

Prediction for the Cosmological Constant and Constraints on SUSY GUTS: Status Report for Resummed Quantum Gravity

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- Overview of Resummed Quantum Gravity
- Planck Scale Cosmology
- An Estimate of Λ
- An Open Question?
- Einstein-Heisenberg Consistency Condition
- Constraints on SUSY GUTs
- Outlook

- WHAT IS RESUMMATION?

- FAMILIAR SUMMATION: $\frac{1}{1-x} = \sum_{n=0}^{\infty} x^n$

- RESUMMATION:

$$\sum_{n=0}^{\infty} C_n \alpha_s^n \left\{ \begin{array}{l} = F_{\text{RES}}(\alpha_s) \sum_{n=0}^{\infty} B_n \alpha_s^n, \text{ EXACT} \\ \cong G_{\text{RES}}(\alpha_s) \sum_{n=0}^{\infty} B'_n \alpha_s^n, \text{ APROX.} \end{array} \right.$$

- A LITTLE HISTORY: 1988 ICHEP-Munich Conference Dinner, ONE YEAR BEFORE LEP DATA TAKING THAT LED, BY PRECISION PHYSICS, TO THE 't HOOFT-VELTMAN (1999) EW AND GROSS-WILCZEK-POLITZER (2004) QCD NOBEL PRIZES IN PHYSICS:
F. Berends and I considered, 'How Accurate Can Exponentiation (RESUMMATION) Really Be?'
- Would It Limit or Enhance Precision for a Given Level of Exactness: LO, NLO, NNLO, ?

Introduction


- 'Two' Realizations in Literature:
Jackson-Scharre(JS)(APPROX) vs YFS (EXACT)
- JS → 'limit to precision'
- YFS → 'no limit to precision'
- See 1989 CERN Yellow Book article: Frits was almost convinced, but not completely!
- Today, the analogous discussion continues to a new paradigm: quantum gravity

Introduction

- IS QUANTUM GRAVITY (Einstein-Hilbert Theory) CALCULABLE IN RELATIVISTIC QFT?
- STRING THEORY: NO. You need superstrings, supersymmetric one-dimensional objects of Planck length size, 1.62×10^{-33} cm.

String Theory

- Advantages:
 - UV FINITENESS,
 - THEORY OF EVERYTHING: $SU_{2L} \otimes U_1 \otimes QCD \otimes QG$,
 - APPLIED PHYSICS IN PURE MATHEMATICS, ...
- Disadvantages:
 - So Many Solutions $\sim 10^{500}$ and counting \Rightarrow ANTHROPICITY, ...
- quark:

QFT -	•	SST -	
0 size			$\ell_{Pl} = 1.62 \times 10^{-33} \text{ cm}$

- LOOP QUANTUM GRAVITY: NO. You need Planck length size loops that are the fundamental constructs for quantum gravity.

Loop Quantum Gravity

- Advantages:
Suspected UV FINITENESS,
LINEARIZED QG Constraints, ...
- Disadvantages:
UNKNOWN SEMI-CLASSICAL LIMIT,
QUESTIONS OF PRINCIPLE, ...

graviton: QFT - •
0 size

LQG -



....

$$\ell_{Pl} = 1.62 \times 10^{-33} \text{ cm}$$

- HORAVA-LIFSHITZ THEORY: NO. You need anisotropic scaling at Planck length scales:
 - Time and space differ by a factor of z in how they scale at Planck length distances with $z = 3$ in the original proposal—this violates local Lorentz invariance.

HORAVA-LIFSHITZ Quantum Gravity

- Advantages:
POWER COUNTING UV RENORMALIZABLE,
CONNECTIONS TO CONDENSED MATTER, ...
- Disadvantages:
VIOLATES LOCAL LORENTZ INVARIANCE,
SCALAR GRAVITON (PATHOLOGY?), ...
- graviton: RQFT - • : $1/(\omega^2 - \vec{k}^2)$ 0 size propagator HLQG - • : $1/(\omega^2 - \vec{k}^2 - G\vec{k}^6)$ 0 size propagator ,
 $\ell_{Pl} = 1.62 \times 10^{-33} \text{ cm} = \sqrt{G}$

- Our Response: Exact Amplitude-Based Resummation of Feynman's Formulation of Einstein's Theory – Resummed Quantum Gravity (RQG)
- RESULT (1): UV Finiteness!
- RESULT (2): Constraints on SUSY GUT's
- RESULT (3): Prediction for the Cosmological Constant Λ with Relatively Small Theoretical Uncertainty.
- RESULT (4): Consistent with Weinberg's Asymptotic Safety Ansatz, as realized by Exact Field Space Renormalization Group Program of Reuter *et al.*

- RESULT (5): Consistent with Kreimer's Leg Renormalizability Results ...
- Today we give a report on the status and outlook for this new RQG approach.

