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high Energy PHYSICS

Experimental investigation of $\pi^+\pi^-$ and K^+K^- atom

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DIRAC collaboration

DImeson Relativistic Atomic Complexes



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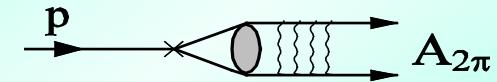
totally 68 physicists from 20 Institutes

Production of pionium

Atoms are Coulomb bound state of two pions produced in one proton-nucleus collision

$$\frac{d\sigma_{nlm}^A}{d\vec{P}} = (2\pi)^3 \frac{E_A}{M_A} \left| \psi_{nlm}^{(C)}(0) \right|^2 \frac{d\sigma_s^0}{d\vec{p}_+ d\vec{p}_-} \Bigg|_{\vec{p}_+ = \vec{p}_-} \quad \sigma_A = k \sigma_C (Q < Q_0)$$

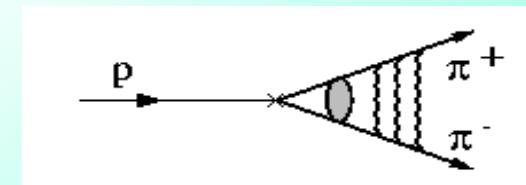
Nemenov 1985



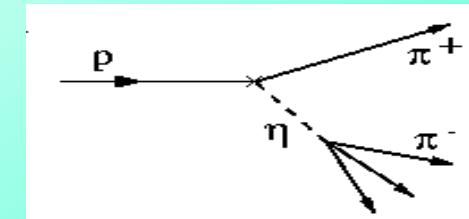
Background processes:

Coulomb pairs. They are produced in one proton nucleus collision from fragmentation or short lived resonances and exhibit Coulomb interaction in the final state

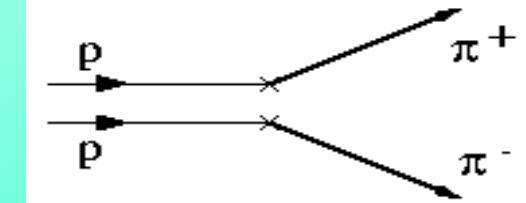
$$\frac{d^2\sigma_C}{d\vec{p}_+ d\vec{p}_-} = A_C(q) \frac{d\sigma_s^0}{d\vec{p}_+ d\vec{p}_-}, \quad A_C(q) = \frac{2\pi m_\pi \alpha / q}{1 - \exp(-2\pi m_\pi \alpha / q)}$$



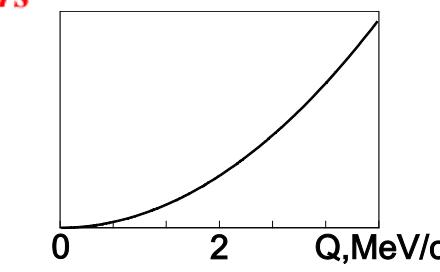
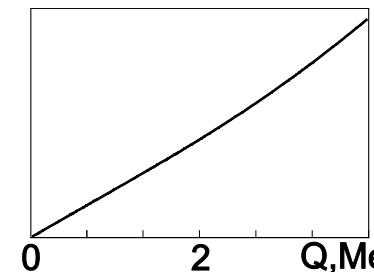
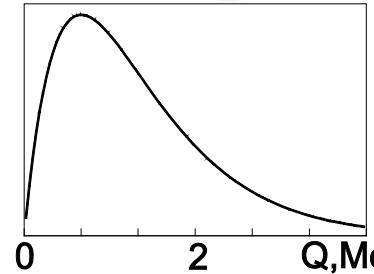
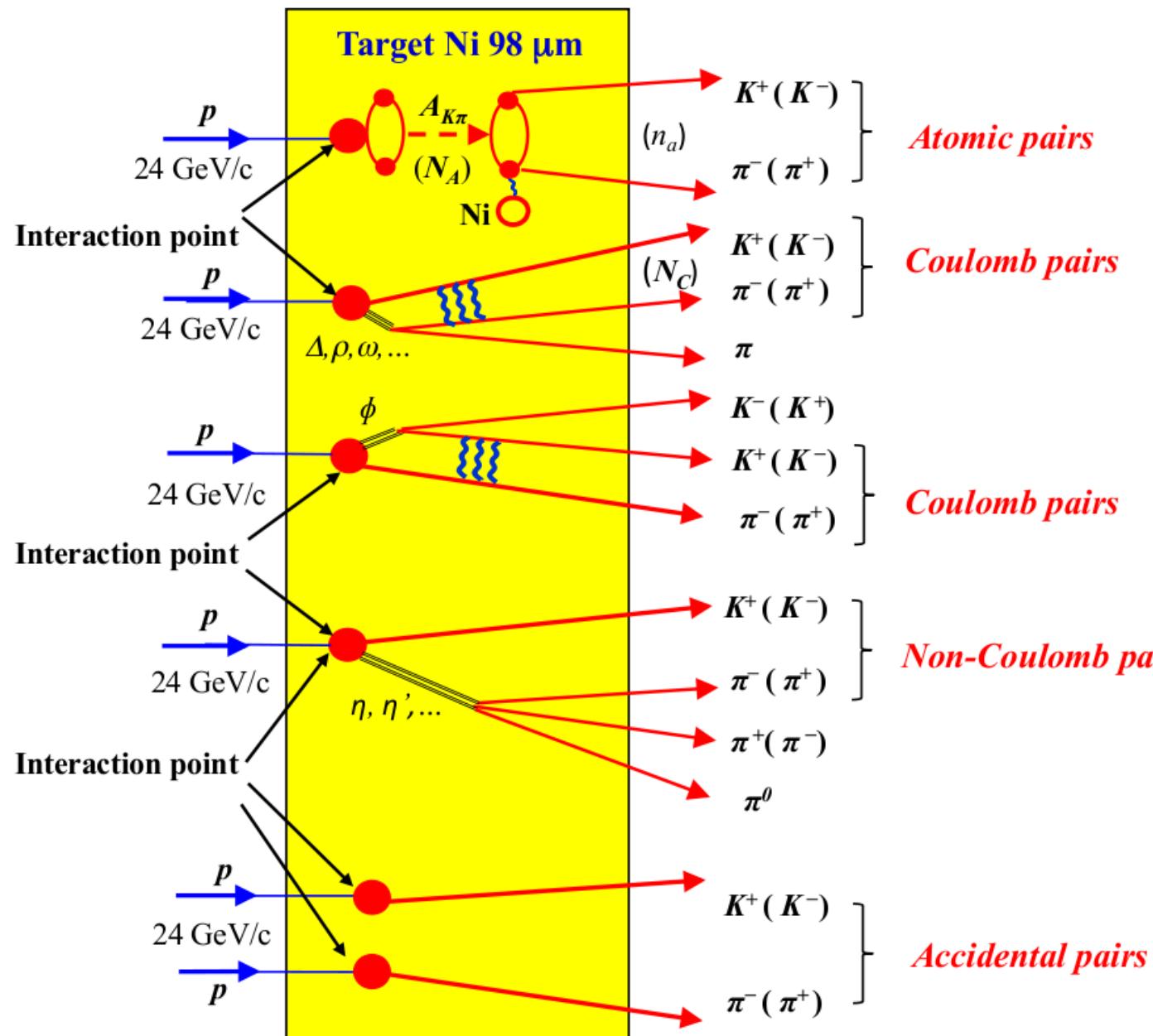
Non-Coulomb pairs. They are produced in one proton nucleus collision. At least one pion originates from a long lived resonance. No Coulomb interaction in the final state



Accidental pairs. They are produced in two independent proton nucleus collision. They do not exhibit Coulomb interaction in the final state



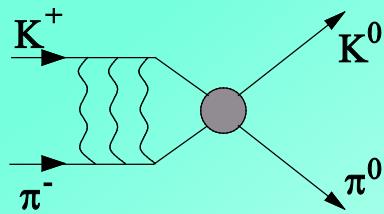
Method of πK ($\pi\pi$) atom observation and investigation



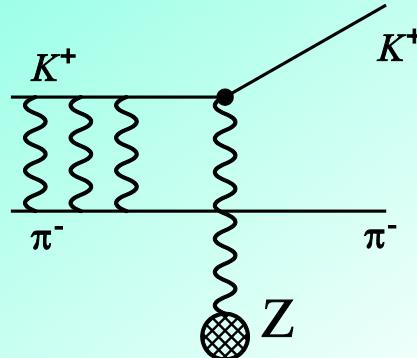
Break-up probability

During propagation in matter atoms:

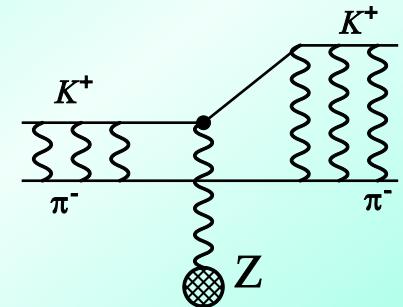
annihilate



break up(ionized)



excite



result in production of πK ($\pi^+\pi^-$) atomic pairs n_A

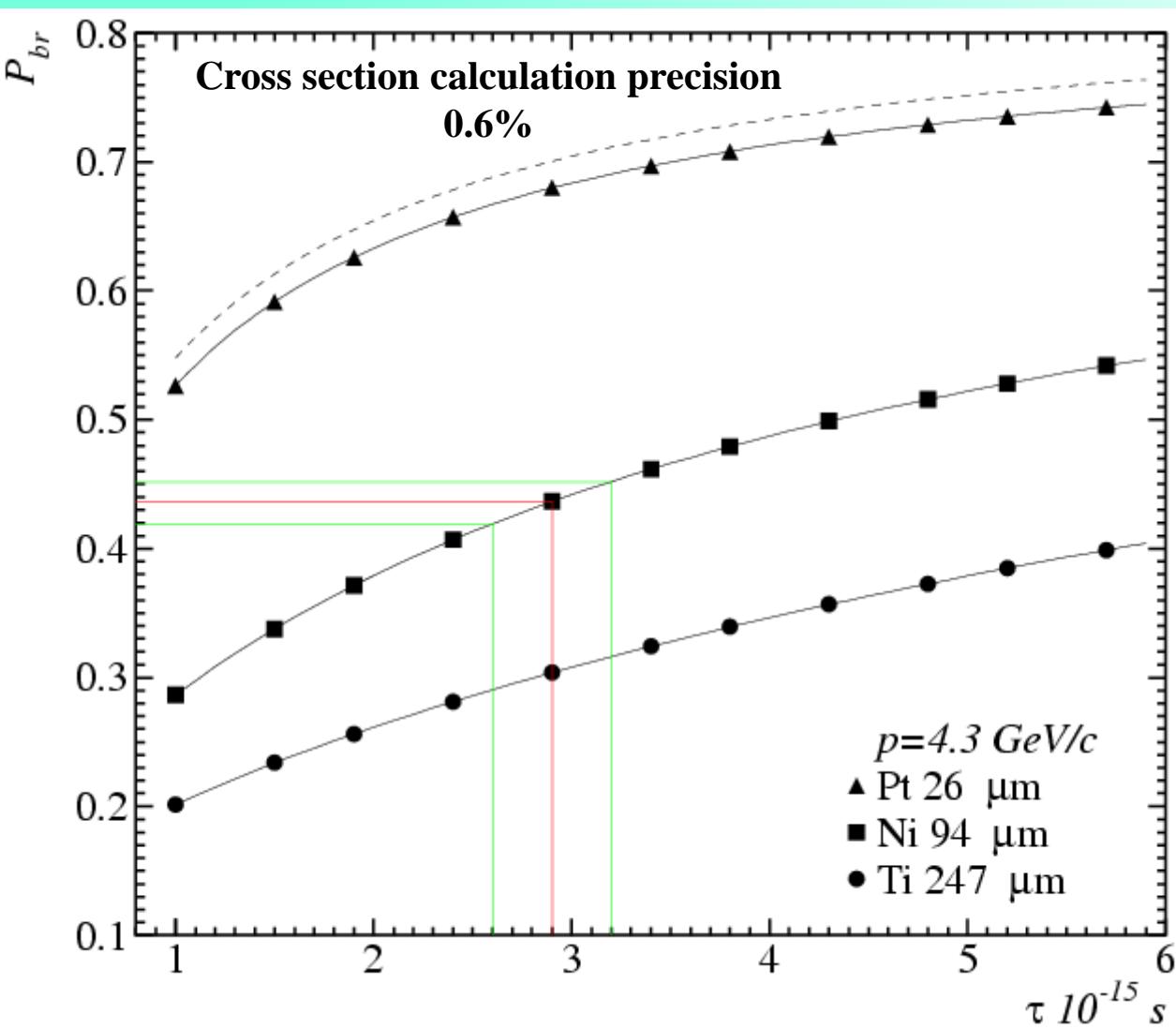
$$\frac{\sigma_A}{\sigma_C} = \frac{|\psi_{nlm}^{(C)}|^2}{|\psi_{\vec{q}}^{(C)}|^2}$$

$$N_A = K(Q_0)N_C(Q \leq Q_0), \frac{\delta K(Q_0)}{K(Q_0)} \leq 10^{-2}$$

$$n_A - \text{atomic pairs number}, \quad P_{br} = \frac{n_A}{N_A}$$

Break-up probability

Solution of the transport equations provides one-to-one dependence of the measured break-up probability (P_{br}) on pionium lifetime τ



$$\delta\tau=10\% \rightarrow \delta P_{br}=4\%$$

All targets have the same thickness in radiation lengths $6.7 \cdot 10^{-3} X_0$

There is an optimal target material for a given lifetime

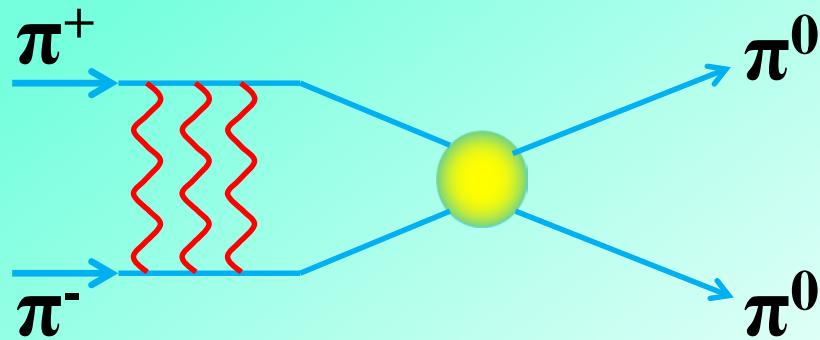
$\pi^+\pi^-$ atom lifetime

$\pi^+\pi^-$ atom (pionium) is a hydrogen-like atom consisting of π^+ and π^- mesons:

$$E_B = -1.86 \text{ keV},$$

$$r_B = 387 \text{ fm},$$

$$p_B \approx 0.5 \text{ MeV/c}$$



The lifetime of $\pi^+\pi^-$ atom is dominated by the decay into $\pi^0\pi^0$ mesons:

$$\Gamma = \frac{1}{\tau} = \Gamma_{2\pi^0} + \Gamma_{2\gamma} \quad \frac{\Gamma_{2\gamma}}{\Gamma_{2\pi^0}} \approx 4 \times 10^{-3}$$

$$\Gamma_{ns \rightarrow 2\pi^0} = R |\psi_{ns}(0)|^2 |a_0 - a_2|^2$$

$$\tau_{1s} = (2.9 \pm 0.1) \times 10^{-15} \text{ s}$$

a_0 and a_2 are the $\pi\pi$ s-wave scattering lengths for isospin $I=0$ and

$I=2$.

