Theory interpretations and challenges of DM searches

Felix Kahlhoefer LHCP 2018 Bologna 4-9 June 2018



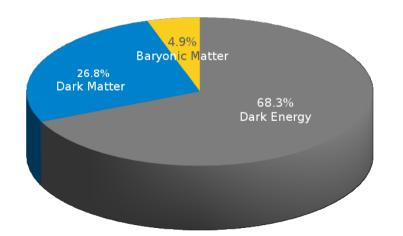


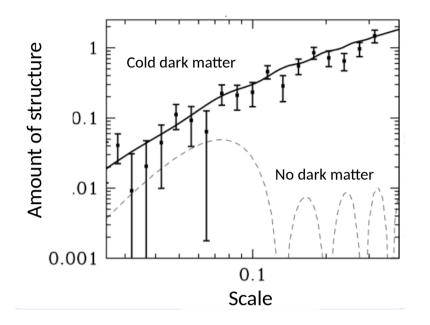


Why dark matter?

- Dark matter (DM) is an essential ingredient to describe Early Universe cosmology
 - Acts as the seed for **structure formation**
 - Creates the **potential wells** for stars and galaxies
- DM explains the amount and distribution of structure that we observe today

Note: Substantial evidence for DM also in the present Universe (galactic rotation curves, gravitational lensing...)





- A wealth of **successful predictions** from a very simple model
- Only draw-back: We understand only 5% of the Universe!

2



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Dark matter models

- **Particle physics** is the language of the early Universe
 - Example: Cosmology depends on the number, mass and interaction strength of neutrinos
- Likewise, we would like to understand DM in terms of particle physics
 - Require a **new stable particle** beyond the Standard Model
- At first sight, no obvious connection to LHC physics
 - There are viable DM models that are fundamentally unobservable at the LHC (too heavy, too weakly coupled)
 FK. arXiv:1801.07621
- Still, the LHC can probe many of the most attractive and most predictive DM models
 - Attractive: DM models connected to the hierarchy problem (e.g. SUSY)
 - **Predictive:** DM models where the observed relic abundance is obtained naturally







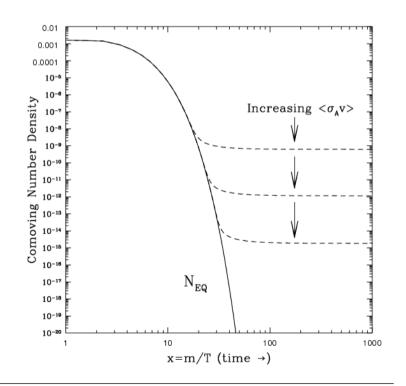


Dark matter production

• The one thing we know about dark matter is how much there is in the Universe:

 $\Omega h^2 = 0.1199 \pm 0.0027$

- Any model of dark matter must provide a mechanism to explain this number
- Most widely studied paradigm: Thermal freeze-out
 - Fundamental assumption: Thermal equilibrium between DM and other particles in the early Universe
 - New particles with weak interactions and weak-scale mass (so-called WIMPs) freeze out with the correct relic density
 - WIMPs can neither be too heavy nor too weakly coupled, so we can search for them at the LHC



4

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How does dark matter interact?

- Simplest assumption: DM particles couple directly to SM particles (e.g. gauge bosons or Higgs bosons)
 - Many **interesting models** (Higgs portal, minimal DM, ...)
 - Strong **experimental constraints** (e.g. from direct detection)
- What if there are **no direct interactions** between DM and SM particles?
 - Assume dark matter (DM) is uncharged under the Standard Model (SM) gauge group
 - Focus on models where the DM particle is exactly stable (e.g. due to a Z_2 symmetry)

Two possible directions:

Consider **non-renormalisable** interactions, i.e. higher-dimension effective operators Insist on **renormalisable interactions** but allow additional particles interacting with SM particles and DM (so-called **mediators**)









