Precision Higgs Measurements at the 250 GeV ILC

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After the Discovery of the Higgs Boson

A major theme in particle physics is to discover
 new phenomena or principles which can explain
 1) EWSP filling the Higgs field in the year

- 1). EWSB filling the Higgs field in the vacuum,
- 2). A dark matter candidate
- 3). matter-antimatter asymmetry
- ... neutrino mass ... fine tuning

Standard Model can not explain these facts.





ttps://www.symmetrymagazine.org/article/december-2013/four-things-you-might-not-know-about-dark-matte

The Higgs boson is a window to new physics beyond the SM.

Required Higgs Precision

O Expected deviations for typical BSM scenarios.

Super Symmetry:
$$\frac{g_{hbb}}{g_{h_{SM}bb}} = \frac{g_{h\tau\tau}}{g_{h_{SM}\tau\tau}} \simeq 1 + 1.7\% \left(\frac{1 \text{ TeV}}{m_A}\right)^2$$

 m_A : CP-odd Higgs A⁰

Composite Higgs:
$$\frac{g_{hVV}}{g_{h_{\rm SM}VV}} \simeq 1 - 3\% (1 \text{ TeV}/f)^2$$
$$\frac{g_{hff}}{g_{h_{\rm SM}ff}} \simeq \begin{cases} 1 - 3\% (1 \text{ TeV}/f)^2 & (\text{MCHM4})\\ 1 - 9\% (1 \text{ TeV}/f)^2 & (\text{MCHM5}) \end{cases}$$

f: a compositeness scale

Mixing with
a singlet scalar:
$$\frac{g_{hVV}}{g_{h_{\rm SM}VV}} = \frac{g_{hff}}{g_{h_{\rm SM}ff}} = \cos\theta \simeq 1 - \frac{\delta^2}{2}$$

 $h = h_{\rm SM} \cos \theta + S \sin \theta$

○ Predicted deviations are small, a few % ~ 10 % level.

O(1%) precision (model independent) is needed to see the deviations.

Fingerprinting BSM models

O Different BSM models predict different deviation patterns.

2HDM model (m_A =600 GeV, tan β =7)

Composite Higgs (*f*=1.2 TeV)



The ILC250 has the capability to tell the nature of the BSM from its deviation patterns !

The ILC250 Higgs production

ILC-Technical Design Repor arXiv1306.6352 [hep-ex]



• The Higgs-strahlung production is maximum at $\sqrt{s} = 250$ GeV.

Beam Polarization is available.
 Pol(e-,e+)=(-80%,+30%) & (+80%,-30%)

O Accumulate data for 10 years, 2000 fb⁻¹.

Zh : ~ 500 K Higgs boson. ww-f : ~ 15 K Higgs boson.

The ILC250 key measurement σ_{zh}

Phys. Rev. D 94, 113002 (2016) Eur. Phys. J. C (2016) 76:72

 \bigcirc Unique measurement at lepton colliders is absolute σ_{zh} .

→ A key for determining Higgs couplings model independently.



Other direct Higgs Observables at the ILC250

 $\sigma_{ZH} \times Br(H \longrightarrow bb), \sigma_{\nu\nu H} \times Br(H \longrightarrow bb)$ $\sigma_{ZH} \times Br(H \rightarrow cc)$ $\sigma_{ZH} \times Br(H \rightarrow gg)$ $\sigma_{ZH} \times Br(H \longrightarrow WW^*)$ $\sigma_{ZH} \times Br(H \longrightarrow ZZ^*)$ $\sigma_{ZH} \times Br(H \longrightarrow \tau\tau)$ $\sigma_{ZH} \times Br(H \longrightarrow \gamma \gamma)$ $\sigma_{ZH} \times Br(H \longrightarrow \mu\mu)$ $\sigma_{ZH} \times Br(H \longrightarrow Invisible)$ (H $\longrightarrow exotic$) : Observation is difficult at LHC

+ Differential cross-section

*all the observables are estimated based on full detector simulation of ILD & SiD



к-formalism to EFT-formalism

О A problem of the к-formalism

Higher dimensional operators (gauge invariant) are possible to induce new structures in hZZ.

$$\delta \mathcal{L} = \frac{m_Z^2}{v} (1 + \eta_Z) h Z_\mu Z^\mu + \frac{1}{2v} \zeta_Z h Z_{\mu\nu} Z^{\mu\nu}$$
New structure

Under the κ -formalism $\rightarrow \zeta z$ is assumed to be 0 \rightarrow model dependent !

