

PAUL SCHERRER INSTITUT



Andreas Crivellin

Theory Group of the Laboratory for Particle Physics

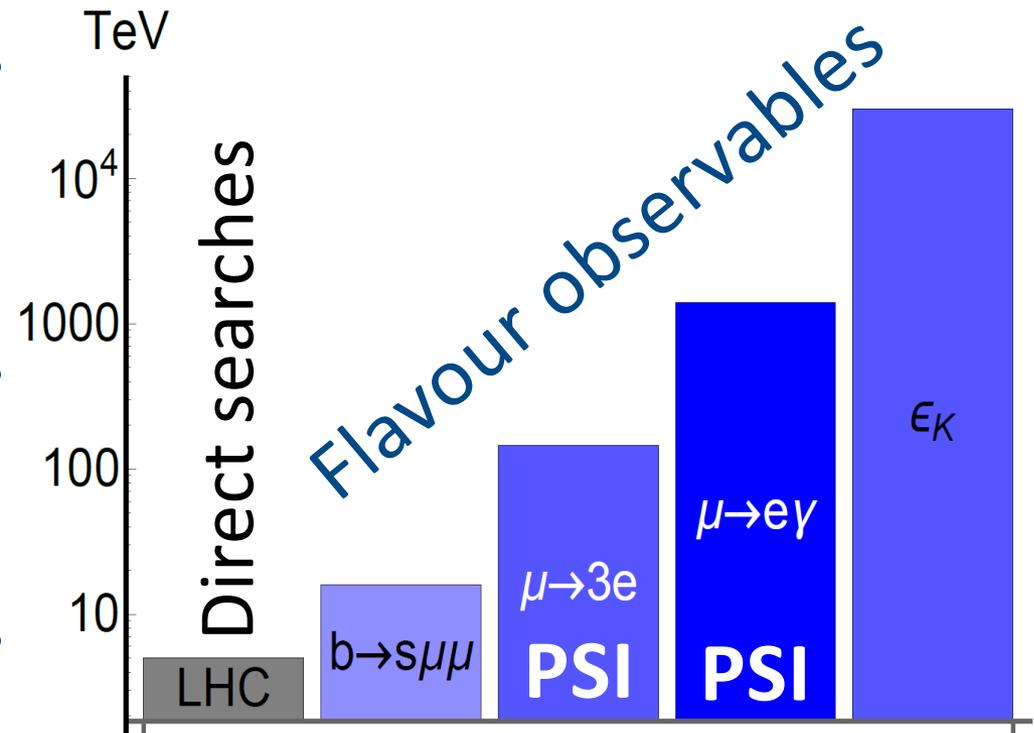
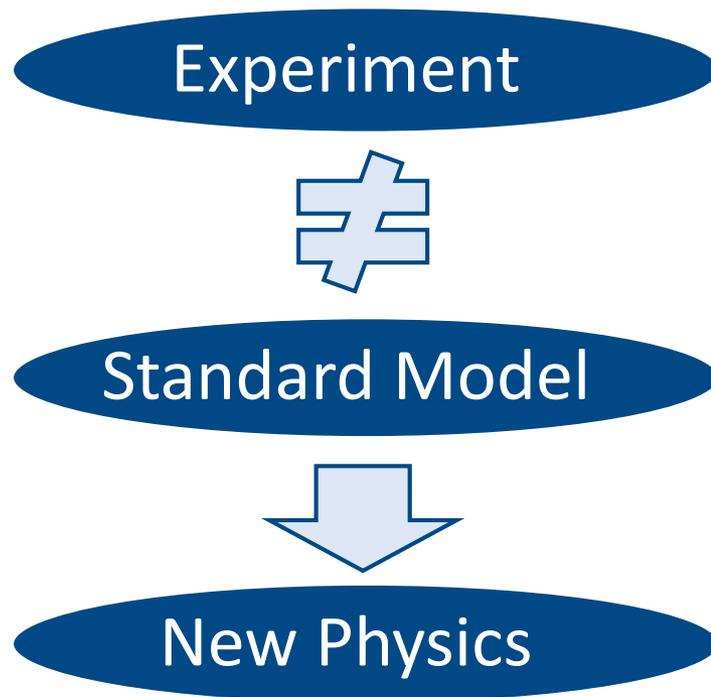
Flavour Anomalies

Bologna, 04.06.2018

- Introduction: Flavour anomalies
 - $b \rightarrow s \mu^+ \mu^-$
 - $b \rightarrow c \tau \nu$
 - a_μ (anomalous magnetic moment of the muon)
- New Physics explanations for the anomalies
 - Z', W'
 - Leptoquarks,
 - MSSM, 2HDMs, extra dimensions,...
- Simultaneous explanations with leptoquarks
- Conclusions and outlook

Finding New Physics with Flavour

- At colliders one produces many (up to 10^{14}) heavy quarks or leptons and measures their decays into light flavours



Flavour observables are sensitive to higher energy scales than collider searches

Sign In | Register  1

SCIENTIFIC AMERICAN™

Search ScientificAmerican.com 

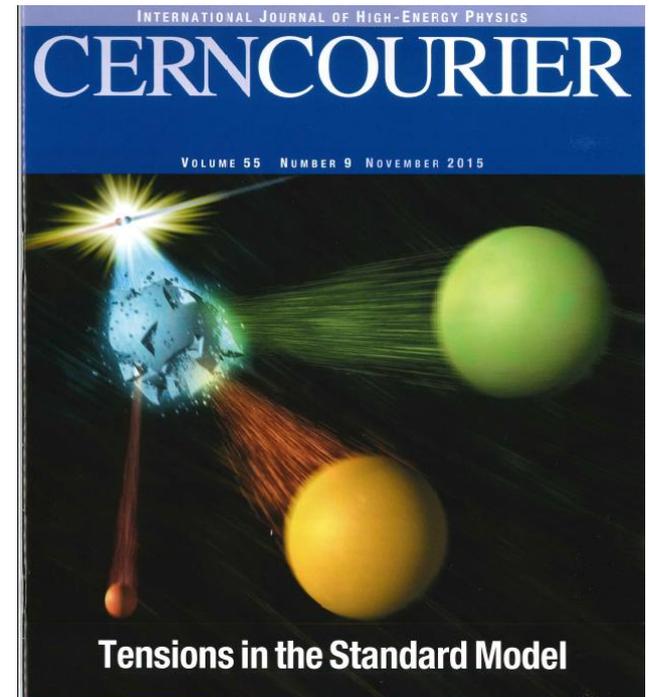
[Subscribe](#) [News & Features](#) [Topics](#) [Blogs](#) [Videos & Podcasts](#) [Education](#) [Citizen Science](#)

The Sciences » News 13 ::  Email ::  Print

2 Accelerators Find Particles That May Break Known Laws of Physics

The LHC and the Belle experiment have found particle decay patterns that violate the Standard Model of particle physics, confirming earlier observations at the BaBar facility

By Clara Moskowitz | September 9, 2015 | [Véalo en español](#)



physicstoday

[Home](#) [Print Edition](#) [Daily Edition](#) ▾ [About](#) [Jobs](#)  [Subscribe](#)