Effective interactions

- Very attractive idea: Parametrise interactions of DM in terms of unknown scale of new physics Λ
 - Finite number of possible operators
 - Simple parameter space and kinematics
 - Direct connection between different experiments
- In practice: Approach breaks down if Λ is too small
 - Need truncation procedure to restrict search to kinematic regions where effective field theory (EFT) is valid

Busoni et al., arXiv:1307.2253 Bruggisser et al., arXiv:1607.02475

Production

Annihilation

SM

SM

DM

Scattering

- Many interesting features are not captured by effective operators



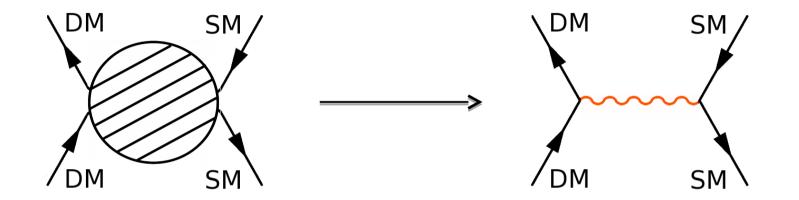






From EFTs to renormalisable models

 We can address the shortcomings of effective interactions by introducing a new mediating particle connecting DM to the SM



- The move to renormalisable models comes at a high price
 - Large number of possible models \rightarrow **loss of generality**
 - Large number of parameters for each model → increase of complexity







Why study renormalisable models?

- **Pragmatist's answer:** Renormalisable models remains valid in all kinematic regions
 - Straight-forward signal generation and calculation of NLO corrections
- Theorist's answer: We can study DM interactions at a more fundamental level
 - Comparison of different models in terms of their plausibility or complexity
- **Phenomenologist's answer:** Renormalisable models predict many new signatures, which can be correlated by exploiting the underlying structure of the theory
 - Understanding of relative importance of different search channels
 - Prediction of unexpected signals or kinematic distributions

Renormalisable models provide an excellent compromise between predictivity and flexibility







New mediators (lightning review)

- The most common extension involves new *s*-channel mediators, which can have
 - Spin 1:

$$egin{aligned} \mathcal{L}_{ ext{vector}} &= -g_{ ext{DM}} Z'_{\mu} ar{\chi} \gamma^{\mu} \chi - g_q \sum_{q=u,d,s,c,b,t} Z'_{\mu} ar{q} \gamma^{\mu} q \,, \ \mathcal{L}_{ ext{axial-vector}} &= -g_{ ext{DM}} Z'_{\mu} ar{\chi} \gamma^{\mu} \gamma_5 \chi - g_q \sum_{q=u,d,s,c,b,t} Z'_{\mu} ar{q} \gamma^{\mu} \gamma_5 q \,, \end{aligned}$$

- Spin 0:

$$egin{aligned} \mathcal{L}_{ ext{scalar}} &= -g_{ ext{DM}} \phi ar{\chi} \chi - g_q \, rac{\phi}{\sqrt{2}} \sum_{q=u,d,s,c,b,t} y_q ar{q} q \,, \ \mathcal{L}_{ ext{pseudo-scalar}} &= -i g_{ ext{DM}} \phi ar{\chi} \gamma_5 \chi - i g_q \, rac{\phi}{\sqrt{2}} \sum_{q=u,d,s,c,b,t} y_q ar{q} \gamma_5 q \,. \end{aligned}$$

• All these simplified models appear renormalisable (all couplings are dimensionless), but there are a lot of hidden complications

9





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10





Renormalisable models of spin-1 mediators

- Adding a massive spin-1 mediator to the SM yields a **non-renormalisable theory**
- Mediators with longitudinal polarisation violate unitarity at high energies

$$\langle Z'^{\mu}(k) Z'^{\nu}(-k) \rangle = \frac{1}{k^2 - m_{Z'}^2} \left(g^{\mu\nu} - \frac{k^{\mu}k^{\nu}}{m_{Z'}^2} \right)$$