Overview of Resummed Quantum Gravity

SM \Leftrightarrow Many Massive Point Particles.

Feynman: spin is an inessential complication – checked. We replace $L_{SM}^G(x)$ with that of a free physical Higgs field, $\varphi(x)$, with a rest mass 125 GeV (ATLAS, CMS) \Rightarrow the representative model {R.P. Feynman, *Acta Phys. Pol.* 24 (1963) 697; Feynman *Lectures on Gravitation*, eds. F.B. Moringo and W.G. Wagner, (Caltech, Pasadena, 1971). }

$$\begin{aligned}\mathcal{L}(x) &= \frac{1}{2\kappa^2} R\sqrt{-g} + \frac{1}{2} (g^{\mu\nu} \partial_\mu \varphi \partial_\nu \varphi - m_o^2 \varphi^2) \sqrt{-g} \\ &= \frac{1}{2} \left\{ h^{\mu\nu, \lambda} \bar{h}_{\mu\nu, \lambda} - 2\eta^{\mu\mu'} \eta^{\lambda\lambda'} \bar{h}_{\mu\lambda, \lambda'} \eta^{\sigma\sigma'} \bar{h}_{\mu'\sigma, \sigma'} \right\} \\ &\quad + \frac{1}{2} \left\{ \varphi_{, \mu} \varphi^{, \mu} - m_o^2 \varphi^2 \right\} - \kappa h^{\mu\nu} \left[\overline{\varphi_{, \mu} \varphi_{, \nu}} + \frac{1}{2} m_o^2 \varphi^2 \eta_{\mu\nu} \right] \\ &\quad - \kappa^2 \left[\frac{1}{2} h_{\lambda\rho} \bar{h}^{\rho\lambda} (\varphi_{, \mu} \varphi^{, \mu} - m_o^2 \varphi^2) - 2\eta_{\rho\rho'} h^{\mu\rho} \bar{h}^{\rho'\nu} \varphi_{, \mu} \varphi_{, \nu} \right] + \dots\end{aligned}\tag{1}$$

Overview of Resummed Quantum Gravity

where $\varphi_{,\mu} \equiv \partial_\mu \varphi$ and we have

- $g_{\mu\nu}(x) = \eta_{\mu\nu} + 2\kappa h_{\mu\nu}(x)$,
 $\eta_{\mu\nu} = \text{diag}\{1, -1, -1, -1\}$
- $\bar{y}_{\mu\nu} \equiv \frac{1}{2} (y_{\mu\nu} + y_{\nu\mu} - \eta_{\mu\nu} y_\rho{}^\rho)$ for any tensor $y_{\mu\nu}$
- Feynman rules already worked-out by Feynman (*op. cit.*), where we use his gauge, $\partial^\mu \bar{h}_{\nu\mu} = 0$

\Leftrightarrow Quantum Gravity is just another quantum field theory where the metric now has quantum fluctuations as well.

Overview of Resummed Quantum Gravity

For example, the one-loop corrections to the graviton propagator due to matter loops is just given by the diagrams in the Figure.

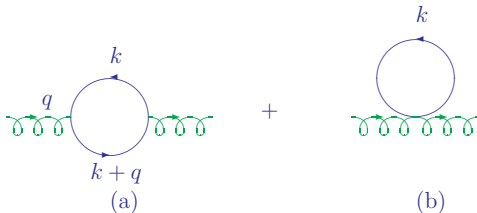


Figure : The scalar one-loop contribution to the graviton propagator.
 q is the 4-momentum of the graviton.

QG's BAD UV behavior – unrenormalizable.

