CMB constraints on (evaporating) primordial black holes

An application of ExoCLASS

Patrick Stöcker^a

in collaboration with: Vivian Poulin^b, Michael Krämer^a, Julien Lesgourgues^a

> ^aInstitute for Theoretical Particle Physics and Cosmology RWTH Aachen University

> > ^bDepartment of Physics and Astronomy Johns Hopkins University, Baltimore

> > > December 14, 2017

Outline

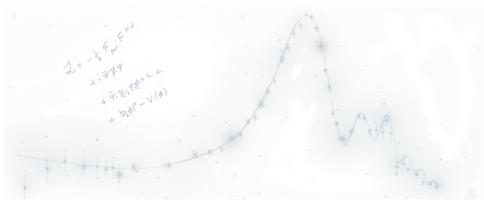
1 How does energy injected by dark matter influence the CMB? **Recombination in a nutshell** How to parametrize the model-specific influence on recombination

2 Constraining evaporating primordial black holes with the CMB An alternative to the "DM is a particle"-paradigm The influence of evaporating PBHs on recombination and the CMB Results: Can DM be made up of evaporating primordial black holes?

Effect of DM induced energy injection on the early IGM

Without loss of generality consider a model of DM χ which annihilates into the standard model χχ → SM SM, most of these SM-particles decay and additional radiation is induced.

$$\chi\chi \rightarrow \text{SM SM} \rightarrow e^+, e^-, \gamma, \nu, p$$





Effect of DM induced energy injection on the early IGM

Without loss of generality consider a model of DM χ which annihilates into the standard model χχ → SM SM, most of these SM-particles decay and additional radiation is induced.

$$\chi\chi
ightarrow {
m SM} \ {
m SM} \
ightarrow e^+, e^-, \gamma,
u, p$$

Evolution of the ionization fraction x_e and the matter temperature T_M

Evolution with/without DM-induced energy injection:

$$\frac{\mathrm{d}x_e(z)}{\mathrm{d}z} = \frac{1}{(1+z)H(z)}\left(R(z) - I(z) - I(z)_X\right)$$

and

$$\frac{\mathrm{d}T_M}{\mathrm{d}z} = \frac{1}{1+z} \left[2 \cdot T_M + \gamma(x_e) \left(T_M - T_{\mathrm{CMB}} \right) + \mathbf{K}_h \right]$$

All additional contributions, ionization $I(z)_X$ and heating K_h , depend on the rate of energy deposition $\frac{dE(z,x_e)}{dVdt}\Big|_{dep.}$

How does this rate relate to the rate of energy injection?

Relating the rate of energy deposition with the rate of energy injection

- ► The deposited energy rate $\frac{dE(z,x_e)}{dVdr}\Big|_{dep.}$ and injected energy rate $\frac{dE(z)}{dV dr}\Big|_{inj.}$ describe physics on two different energy scales. What happens in between?
- ► To infer the deposited energy rate (into a certain channel *c*) $\frac{dE(z,x_e)}{dV dt}\Big|_{dep. (ch.)}$, introduce the efficiency factor $f_c(z, x_e)$

$$\frac{\mathrm{d}E(z, x_e)}{\mathrm{d}V \,\mathrm{d}t} \bigg|_{\mathrm{dep. (ch.)}} = f_c(z, x_e) \cdot \frac{\mathrm{d}E(z)}{\mathrm{d}V \,\mathrm{d}t} \bigg|_{\mathrm{inj.}}$$



Relating the rate of energy deposition with the rate of energy injection

- ► The deposited energy rate $\frac{dE(z,x_e)}{dVdt}\Big|_{dep.}$ and injected energy rate $\frac{dE(z)}{dV dt}\Big|_{inj.}$ describe physics on two different energy scales. What happens in between?
- ► To infer the deposited energy rate (into a certain channel *c*) $\frac{dE(z,x_e)}{dV dt}\Big|_{dep. (ch.)}$, introduce the efficiency factor $f_c(z, x_e)$

$$\left. \frac{\mathrm{d}E(z, x_e)}{\mathrm{d}V \,\mathrm{d}t} \right|_{\mathrm{dep.}\,(\mathrm{ch.})} = f_c(z, x_e) \cdot \left. \frac{\mathrm{d}E(z)}{\mathrm{d}V \,\mathrm{d}t} \right|_{\mathrm{inj}}$$

