

Recent results in precision Higgs calculations

Giulia Zanderighi, CERN, University of Oxford & ERC

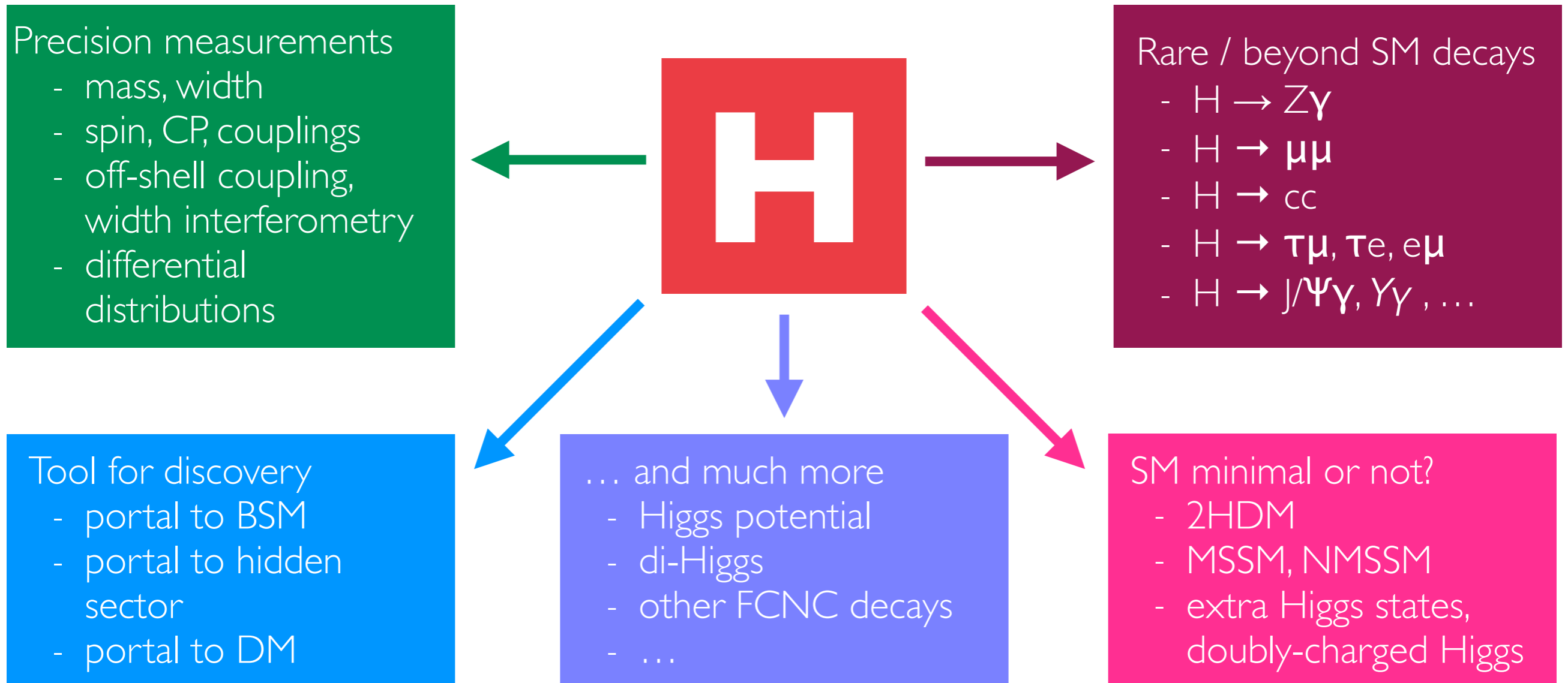


LHCP, Bologna June 2018

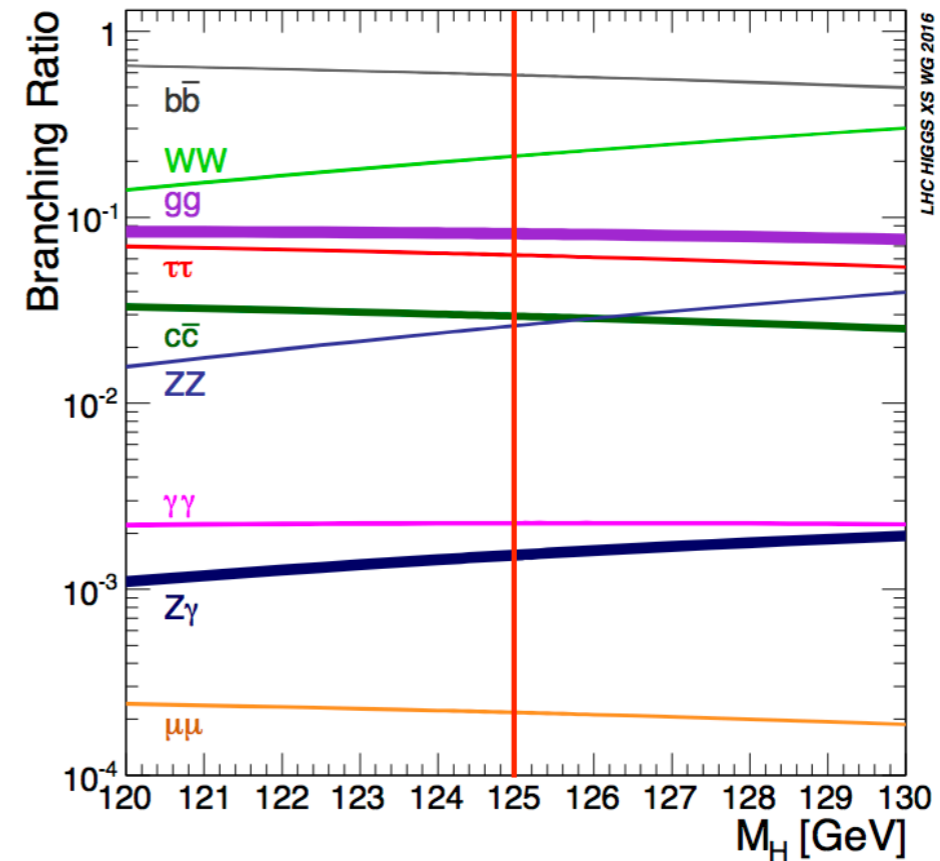
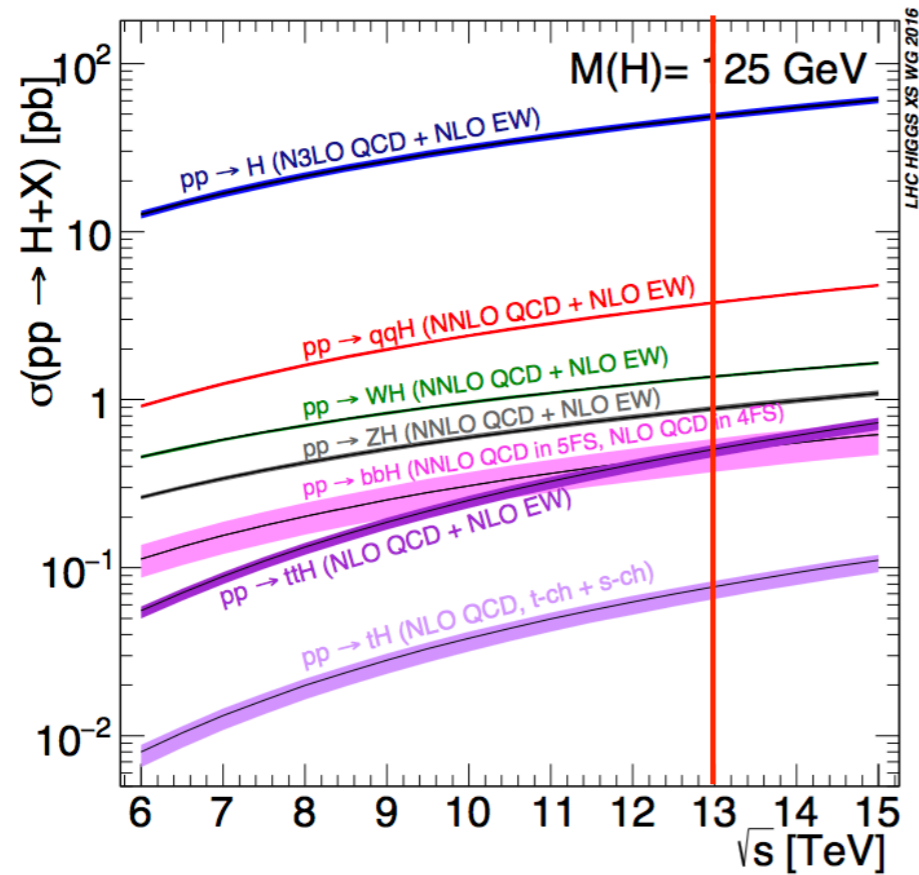
Higgs studies at the LHC

- The discovery of the Higgs boson at the LHC was a milestone in particle physics
- Higgs boson is the only fundamental scalar particle ever discovered. Its study at the LHC is new territory
- It is clear that this will be a long research program at the LHC [in comparison the b-quark was discovered forty years ago and, Belle II at SuperKEK, will now further study hadrons containing b-quarks]

An extremely rich program



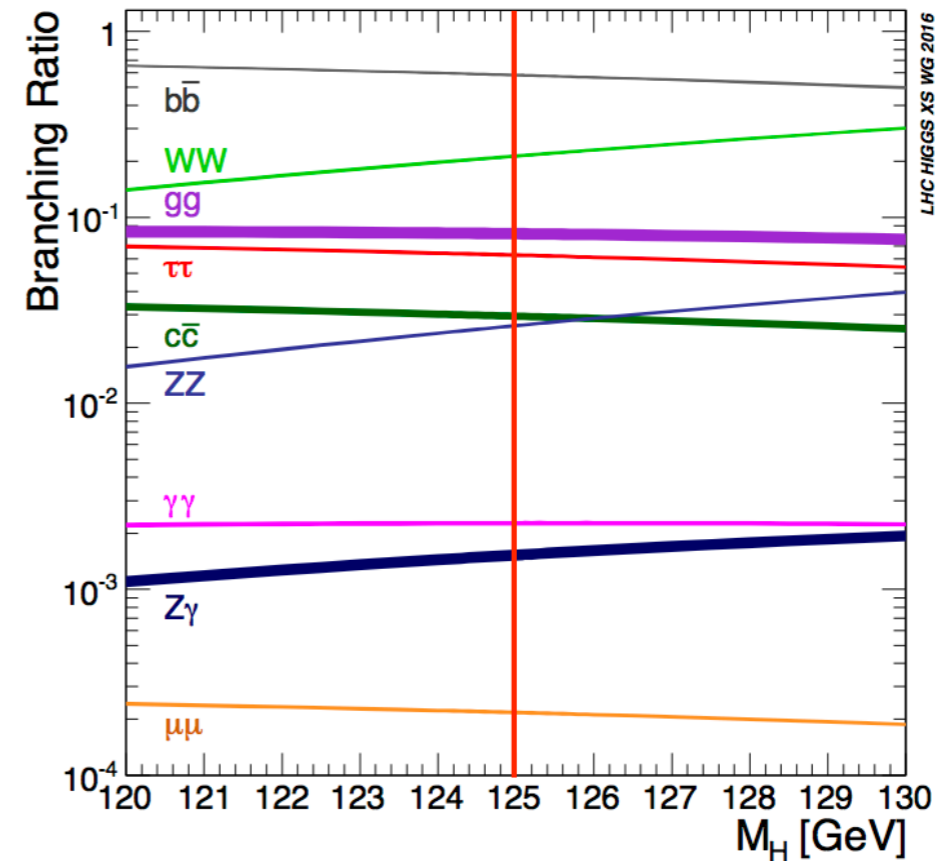
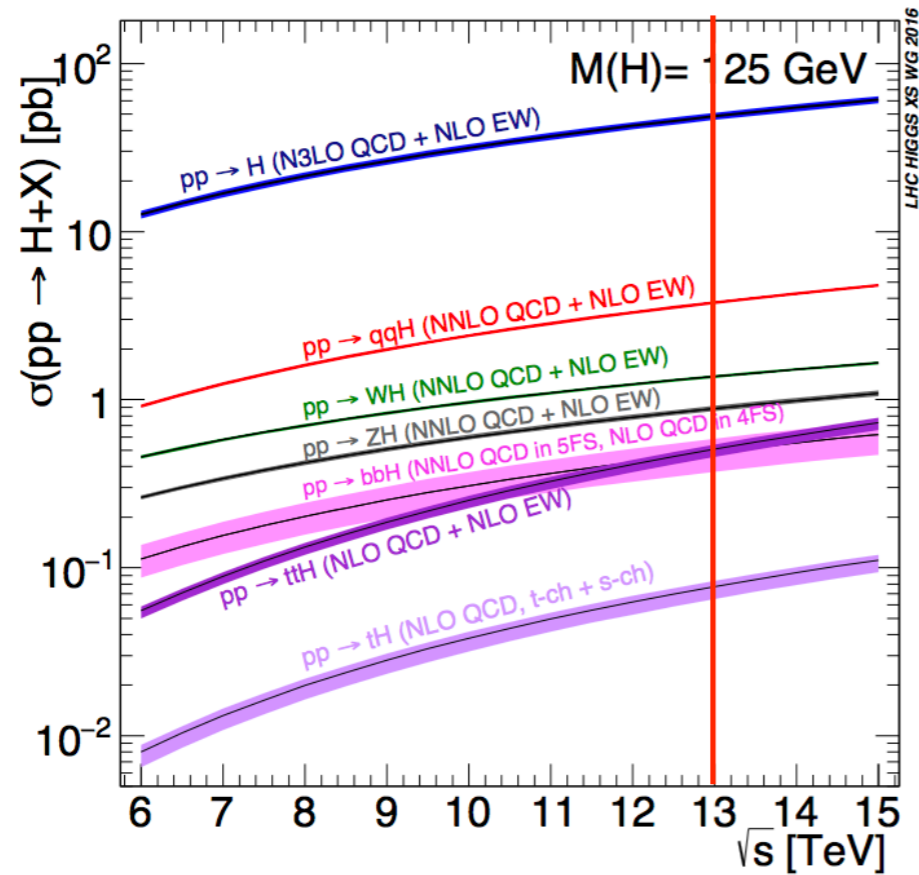
Production and decay