Overview of Resummed Quantum Gravity

YFS resum the propagators in the **NON-ABELIAN** gauge theory of QG:

⇒ from the YFS formula

$$iS'_F(p) = \frac{ie^{-\alpha B''_\gamma}}{S_F^{-1}(p) - \Sigma'_F(p)}, \quad (2)$$

we find for Quantum Gravity, proceeding as above, the analogue of

$$\alpha B''_\gamma = \int \frac{d^4 \ell}{(2\pi)^4} \frac{-i\eta^{\mu\nu}}{(\ell^2 - \lambda^2 + i\epsilon)} \frac{-ie(2ik_\mu)}{(\ell^2 - 2\ell k + \Delta + i\epsilon)} \frac{-ie(2ik'_\nu)}{(\ell^2 - 2\ell k' + \Delta' + i\epsilon)} \Big|_{k=k'} \quad (3)$$

as $-B''_g(k)$ with

$$B''_g(k) = -2i\kappa^2 k^4 \int \frac{d^4 \ell}{16\pi^4} \frac{1}{\ell^2 - \lambda^2 + i\epsilon} \frac{1}{(\ell^2 + 2\ell k + \Delta + i\epsilon)^2} \quad (4)$$

for $\Delta = k^2 - m^2 \Rightarrow$ for a scalar field

$$i\Delta'_F(k)|_{\text{YFS-resummed}} = \frac{ie^{B''_g(k)}}{(k^2 - m^2 - \Sigma'_S + i\epsilon)}.$$

Overview of Resummed Quantum Gravity

⇒

Expand theory with the 'improved Born' propagators

$$iP_{\alpha_1 \dots; \alpha'_1 \dots} \Delta'_F(k) |_{YFS\text{-resummed}, \Sigma'_s=0} = \frac{iP_{\alpha_1 \dots; \alpha'_1 \dots} e^{B''_g(k)}}{(k^2 - m^2 + i\epsilon)} \quad (5)$$

where in the DEEP UV we get

$$B''_g(k) = \frac{\kappa^2 |k^2|}{8\pi^2} \ln \left(\frac{m^2}{m^2 + |k^2|} \right), \quad (6)$$

⇒ ALL PROPAGATORS FALL FASTER THAN ANY POWER
OF $|k^2|$ ⇒ QG IS FINITE (SEE MPLA17 (2002)
2371;hep-ph/0607198)!

CONTACT WITH ASYMPTOTIC SAFETY APPROACH

- OUR RESULTS IMPLY

$$G(k) = G_N / (1 + \frac{k^2}{a^2})$$

⇒ FIXED POINT BEHAVIOR FOR

$$k^2 \rightarrow \infty,$$

IN AGREEMENT WITH THE PHENOMENOLOGICAL
ASYMPTOTIC SAFETY APPROACH OF **BONANNO &
REUTER** IN PRD**62**(2000) 043008.

- OUR RESULTS ⇒ AN ELEMENTARY PARTICLE HAS
NO HORIZON. THIS AGREES WITH **BONANNO & REUTER**
THAT A BLACK HOLE WITH A MASS LESS THAN
 $M_{cr} \sim M_{Pl}$
HAS NO HORIZON.

BASIC PHYSICS:

$G(k)$ VANISHES FOR $k^2 \rightarrow \infty$.

Planck Scale Cosmology

- Bonanno and Reuter [see arXiv.org:0803.2546](https://arxiv.org/abs/0803.2546), and refs. therein
– phenomenological approach to Planck scale cosmology:
STARTING POINT IS THE EINSTEIN-HILBERT THEORY

$$\mathcal{L}(x) = \frac{1}{2\kappa^2} \sqrt{-g} (R - 2\Lambda) \quad (7)$$

PHENOMENOLOGICAL EXACT RENORMALIZATION GROUP
FOR THE WILSONIAN COARSE GRAINED EFFECTIVE
AVERAGE ACTION IN FIELD SPACE \Rightarrow RUNNING NEWTON
CONSTANT $G_N(k)$ AND COSMOLOGICAL CONSTANT $\Lambda(k)$
APPROACH UV FIXED POINTS AS k GOES TO ∞ IN THE
DEEP EUCLIDEAN REGIME – $k^2 G_N(k) \rightarrow g_*$, $\Lambda(k) \rightarrow \lambda_* k^2$.

- Due to the thinning of the degrees of freedom in Wilsonian field space renormalization theory, the arguments of Foot et al. ([PLB664\(2008\)199](#)) are obviated. – See also [MPLA 25\(2010\)607](#); [SHAPIRO&SOLA](#), [PLB682\(2009\)105](#)