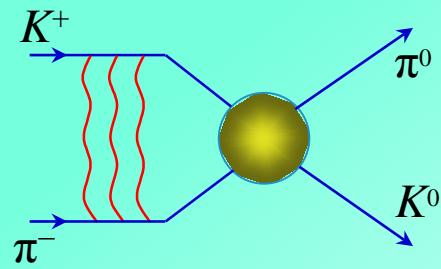
$$\psi_{nl} \begin{cases} \neq 0 \text{ for } l=0 & A_{2\pi}(1s, 2s, \dots, (n-1)s) \xrightarrow{\gamma} \pi^0\pi^0 \\ = 0 \text{ for } l \neq 0 & A_{2\pi}(np) \xrightarrow{\gamma} A_{2\pi}(1s, 2s, \dots, (n-1)s) \xrightarrow{\gamma} \pi^0\pi^0 \end{cases}$$

The lifetime of np states depends on transition $np \xrightarrow{\gamma} 1s, 2s, \dots, (n-1)s$ probability
 This probability is about three orders less than $ns \xrightarrow{\gamma} \pi^0\pi^0$ decay into $\pi^0\pi^0$

$K^{\dagger}\pi$ and $K\pi^{\dagger}$ atoms lifetime

$K\pi$ -atom ($A_{K\pi}$) is a hydrogen-like atom consisting of K^{\pm} and π^{\mp} mesons:

$$E_B = -2.9 \text{ keV} \quad r_B = 249 \text{ fm} \quad p_B = 0.79 \text{ MeV}$$



$$\frac{1}{\tau} = \frac{8}{9} \alpha^3 \mu^2 p^* (a_{1/2} - a_{3/2})^2 (1 + \delta_K)$$

The $K\pi$ -atom lifetime ground state $1S$, $\tau = 1/\Gamma$ is dominated by the annihilation process into $K^0\pi^0$:

$$A_{K^+\pi^-} \rightarrow \pi^0 K^0$$

$$A_{\pi^+K^-} \rightarrow \pi^0 \bar{K}^0$$

$$\mu = 109 \text{ MeV}/c^2$$

$$p^* = 11.8 \text{ MeV}/c$$

$$\delta_k = 0.040 \pm 0.022$$

[S.Bilenky et al., Sov. J. Nucl. Phys. 10 (1969) 469]

[J. Schweizer, Phys. Lett. B 587 (2004) 33]

SU(3) ChPT predictions [J. Bijnens et al. JHEP 0405 (2004) 036]

$$\frac{1}{3} M_\pi (a_{1/2} - a_{1/3}) = M_\pi a_0^- = 0.071(CA) \rightarrow 0.079(1l) \rightarrow 0.89(2l) \quad [\text{P. Buttiker et al., Eur. Phys. J. C33 (2004) 409}]$$

$$\rightarrow 0.090 \pm 0.005(\text{dispersion}) \rightarrow \tau = (3.5 \pm 0.4) \times 10^{-15} \text{ s}$$

Lattice QCD calculations of ChPT low energy constant

[NPLQCD, Phys. Rev. D74 (2006) 114503]

[Z.Fu, Phys. Rev. D85 (2012) 074501]

[C.B. Lang et al., Phys. Rev. D86 (2012) 054508]

$$M_\pi a_0^- = 0.077 \pm 0.001^{+0.002}_{-0.005}$$

$$M_\pi a_0^- = 0.0777 \pm 0.0013 \pm ?$$

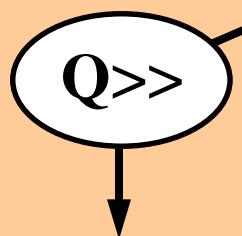
$$M_\pi a_0^- = 0.0811 \pm 0.0143$$

Theoretical motivation



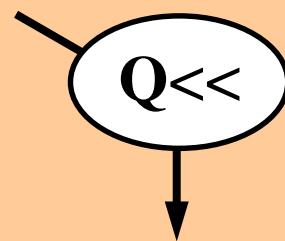
Strong interaction: $L_{QCD} = L_{sym} + L_{sym-break}$ ($m_q \neq 0$)

HIGH energy
(small distance)



(chiral symmetry)

LOW energy
(large distance)



perturbative QCD:

$$L_{QCD}(q,g)$$

Interaction → „weak“ (asympt. freedom)
Method: expansion in coupling

Checks only L_{sym} ($m_q \ll 0$) !

non-perturbative QCD:

$$L_{eff}(\text{GB: } \pi, K, \eta); L_{lattice}(q,g)$$

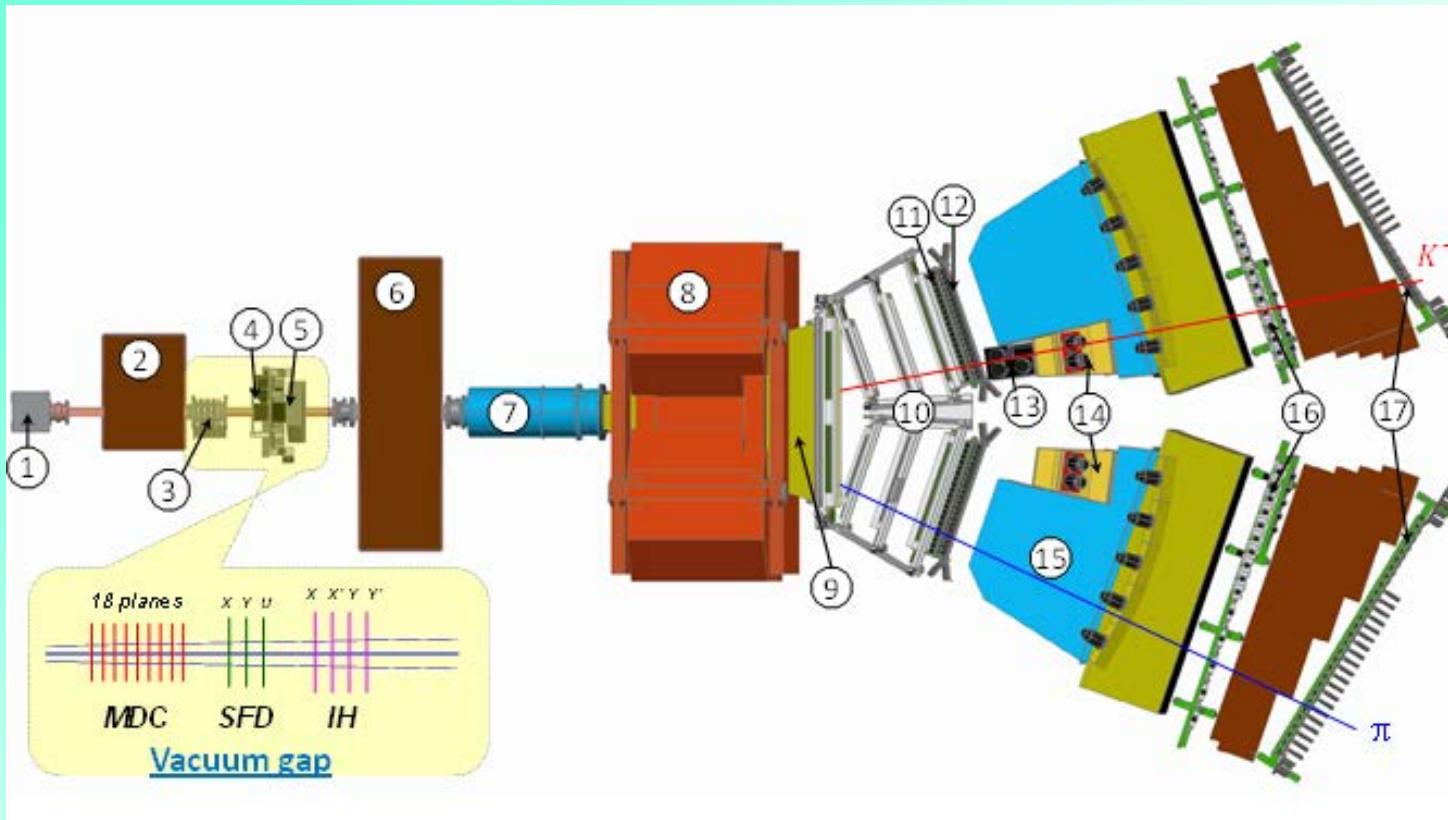
Interaction → „strong“ (confinement)
Methods: 1) Chiral Perturbation Theory
2) Lattice Gauge Theory

Checks L_{sym} as well as $L_{sym-break}$!

spontaneously broken symmetry

quark-condensate

Experimental setup at CERN PS



1 Target station with Ni foil; 2 First shielding; 3 Micro Drift Chambers;
4 Scintillating Fiber Detector; 5 Ionization Hodoscope; 6 Second Shielding;
7 Vacuum Tube; 8 Spectrometer Magnet; 9 Vacuum Chamber; 10 Drift
Chambers; 11 Vertical Hodoscope; 12 Horizontal Hodoscope; 13 Aerogel
Čerenkov; 14 Heavy Gas Čerenkov; 15 Nitrogen Čerenkov; 16 Preshower;
17 Muon Detector

Spectrometer resolutions

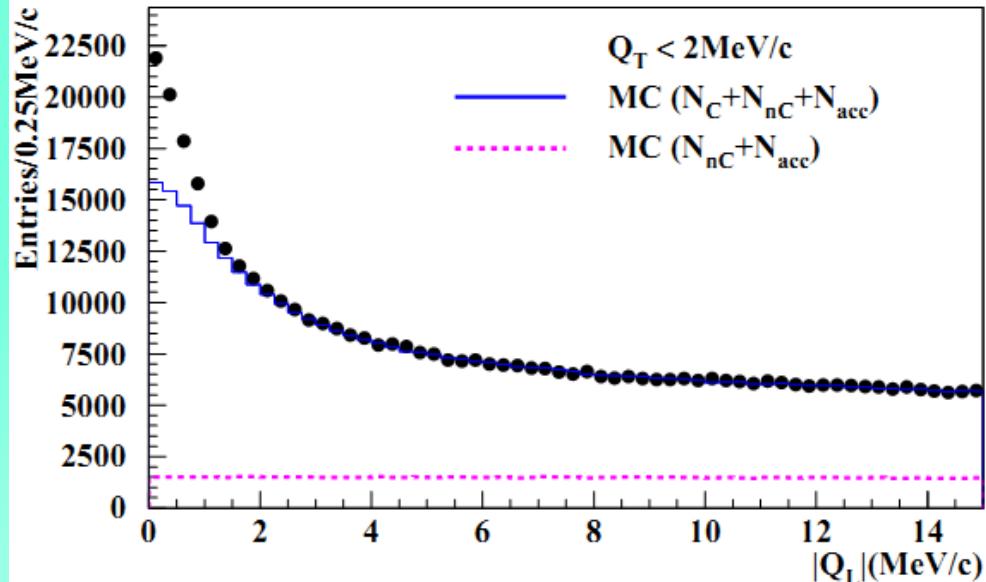
SFD			
Coordinate precision	$\sigma_X = 60 \mu\text{m}$	$\sigma_Y = 60 \mu\text{m}$	$\sigma_W = 120 \mu\text{m}$
Time precision	$\sigma^t_X = 380 \text{ ps}$	$\sigma^t_Y = 512 \text{ ps}$	$\sigma^t_W = 522 \text{ ps}$

DC	VH
Coordinate precision $\sigma = 85 \mu\text{m}$	Time precision $\sigma = 100 \text{ ps}$

Spectrometer		
Relative resolution on the particle momentum in L.S.		$3 \cdot 10^{-3}$
Precision on Q-projections	$\sigma_{QX} = \sigma_{QY} = 0.5 \text{ MeV/c}$	$\sigma_{QL} = 0.5 \text{ MeV/c } (\pi\pi)$ $\sigma_{QL} = 0.9 \text{ MeV/c } (\pi K)$

