

Status of V_{us} determination from Kaon decays

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Outline :

1. Introduction and Motivation
2. V_{us} from $K_{\ell 3}$ decays
3. V_{us}/V_{ud} from $K_{\ell 2}/\pi_{\ell 2}$ decays
4. V_{us} and Unitarity of the CKM matrix
5. Conclusion and outlook

1. Introduction and Motivation

1.1 Test of the Standard Model: V_{us} and CKM unitarity

- Extraction of the Cabibbo-Kobayashi-Maskawa matrix element V_{us}
 - Fundamental parameter of the Standard Model

Description of the **weak interactions**:

$$\mathcal{L}_{EW} = \frac{g}{\sqrt{2}} W_{\alpha}^{+} \left(\bar{D}_L V_{CKM} \gamma^{\alpha} U_L + \bar{e}_L \gamma^{\alpha} \nu_{e_L} + \bar{\mu}_L \gamma^{\alpha} \nu_{\mu_L} + \bar{\tau}_L \gamma^{\alpha} \nu_{\tau_L} \right) + \text{h.c.}$$

	I	II	III		
Quarks	u	c	t	γ	H
	d	s	b	g	
	ν_e	ν_{μ}	ν_{τ}	Z	
Leptons	e	μ	τ	W	
	3 generations			Forces	Higgs

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Unitary
matrix

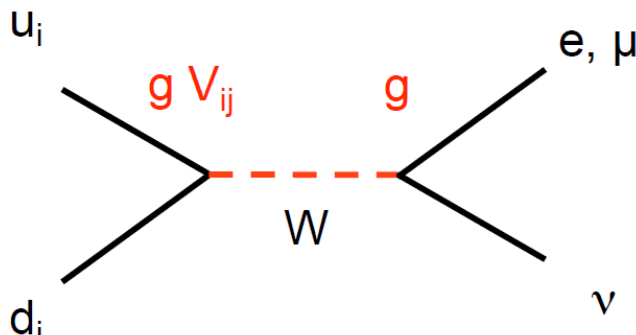
- Check unitarity of the first row of the CKM matrix:



Cabibbo Universality:

$$|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 1$$

Negligible $\sim 2 \times 10^{-5}$
(B decays)



	I	II	III		
Leptons	u	c	t	γ	H
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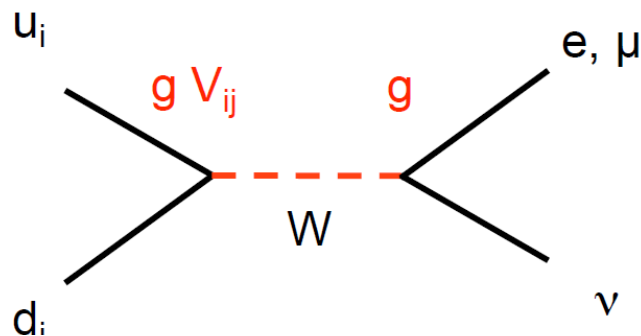
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Gauge coupling

- Universality: Is G_F from μ decay equals to G_F from π , K, nuclear β decay?

$$G_{\mu}^2 = (g_{\mu} g_e)^2 / M_W^4 \stackrel{?}{=} G_{CKM}^2 = (g_q g_{\ell})^2 (|V_{ud}|^2 + |V_{us}|^2) / M_W^4$$



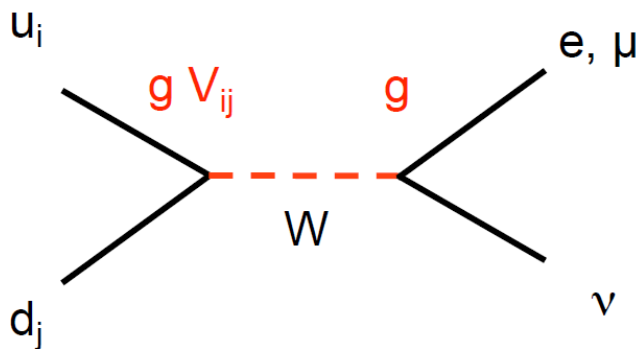
1.2 Constraining New Physics

- Extraction of the Cabibbo-Kobayashi-Maskawa matrix element V_{us}
 - Fundamental parameter of the Standard Model

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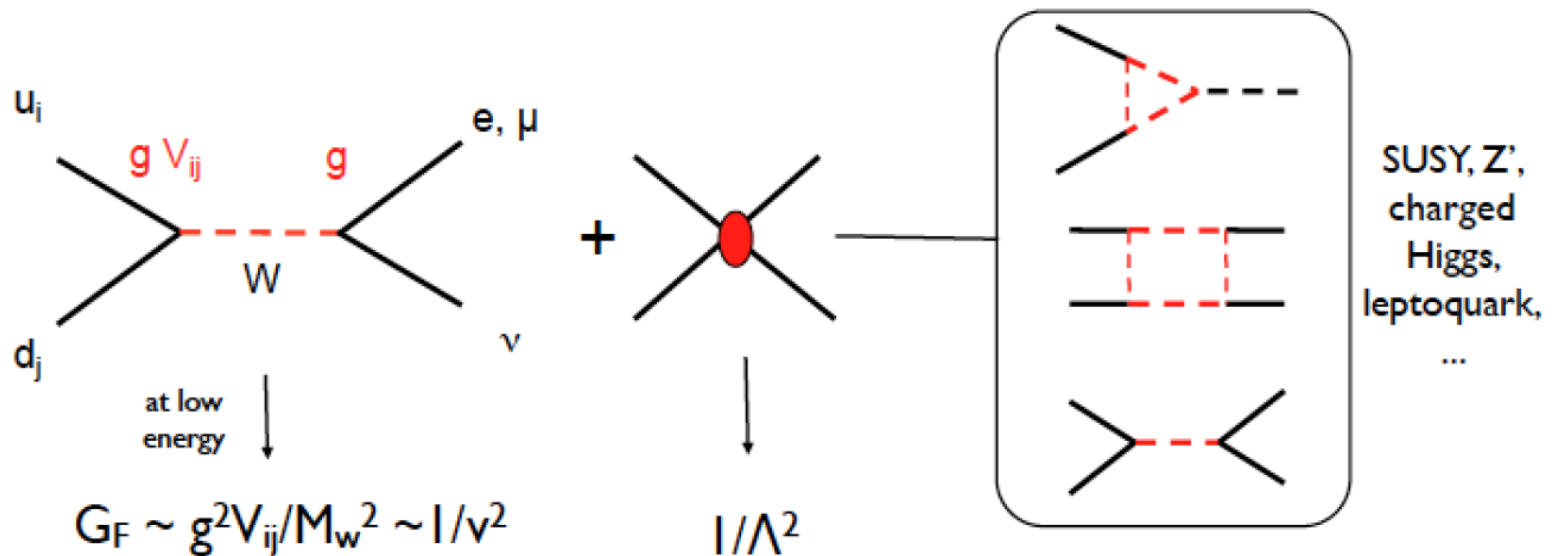
- Look for **new physics**
 - In the Standard Model : W exchange ➡ only V-A structure



1.2 Constraining New Physics

- BSM: sensitive to tree-level and loop effects of a large class of models

➔ $|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 1 + \Delta_{CKM}$

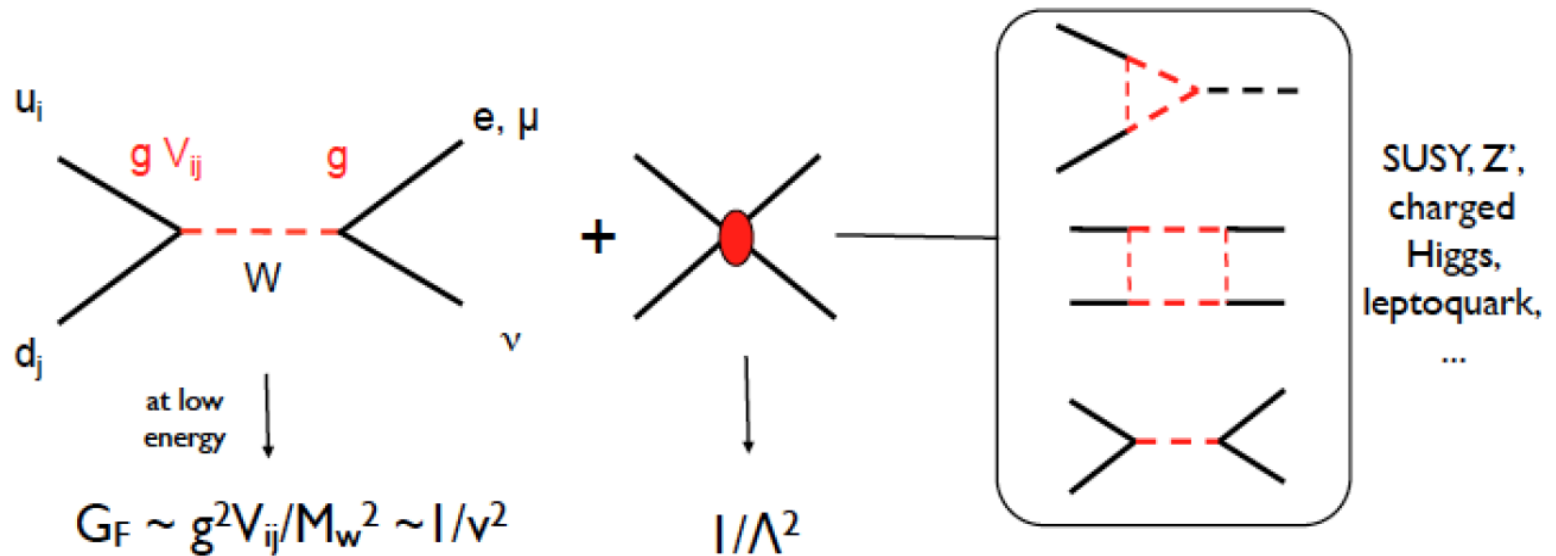


➔ BSM effects : $\Delta \sim \frac{c_n}{g^2} \frac{M_W^2}{\Lambda^2} \leq 10^{-2} - 10^{-3} \longleftrightarrow \Lambda \sim 1-10 \text{ TeV}$

1.2 Constraining New Physics

- BSM: sensitive to tree-level and loop effects of a large class of models

➔ $|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 1 + \Delta_{CKM}$



- Look for new physics by comparing the extraction of V_{us} from different processes: helicity suppressed $K_{\mu 2}$, helicity allowed $K_{l 3}$, hadronic τ decays

1.3 Some history

- 2002: Old $K \rightarrow \pi l \nu_l$ data give $\Delta_{CKM} = 1 - |V_{ud}|^2 - |V_{us}|^2 = 0.0035(15)$
 - ➡ PDG 2004: a 2.3σ hint of *unitarity violation*?
- 2003 BNL 865 measures $BR(K^+ \rightarrow \pi^0 e^+ \nu) = 5.13(10)\%$
 - ➡ value of V_{us} consistent with unitarity
- 2004 – present: Many new measurements from **KTeV**, **ISTRA+**, **KLOE**, **NA48**
 - BRs, lifetimes, form-factors
 - Much higher statistics than older measurements
 - Proper account of correlations between measurements
 - ➡ Isospin breaking, radiative corrections start to matter: computed within ChPT
- 2008 – beyond: Progress in the computation of hadronic elements from lattice QCD
- Value of V_{us} used in precise test of the SM

1.4 Experiment, Theory & Evaluation

V_{us} from $K_{\ell 3}$ & $K_{\ell 2}$ { ~ 100 measurements of ~ 10 experimental parameters
50+ (and counting!) lattice results for 2 hadronic matrix elements
Radiative and SU(2)-breaking corrections, ChPT results, etc.

FlaviA
net **Kaon WG**
2006-2010 (EU 6FP)

Experimental averages, fits, etc

Selection of results (experiments, corrections)

Evaluation, discussion and interpretation

Final report: EPJC 69 (2010) 399

This talk is an attempt at an update to 2018

Corresponding effort to synthesize results from **lattice QCD**:

**Flavor Lattice
Averaging Group
(FLAG):**

<http://itpwiki.unibe.ch/flag>

Participation by all major lattice collaborations

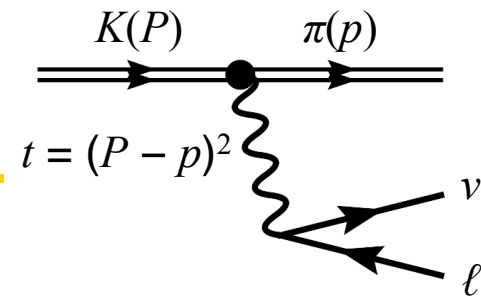
Biannual review of lattice results for π , K , B , D physics

2013 review: EPJC 74 (2014) 2890

2016 review: EPJC77 (2017) no.2, 112

2. V_{us} from $K_{\ell 3}$ decays

2.1 V_{us} from K_{l3} decays



- Master formula for $K \rightarrow \pi l \nu$: $K = \{K^+, K^0\}$, $l = \{e, \mu\}$

$$\Gamma(K \rightarrow \pi l \nu [\gamma]) = Br(K_{l3}) * \tau = C_K^2 \frac{G_F^2 m_K^5}{192 \pi^3} S_{EW}^K |V_{us}|^2 \left| f_+^{K^0 \pi^-}(0) \right|^2 I_{KI} \left(1 + 2\Delta_{EM}^{KI} + 2\Delta_{SU(2)}^{K\pi} \right)$$

Experimental inputs:

$\Gamma(K_{l3})$ Rates with well-determined treatment of radiative decays

- Branching ratios
- Kaon lifetimes

$I_{KI}(\lambda_{KI})$ Integral of form factor over phase space: λ s parametrize evolution in $t=q^2$

Inputs from theory:

S_{EW}^K Universal short distance EW corrections

$f_+^{K^0 \pi^-}(0)$ Hadronic matrix element (form factor) at zero momentum transfer ($t=0$)

Δ_{EM}^{KI} Form-factor correction for long-distance EM effects

$\Delta_{SU(2)}^{K\pi}$ Form-factor correction for SU(2) breaking

2.2 Modern experimental data for V_{us} from $K_{\ell 3}$

Experiment	Measurement	Year
BNL865	$\text{BR}(K^+ \rightarrow \pi^0_{\text{D}} e^+ \nu) / \text{BR}(K^+ \rightarrow \pi^0_{\text{D}} X^+)$	2003
KTeV	$\tau(K_S)$	2003
	$\text{BR}(K_{Le3}), \text{BR}(K_{L\mu 3}), \lambda_+(K_{Le3}), \lambda_{+,0}(K_{L\mu 3})$	2004
ISTRA+	$\lambda_+(K^-_{e3}), \lambda_{+,0}(K^-_{e3})$	2004
KLOE	$\tau(K_L)$	2005
	$\text{BR}(K_{Le3}), \text{BR}(K_{L\mu 3}), \text{BR}(K_{Se3}), \lambda_+(K_{Le3})$	2006
	$\lambda_{+,0}(K_{L\mu 3})$	2007
	$\tau(K^\pm), \text{BR}(K_{Le3}), \text{BR}(K_{L\mu 3})$	2008
NA48	$\tau(K_S)$	2002
	$\text{BR}(K_{Le3}/2 \text{ tracks}), \lambda_+(K_{Le3})$	2004
	$\Gamma(K_{Se3}/K_{Le3}), \lambda_{+,0}(K_{L\mu 3})$	2007
NA48/2	$\text{BR}(K^+_{e3}/\pi^+\pi^0), \text{BR}(K^+_{\mu 3}/\pi^+\pi^0)$	2007

Above data set used for 2010 FlaviaNet review (fits, averages, etc.)

Updated fit to K_S rate data

6 input measurements:

KLOE BR $\pi^0\pi^0/\pi^+\pi^-$

KLOE BR $\pi e\nu/\pi^+\pi^-$

NA48 $\Gamma(K_S \rightarrow \pi e\nu)/\Gamma(K_L \rightarrow \pi e\nu)$, τ_S

KLOE'11 τ_S

KTeV'11 τ_S

2 constraints:

- $\Sigma \text{BR} = 1$
- $\text{BR}(K_{e3})/\text{BR}(K_{\mu3}) = 0.66492(137)$

From ratio of phase-space integrals from current fit to dispersive $K_{\ell3}$ form factor parameters

Parameter	Value
$\text{BR}(\pi^+\pi^-(\gamma))$	69.20(5)%
$\text{BR}(\pi^0\pi^0)$	30.69(5)%
$\text{BR}(K_{e3})$	$7.05(8) \times 10^{-4}$
$\text{BR}(K_{\mu3})$	$4.69(6) \times 10^{-4}$
τ_S	89.58(4) ps

$\chi^2/\text{ndf} = 0.20/3$ (Prob = 98%)

$\rho(\text{BR}(\pi^+\pi^-), \text{BR}(\pi^0\pi^0)) = -0.998$

Little freedom in fit

Largest effect of **2011 τ_S data**:

FlaviaNet 2010
 $\tau_S = 89.59(6)$ ps



Update
 $\tau_S = 89.58(4)$ ps

Updated fit to K_L rate data

21 input measurements:

5 KTeV ratios

NA48 $\text{BR}(K_{e3}/2 \text{ track})$

4 KLOE BRs

with dependence on τ_L

KLOE, NA48 $\text{BR}(\pi^+\pi^-/K_{\ell 3})$

KLOE, NA48 $\text{BR}(\gamma\gamma/3\pi^0)$

$\text{BR}(2\pi^0/\pi^+\pi^-)$ from K_S fit, $\text{Re } \varepsilon'/\varepsilon$

KLOE τ_L from $3\pi^0$

Vosburgh '72 τ_L

KTeV $\text{BR}(\pi^+\pi^-\gamma/\pi^+\pi^-(\gamma))$

E731, 2 KTeV $\text{BR}(\pi^+\pi^-\gamma_{\text{DE}}/\pi^+\pi^-\gamma)$

Parameter	Value	S
$\text{BR}(K_{e3})$	0.4056(9)	1.3
$\text{BR}(K_{\mu 3})$	0.2704(10)	1.5
$\text{BR}(3\pi^0)$	0.1952(9)	1.2
$\text{BR}(\pi^+\pi^-\pi^0)$	0.1254(6)	1.3
$\text{BR}(\pi^+\pi^-(\gamma_{\text{IB}}))$	$1.967(7) \times 10^{-3}$	1.1
$\text{BR}(\pi^+\pi^-\gamma)$	$4.15(9) \times 10^{-5}$	1.6
$\text{BR}(\pi^+\pi^-\gamma_{\text{DE}})$	$2.84(8) \times 10^{-5}$	1.3
$\text{BR}(2\pi^0)$	$8.65(4) \times 10^{-4}$	1.4
$\text{BR}(\gamma\gamma)$	$5.47(4) \times 10^{-4}$	1.1
τ_L	51.16(21) ns	1.1

$\chi^2/\text{ndf} = 19.8/12$ (Prob = 7.0%)

Essentially same result as 2010 fit

Current PDG ('09): 37.4/17 (0.30%)

1 constraint: $\Sigma \text{BR} = 1$

Updates: K^\pm BRs and lifetimes

KLOE-2
PLB 738 (2014)

$$\text{BR}(\pi^+\pi^+\pi^-) = 0.05565(31)(25) \quad (0.7\%)$$

- **No good measurements of $\text{BR}(\pi^+\pi^+\pi^-)$ in 2010 fit**
- Reconstruct 2 tracks in small fiducial volume near interaction region; evaluate missing mass for 3rd track
- Fully inclusive of radiation, but radiative corrections handled differently from other KLOE measurements
- Significant impact on value of $\text{BR}(\mu\nu)$ **from fit**
Correlation between $\text{BR}(\mu\nu)$, $\text{BR}(\pi^+\pi^+\pi^-) = -0.75$

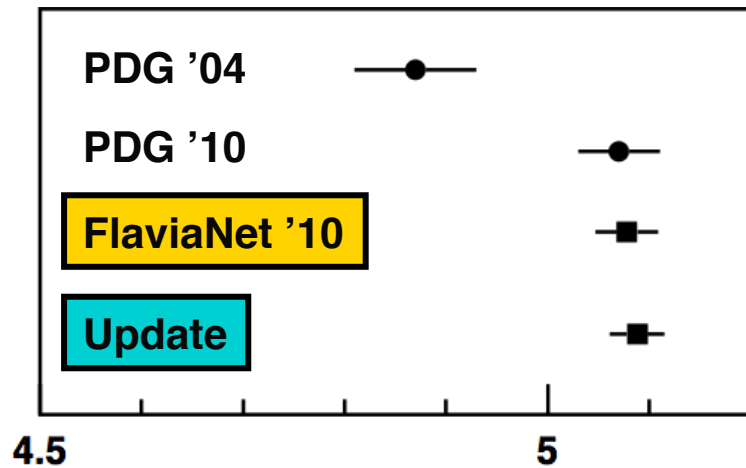
ISTRA+
PAN 77 (2014)

$$\text{BR}(K_{e3}^-/\pi^-\pi^0) = 0.2423(15)(37) \quad (1.6\%)$$

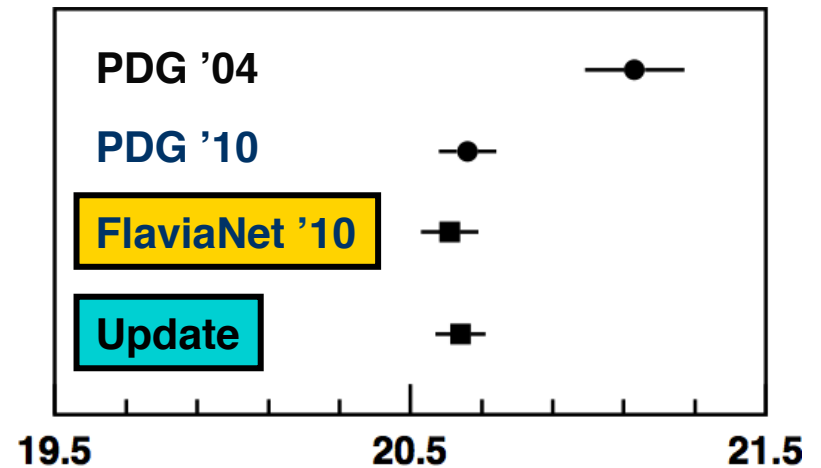
- Claimed to be fully inclusive for $K_{e3\gamma}$
 - No mention of radiative corrections
 - Many cuts, mainly topological
 - 3 different selections, at least 1 may be largely inclusive
- Included in PDG '15 fit
- **Treated as preliminary here (not in K^\pm BR fit)**

Updates: K^\pm BRs and lifetimes

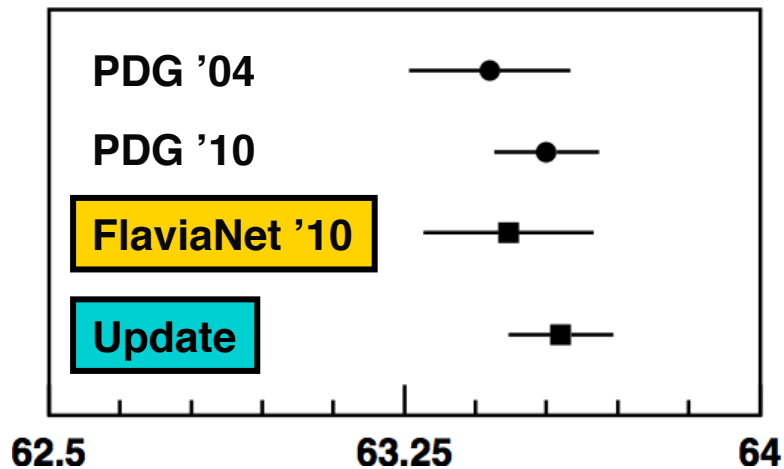
$\text{BR}(K^\pm \rightarrow \pi^0 e \nu)$



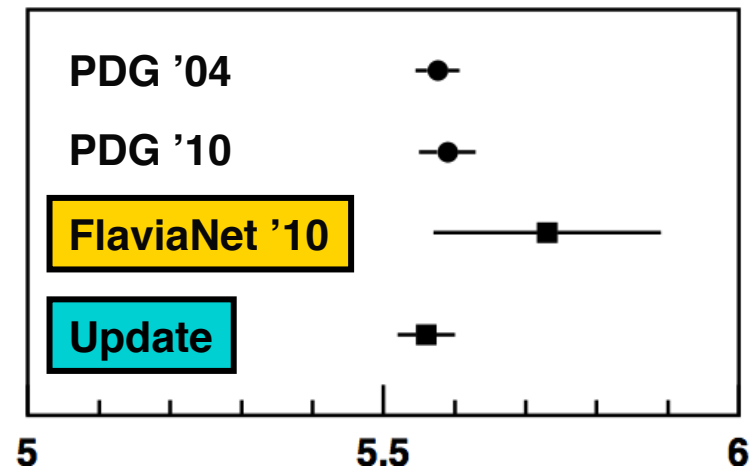
$\text{BR}(K^\pm \rightarrow \pi \pi^0)$



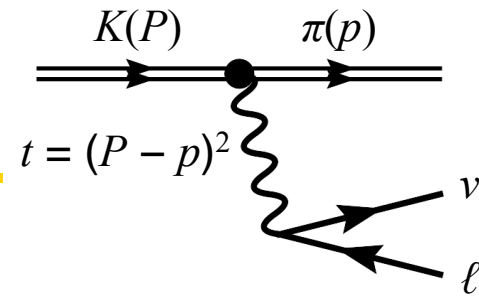
$\text{BR}(K^\pm \rightarrow \mu \nu)$



$\text{BR}(K^\pm \rightarrow \pi \pi \pi)$



2.3 Electroweak corrections



- Master formula for $K \rightarrow \pi l \nu_l$: $K = \{K^+, K^0\}$, $l = \{e, \mu\}$

$$\Gamma(K \rightarrow \pi l \nu [\gamma]) = Br(K_{l3}) * \tau = C_K^2 \frac{G_F^2 m_K^5}{192 \pi^3} \boxed{S_{EW}^K} |V_{us}|^2 \left| f_+^{K^0 \pi^-}(0) \right|^2 I_{KI} \left(1 + 2 \Delta_{EM}^{KI} + 2 \Delta_{SU(2)}^{K\pi} \right)$$

- S_{ew} : Short distance electroweak correction

$$\boxed{S_{ew} = 1 + \frac{2\alpha}{\pi} \left(1 + \frac{\alpha_s}{4\pi} \right) \log \frac{m_Z}{m_\rho} + O\left(\frac{\alpha\alpha_s}{\pi^2}\right)} \Rightarrow \boxed{S_{ew} = 1.0232(3)}$$

Sirlin'82

$$\text{Tree level} + \text{One loop} + \dots \Rightarrow G_\mu V_{ij} \sqrt{S_{ew}}$$

2.4 $K\pi$ form factors

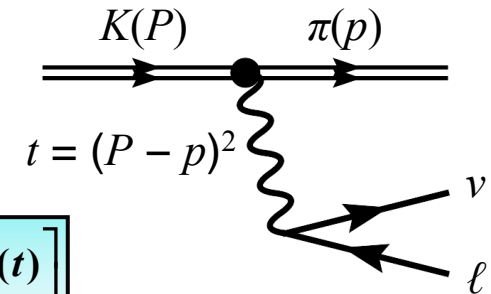
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- $f_+(0)$: vector form factor at zero momentum transfer:

Hadronic matrix element:

$$\langle \pi^-(p) | \bar{s} \gamma_\mu u | K^0(P) \rangle = f_+^{K^0 \pi^-}(0) \left[(P + p)_\mu \bar{f}_+^{K^0 \pi^-}(t) + (P - p)_\mu \bar{f}_-^{K^0 \pi^-}(t) \right]$$



$f_+(0)$ key hadronic quantity: In $SU(3)_V$ limit ($m_u = m_d = m_s$), CVC $\Rightarrow f_+(0) = 1$
 Need to compute corrections in second order in $SU(3)$ breaking

\Rightarrow see later

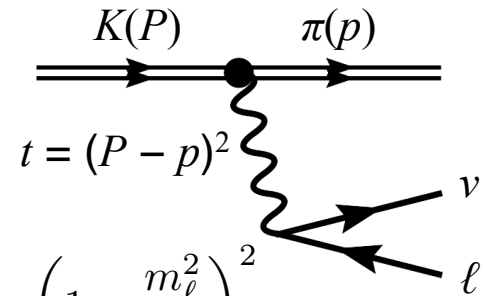
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Hadronic matrix element:

$$\langle \pi^-(p) | \bar{s} \gamma_\mu u | K^0(P) \rangle = f_+^{K^0 \pi^-}(0) \left[(P+p)_\mu \bar{f}_+^{K^0 \pi^-}(t) + (P-p)_\mu \bar{f}_-^{K^0 \pi^-}(t) \right]$$



- Phase space integrals:
$$I_{K\ell} = \frac{2}{3} \int_{m_\ell^2}^{t_0} \frac{dt}{M_K^8} \bar{\lambda}^{3/2} \left(1 + \frac{m_\ell^2}{2t}\right) \left(1 - \frac{m_\ell^2}{2t}\right)^2 \times \left(\bar{f}_+^2(t) + \frac{3m_\ell^2 \Delta_{K\pi}^2}{(2t + m_\ell^2) \bar{\lambda}} \bar{f}_0^2(t) \right),$$
- In K_{e3} decays: only vector FF $\bar{f}_+^{K^0 \pi^-}(t)$
- In $K_{\mu 3}$ decays, also need the scalar FF
$$\bar{f}_0(t) = \bar{f}_+(t) + \frac{t}{m_K^2 - m_\pi^2} \bar{f}_-(t)$$
- For V_{us} , need integral over phase space of squared matrix element: Parameterize form factors and fit distributions in t (or related variables)

$K_{\ell 3}$ form factors parametrizations

- Parametrizations based on Taylor expansion:

$$\bar{f}_{+,0}(t) = 1 + \lambda_{+,0} \left(\frac{t}{m_{\pi^\pm}^2} \right) \quad \text{or} \quad \bar{f}_{+,0}(t) = 1 + \lambda'_{+,0} \left(\frac{t}{m_{\pi^\pm}^2} \right) + \lambda''_{+,0} \left(\frac{t}{m_{\pi^\pm}^2} \right)^2$$

Very simple parametrization but limited in energy range and not physically motivated: many parameters and strong correlations between them

➡ unstable fits

- Physically motivated parametrizations:
 - Pole parametrization

$$\bar{f}_{+,0}(t) = \left(\frac{M_{V,S}^2}{M_{V,S}^2 - t} \right)$$

Well motivated for the vector (K^* resonance)
But for the scalar M_S ?