Planck Scale Cosmology

CONTACT WITH COSMOLOGY PROCEEDS AS FOLLOWS:
PHENOMENOLOGICAL CONNECTION BETWEEN THE
MOMENTUM SCALE k CHARACTERIZING THE
COARSENESS OF THE WILSONIAN GRAININESS OF THE
AVERAGE EFFECTIVE ACTION AND THE COSMOLOGICAL
TIME t , B-R SHOW STANDARD COSMOLOGICAL
EQUATIONS ADMIT (see also Bonanno et al., 1006.0192) THE
FOLLOWING EXTENSION:

$$\begin{aligned}\left(\frac{\dot{a}}{a}\right)^2 + \frac{K}{a^2} &= \frac{1}{3}\Lambda + \frac{8\pi}{3}G_N\rho \\ \dot{\rho} + 3(1 + \omega)\frac{\dot{a}}{a}\rho &= 0 \\ \dot{\Lambda} + 8\pi\rho\dot{G}_N &= 0 \\ G_N(t) &= G_N(k(t)) \\ \Lambda(t) &= \Lambda(k(t))\end{aligned}$$

(8)



FOR DENSITY ρ AND SCALE FACTOR $a(t)$



WITH ROBERTSON-WALKER METRIC REPRESENTATION

$$ds^2 = dt^2 - a(t)^2 \left(\frac{dr^2}{1 - Kr^2} + r^2(d\theta^2 + \sin^2 \theta d\phi^2) \right) \quad (9)$$

$K = 0, 1, -1 \Leftrightarrow$ RESPECTIVELY FLAT, SPHERICAL AND
PSEUDO-SPHERICAL 3-SPACES FOR CONSTANT TIME t
FOR A LINEAR RELATION BETWEEN THE PRESSURE p
and ρ (EQN. OF STATE)

$$p(t) = \omega \rho(t). \quad (10)$$

Planck Scale Cosmology

FUNCTIONAL RELATIONSHIP BETWEEN MOMENTUM SCALE k AND COSMOLOGICAL TIME t DETERMINED PHENOMENOLOGICALLY VIA

$$k(t) = \frac{\xi}{t} \quad (11)$$

WITH POSITIVE CONSTANT ξ .

Using the UV fixed points for $k^2 G_N(k) = g_*$ and $\Lambda(k)/k^2 = \lambda_*$ B-R SHOW THAT (8) ADMITS, FOR $K = 0$, A SOLUTION IN THE PLANCK REGIME ($0 \leq t \leq t_{\text{class}}$, with t_{class} a few times the Planck time t_{Pl}), WHICH JOINS SMOOTHLY ONTO A SOLUTION IN THE CLASSICAL REGIME ($t > t_{\text{class}}$) **which agrees with standard Friedmann-Robertson-Walker phenomenology but with the horizon, flatness, scale free Harrison-Zeldovich spectrum, and entropy problems solved by Planck scale quantum physics.**



PHENOMENOLOGICAL NATURE OF THE ANALYSIS: THE fixed-point results g_* , λ_* depend on the cut-offs used in the Wilsonian coarse-graining procedure.

KEY PROPERTIES OF g_* , λ_* USED FOR THE B-R ANALYSES: they are both positive and the product $g_*\lambda_*$ is cut-off/threshold function independent.

An Estimate of Λ

- In Phys. Dark Univ. **2** (2013) 97, using (5) and (5) we get rigorous cut-off independent values for the fixed points g_* , λ_* and the following estimate of Λ :

$$\begin{aligned}
 \rho_\Lambda(t_0) &\cong \frac{-M_{Pl}^4(1 + c_{2,eff}k_{tr}^2/(360\pi M_{Pl}^2))^2}{64} \sum_j \frac{(-1)^{F_j} n_j}{\rho_j^2} \\
 &\quad \times \frac{t_{tr}^2}{t_{eq}^2} \times \left(\frac{t_{eq}^{2/3}}{t_0^{2/3}}\right)^3 \\
 &\cong \frac{-M_{Pl}^2(1.0362)^2(-9.194 \times 10^{-3})(25)^2}{64 t_0^2} \\
 &\cong (2.4 \times 10^{-3} \text{eV})^4,
 \end{aligned} \tag{12}$$

where the age of the universe is $t_0 \cong 13.7 \times 10^9$ yrs.

- Compare: $\rho_\Lambda(t_0)|_{\text{expt}} \cong ((2.37 \pm 0.05) \times 10^{-3} \text{eV})^4$.



An Open Question

- A MAIN UNCERTAINTY: t_{tr}
- B-R: NUMERICAL STUDIES $\Rightarrow t_{\text{tr}} \cong 25/M_{Pl}$
- IN GENERAL, A FACTOR of $\mathcal{O}(100)$ IS ALLOWED
- CAN WE DO BETTER?