- If the mediator has axial couplings, the longitudinal polarisations do not decouple FK et al., arXiv:1510.02110
- Solution: Generate mediator mass via the vev of an **additional Higgs boson**
 - Consistent theory requires two mediators, each of which can give rise to new signatures

Duerr, FK et al., arXiv:1606.07609

 To fully remove unitarity violation, coupling structure must respect gauge structure of unbroken electroweak symmetry
 Bell et al., arXiv:1512.00476

Bell et al., arXiv:1512.00476 Haisch, FK et al., arXiv:1603.01267







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Renormalisable models of spin-0 mediators

- Scalar mediators can only couple to SM quarks via **mixing** with the SM Higgs
 - Modification of **branching ratios** of the SM-like Higgs

Albert et al., arXiv:1607.06680 Bell et al., arXiv:1612.03475

- Scalar mediator picks up couplings to gauge bosons







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- Scalar mediator picks up couplings to gauge bosons
- Pseudoscalar mediators cannot mix with the SM Higgs (without CP violation)
 - Simplest extension requires a second Higgs doublet
 - The pseudoscalar mediator then mixes with the pseudoscalar component of the second Higgs doublet
 - New diagrams restore unitarity (e.g. in single-top production)

No. arXiv:1509.01110 Ipek et al., arXiv:1404.3716 Goncalves et al., arXiv:1611.04593 Bauer, Haisch & FK, arXiv:1701.07427

Pani & Polesello, arXiv:1712.03874



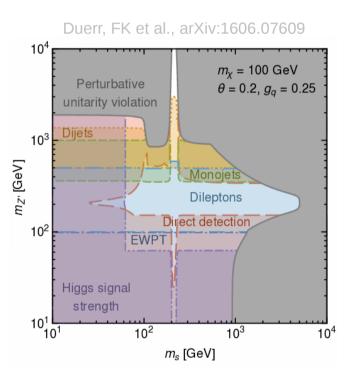
14





Central predictions of renormalisable models (I)

- In effective (or simplified) DM models, the properties (branching ratios, masses) of SM particles are not modified
- In consistent renormalisable models, **modifications occur copiously**:
 - Kinetic mixing and radiative corrections modify gauge boson masses (constrained by EWTP)
 - Higgs mixing reduces the overall Higgs signal strength
 - New decay modes lead to non-standard Higgs decays
- Moreover, we can observe new signatures in SM final states (dijet resonances, dilepton resonances, ...)





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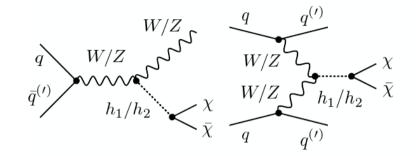


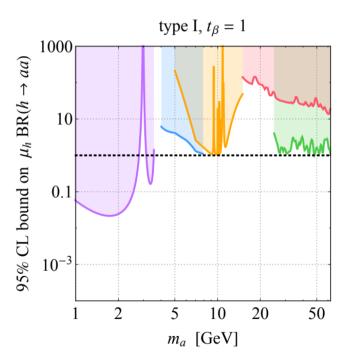


A closer look: Higgs decay modes

- Simplest example: Mixing between SM Higgs (*h*) and new "dark" Higgs boson (*s*)
 - **Invisible decay modes** open up if $m_{\text{DM}} < m_{\text{h}}/2$ or $m_s < m_{\text{h}}/2$ (and *s* long-lived)
 - Direct searches for invisible decays in VBF, h+V and h+j channels
 - Higgs signal strength reduced by factor cos²θ
 - Constraints from global Higgs fits
 - Exotic decay modes if s decays back into SM particles
 - Relevant constraints from searches for $h \rightarrow 4I$

See e.g. Haisch et al., arXiv:1802.02156











Central predictions of renormalisable models (II)

• In effective (or simplified) DM models, the signatures of DM production at the LHC stem from the initial state radiation of SM particles (mono-X signatures)