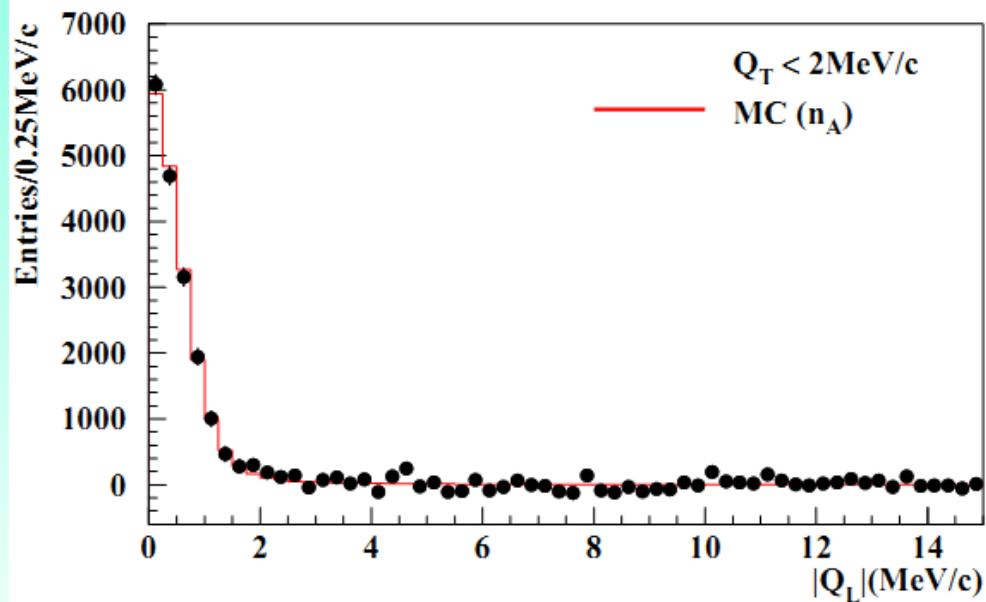
Trigger efficiency 98 % for pairs with $Q_L < 28 \text{ MeV/c}$
 $Q_X < 6 \text{ MeV/c}$
 $Q_Y < 4 \text{ MeV/c}$

$\pi^+\pi^-$ atoms 2001-2003



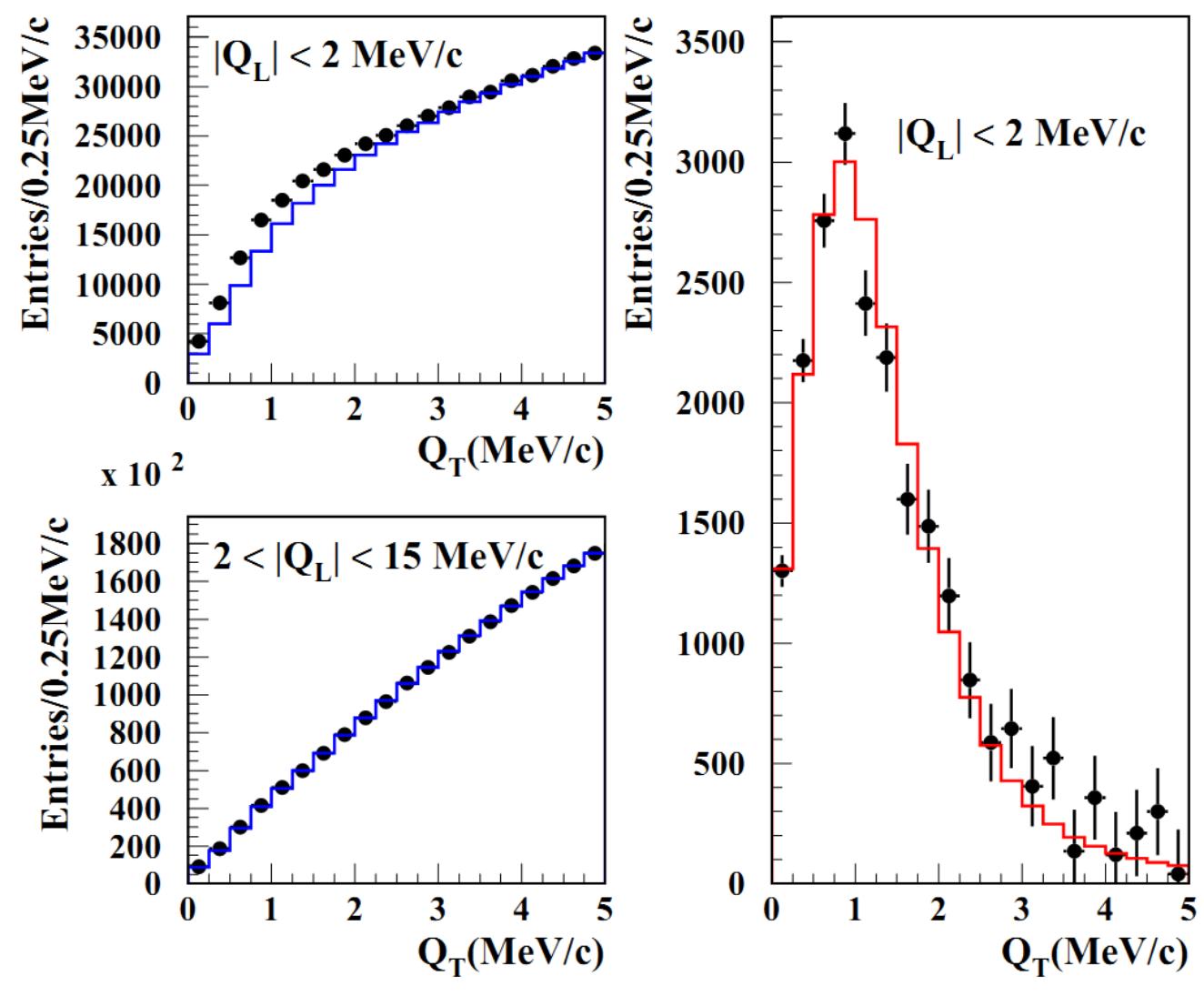
Q_L distribution

← All events



← After background subtraction

$\pi^+\pi^-$ atoms 2001-2003



Q_T distribution

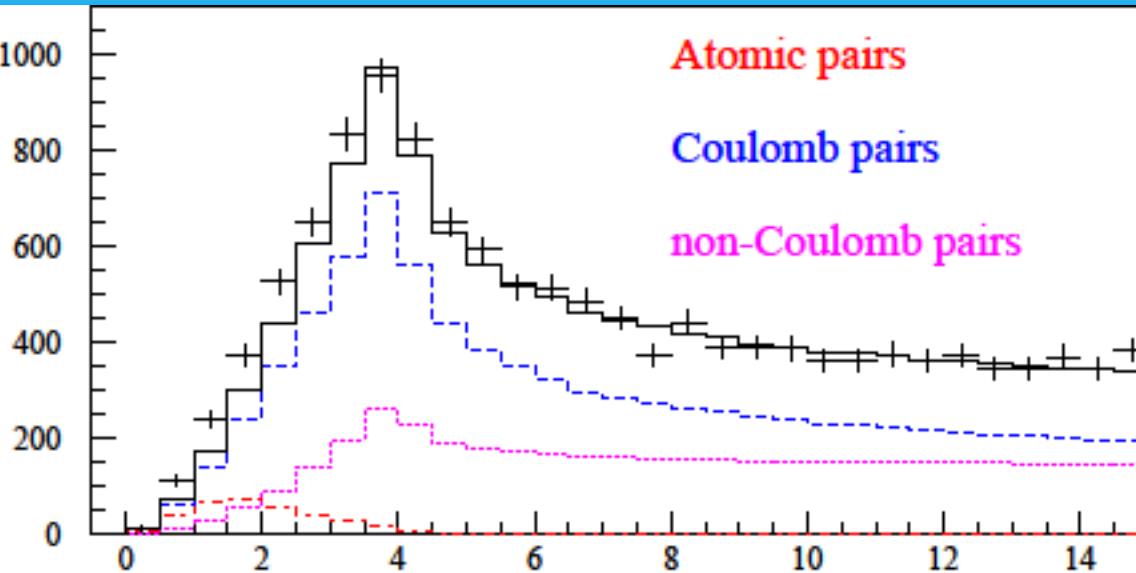
← After
background
subtraction
for
 $Q_L < 2 \text{ MeV}/c$

$\pi^+\pi^-$ atoms 2001-2003

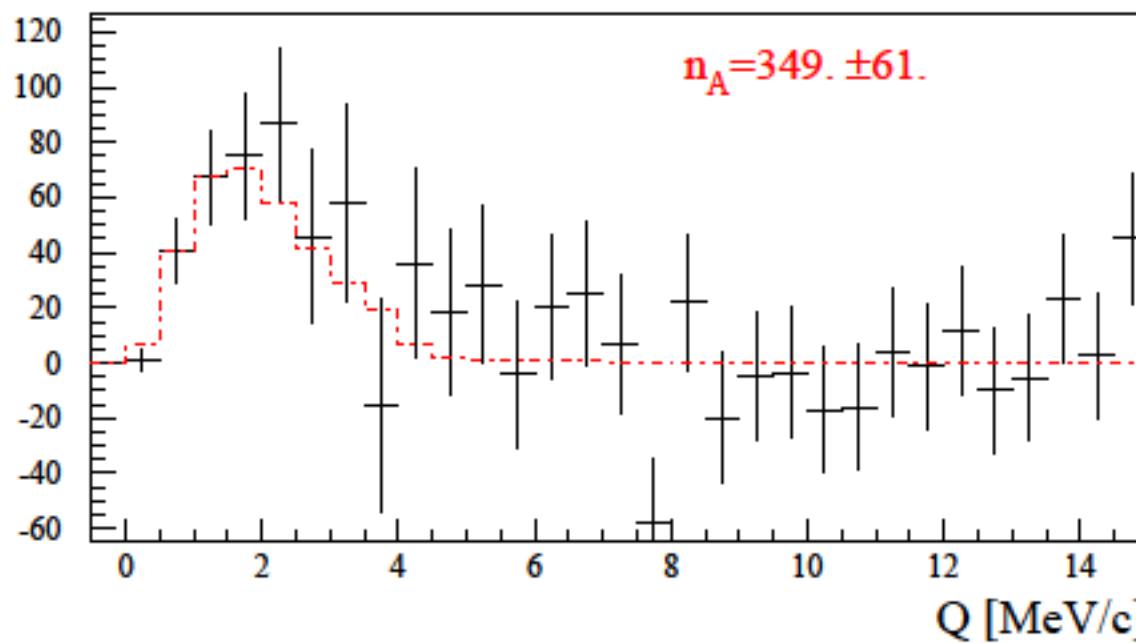
Ni, p_{beam}	χ^2/ndf	n_A	N_C	N_{nC}	N_{acc}	P_{br}
94 μm , 24 GeV/c	2127/2079	6020 ± 216	546003 ± 4549	45624 ± 4501	63212 ± 208	0.441 ± 0.018
98 μm , 24 GeV/c	4288/4149	9321 ± 274	828554 ± 5811	93148 ± 5754	98499 ± 255	0.452 ± 0.015
98 μm , 20 GeV/c	4257/4144	5886 ± 210	496820 ± 4441	60867 ± 4397	59392 ± 144	0.472 ± 0.020
combined samples		21227 ± 407	1871377 ± 8613	199639 ± 8526	221103 ± 359	

DIRAC data	$\tau_{1s} (10^{-15} \text{ s})$					$ a_0 - a_2 $					Reference
	value	stat	syst	<i>theo</i> *	tot	value	stat	syst	<i>theo</i> *	tot	
2001	2.91 $[+0.45 \quad +0.19 \quad [+0.49]$ $-0.38 \quad -0.49 \quad [-0.62]$					0.264 $+0.017 \quad +0.022 \quad [+0.033]$ $-0.020 \quad -0.009 \quad [-0.020]$					PL B 619 (2005) 50
2001-03	3.15 $[+0.20 \quad +0.20 \quad [+0.28]$ $-0.19 \quad -0.18 \quad [-0.26]$					0.2533 $+0.0078 \quad +0.0072 \quad [+0.0106]$ $-0.0080 \quad -0.0077 \quad [-0.0111]$					PL B 704 (2011) 24

πK atoms observation



All data Platinum
and Nickel targets



$K^+\pi^-$ and $K^-\pi^+$ atoms
 Q_L distribution for
 $Q_T < 4$ MeV/c
 $\chi^2 / \text{ndf} = 41/37$
In absense of “atomic
pairs” $\chi^2 / \text{ndf} = 73/38$

$K^+ \pi^-$ and $K^- \pi^+$ pairs analysis

Analysis	$\pi^- K^+$	$\pi^+ K^-$	$\pi^+ K^- + \pi^- K^+$
Q	$243 \pm 51 (4.7\sigma)$	$106 \pm 32 (3.3\sigma)$	$349 \pm 61 (5.7\sigma)$
$ Q_L $	$164 \pm 79 (2.1\sigma)$	$67 \pm 47 (1.4\sigma)$	$230 \pm 92 (2.5\sigma)$
$ Q_L , Q_T$	$237 \pm 50 (4.7\sigma)$	$78 \pm 32 (2.5\sigma)$	$314 \pm 59 (5.3\sigma)$

Analysis with $|Q_L|, Q_T$

$$n_A = 314 \pm 59(\text{stat}) \pm 10(\text{syst}) = 314 \pm 60(\text{tot})$$

5.2 standard deviations

Analysis with Q

$$n_A = 349 \pm 61(\text{stat}) \pm 9(\text{syst}) = 349 \pm 62(\text{tot})$$

5.6 standard deviations

[DIRAC, Physical Review Letters 117, 112001 (2016)
CERN-EP-2016-128 ; arXiv:1605.06103]