\sqrt{s} (TeV)	Production cross section (in pb) for $m_H = 125$ GeV					
	ggF	VBF	WH	ZH	$t\bar{t}H$	total
1.96	$0.95^{+17\%}_{-17\%}$	$0.065^{+8\%}_{-7\%}$	$0.13^{+8\%}_{-8\%}$	$0.079^{+8\%}_{-8\%}$	$0.004^{+10\%}_{-10\%}$	1.23
7	$16.9^{+5\%}_{-5\%}$	$1.24^{+2\%}_{-2\%}$	$0.58^{+3\%}_{-3\%}$	$0.34^{+4\%}_{-4\%}$	$0.09^{+8\%}_{-14\%}$	19.1
8	$21.4^{+5\%}_{-5\%}$	$1.60^{+2\%}_{-2\%}$	$0.70^{+3\%}_{-3\%}$	$0.42^{+5\%}_{-5\%}$	$0.13^{+8\%}_{-13\%}$	24.2
13	$48.6^{+5\%}_{-5\%}$	$3.78^{+2\%}_{-2\%}$	$1.37^{+2\%}_{-2\%}$	$0.88^{+5\%}_{-5\%}$	$0.50^{+9\%}_{-13\%}$	55.1
14	$54.7^{+5\%}_{-5\%}$	$4.28^{+2\%}_{-2\%}$	$1.51^{+2\%}_{-2\%}$	$0.99^{+5\%}_{-5\%}$	$0.60^{+9\%}_{-13\%}$	62.1

Decay channel	Branching ratio	Rel. uncertainty
$H \rightarrow \gamma\gamma$	2.27×10^{-3}	+5.0% -4.9%
$H \rightarrow ZZ$	2.62×10^{-2}	+4.3% -4.1%
$H \rightarrow W^+W^-$	2.14×10^{-1}	+4.3% -4.2%
$H \rightarrow \tau^+\tau^-$	6.27×10^{-2}	+5.7% -5.7%
$H \rightarrow b\bar{b}$	5.84×10^{-1}	+3.2% -3.3%
$H \rightarrow Z\gamma$	1.53×10^{-3}	+9.0% -8.9%
$H \rightarrow \mu^+\mu^-$	2.18×10^{-4}	+6.0% -5.9%

Production and decay

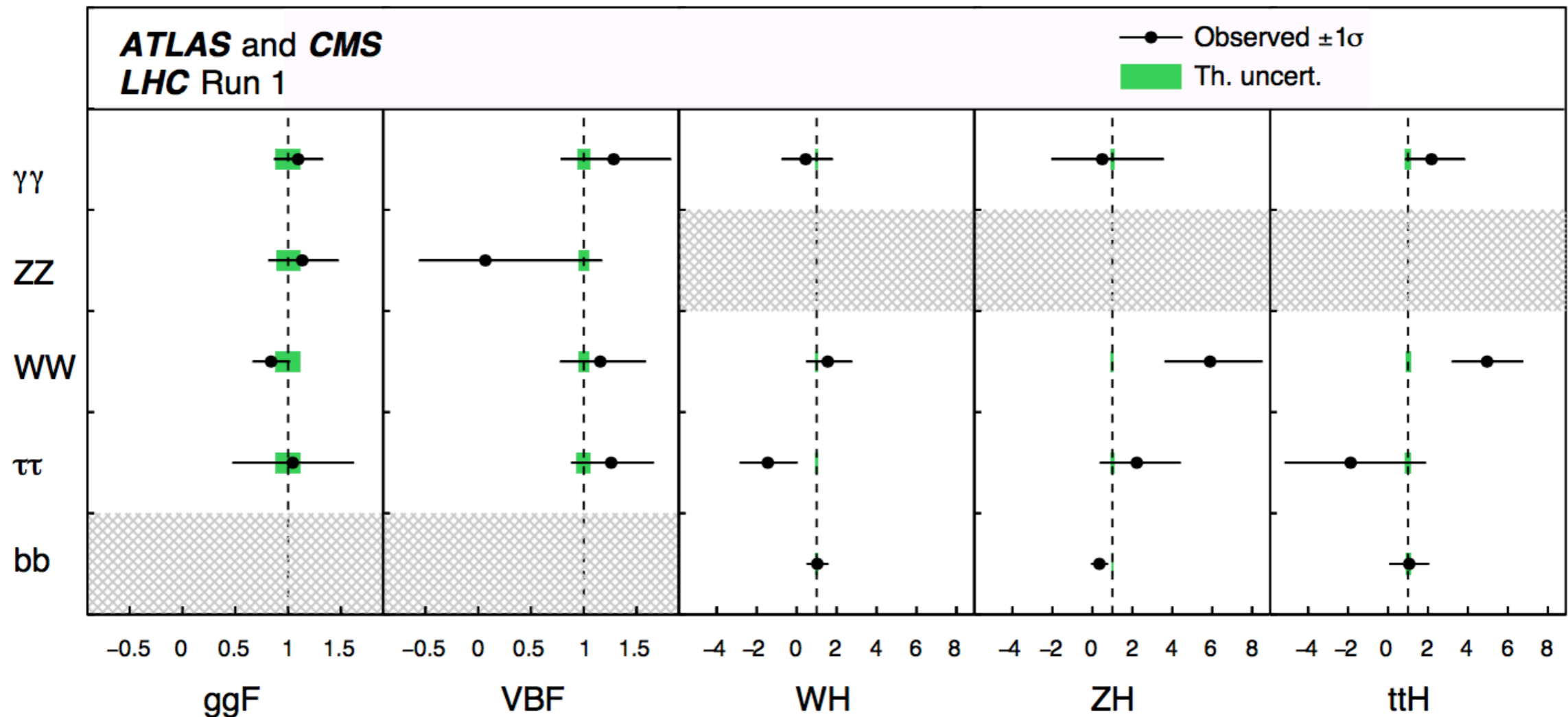


\sqrt{s} (TeV)	Production cross section (in pb) for $m_H = 125$ GeV					
	ggF	VBF	WH	ZH	$t\bar{t}H$	total
1.96	$0.95^{+17\%}_{-17\%}$	$0.065^{+8\%}_{-7\%}$	$0.13^{+8\%}_{-8\%}$	$0.079^{+8\%}_{-8\%}$	$0.004^{+10\%}_{-10\%}$	1.23
7	$16.9^{+5\%}_{-5\%}$	$1.24^{+2\%}_{-2\%}$	$0.58^{+3\%}_{-3\%}$	$0.34^{+4\%}_{-4\%}$	$0.09^{+8\%}_{-14\%}$	19.1
8	$21.4^{+5\%}_{-5\%}$	$1.60^{+2\%}_{-2\%}$	$0.70^{+3\%}_{-3\%}$	$0.42^{+5\%}_{-5\%}$	$0.13^{+8\%}_{-13\%}$	24.2
13	$48.6^{+5\%}_{-5\%}$	$3.78^{+2\%}_{-2\%}$	$1.37^{+2\%}_{-2\%}$	$0.88^{+5\%}_{-5\%}$	$0.50^{+9\%}_{-13\%}$	55.1
14	$54.7^{+5\%}_{-5\%}$	$4.28^{+2\%}_{-2\%}$	$1.51^{+2\%}_{-2\%}$	$0.99^{+5\%}_{-5\%}$	$0.60^{+9\%}_{-13\%}$	62.1

Decay channel	Branching ratio	Rel. uncertainty
$H \rightarrow \gamma\gamma$	2.27×10^{-3}	+5.0% -4.9%
$H \rightarrow ZZ$	2.62×10^{-2}	+4.3% -4.1%
$H \rightarrow W^+W^-$	2.14×10^{-1}	+4.3% -4.2%
$H \rightarrow \tau^+\tau^-$	6.27×10^{-2}	+5.7% -5.7%
$H \rightarrow \mu^+\mu^-$	2.18×10^{-4}	-5.9%

⇒ accordingly 90% of the talk on ggF!

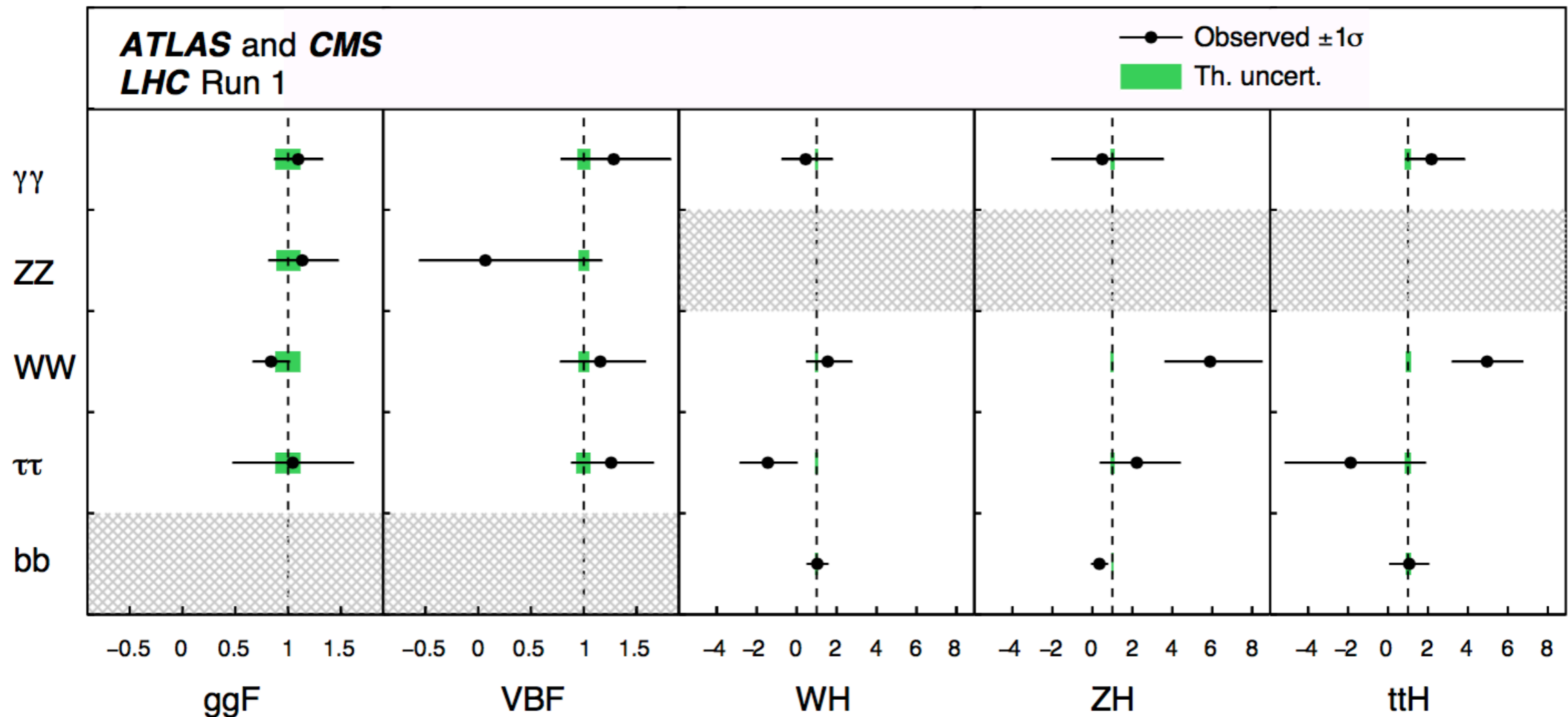
Theory vs data



Conclusion 1: a coherent picture emerges, with good consistency between data and theory

$$\mu = 1.09 \pm 0.07 \text{ (stat)} \pm 0.04 \text{ (expt)} \pm 0.03 \text{ (th. bkg)} \pm 0.07 \text{ (th. sig)}$$

