

# Extra Dimensions Working Group Report

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# Outline

## ① Randall-Sundrum model

# Introduction

- An area that has been actively researched into in this working group has been Bulk Randall-Sundrum Models.
- It is important to go through the basics of the model before discussing the projects completed.

## Bulk RS Model basics

- Start with the action for a 5-d scalar field  $\phi$  with a bulk mass  $ak$ , boundary mass term  $bk$ ,  $k$ : bulk curvature.
- Separation of variables in the KK decomposition gives

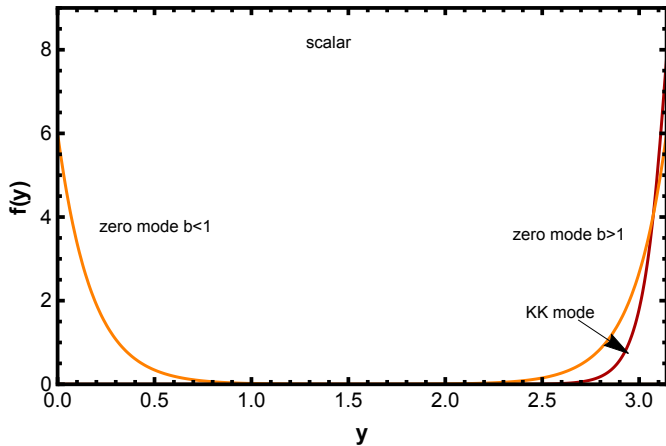
$$\phi(x^\mu, y) = \sum_{n=0}^{\infty} \phi^{(n)}(x^\mu) f_\phi^{(n)}(y) \quad (1)$$

- We can solve for the profiles  $f_\phi^{(n)}(y)$ .

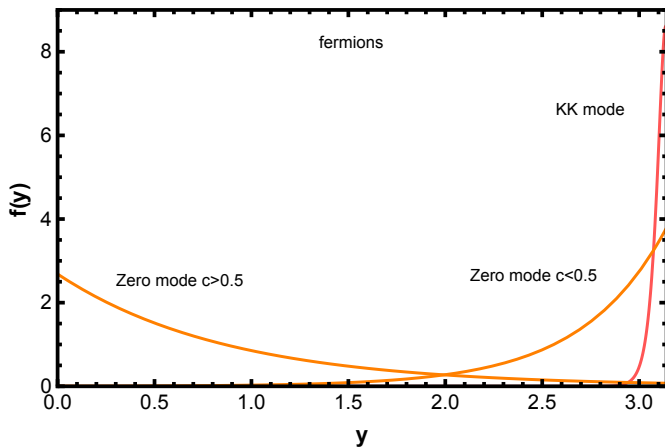
## Bulk RS Model basics II

- By tuning  $b$ , scalar zero-mode can be localised anywhere in the bulk.
- KK modes are given by combinations of Bessel functions and are localised towards the IR brane and have masses of the order of the infrared scale  $ke^{-\pi kR}$ .
- Follows similarly for fermion, gauge and graviton fields.

# Scalar modes

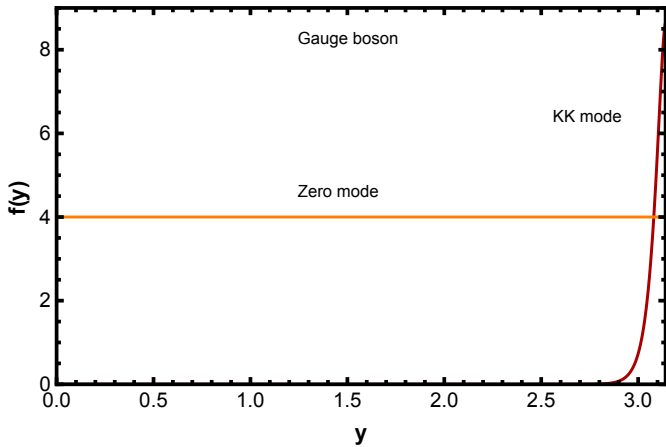


# Fermion modes





# Gauge modes



## AdS/CFT in a segment

- AdS/CFT works very differently in a segment of space-time i.e. between the UV and the IR branes.
- The UV brane in 5d shows up as a UV cut-off in the dual theory. Schwinger source field  $\phi_0$  becomes dynamical.
- The IR brane leads to a spontaneous breakdown of 4d-conformal symmetry, gives rise to CFT bound states.
- The mixing  $\phi_0\mathcal{O}$  between the source field (elementary) and the CFT bound states (composites) gives rise to the mass eigenstates which are from a 5-d perspective KK excitations of the bulk fields.

# Holographic Expansion

- Instead of a KK expansion of the bulk field, expand it in terms of the source field and CFT composites  $\rightarrow$  Holographic Expansion
- In general, there will be mixing but after diagonalisation and matching with the KK decomposition, one can identify the degree of compositeness of any bulk field.

## Elementary or composite?

- Graviton zero mode is UV localised and is pure elementary (tiny mixing with composites). KK modes are purely composite.
- Gauge field zero mode has a flat profile. So mixing is expected to be large. But coupling is marginal and the zero-modes are elementary. KK modes are purely composite.