πK atoms lifetime

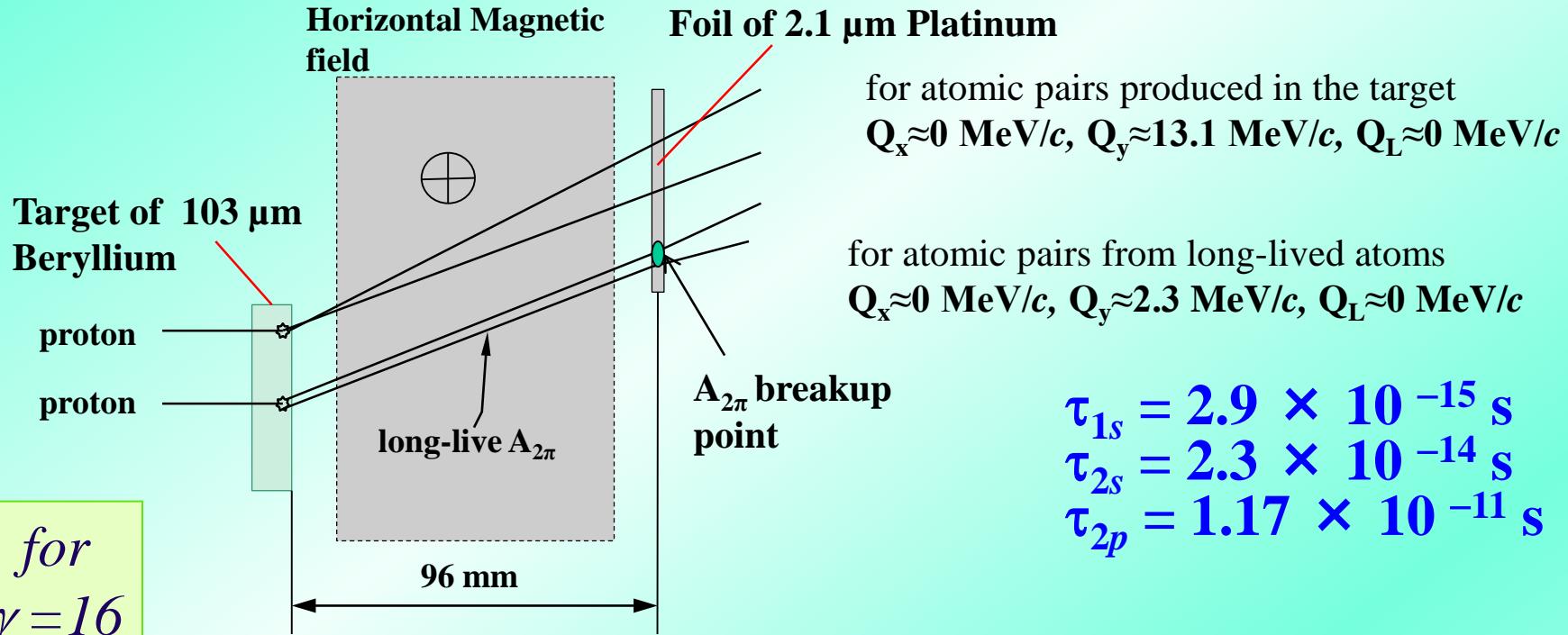
$$\tau = (5.5^{+5.0}_{-2.8}|\text{tot}) \times 10^{-15} \text{ s}$$

$$1/3 |a_{1/2} - a_{3/2}| = (0.072^{+0.031}_{-0.020}) M_{\pi^+}^{-1}$$

[DIRAC, Physical Review D 96, 052002 (2017)
CERN-EP-2017-137, arXiv:1707.02184]

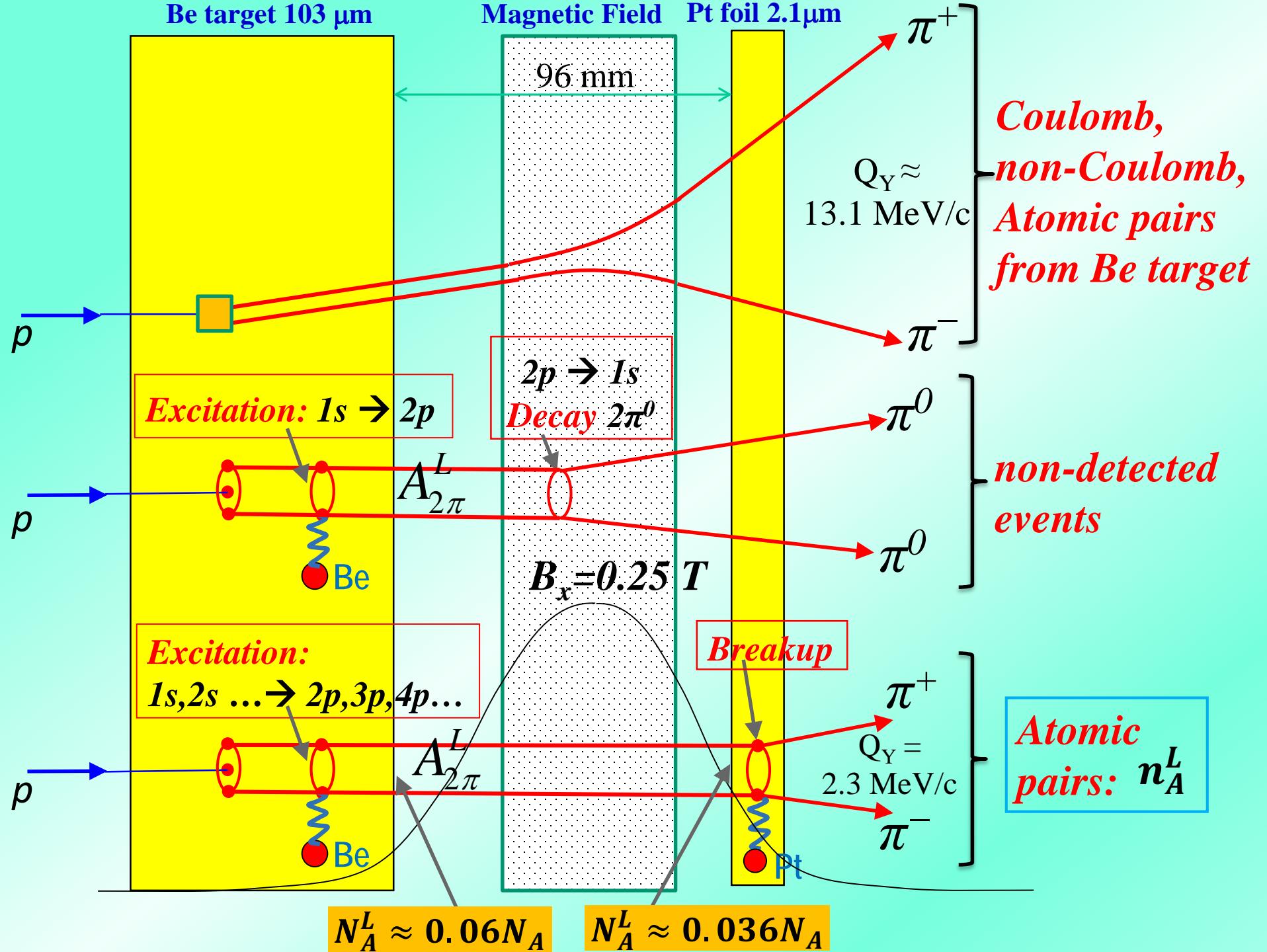
Search for long-lived states of $\pi^+\pi^-$ atoms

During 2011-2012 the data were collected for observation of the long-lived states of $\pi^+\pi^-$ atom. This observation opens the future possibility to measure the energy difference between ns and np states $\Delta E(ns-np)$ and the value of $\pi\pi$ scattering length combination $|2a_0+a_2|$.

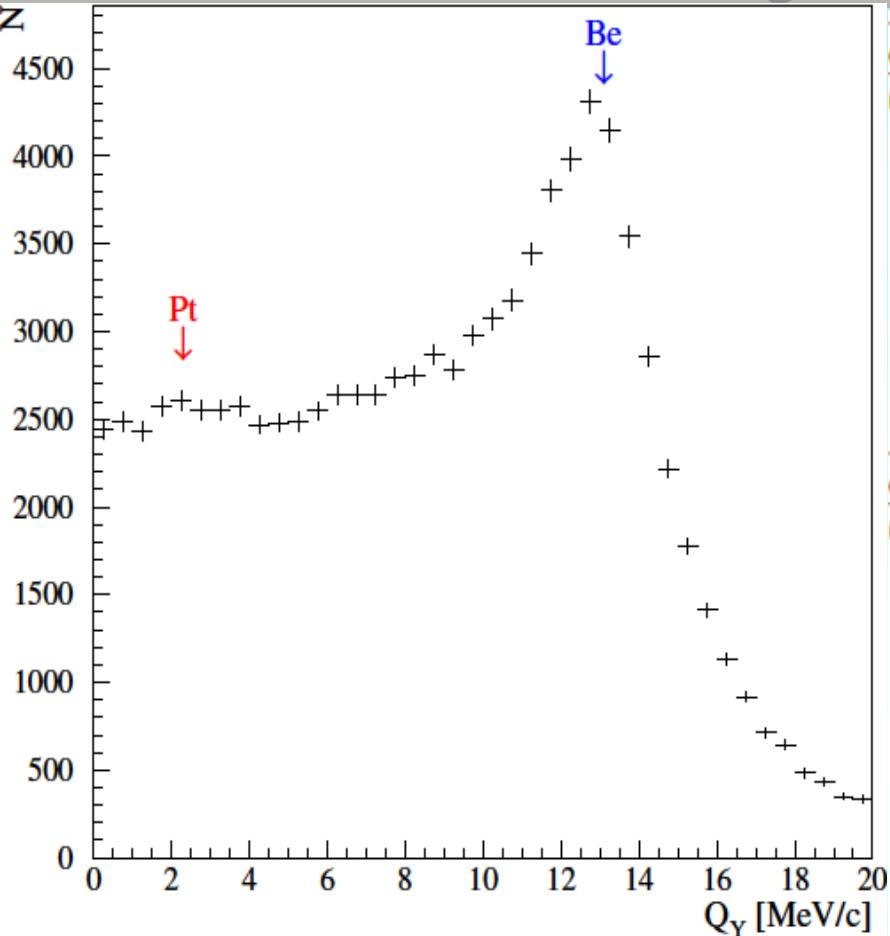


$$l(2p) = 56 \text{ mm}, l(3p) = 190 \text{ mm}, l(4p) = 430 \text{ mm}, l(5p) = 840 \text{ mm}$$

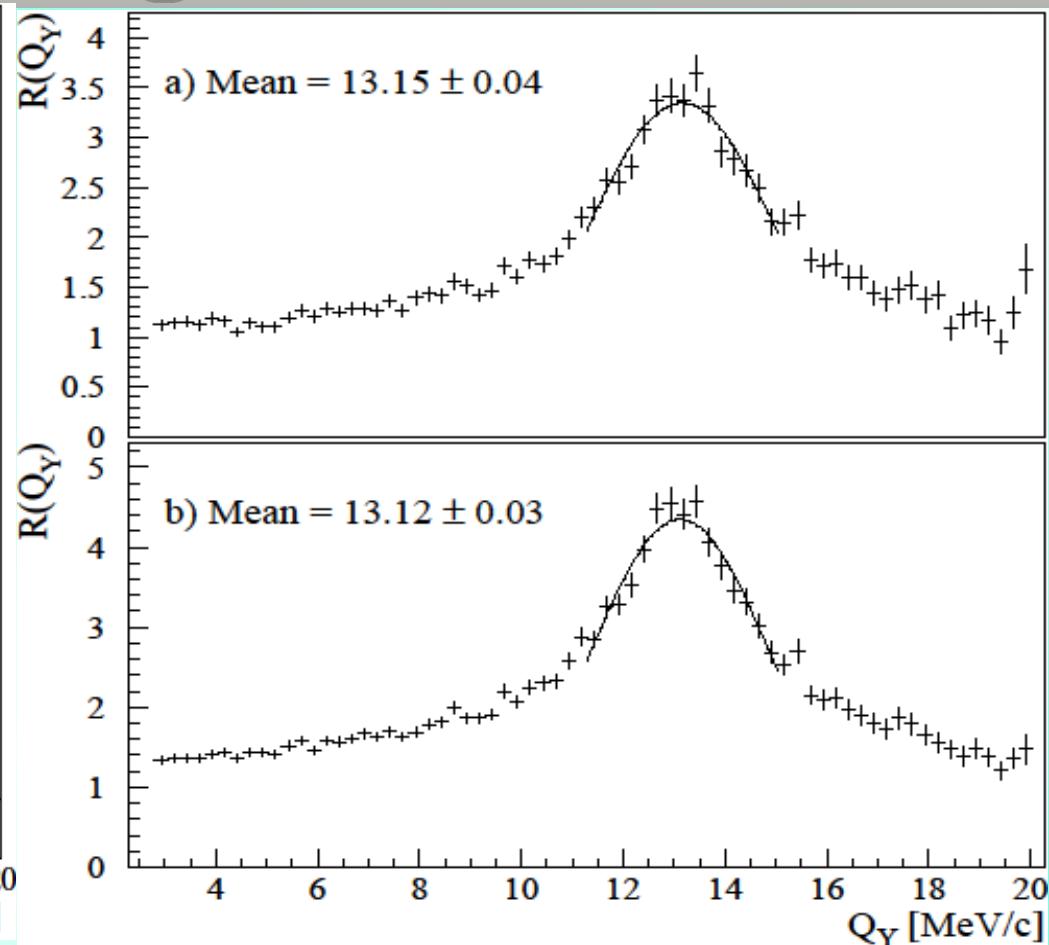
$$l(1s) = 15 \mu\text{m}, l(2s) = 0.11 \text{ mm}, l(3s) = 0.38 \text{ mm}, l(4s) = 0.89 \text{ mm}, l(5s) = 1.74 \text{ mm}$$



Observation of long-lived $\pi^+\pi^-$ atoms



Experimental distribution of the $\pi^+\pi^-$ pairs over Q_Y . Data selected with criteria $|Q_X| < 2$ MeV/c and $|Q_L| < 2$ MeV/c. The indicated peaks are due to $\pi^+\pi^-$ pairs produced in the Pt foil and Be target.