Einstein-Heisenberg Consistency Condition

- Recently, arXiv:1507.00661, we use the de Sitter space solutions of Duerr et al. to get the Einstein-Heisenberg consistency condition

$$k \geq \frac{\sqrt{5}}{2w_0} = \frac{\sqrt{5}}{2} \frac{1}{\sqrt{3/\Lambda(k)}} \quad (13)$$

from the Heisenberg uncertainty relation $\Delta p \Delta q \geq \frac{1}{2}$, with $\Delta p = k$ and $(w_0 = \sqrt{3/\Lambda})$

$$(\Delta q)^2 \cong \frac{\int_0^{w_0} dw w^2 w^2 \langle \cos^2 \theta \rangle}{\int_0^{w_0} dw w^2} = \frac{1}{5} w_0^2. \quad (14)$$

- Violation of (13) ends Planck scale inflation: solving for $k_{\text{tr}} \Rightarrow k_{\text{tr}} \cong M_{\text{Pl}}/25.3$, in agreement with what Bonnano and Reuter suggested from numerical studies.
- \Rightarrow uncertainty on our estimate of ρ_Λ is $\mathcal{O}(10)$.



Note

$$\langle 0|\mathcal{H}|0\rangle \sim \int^{M_{Pl}} \frac{d^3k}{(2\pi)^3} \frac{1}{2} \omega(k) = \int^{M_{Pl}} \frac{d^3k}{(2\pi)^3} \frac{1}{2} \sqrt{k^2 + m_i^2}$$

Raises the question of GUTS: Use SO(10)
SUSY GUT Approach of Dev & Mohapatra
(PRD82(2010)035014):

Intermediate Stage:

$SU_{2L} \times SU_{2R} \times U_1 \times SU(3)^c$

SM Stage at $\sim 2\text{TeV} = M_R$:

$SU_{2L} \times U_1 \times SU(3)^c$

SUSY Breaking at EW scale M_S :

$U_1 \times SU(3)^c$

- Possible spectrum(?)

$$m_{\tilde{g}} \cong 1.5(10)\text{TeV}$$

$$m_{\tilde{G}} \cong 1.5\text{TeV}$$

$$m_{\tilde{q}} \cong 1.0\text{TeV}$$

$$m_{\tilde{t}} \cong 0.5\text{TeV}$$

$$m_{\tilde{\chi}_i^0} \cong \begin{cases} 0.4\text{TeV}, & i = 1 \\ 0.5\text{TeV}, & i = 2, 3, 4 \end{cases}$$

$$m_{\tilde{\chi}_i^\pm} \cong 0.5\text{TeV}, \quad i = 1, 2$$

$$m_S = .5\text{TeV}, \quad S = A^0, H^\pm, H_2.$$



$$\Delta_{\text{GUT}} = \sum_{j \in \{\text{MSSM low energy susy partners}\}} \frac{(-1)^F n_j}{\rho_j^2} \\ \cong 1.13(1.12) \times 10^{-2}$$

Constraints on SUSY GUTS



- Compensate by either (A) adding new susy families with scalars lighter than fermions or (B) allowing the gravitino mass to go to $\sim .05 M_{\text{GUT}} \sim 2 \times 10^{15} \text{ GeV}$.
- For approach (A),
new quarks and leptons at $M_{\text{High}} \sim 3.4(3.3) \times 10^3 \text{ TeV}$,
scalar partners at $\sim .5 \text{ TeV} = M_{\text{Low}}$

What About EW, QCD, GUT Symmetry Breaking Scales?

Consider GUT symmetry breaking:

It gives a $M_{\text{GUT}}^4 / (.01 M_{\text{Pl}}^4 / 64) < 10^{-6}$ correction, which we drop here.

The other breaking scales are even smaller and hence their corrections are even less significant in our result for

ρ_Λ .

Constraints on SUSY GUTS

Covariance Issues for $\dot{\Lambda}$, $\dot{G}_N \neq 0$:

Bianchi's Identity,

$$D^\nu(\Lambda g_{\nu\mu} + 8\pi G_N T_{\nu\mu}) = 0,$$

allows

$$\dot{\rho} + 3\frac{\dot{a}}{a}(1+\omega)\rho = -\frac{\dot{\Lambda} + 8\pi\rho\dot{G}_N}{8\pi G_N},$$

a more general form of the new Friedmann eqns:
qualitatively the same but details differ -- see
[arXiv:0907.4555,1103.4632,1202.5097....](#)

Our estimate uses the more general form.

- Precision Quantum Field Theory: EW, QCD, QG \equiv Control all limits:

IR ($z \rightarrow 1$)

and

Collinear ($p_T \rightarrow 0$)

UV limit

- We now have control over all aspects of the QG corrections.
- Toward quantitative understanding of ρ_Λ .