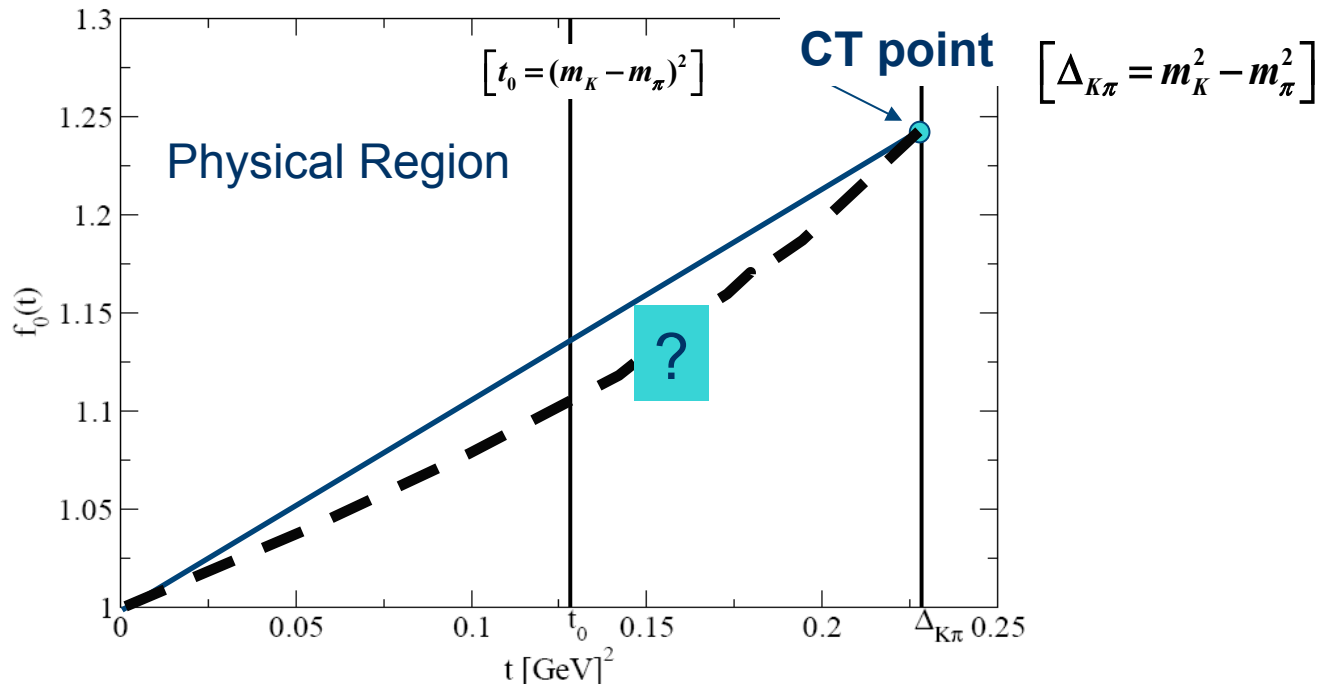
- Dispersive parametrization

Bernard, Oertel, E.P., Stern'06,'09

$$\bar{f}_+(t) = \exp \left[\frac{t}{m_\pi^2} \left(\Lambda_+ - H(t) \right) \right] \quad \text{and} \quad \bar{f}_0(t) = \exp \left[\frac{t}{m_K^2 - m_\pi^2} \left(\ln C - G(t) \right) \right]$$

Dispersive representation for the form factors

- Take the $K\pi$ rescattering into account *Bernard, Oertel, E.P., Stern'06, '09*
- Allow to determine the slope and *curvature* of the form factors: only 2 param.



- Use the CT theorem for the scalar FF ➡ Write a twice subtracted dispersion relation for $\ln f(t)$ at $t=0$ and at the CT point for the scalar FF
- Does it improve the agreement with data?

2.4 Dispersive representation for the form factors

- Use dispersion relations to parametrize the FFs

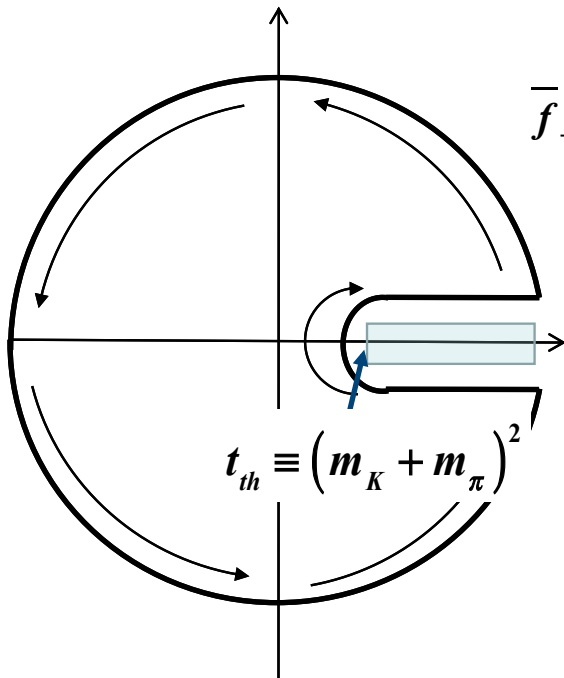
Unitarity: $\text{disc}[\bar{f}_{0,+}(s)] \propto t_\ell^{I*}(s) \bar{f}_{0,+}(s)$

Bernard, Oertel, E.P., Stern'06,'09

- Omnès representation:

$$\bar{f}_0(t) = \exp \left[\frac{t}{m_K^2 - m_\pi^2} \left(\ln C - \frac{\Delta_{K\pi}(\Delta_{K\pi} - t)}{\pi} \int_{t_{th}}^{\infty} \frac{ds}{s} \frac{\phi_0(s)}{(s - \Delta_{K\pi})(s - t - i\epsilon)} \right) \right]$$

$$\bar{f}_+(t) = \exp \left[\frac{t}{m_\pi^2} \left(\Lambda_+ + \frac{m_\pi^2 t}{\pi} \int_{t_{th}}^{\infty} \frac{ds}{s^2} \frac{\phi_+(s)}{(s - t - i\epsilon)} \right) \right]$$



$\phi_{+,0}(s)$: phase of the form factor

- $s < s_{in}$: $\phi_{+,0}(s) = \delta_{K\pi}(s)$

\swarrow
K π scattering phase

- $s \geq s_{in}$: $\phi_{+,0}(s)$ unknown

$\Rightarrow \phi_{+,0}(s) = \phi_{+,0as}(s) = \pi \pm \pi \quad (\bar{f}_{+,0}(s) \rightarrow 1/s)$

Brodsky & Lepage

- A large error turns out in a small uncertainty in the physical region

Callan-Treiman Low Energy Theorem

- Callan-Treiman theorem:

Bernard, Oertel, E.P., Stern'06, '08

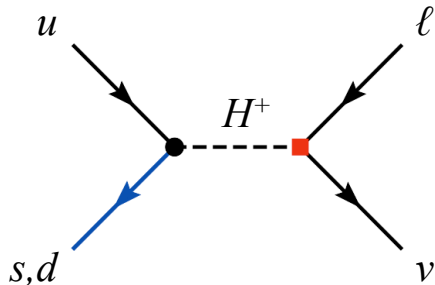
$$C = \overline{f}_0(\Delta_{K\pi}) = \frac{F_K}{F_\pi f_+(0)} + \Delta_{CT} = \underbrace{\frac{F_K |V^{us}|}{F_\pi |V^{ud}|} \frac{1}{f_+(0) |V^{us}|} |V^{ud}|}_{\text{Very precisely known from Br(Kl2/\pi l2), } \Gamma(\text{Ke3}) \text{ and } |V_{ud}|} r + \Delta_{CT}$$

\nearrow $m_K^2 - m_\pi^2$

$$B_{\text{exp}} = 1.2446(41)$$

- In the Standard Model : $r = 1$ $(\ln C_{SM} = 0.2141(73))$ $\Delta_{CT} = (-3.5 \pm 8) \cdot 10^{-3}$
 NLO value + large error bars in agreement with *Bijnens&Ghorbani'07* *Kastner & Neufeld'08*
- In presence of new physics, new couplings : $r \neq 1$

- Ex:



2.4 $K_{\ell 3}$ form factor data

- Form-factor parameter measurements in FlaviaNet 2010 fit:

K_L : **KTeV**, **KLOE**, **NA48** (K_{e3} only)

K^- : **ISTRA+**

- Even if not in the original publications, all experiments have:
 - Obtained results for Taylor, pole, and dispersive parameterizations
 - Supplied parameter correlation coefficients

New measurements:

NA48/2
1808.09041

$2.3 \times 10^6 K_{\mu 3}^{\pm}$
 $4.4 \times 10^6 K_{e3}^{\pm}$

Updating 2012 preliminary
See talk by M. Piccini

K^+ and K^- simultaneously acquired in dedicated minimum-bias run

Taylor, pole, and dispersive fits with complete investigation of systematics

OKA
JETPL 107 (2018)

$5.25 \times 10^6 K_{e3}^+$

Described as preliminary

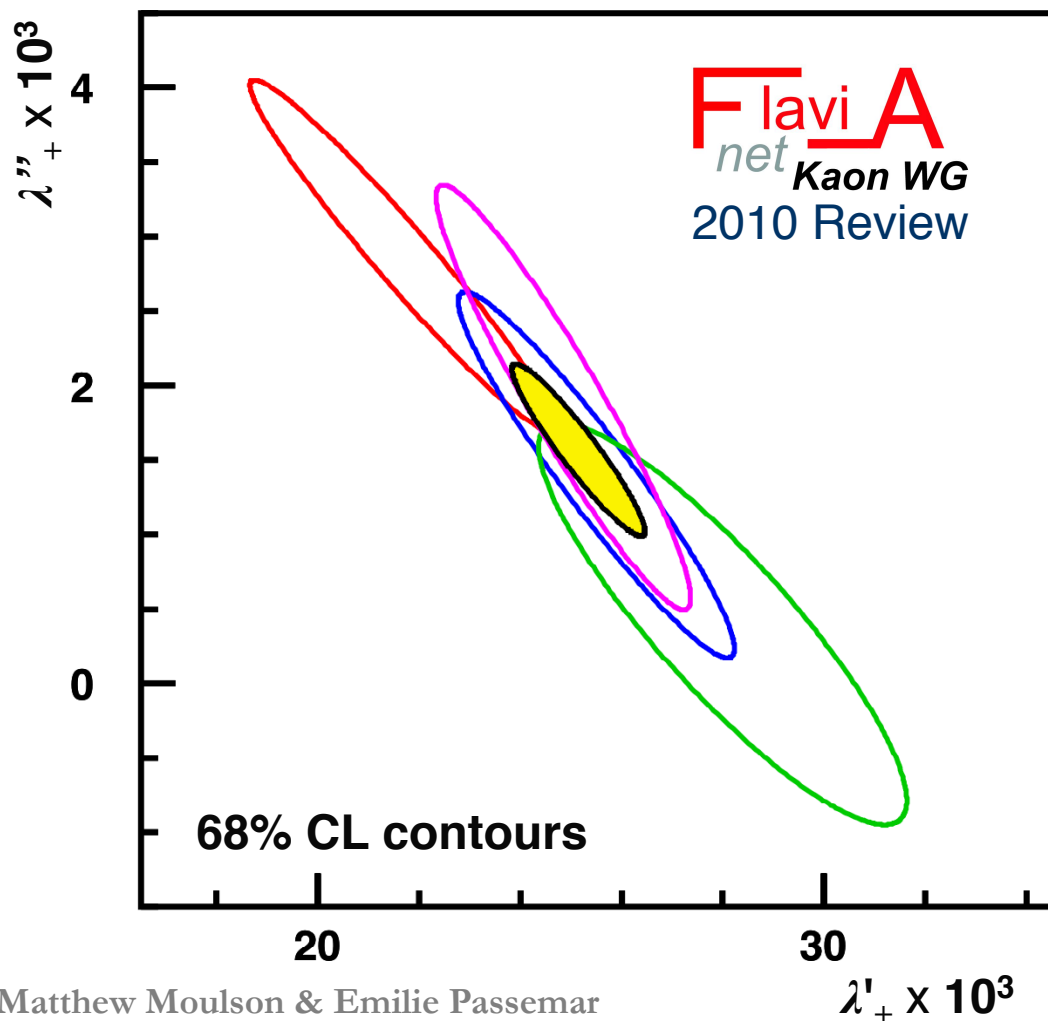
Extraordinarily high precision claimed, esp. for λ_+' , λ_+''

Rudimentary discussion of systematics

Not yet included in updated K_{e3} fit

2.4 Fit to K_{e3} form-factor slopes: 2010

Slopes from **KTeV** **KLOE** **ISTRA+** **NA48** **2010 fit**



Slope parameters $\times 10^3$

$$\lambda'_+ = 25.15 \pm 0.87$$

$$\lambda''_+ = 1.57 \pm 0.38$$

$$\rho(\lambda'_+, \lambda''_+) = -0.941$$

$$\chi^2/\text{ndf} = 5.3/6 \text{ (51\%)}$$

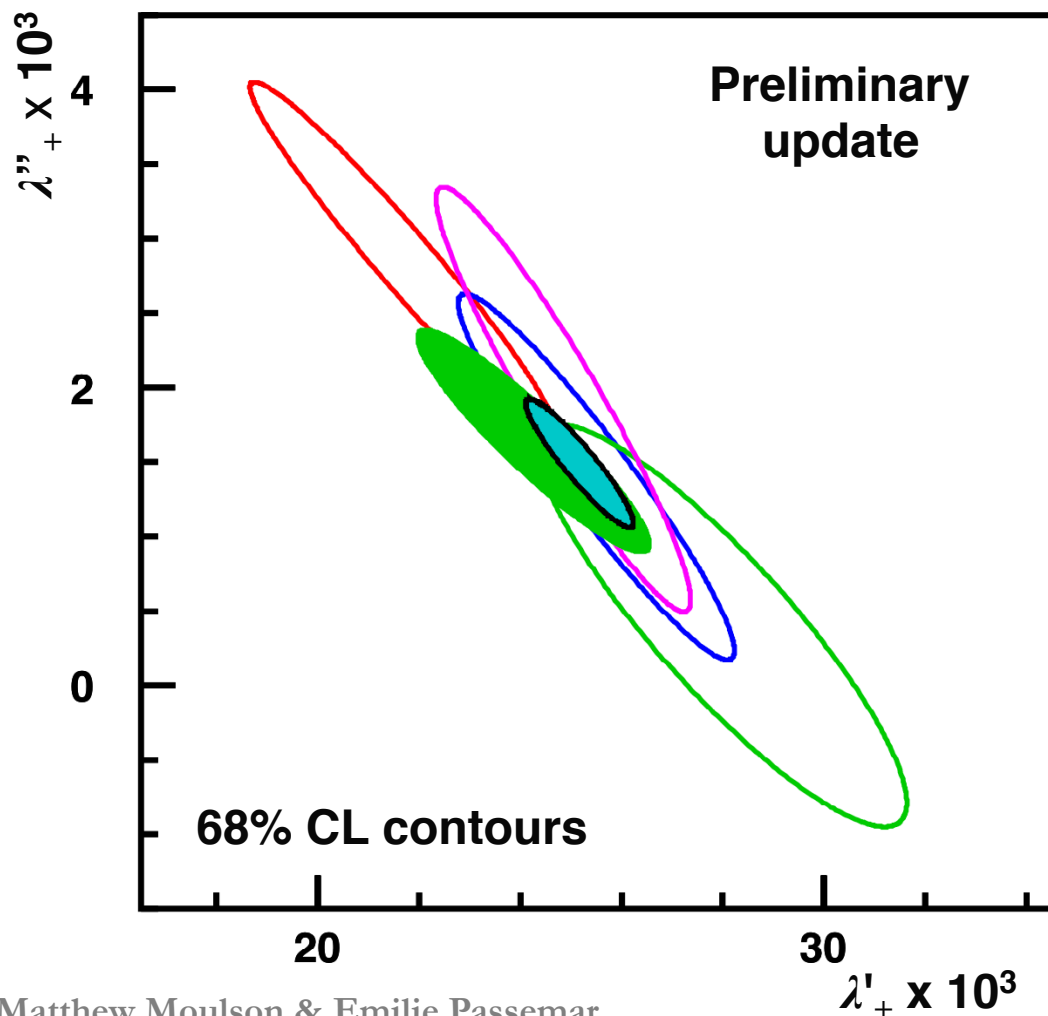
Excellent compatibility
Significance of $\lambda''_+ > 4\sigma$

$$I(K^0_{e3}) = 0.15463(21)$$

$$I(K^+_{e3}) = 0.15900(22)$$

2.4 Fit to K_{e3} form-factor slopes: Update

Slopes from **KTeV** **KLOE** **ISTRA+** **NA48** **NA48/2** **Update**



Slope parameters $\times 10^3$

$$\lambda'_+ = 25.17 \pm 0.70$$

$$\lambda''_+ = 1.49 \pm 0.29$$

$$\rho(\lambda'_+, \lambda''_+) = -0.929$$

$$\chi^2/\text{ndf} = 6.4/10 \text{ (61\%)}$$

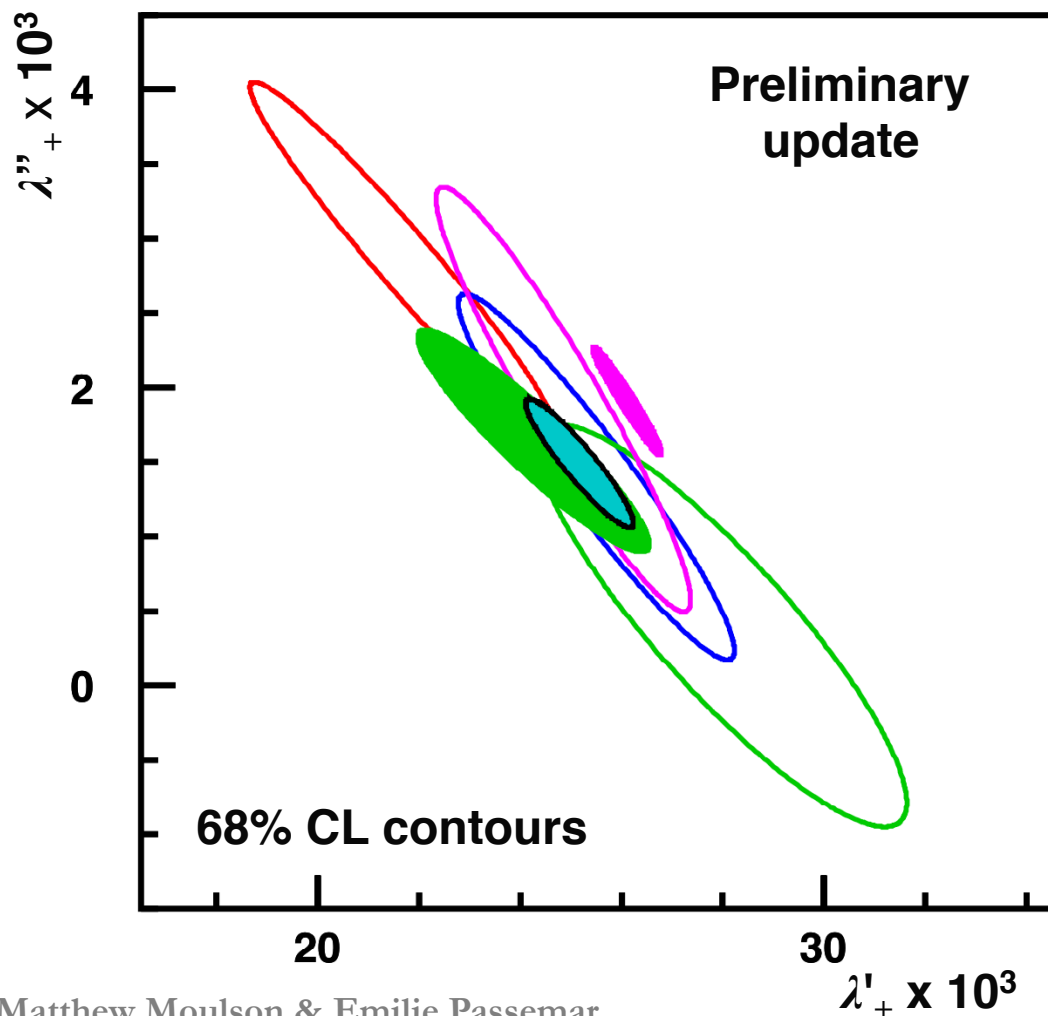
Excellent compatibility
Very small change in λ'_+

$$I(K^0_{e3}) = 0.15463(21)$$

$$I(K^+_{e3}) = 0.15900(22)$$

2.4 Fit to K_{e3} form-factor slopes: Update

Slopes from **KTeV** **KLOE** **ISTRA+** **NA48** **NA48/2** **Update**



Slope parameters $\times 10^3$

$$\lambda'_+ = 25.17 \pm 0.70$$

$$\lambda''_+ = 1.49 \pm 0.29$$

$$\rho(\lambda'_+, \lambda''_+) = -0.929$$

$$\chi^2/\text{ndf} = 6.4/10 \text{ (61\%)}$$

OKA

JETPL 107 (2018)