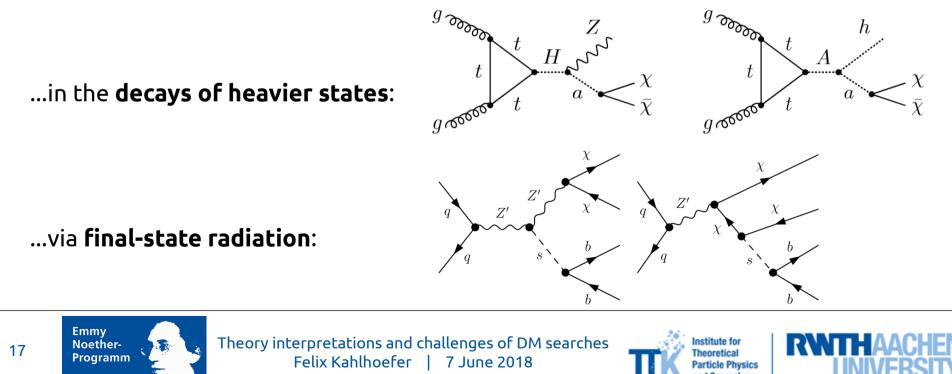
Bauer et al., arXiv:1712.06597

• One therefore expects a **very simple pattern**:

DFG

mono-jet > mono-photon > mono-Z/W > mono-Higgs

 In renormalisable models, SM particles can also be produced together with invisible particles...



Central predictions of renormalisable models (II)

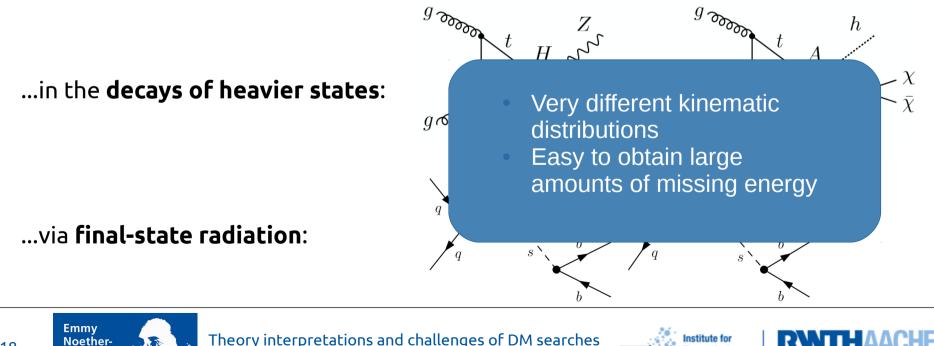
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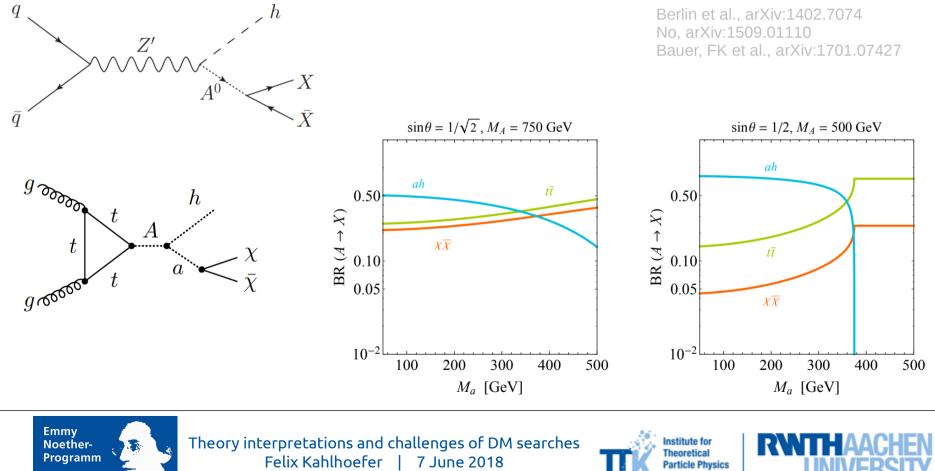


