overall
 - Rising behavior -> gluon resummation shifted to higher p_T





Jet Multiplicity in Drell-Yan +Jets

- Average number of jets in forward Drell-Yan production as function of rapidity separation
- Δy(µµ,j₁) DY dimuon and leading jet separation
 - increase in jet multiplicity in data is slower than predicted from the simulation



Conclusions

- CMS measured diffractive and EM processes at the LHC
- Exclusive processes
 - Standard candle $\gamma\gamma \rightarrow l^+l^-$ used to correct for proton dissociation
 - Evidence for $\gamma\gamma \rightarrow WW \rightarrow 13$ candidates \rightarrow agreement with SM

$$\sigma (pp \to p^{(*)}W^+W^- p^{(*)} \to p^{(*)}\mu^{\pm}e^{\mp} p^{(*)}) = 12.3^{+5.5}_{-4.4} \text{ fb}$$

- AQGC limits two orders of magnitude more stringent than LEP and Tevatron
- Diffractive cross sections measured at 7 TeV
 - $\Box \ \sigma_{_{vis}}^{_{SD}} = 4.27 \pm 0.04(stat.) + 0.65/-0.58(syst.) \text{ mb for } -5.5 < \log \xi < -2.5$
 - $\Box \ \sigma_{_{vis}}^{^{DD}} = 0.93 \pm 0.01(stat.) + 0.26/-0.22(syst.) \text{ mb for } \Delta\eta > 3, M_X > 10 \text{ GeV}, M_y > 10 \text{ GeV}$
 - Good agreement with ATLAS on LRG cross section
- $Z+jets | \gamma + jets$ Measurements
 - Ratio overestimated by MC

$$R_{\rm dilep} = \frac{\sigma_{\rm Z \to \ell^+ \ell^-}(p_{\rm T}^{\rm Z} > 314\,{\rm GeV})}{\sigma_{\gamma}(p_{\rm T}^{\gamma} > 314\,{\rm GeV})} = 0.0322 \pm 0.0008\,{\rm (stat)} \pm 0.0020\,{\rm (syst)}.$$

- Drell-Yan in association with jets
 - increase in jet multiplicity in data is slower than predicted from the simulation
- More results coming soon

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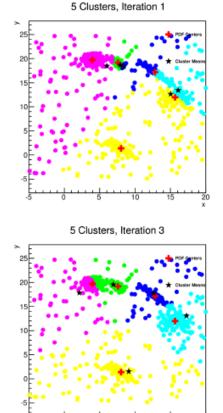


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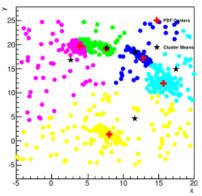
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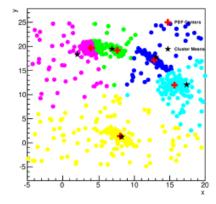








5 Clusters, Iteration 4



- Collinear- and Infrared-Safe
 - ★ collinear splitting shouldn't change jets
 - \star soft emissions shouldn't change jets

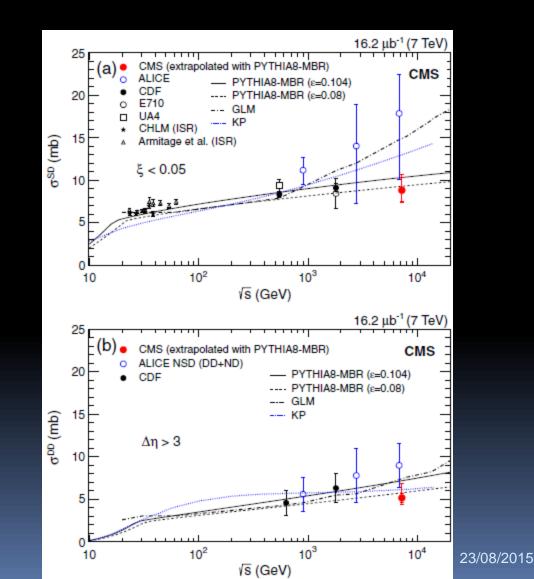
Collinear-Safety

Infrared-Safety



Measurements: SD & DD cross sections





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Luminosity and Pile-up



- The integrated luminosity (L) is based on the Van der Meer scans
- The uncertainty of the luminosity is 4%: dominates the systematic uncertainties of this analysis
- Number of collisions per bunch crossing follows Poisson
 Average λ (pile-up)

$$F_{\text{pileup}} = \frac{\sum_{i=1}^{\infty} iP(i,\lambda)}{\sum_{i=1}^{\infty} (1 - (1 - \epsilon_{\text{inel}})^i)P(i,\lambda)} \cdot \epsilon_{\text{inel}} = \frac{\epsilon_{\text{inel}}\lambda}{\sum_{i=1}^{\infty} (1 - (1 - \epsilon_{\text{inel}})^i)P(i,\lambda)} = 1 + \frac{1}{2}\lambda\epsilon_{\text{inel}} + \frac{1}{12}\lambda^2\epsilon_{\text{inel}}^2 + \mathcal{O}(\lambda^3)$$

 Correction factor – accounts for multiple collisions being counted as one.

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HC as a small x machine