► $f_c(z)$ can be expressed by the spectrum of injected particles $\frac{d\dot{N}(E,t(z))}{dE}$, through convolution with precomputed transfer functions [Slatyer '12 & '15, 1211.0283, 1506.03811, 1506.03812] ($\ell = e^+e^-, \gamma$):

$$f_{c}(z) = \frac{\int_{z}^{\infty} d\ln (1+z') \frac{(1+z')^{3}}{H(z')}}{\frac{z}{H(z)} \int_{0}^{m} dE \ T_{c}^{(\ell)}(z',z,E) E \ \frac{d\dot{N}(E,t(z))}{dE} \Big|_{\text{inj.}}^{(\ell)}}{\int_{0}^{\frac{(1+z)^{3}}{H(z)}} \int_{0}^{m} dE \ E \ \frac{d\dot{N}(E,t(z))}{dE} \Big|_{\text{inj.}}^{\text{tot.}}}$$

Relating the rate of energy deposition with the rate of energy injection

- ► To infer the deposited energy rate (into a certain channel *c*) $\frac{dE(z,x_e)}{dV dt}\Big|_{dep. (ch.)}$, introduce the efficiency factor $f_c(z, x_e)$

$$\left. \frac{\mathrm{d}E(z, x_e)}{\mathrm{d}V \,\mathrm{d}t} \right|_{\mathrm{dep.}\,(\mathrm{ch.})} = f_c(z, x_e) \cdot \left. \frac{\mathrm{d}E(z)}{\mathrm{d}V \,\mathrm{d}t} \right|_{\mathrm{inj.}}$$

► $f_c(z)$ can be expressed by the spectrum of injected particles $\frac{d\dot{N}(E,t(z))}{dE}$, through convolution with precomputed transfer functions [Slatyer '12 & '15, 1211.0283, 1506.03811, 1506.03812] ($\ell = e^+e^-, \gamma$):

$$f_{c}(z) = \frac{\int_{\ell}^{\infty} d\ln(1+z') \frac{(1+z')^{3}}{H(z')} \sum_{\ell} \int_{0}^{m} dE \ T_{c}^{(\ell)}(z',z,E) E \ \frac{d\dot{N}(E,t(z))}{dE} \Big|_{\text{inj.}}^{(\ell)}}{\frac{(1+z)^{3}}{H(z)} \int_{0}^{m} dE \ E \ \frac{d\dot{N}(E,t(z))}{dE} \Big|_{\text{inj.}}^{\text{tot.}}}$$

→Development of ExoCLASS, an extension of the CMB-code CLASS, to incorporate the effect on the CMB for every exotic scenario of energy injection. (Publication in preparation)

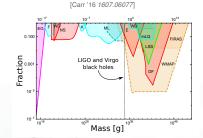
The fuzz about primordial black holes as (subdominant) dark matter

- PBHs are black holes produced through density fluctuations in the early universe.
- ► Broad range of possible masses. → Broad phenomenology.

(e.g. Constraints from CMB (accretion and evaporation), extragalactic $\gamma\text{-background},$ lensing and LSS)

 Possible origin of black hole-binary mergers measured by LIGO and VIRGO.

+ 12:31:124 - V (4)

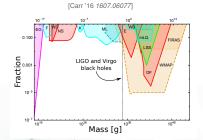


The fuzz about primordial black holes as (subdominant) dark matter

- PBHs are black holes produced through density fluctuations in the early universe.
- ► Broad range of possible masses. → Broad phenomenology.

(e.g. Constraints from CMB (accretion and evaporation), extragalactic $\gamma\text{-background},$ lensing and LSS)

 Possible origin of black hole-binary mergers measured by LIGO and VIRGO.



- ▶ Light PBHs ($M \sim 2 \times 10^{13} 10^{17}$ g) have a temperature ($T_{\rm BH} \sim 1/M_{\rm BH}$) high enough to evaporate and to sizeably inject energy into the IGM through Hawking-radiation.
- ► Injected spectrum (for particle species with spin *s*) follows thermal distribution:

$$rac{\mathrm{d}\dot{N}_s}{\mathrm{d}E}(E,z) \propto rac{E/T_{\mathrm{BH}}(z)}{\exp{(E/T_{\mathrm{BH}}(z))} - (-1)^{2s}}$$

Evaporation in turn reduces the mass of the black hole

$$\frac{\mathrm{d}M}{\mathrm{d}t} \propto -\mathcal{F}(M)M^{-2}$$

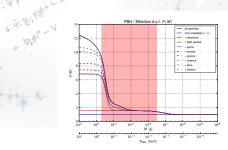
with $\mathcal{F}(M)$ being the effective degrees of freedom for a black hole with mass M.