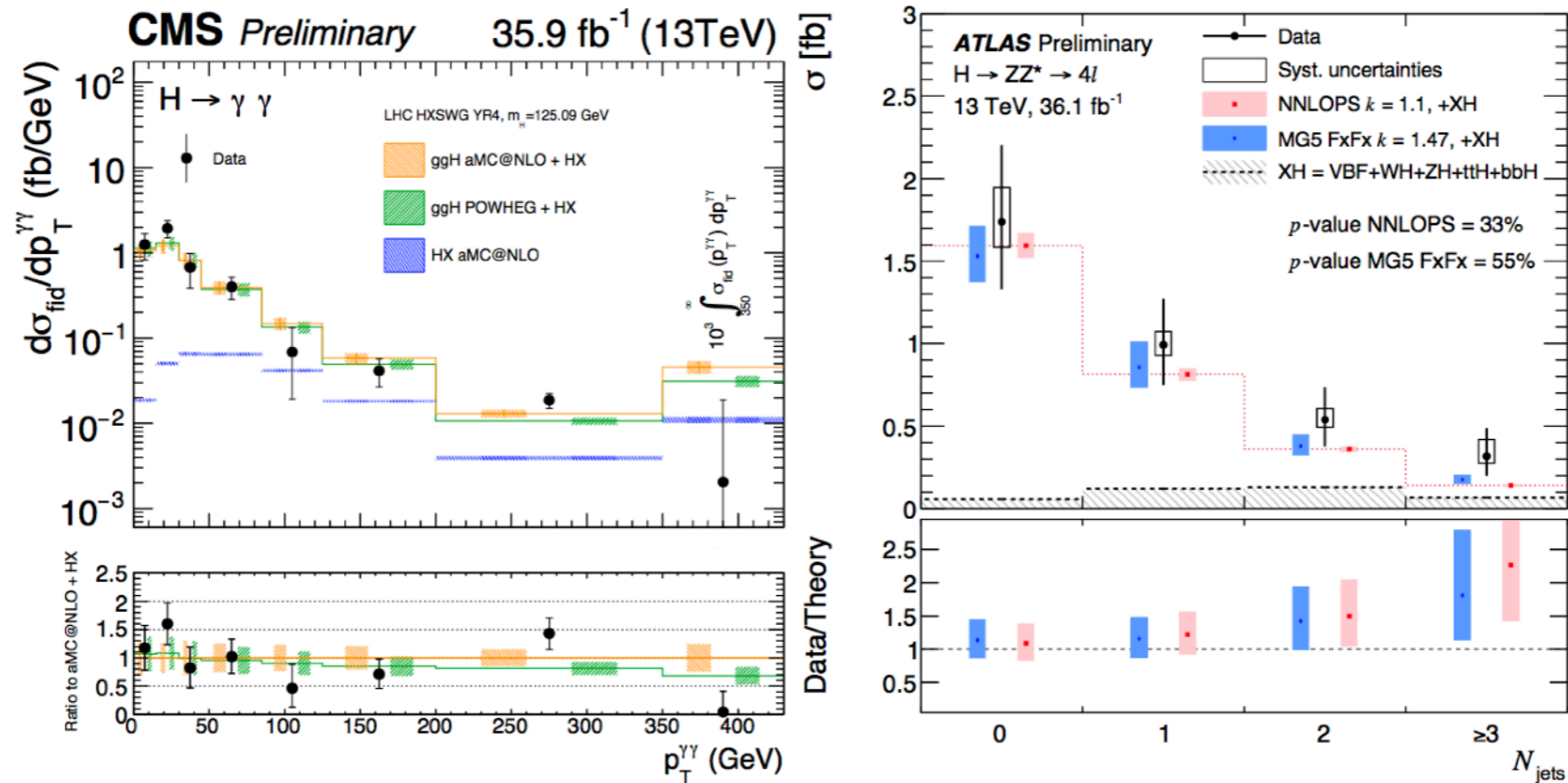
Theory vs data



Conclusion 2: theory is good enough ($\text{err}_{\text{TH}} \ll \text{err}_{\text{EXP.}}$), any further improvement does not bring in much ...

Why theory needs to improve

Take more exclusive measurements ($p_{t,H}$, jet-binned cross sections, ...):



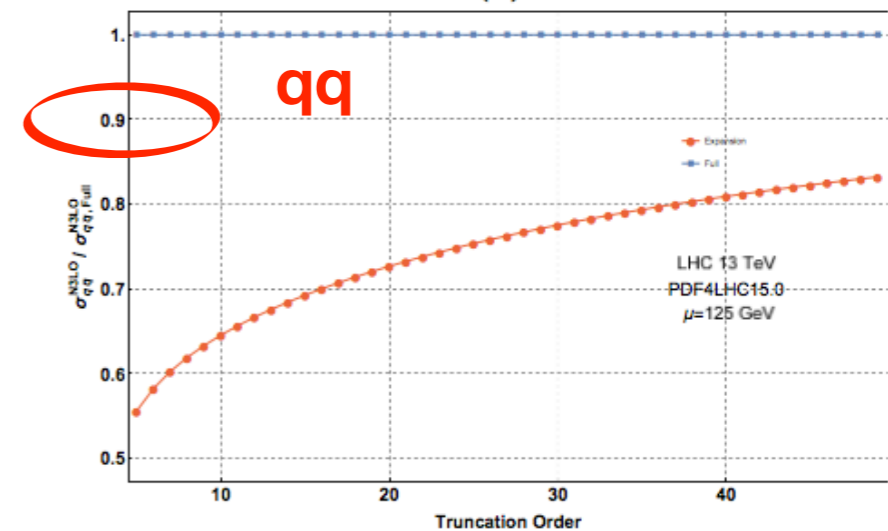
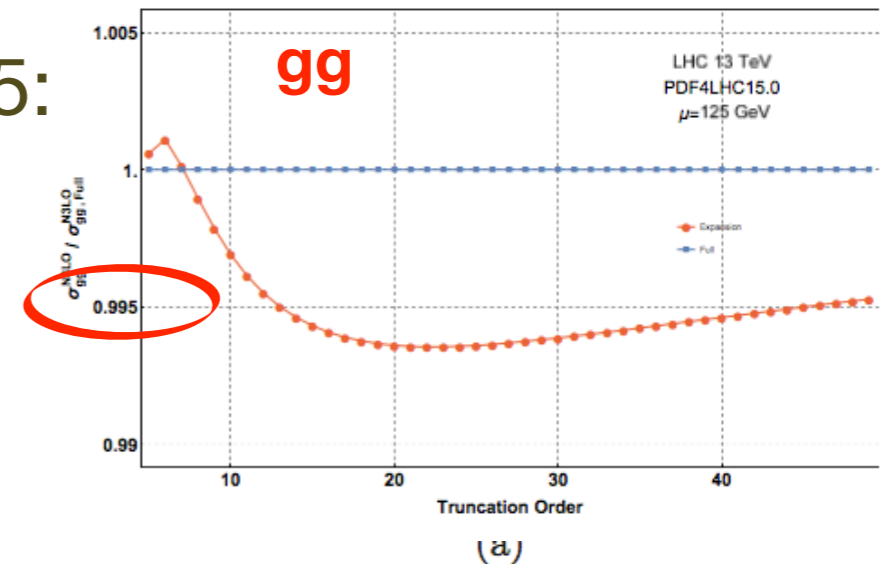
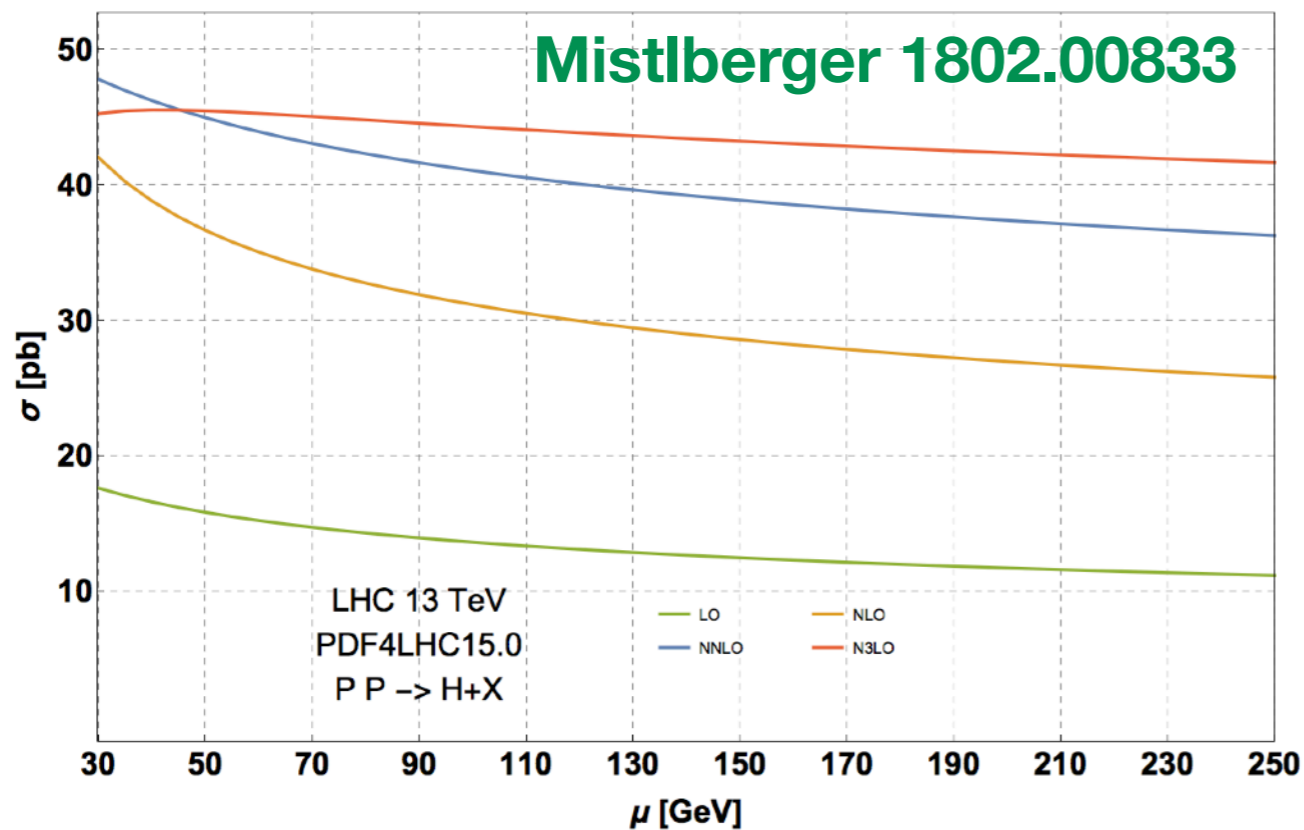
- err_{TH} already closer to err_{EXP}
- experimental error will be reduced considerably, systematic error will decrease too
- often total theory uncertainty was underestimated in the past

Since the Higgs mechanism was proposed in the 60s', the Higgs boson production and decay has been at the center of theorists attention. Accordingly, the literature on the subject is immensely vast. In this talk I will concentrate on presenting just

A few very recent (2018) highlights

N³LO Higgs

HEFT *beyond* the threshold expansion of 2015:



- threshold expansion works very well for gluon-initiated channels, less so for quark-initiated ones
- the estimate of the missing higher-order terms in the threshold approximation covers the difference to the exact result well (3% of the corrections)

NLO QCD to QCD-EW $gg \rightarrow H$

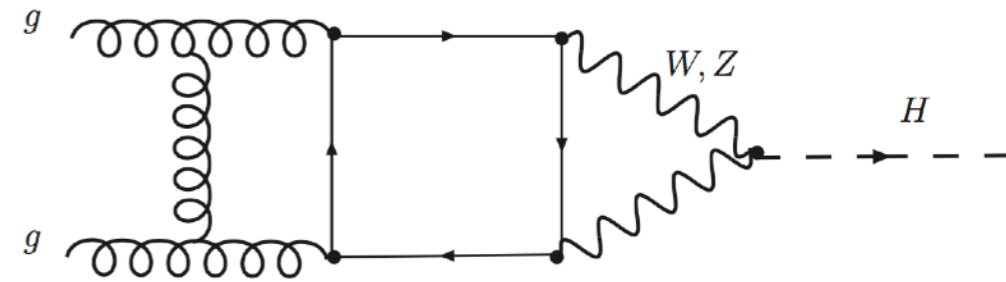
NLO QCD corrections to mixed QCD-EW contributions to gluon fusion production in the soft-gluon approximation

Bonetti, Melnikov, Tancredi 1801.10403

$$\begin{array}{ll} \sigma_{\text{QCD}}^{\text{LO}} = 20.6 \text{ pb}, & \sigma_{\text{QCD/EW}}^{\text{LO}} = 21.7 \text{ pb}, \\ \sigma_{\text{QCD}}^{\text{NLO}} = 32.66 \text{ pb}, & \end{array}$$

+5% (from LO to LO+EW)

+??% (from LO to NLO)



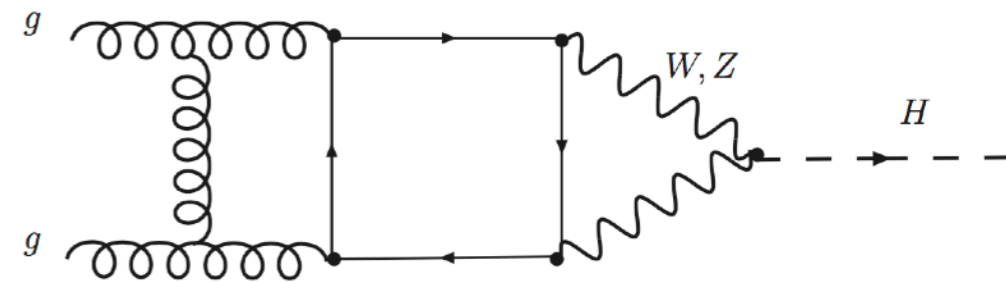
NLO QCD to QCD-EW $gg \rightarrow H$

NLO QCD corrections to mixed QCD-EW contributions to gluon fusion production in the soft-gluon approximation

Bonetti, Melnikov, Tancredi 1801.10403

$$\begin{array}{ll} \sigma_{\text{QCD}}^{\text{LO}} = 20.6 \text{ pb}, & \sigma_{\text{QCD/EW}}^{\text{LO}} = 21.7 \text{ pb}, \\ \sigma_{\text{QCD}}^{\text{NLO}} = 32.66 \text{ pb}, & \sigma_{\text{QCD/EW}}^{\text{NLO}} = 34.41 \text{ pb}. \end{array}$$

+5% (from LO to NLO for both)



Findings consistent with estimation in Anastasiou et al 0811.3458 (relied on unphysical approximation $M_W \gg M_H$)

Result further supports the factorisation of QCD and EW corrections

Further improvements require exact calculation of real corrections

Higgs p_t

Beyond inclusive cross-sections, **accurate predictions for differential distributions crucial for Run II**