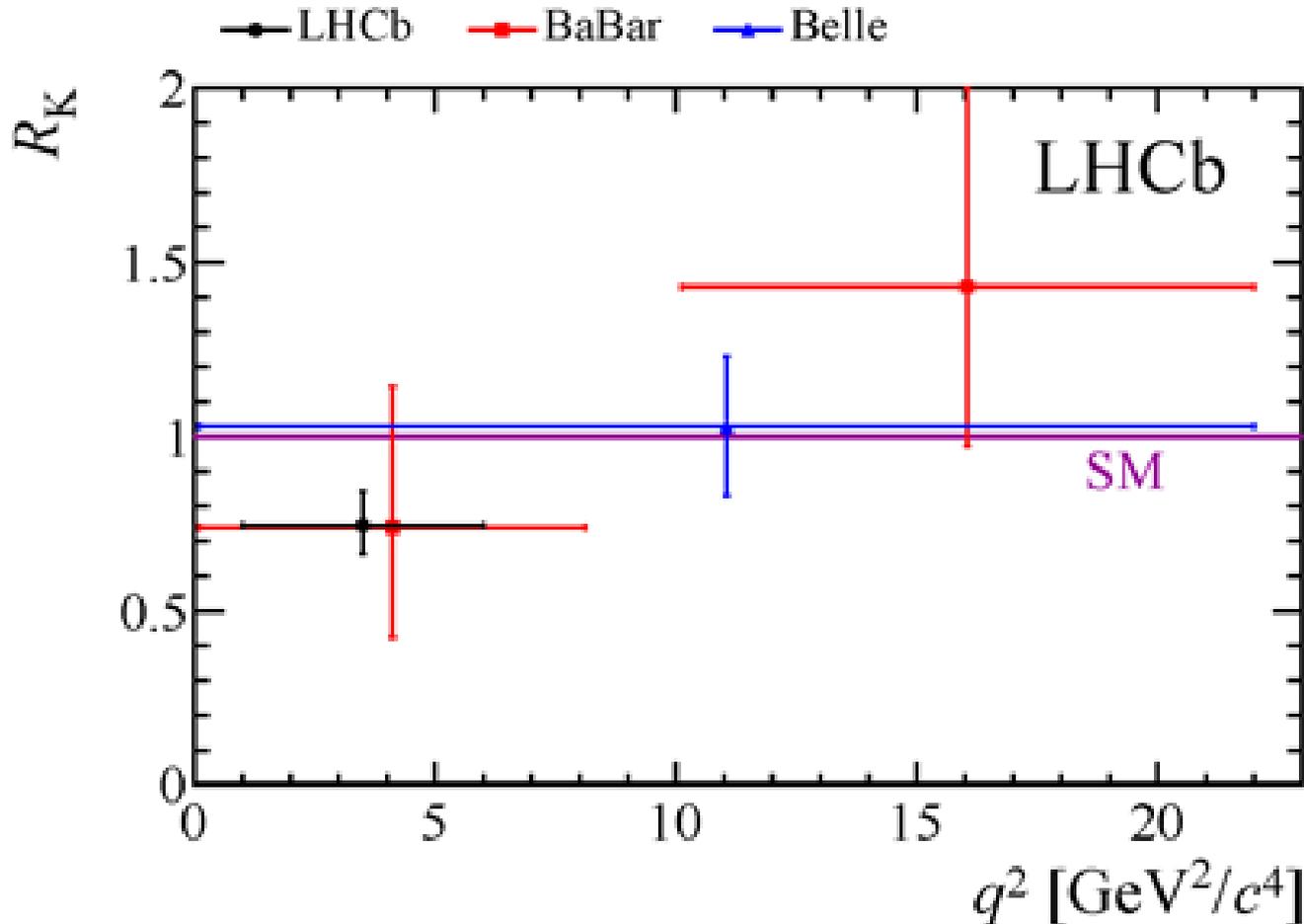
Democracy suffers a blow—in particle physics

Three independent B-meson experiments suggest that the charged leptons may not be so equal after all.

Steven K. Blau 17 September 2015

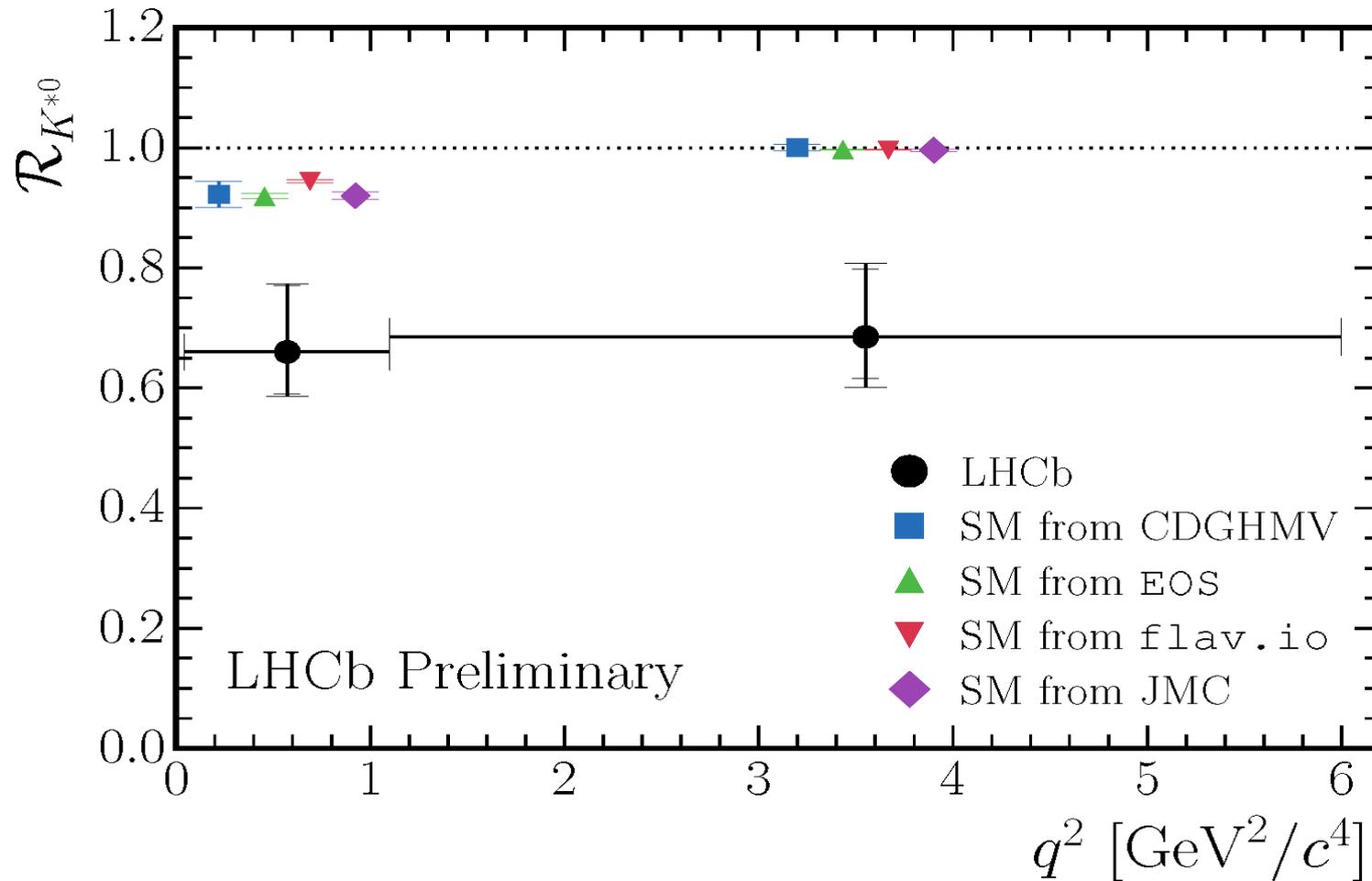
Hints for
New Physics
in flavour
observables

$$R(K^{(*)}) = \mathcal{B}(B \rightarrow K^{(*)} \mu^+ \mu^-) / \mathcal{B}(B \rightarrow K^{(*)} e^+ e^-)$$



Combined $\approx 4\sigma$ evidence for LFUV

$$R(K^{(*)}) = B \rightarrow K^{(*)} \mu^+ \mu^- / B \rightarrow K^{(*)} e^+ e^-$$



Combined $\approx 4\sigma$ evidence for LFUV

Global fit to $b \rightarrow s \mu^+ \mu^-$ data

- Global analyses of all $b \rightarrow s \mu^+ \mu^-$ data gives a very good fit to data
- Good fit to data:

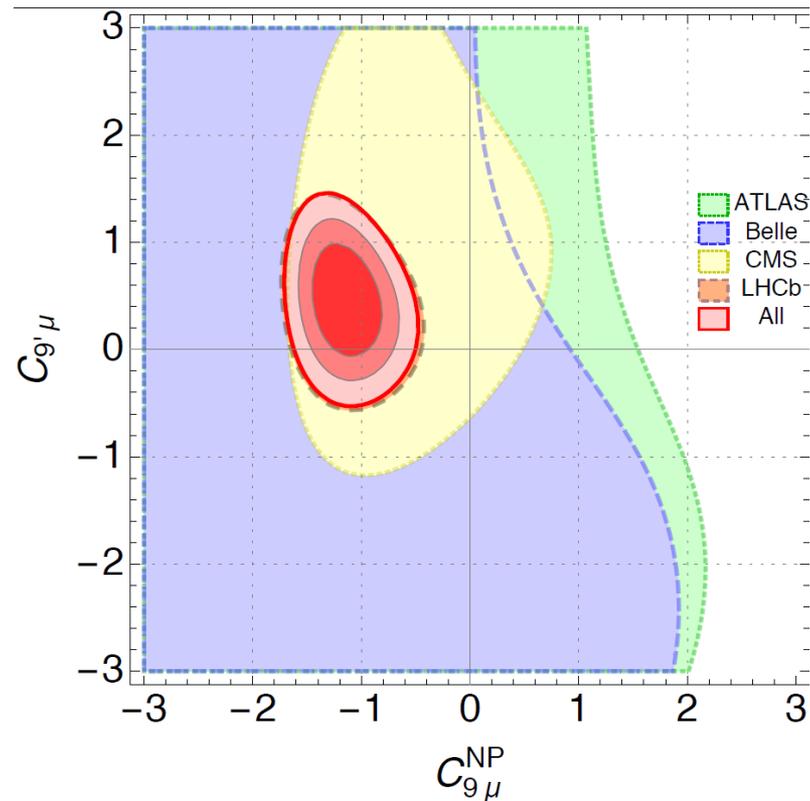
- C_9

- $C_9 = -C_{10}$

- $C_9 = -C'_9$

$$O_9 = \bar{s} \gamma^\mu P_L b \bar{\ell} \gamma_\mu \ell$$

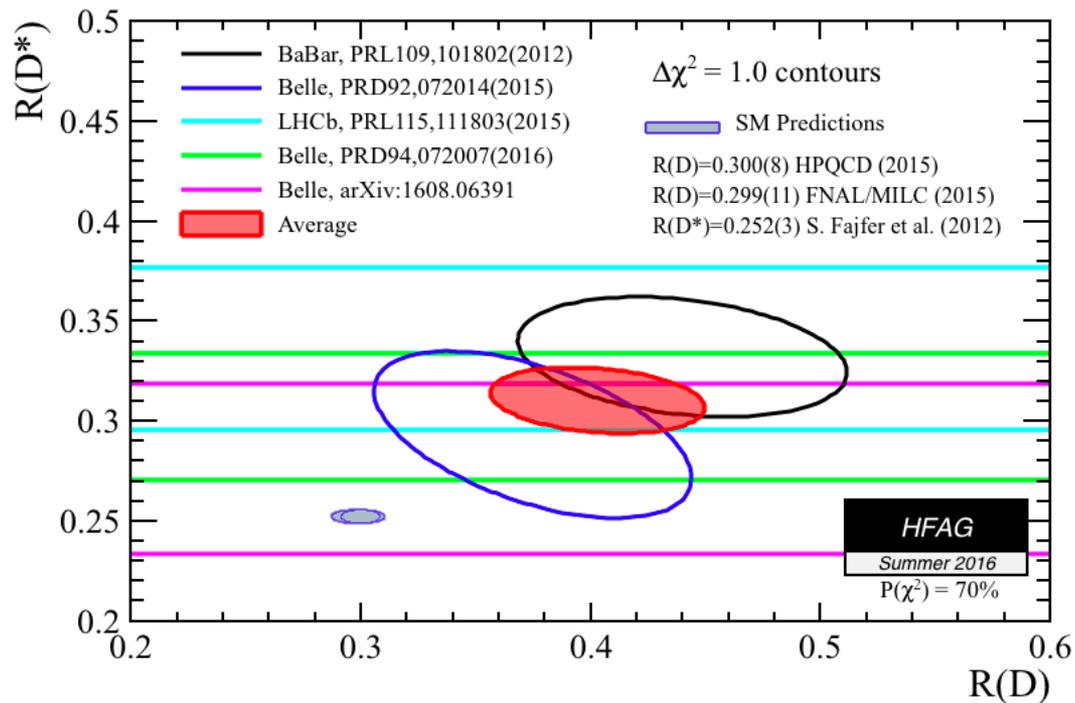
$$O_{10} = \bar{s} \gamma^\mu P_L b \bar{\ell} \gamma_\mu \gamma^5 \ell$$



