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# Low-Temperature Enhanced Semi-Annihilation and AMS-02

Andrew Spray & Yi Cai JHEP 1601 (2016) 087 JHEP 1606 (2016) 156 JHEP 1702 (2017) 120 1807.00832

## Motivation

#### Bounds on thermal DM starting to get quite strong

LPT-Orsay-17-09, CPHT-RR009.032017, SCIPP 17/03

The Waning of the WIMP? A Review of Models, Searches, and Constraints

Giorgio Arcadi<sup>*a*</sup>, \* Maíra Dutra<sup>*b*</sup>, † Pradipta Ghosh<sup>*b,c*</sup>, † Manfred Lindner <sup>*a*</sup>, § Yann Mambrini<sup>*b*</sup>, ¶ Mathias Pierre<sup>*b*</sup>, \*\* Stefano Profumo<sup>*d,e*</sup>, †† and Farinaldo S. Queiroz<sup>*a*‡‡</sup>

- Successful test of this idea!
- But we should be diligent in checking for loopholes
- What are our assumptions?
   What if we relax them?





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- Successful test of this idea!
- But we should be diligent in checking for loopholes
- What are our assumptions?
   What if we relax them?
- Very basic assumption:
   DM stabilised by Z<sub>2</sub> symmetry







- Implies this familiar diagram
- Detection rates related to relic density calculation
- Leads to these strong bounds



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- Detection rates related to relic density calculation
- Leads to these strong bounds



- \* Not Generic! (D'Eramo & Thaler, 2010)
- \* Non- $Z_2$  syms  $\rightarrow$  Semi-Annihilation:
  - Non-decay processes
  - Odd number of external dark states
- Irrelevant for colliders & DD



Leads to these strong bounds

Irrelevant for colliders & DD

BUT constraints weakened, not removed



 Semi-annihilation leads to cosmic ray signals



\* BUT constraints weakened, not removed



\* Semi-annihilation leads to cosmic ray signals  $\chi \qquad \sigma_{RD/ID} \qquad \chi'$  BUT constraints weakened, not removed



- Important to understand model space, phenomenology and thus constraints
- \* Initial effort in EFT language, JHEP 1702 (2017) 120

### \* $2 \rightarrow 2$ SA phenomenologically limited





- \*  $2 \rightarrow 2$  SA phenomenologically limited
- \*  $2 \rightarrow 3+$  SA phase space suppressed... ... unless it is really  $2 \rightarrow 2$











- EFT approach assumes cross section simple function of velocity
- Exceptions (enhanced at low v) are well-known:
  - Sommerfeld, bound states, Breit-Wigner resonance
- Bigger signals today >>
   phenomenologically interesting
- Thermal history becomes sensitive to DM temperature



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## Dark Matter Temperature and Semi-annihilation

\* Definition of DM temperature  $T_{\chi}$ : (Binder et al, 1706.07433)

Maxwellian

$$f(E) = e^{-E/T_{\chi}}$$

$$T_{\chi} = \frac{g_{\chi}}{n_{\chi}} \int \frac{d^3 p_{\chi}}{(2\pi)^3} \frac{\vec{p}_{\chi}^2}{3E_{\chi}} f(p_{\chi})$$

Usual behaviour:



Details require solving an extra Boltzmann equation



$$\langle \sigma v \rangle_2 = \frac{g_{\chi}^2}{n_{\chi}^3} \int \frac{d^3 p_1}{(2\pi)^3} \, \frac{d^3 p_2}{(2\pi)^3} \, \frac{p_{\chi}^2}{\mathbf{3}T_{\chi}E_{\chi}} \, f_1(p_1) \, f_2(p_2) \, \sigma v$$