- ➔ signal significance optimized by categorizing events according to kinematic properties (e.g. jet bins, Higgs p_t ...)
- ➔ a large fraction (30-40%) of Higgs events come with at least one jet
- ➔ kinematical distributions used to extract/constraint couplings and quantum numbers

The most basic distribution: transverse momentum of the Higgs boson

It is inclusive on radiation, not sensitive to definition of jets, relatively insensitive to hadronization effects

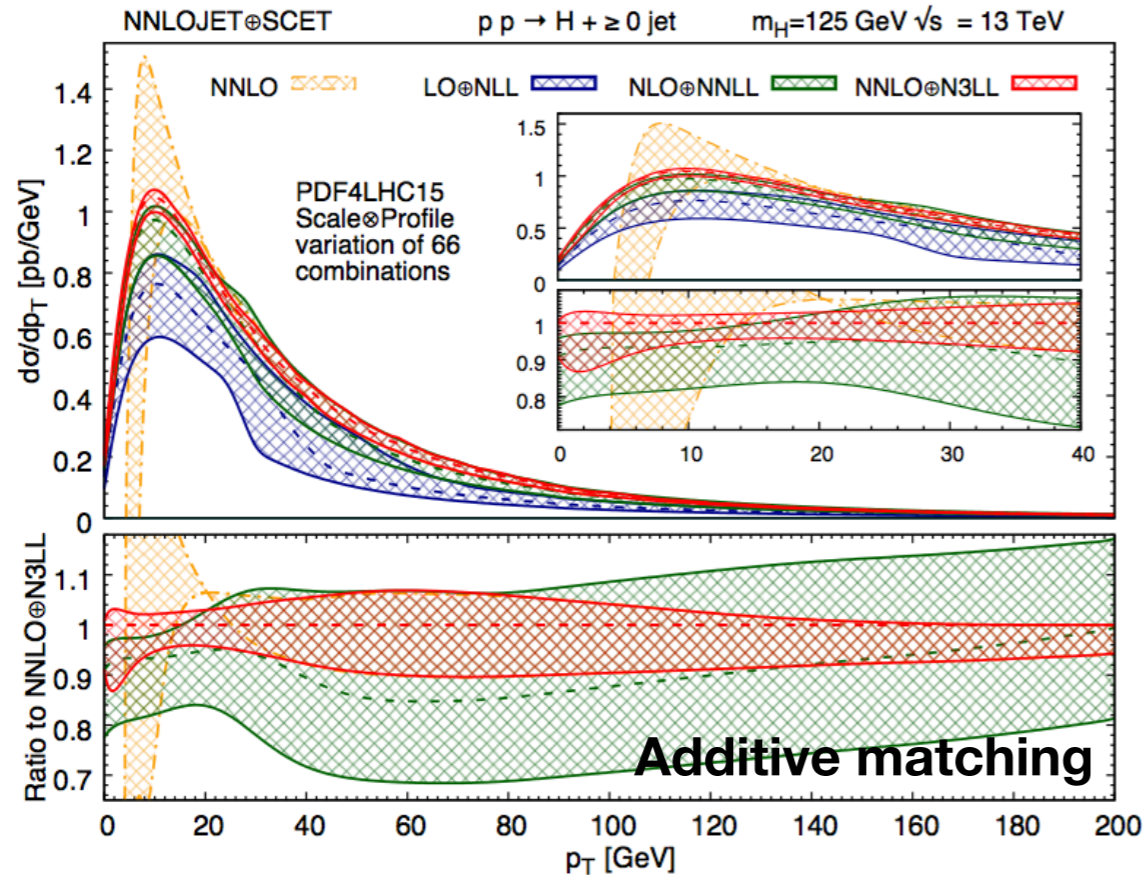
Why is Higgs p_t so difficult

Higgs p_t is shaped by many different and competing effects

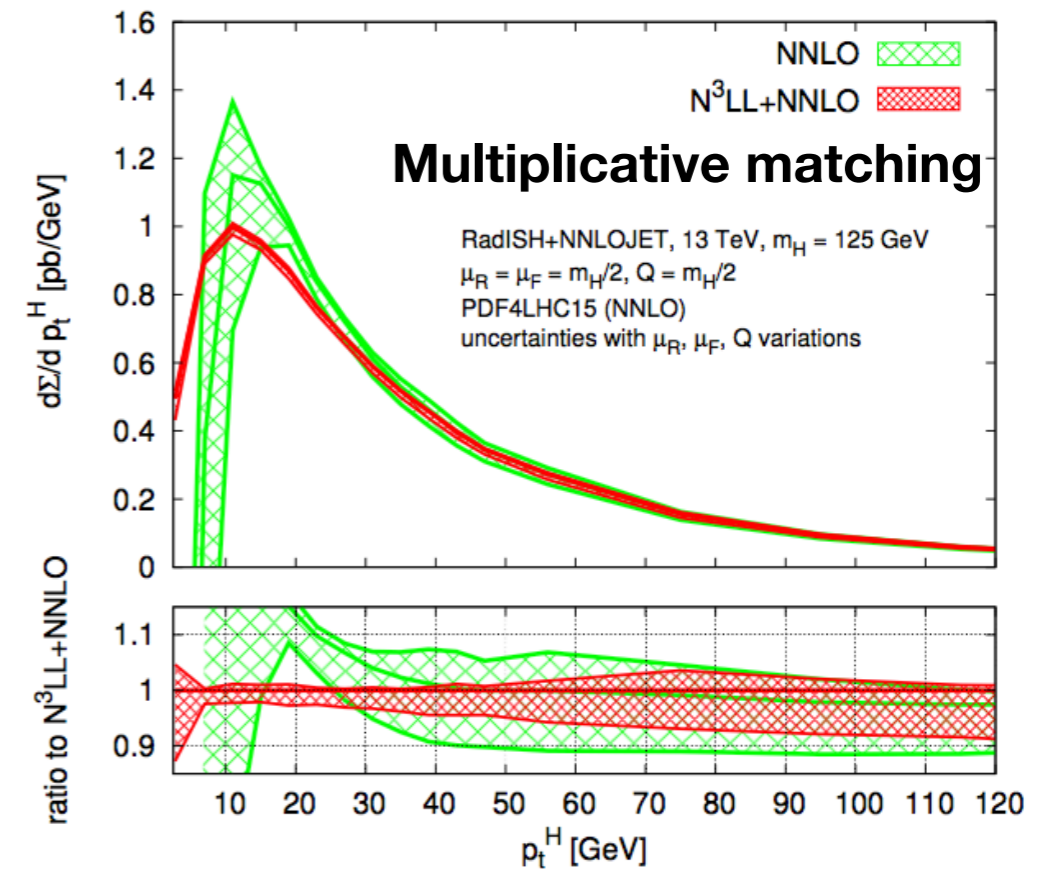
- ➔ dominant contribution due to **top-quark loops**. HEFT employed at small $p_{t,H}$, but top-loop must be resolved at high $p_{t,H}$ (*leading order description of high- $p_{t,H}$ involves already a one-loop 2 to 2 process*)
- ➔ non-negligible shape-distortions due to **interference effects** of amplitudes with top and bottom or charm quarks
- ➔ **new types of logarithms** when light quarks are present
- ➔ resummation of **soft logarithms** at small $p_{t,H}$ and combination of **resummation with bottom-quark effects**

Higgs at $N^3LL+NNLO$

Chen et al 1805.00736



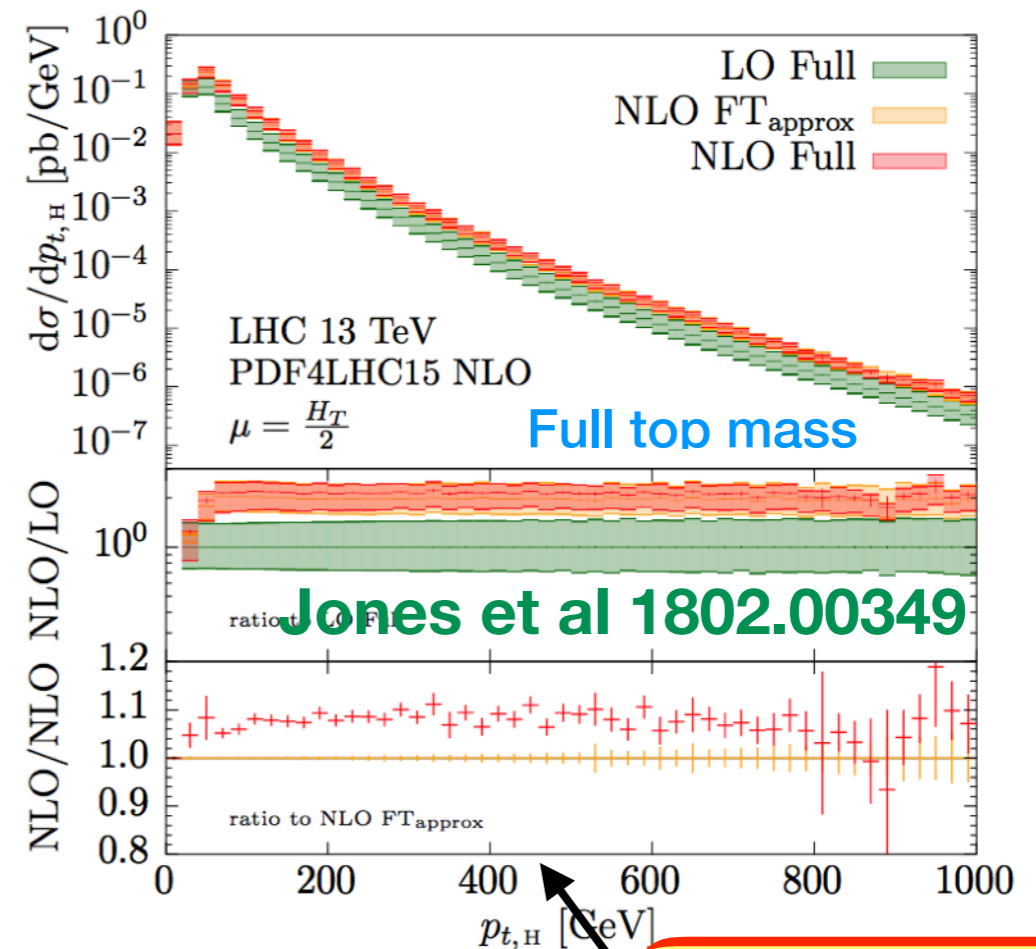
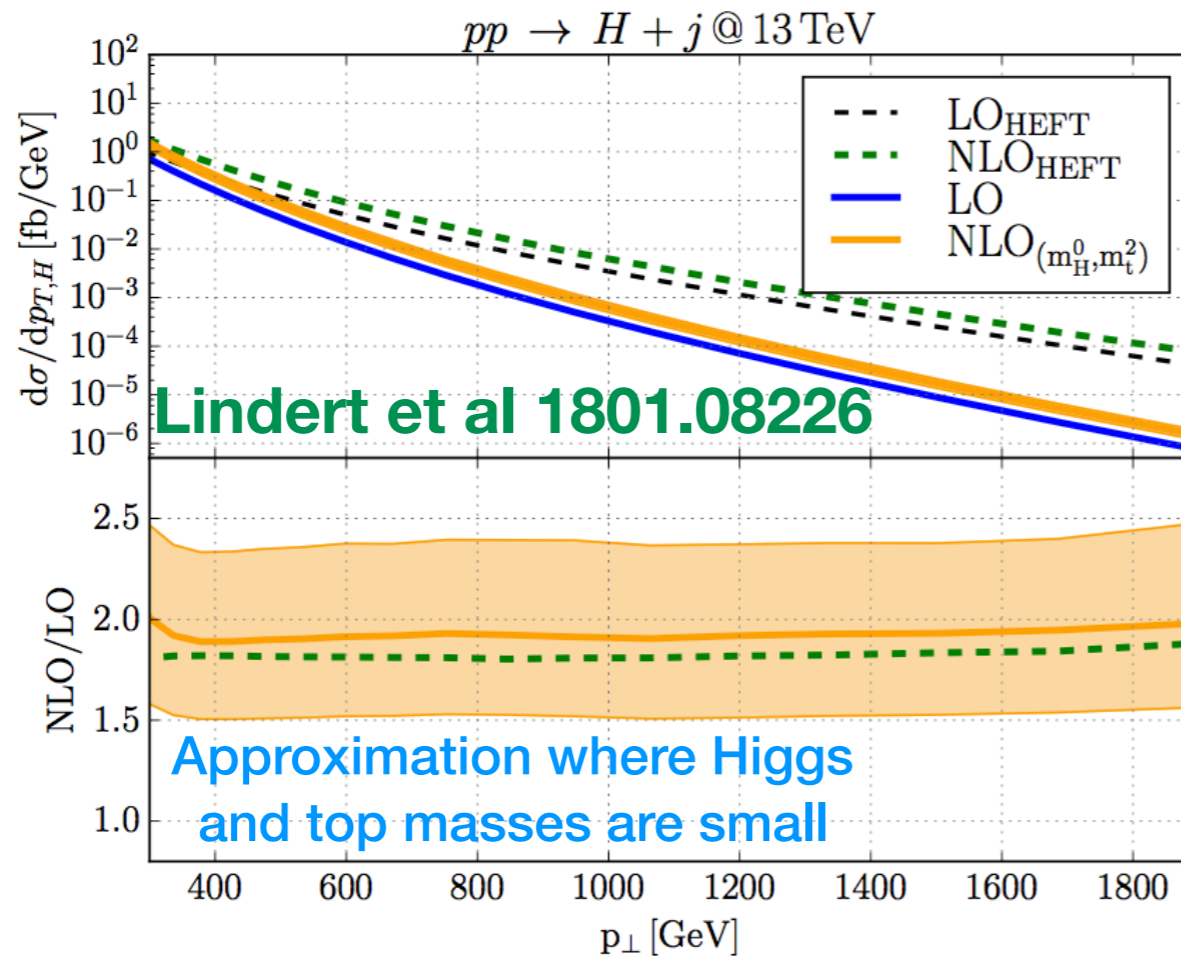
Bizon et al 1805.05916