Not included in the fit

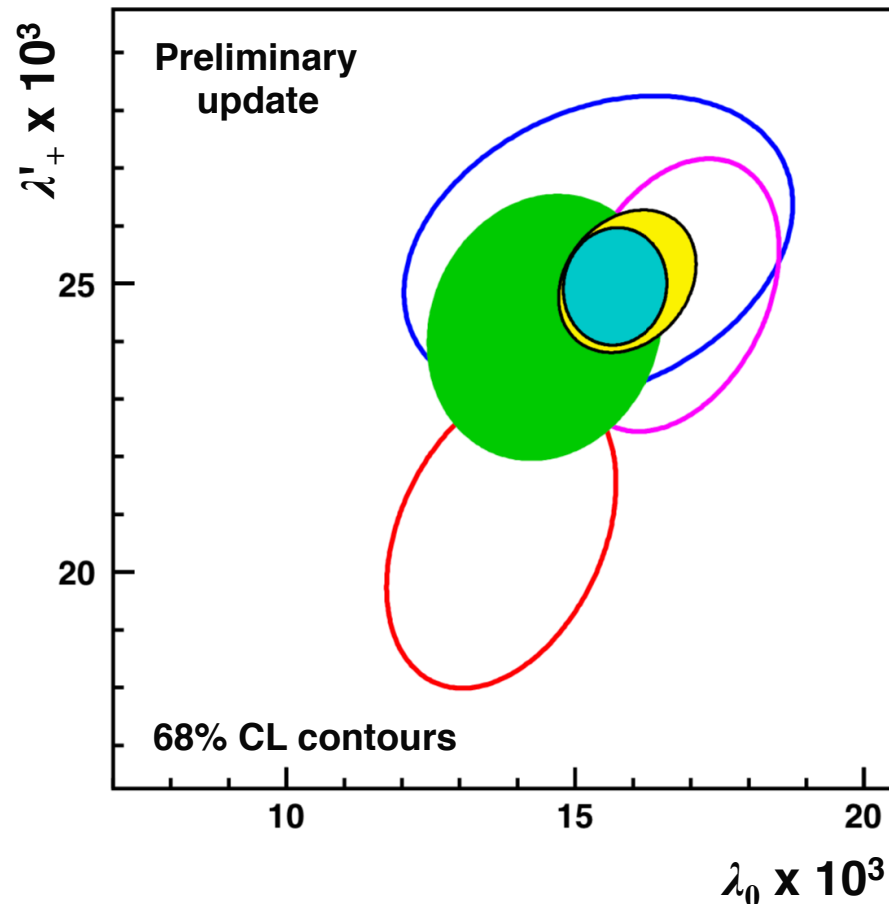
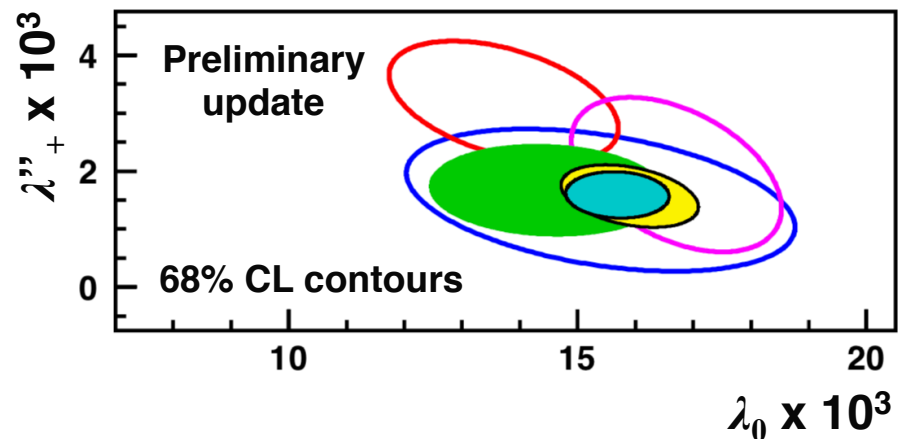
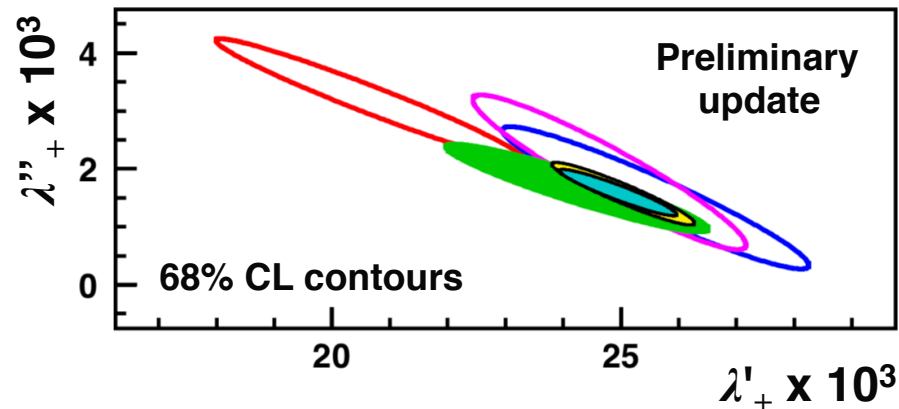
- Stated as preliminary
- If included: $\chi^2/\text{ndf} \rightarrow 45/10$ ($P \sim 10^{-6}$)

2.4 Fits to $K_{e3} + K_{\mu3}$ form-factor slopes: Update

KTeV **KLOE** **ISTRA+** **NA48/2**

NA48 K_{e3} data included in fits but not shown

2010 fit **Update**



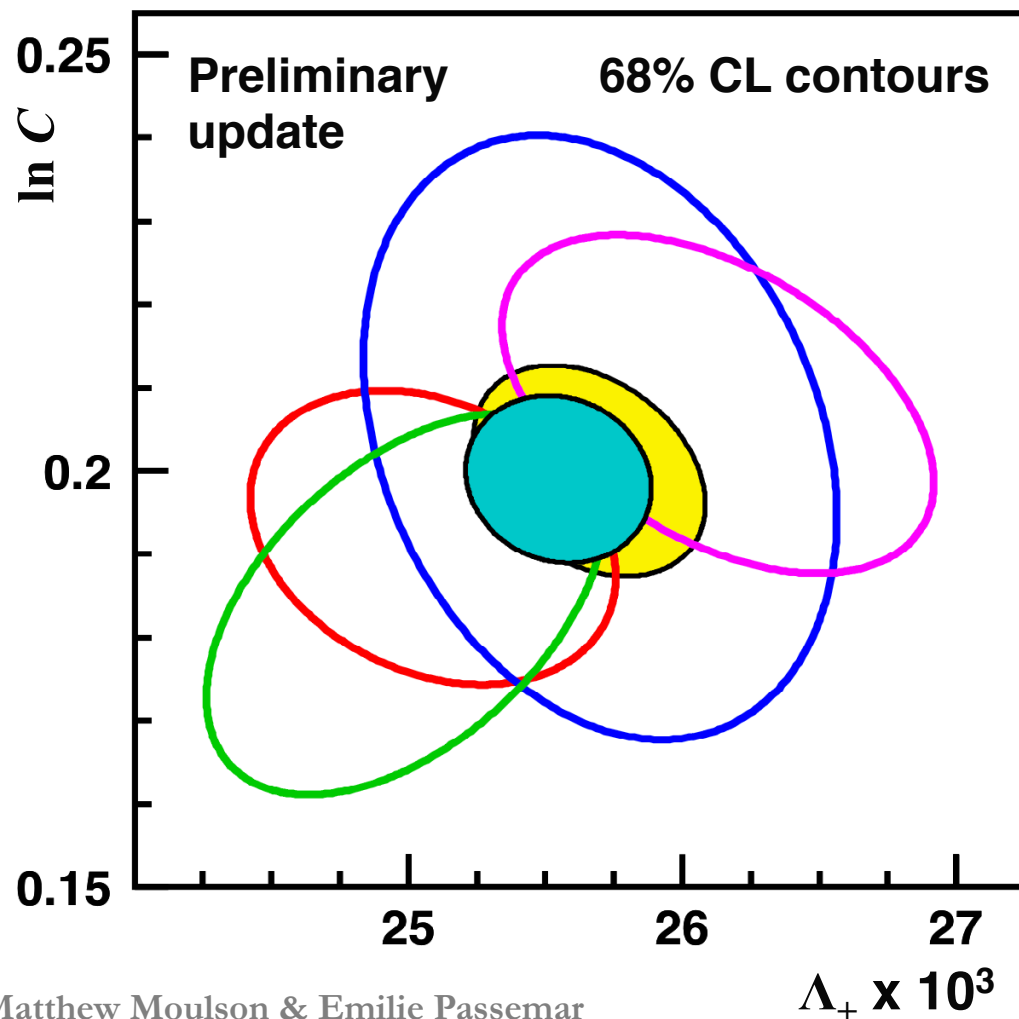
2010: $\chi^2 = 12.1/8$ ($P = 14.5\%$)

Update: $\chi^2 = 13.4/11$ ($P = 26.8\%$)

2.4 Dispersive parameters for $K_{\ell 3}$ form factors

$K_{\ell 3}$ avgs from **KTeV** **KLOE** **ISTRA+** **NA48/2**
 NA48 K_{e3} data included in fits but not shown

2010 fit **Update**



$$\Lambda_+ \times 10^3 = 25.55 \pm 0.38$$

$$\ln C = 0.1992(78)$$

$$\rho(\Lambda_+, \ln C) = -0.110$$

$$\chi^2/\text{ndf} = 7.5/7 \text{ (38\%)}$$

Integrals

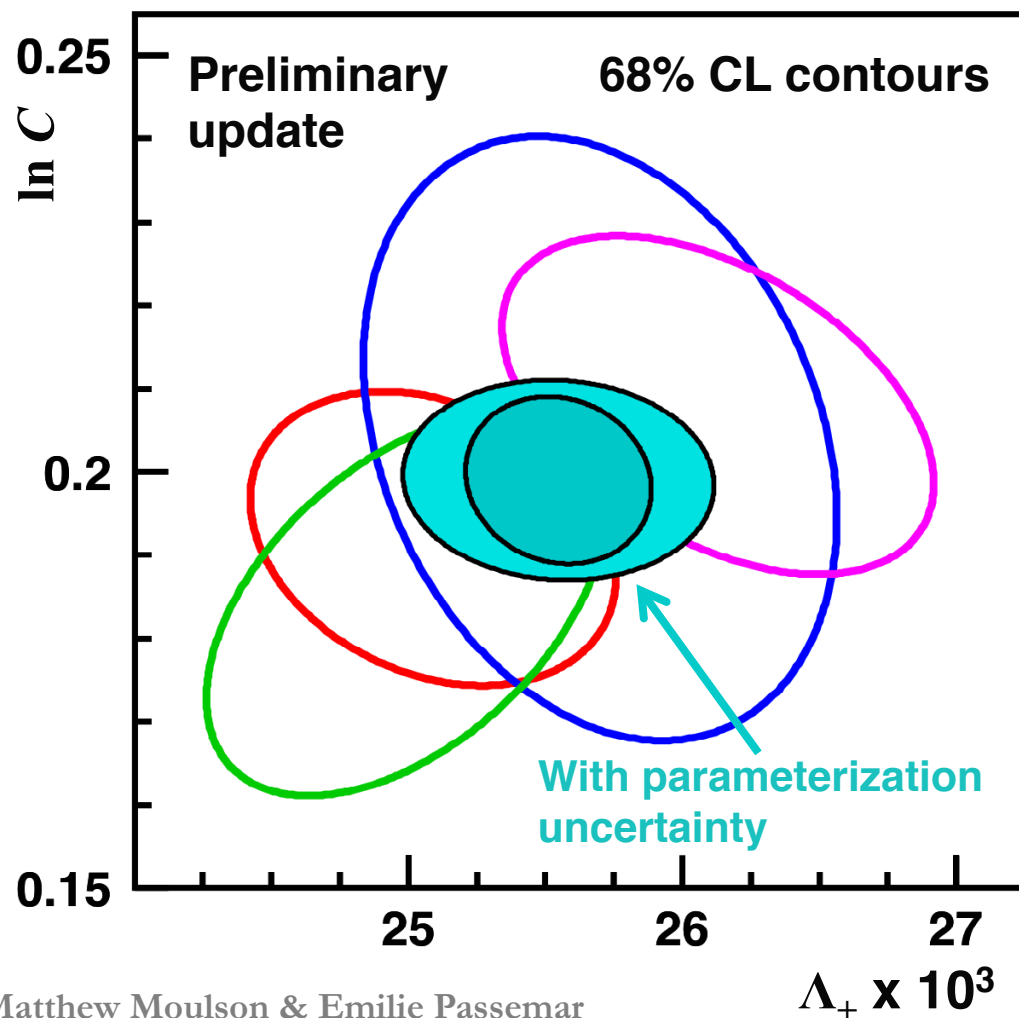
Mode	Update	2010
K^0_{e3}	0.15470(15)	0.15476(18)
K^+_{e3}	0.15915(15)	0.15922(18)
$K^0_{\mu 3}$	0.10247(15)	0.10253(16)
$K^+_{\mu 3}$	0.10553(16)	0.10559(17)

Only tiny changes in central values

2.4 Dispersive parameters for $K_{\ell 3}$ form factors

$K_{\ell 3}$ avgs from **KTeV** **KLOE** **ISTRA+** **NA48/2**
 NA48 K_{e3} data included in fits but not shown

2010 fit **Update**



$$\Lambda_+ \times 10^3 = 25.55 \pm 0.38$$

$$\ln C = 0.1992(78)$$

$$\rho(\Lambda_+, \ln C) = -0.110$$

$$\chi^2/\text{ndf} = 7.5/7 \text{ (38\%)}$$

Fit results include common uncertainty from $H(t)$, $G(t)$:

$$\sigma_{\text{param}}(\Lambda_+) = 0.3 \times 10^{-3}$$

$$\sigma_{\text{param}}(\ln C) = 0.0040$$

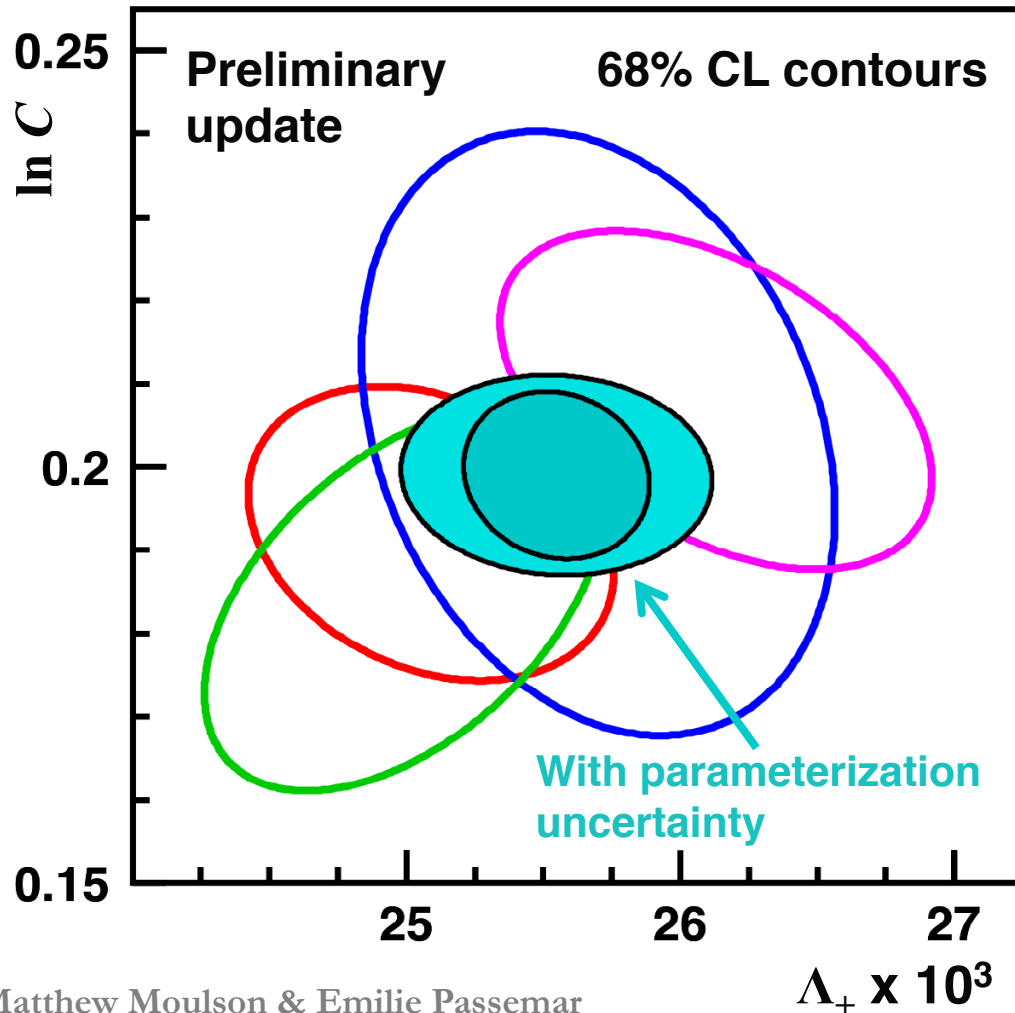
KTeV, Bernard et al.'09

Confidence ellipses shown **without** common uncertainty (except as indicated)

2.4 Dispersive parameters for $K_{\ell 3}$ form factors

$K_{\ell 3}$ avgs from **KTeV** **KLOE** **ISTRA+** **NA48/2**
 NA48 K_{e3} data included in fits but not shown

2010 fit **Update**



$$\Lambda_+ \times 10^3 = 25.55 \pm 0.38$$

$$\ln C = 0.1992(78)$$

$$\rho(\Lambda_+, \ln C) = -0.110$$

$$\chi^2/\text{ndf} = 7.5/7 \text{ (38\%)}$$

Fit results include common uncertainty from $H(t)$, $G(t)$.

Without common uncertainty:

$$\sigma(\Lambda_+) \quad (0.38 \rightarrow 0.22) \times 10^{-3}$$

$$\sigma(\ln C) \quad 0.0078 \rightarrow 0.0067$$

$$\sigma(K_{e3} \text{ int}) \quad 0.10\% \rightarrow 0.09\%$$

$$\sigma(K_{\mu 3} \text{ int}) \quad 0.15\% \rightarrow 0.11\%$$

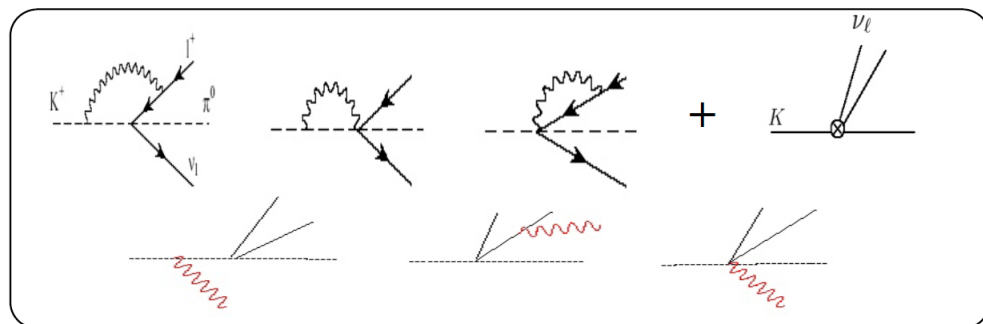
2.5 Long distance electromagnetic corrections

- Master formula for $K \rightarrow \pi l \nu_l$: $K = \{K^+, K^0\}$, $l = \{e, \mu\}$

$$\Gamma(K \rightarrow \pi l \nu [\gamma]) = Br(K_{l3}) * \tau = C_K^2 \frac{G_F^2 m_K^5}{192 \pi^3} S_{EW}^K |V_{us}|^2 \left| f_+^{K^0 \pi^-}(0) \right|^2 I_{KI} \left(1 + 2\Delta_{EM}^{KI} + 2\Delta_{SU(2)}^{K\pi} \right)$$

- Long distance EM corrections: δ_{EM}^{KI}

Cirigliano, Giannotti, Neufeld '08



Mode	$\delta_{EM}^{K\ell}$ (%)
K_{e3}^0	0.495 ± 0.110
K_{e3}^{\pm}	0.050 ± 0.125
$K_{\mu 3}^0$	0.700 ± 0.110
$K_{\mu 3}^{\pm}$	0.008 ± 0.125

- ChPT to $O(p^2 e^2)$
- Fully inclusive prescription for real photons
- Uncertainties: LECs (100%)

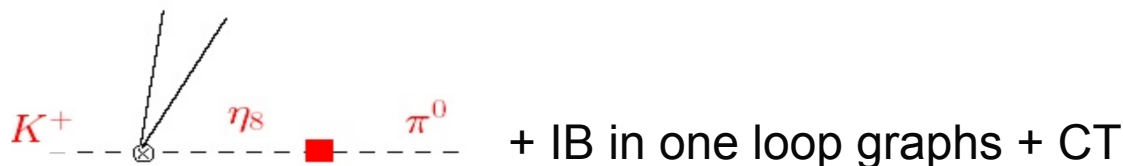
2.6 Isospin breaking corrections

- Master formula for $K \rightarrow \pi l \nu$: $K = \{K^+, K^0\}$, $l = \{e, \mu\}$

$$\Gamma(K \rightarrow \pi l \nu [\gamma]) = Br(K_{l3}) * \tau = C_K^2 \frac{G_F^2 m_K^5}{192 \pi^3} S_{EW}^K |V_{us}|^2 \left| f_+^{K^0 \pi^-}(0) \right|^2 I_{KI} \left(1 + 2 \Delta_{EM}^{KI} + 2 \Delta_{SU(2)}^{K\pi} \right)$$

- Isospin breaking corrections: $\Delta_{SU(2)}^{K\pi}$

$$\Delta_{SU(2)}^{K\pi} = \frac{f_+^{K^+ \pi^0}(0)}{f_+^{K^0 \pi^-}(0)} - 1$$



- In ChPT at $O(p^4)$:

$$\Delta_{SU(2)}^{K\pi} = \frac{3}{4} \frac{1}{Q^2} \left[\frac{m_K^2}{m_\pi^2} + \frac{\chi_{p^4}}{2} \left(1 + \frac{m_s}{\widehat{m}} \right) \right]$$

$$Q^2 \equiv \frac{m_s^2 - \widehat{m}^2}{m_d^2 - m_u^2}$$

$$\left[\widehat{m} \equiv \frac{m_u + m_d}{2} \right]$$

Gasser & Leutwyler '85

2.5 Isospin breaking corrections

$$\Delta^{SU(2)} \equiv \frac{f_+(0)^{K^+\pi^0}}{f_+(0)^{K^0\pi^-}} - 1$$

Strong isospin breaking
Quark mass differences, η - π^0 mixing in $K^+\pi^0$ channel

$$= \frac{3}{4} \frac{1}{Q^2} \left[\frac{\overline{M}_K^2}{\overline{M}_\pi^2} + \frac{\chi_{p^4}}{2} \left(1 + \frac{m_s}{\hat{m}} \right) \right] \quad Q^2 = \frac{m_s^2 - \hat{m}^2}{m_d^2 - m_u^2}$$

$\chi_p^4 = 0.252$
NLO in strong interaction
 $O(e^2 p^2)$ term $\varepsilon_{\text{EM}}^{(4)} \sim 10^{-6}$

= **+2.61(17)%** Calculated using

$$\begin{array}{ll} Q = 22.1(7) & M_K = 494.2(3) \\ m_s/\hat{m} = 27.43(13)(27) & M_\pi = 134.8(3) \end{array}$$

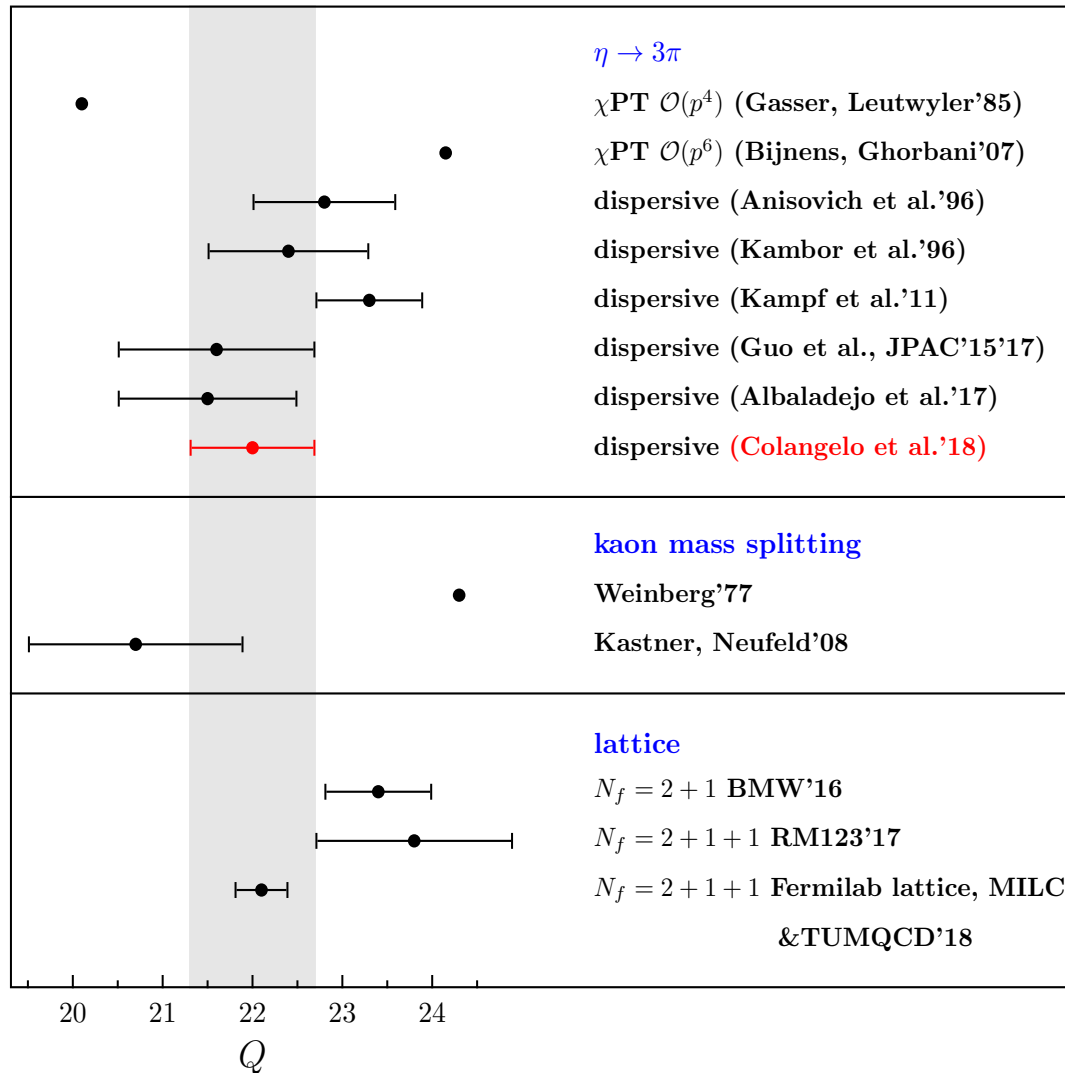
Isospin-limit
meson masses

- Calculation scheme of Kastner & Neufeld '08, Cirigliano et al. '02
- LECs from Bijmans & Ecker '14

Test by evaluating V_{us} from K^\pm and K^0 data with **no** corrections:
Equality of V_{us} values would require $\Delta^{SU(2)} = \mathbf{2.82(38)\%}$

2.5 Isospin breaking corrections

Previous to new results on Q , uncertainty on $\Delta^{SU(2)}$ leading contributor to uncertainty on V_{us} from K^\pm decays — **can it be reduced?**



Continuing progress + systematic review of existing results for light-quark masses may help