## Elementary or composite? -II

- Light fermions like electrons ( $c > 1/2$ ) are UV localised and so elementary. For  $(t, b)_L$ ,  $c$  is  $0.3 - 0.4$ . Zero modes contain significant fraction of the composite CFT states though it is largely elementary. Massive modes are elementary
- Right handed top ( $c < -1/2$ ) has a large composite component.
- Higgs is localised on or close to the brane so fully composite or very nearly so.

# Custodial model

- Bulk RS models have large custodial symmetry violations because of interactions between Higgs and KK excitations of gauge bosons.
- Leads to large effects in EW parameters.
- Gauge KK modes mass  $\rightarrow$  15 -20 TeV.
- Custodial models extend the gauge group to  $SU(3)_c \times SU(2)_L \times SU(2)_R \times U(1)_{B-L}$ . Reduces bound to 3 TeV on gauge boson masses.

## Deformed Metric

- A metric singular at  $y = y_s = L + \Delta$  :

$$A(y) = ky - \frac{1}{\nu^2} \log\left(1 - \frac{y}{y_s}\right) \quad (2)$$

- Becomes the usual R-S metric for  $\nu \rightarrow \infty$ .
- The singularity at  $y = y_s$  deforms the metric near the TeV brane.
- The Higgs is also moved into the bulk and alongwith the deformed metric keeps  $T$  under control.

## KK gluons

- Interesting signal – KK gluon production.
- The KK gluon coupling to  $t_R$  is enhanced by a factor  $\xi$  compared to  $\alpha_s$  where  $\xi \equiv \sqrt{\log(M_{pl}/\text{TeV})} \sim 5$ .
- Consequently, it decays predominantly to tops if produced.
- To the  $(t, b)_L$  doublet its coupling is  $\alpha_s$  and to the light quarks its couplings are suppressed.



# KK Gluon Production

- However, the  $ggg_{KK}$  vanishes because of the the orthogonality of the profiles of these particles.
- So  $gg$  initial state does not contribute, only  $q\bar{q}$  does.
- At the LHC, KK gluon masses of the order of 2 TeV can be probed with this process.

# Associated Jets Production

- It is useful to consider other production mechanisms:  $g_{KK} + \text{jets}$  is interesting because  $qg$  and  $gg$ -initial states contribute.
- Hadron-level simulation of the signal and background shows that the jet-shape variable N-subjettiness can be very effective in suppressing the QCD background.
- The irreducible  $t\bar{t}$  background can be suppressed using transverse momentum cuts.

*A.M. Iyer, F. Mahmoudi, N. Manglani, K. Sridhar, Phys.Lett. B759 (2016) 342-348.*

# KK gluon + jet: $p_T$ distributions

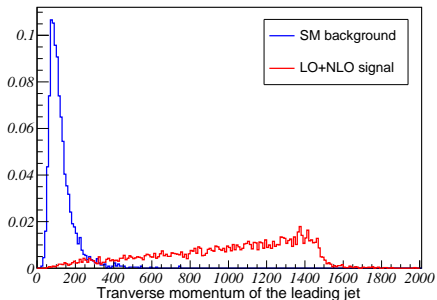


Figure:  $p_T$  of the leading jet

## KK gluon + jet: Cross-section

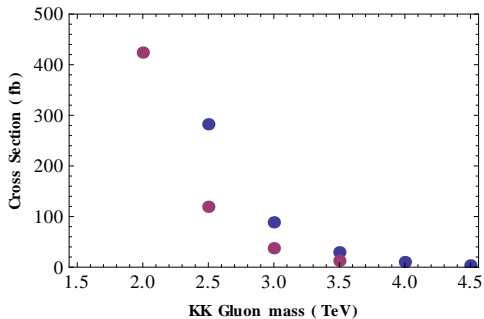


Figure: Cross-section as a function of  $g_{KK}$  mass

## Higgs KK mode: Custodial

- We have shown that in Bulk Higgs models, the Higgs KK modes can be light consistent with naturalness, electroweak precision and Higgs mass constraint.
- The first KK mode of the Higgs ( $h_1$ ) could lie in the 1-2 TeV range in the custodial model.
- $h_1$  decays dominantly into a boosted  $t\bar{t}$  pair.
- HEPTopTagger is efficient in taming the large QCD background.
- The irreducible  $t\bar{t}$  background can be controlled by using a cut on  $\Delta\eta$  and  $p_T$ .

*F. Mahmoudi, U. Maitra, N. Manglani, K. Sridhar, JHEP 1611 (2016) 075.*

## Higgs KK mode: Deformed

- The first KK mode in the deformed model is even lighter (800 GeV - 1.2 TeV) but the cross-sections are small.
- To handle the QCD background, the HEPToptagger needs to be supplemented with  $b$ -tagging inside the top jet.
- We need to go to the high-luminosity option of the LHC to be able to fully study the  $h_1$ -mass range.

*F. Mahmoudi, N. Manglani, K. Sridhar. arXiv:1712.04966*

# Vector-Like Quarks

- One novel way of searching for exotic vector-like quarks is to look for their effects in electroweak loops.
- We have begun the study of  $gg \rightarrow ZZ/\gamma\gamma$  through a box involving a vector-like quark.

*G. Cacciapaglia, D. Choudhury, N. Gaur, K. Sridhar – Work in progress*