- Calculations in HEFT
- Resummation important below about 40 GeV
- At this level of precision results with scales m_H or $m_H/2$ are very similar (but $m_H/2$ might suffer from unphysical cancelations)

NEW: includes decay to photons and fiducial cuts

Boosted Higgs: full NLO



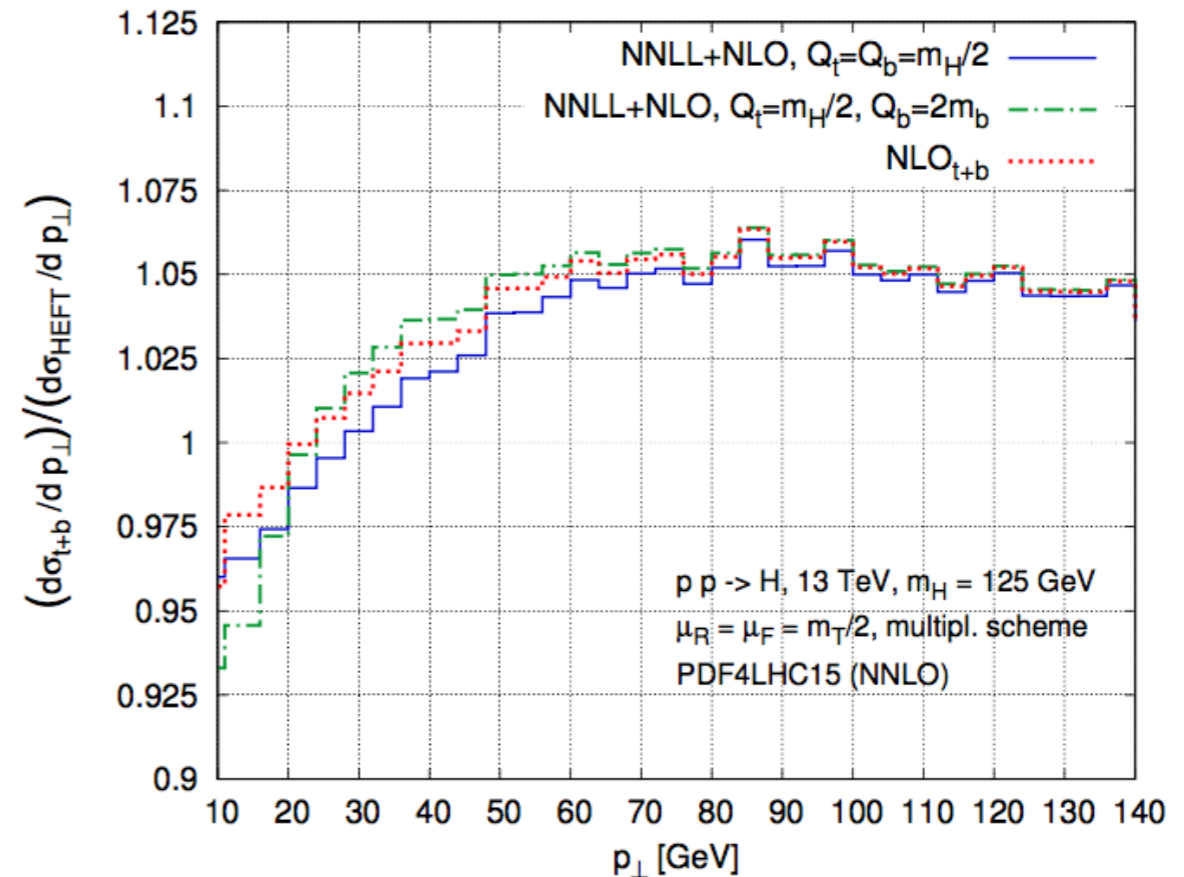
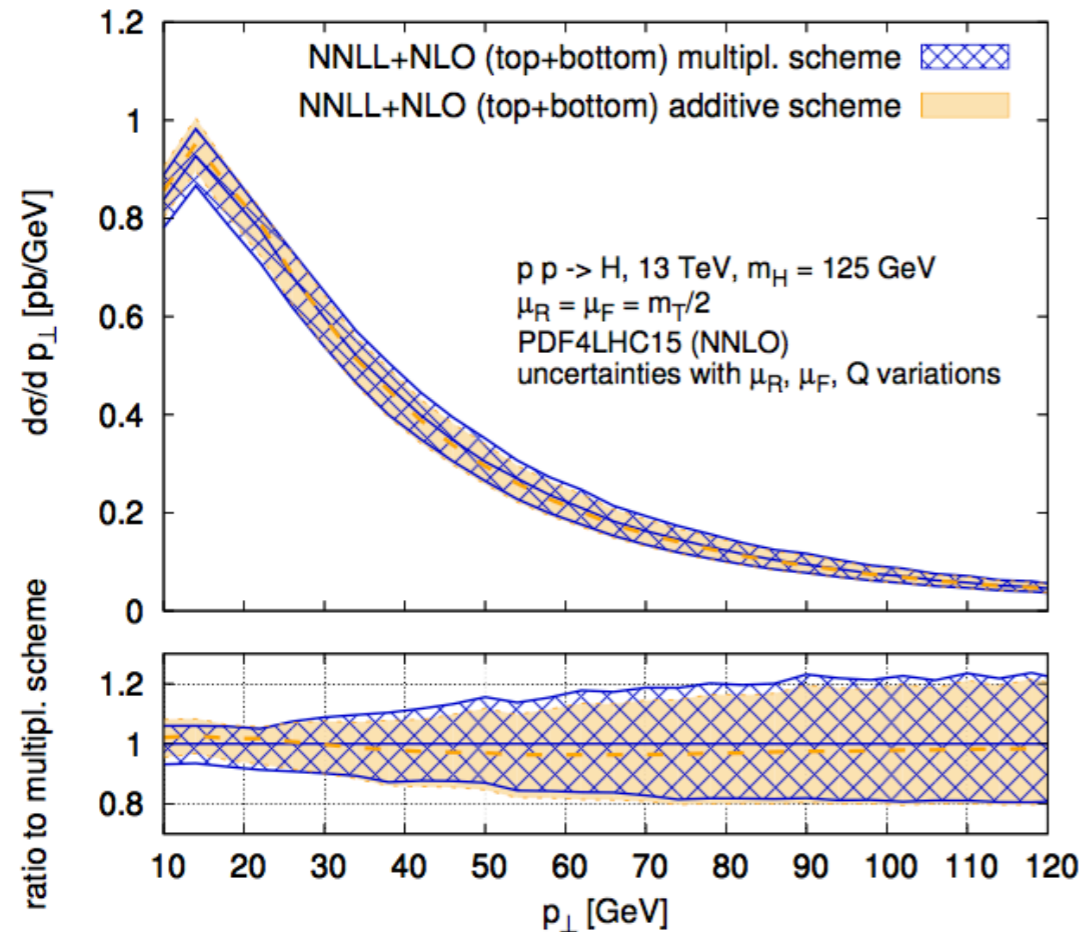
- NLO-approx OK in describing the shape (to within 5%)
- residual 10% normalisation effect
- **one of the largest sources of uncertainties in the description of the high-pt region finally removed.** Opens up the possibility to use high-pt region to search for BSM (e.g. disentangle point-like to top-induced ggH coupling)

Current CMS data up to here

Banfi, Martin, Sanz 1308.4771

Bottom mass effects (intermediate p_t)

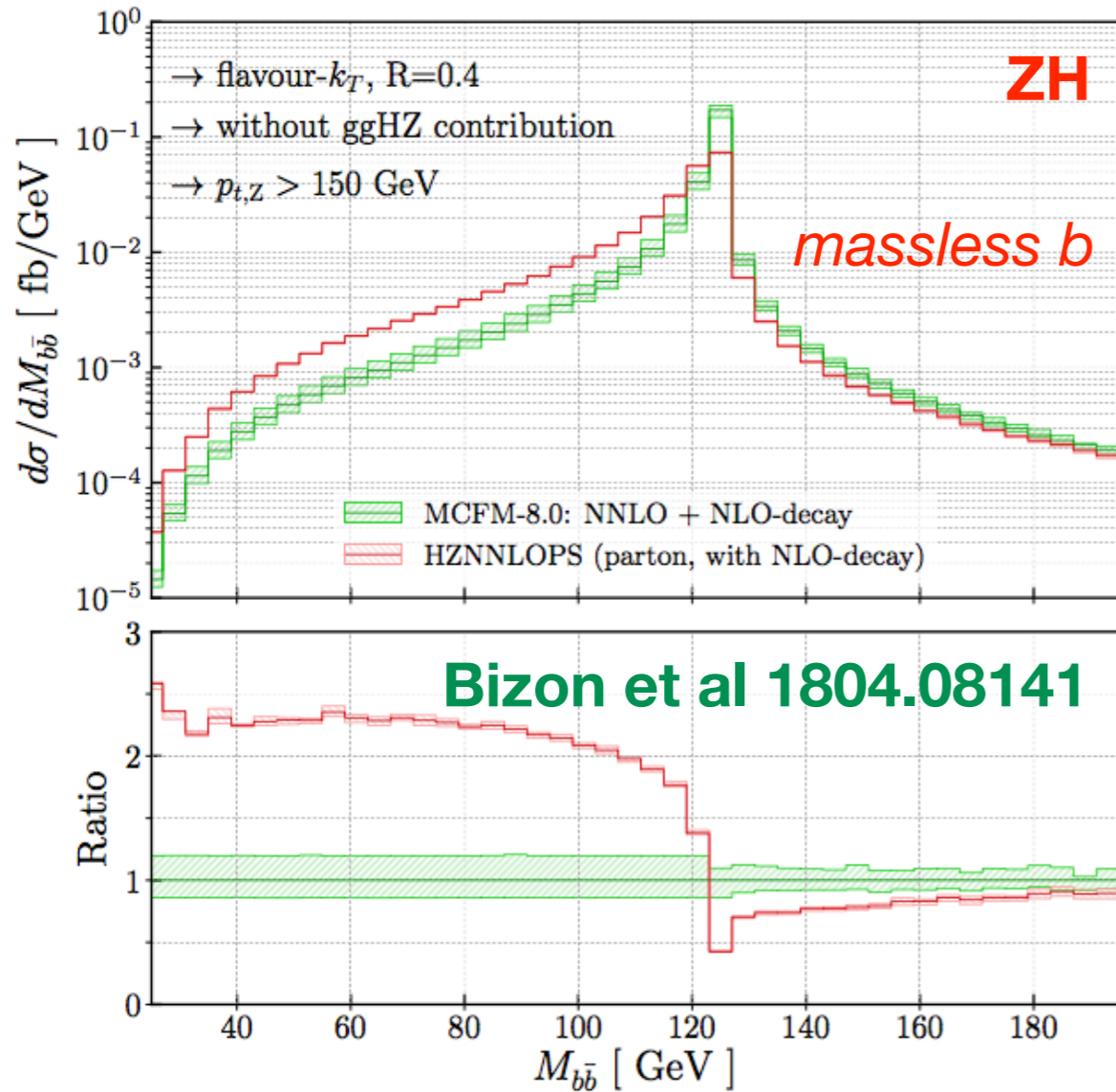
Caola et al 1804.07632 (see also Lindert et al 1703.03886)



- top-bottom interference effects known to NNLL+NLO can lead to distortions up to 5%
- about 20% residual uncertainty

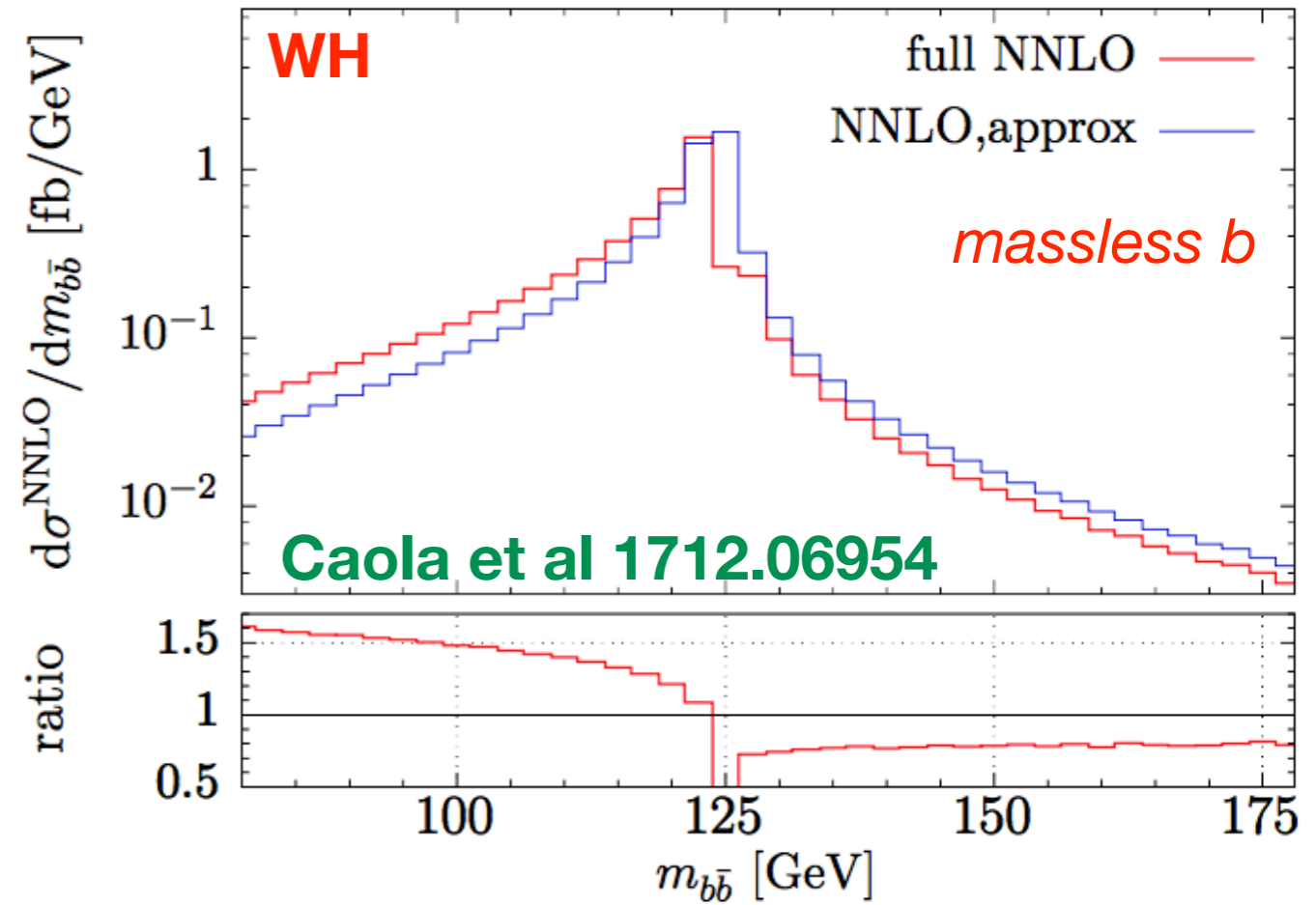
Improvements in VH (\rightarrow bb)