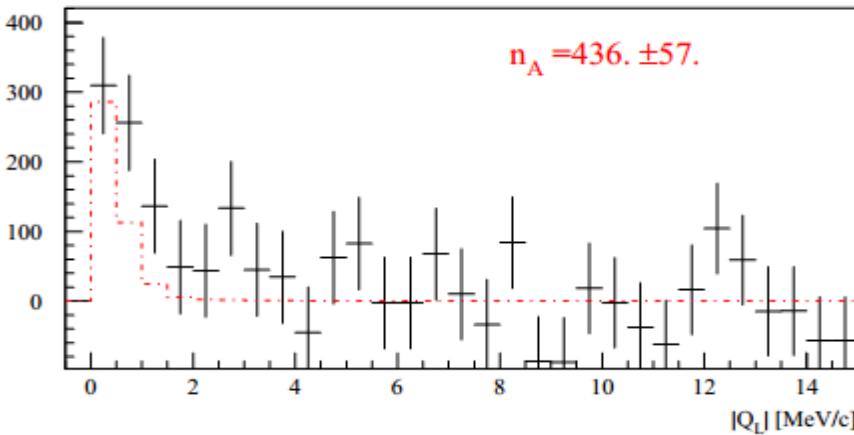
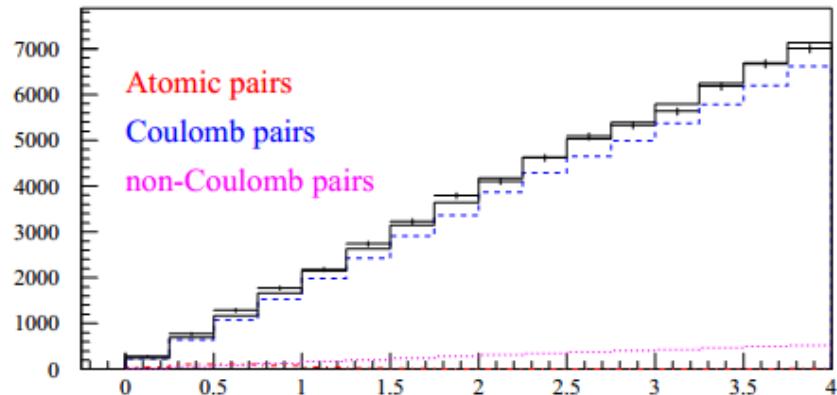
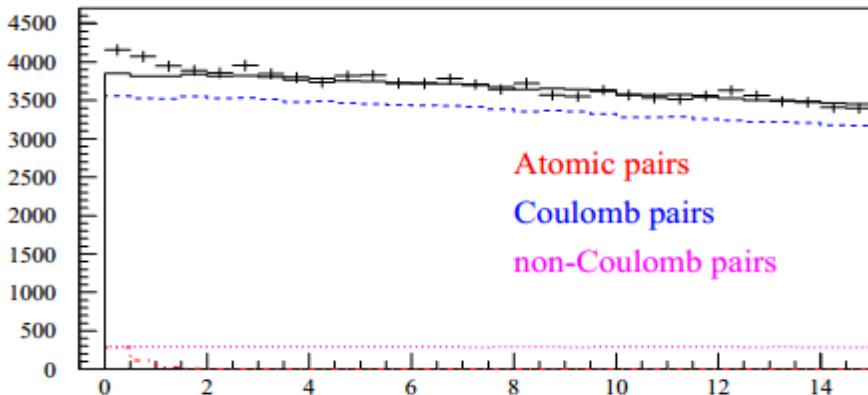
Ratio of the prompt to accidental $\pi^+\pi^-$ pairs over Q_Y projection.

Top: Experimental distribution, the peak at $Q_Y = 13.15$ MeV/c corresponds to the Coulomb pairs produced in the Be target.

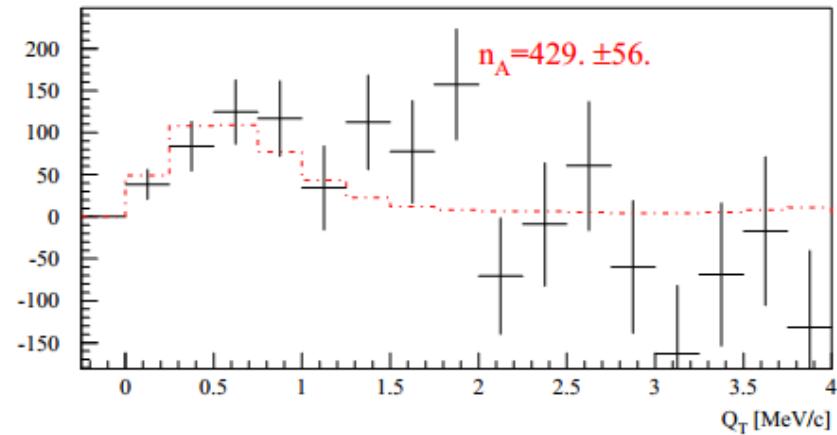
Bottom: Simulated distribution.

Observation of long-lived $\pi^+\pi^-$ atoms

Two-dimensional distribution over $|Q_L| Q_T$, have been fitted with $\chi^2/\text{ndf} = 138/140$. Projections to $|Q_L|$ and Q_T are presented.



$|Q_L|$ for $Q_T < 2.0 \text{ MeV}/c$



Q_T for $|Q_L| < 2.0 \text{ MeV}/c$

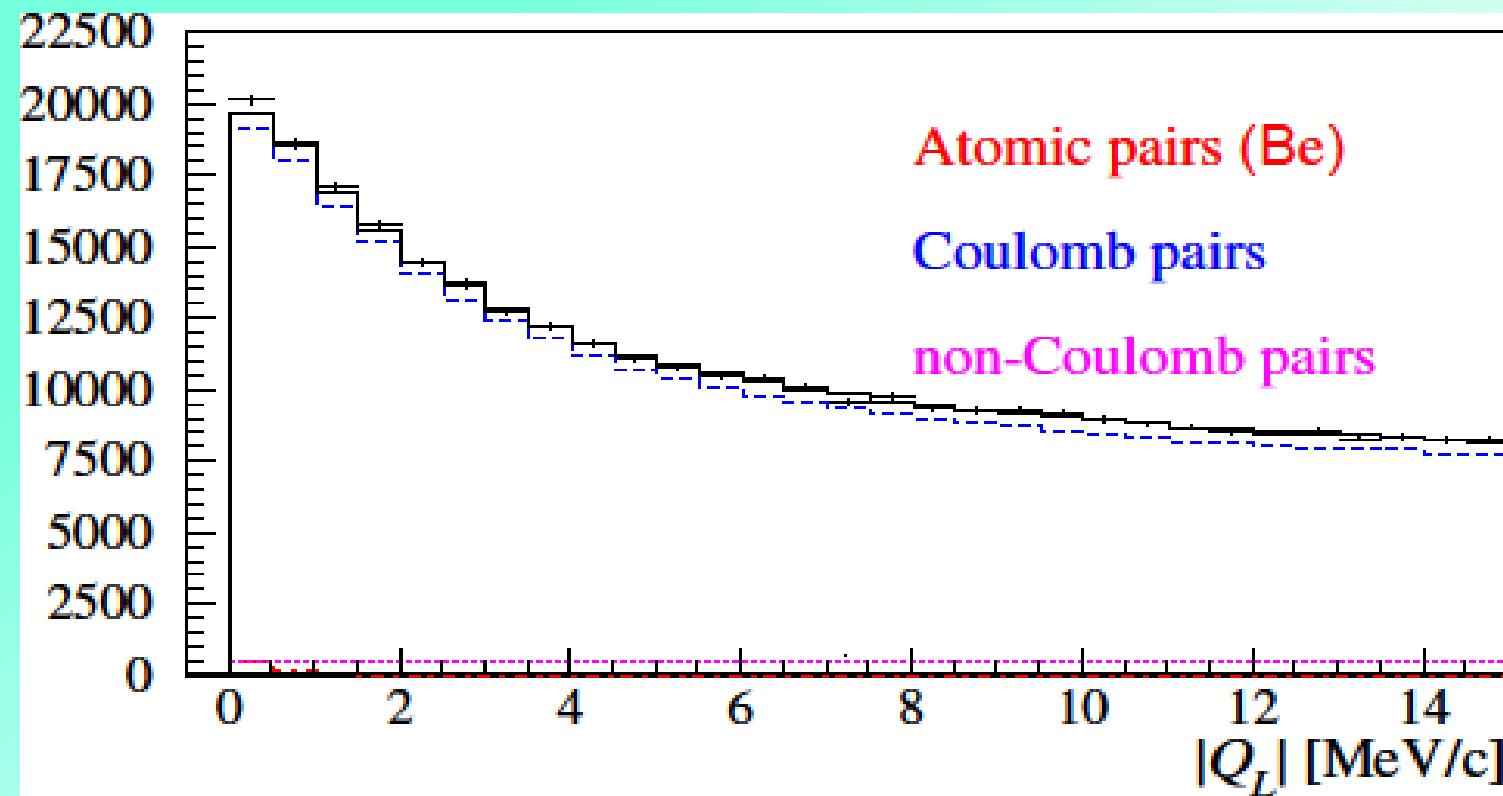
Observation of long-lived $\pi^+\pi^-$ atoms

Q_T cut	n_A	n_A^{tot}	Background	χ^2/ndf
$Q_T < 2.0 \text{ MeV}/c$	436 ± 57 $(\sim 7.6\sigma)$	488 ± 64	16719	138/140

$$n_A^L = 436 \pm 57(\text{stat}) \pm 23(\text{syst}) = 436 \pm 61(\text{tot})$$

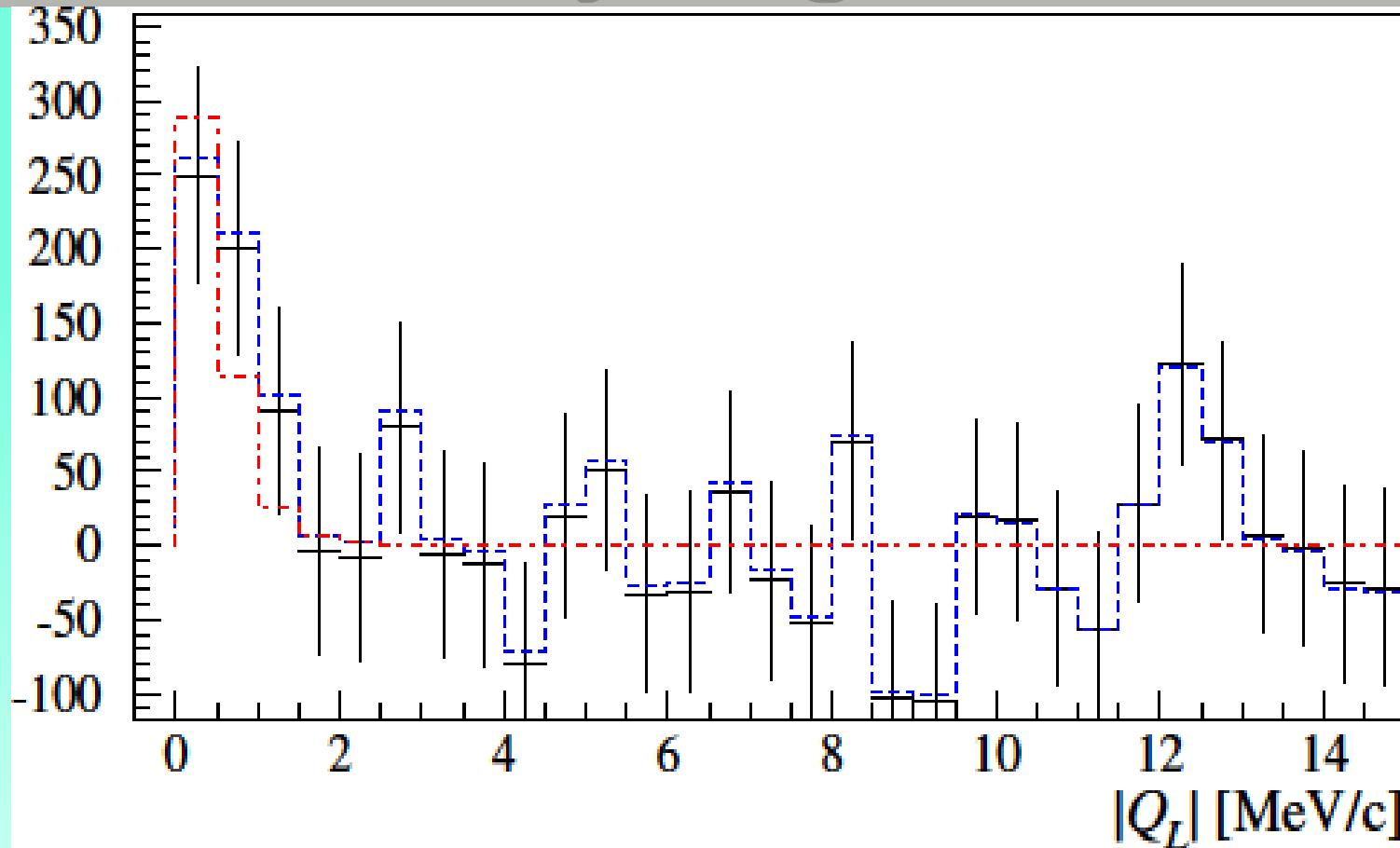
B.Adeva et al., Phys. Lett. B 751 (2015) 12

Lifetime of long-lived $\pi^+\pi^-$ atoms



$$N_C^{Be} = 319890 \pm 2610 \longrightarrow N_A = 16960 \pm 130 \longrightarrow N_A^L = 1021^{+20}_{-18}$$

Observation of long-lived $\pi^+\pi^-$ atoms



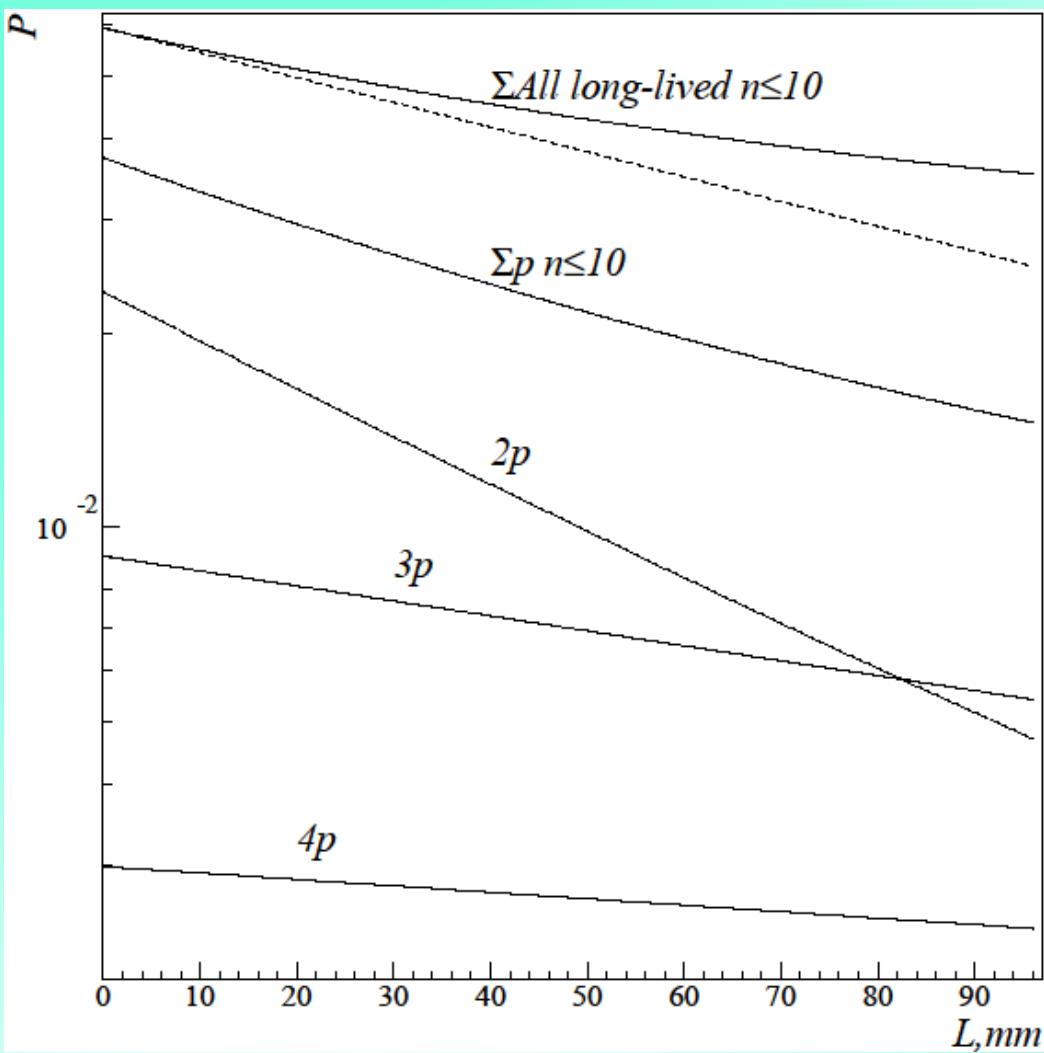
$|Q_L|$ experimental distribution after subtraction of background obtained with 3 parameter fit (black points with statistical error) and after subtraction of background obtained with 2 parameter fit (blue dashed line), comparing to the simulated distribution of atomic pairs (red dotted-dashed line).