The choice of $\mathcal{F}(M)$

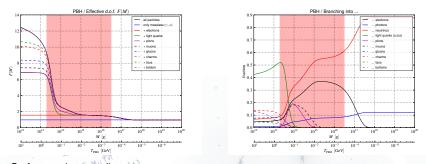
- ▶ In principle, take a sum of step-functions which kick in at $T_{BH} = m_i$, weighted with the d.o.f. of the particle.
- We follow the prescription given in [J. H. MacGibbon '91 (PRD 44, 376)] and include the QCD-phase transition.

$$\mathcal{F}(M) = \sum_{\text{part. } i} \prod_{i} \cdot f_{s,q} \cdot \exp\left(-\frac{M}{\beta_s \tilde{M}_i}\right) \cdot Q_i(T(M))$$

with Π_i being the internal degrees of freedom of the particle *i*, \tilde{M}_i being the mass of a black hole with a temperature $T_{BH} = m_i$ and $f_{s,q}$, β_s encoding the peak height and position of a black-body spectrum given the spin *s* and the charge *q* of the particle. (Takes the tail of the thermal distribution into account.)



CMB-constraints on evaporating primordial black holes I

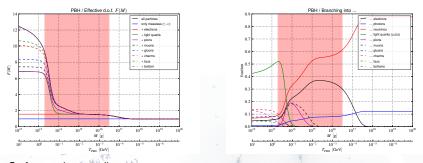


So far, previous studies [Poulin ,Lesgourgues ,Serpico '16, 1610.1005] ...

- ... did not include the variation of the PBH-mass for initial masses below $\sim 10^{15}$ g.
- ... only considered the primary electrons/positrons and photons. All other species are "inefficient".



CMB-constraints on evaporating primordial black holes I

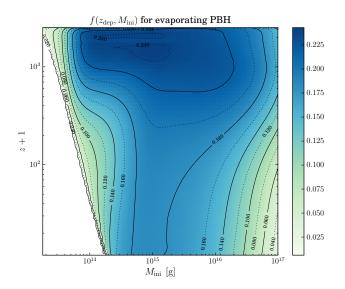


So far, previous studies [Poulin ,Lesgourgues ,Serpico '16, 1610.1005] ...

- ... did not include the variation of the PBH-mass for initial masses below $\sim 10^{15}$ g.
- ... only considered the primary electrons/positrons and photons. All other species are "inefficient".
 But in the mass range of our interest pions, muons and light quarks can have a non-negligible contribution to the energy injection and should be included
- The implementation in the ExoCLASS-package addresses both points above (Still excluding the secondary of the quarks).
 - →First precise analysis of CMB constraints on evaporating PBHs

CMB-constraints on evaporating primordial black holes II

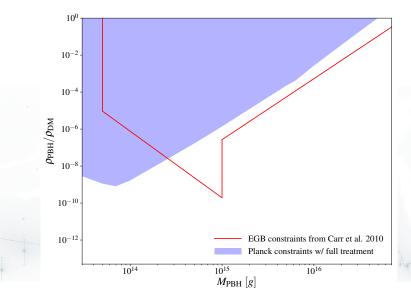
f(z) as function of the initial mass



8/11

CMB-constraints on evaporating primordial black holes III

Constraint on the allowed DM-fraction



/11

Conclusions

- Development of ExoCLASS, a package to calculate the influence of dark matter induced energy injection on the IGM during the "dark ages" for any generic model (for any given injection spectrum) and for any injection history (for any z-dependence of the energy rate)
- First detailed analyses on evaporating primordial black holes lead to the strongest results in the range of "short-lived" black holes ($M < 10^{15}$ g), where the PBH-mass significantly change with time and the black hole fully evaporates within the redshift-window the CMB is sensitive to.
- Lowest possible mass which has an effect on the CMB:

 $M_{
m PBH}>2 imes10^{13}~
m g$

 For "long-lived" black holes, constraints from extragalactic gamma-ray background are stronger by an order of magnitude.