Programm

DFG

A closer look: mono-Higgs

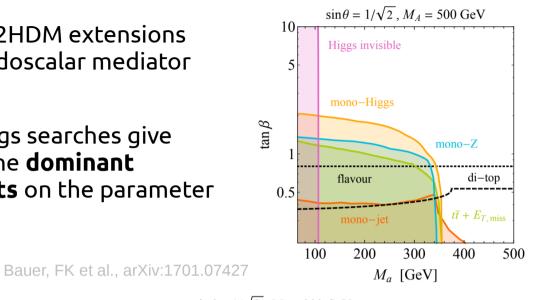
- Even if DM couples dominantly to heavy quarks, the cross section for mono-Higgs events from initial state radiation or from top-quark loops are completely negligible
- Much stronger signals if Higgs can be produced in **decays of a new heavy particle**

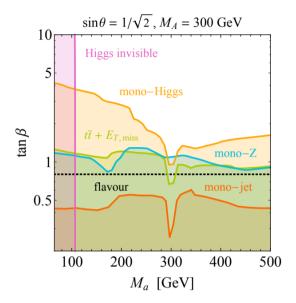


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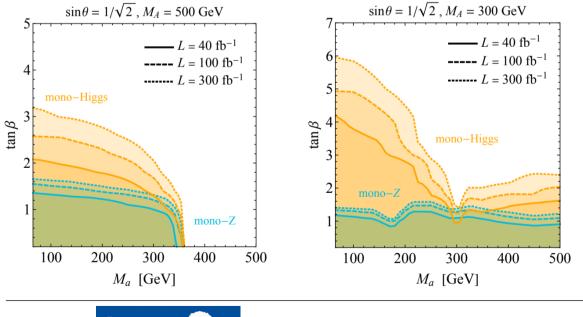
A closer look: mono-Higgs

- Example: 2HDM extensions with pseudoscalar mediator
- Mono-Higgs searches give some of the **dominant** constraints on the parameter space





- Sensitivity (at least for $h \rightarrow \gamma \gamma$) is still statistics limited
- Substantial room for improvement in future searches







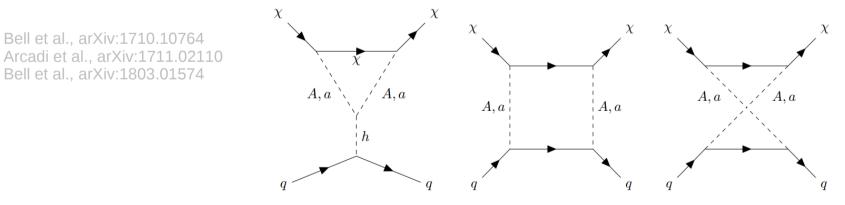
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Outlook

- Current focus of the attention: 2HDM + pseudoscalar mediator
 - Very rich model with many different signatures
 - Non-trivial constraints on the parameter space from EWPT and flavour physics
- More work needed to **understand the allowed regions** of parameter space
 - LHC DM working group: Detailed study of search channels and kinematic distributions
 - To appear soon: Whitepaper with recommended benchmark models and parameter grids
- Ongoing effort: Study direct detection and relic density constraints, including loop effects











22

Emmy **DFG**

Instead of conclusions From DM models to dark sector models

- Renormalisable DM models typically require **several new particles** in addition to the DM particle, which cannot be arbitrarily heavy
- These new particles can give rise to a **wide range of exotic signatures**:
 - Missing energy + exotic resonances (mono-Z', mono-dark-Higgs, ...)
 - Missing energy + long-lived particles
 - Emerging (or semi-visible) jets

Duerr, FK et al., arXiv:1701.08780

Buchmueller et al., arXiv:1704.06515

Many unexplored signatures – much work remains to be done!