B. Capdevila, AC, S. Descotes-Genon, J. Matias and J. Virto, arXiv:1704.05340 [hep-ph].

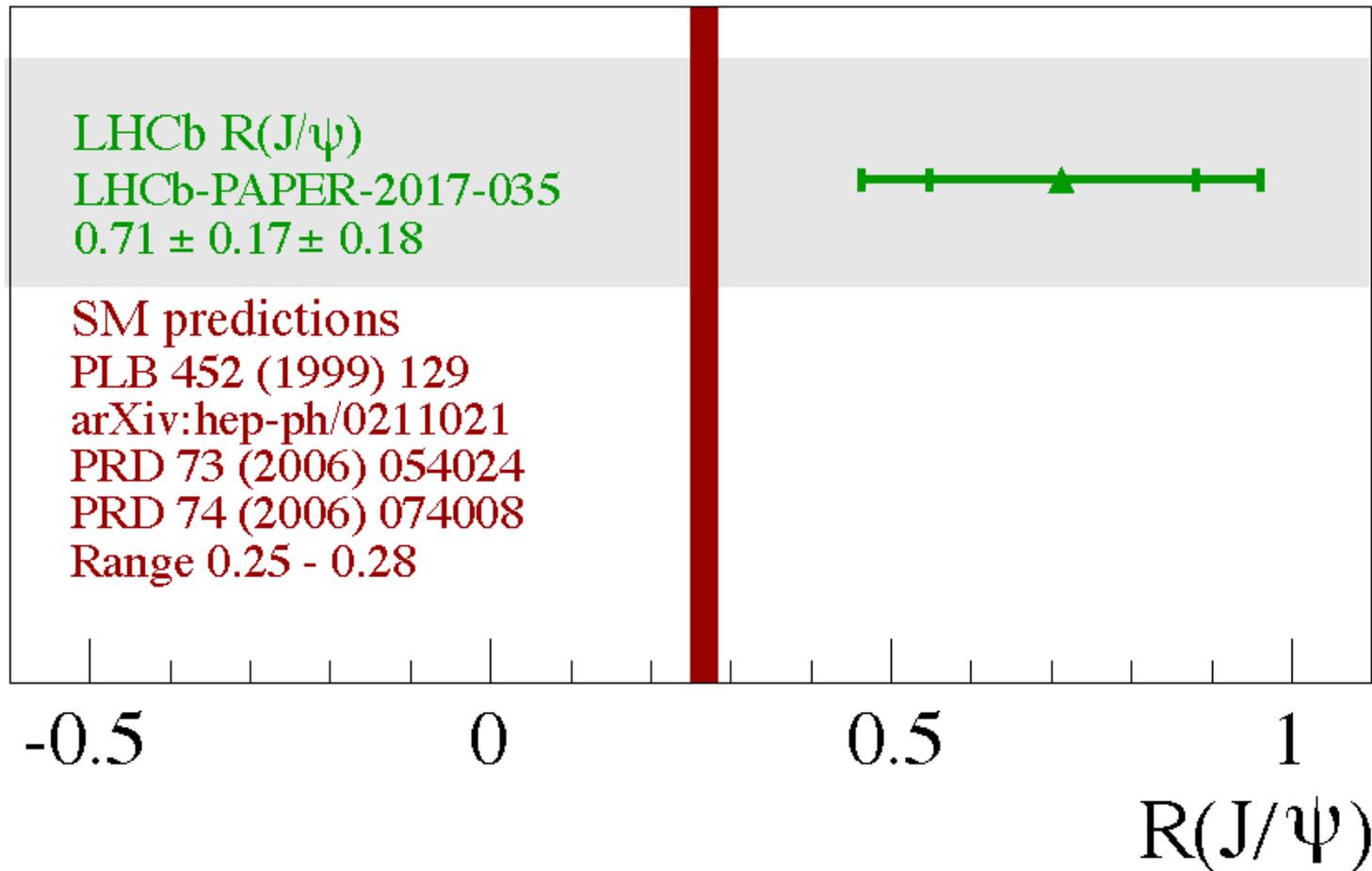
Fit is $>5 \sigma$ better than the SM

$$R(D^{(*)}) = B \rightarrow D^{(*)} \tau \nu / B \rightarrow D^{(*)} \ell \nu$$



All measurements above the SM prediction
4 σ deviation

$b \rightarrow c\tau\nu$ processes

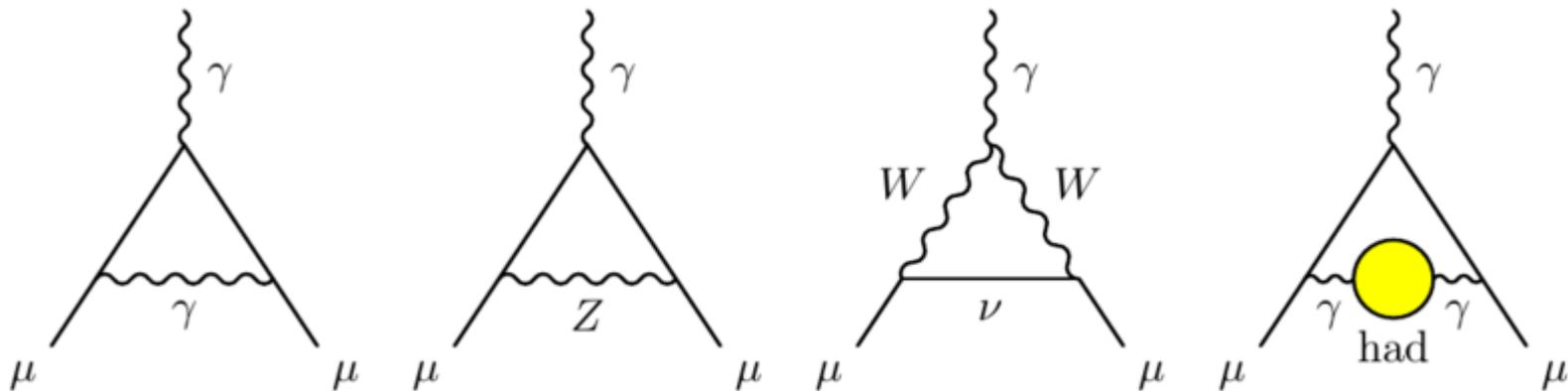


All measurements above the SM prediction
4 σ deviation

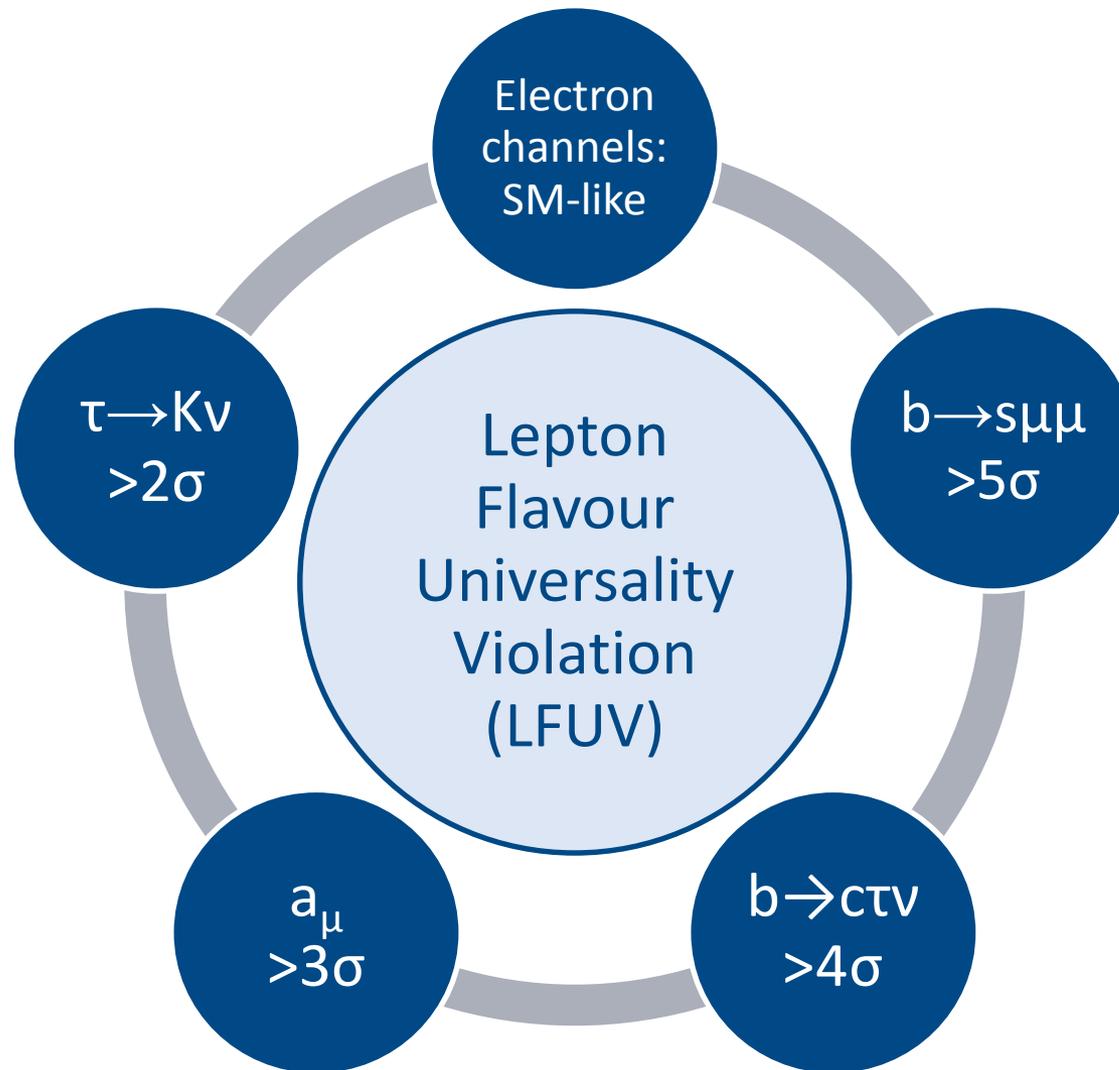
- Single measurement from BNL
- Theory prediction sound but challenging because of hadronic effects.

$$\Delta a_{\mu} = (236 \pm 87) \times 10^{-11}$$

- Soon new experimental results from Fermilab



3 σ deviation (order of SM-EW contribution)



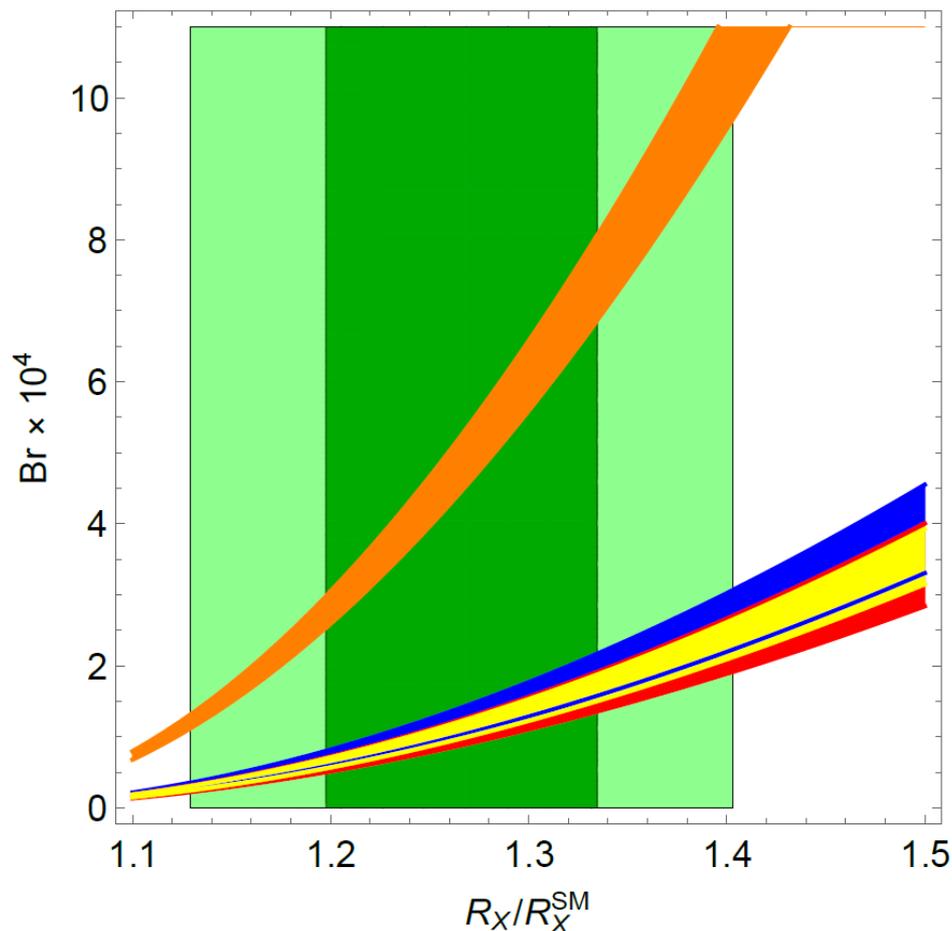
Probability
for
statistical
fluctuation
 $< 0.0001\%$