Recent dispersion relation analyses of $\eta \rightarrow 3\pi$ Dalitz plot

e.g. Colangelo, Lanz, Leutwyler, E.P'18

1.6 fb⁻¹ KLOE '04 -'05 data

Continuing progress on lattice

E.g. BMW '16 PRL 117

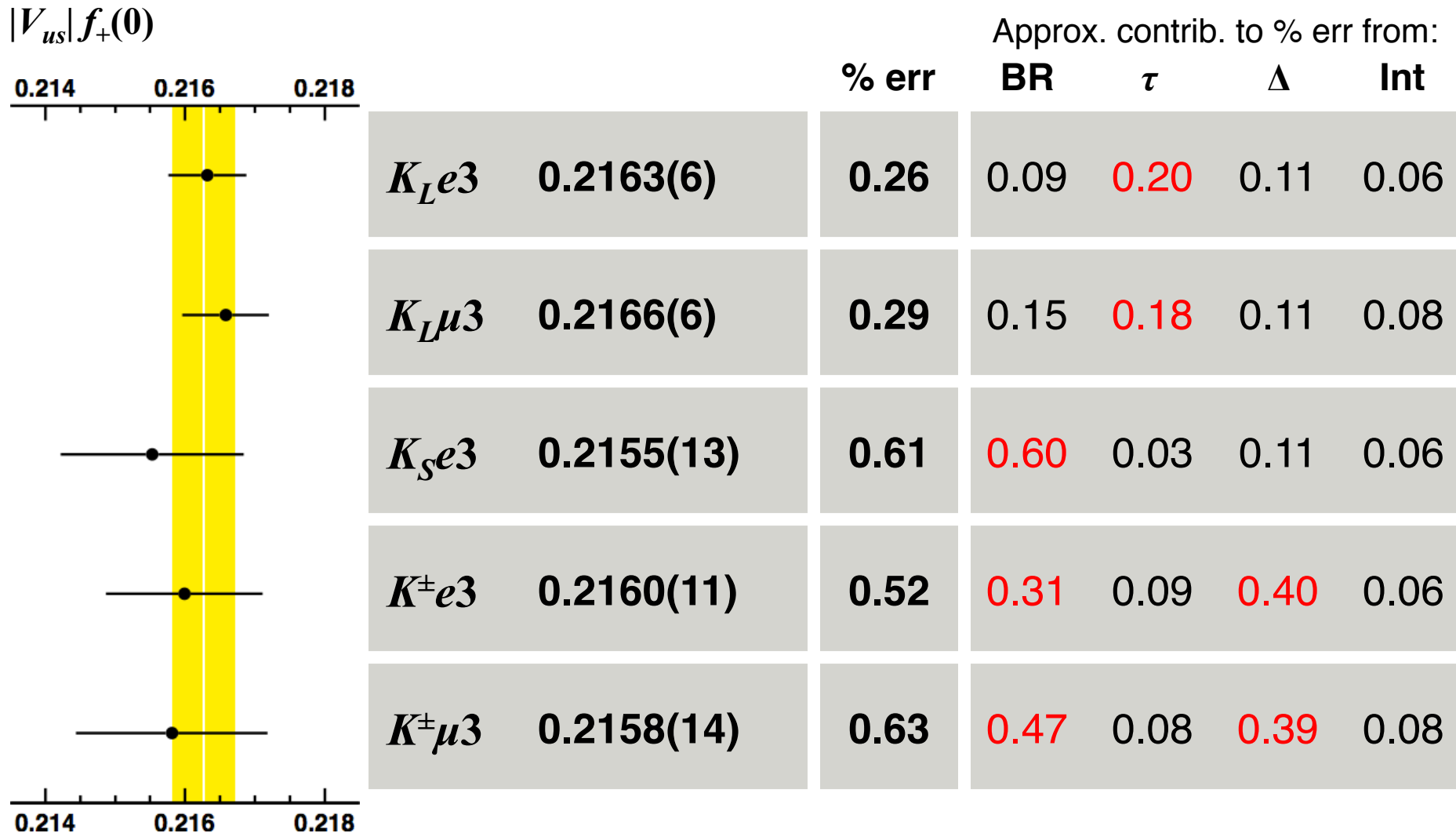
$N_f = 2+1$ QCD, 5sp, m_π phys

Partially quenched QED

$Q = 23.4(4)_{\text{st}}(3)_{\text{sy}}(4)_{\text{QED}}$

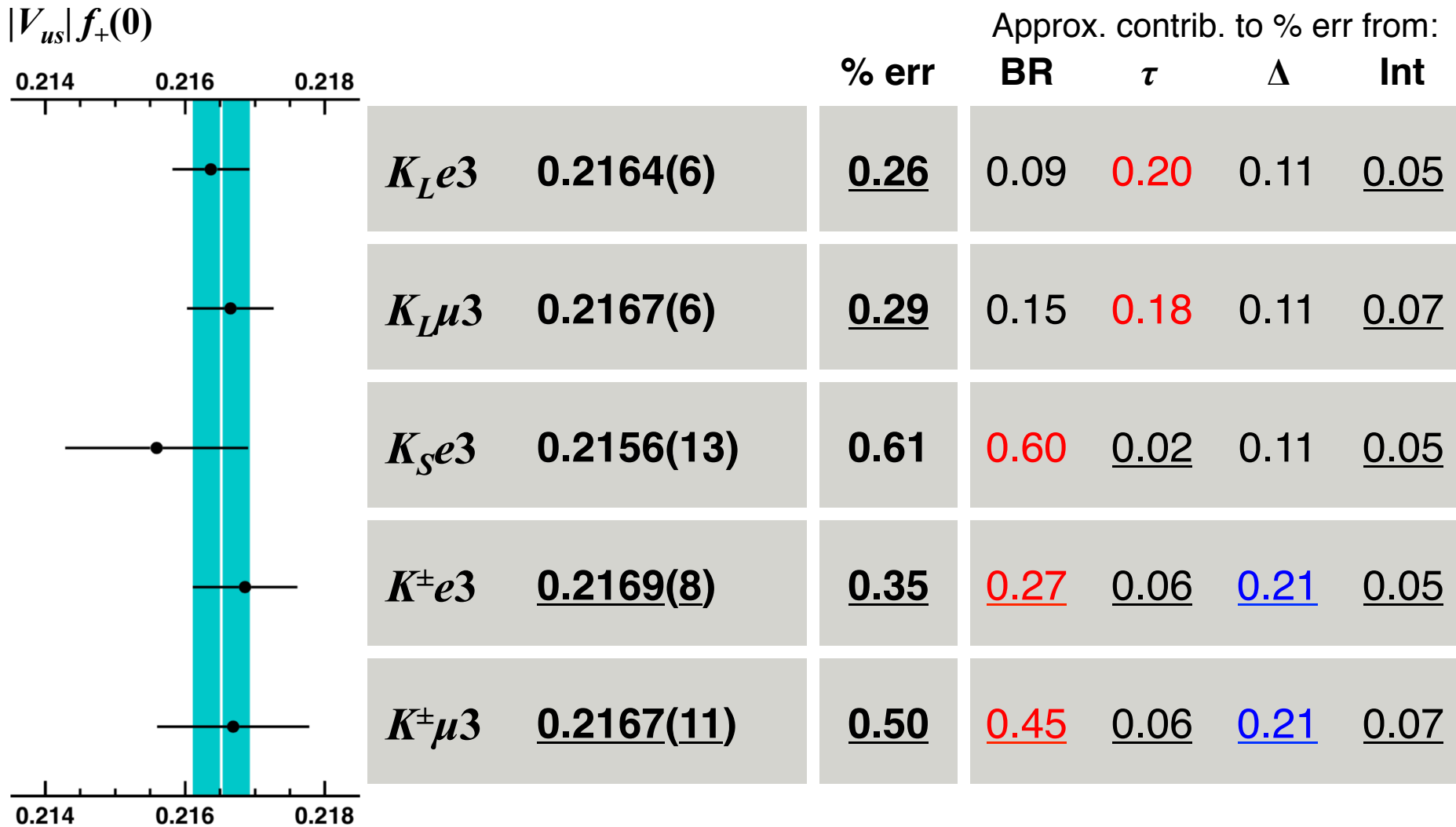
but some tension

2.6 $|V_{us}|f_+(0)$ from world data: 2010



Average: $|V_{us}|f_+(0) = 0.2163(5)$ $\chi^2/\text{ndf} = 0.77/4$ (94%)

2.6 $|V_{us}|f_+(0)$ from world data: Update



Average: $|V_{us}|f_+(0) = 0.21652(41)$ $\chi^2/\text{ndf} = 0.98/4$ (91%)

2.7 Determination of $f_+(0)$

- SU(3) breaking in $f_+(0)$
 - CVC + Ademollo-Gatto theorem: $f_+^{K^0\pi^-}(0) - 1 = O((m_s - m_u)^2)$

$$f_+^{K^0\pi^-}(0) = 1 + \underbrace{f_{p^4}}_{O(m_q)} + \underbrace{f_{p^6}}_{O(m_q^2)} + \dots$$

chiral expansion

- f_{p^4}

Gasser & Leutwyler'85

→ One loop graph :



→ First order in m_q , 2nd order in $(m_s - m_u)$ $\Rightarrow f_{p^4} \sim \frac{(m_s - m_u)^2}{m_s}$

→ No local operators, UV finite, free of uncertainties



$$f_{p^4} = -0.0227$$

2.7 Determination of $f_+(0)$

- SU(3) breaking in $f_+(0)$
 - CVC + Ademollo-Gatto theorem: $f_+^{K^0\pi^-}(0) - 1 = O((m_s - m_u)^2)$

$$f_+^{K^0\pi^-}(0) = 1 + \underbrace{f_{p^4}}_{O(m_q)} + \underbrace{f_{p^6}}_{O(m_q^2)} + \dots$$

chiral expansion

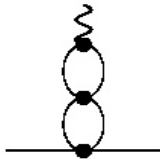
– f_{p^6} :

Bijnens & Talavera'02

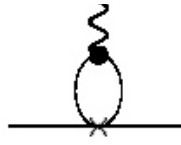
$$f_{p^6} = f_{p^6}^{2\text{-loops}}(\mu) + f_{p^6}^{L_i \times \text{loop}}(\mu) + f_{p^6}^{\text{tree}}(\mu)$$

$$f_{p^6}^{2\text{-loops}}(M_\rho) = 0.0113$$

$$f_{p^6}^{L_i \times \text{loop}}(M_\rho) = -0.0020$$



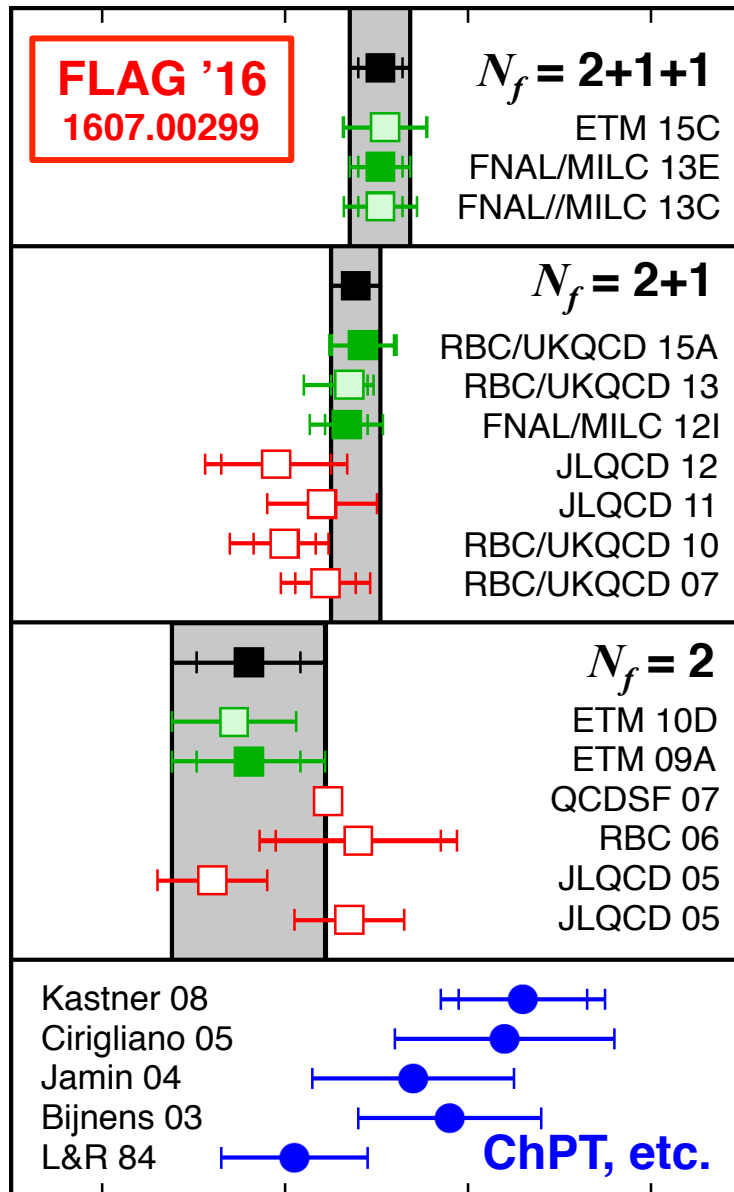
Large positive
chiral loop cont.



$$8 \frac{(M_K^2 - M_\pi^2)^2}{F_\pi^2} \left[\frac{(L_5^r(M_\rho))^2}{F_\pi^2} - C_{12}^r(M_\rho) - C_{34}^r(M_\rho) \right]$$

LECs not fixed by chiral symmetry:
quark model, large-Nc estimates, **LQCD**

2.7 Determination of $f_+(0)$



FLAG'16 averages:

$N_f = 2+1$

$$f_+(0) = 0.9677(27)$$

Uncorrelated average of:

RBC/UKQCD 15A: DWF, $m_\pi \rightarrow 139$ MeV

FNAL/MILC 12I: HISQ, $m_\pi \sim 300$ MeV

$N_f = 2+1+1$

$$f_+(0) = 0.9704(32)$$

FNAL/MILC 13E: HISQ, $m_\pi \rightarrow 135$ MeV

Recent updates:

$N_f = 2+1$

$$f_+(0) = 0.9636(^{+62}_{-65}) \text{ PRD96 (2017)}$$

JLQCD: Overlap, $m_\pi \rightarrow 300$ MeV

Exact chiral symmetry, one lattice spacing

$N_f = 2+1+1$

$$f_+(0) = 0.9709(44)(9)(11) \text{ PRD 93 (2016)}$$

ETM 16: TwMW, 3sp, $m_\pi \rightarrow 210$ MeV

Full q^2 dependence of f_+ , f_0

$$f_+(0) = 0.9696(15)(11) \quad 1809.02827$$

FNAL/MILC update 13E

ChPT:

$N_f = 2+1$

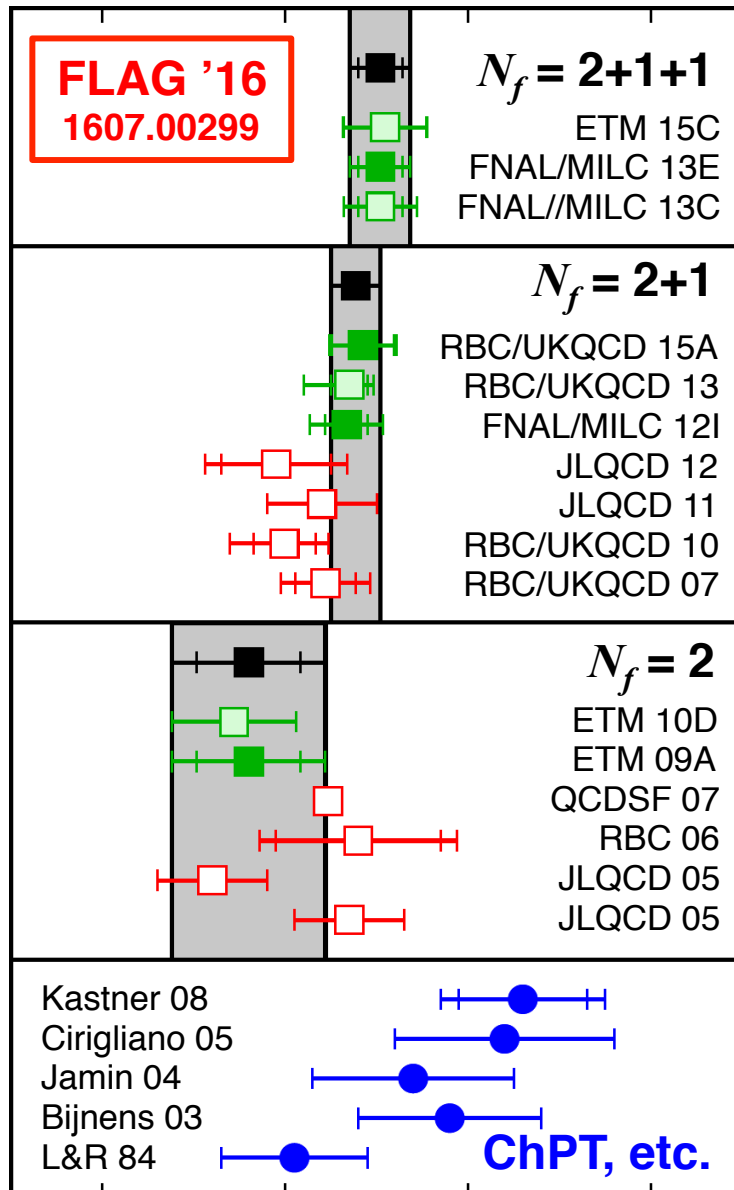
$$f_+(0) = 0.970(8)$$

Chiral Dynamics 15

Ecker 15: According to Bijns 03

New LECs from Bijns, Ecker 14

2.7 Determination of $f_+(0)$



FLAG'16 averages:

$$N_f = 2+1$$

$$f_+(0) = 0.9677(27)$$

Uncorrelated average of:

RBC/UKQCD 15A: DWF, $m_\pi \rightarrow 139$ MeV

FNAL/MILC 12I: HISQ, $m_\pi \sim 300$ MeV

$$N_f = 2+1+1$$

$$f_+(0) = 0.9704(32)$$

FNAL/MILC 13E: HISQ, $m_\pi \rightarrow 135$ MeV

Our averages:

$$N_f = 2+1$$

$$f_+(0) = 0.9677(27)$$

FLAG average, Nov 2016 update
JLQCD17 not included because
only 1 lattice spacing used

RBC/UKQCD15A 0.9685(34)(14)

FNAL/MILC12I 0.9667(23)(33)

$$N_f = 2+1+1$$

$$f_+(0) = 0.9709(44)(9)(11)$$

FNAL/MILC18 preliminary replaces
FNAL/MILC13E in FLAG average

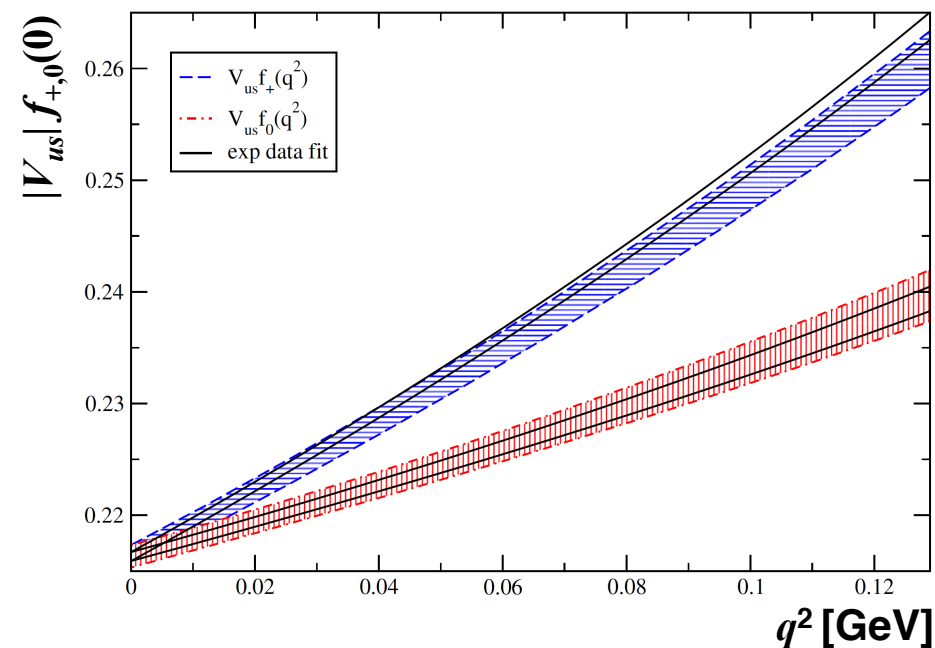
FNAL/MILC18 0.9696(15)(11)

ETM16 0.9709(44)(9)(11)_{ext}

q^2 dependence of $K\pi$ form factors

ETM
PRD 93 (2016)

$N_f = 2+1+1$ Twisted-mass Wilson fermions
3 lattice spacings, smallest $m_\pi \rightarrow 210$ MeV
Results for full q^2 dependence of f_+, f_0



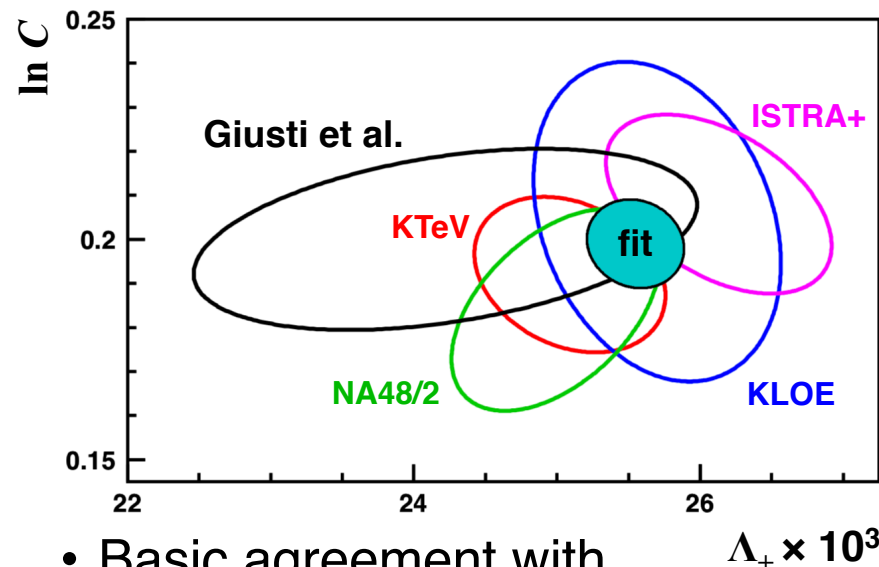
Fit synthetic data points with
dispersive parameterization

$$\Lambda_+ = 24.22(1.16) \times 10^{-3} \quad \rho(\Lambda_+, f_+(0)) = -0.228$$

$$\ln C = 0.1998(138) \quad \rho(\ln C, f_+(0)) = -0.719$$

$$\rho(\Lambda_+, \ln C) = +0.376$$

$$f_+(0) = 0.9709(44)_{\text{st}}(9)_{\text{sy}}(11)_{\text{ext}}$$



- Basic agreement with experimental results
- Confirms basic correctness of lattice calculations for $f_+(0)$
- In the near future FF parameters will be obtained on lattice?

2.8 $|V_{us}|(K_{\ell 3})$ and $|V_{ud}|(0^+ \rightarrow 0^+)$: Update

Hardy & Towner, CIPANP '18

$$|V_{ud}| = 0.97420(21)$$

World data set very robust

14 transitions with compatible measurements at 0.1% precision or better

From FlaviaNet 2010 $K_{\ell 3}$ analysis

$$|V_{us}| f_+(0) = 0.2163(5) \quad |V_{us}| = 0.2254(13)$$

$$\text{with } f_+(0) = 0.959(5) \quad \text{with } |V_{ud}| = 0.97425(22)$$

$$\Delta_{\text{CKM}} = +0.0000(8)$$

Update with $|V_{us}| f_+(0) = 0.21652(41)$ and $|V_{ud}| = 0.97420(21)$

$$N_f = 2+1$$

$$f_+(0) = 0.9677(27)$$

$$V_{us} = 0.22375(43)_{\text{exp}}(62)_{\text{lat}}$$

$$\Delta_{\text{CKM}} = -0.00085(19)_{\text{exp}}(28)_{\text{lat}}(41)_{ud} = -1.6\sigma$$

$$N_f = 2+1+1$$

$$f_+(0) = 0.9698(17)$$

$$V_{us} = 0.22326(43)_{\text{exp}}(39)_{\text{lat}}$$

$$\Delta_{\text{CKM}} = -0.00107(19)_{\text{exp}}(17)_{\text{lat}}(41)_{ud} = -2.2\sigma$$

1.5-2 σ inconsistency with unitarity first seen with 2014-era lattice results

Relative to 2014 slightly better agreement between $N_f = 2+1$ and $N_f = 2+1+1$

V_{ud} from $0^+ \rightarrow 0^+$

$$\frac{1}{t} = \frac{G_\mu^2 |V_{ud}|^2 m_e^5}{\pi^3 \log 2} f(Q) (1 + RC) \longrightarrow ft (1 + RC) = \frac{2984.48(5) s}{|V_{ud}|^2}$$

$$(1 + RC) = (1 - \delta_C) (1 + \delta_R) (1 + \Delta_C)$$

$$\langle f | \tau_+ | i \rangle = \sqrt{2} (1 - \delta_C/2)$$

Coulomb distortion
of wave-functions

$$\delta_C \sim 0.5\%$$

Towner-Hardy
Ormand-Brown

Nucleus-dependent
rad. corr.
(Z, E^{\max} , nuclear structure)

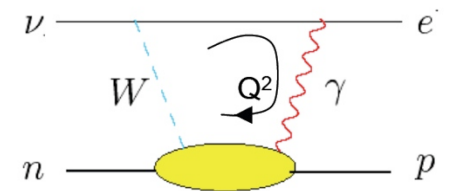
$$\delta_R \sim 1.5\%$$

Sirlin-Zucchini '86
Jaus-Rasche '87

Nucleus-independent
short distance rad. corr.

$$\Delta_R \sim 2.4\%$$

Marciano-Sirlin '06

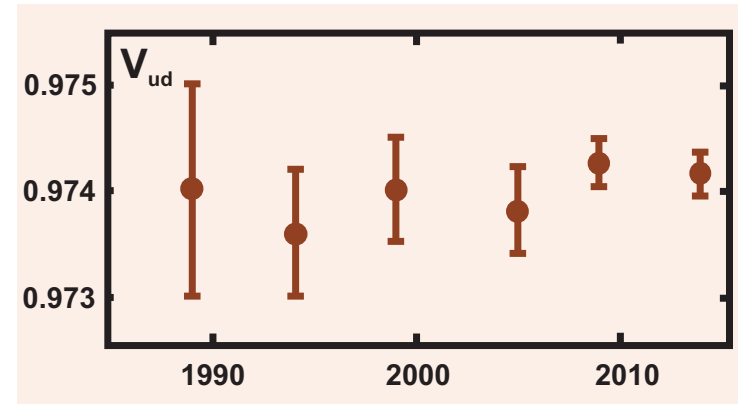


V_{ud} from $0^+ \rightarrow 0^+$

In 2010 :