The fit procedures have been applied to the 1-dimensional $|Q_L|$ distribution.

The atomic pairs number in the region $|Q_L| < 2, Q_T < 4 \text{ MeV}/c$ obtained with 3 parameter fit is $n_A^L = 435 \pm 103$ and with 2 parameter fit is $n_A^L = 579 \pm 164$.

Lifetime of long-lived $\pi^+\pi^-$ atoms

Alteration of atomic states populations at the gap between Be and Pt

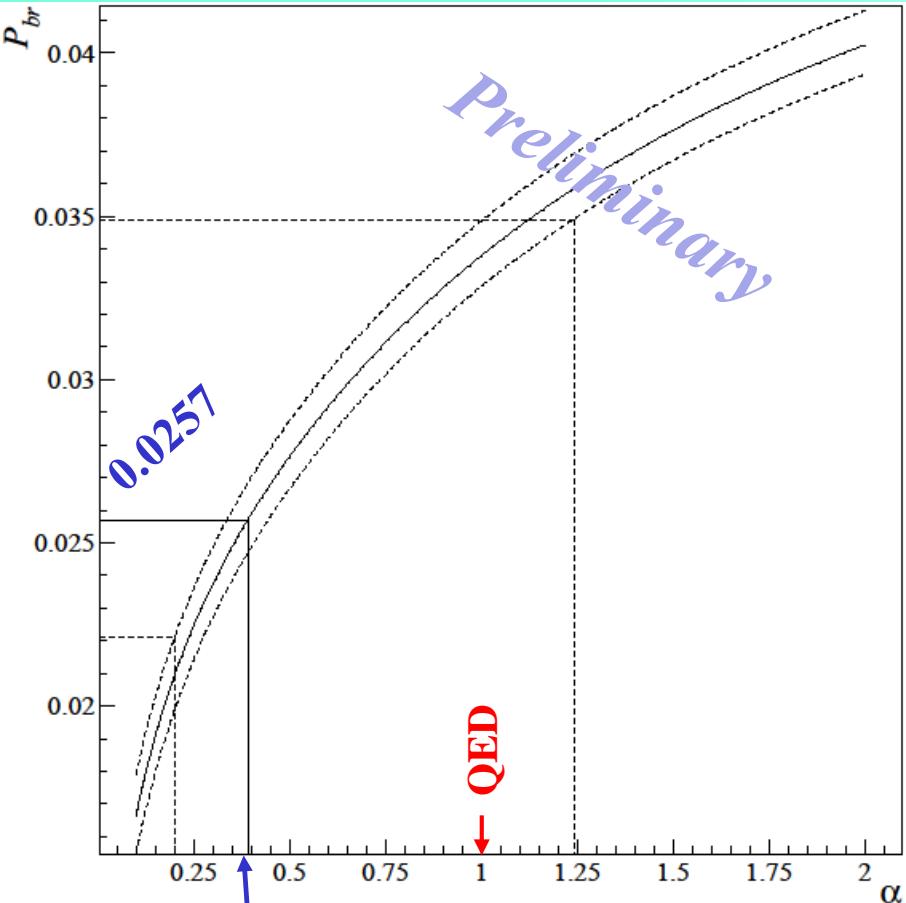


Simplified exponential decay

$$\tau^L = 2.4^{+1.0}_{-0.4} \Big|_{tot} \times 10^{-11} \text{ s}$$

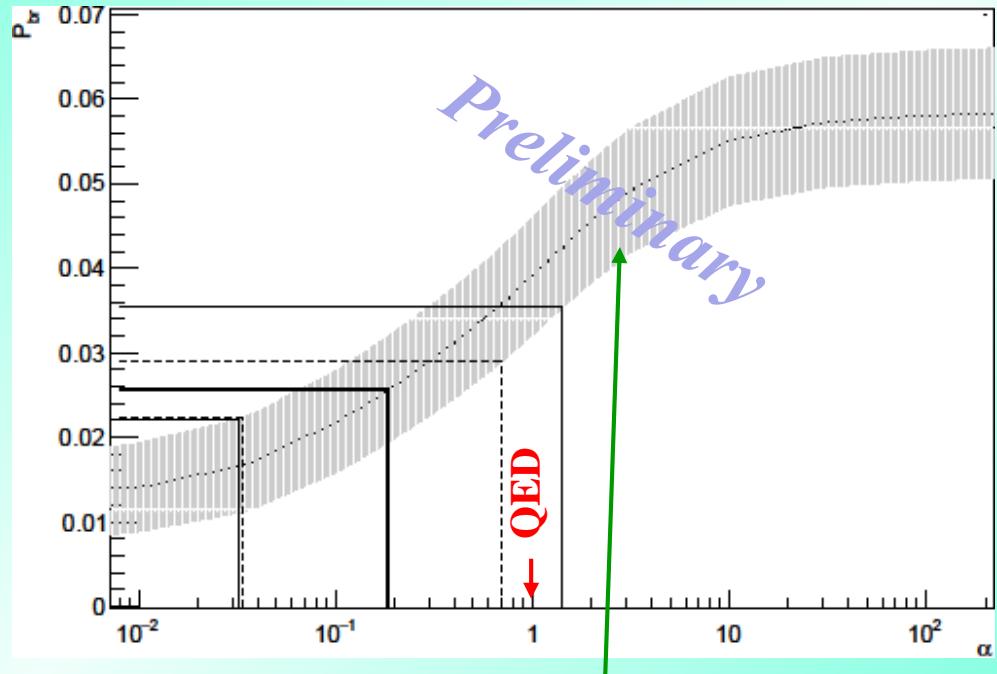
$$P_{br}^L = \frac{n_A^L}{N_A} = 0.0257 \pm 0.0034 \Big|_{stat} \begin{array}{l} +0.0086 \\ -0.0014 \end{array} \Big|_{syst} = 0.0257 \begin{array}{l} +0.0092 \\ -0.0036 \end{array} \Big|_{tot}$$

Lifetime of long-lived $\pi^+\pi^-$ atoms



$$\alpha = 0.39^{+0.85}_{-0.19}$$

$$\tau_{2p}^{free} = 0.46^{+0.99}_{-0.22} \times 10^{-11} \text{ s}$$



$$\alpha = \frac{\tau_i^{free}}{\tau_i^{QED}}$$

$$\alpha = 0.185^{+1.22}_{-0.15}$$

$$\tau_{2p}^{free} = 0.22^{+1.42}_{-0.18} \times 10^{-11} \text{ s}$$

$$\tau_{2p}^{QED} = 1.17 \times 10^{-11} \text{ s}$$

$$\tau_{1s}^{DIRAC} = 3.15^{+0.28}_{-0.26} \times 10^{-15} \text{ s}$$

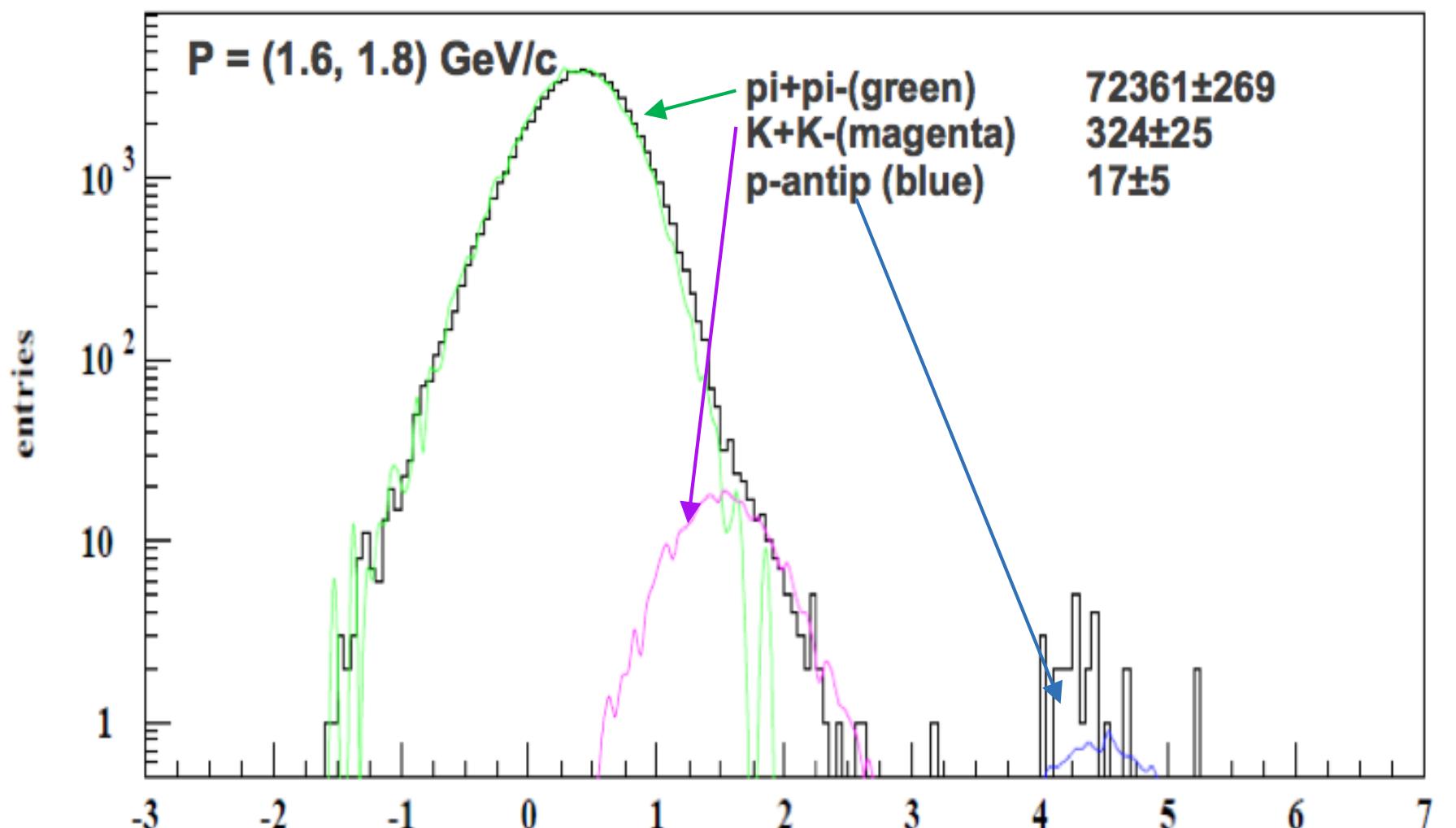
K⁺K⁻ atom and its lifetime

The A_{2K} lifetime is strongly reduced by strong interaction (OBE, scalar meson f₀ and a₀) as compared to the annihilation of a purely Coulomb-bound system (K⁺K⁻).