- Charged scalars
 - Problems with q^2 distributions and B_c lifetime
- W'
 - Strong constraints from direct LHC searches
- Leptoquark (also in the RPV MSSM)
 - EW precision constraints
 - Strong signals in $qq \rightarrow \tau\tau$ searches

Explanation difficult but possible with
Leptoquarks

$R(D^{(*)})$ and $b \rightarrow s\tau\tau$ (model-independent)

- Large couplings to the second generation
- Cancellation in $b \rightarrow svv$ needed: $C^{(1)}=C^{(3)}$



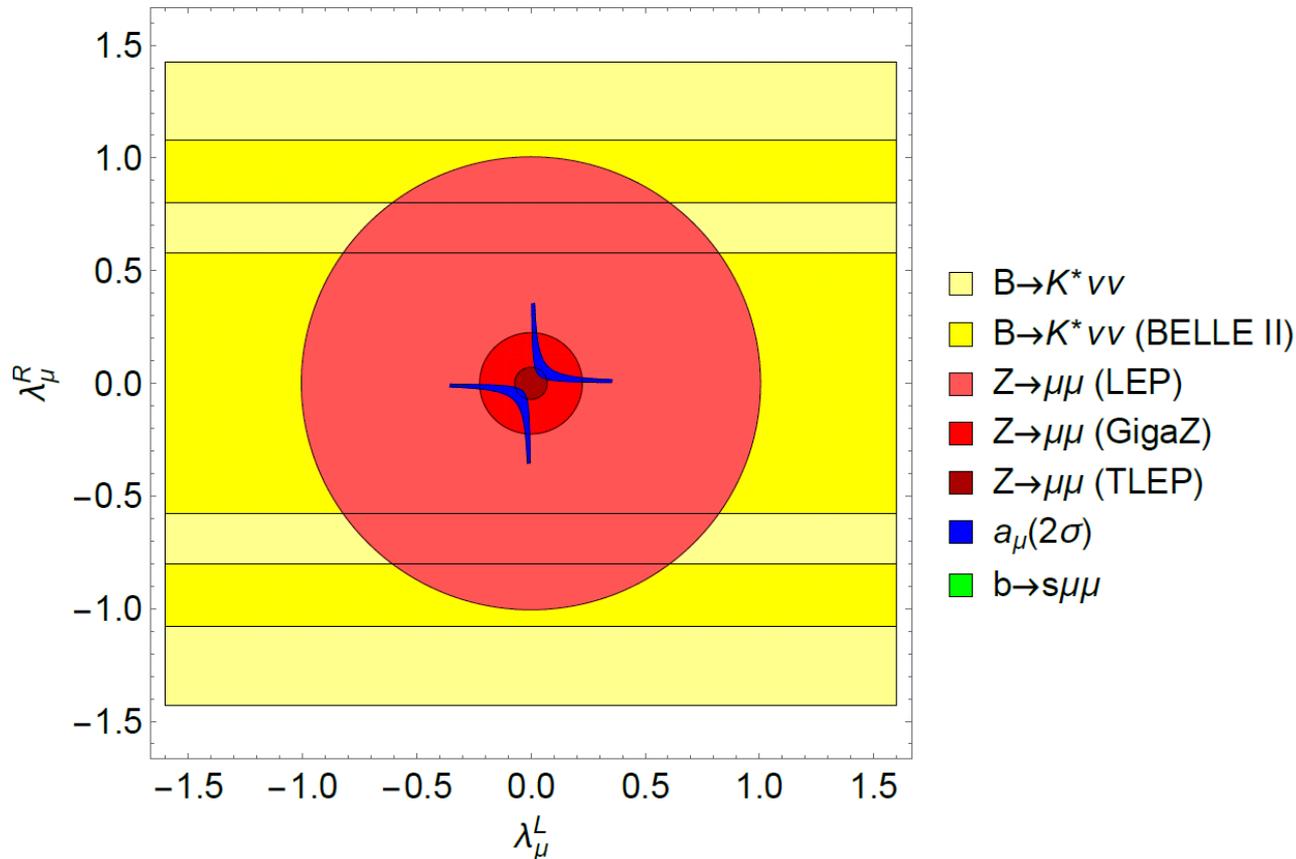
$b \rightarrow s\tau\tau$
very
strongly
enhanced

B. Capdevila, A.C., S. Descotes-Genon,
L. Hofer and J. Matias, PRL.120.181802

- MSSM
 - $\tan(\beta)$ enhanced slepton loops
- Scalars
 - Light scalars with enhanced muon couplings
- Z'
 - Very light with $\tau\mu$ couplings (m_τ enhancement)
- Leptoquarks
 - m_t enhanced effects

Chiral enhancement or very light particles

■ Chirally enhanced effects via top-loops



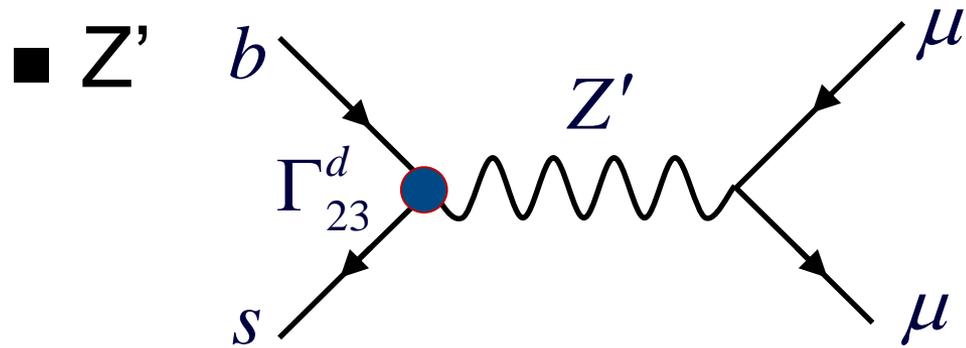
$\lambda_\mu^{L,R}$

Left-, right-
handed
muons-top
coupling

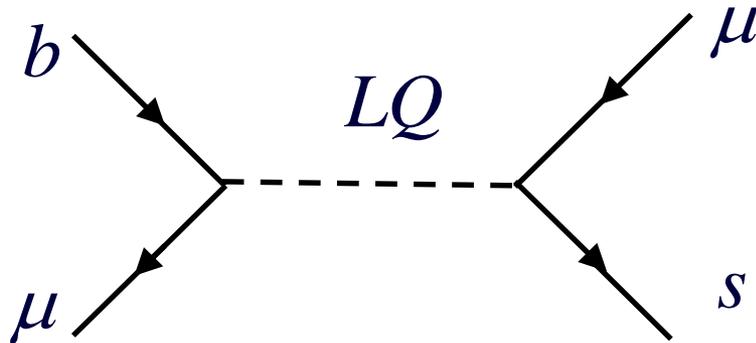
E. Leskow, A.C.,
G. D'Ambrosio,
D. Müller
arXiv:1612.06858