Conclusions

- Development of ExoCLASS, a package to calculate the influence of dark matter induced energy injection on the IGM during the "dark ages" for any generic model (for any given injection spectrum) and for any injection history (for any z-dependence of the energy rate)
- First detailed analyses on evaporating primordial black holes lead to the strongest results in the range of "short-lived" black holes ($M < 10^{15}$ g), where the PBH-mass significantly change with time and the black hole fully evaporates within the redshift-window the CMB is sensitive to.
- Lowest possible mass which has an effect on the CMB:

 $M_{\rm PBH} > 2 \times 10^{13} \, {\rm g}$

- For "long-lived" black holes, constraints from extragalactic gamma-ray background are stronger by an order of magnitude.
- The ExoCLASS-package and the PBH-results will be published soon (Together with the validation of the package in the scope of the Higgs-Portal model)

Thank You for Your Attention

Any Open Questions?



2"





Recombination in a slightly bigger nutshell

Evolution of the ionization fraction x_e and the matter temperature T_M

- The easiest recombination-model assumes a three level atom (Ground state, excited state, continuum) [Peebles '68, Zeldovich et al. '69]
- Evolution with/without DM-induced energy injection:

$$\frac{\mathrm{d}x_e(z)}{\mathrm{d}z} = \frac{1}{(1+z)H(z)} \left(R(z) - I(z) - I(z)_X\right)$$

and

$$\frac{\mathrm{d}T_M}{\mathrm{d}z} = \frac{1}{1+z} \left[2 \cdot T_M + \gamma(x_e) \left(T_M - T_{\mathrm{CMB}} \right) + \mathbf{K}_h \right]$$

• The additional ionization term $I(z)_X = I_{X_i} + I_{X_{\alpha}}$ is given by:

$$I_{X_i} = \frac{1}{n_H(z)E_i} \left. \frac{\mathrm{d}E(z, x_e)}{\mathrm{d}V\mathrm{d}t} \right|_{\mathrm{dep. (direct ion.)}} \text{ and } I_{X_\alpha} = \frac{1-C}{n_H(z)E_\alpha} \left. \frac{\mathrm{d}E(z, x_e)}{\mathrm{d}V\mathrm{d}t} \right|_{\mathrm{dep. (excit. + ion.)}}$$

The contribution to the heating term K_h is:

$$K_{h} = \frac{2}{H(z)3k_{B}n_{H}(z)(1+f_{He}+x_{e})} \left.\frac{dE(z,x_{e})}{dVdt}\right|_{dep. (heat.)}$$

Under the bonnet of ExoCLASS

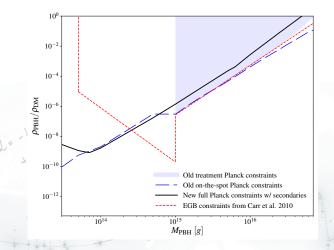
- Cosmological constraints on a given DM-model can in principle be derived in three steps
 - Calculate, given the parameters of the model, the injected spectrum of e[±], γ at given energy E (and at given redshift z').
 - 2. From this $f_c(z)$ is calculated by convolution of the spectra with the transfer function $T_c^{(\ell)}(z', z, E)$.
 - 3. Feed $f_c(z)$ into *CLASS* for the calculation of C_ℓ

In practice this looks like:

- The f(z)-backend of the ExoCLASS-package (written in Python) calculates f_c(z) from a given injection spectra (e.g. from MadGraph / PYTHIA) (steps 1 and 2)
 - Possibility to derive f(z) for different injection histories: 'Annihilation of particle DM', 'Decay of particle DM', 'Injection by black hole evaporation', ...
 - Spectra are automatically read and processed. (Automatic interpolation if spectra for different points in parameter space are given)
 - Inclusion of the cosmological background H_0, Ω_m, Ω_r into the convolution with the transfer functions
 - \rightarrow Can be also used as standalone-package.
- This backend is interfaced with the CMB-anisotropy solver CLASS. (step 3)
 - DM model-parameters are input parameters of (Exo)CLASS. (Exo)CLASS sets up the call to the backend and processes the output.

CMB-constraints on evaporating primordial black holes

Comparison to previous analyses



 Previous analysis overestimated the relative impact of electrons and photons to other (ineffective) particles. Hence the constraints were too strong.

Spectra from PBH evaporation (Primary emission vs. secondary emission)