New: NNLOPS for VH + NLO H \rightarrow bb



NNLO_{approx} = NNLO prod.+NLO decay

New: NNLO production + NNLO decay



Large higher-order corrections in H \rightarrow bb decay in m_{bb} and other observables

Theory issue in $H \rightarrow bb$ @ NNLO

Caola et al 1712.06954

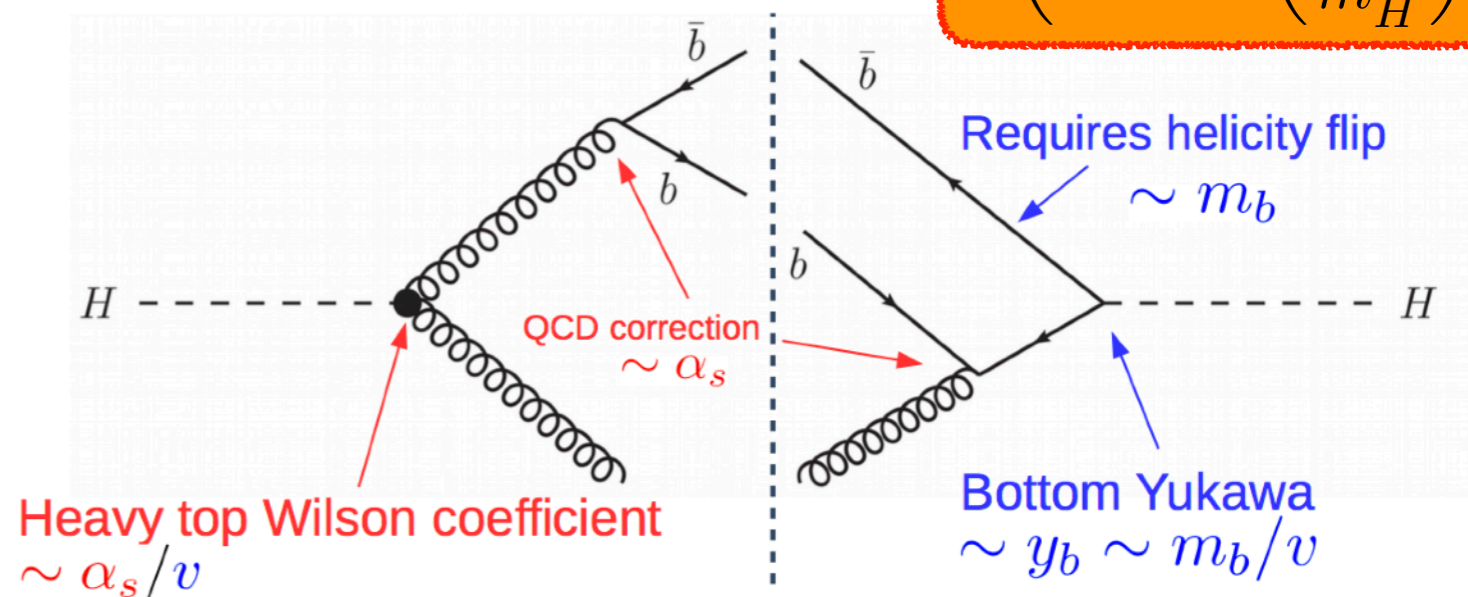
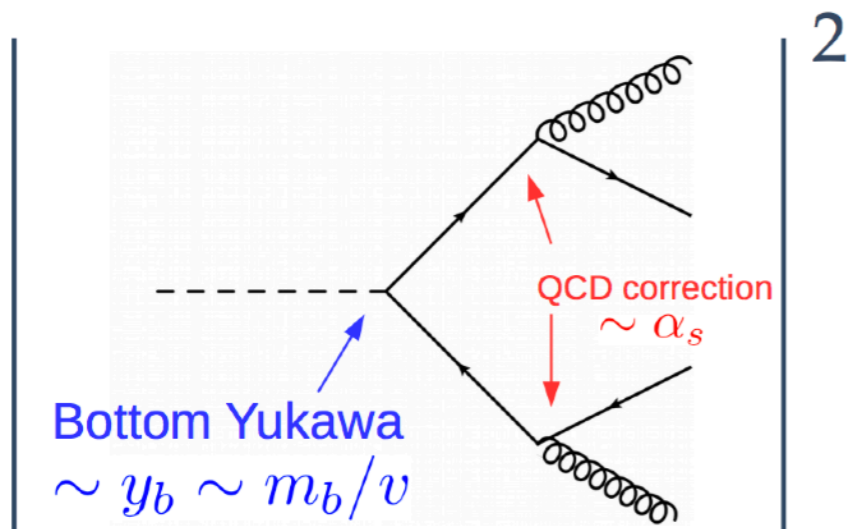
Decay in the massless approximation: extract Yukawa and then set $m_b = 0$

$$\Gamma_{H \rightarrow bb} = y_b^2 \left(\Gamma_{m_b=0} + \mathcal{O} \left(\frac{m_b^2}{m_H^2} \right) \right)$$

Above is OK at LO and NLO, but fails at NNLO as it neglects contributions that are of the same order, i.e. $y_b^2 \alpha_s^2$, that arise in diagrams with a helicity flip (and hence a mass insertion)

Standard NNLO correction: $\mathcal{O}(y_b^2 \alpha_s^2)$

New NNLO contribution: $\mathcal{O} \left(y_b^2 \alpha_s^2 \ln \left(\frac{m_b^2}{m_H^2} \right) \right)$

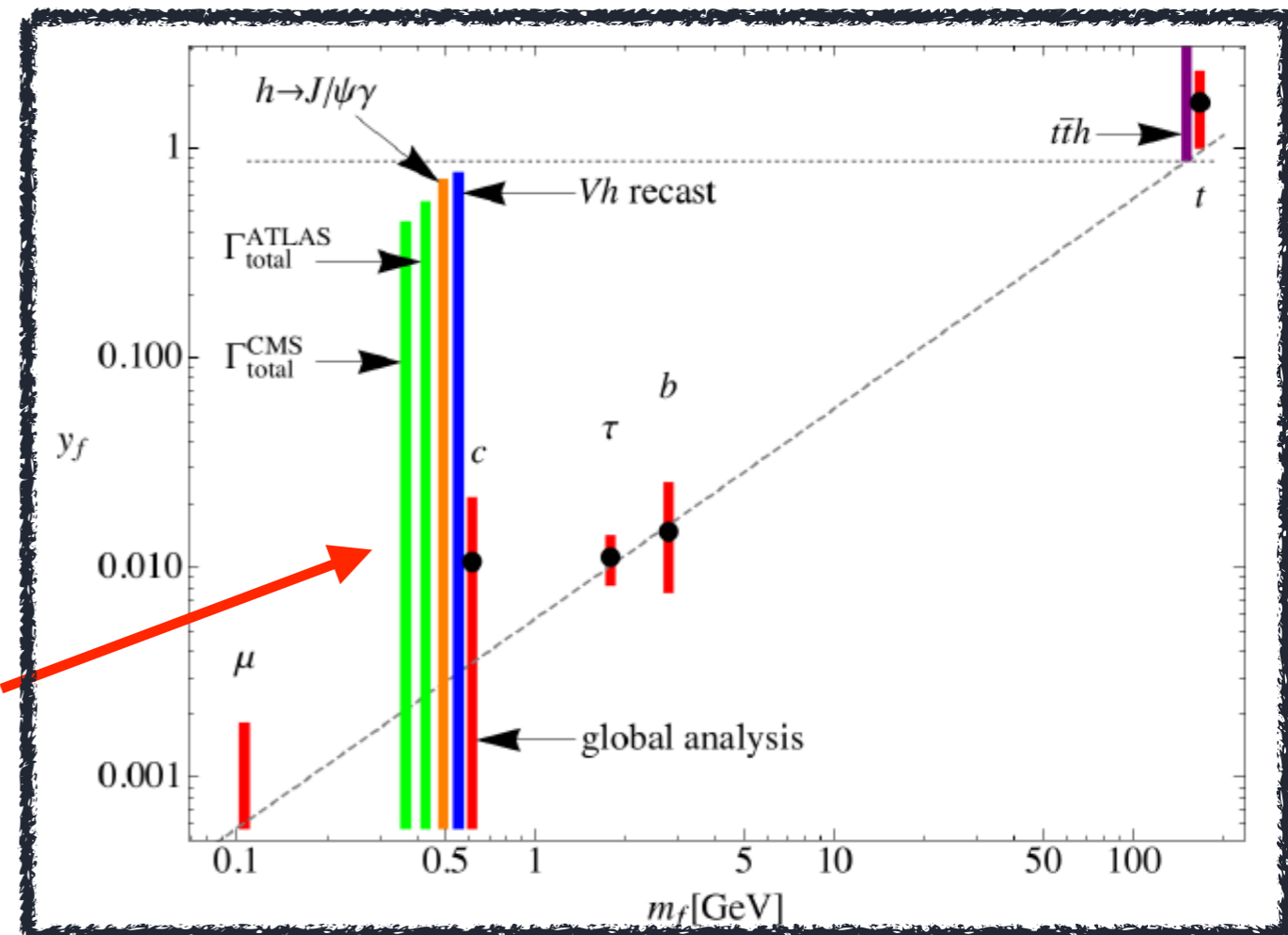


Two examples, out of many, where theoretical precision brings new opportunities in the Higgs sector

1. H coupling to light quarks

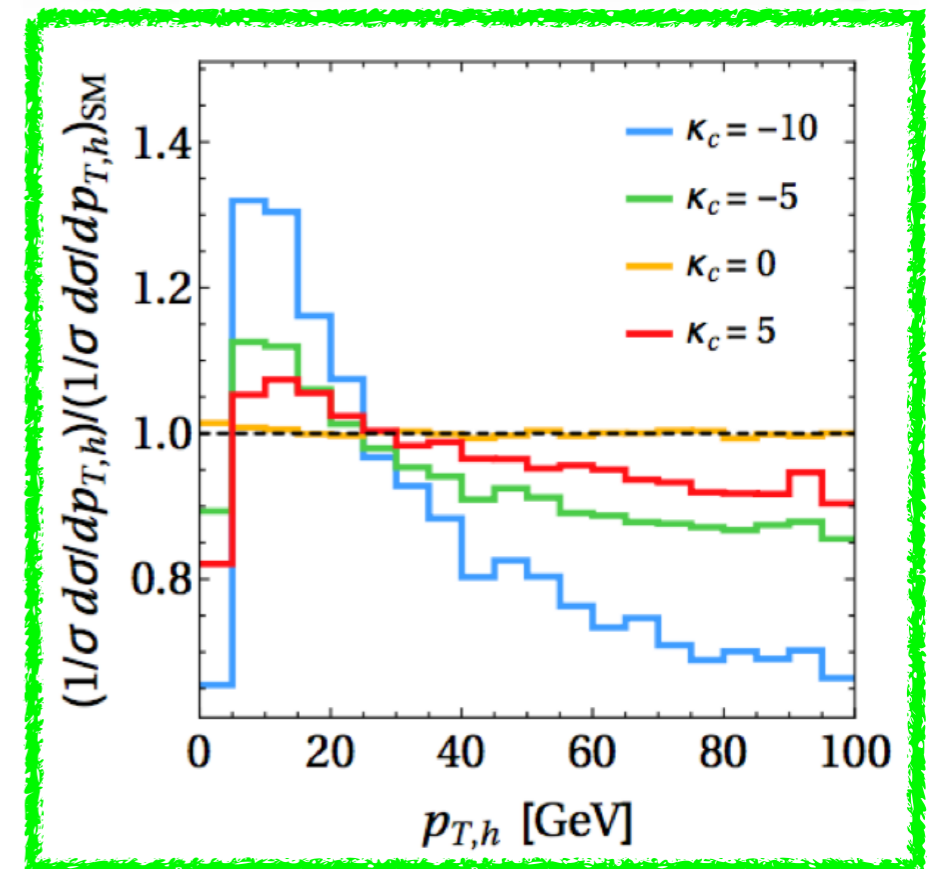
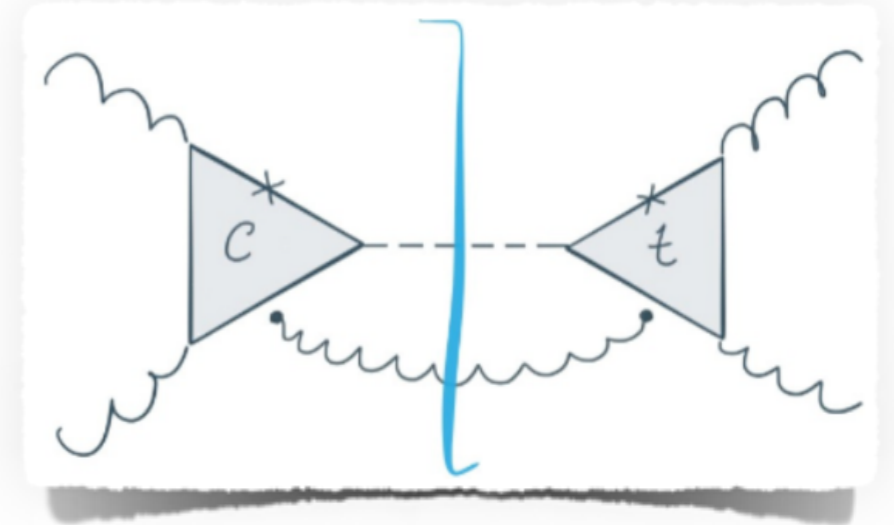
- couplings to 2nd (and 1st) generation notoriously very difficult because they are very small
- a number of ways to constraint the coupling of Higgs to charm:
 - ▶ rare exclusive Higgs decays
 - ▶ Higgs + charm production
 - ▶ constraint from VH (H → bb) including charm mis-tagging
 - ▶ constraint from Higgs width