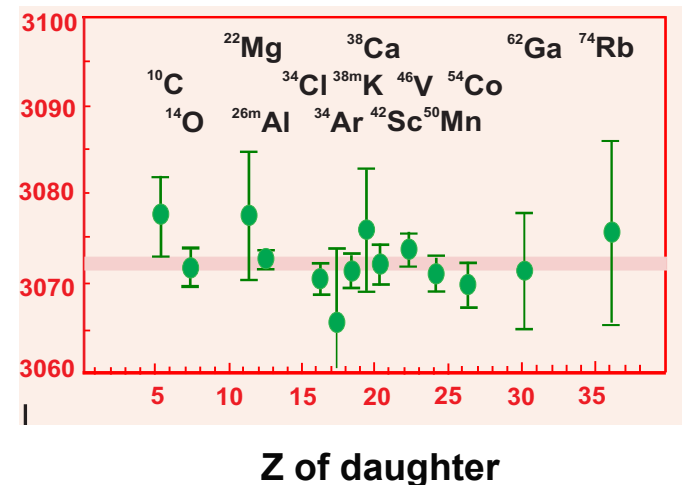
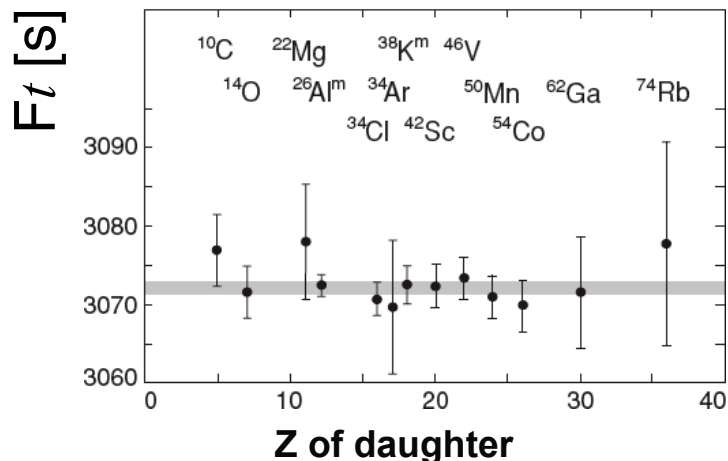
- Survey of 150 measurements of 13 different $0^+ \rightarrow 0^+$ β decays
- 27 new ft measurements including Penning-trap measurements for QEC
- Some old measurements dropped
- Improved EW radiative corrections
Marciano & Sirlin'06
- New $SU(2)$ -breaking corrections
Towner & Hardy'08

Evolution of V_{ud} over years



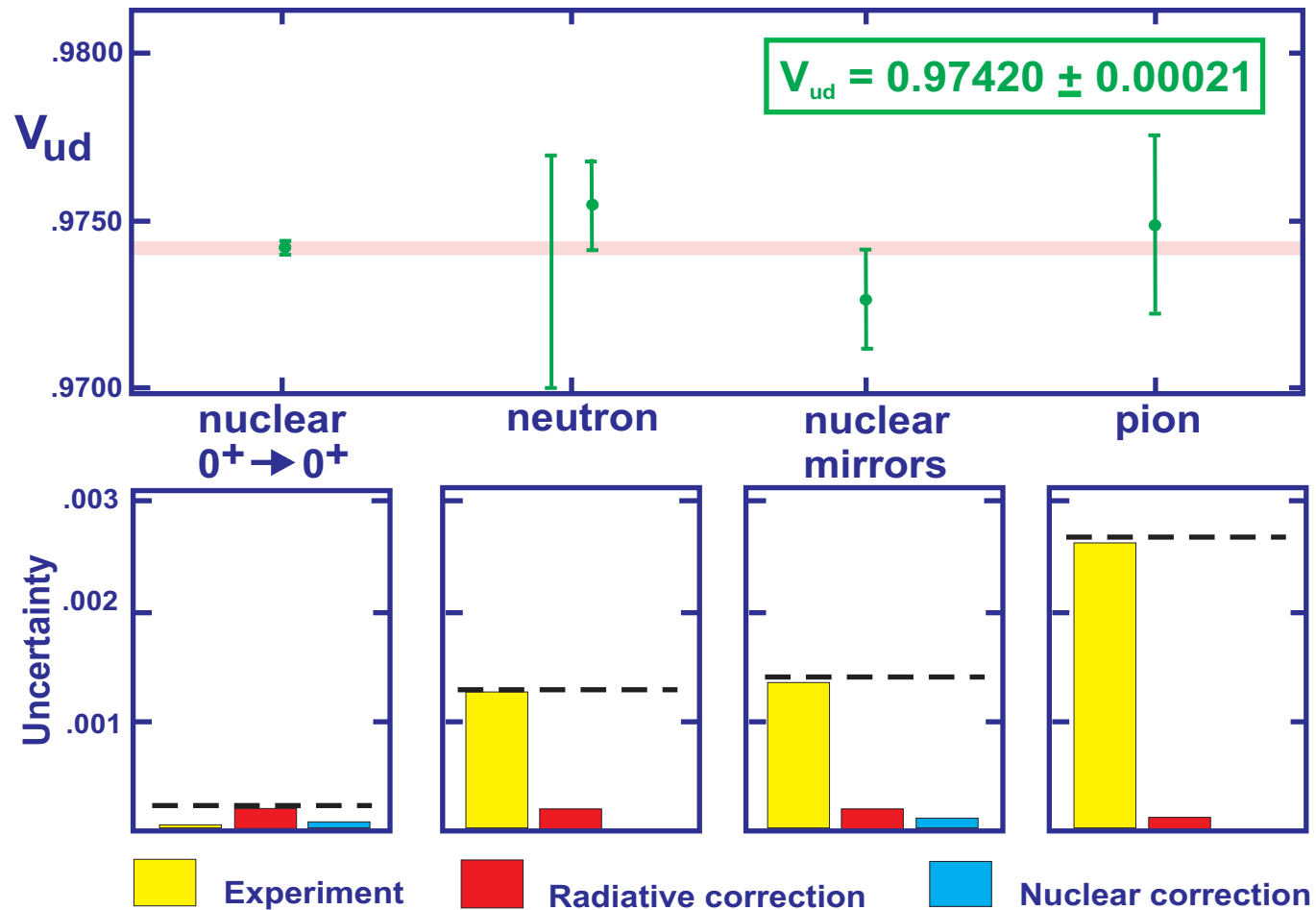
Towner@CIPANP'18

Since then 24 new measurements, critical review of IB correction and test of CVC



Current Status of V_{ud} determination

Towner@CIPANP'18



2.8 $|V_{us}|(K_{\ell 3})$ and $|V_{ud}|(0^+ \rightarrow 0^+)$: Update

**Seng, Gorchtein, Patel &
Ramsey-Musolf**
arXiv:1807.10197

$|V_{ud}| = 0.97366(15)$

-1.5σ shift in V_{ud}

New calculation γW -box contribution to universal radiative correction using dispersion relations and DIS structure functions

- Contribution possibly already in part included in structure-dependent radiative corrections
- **Needs verification!**

Update with $|V_{us}|f_+(0) = 0.21652(41)$ and $|V_{ud}| = 0.97366(15)$

Choice of $f_+(0)$		V_{us}	$\Delta_{\text{CKM}} = V_{ud}^2 + V_{us}^2 - 1$	
$N_f = 2+1$	0.9677(27)	0.2238(8)	-0.0019(5)	$= -4.2\sigma$
$N_f = 2+1+1$	0.9698(17)	0.2233(6)	-0.0021(4)	$= -5.4\sigma$

If correct, 4-5 σ unitarity violation in first row!

Calculation is attracting interest and requires better understanding

3. V_{us}/V_{ud} from $K_{\ell 2}/\pi_{\ell 2}$ decays

3.1 Master formula for V_{us}/V_{ud} from $K_{\ell 2}/\pi_{\ell 2}$ decays

- From $K_{\ell 2}/\pi_{\ell 2}$:

$$\frac{|V_{us}|}{|V_{ud}|} \frac{f_K}{f_\pi} = \left(\frac{\Gamma_{K_{\mu 2}(\gamma)} m_{\pi^\pm}}{\Gamma_{\pi_{\mu 2}(\gamma)} m_{K^\pm}} \right)^{1/2} \frac{1 - m_\mu^2/m_{\pi^\pm}^2}{1 - m_\mu^2/m_{K^\pm}^2} \left(1 - \frac{1}{2} \delta_{\text{EM}} - \frac{1}{2} \delta_{SU(2)} \right)$$

Inputs from theory:

Cirigliano, Neufeld '11

$$\delta_{\text{EM}} = -0.0069(17)$$

Long-distance EM corrections

$$\delta_{SU(2)} = -0.0043(5)(11)$$

Strong isospin breaking

$$f_K/f_\pi \rightarrow f_{K^\pm}/f_{\pi^\pm}$$

Lattice: f_K/f_π

Cancellation of lattice-scale uncertainties from ratio

NB: Most lattice results already corrected for $SU(2)$ -breaking: f_{K^\pm}/f_{π^\pm}

Inputs from experiment:

Updated K^\pm BR fit:

$$\text{BR}(K_{\mu 2}^\pm) = 0.6358(11)$$

$$\tau_{K^\pm} = 12.384(15) \text{ ns}$$

PDG:

$$\text{BR}(\pi_{\mu 2}^\pm) = 0.9999$$

$$\tau_{\pi^\pm} = 26.033(5) \text{ ns}$$

$$|V_{us}/V_{ud}| \times f_{K^\pm}/f_{\pi^\pm} = 0.27599(37)$$

No $SU(2)$ -breaking correction

3.2 Electromagnetic corrections

Giusti et al.
PRL 120 (2018)

First lattice calculation of EM corrections to P_{l2} decays

- Ensembles from ETM
- $N_f = 2+1+1$ Twisted-mass Wilson fermions

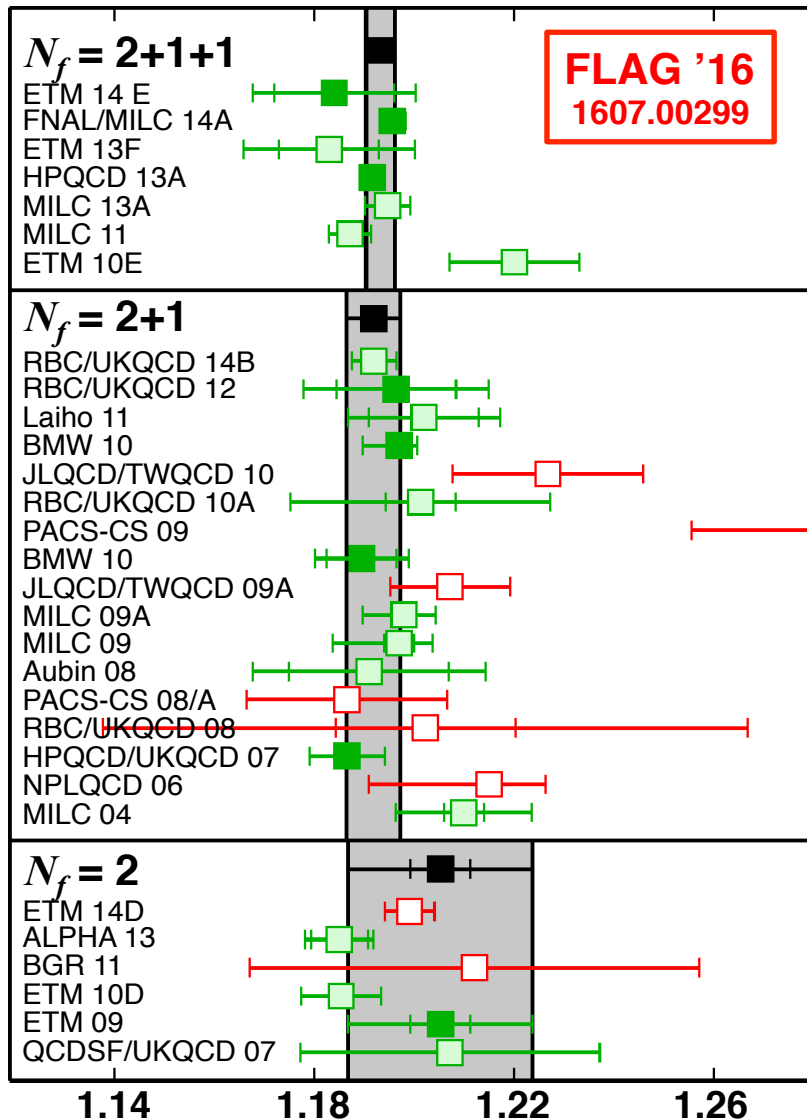
$$\delta_{SU(2)} + \delta_{EM} = -0.0122(16)$$

- Uncertainty from quenched QED included (0.0006)

Compare to ChPT result from Cirigliano, Neufeld '11:

$$\delta_{SU(2)} + \delta_{EM} = -0.0112(21)$$

3.3 Lattice results for f_K/f_π



FLAG '16 averages:

$N_f = 2+1$ $f_{K^\pm}/f_{\pi^\pm} = 1.192(5)$
Unchanged from FLAG '13 average

$N_f = 2+1+1$ $f_{K^\pm}/f_{\pi^\pm} = 1.1933(29)$
ETM 14E: TwM, 3sp, $m_\pi = 210$ -450 MeV
FNAL/MILC 14A: HISQ, 4sp, m_π phys
Updates MILC 13A
HPQCD 13A: HISQ, 3sp, m_π phys,
Same ensembles as FNAL/MILC 14A

Recent updates:

$N_f = 2+1$ $f_K/f_\pi = 1.1945(45)$
RBC/UKQCD '14: DWF, $m_\pi = 139$ MeV
 f_K and f_π separately (isospin limit)
Recently published

$f_{K^\pm}/f_{\pi^\pm} = 1.1978(28)$
BMW '16: Clover, 5sp, $m_\pi \rightarrow 139$ MeV

3.3 Lattice results for f_K/f_π

Our updates of FLAG averages for results without $SU(2)$ -breaking

$N_f = 2+1+1$

$f_K/f_\pi = 1.1960(40)$

HPQCD13A	1.1948(15)(18)	} Correlated uncertainties
FNAL/MILC14A*	1.1983($^{+28}_{-21}$)	
ETM14E	1.188(15)	← Uncorrelated uncertainty

$N_f = 2+1$

$f_K/f_\pi = 1.1927(38)$

HPQCD/UKQCD07	1.198(7)
RBC/UKQCD14B	1.1945(45)
BMW16	1.182(10)(26)

* Corrected using Cirigliano, Neufeld '11 with updated values:

$$\frac{f_{K^\pm}}{f_{\pi^\pm}} = \frac{f_K}{f_\pi} \sqrt{1 + \delta_{SU(2)}} \quad \delta_{SU(2)} \approx \frac{3}{4R} \left[-\frac{4}{3}(f_K/f_\pi - 1) + \frac{2}{3(4\pi)^2 f_0^2} \left(M_K^2 - M_\pi^2 - M_\pi^2 \ln \frac{M_K^2}{M_\pi^2} \right) \right]$$

$$R = 34.4(2.1) \quad \text{Colangelo et al. '18} \quad M_K = 494.2(3) \quad \text{FLAG '17}$$

3.3 Lattice results for f_K/f_π

$$|V_{us}/V_{ud}| \times f_{K^\pm}/f_{\pi^\pm} = 0.27599(37) \text{ and } |V_{ud}| = 0.97420(21)$$

$$\delta_{SU(2)} + \delta_{EM} = -0.0122(16) \text{ from Giusti et al. '18}$$

$$N_f = 2+1$$

$$f_{K^\pm}/f_{\pi^\pm} = 1.1927(38)$$

$$V_{us} = 0.22604(29)_{\text{exp}}(72)_{\text{lat}}(05)_{ud}$$

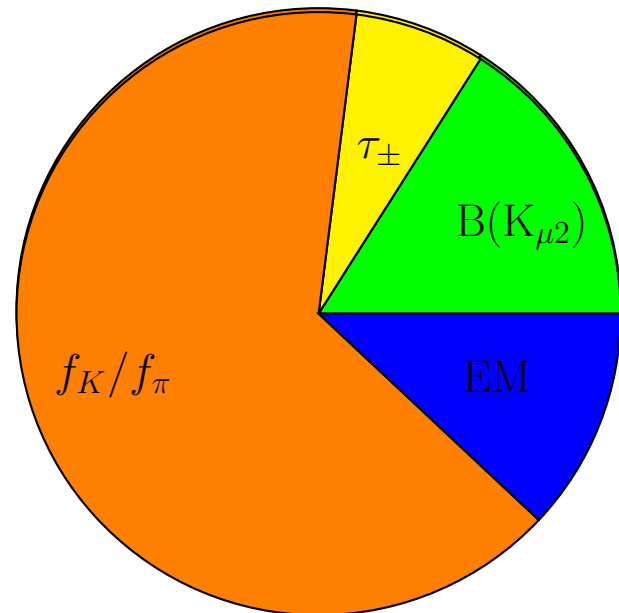
$$\Delta_{\text{CKM}} = +0.00018(13)_{\text{exp}}(33)_{\text{lat}}(43)_{ud} = +0.3\sigma$$

$$N_f = 2+1+1$$

$$f_{K^\pm}/f_{\pi^\pm} = 1.1960(40)$$

$$V_{us} = 0.22542(29)_{\text{exp}}(75)_{\text{lat}}(05)_{ud}$$

$$\Delta_{\text{CKM}} = -0.00011(13)_{\text{exp}}(34)_{\text{lat}}(43)_{ud} = -0.2\sigma$$



$$\frac{|V_{us}|}{|V_{ud}|} \frac{f_K}{f_\pi} = \left(\frac{\Gamma_{K_{\mu 2}(\gamma)} m_{\pi^\pm}}{\Gamma_{\pi_{\mu 2}(\gamma)} m_{K^\pm}} \right)^{1/2} \frac{1 - m_\mu^2/m_{\pi^\pm}^2}{1 - m_\mu^2/m_{K^\pm}^2} \left(1 - \frac{1}{2} \delta_{EM} - \frac{1}{2} \delta_{SU(2)} \right)$$

3.3 Lattice results for f_K/f_π

$$|V_{us}/V_{ud}| \times f_{K^\pm}/f_{\pi^\pm} = 0.27599(37) \text{ and } |V_{ud}| = 0.97420(21) \\ \delta_{SU(2)} + \delta_{EM} = -0.0122(16) \text{ from Giusti et al. '18}$$

$N_f = 2+1$ $f_{K^\pm}/f_{\pi^\pm} = 1.1927(38)$	$V_{us} = 0.22604(29)_{\text{exp}}(72)_{\text{lat}}(05)_{ud}$ $\Delta_{\text{CKM}} = +0.00018(13)_{\text{exp}}(33)_{\text{lat}}(43)_{ud} = +0.3\sigma$
$N_f = 2+1+1$ $f_{K^\pm}/f_{\pi^\pm} = 1.1960(40)$	$V_{us} = 0.22542(29)_{\text{exp}}(75)_{\text{lat}}(05)_{ud}$ $\Delta_{\text{CKM}} = -0.00011(13)_{\text{exp}}(34)_{\text{lat}}(43)_{ud} = -0.2\sigma$

$K_{\ell 2}$ results give better agreement with unitarity via V_{ud} than $K_{\ell 3}$ results (-2σ)

Exercise:

- Assume $|V_{ud}|$, $|V_{us}/V_{ud}| \times f_{K^\pm}/f_{\pi^\pm}$, and f_{K^\pm}/f_{π^\pm} all correct
- In $K_{\ell 3}$ does the discrepancy arise from data or from lattice results for $f_+(0)$?

4. V_{us} and Unitarity of the CKM matrix

4.1 Looking for New Physics with K_{12} and K_{13}

- Callan-Treiman theorem:

$$C = \overline{f}_0(\Delta_{K\pi}) = \frac{F_K}{F_\pi f_+(0)} + \Delta_{CT} = \underbrace{\frac{F_K |V^{us}|}{F_\pi |V^{ud}|} \frac{1}{f_+(0) |V^{us}|} |V^{ud}|}_{\text{Very precisely known from Br(Kl2/\pi l2), \Gamma(Ke3) and |V_{ud}|}} r + \Delta_{CT}$$

$B_{\text{exp}} = 1.2446(41)$

- In the Standard Model : $r = 1$ $(\ln C_{SM} = 0.2141(73))$ $\Delta_{CT} = (-3.5 \pm 8) \cdot 10^{-3}$

- In presence of new physics, new couplings : $r \neq 1$

NLO value + large error bars in agreement with

Experiment $K_{e3}+K_{\mu3}$	$\ln C$
<i>NA48'07</i> ($K_{\mu3}$ alone)	0.144(14)
<i>KLOE'08</i>	0.204(25)
<i>KTeV'10</i>	0.192(12)
<i>NA48/2, previous talk</i>	0.184(15)

4.1 Form factors & the Callan-Treiman relation

Callan-Treiman relation:

$$\tilde{f}_0(t_{\text{CT}}) = \frac{f_K}{f_\pi} \frac{1}{f_+(0)} + \Delta_{\text{CT}}$$

Use ChPT & form-factor data to test $N_f = 2+1+1$ lattice consistency:

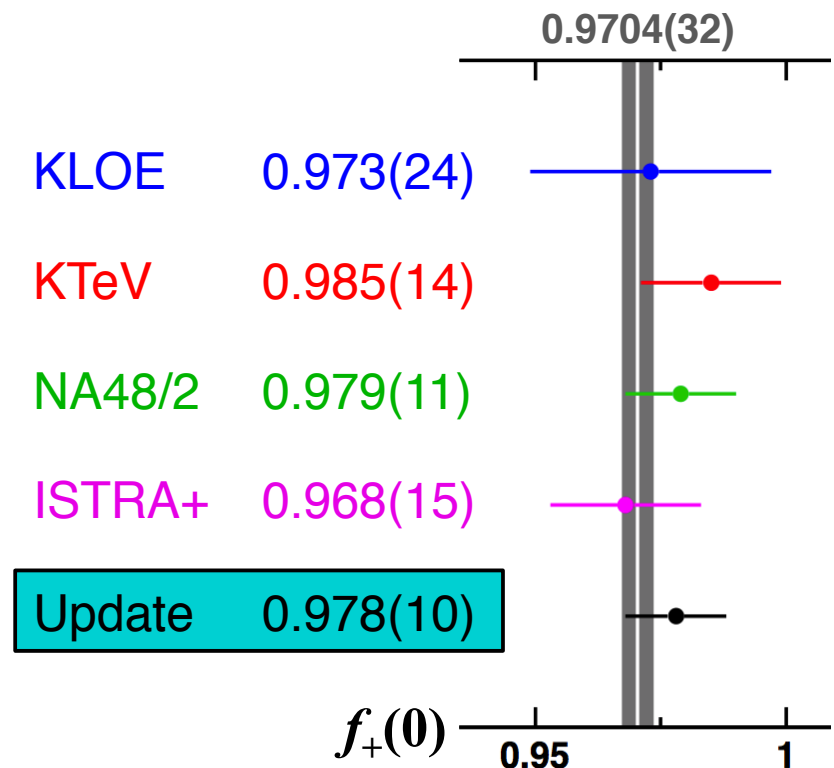
- Use lattice reference value
 $f_K/f_\pi = 1.1933(29)$
- Obtain $f_+(0)$ corresponding to each result for $\ln C$
- Compare to lattice reference value
 $f_+(0) = 0.9704(32)$
- Basic consistency (0.7σ) between lattice values for f_K/f_π and $f_+(0)$ and measurements of $\ln C$
- **Uses no experimental information on decay widths**

$$t_{\text{CT}} = m_K^2 - m_\pi^2$$

$$\Delta_{\text{CT}} = (-3.5 \pm 0.8) \times 10^{-3} \sim \mathcal{O}(m_u, m_d)$$

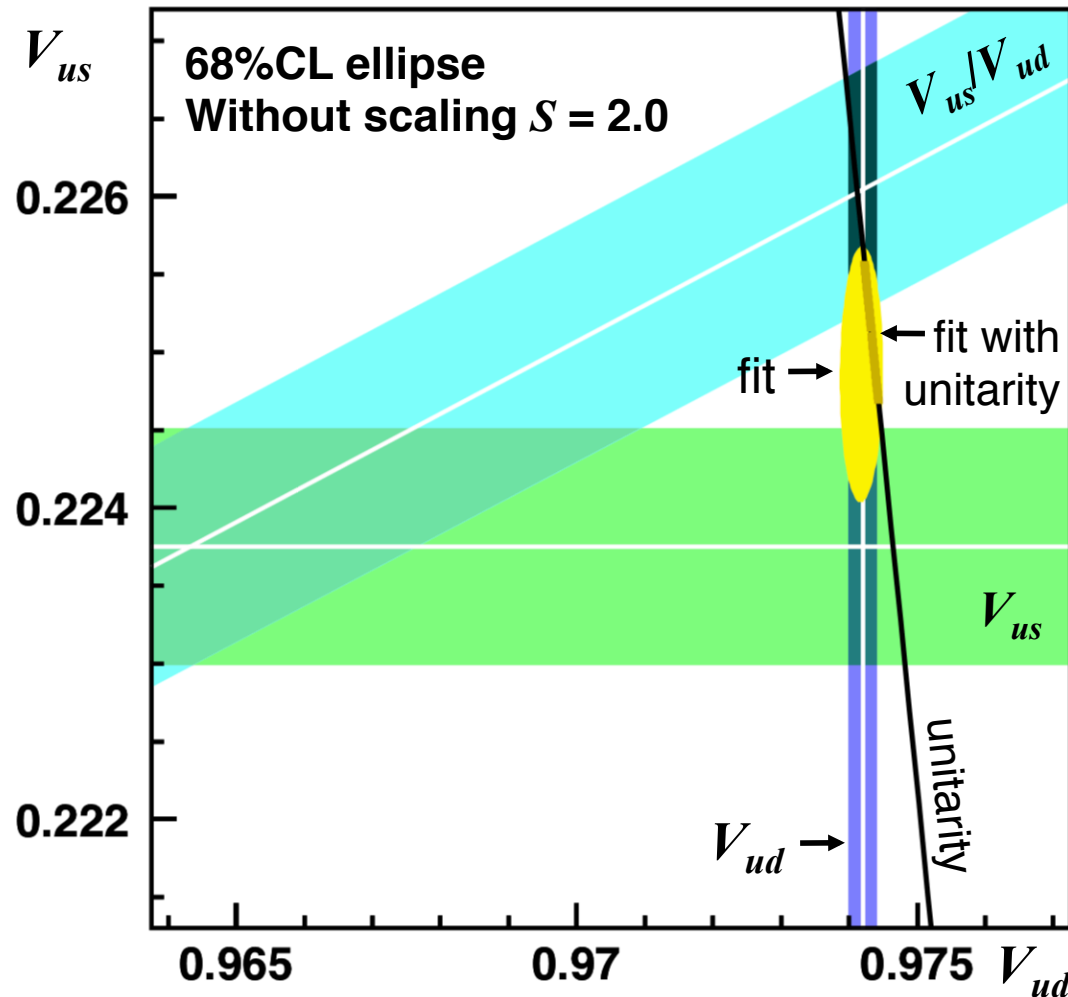
Gasser, Leutwyler '85

Dispersive representation: $f_0(t_{\text{CT}}) \equiv C$



4.2 V_{us} and CKM unitarity: All data, $N_f=2+1$

$N_f = 2+1$: Fit to results for $|V_{ud}|$, $|V_{us}|$, $|V_{us}|/|V_{ud}|$
 $f_+(0) = 0.9677(27)$, $f_K/f_\pi = 1.1927(38)$



$$\begin{aligned} |V_{ud}| &= 0.97420(21) \\ |V_{us}| &= 0.2238(8) \\ |V_{us}|/|V_{ud}| &= 0.2320(8) \end{aligned}$$