$Z \rightarrow \mu \mu$ at future colliders

$b \rightarrow s \mu \mu$ explanations



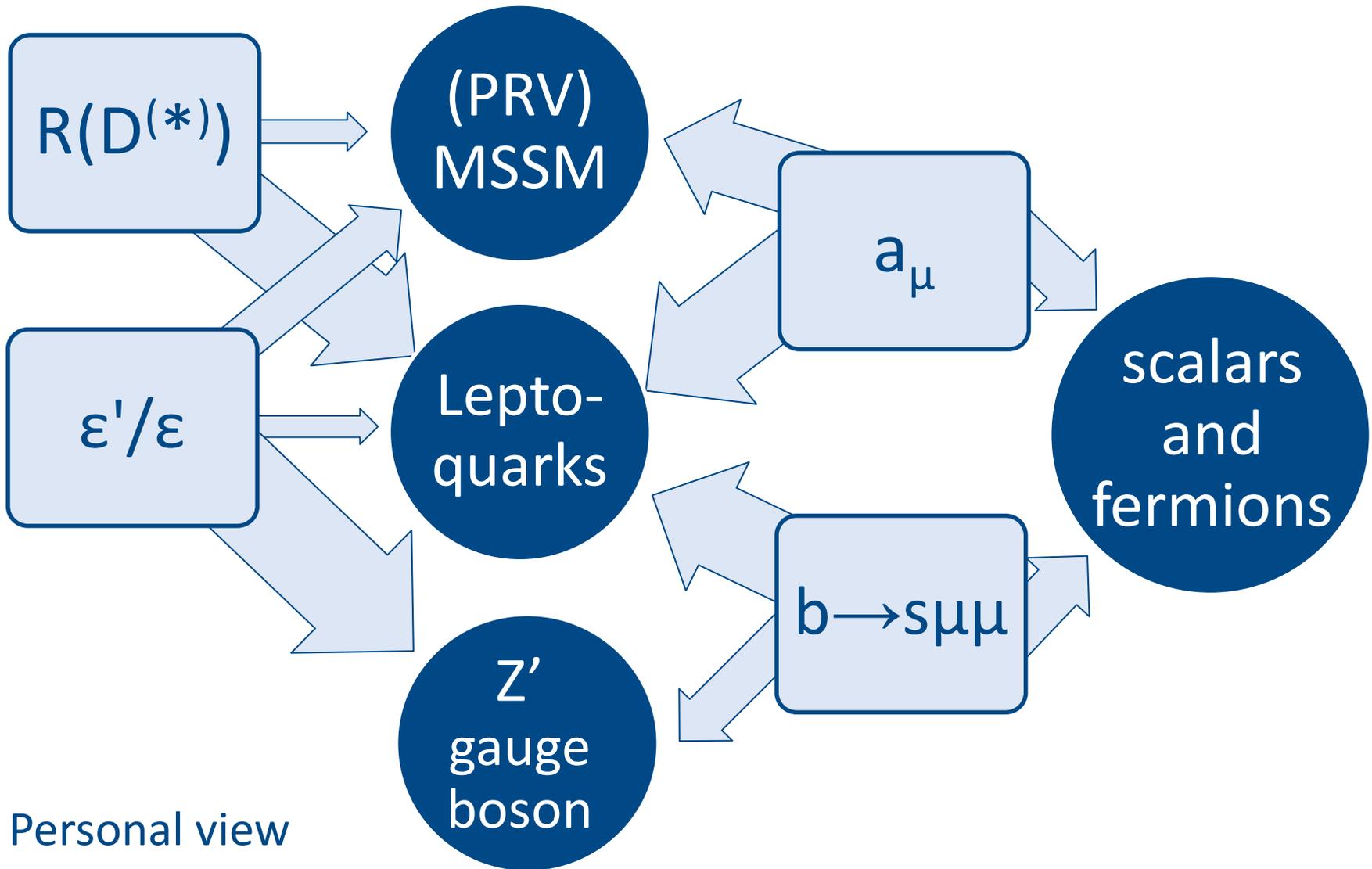
■ Leptoquarks



■ Loop effects of scalars and fermions

Even high scale NP explanations possible

Implications for New Particles



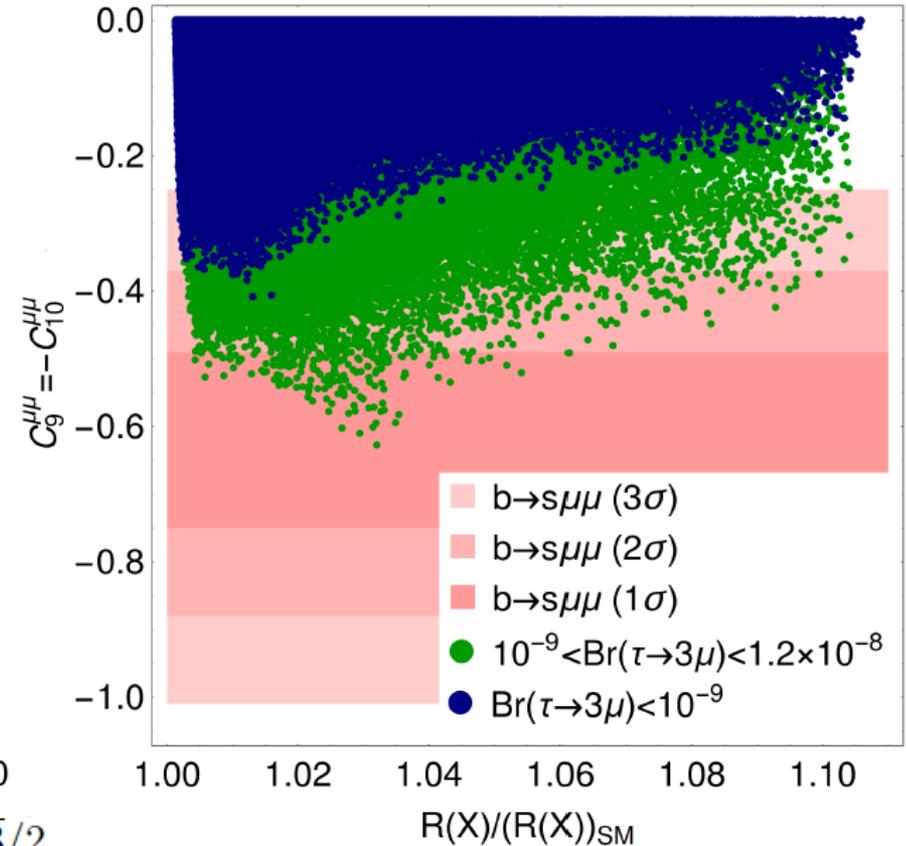
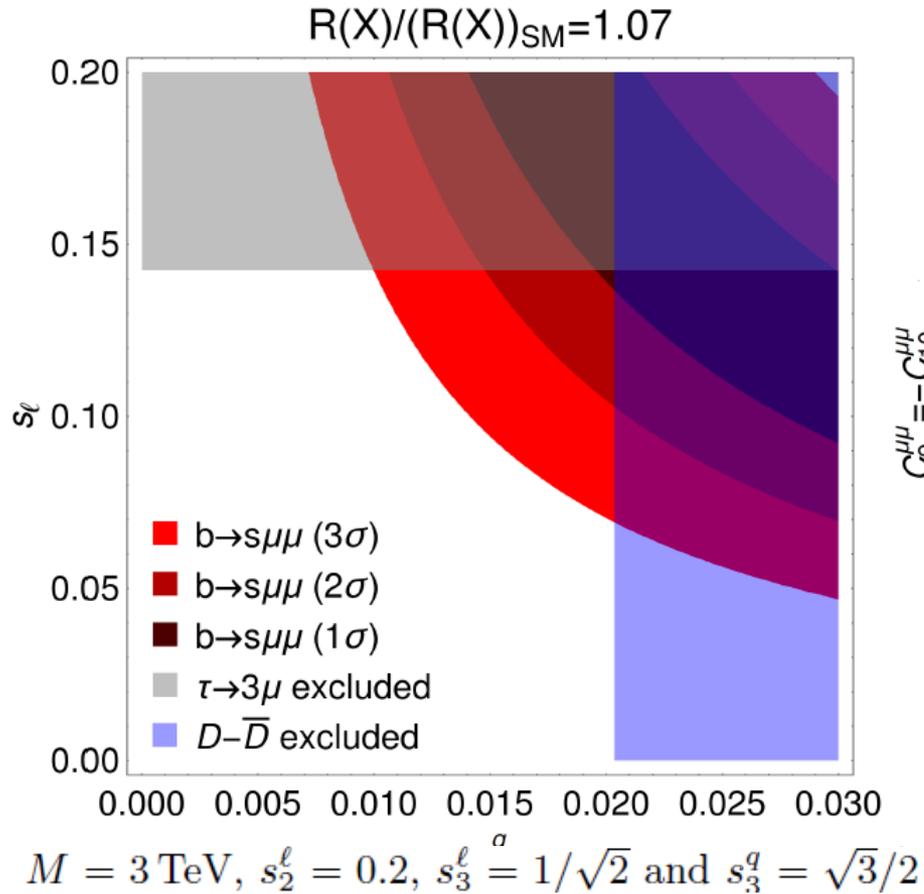
- $C_9 = -C_{10}$ effect in $b \rightarrow s \mu \mu$
- Left handed vector current in $R(D)$ and $R(D^*)$
- No effect in $b \rightarrow s \nu \nu$
- No proton decay
- Contained within the Pati-Salam model
- Massive vector bosons
 - Non-renormalizable without Higgs mechanism
 - Pati Salam not possible at the TeV scale because of $K_L \rightarrow \mu e$ and $K \rightarrow \pi \mu e$