τ (A _{2K} → ππ, πη)	K ⁺ K ⁻ interaction
1.0×10^{-18} s [1]	Coulomb-bound
8.5×10^{-18} s [3]	momentum dependent potential
3.2×10^{-18} s [2]	+ one-boson exchange (OBE)
1.1×10^{-18} s [2]	+ f'_0 (I=0) + πη-channel (I=1)
2.2×10^{-18} s [4]	ChPT

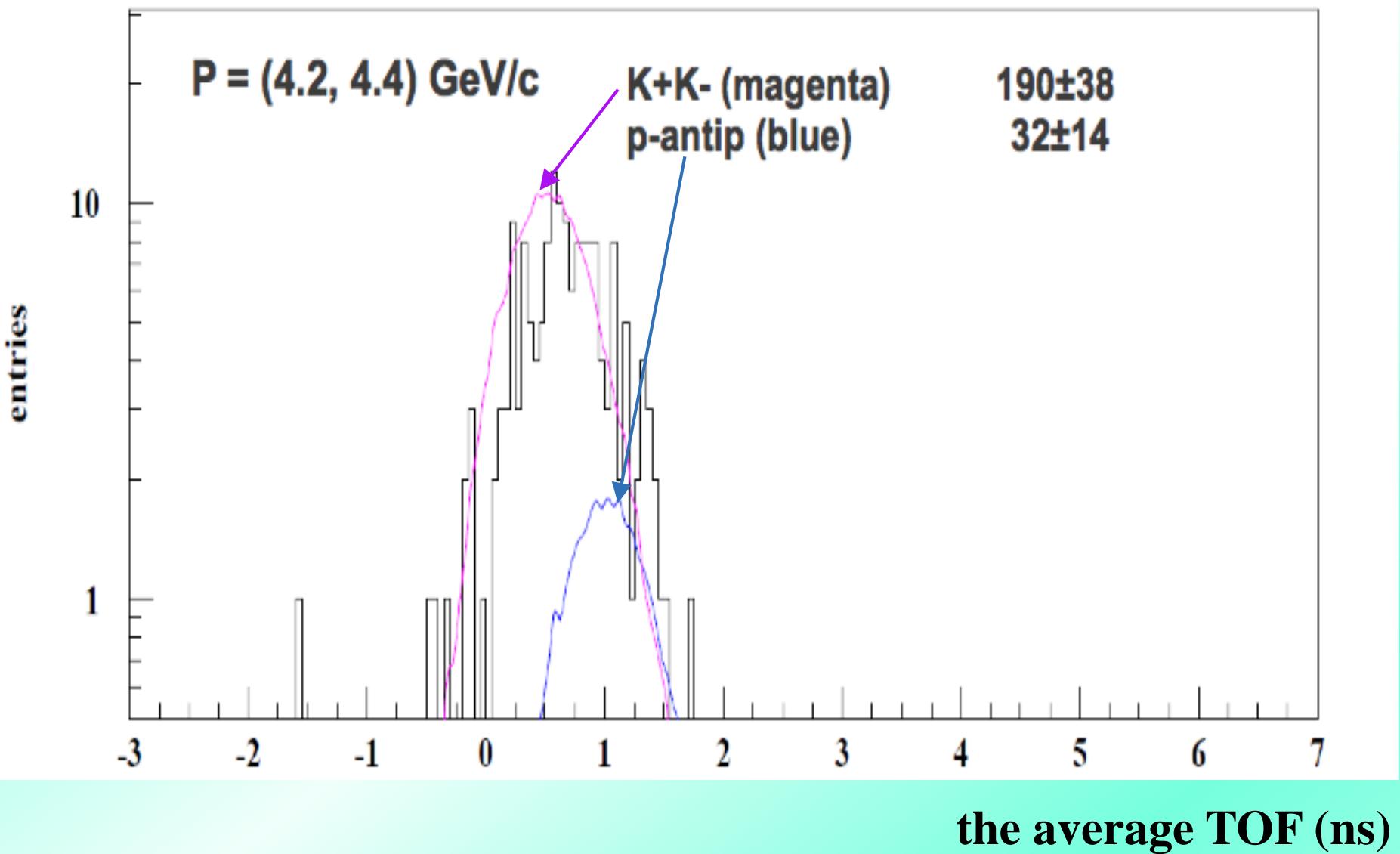
- References:** [1] S. Wycech, A.M. Green, Nucl. Phys. A562 (1993), 446;
[2] S. Krewald, R. Lemmer, F.P. Sasson, Phys. Rev. D69 (2004), 016003;
[3] Y-J Zhang, H-C Chiang, P-N Shen, B-S Zou, PRD74 (2006) 014013;
[4] S.P. Klevansky, R.H. Lemmer, PLB702 (2011) 235.

Search of K^+K^- pair with TOF (low momentum)

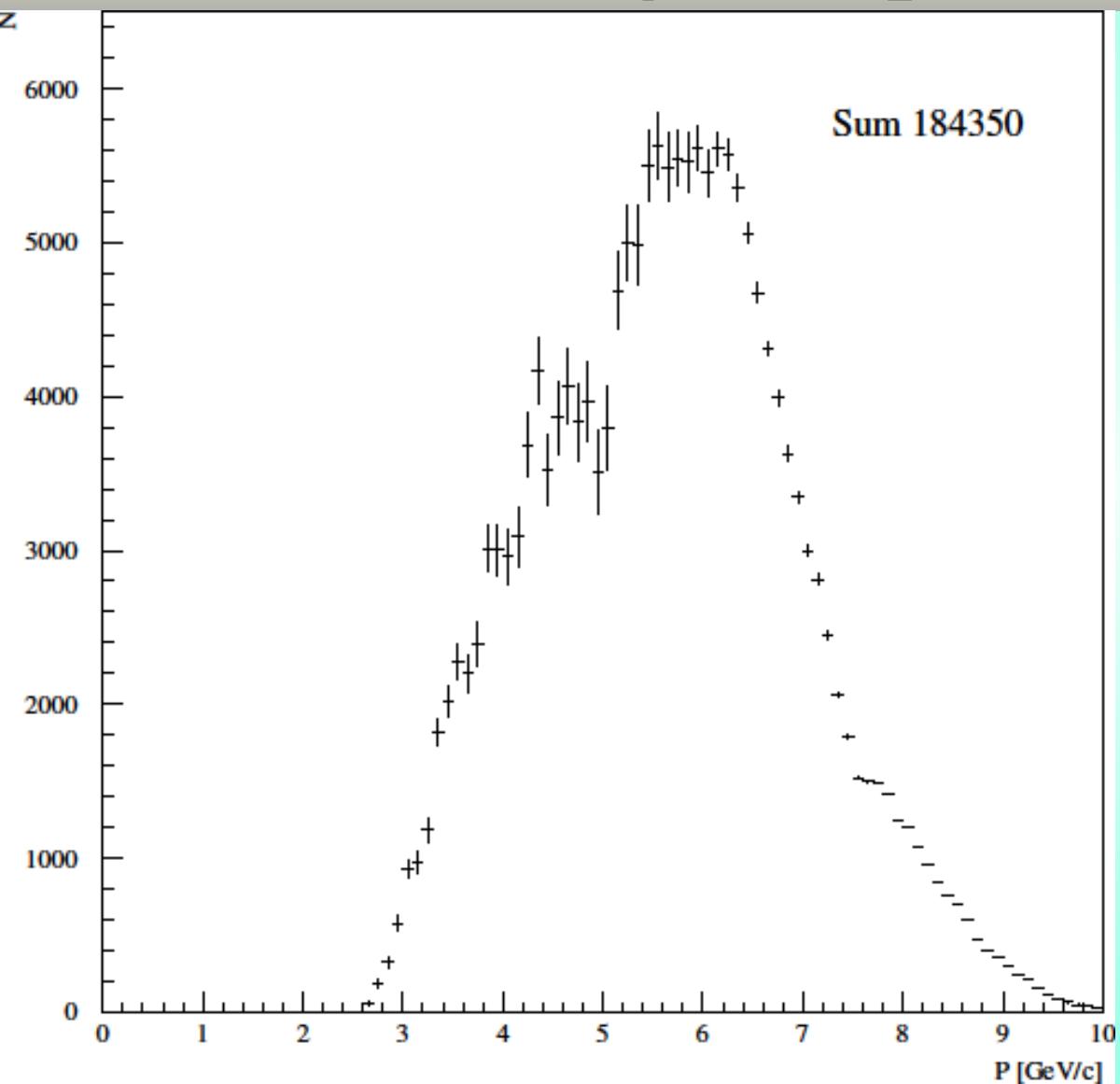


the average TOF (ns)

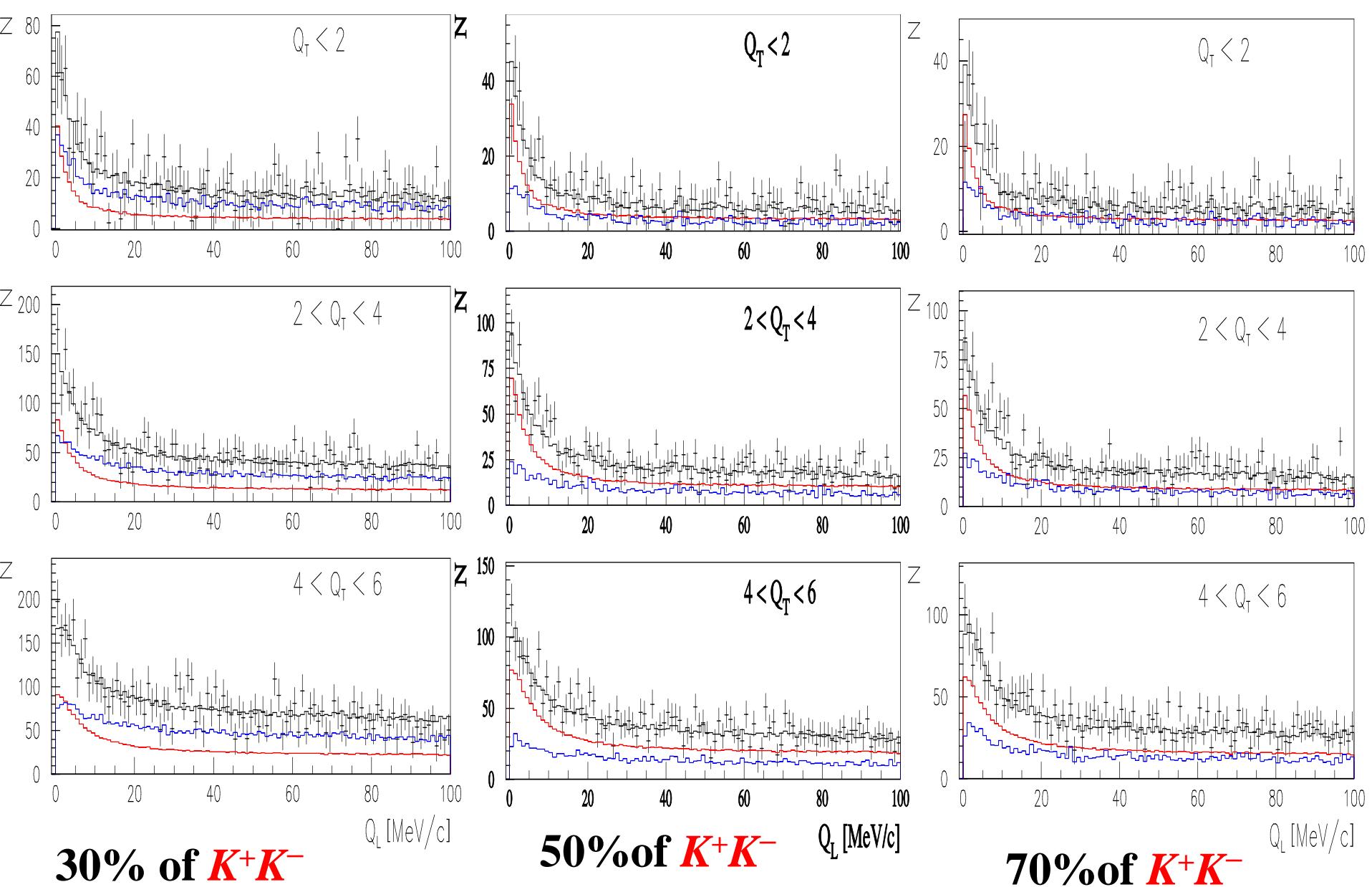
Search of K^+K^- pair with TOF (high momentum)



Search of K^+K^- pair with TOF



Distribution over the full pair momentum in lab system with contamination 50% of K^+K^- pairs