Fit results, no constraint

$$\begin{aligned} V_{ud} &= 0.97418(21) \\ V_{us} &= 0.2249(5) \\ \chi^2/\text{ndf} &= 4.5/1 \quad (3.5\%) \\ \Delta_{\text{CKM}} &= -0.0004(5) \\ &\quad -0.8\sigma \end{aligned}$$

With scale factor $S = 2.1$

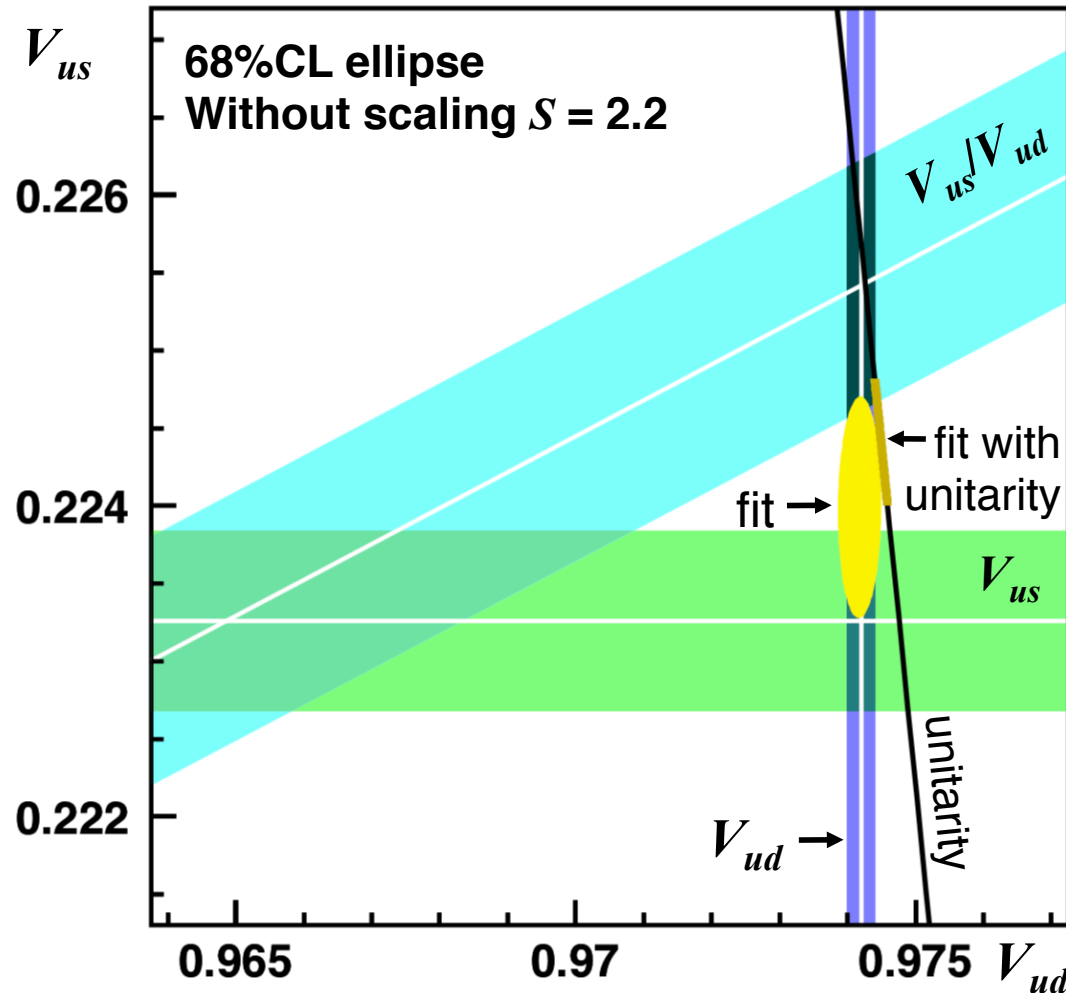
$$\begin{aligned} V_{ud} &= 0.97418(45) \\ V_{us} &= 0.2249(12) \end{aligned}$$

4.2 V_{us} and CKM unitarity: All data, $N_f=2+1+1$

$N_f = 2+1+1$: Fit to results for $|V_{ud}|$, $|V_{us}|$, $|V_{us}|/|V_{ud}|$
 $f_+(0) = 0.9698(17)$, $f_K/f_\pi = 1.1960(40)$



$$\begin{aligned} |V_{ud}| &= 0.97420(21) \\ |V_{us}| &= 0.2233(6) \\ |V_{us}|/|V_{ud}| &= 0.2314(8) \end{aligned}$$



Fit results, no constraint

$$\begin{aligned} V_{ud} &= 0.97418(21) \\ V_{us} &= 0.2240(5) \\ \chi^2/\text{ndf} &= 4.7/1 \quad (3.1\%) \\ \Delta_{\text{CKM}} &= -0.0008(5) \\ &\quad -1.7\sigma \end{aligned}$$

With scale factor $S = 2.2$

$$\begin{aligned} V_{ud} &= 0.97418(46) \\ V_{us} &= 0.2240(10) \end{aligned}$$

5. Conclusion and outlook

5.1 Preliminary conclusions

Experimental results

$$|V_{us}|f_+(0) = 0.21654(41)$$
$$|V_{us}/V_{ud}| \times f_{K^\pm}/f_{\pi^\pm} = 0.27599(37)$$

With $N_f = 2+1+1$ lattice inputs

$$V_{ud} = 0.97418(21) \pm 0.02\%$$
$$V_{us} = 0.2246(5) \pm 0.22\%$$
$$\Delta_{\text{CKM}} = -0.0005(5) = -1.1\sigma$$

Good agreement with unitarity for $K_{\ell 2}$

Previous excellent consistency for $K_{\ell 3}$ no longer observed

- Change occurred after 2014-era more precise evaluations of $f_+(0)$
- Experimental results for $K_{\ell 3}$ have changed little since 2010

Are residual systematics in the data and/or calculations becoming important as stated uncertainties shrink?

- Evaluation of $|V_{us}|f_+(0)$ from $K_{\ell 3}$ data set based on some creaky BR fits, but errors are scaled and consistency between modes is good (K_L , K_S , K^\pm)
- Lots of redundancy in $K_{\ell 3}$ data set. Adding or eliminating individual measurements doesn't change $|V_{us}|f_+(0)$ much.

5.2 Prospects for new measurements

NA48/2



NA62

Can measure BRs and form-factor parameters for K^+

NA48/2 (2003-2004) recently measured $K_{\ell 3}$ form factors

NA62-RK (2007) has O(10M) $K_{\ell 3}$ decays

NA62 has O(few M) K_{e3} from minimum bias runs (2015-16)

Relative to NA48/2, NA62 has

- Better particle identification π/μ
- Better systematics for t reconstruction:
 - full beam tracking, better σ_p in spectrometer

ISTRA+



OKA

Fixed target experiment at U-70 (Protvino), like ISTRA+

- New beamline with RF-separated K^+ beam

Can measure BRs and form-factor parameters

- Need more analysis of systematics for K_{e3} form factors

Runs from 2010-2013: $\sim 17\text{M}$ K_{e3}^+ events

- Additional runs in 2016-2018; more planned in future

5.2 Prospects for new measurements

KLOE



KLOE-2

Can measure all observables: BRs, τ s, FFs: K^\pm, K_L, K_S

5.5 fb⁻¹ of data from KLOE-2 running (2015-2018)

- +2 fb⁻¹ of original KLOE data not yet analyzed for V_{us}

Measurements that can be improved with KLOE-2 statistics:

- K_S BRs ($K_S \rightarrow \pi e \nu$, but also $K_S \rightarrow \pi \mu \nu$)

See e.g. KLOE-2 measurement of A_S 1806.08654

70k $K_S \rightarrow \pi e \nu$ decays

- K^\pm, K_L form factors (particularly $K_{\ell 3}$), K_L mean life?

LHCb

Proven capability to measure K_S decays to muons

- 10¹³ K_S /fb⁻¹ produced
- EPJC 77 (2017): BR($K_S \rightarrow \mu \mu$) < 1.0 × 10⁻⁹ 95%CL

Limited by hardware trigger efficiency ($\epsilon_{\text{trig}} \sim 1\%$)

Can LHCb measure BR($K_S \rightarrow \pi \mu \nu$) to < 1% in Run III?

- Would require dedicated software HLT line

$K_S \rightarrow \pi \mu \nu$ never yet measured – a new channel for V_{us}

- τ_S known to 0.04% (vs 0.41% for τ_L , 0.12% for τ_\pm)

5.2 Prospects for new measurements

KEK-246



TREK E36

Primary focus is $\text{BR}(K_{e2}/K_{\mu2})$ to 0.25%

+ Invisible heavy neutrino searches

+ T violation in $K_{\mu3}$ (as E06)

Upgraded KEK-246 setup, moved to J-PARC

- Stopped K^+ in active target
- Toroidal spectrometer surrounding target
- e/μ particle ID by time of flight, Cerenkov counters, lead-glass calorimetry

KEK-246 measured $\text{BR}(K_{\mu3}/K_{e3})$ and K_{e3} FF, so TREK could potentially measure at least some BRs and FFs of interest for V_{us}

5.3 Progress on V_{us} from kaons: Final notes

- **$K_{\ell 3}$ FFs do not directly contribute significantly to uncertainty on V_{us}**
 - However, uncertainties on high-statistics BR ratio measurements may be so low that FFs become a major systematic
 - e.g. $\text{BR}(K_{\mu 3}/\pi\pi^0)$, $\text{BR}(K_{\mu 3}/K_{e 3})$
- **Uncertainties from parameterization of $K\pi$ phase shift data now limit precision for $K_{\ell 3}$ FFs and phase space integrals**
 - Better parameterization will require old data to be re-fit!
 - Imperative for future averages that experiments publish full FF data so that it can be re-fit as parameterizations improve
 - Direct lattice calculation of $K_{\ell 3}$ FFs may help
- **For K^\pm , normalization BRs have significant uncertainties**
 - Effect of any precise new $\text{BR}(K_{e 3}/\pi\pi^0)$ results will be limited by uncertainty on $\text{BR}(\pi\pi^0)$
 - Very important to measure absolute BRs or ratios involving BRs of other modes, e.g. $\pi\pi^0/\mu\nu$, $\pi\pi\pi/\pi\pi^0$, $\pi\pi\pi/\mu\nu$

5.4 Summary and conclusions

Experimental results from kaons

$$|V_{us}|f_+(0) = 0.21652(41)$$

$$|V_{us}/V_{ud}| \times f_{K^\pm}/f_{\pi^\pm} = 0.27599(37)$$

With $|V_{ud}|(0^+ \rightarrow 0^+)$ and $N_f = 2+1+1$ lattice

$$\Delta_{\text{CKM}} = -0.0011(5) = -2.2\sigma$$

$$\Delta_{\text{CKM}} = -0.0001(6) = -0.2\sigma$$

2σ inconsistency between $K_{\ell 3}$ and $K_{\ell 2}$ results for V_{us}

- $K_{\ell 2}$ result shows good agreement with unitarity and V_{ud}
- $K_{\ell 3}$ result 2σ smaller than expected from unitarity and V_{ud}
 - Change occurred after 2014-era more precise evaluations of $f_+(0)$
 - Experimental result for $|V_{us}|f_+(0)$ has changed little since 2010

Continuing to see impressive progress on the lattice

- Not only $f_+(0)$ and f_{K^\pm}/f_{π^\pm} , but also full t -dependence of FFs, EM corrections, etc.

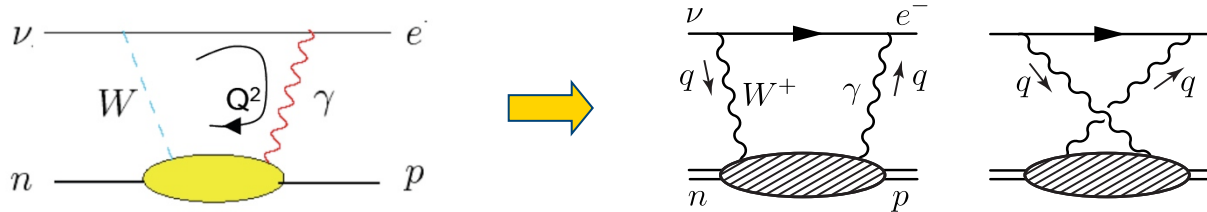
Good prospects for new round of measurements to reduce uncertainty on $|V_{us}|f_+(0)$ from current 0.18% to $\sim 0.12\%$ within next few years:

NA62, OKA, KLOE-2, LHCb, TREK...

6. Back-up

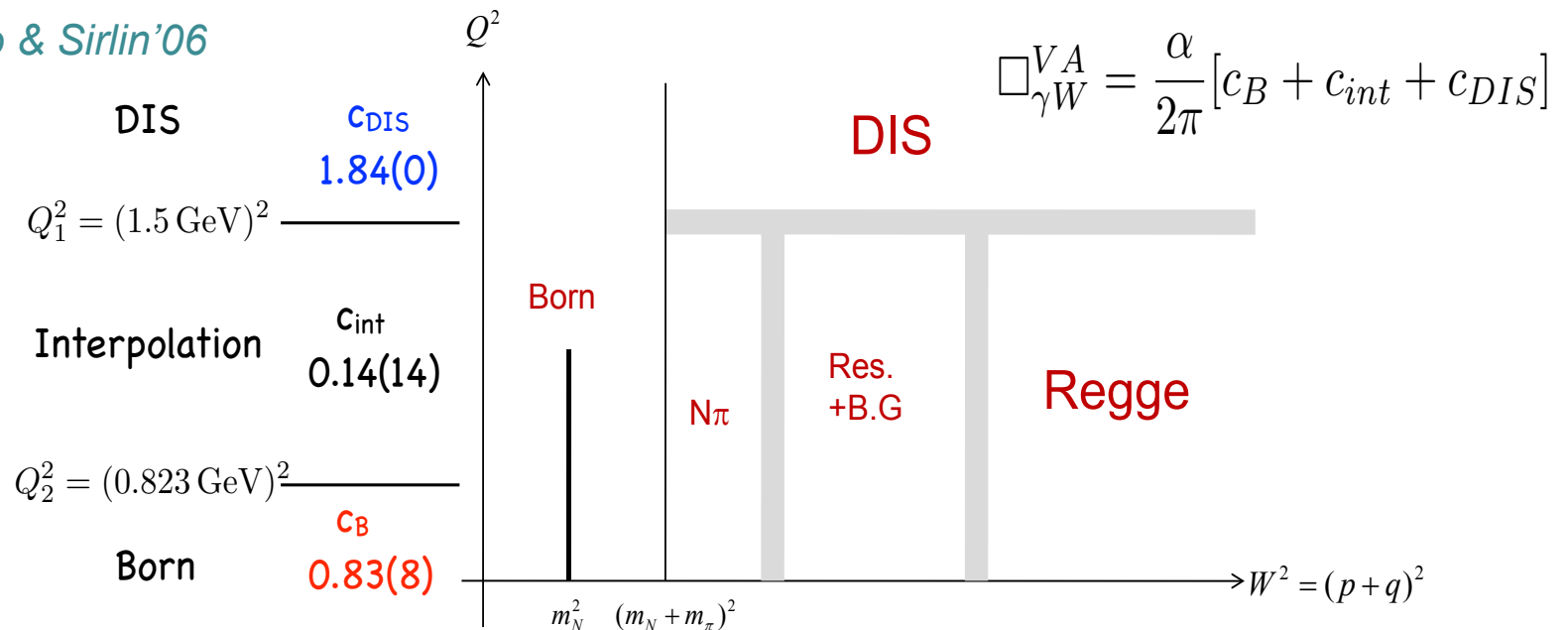
New Radiative Corrections for free neutron

Gorchtein@CIPANP'18



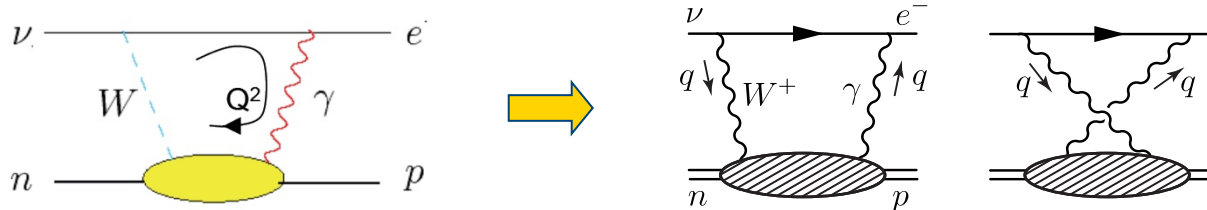
$$\square_{\gamma W}^{VA(0)} = \frac{\alpha}{\pi M} \int_0^\infty \frac{dQ^2 M_W^2}{M_W^2 + Q^2} \int_0^\infty d\nu \frac{(\nu + 2q)}{\nu(\nu + q)^2} F_3^{(0)}(\nu, Q^2)$$

Marciano & Sirlin'06



New Radiative Corrections for free neutron

Gorshtein@CIPANP'18



$$\square_{\gamma W}^{VA(0)} = \frac{\alpha}{\pi M} \int_0^\infty \frac{dQ^2 M_W^2}{M_W^2 + Q^2} \int_0^\infty d\nu \frac{(\nu + 2q)}{\nu(\nu + q)^2} F_3^{(0)}(\nu, Q^2)$$

Marciano & Sirlin'06 $\square_{\gamma W}^{VA} = \frac{\alpha}{2\pi} [c_B + c_{int} + c_{DIS}] = \frac{\alpha}{2\pi} [0.83(8) + 0.14(14) + 1.84(0)]$

$$\square_{\gamma W}^{MS} = \frac{\alpha}{2\pi} 2.79(17) = 3.24(20) \times 10^{-3}$$

New evaluation: *Seng, Gorshtein, Patel & Ramsey-Musolf arXiv:1807.10197*

$$\square_{\gamma W}^{VA} = \frac{\alpha}{2\pi} [c_B + c_{piN} + c_{Res} + c_{Regge} + c_{DIS}] = \frac{\alpha}{2\pi} [0.91(5) + 0.044(5) + 0.01(1) + 0.238(14) + 1.84(0)]$$

$$\square_{\gamma W}^{New} = \frac{\alpha}{2\pi} 3.03(5) = 3.51(6) \times 10^{-3}$$



V_{ud} from free n: **$\sim 1\sigma$ smaller**

2.6 $K_{\ell 3}$ data and lepton universality

- For each state of kaon charge, evaluate:

$$r_{\mu e} = \frac{(R_{\mu e})_{\text{obs}}}{(R_{\mu e})_{\text{SM}}} = \frac{\Gamma_{\mu 3}}{\Gamma_{e 3}} \cdot \frac{I_{e 3} (1 + \delta_{e 3})}{I_{\mu 3} (1 + \delta_{\mu 3})} = \frac{[|V_{us}| f_+(0)]_{\mu 3, \text{obs}}^2}{[|V_{us}| f_+(0)]_{e 3, \text{obs}}^2} = \frac{g_{\mu}^2}{g_e^2}$$

Modes	2004 BRs ^{*,†}	Current [†]
K_L	1.054(14)	1.003(5)
K^{\pm}	1.014(12)	0.999(9)
Avg	1.030(9)	1.002(5)

← Was 0.998(9)
for 2010

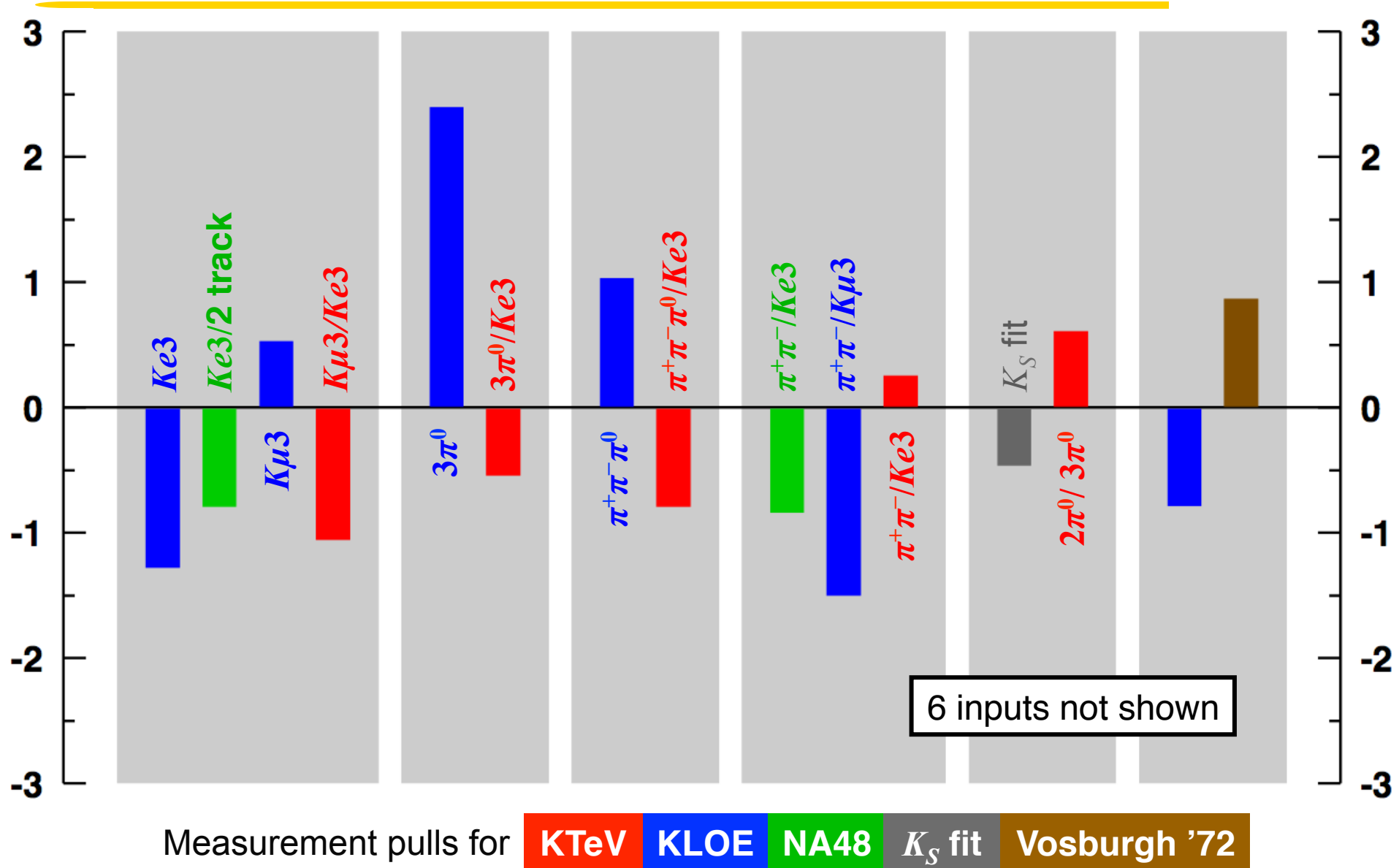
*Assuming current values for form-factor parameters and Δ^{EM} $^{\dagger}K_S$ not included

- Compare to other precise tests:

$\pi \rightarrow \ell \nu$ **$(r_{\mu e}) = 1.0020(19)$**
PDG '16 with PIENU '15 result

$\tau \rightarrow \ell \nu \nu$ **$(r_{\mu e}) = 1.0038(28)$**
HFLAV May '17 web update

Comparison: K_L fit result vs. input data



Updated fit to K^\pm rate data

17 input measurements:

3 old τ values in PDG

KLOE τ

KLOE BR $\mu\nu, \pi\pi^0$

KLOE BR $K_{e3}, K_{\mu3}$

with dependence on τ

NA48/2 BR $K_{e3}/\pi\pi^0, K_{\mu3}/\pi\pi^0$

E865 BR $K_{e3}/K\text{Dal}$

3 old BR $\pi\pi^0/\mu\nu$

KEK-246 $K_{\mu3}/K_{e3}$

KLOE BR $\pi\pi\pi, \pi\pi^0\pi^0$

(**Bisi '65** BR $\pi\pi^0\pi^0/\pi\pi\pi$ removed)

1 constraint: $\Sigma \text{BR} = 1$

Much more selective than PDG fit
PDG '16: 35 inputs, 8 parameters

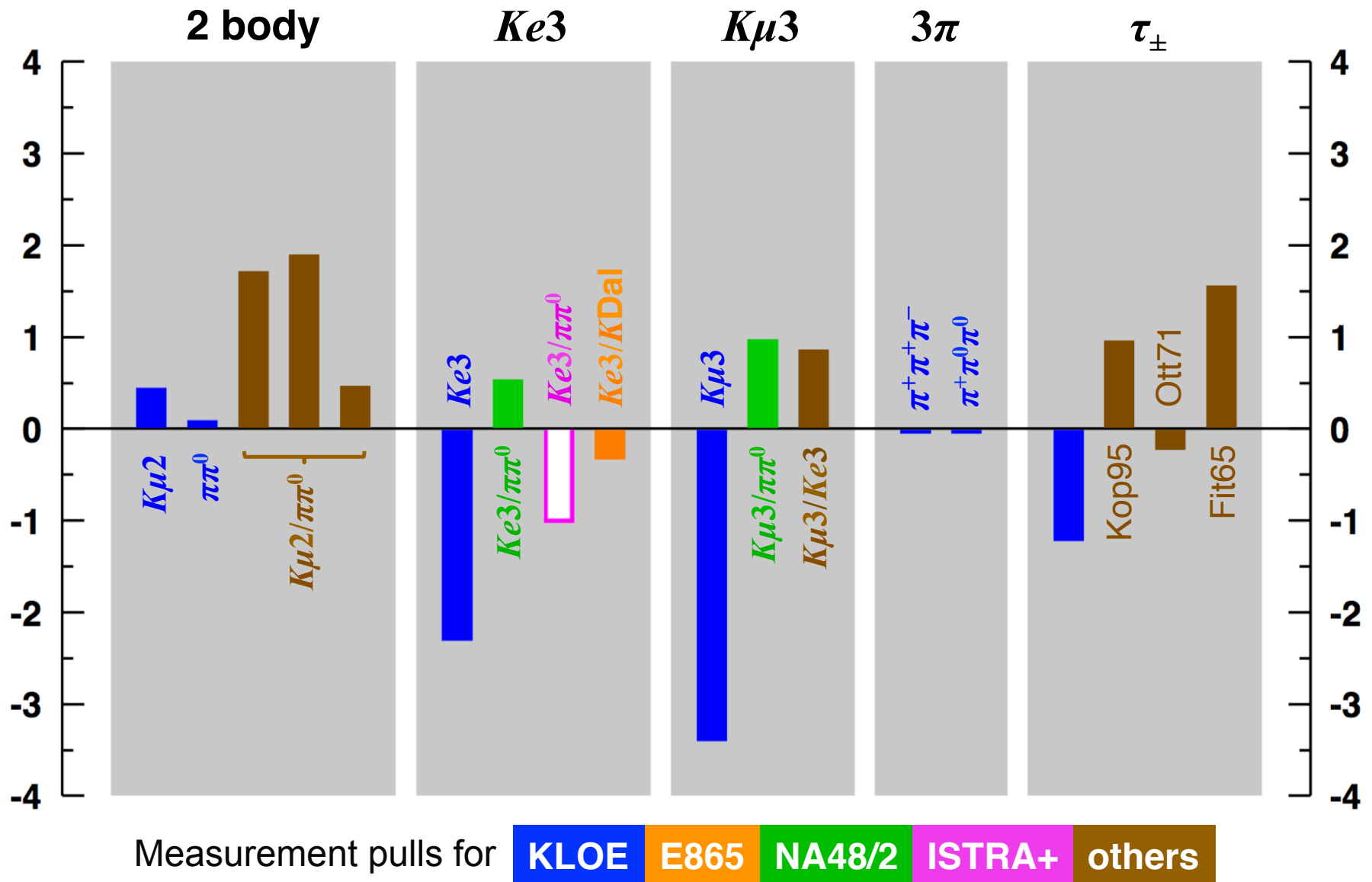
Parameter	Value	S
BR($\mu\nu$)	63.58(11)%	1.1
BR($\pi\pi^0$)	20.64(7)%	1.1
BR($\pi\pi\pi$)	5.56(4)%	1.0
BR(K_{e3})	5.088(27)%	1.2
BR($K_{\mu3}$)	3.366(30)%	1.9
BR($\pi\pi^0\pi^0$)	1.764(25)%	1.0
τ_\pm	12.384(15) ns	1.2