Good solution, but difficult UV completion

M. Blanke, AC, arXiv:1801.07256

- $SU(4) \otimes SU(2)_L \otimes SU(2)_R$ broken to the SM via boundary conditions on a compact extra dimension
- Zero modes: SM fermions
- KK modes: Vector-like fermions and massive gauge bosons
- No zero mode for the Leptoquark
- Flavour alignment to the down-sector

PS + RS naturally accounts for a vector LQ + VLFs

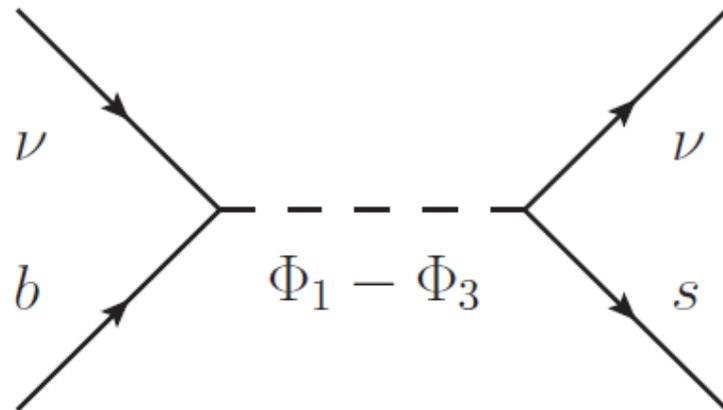
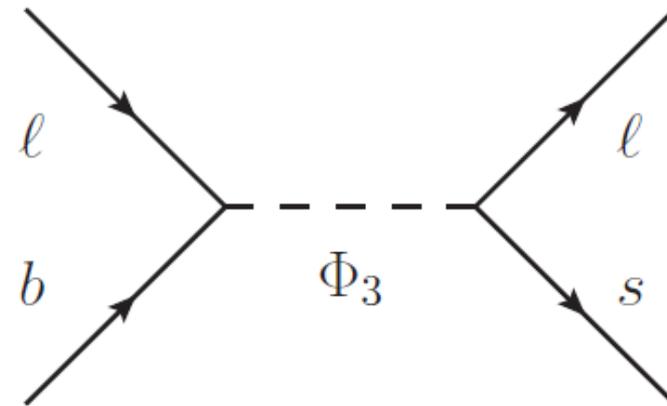
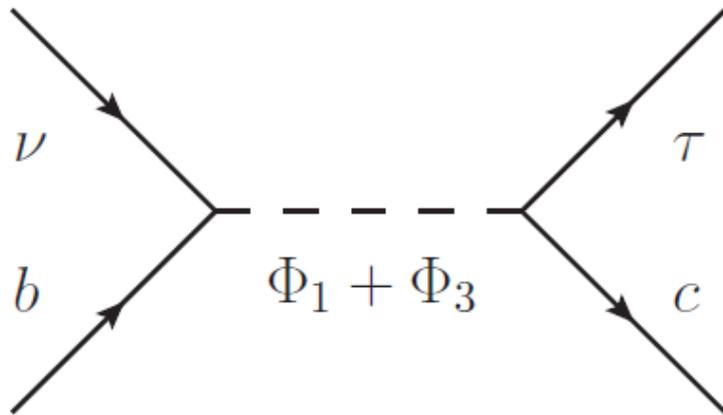


Modell well motivated + limited but sizable effect

- P5'  $b \rightarrow d\mu\mu$
- R(D) & R(D*)  $b \rightarrow s\tau\tau$
- R(K) & R(K*)  $\mu \rightarrow e\gamma$
- R(D), R(D*) & a_μ  $\tau \rightarrow \mu\gamma$
- R(D), R(D*) & $b \rightarrow s\mu\mu$  $b \rightarrow s\tau\mu$

Exciting times in particle physics are ahead of us!

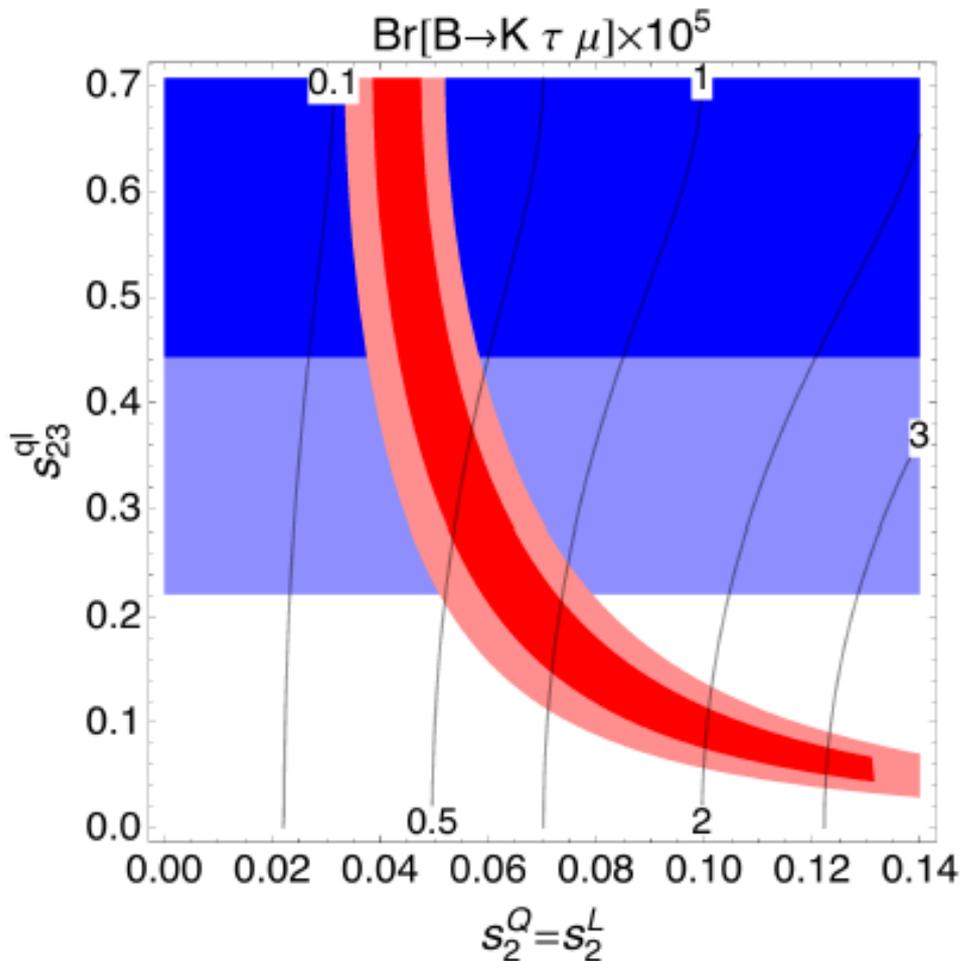
- Φ_1 scalar leptoquark singlet with $Y=-2/3$
- Φ_3 scalar leptoquark triplet with $Y=-2/3$