- \* Average of  $p^2/E$  for external DM particle
- Effect of SA first discussed in Kamada et al [1707.09238]:
  - Must include forward and backward contributions
  - \* Self-heating: fraction  $q_{\chi}$  of mass converted to DM kinetic energy

$$\frac{dT_{\chi}}{dT} \sim \frac{1}{n_{\chi}} \frac{dt}{dT} \frac{dE_{\chi}}{dt} \sim \frac{1}{n_{\chi}} \frac{1}{HT} q_{\chi} m_{\chi} n_{\chi}^2 \langle \sigma v(SA) \rangle$$
$$\sim \frac{s}{HT} q_{\chi} m_{\chi} Y_{\chi} \langle \sigma v(SA) \rangle$$



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Constant (rad. dom.)



$$\langle \sigma v \rangle_2 = \frac{g_{\chi}^2}{n_{\chi}^3} \int \frac{d^3 p_1}{(2\pi)^3} \, \frac{d^3 p_2}{(2\pi)^3} \, \frac{p_{\chi}^2}{\mathbf{3}T_{\chi}E_{\chi}} \, f_1(p_1) \, f_2(p_2) \, \sigma v$$

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Constant (rad. dom.) Constant

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Constant (rad. dom.) Constant (relic dens.)



$$\langle \sigma v \rangle_2 = \frac{g_{\chi}^2}{n_{\chi}^3} \int \frac{d^3 p_1}{(2\pi)^3} \frac{d^3 p_2}{(2\pi)^3} \frac{p_{\chi}^2}{\mathbf{3}T_{\chi}E_{\chi}} f_1(p_1) f_2(p_2) \sigma v$$

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$$\frac{dT_{\chi}}{dT} \sim \frac{s}{HT} q_{\chi} m_{\chi} Y_{\chi} \langle \sigma v(SA) \rangle$$
$$T_{\chi} \propto T \qquad \text{in far IR}$$



- Situation with dark partners more complex
- Dark partner number density?



Does dark partner scatter before decaying?
 Do we need to track dark partner temperature?

$$\frac{dy_{\chi}}{dx} \supset \frac{1}{xHZ} \left\langle \frac{p_{\chi}^2}{3E_{\chi}T_{\chi}} \right\rangle \Gamma_{\Psi} \left( \frac{Y_{\Psi}}{Y_{\chi}} - \frac{Y_{\Psi}^{eq}}{Y_{\chi}^{eq}} \mathcal{D}_{\mathcal{T}}(T, T_{\chi}) \right)$$

General conclusion is unchanged

 $T_{\chi} \propto T$  in far IR



- SADM vs annihilating DM with low-temperature enhancements:
  - Warmer at late times
  - Smaller rates after kinetic decoupling
  - Relic density for larger couplings
  - Larger possible signals today
- Annihilating DM max signal ~100 times thermal cross section
- What about SADM?





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# Case Study: AMS-02

- Positron anomaly
  - Old excess in  $e^+$  at E > 10 GeV
  - AMS-02 most recent, best precision
  - Could be astrophysics, esp. pulsars (HAWC observations of Geminga)
- Particle DM explanations:
  - Dominantly produce leptons (no antiproton excess)
  - \* Direct  $e^+$  typically a poor fit
  - \* Need large cross sections  $\sigma \sim 10^3\,{\rm pb}$





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Simplified model



- \* Coupling structure could derive from  $U(1)_{L\mu-L\tau}$
- Two dark partner decay modes:



Sequential 2-body





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- Can fit excess well, though constrained by CMB
- Very large signals possible
  - 10<sup>5</sup> enhancement consistent with relic density
  - Smaller enhancement requires less tuning of parameters







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## Conclusions

- Semi-annihilation is a generic feature of dark matter stabilised by any symmetry other than a Z<sub>2</sub>
- It eases the bounds from colliders and direct detection
- Cosmic ray observations are relevant, motivating models where SA today is enhanced
- \* All SADM models redshift like radiation,  $T_{\chi} \sim T$ , in far IR
- Warmer DM allows signal enhanced by up to 10<sup>5</sup>, much greater than possible for annihilating DM
- We have used this to explain AMS-02 positron excess

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   Thank You!