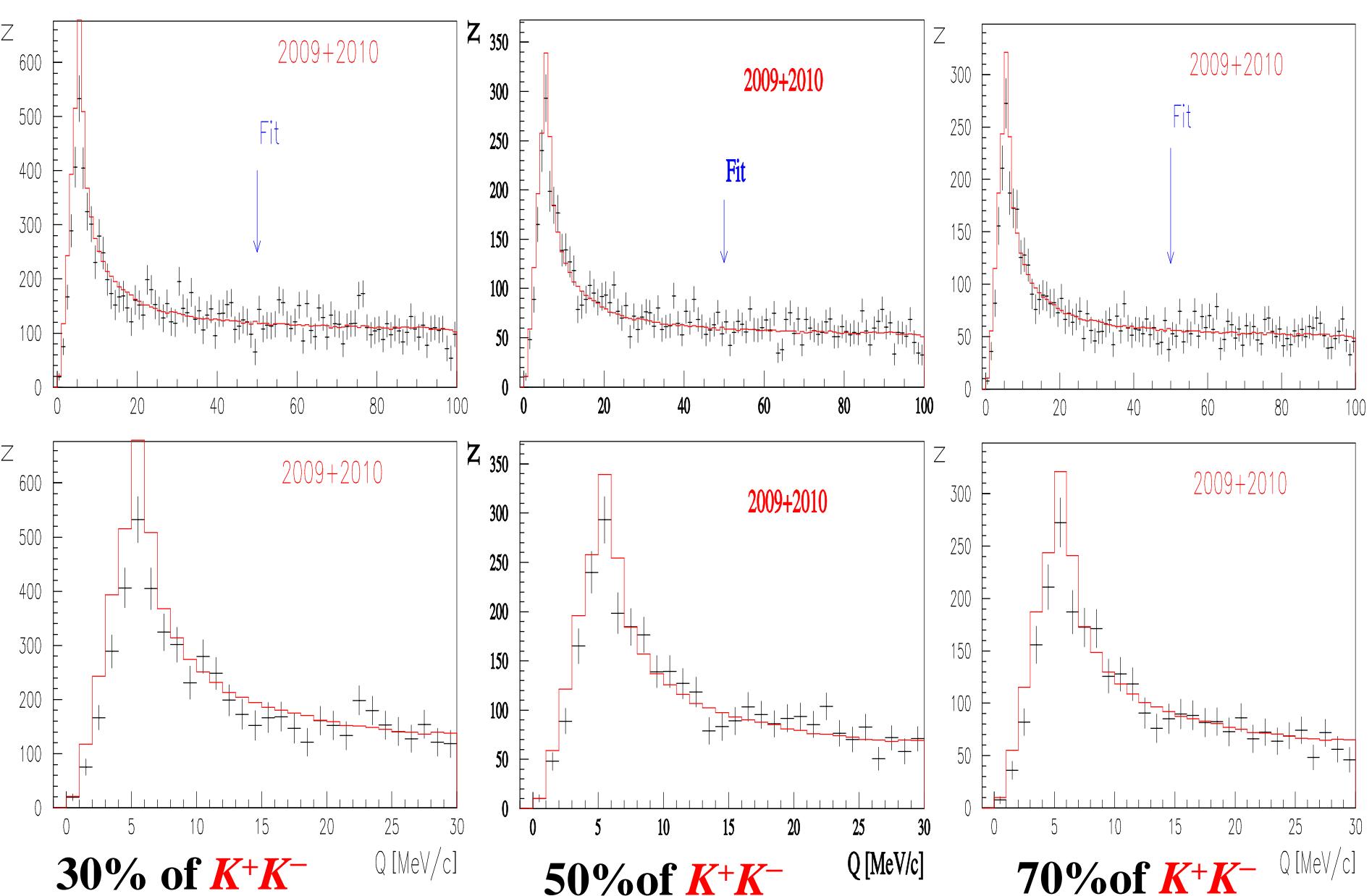
30% of K^+K^-

50% of K^+K^-

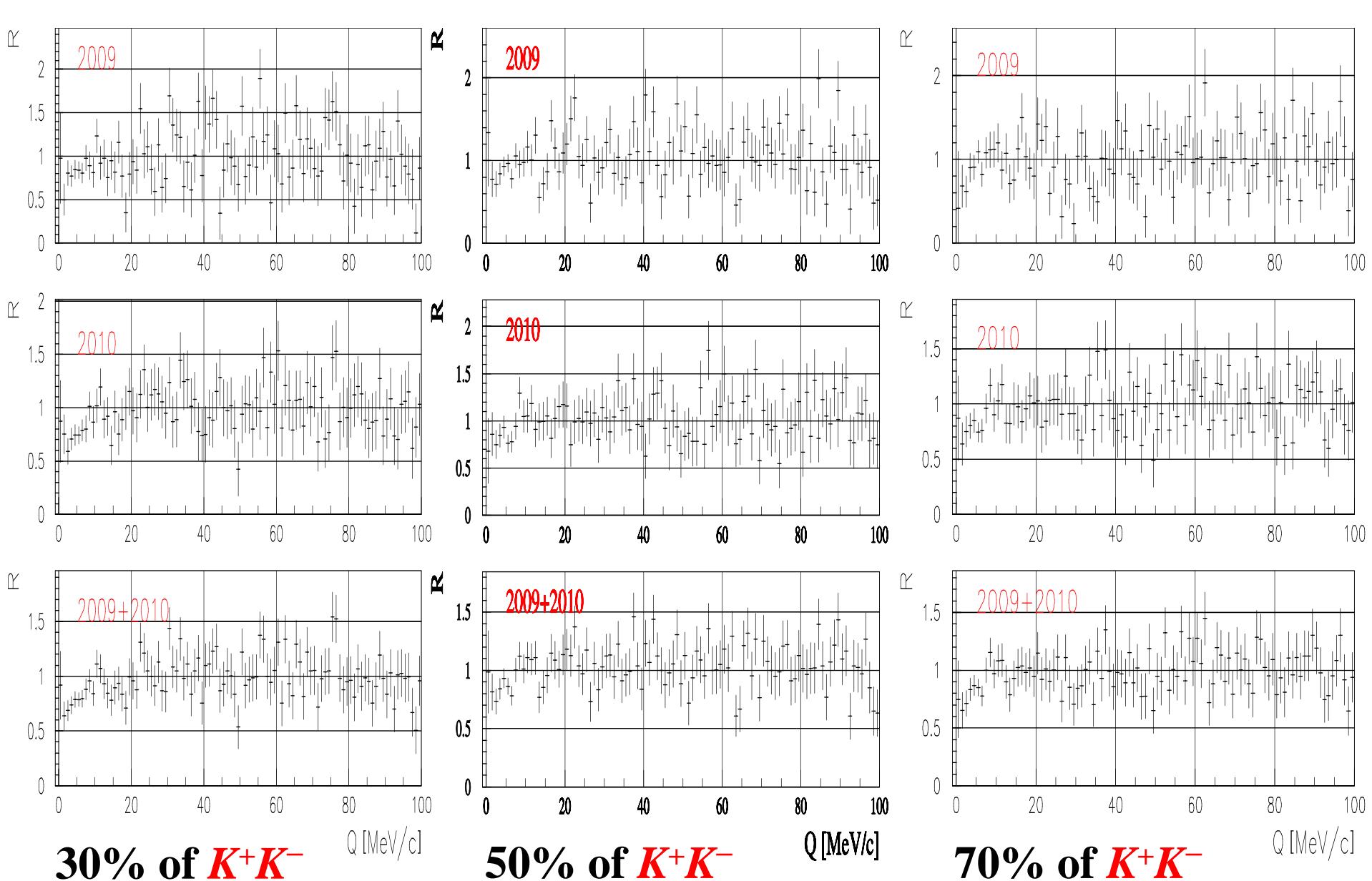
Q_L [MeV/c]

70% of K^+K^-

**Fit with simulated mixture
of K^+K^- and $\pi^+\pi^-$ pairs**

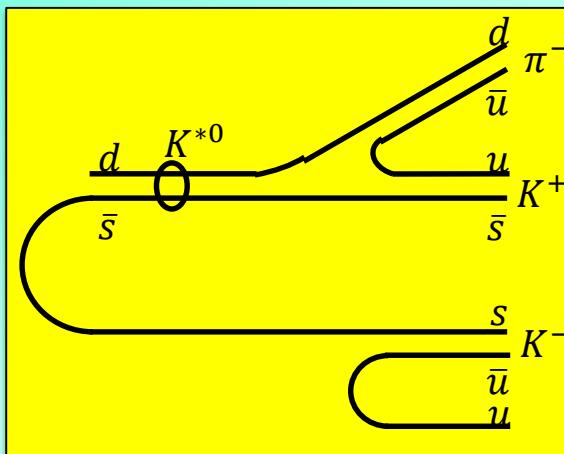
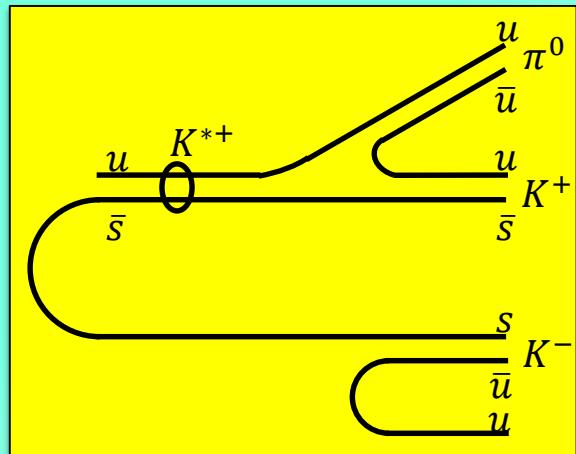


**Fit with simulated point-like Coulomb pairs,
normalized at $Q>50$ MeV/c**



Ratio of experimental to simulated point-like Coulomb pairs
Deep at origin is due to presence of non point-like sources of K^+K^-

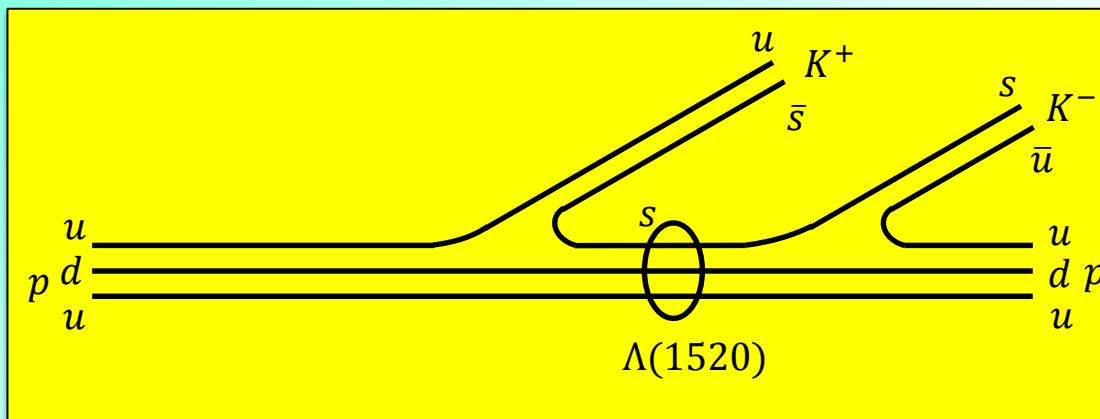
Scale of Coulomb interaction is K^+K^- Bohr radius 109 fm



$K^{*+}(892)$ $K^{*0}(892)$

$c\tau \approx 4.2$ fm

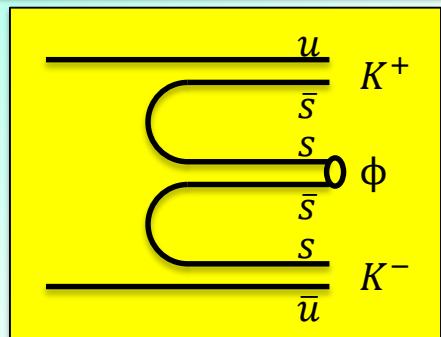
$(25.3 \pm 5.5)\%$



$\Lambda(1520)$

$c\tau \approx 12.6$ fm

$(8.1 \pm 2.0)\%$



ϕ

$c\tau \approx 46$ fm

$(1.8 \pm 0.4)\%$

**Thank you
for your attention!**

Theoretical motivation

$\pi\pi$ scattering length

In ChPT the effective Lagrangian, which describes the $\pi\pi$ interaction, is an expansion in terms:

$$L_{eff} = L^{(2)} + L^{(4)} + L^{(6)} + \dots$$

(tree) (1-loop) (2-loop)

G. Colangelo, J. Gasser and H. Leutwyler, Nucl. Phys. B603 (2001) 125,
using ChPT (2-loop) & Roy equations:

$$\left. \begin{array}{l} a_0 = 0.220 \pm 2.3\% \\ a_2 = -0.0444 \pm 2.3\% \end{array} \right\} a_0 - a_2 = 0.265 \pm 1.5\%$$

These results (precision) depend on the low-energy constants (LEC) \overline{l}_3 and \overline{l}_4 :
Lattice gauge calculations from 2006 provided values for these \overline{l}_3 and \overline{l}_4 .

Theoretical motivation

Lattice calculations of \bar{l}_3, \bar{l}_4

- 2006: \bar{l}_3, \bar{l}_4 ... first lattice calculations
- 2012: 10 collaborations: 3 in USA, 5 in Europe and 2 in Japan
- J. Gasser, H. Leutwyler: model calculation (1985)
 $\bar{l}_3 = 2.9 \pm 2.4, \quad \bar{l}_4 = 4.3 \pm 0.9$
- Lattice calculations of these constants have been done in 20 works.
Best result: $l_3 = 2.6 \pm 0.5 \text{st} \pm 0.4 \text{sy}, \quad l_4 = 3.8 \pm 0.4 \text{st} \pm 0.2 \text{sy}$

Therefore, the theoretical pion-pion scattering length precision can be improved.

The best experimental results on the scattering length have a precision not better than 4%.

Lifetime and breakup probability

The P_{br} value depends on the lifetime value, τ . To obtain the precise $P_{br}(\tau)$ curve a large differential equation system must be solved:

$$\frac{dp_{nlm}(s)}{ds} = \sum_{n'l'm'} a_{nlm}^{n'l'm'} p_{n'l'm'}(s)$$

where s is the position in the target, p_{nlm} is the population of a definite hydrogen-like state of pionium. The $a_{nlm}^{n'l'm'}$ coefficients are given by:

$$a_{nlm}^{n'l'm'} = \frac{\sigma_{nlm}^{n'l'm'} \rho N_0}{A} \quad \text{if } nlm \neq n'l'm', \quad a_{nlm}^{nlm} = -\frac{\sigma_{nlm}^{tot} \rho N_0}{A} - \begin{cases} 2M_\pi / P c \tau_n & l=0, \\ 0 & l \neq 0. \end{cases}$$

$\sigma_{nlm}^{n'l'm'}$ being the pionium-target atom cross section, N_0 the Avogadro Number, ρ the material density and A its atomic weight.

The detailed knowledge of the cross sections (Afanashev&Tarasov, Trautmann et al) (Born and Glauber approach) together with the accurate solution of the differential equation system permits us to know the curves within 1%.

$\pi^+\pi^-$ experimental results

K \rightarrow 3 π :

(scattering length in m_π^{-1})

2009 NA48/2 (EPJ C64, 589)

$$\Rightarrow a_0 - a_2 = 0.2571 \pm 0.0048 \Big|_{stat} \pm 0.0025 \Big|_{syst} \pm 0.0014 \Big|_{ext} = \dots \pm 2.2\%$$

Ke4:

plus additional 3.4% theory uncertainty

2010 NA48/2 (EPJ C70, 635)

$$\Rightarrow a_0 = 0.2220 \pm 0.0128 \Big|_{stat} \pm 0.0050 \Big|_{syst} \pm 0.0037 \Big|_{theo} = \dots \pm 6.4\%$$

$$\Rightarrow a_2 = -0.0432 \pm 0.0086 \Big|_{stat} \pm 0.0034 \Big|_{syst} \pm 0.0028 \Big|_{theo} = \dots \pm 22\%$$

$\pi^+\pi^-$ atom:

2011 DIRAC (PLB 704, 24)

$$\Rightarrow |a_0 - a_2| = 0.2533 \begin{array}{l} +0.0078 \\ -0.0080 \end{array} \Big|_{stat} \begin{array}{l} +0.0072 \\ -0.0077 \end{array} \Big|_{syst} = \dots \begin{array}{l} +4.2\% \\ -4.4\% \end{array}$$

Experiment DIRAC at SPS CERN

In 2013 DIRAC setup has been dismantled from the experimental hall of PS CERN. All detectors are stored for using in the future experiment.

DIRAC collaboration is planning to continue investigation of π^-K^+ , π^+K^- and $\pi^+\pi^-$ atoms at SPS accelerator at CERN. The correspondent gains in production rates of these atoms at SPS relative to PS (450 GeV vs. 24 GeV) are 18, 24 and 12. This allows to increase significantly the collected data and to check the precise prediction of Low-Energy QCD at a higher accuracy. Now the collaboration is planning to submit the Letter of Intend for study πK and $\pi^+\pi^-$ atoms at SPS to SPSC CERN.