$\chi^2/\text{ndf} = 25.5/11$ (Prob = 0.78%)
compare PDG '16: 53/28 (0.26%)

With **ISTRA+ '14** BR($K_{e3}^-/\pi^-\pi^0$)

- **BR(K_{e3}) = 5.083(27)%**
- Negligible changes in other parameters, fit quality

Comparison: K^\pm fit result vs. input data



$|V_{us}|(K_{\ell 3})$ and $|V_{ud}|(0^+ \rightarrow 0^+)$: 2010

Moulson@CKM2014

$$|V_{us}| f_+(0) = 0.2163(5)$$

$$f_+(0) = 0.959(5)$$

$$|V_{us}| = 0.2254(13)$$

Hardy & Towner '10

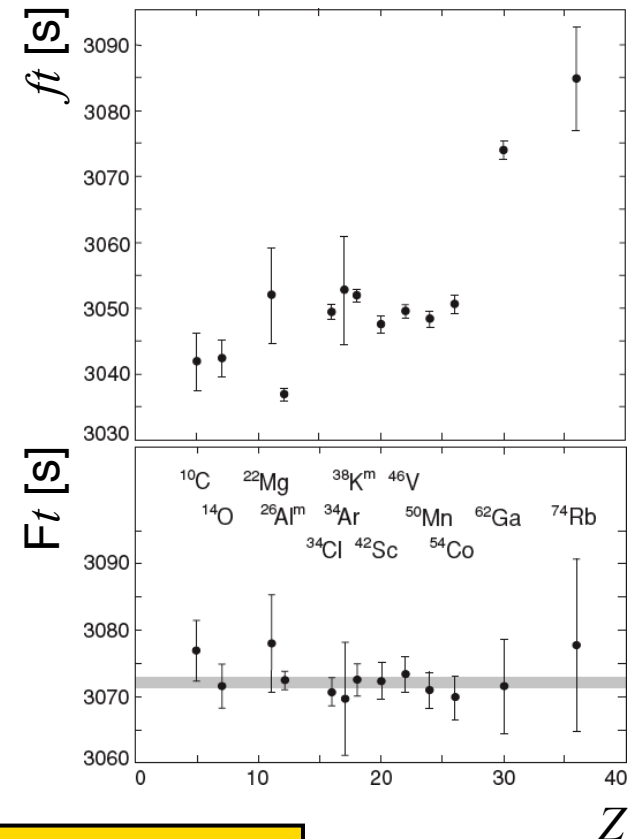
$$|V_{ud}| = 0.97425(22)$$

Survey of 150 measurements of 13 different $0^+ \rightarrow 0^+ \beta$ decays

27 new ft measurements including Penning-trap measurements for Q_{EC}

Some old measurements dropped
Improved EW radiative corrections
[Marciano & Sirlin '06]

New $SU(2)$ -breaking corrections
[Towner & Hardy '08]



$$\Delta_{\text{CKM}} = V_{ud}^2 + V_{us}^2 - 1 = +0.0000(8)$$

Exact compatibility with unitarity

$|V_{us}|(K_{\ell 3})$ and $|V_{ud}|(0^+ \rightarrow 0^+)$: Update

M. Moulson@CKM'16

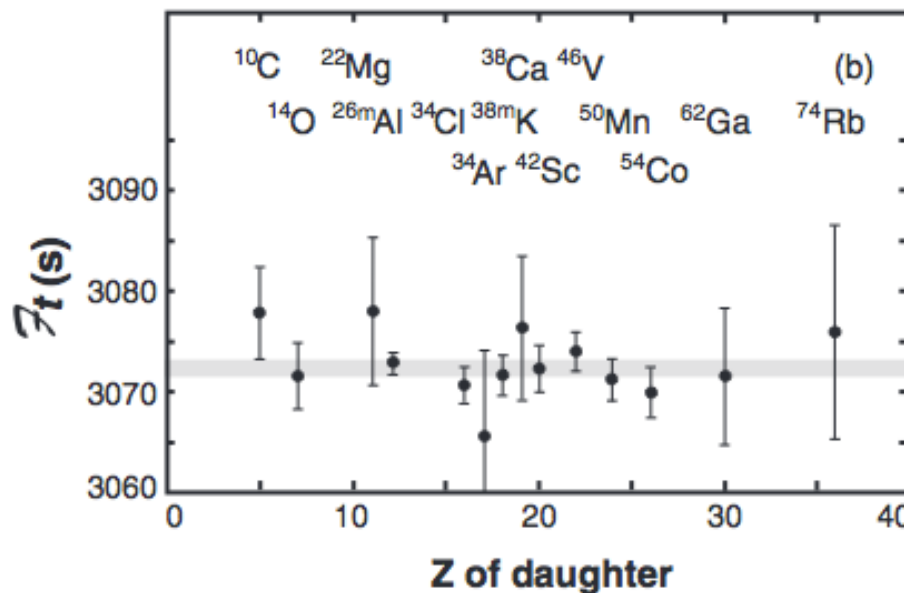
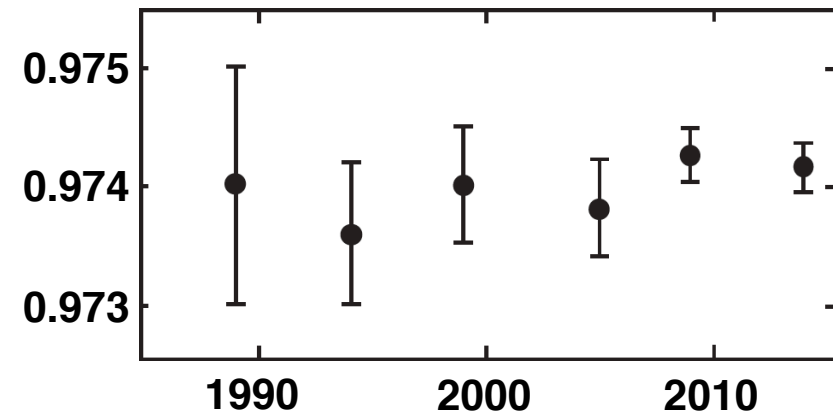
Hardy, CKM '16 preliminary

$|V_{ud}| = 0.97420(21)$

- 24 new measurements
- Critical review of IB correction schemes as per PRC 82 (2010)
- Rejection of results with IB corrections giving results in conflict with CVC

V_{ud} vs analysis year

Courtesy of J. Hardy

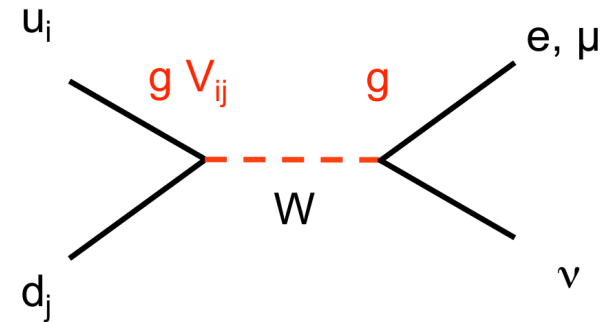


Hardy & Towner'15

1.2 Paths to V_{ud} and V_{us}

- From kaon, pion, baryon and nuclear decays

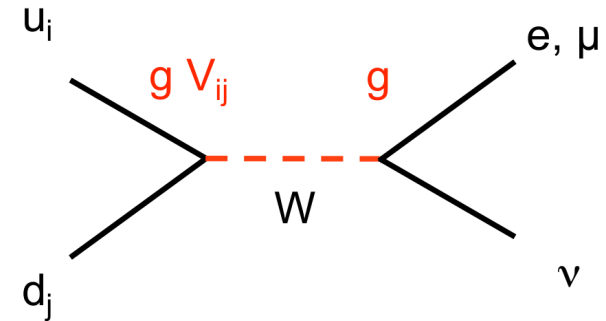
V_{ud}	$0^+ \rightarrow 0^+$ $\pi^\pm \rightarrow \pi^0 e \nu_e$	$n \rightarrow p e \bar{\nu}_e$	$\pi \rightarrow l \nu_l$
V_{us}	$K \rightarrow \pi l \nu_l$	$\Lambda \rightarrow p e \bar{\nu}_e$	$K \rightarrow l \nu_l$



1.2 Paths to V_{ud} and V_{us}

- From kaon, pion, baryon and nuclear decays

V_{ud}	$0^+ \rightarrow 0^+$ $\pi^\pm \rightarrow \pi^0 e \bar{\nu}_e$	$n \rightarrow p e \bar{\nu}_e$	$\pi \rightarrow l \nu_l$
V_{us}	$K \rightarrow \pi l \nu_l$	$\Lambda \rightarrow p e \bar{\nu}_e$	$K \rightarrow l \nu_l$



- These are the *golden modes* to extract V_{ud} and V_{us}
 - Only the *vector current* contributes $\langle A(p_A) | \bar{q}^i \gamma_\mu q^j | B(p_B) \rangle$
 - Normalization known in SU(2) [SU(3)] symmetry limit
 - Corrections start at 2nd order in SU(2) [SU(3)] breaking

Ademollo & Gato, Berhands & Sirlin

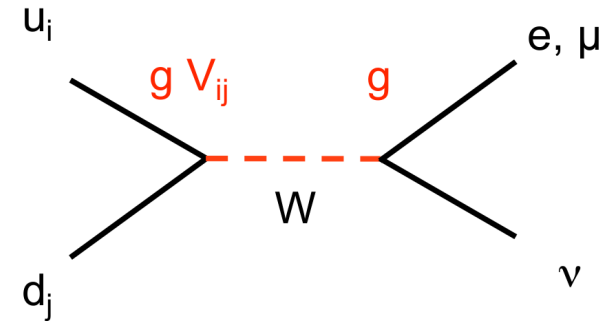
- Currently the most precise determination of V_{ud} and V_{us}

➡ V_{ud} (0.02 %) and V_{us} (0.5 %)

1.2 Paths to V_{ud} and V_{us}

- From kaon, pion, baryon and nuclear decays

V_{ud}	$0^+ \rightarrow 0^+$ $\pi^\pm \rightarrow \pi^0 e \nu_e$	$n \rightarrow p e \bar{\nu}_e$	$\pi \rightarrow l \nu_l$
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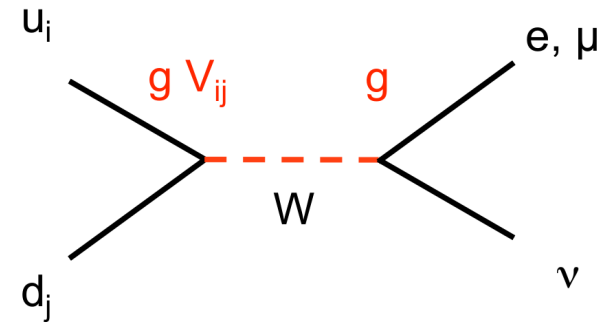


- Both V and A currents contribute: need experimental input on $\langle A \rangle$ (e.g. β -asymmetry)
- Free of nuclear structure uncertainties
- Probe different combinations of BSM operators

1.2 Paths to V_{ud} and V_{us}

- From kaon, pion, baryon and nuclear decays

V_{ud}	$0^+ \rightarrow 0^+$ $\pi^\pm \rightarrow \pi^0 e \bar{\nu}_e$	$n \rightarrow p e \bar{\nu}_e$	$\pi \rightarrow l \nu_l$
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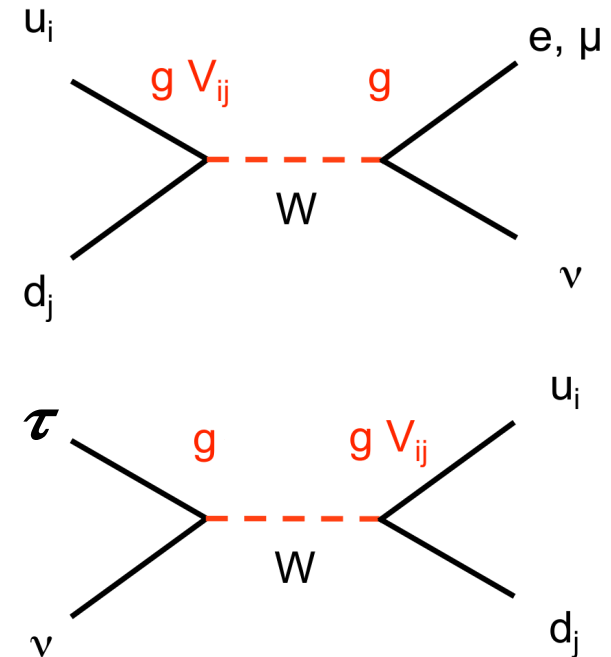


- K_{l2}/π_{l2}
 - Only the *axial current* contributes
 - Need to know the decay constants F_K, F_π
➡ *Lattice QCD*
 - Probe different BSM operators than from the vector case
- Input on F_K/F_π ➡ V_{us}/V_{ud} very precisely

1.2 Paths to V_{ud} and V_{us}

- From kaon, pion, baryon and nuclear decays

V_{ud}	$0^+ \rightarrow 0^+$ $\pi^\pm \rightarrow \pi^0 e \nu_e$	$n \rightarrow p e \bar{\nu}_e$	$\pi \rightarrow l \nu_l$
V_{us}	$K \rightarrow \pi l \nu_l$	$\Lambda \rightarrow p e \bar{\nu}_e$	$K \rightarrow l \nu_l$



- From τ decays (crossed channel) + test of LU

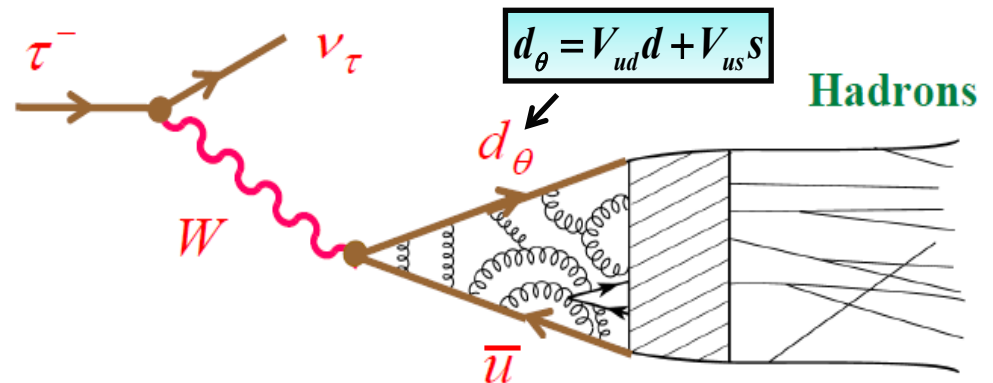
V_{ud}	$\tau \rightarrow \pi \pi \nu_\tau$		$\tau \rightarrow \pi \nu_\tau$	$\tau \rightarrow h_{NS} \nu_\tau$
V_{us}	$\tau \rightarrow K \pi \nu_\tau$		$\tau \rightarrow K \nu_\tau$	$\tau \rightarrow h_S \nu_\tau$ (inclusive)

Paths to V_{ud} and V_{us}

- From τ decays (crossed channel)

V_{ud}	$\tau \rightarrow \pi\pi\nu_\tau$		$\tau \rightarrow \pi\nu_\tau$	$\tau \rightarrow h_{NS}\nu_\tau$
V_{us}	$\tau \rightarrow K\pi\nu_\tau$		$\tau \rightarrow K\nu_\tau$	$\tau \rightarrow h_S\nu_\tau$ (inclusive)

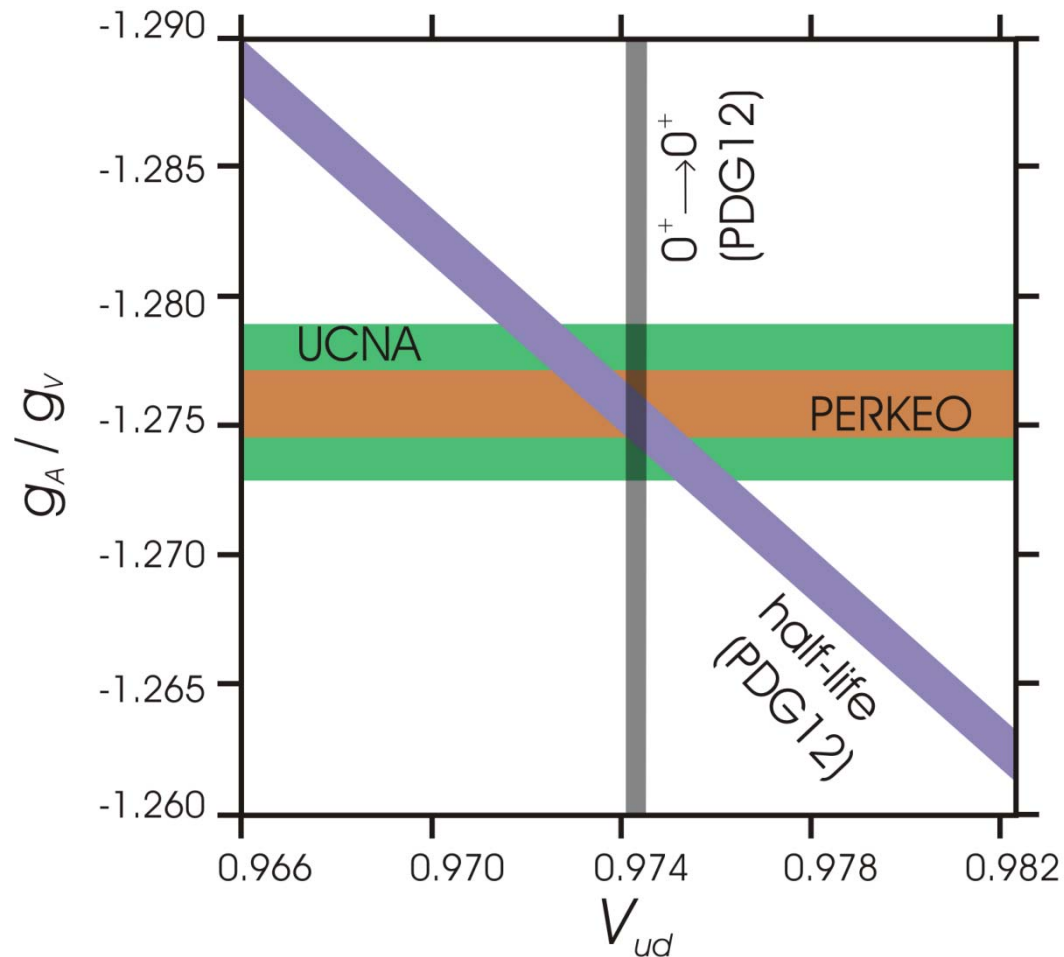
- Use OPE to calculate inclusive BRs
- Information from exclusive modes too



Extraction of V_{ud}

- See also V_{ud} extraction from neutron decay

Mendenhall et al.'13



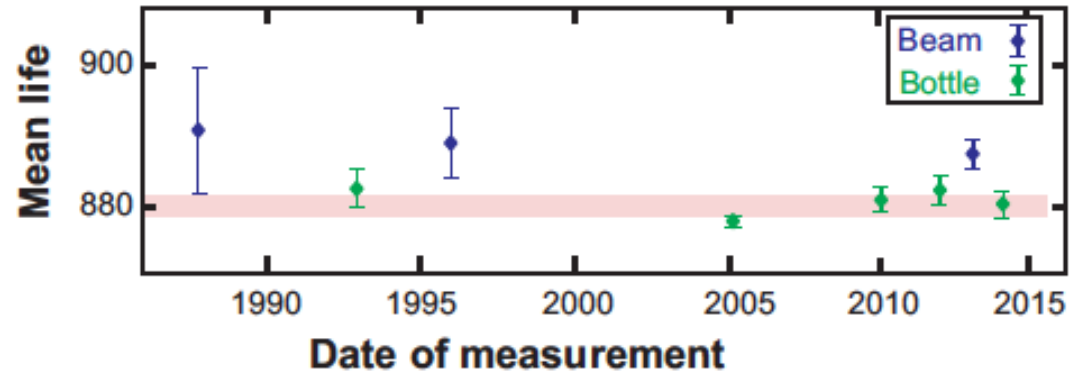
Extraction of V_{ud}

Mean life:

$$\tau = 880.2 \pm 1.0 \text{ s}$$

$$\chi^2/N = 3.7$$

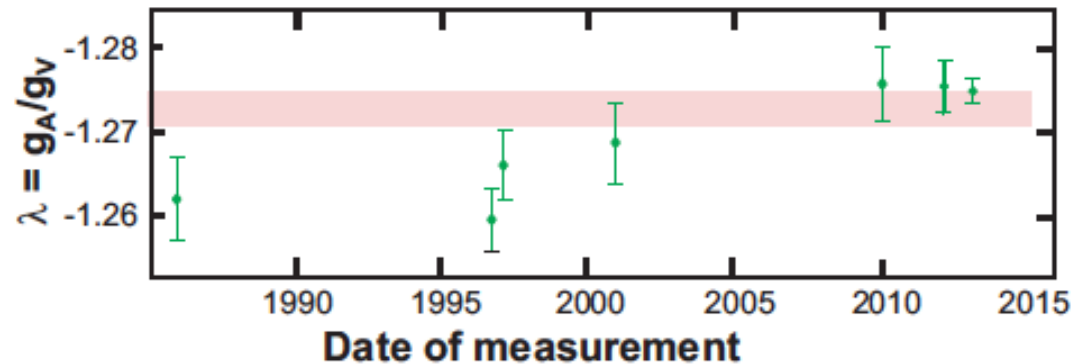
Beam: $888.1 \pm 2.0 \text{ s}$
Bottle: $879.6 \pm 0.7 \text{ s}$



β asymmetry:

$$\lambda = -1.2725 \pm 0.0020$$

$$\chi^2/N = 4.1$$



$$V_{ud} = 0.9757 \pm 0.0014$$

Beam-bottle span
 $0.9701 \leq V_{ud} \leq 0.9767$

nuclear $0^+ \rightarrow 0^+$
 $V_{ud} = 0.9742 \pm 0.0002$

4. Implication of Cabibbo universality tests for new physics

4.1 Looking for New Physics using Δ_{CKM}

- Δ_{CKM} a constraining quantity: $\Delta_{CKM} = 1 - (|V_{ud}|^2 + |V_{us}|^2)$

$$\frac{1 - |V_{uD}|^2}{1 - |U_{\mu N}|^2} \cdot \frac{1}{1 + BR_{\text{exotic}}^{\mu}} \cdot \frac{[G_F^{(\beta)}]^2}{[G_F^{(\mu)}]^2} = 1 + \Delta_{CKM}$$

Heavy fermion
mixing

Exotic
muon decays

Gauge
universality
violations

$$|V_{uD}| \leq 0.03$$

$$|U_{\mu N}| \leq 0.03$$

95% C.L.

$$BR_{\text{exotic}}^{\mu} < 0.001$$

95% C.L.