still largely unconstrained



1. H coupling to light quarks

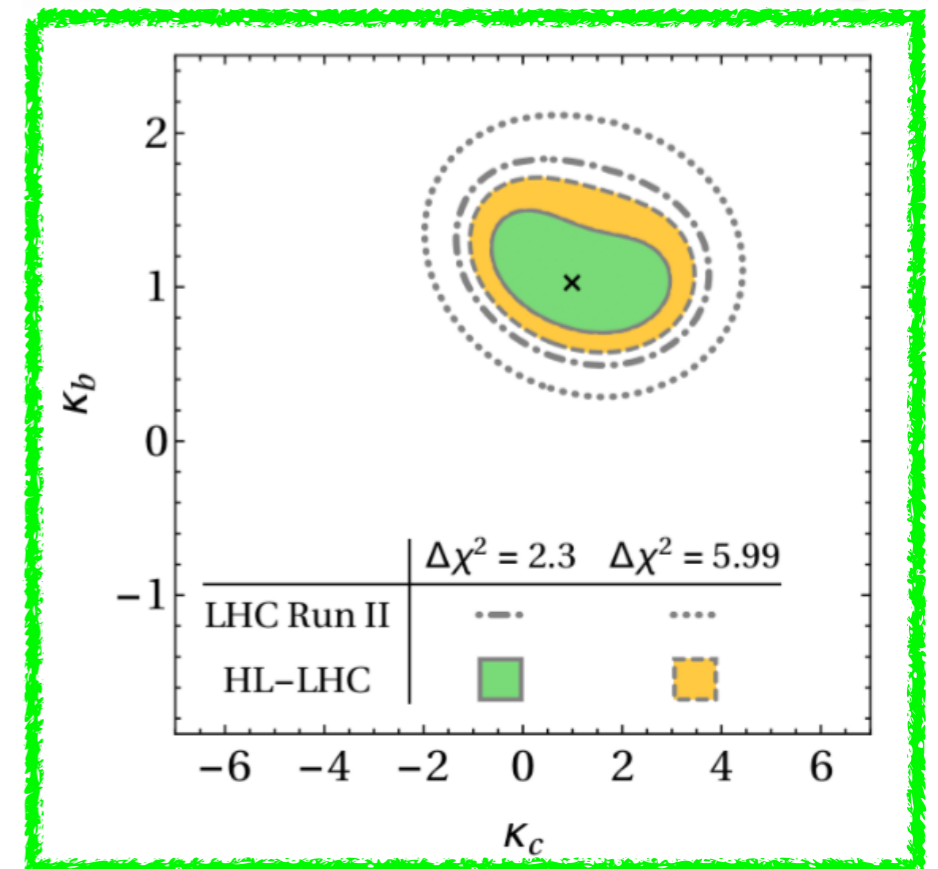
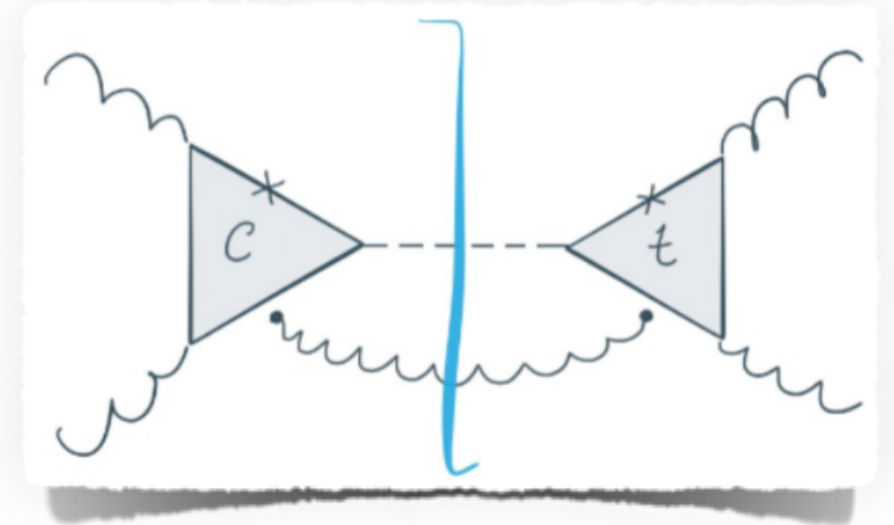
- Higgs produced dominantly via top-quark loop (largest coupling)
- but interference effects with light quarks are not negligible
- provided theoretical predictions are accurate enough (few%?), constraint on charm (and possible strange) Yukawa can be significantly improved



Bishara et al '16

1. H coupling to light quarks

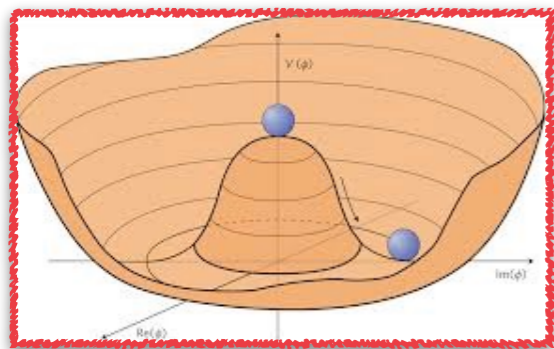
- Higgs produced dominantly via top-quark loop (largest coupling)
- but interference effects with light quarks are not negligible
- provided theoretical predictions are accurate enough (few%?), constraint on charm (and possible strange) Yukawa can be significantly improved



2. The Higgs potential

The Higgs boson is responsible for the masses of all particles we know of. Its potential, linked to the Higgs self coupling, is predicted in the SM, but we have not tested it so far

$$V_{\text{SM}} = \frac{m_h}{2} h^2 + \lambda v h^3 + \frac{\lambda}{4} h^4$$



Single Higgs
done O(45pb)

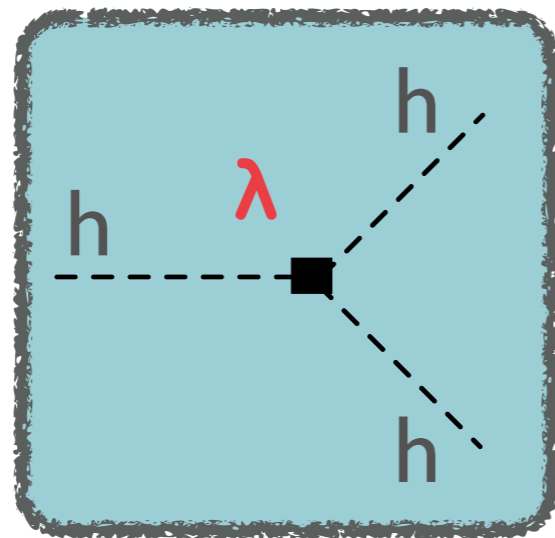
Double Higgs
very hard
O(45fb)

Triple Higgs
out of reach
O(0.1fb)

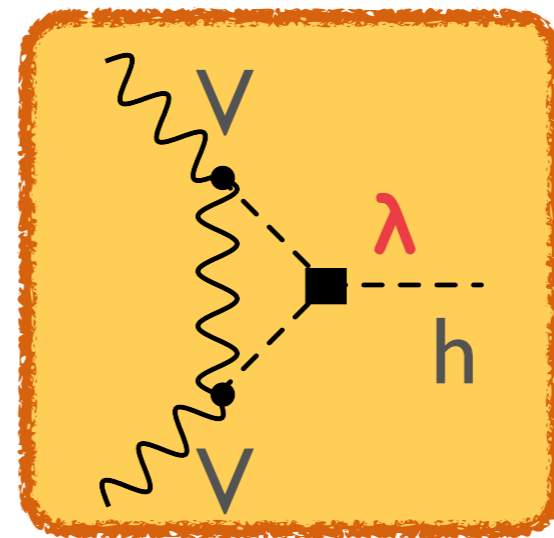
Bounds on λ today from LHC data still very loose (about a factor 10)

2. The Higgs potential

Traditionally: suggested to measure it through the production of two Higgs bosons (but difficult because of very small production rates)



Double Higgs



Single Higgs

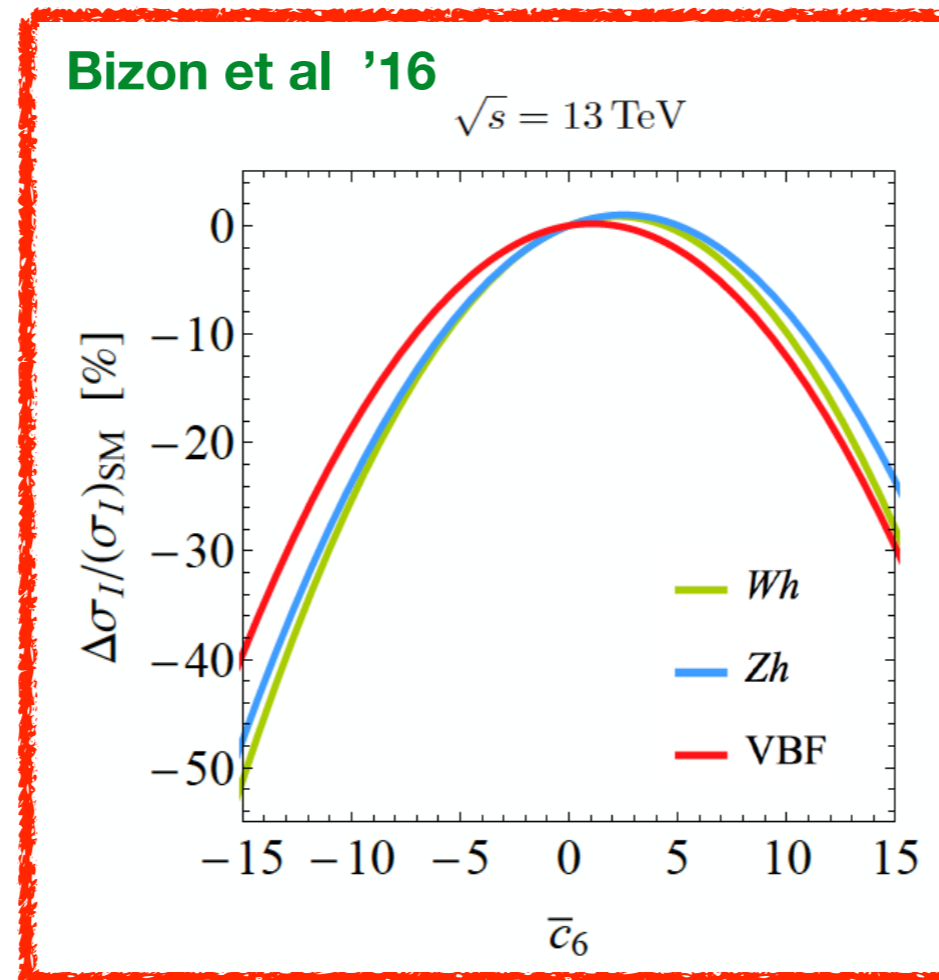
New idea: exploit indirect sensitivity to λ of single Higgs production

Provides a wealth of new measurements (many production processes, many kinematic distributions), but theory and measurements must be accurate enough

2. The Higgs potential

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \frac{\bar{c}_6}{2v^2} \mathcal{O}_6$$

$$\mathcal{O}_6 = -\lambda_{\text{SM}}(H^\dagger H)^3$$



See also
De Grassi et al 1702.01737
Di Vita et al 1704.01953
Maltoni et al 1709.08649
Di Vita et al 1711.03978
[...]