Constructive in $R(D^{(*)})$

Destructive in $b \rightarrow s \mu \mu$

R(D^(*)) and b → s μ μ



$$s_3^Q = s_3^L = \frac{1}{\sqrt{2}}$$

$$M = 1.5 \text{ TeV}$$

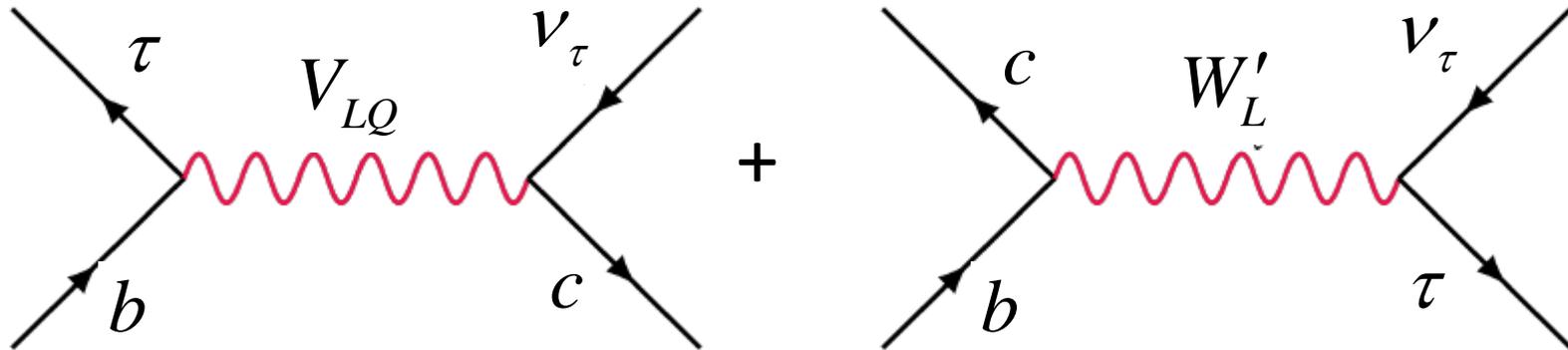
$$s_i^Q = \frac{\frac{m_{ii}^Q}{M_{ii}^Q}}{\sqrt{1 + \frac{m_{ii}^Q}{M_{ii}^Q}}}$$

- R(D^(*)) 2σ
- R(D^(*)) 1σ
- C₉^{μμ} = -C₁₀^{μμ} 2σ
- C₉^{μμ} = -C₁₀^{μμ} 1σ

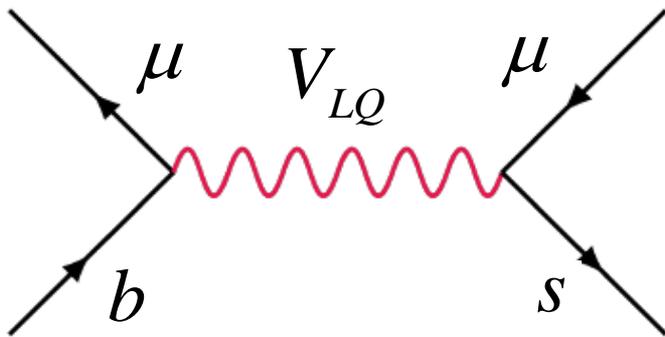
LHCb bounds require additional heavy neutral fermions

Simultaneous explanation possible!
Can also account for the AMM of the muon

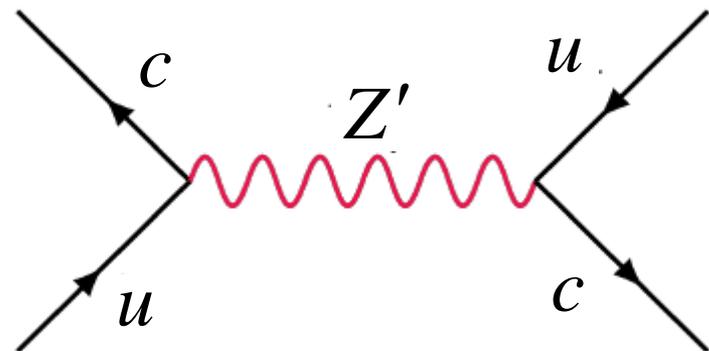
- $b \rightarrow c \tau \nu$



- $b \rightarrow s \mu \mu$

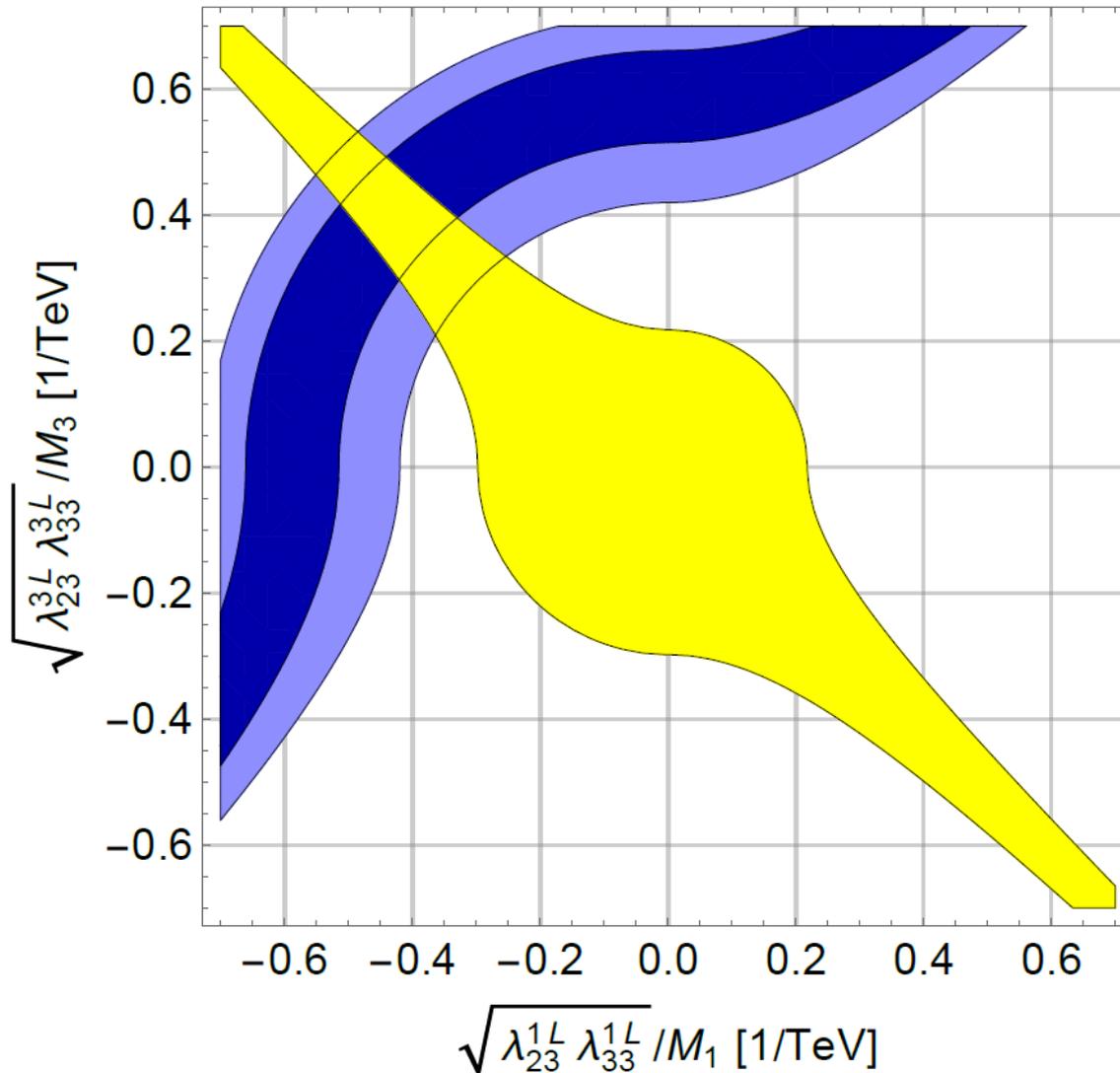


- $\tau \rightarrow \mu \mu \mu$ & D mixing



D mixing $\tau \rightarrow \mu \mu \mu$ and cannot be avoided

$R(D^{(*)})$, $b \rightarrow svv$ with 2 Scalar LQs



- $R(D^{(*)})$ 2σ
- $R(D^{(*)})$ 1σ
- $b \rightarrow svv$ allowed

$$\lambda_{jk}^L \equiv \lambda_{jk}^{1L}$$

$$\lambda_{jk}^{3L} = e^{i\pi j} \lambda_{jk}^L$$