Stronger than direct limits

$$BR(\mu^+ \rightarrow e^+ \bar{\nu}_e \nu_{\mu}) < 0.012$$

Constraints on TeV
scale SM extensions

4.1 Looking for New Physics using Δ_{CKM}

- Effective Theory approach:

$$\mathcal{L} = \mathcal{L}_{SM} + \frac{C^{(5)}}{\Lambda} \mathcal{O}^{(5)} + \sum_i \frac{C_i^{(6)}}{\Lambda^2} \mathcal{O}_i^{(6)} + \dots$$

- Δ_{CKM} a constraining quantity:

$$\boxed{|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 1 + \Delta_{\text{CKM}}}$$

4.1 Looking for New Physics using Δ_{CKM}

Operator		Observable	$K^+ \rightarrow \pi^+ \nu \bar{\nu}$	$K_L \rightarrow \pi^0 \nu \bar{\nu}$	$K_L \rightarrow \pi^0 \ell^+ \ell^-$	$K_L \rightarrow \ell^+ \ell^-$	$K^+ \rightarrow \ell^+ \nu$	$P_T(K^+ \rightarrow \pi^0 \mu^+ \nu)$	Δ_{CKM}	ϵ'/ϵ	ϵ_K	from: SJ, talk at NA62 Handbook workshop 2009
												in MSSM?
$O_{lq}^{(1)}$	$(\bar{D}_L \gamma^\mu S_L)(\bar{L}_L \gamma_\mu L_L)$	✓	✓	✓	hs	—	—	—	—	—	—	✓
$O_{lq}^{(3)}$	$(\bar{D}_L \gamma^\mu \sigma^i S_L)(\bar{L}_L \gamma_\mu \sigma^i L_L)$	✓	✓	✓	hs	hs	✓	✓	—	—	—	✓
O_{qe}	$(\bar{D}_L \gamma^\mu S_L)(\bar{l}_R \gamma_\mu l_R)$	—	—	✓	hs	—	—	—	—	—	—	small
O_{ld}	$(\bar{d}_R \gamma^\mu s_R)(\bar{L}_L \gamma_\mu L_L)$	✓	✓	✓	hs	—	—	—	—	—	—	small
O_{ed}	$(\bar{d}_R \gamma^\mu s_R)(\bar{l}_R \gamma_\mu l_R)$	—	—	✓	hs	—	—	—	—	—	—	small
O_{lq}^\dagger	$(\bar{u}_R S_L) \cdot (\bar{l}_R L_L)$	—	—	—	—	✓	✓	✓	—	—	—	tiny
$(O_{lq}^t)^\dagger$	$(\bar{u}_R \sigma_{\mu\nu} S_L) \cdot (\bar{l}_R \sigma^{\mu\nu} L_L)$	—	—	—	—	—	?	?	—	—	—	tiny
O_{qde}	$(\bar{d}_R S_L)(\bar{L}_L l_R)$	—	—	✓	✓	—	—	—	—	—	—	tiny
O_{qde}^\dagger	$(\bar{D}_L s_R)(\bar{l}_R L_L)$	—	—	✓	✓	✓	✓	✓	—	—	—	large $\tan \beta$
$O_{\varphi q}^{(1)}$	$(\bar{D}_L \gamma^\mu S_L)(H^\dagger D_\mu H)$	✓	✓	✓	hs	—	—	—	✓	(✓)	(✓)	✓
$O_{\varphi q}^{(3)}$	$(\bar{D}_L \gamma^\mu \sigma^i S_L)(H^\dagger D_\mu \sigma^i H)$	✓	✓	✓	hs	hs	✓	✓	✓	(✓)	(✓)	✓
$O_{\varphi d}$	$(\bar{d}_R \gamma^\mu s_R)(H^\dagger D_\mu H)$	✓	✓	✓	hs	—	—	—	✓	(✓)	(✓)	large $\tan \beta$ (non-MFV)

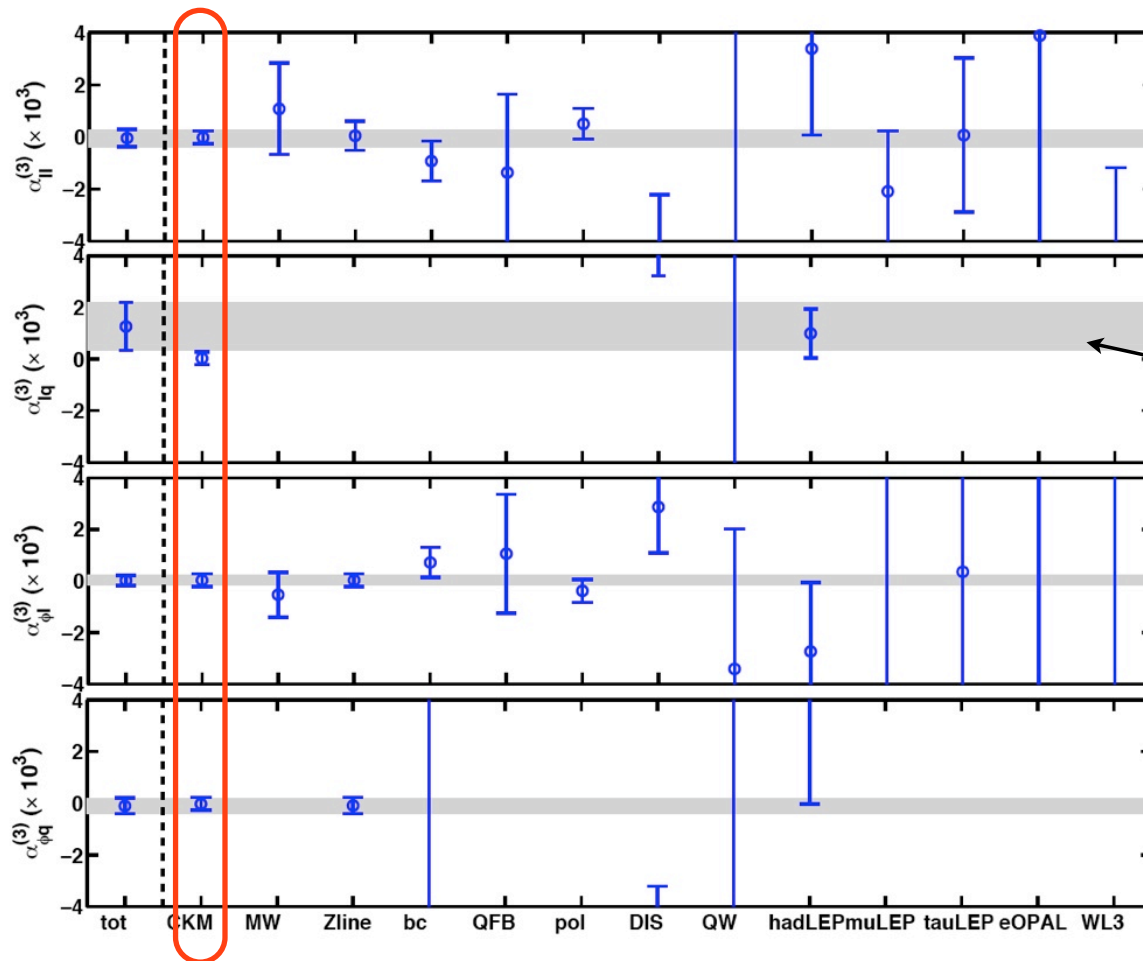
4.1 Looking for New Physics using Δ_{CKM}

- Δ_{CKM} is sensitive to 4 fermion operators:

Cirigliano, Gonzalez-Alonso & Jenkins'09

2) What is the strength of Δ_{CKM} constraint? Same level or better than Z-pole observables (effective scale $\Lambda > 11 \text{ TeV}$ @ 90% CL)

$$\hat{\alpha}_X = \frac{v^2}{\Lambda_X^2}$$



$O_{lq}^{(3)}$

Dramatic improvement over LEP2 and APV

Deviations as large as $\Delta_{\text{CKM}} \sim 0.01$ could be blamed on this operator

4.2 Looking for New Physics with K_{l2} and K_{l3}

- Callan-Treiman theorem:

Bernard, Oertel, E.P., Stern'06, '08

$$C = \frac{\overline{f}_0(\Delta_{K\pi})}{m_K^2 - m_\pi^2} = \frac{F_K}{F_\pi f_+(0)} + \Delta_{CT} = \underbrace{\frac{F_K |V^{us}|}{F_\pi |V^{ud}|} \frac{1}{f_+(0) |V^{us}|} |V^{ud}|}_{\text{Very precisely known from Br}(Kl2/\pi l2), \Gamma(Ke3) \text{ and } |V_{ud}|} r + \Delta_{CT}$$

Very precisely known from $\text{Br}(Kl2/\pi l2)$, $\Gamma(Ke3)$ and $|V_{ud}|$

$$B_{\text{exp}} = 1.2446(41)$$

- In the Standard Model : $r = 1$ $(\ln C_{SM} = 0.2141(73))$ $\Delta_{CT} = (-3.5 \pm 8) \cdot 10^{-3}$

- In presence of new physics, new couplings : $r \neq 1$

NLO value + large error bars in agreement with

Bijnens&Ghorbani'07
Kastner & Neufeld'08

Experiment $K_{e3}+K_{\mu 3}$	$\ln C$
NA48'07 ($K_{\mu 3}$ alone)	0.144(14)
KLOE'08	0.204(25)
KTeV'10	0.192(12)
NA48 (preliminary)	?

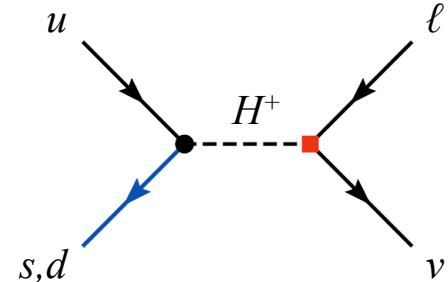
4.3 Constraints on 2HDM

- Ex: Constraints on the aligned 2-Higgs-doublet model:

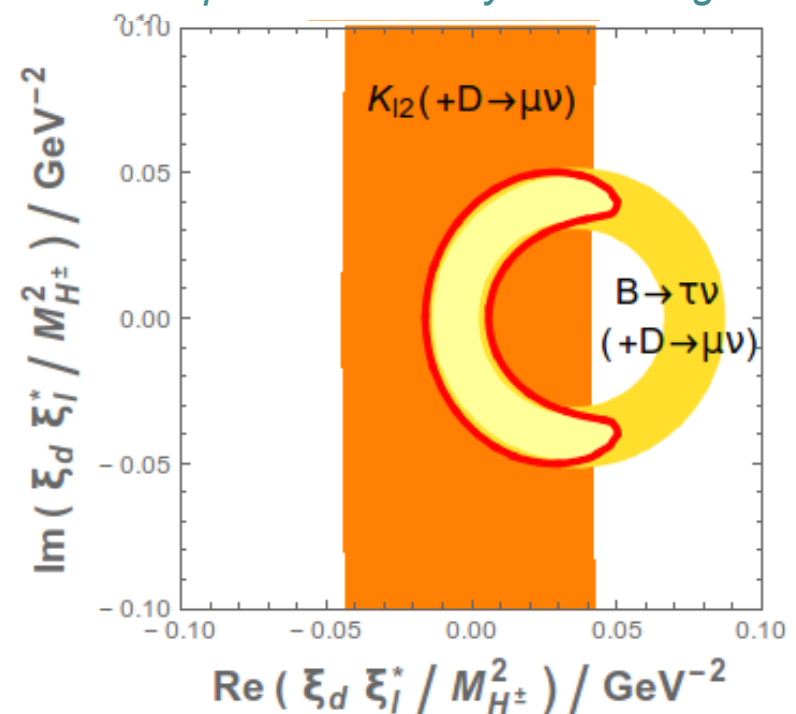
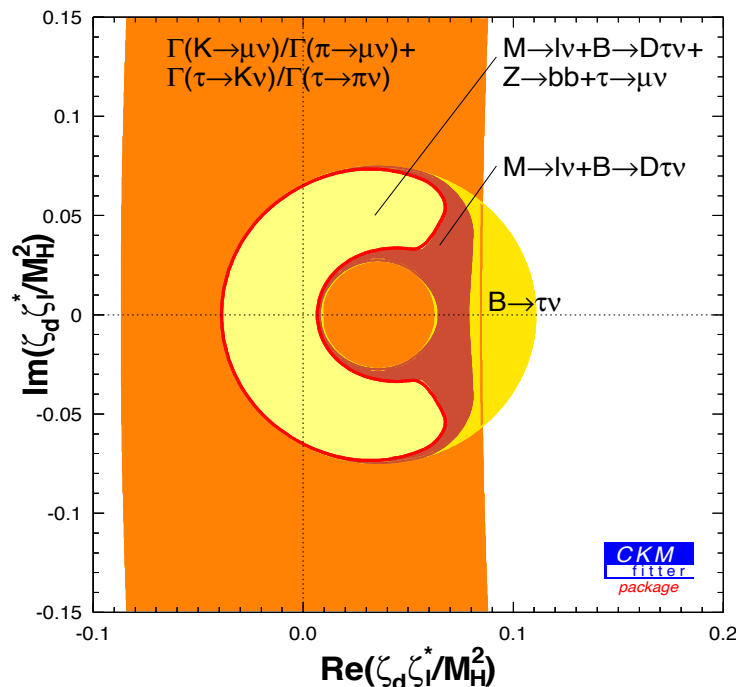
$$\mathcal{L}_Y = -\frac{\sqrt{2}}{v} H^+ \left\{ \bar{u} \left[\zeta_d V_{\text{CKM}} M_d \mathcal{P}_R - \zeta_u M_u^\dagger V_{\text{CKM}} \mathcal{P}_L \right] d + \zeta_l (\bar{\nu} M_l \mathcal{P}_R l) \right\} + \text{h.c.}$$

Jung, Pich, Tuzon'10

Pich@HQL'12

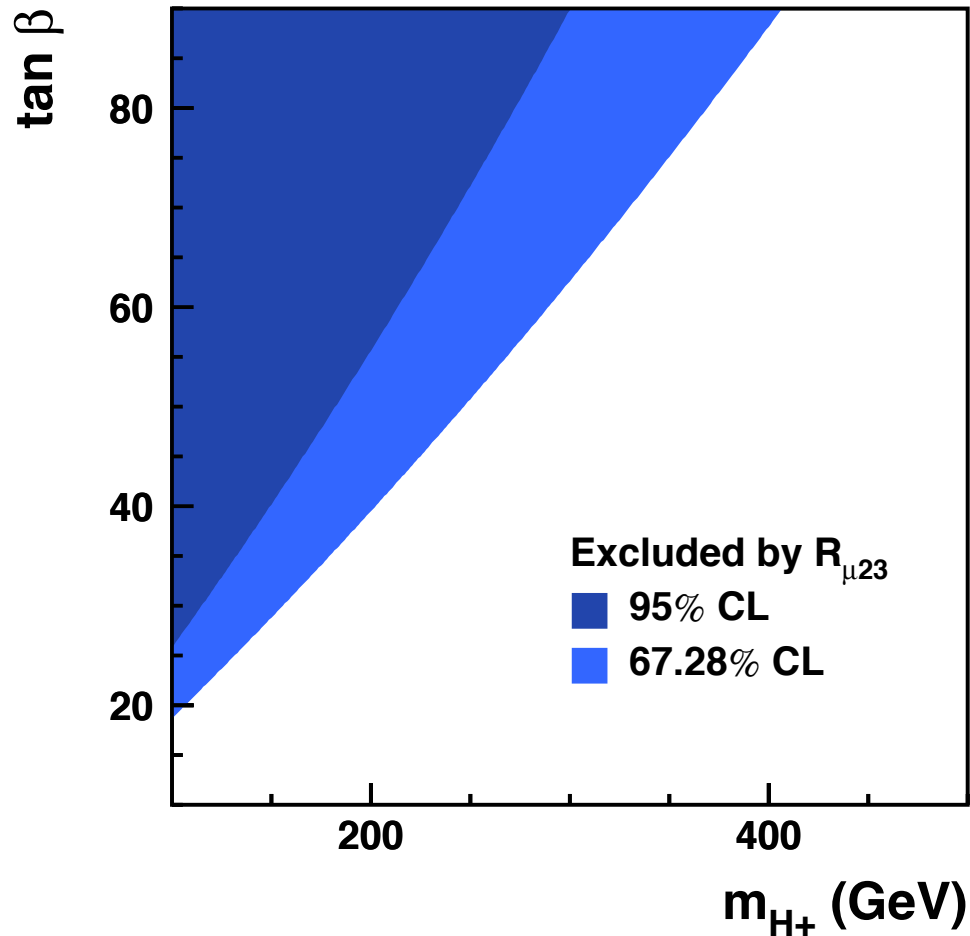


Update: Courtesy of M. Jung



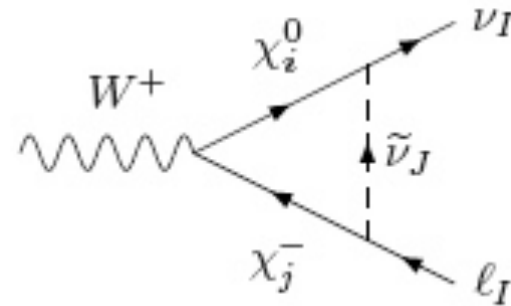
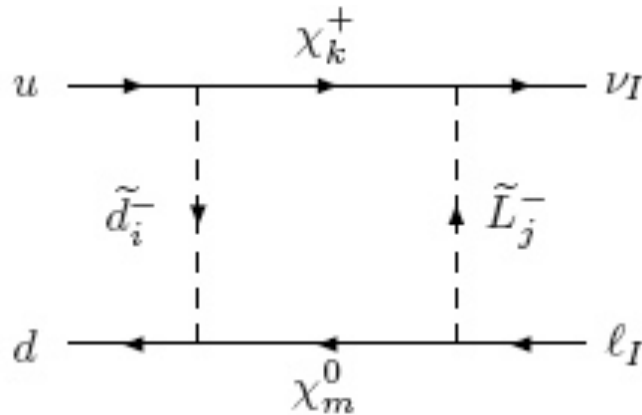
4.3 Constraints on charged Higgs

Antonelli et al.'10



4.4 Universality and SUSY

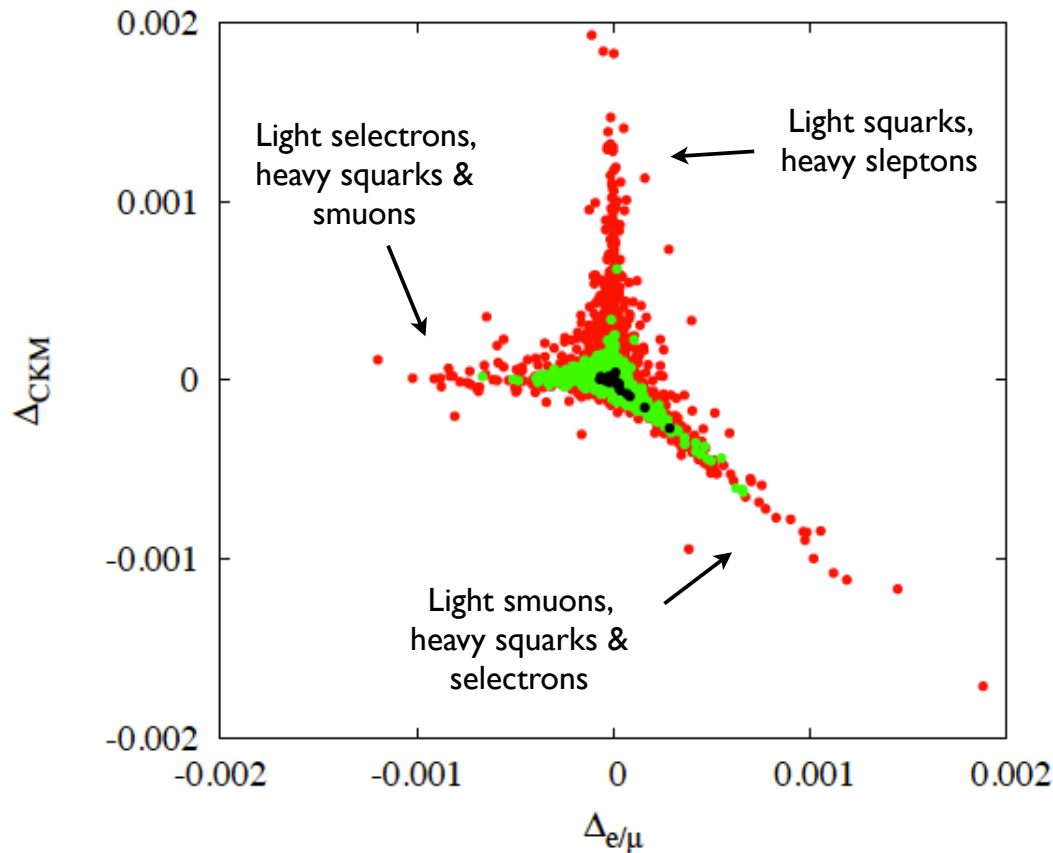
- MSSM: box and vertex corrections induce non-universal corrections to the V-A CC operators
- S,P,T operators suppressed by insertions of Yukawa couplings



*Barbieri et al '85, Hagiwara-Matsumoto-Yamada '95
Ramsey-Musolf Kurylov '01
Bauman, Erler, Ramsey-Musolf '12*

4.4 Universality and SUSY

Bauman, Erler, Ramsey-Musolf '12



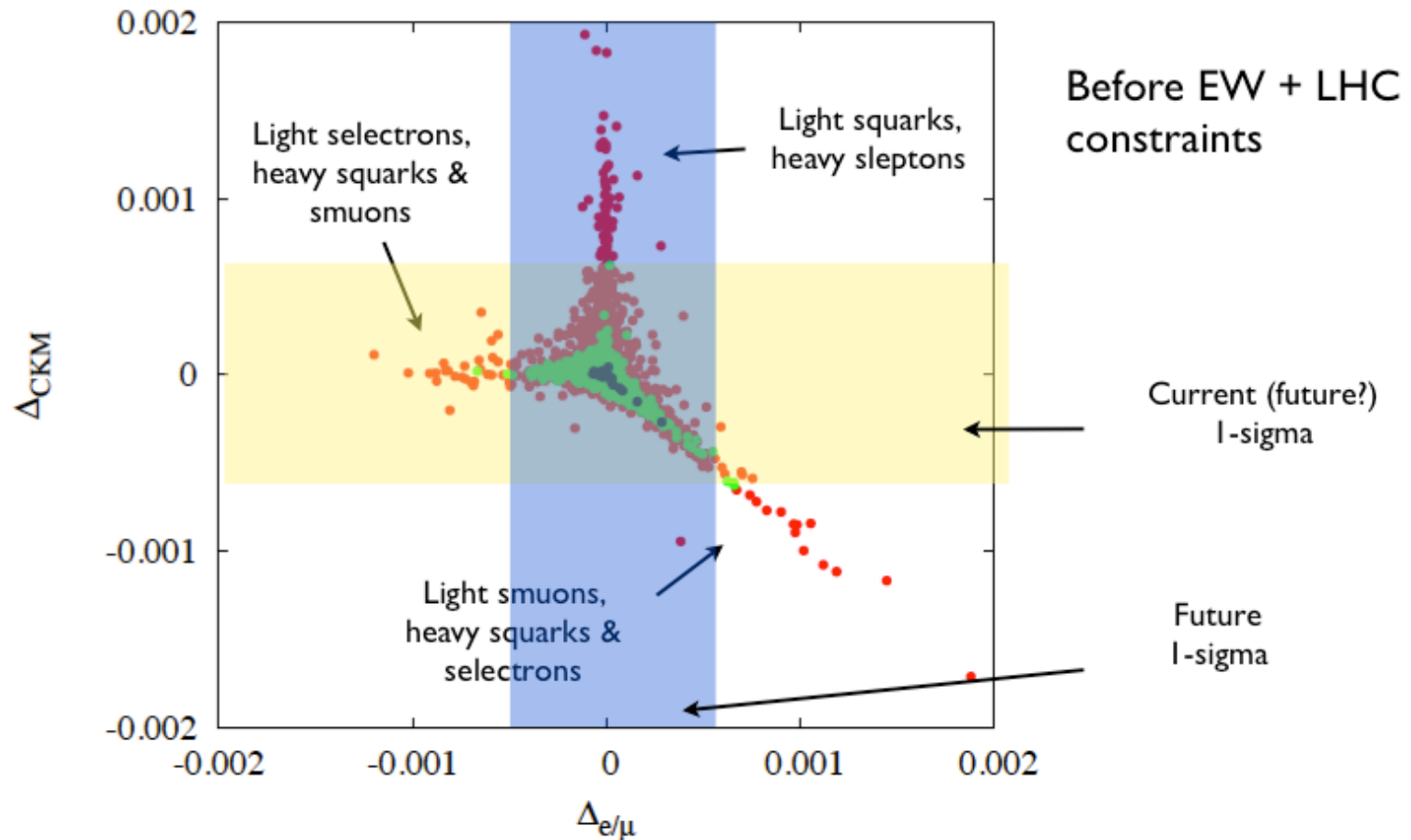
Before EW + LHC constraints

$$\begin{aligned} [G_F]_e/[G_F]_\mu &= 1 + \Delta_{e/\mu} \\ |V_{ud}|^2 + |V_{us}|^2 + |\cancel{V_{ub}}|^2 &= 1 + \Delta_{\text{CKM}} \end{aligned}$$

- Interesting correlation between Cabibbo universality and lepton universality: information on sfermion spectrum
- Essentially squark-slepton and selectron-smuon universality

4.4 Universality and SUSY

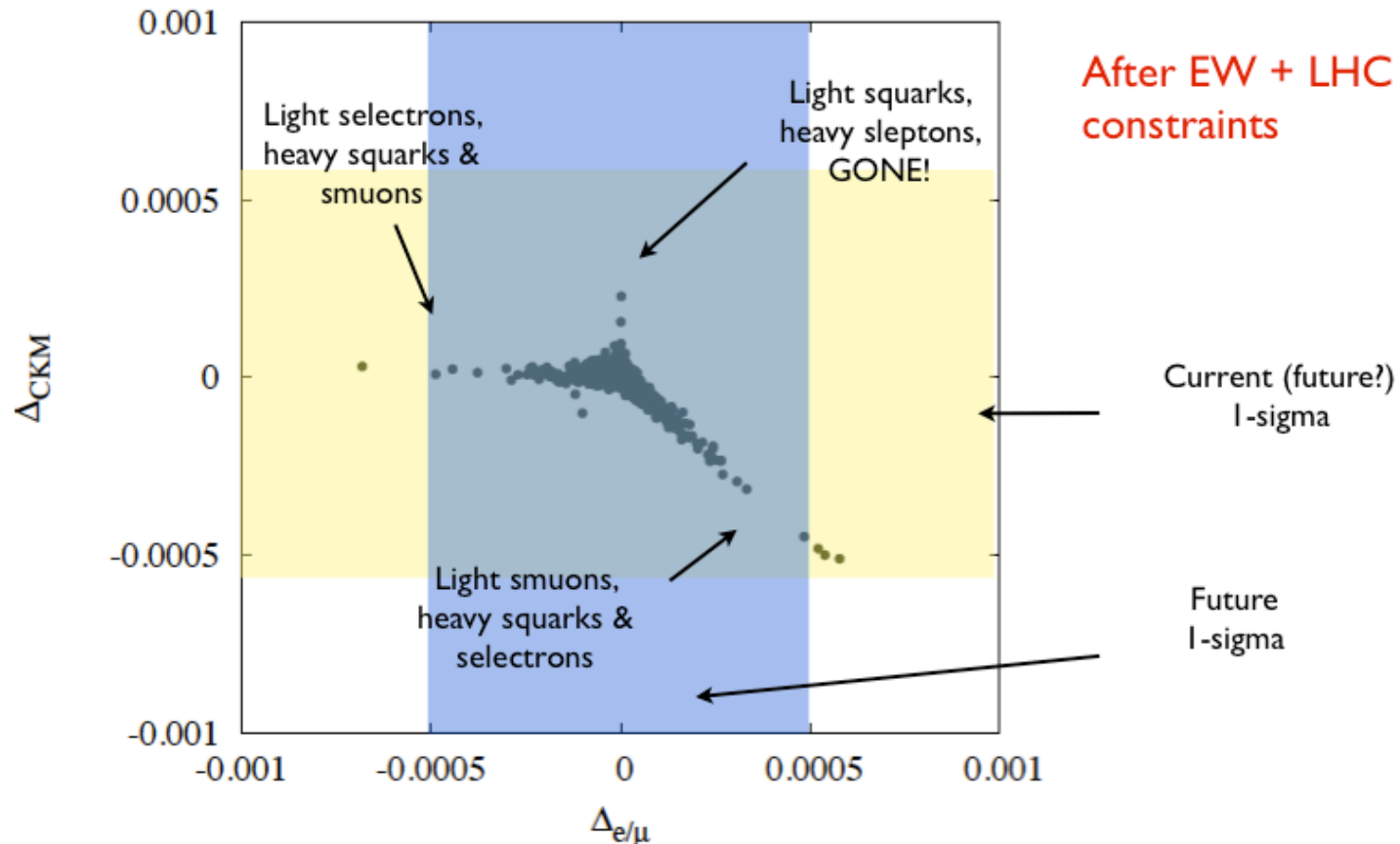
Bauman, Erler, Ramsey-Musolf '12



- Interesting correlation between Cabibbo universality and lepton universality: information on sfermion spectrum
- Essentially squark-slepton and selectron-smuon universality

4.4 Universality and SUSY

Bauman, Erler, Ramsey-Musolf '12



- Effects in the MSSM are small.
- Probing MSSM parameter space requires improved precision