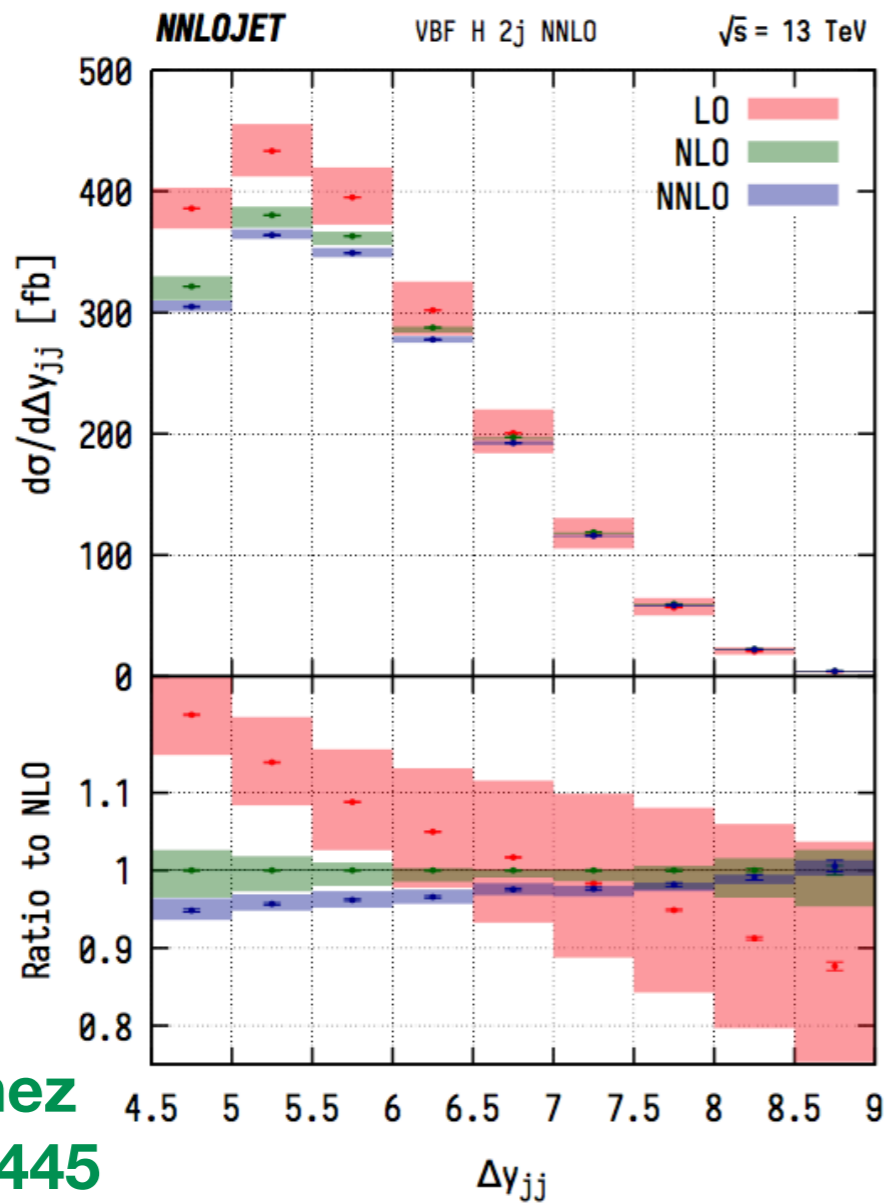
New idea: exploit indirect sensitivity to λ of single Higgs production

Provides a wealth of new measurements (many production processes, many kinematic distributions) to be used in a global fit (but theory must be accurate enough)

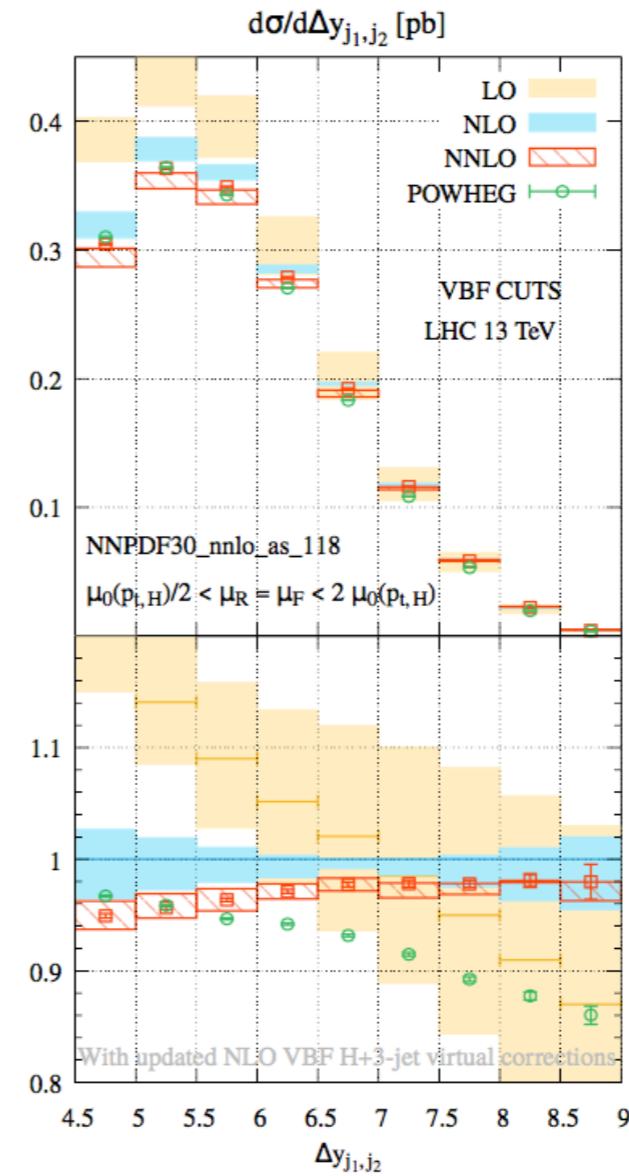
Conclusions

- ▶ Precision physics at hadron colliders is already there
- ▶ Precision Higgs studies in their infancy, much more to come
- ▶ Not just precision measurements of couplings but possibility to address key outstanding questions (Higgs potential, minimal Higgs, fine-tuning, portal to hidden sectors, DM...)
- ▶ Fervid theoretical activity. *This talk focused on selected results of the last few months, many other results that I did not have time to cover ...*

NNLO VBF



Cruz-Martinez et al 1802.02445

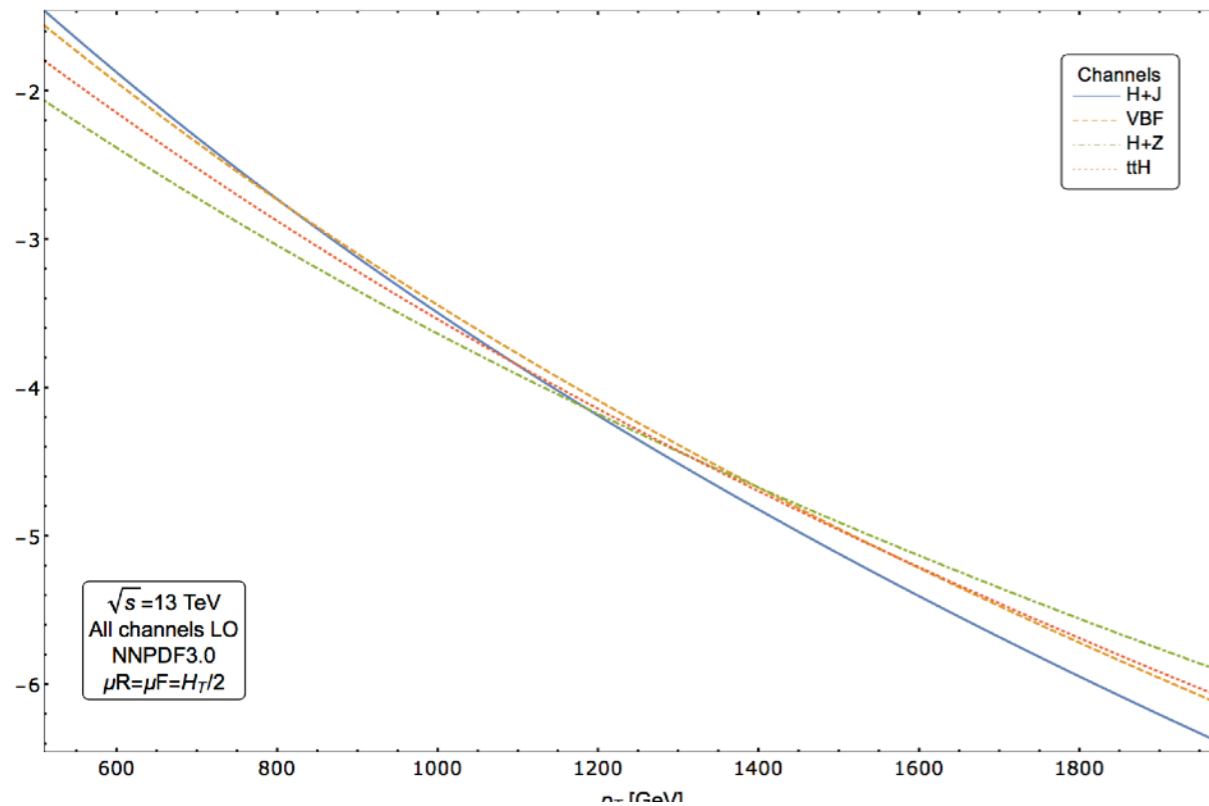


Dryer et al 1506.02660

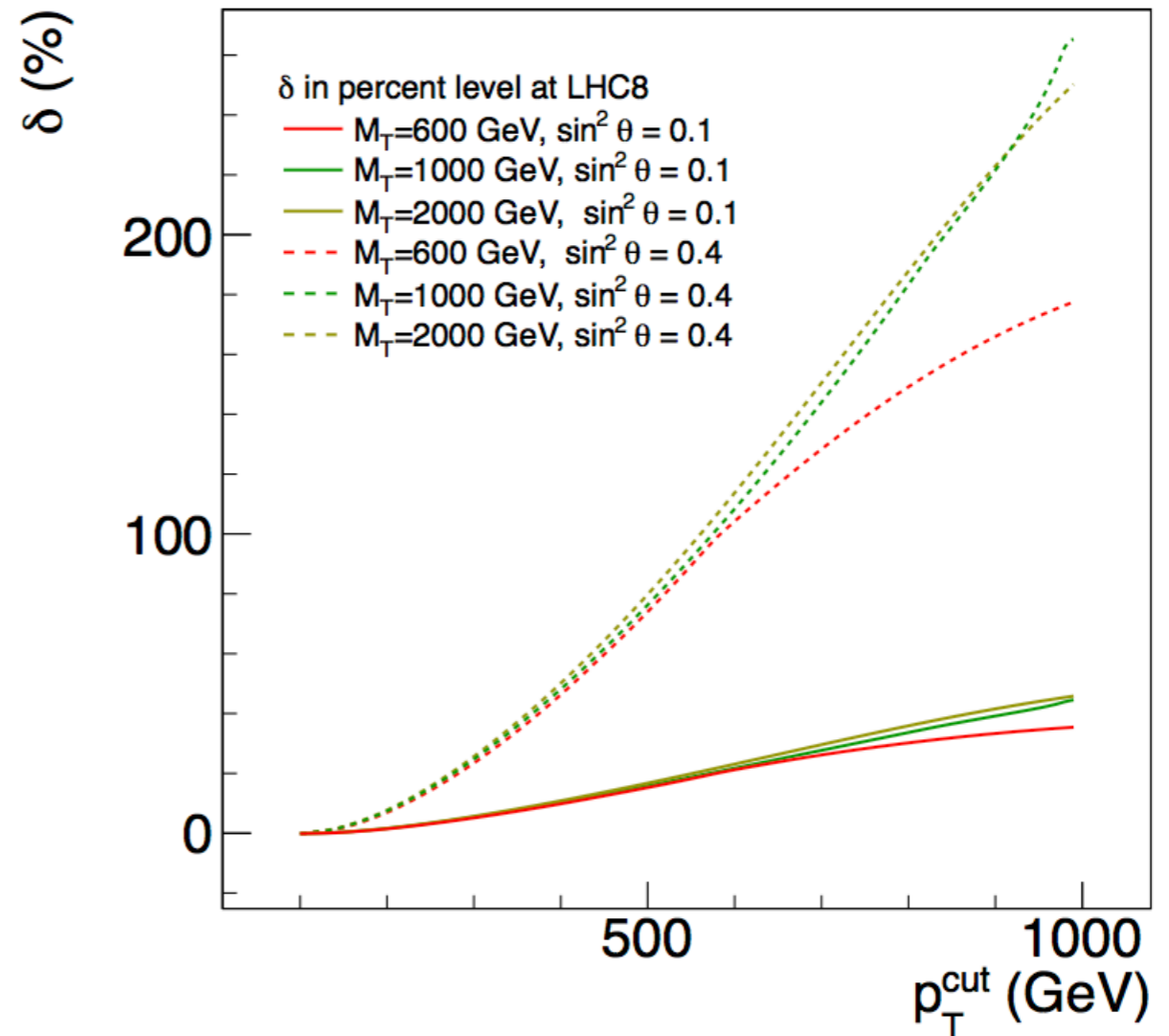
- validation of previous calculation using P2B method
- NNLO important (large, shape distortions) in fiducial regions

Boosted Higgs

- At high p_T hierarchy of modes changes but ggH remains dominant up to about 1 TeV



Banfi, Martin, Sanz 1308.4771



$$\frac{\sigma_H}{dp_{\perp}^2} \sim \frac{\sigma_0}{p_{\perp}^2} \times \begin{cases} (\kappa_g + \kappa_t)^2, & p_{\perp}^2 < 4m_t^2, \\ \left(\kappa_g + \kappa_t \frac{4m_t^2}{p_{\perp}^2}\right)^2, & p_{\perp}^2 > 4m_t^2. \end{cases}$$

- High p_T interesting as it allows to disentangle a point like and a top induced ggH coupling