к-formalism to EFT-formalism

О A problem of the к-formalism

Higher dimensional operators (gauge invariant) are possible to induce new structures in hZZ.

$$\delta \mathcal{L} = \frac{m_Z^2}{v} (1 + \eta_Z) h Z_\mu Z^\mu + \frac{1}{2v} \zeta_Z h Z_{\mu\nu} Z^{\mu\nu}$$
Under the κ -formalism $\rightarrow \zeta_Z$ is assumed to be 0 \rightarrow model dependent !
Consider the 2nd term $\rightarrow \zeta_Z$ term is composed field strength tensors:

$$\stackrel{e^{+}}{\underset{e^{-}}{\overset{Z}}} \xrightarrow{Z} \stackrel{H}{\underset{e^{-}}{\overset{Z}}} \xrightarrow{e^{+}} \xrightarrow{Z} \stackrel{H}{\underset{e^{-}}{\overset{Z}}} \xrightarrow{H} \stackrel{H}{\underset{e^{-}}{\overset{Z}}} \xrightarrow{Z} \stackrel{H}{\underset{e^{-}}{\overset{Z}}} \xrightarrow{H} \stackrel{H}{\underset{e^{-}}{\overset{Z}}} \xrightarrow{Z} \xrightarrow{T} \stackrel{H}{\underset{Z}{\overset{Z}}} \xrightarrow{T} \stackrel{h}{\underset{e^{-}}{\overset{Z}}} \xrightarrow{T} \stackrel{h}{\underset{e^{-}}{\overset{Z}} \xrightarrow{T} \stackrel{h}{\underset{e^{-}}{\overset{Z}}} \xrightarrow{T} \stackrel{h}{\underset{e^{-}}{\overset{Z}}} \xrightarrow{T} \stackrel{h}{\underset{e^{-}}{\overset{Z}}} \xrightarrow{T} \stackrel{h}{\underset{e^{-}}{\overset{Z}}} \xrightarrow{T} \stackrel{h}{\underset{e^{-}}{\overset{Z}} \xrightarrow{T} \stackrel{h}{\underset{e^{-}}{\overset{Z} \xrightarrow{T} \stackrel{h}{\underset{e^{-}}{\overset{Z}} \xrightarrow{T} \stackrel{h}{\underset{e^{-}}{\overset{Z} \xrightarrow{T} \stackrel{h}{\underset{e^{-}}{\overset{Z}} \xrightarrow{T} \stackrel{h}{\underset{e^{-}}{\overset{Z} \xrightarrow{T} \stackrel{h}{\underset{e^{-}}{\overset{Z} \xrightarrow{T} \stackrel{h}{\underset{e^{-}}{\overset{Z} \xrightarrow{T} \stackrel{h}{\underset{e^{-}}{\overset{Z} \xrightarrow{T} \stackrel{h}{\underset{E} \overset{L}{\underset{e^{-}}{\overset{T} \xrightarrow{T} \stackrel{h}{\underset{E} \overset{L}{\underset{E} \overset{L}{\underset$$

formalism

grangian at the ILC

e invariant Lagrangian in addition to the SM. A dedicated talk on EFT by Sunghoon Jung on 6/July



for contact interaction with quarks

for couplings to $\underline{b}, \underline{c}, \underline{\tau}, \underline{\mu}, \underline{g}$

<u>, g', ν, λ</u>

→invisible and exotic

nce

5

- → Improve precision of Higgs couplings
- → The LHC situation has > 50 EFT coefficients, it is not easy to determine them simultaneously.



Observables in EFT-formalism

Phys. Rev. D 97, 053004 (2018)

O ILC250 provides sufficient observables.

23 parameters can be determined simultaneously

1) Higgs-related observables

 $\rightarrow \sigma$ and $\sigma {\times} BR$...

2) Observables from angular distributions

 \rightarrow Test new Lorentz structures...

3) Triple Gauge Couplings from e⁺e⁻ → W⁺W⁻
4) Electroweak precision observables

 \rightarrow Constrain SM parameters ...

5) Beam polarizations double the number of observables
6) HL-LHC Higgs observables, BR(h→γγ, γZ)







Precision in EFT-formalism at the ILC



ILC250 provides O(1%) precisions and that model-independently !

Phys. Rev. D 97, 053003 (2018)

Discrimination of BSM models

 Significance for benchmark points in typical BSM models :
 (escape HL-LHC direct search)

SUSY, Two-Higgs Doublet, Little Higgs, Composite Higgs,

With 2000 fb⁻¹ from 10 years of **ILC250 operation**



Model discrimination in $\boldsymbol{\sigma}$

Most benchmark models are distinguishable by more than 3 sigmas.

Discrimination of BSM models

 Significance for benchmark points in typical BSM models :
 (escape HL-LHC direct search)

SUSY, Two-Higgs Doublet, Little Higgs, Composite Higgs,

- With 2000 fb⁻¹ from 10 years of **ILC250 operation**
- With 4000 fb⁻¹ from additional 10 years running of ILC500 operation



Almost all benchmark models are distinguishable by more than 5 sigmas.

Model discrimination in $\boldsymbol{\sigma}$

Summary

 The Higgs boson is a window to new physics.
 The precision study of the Higgs couplings is the indispensable probe for BSM physics.

2) The ILC250 as a Higgs factory. The ILC250 has capability to achieve
O(1%) Higgs precision model-independently, and perform exploration and verification of BSMs.

3) The ILC for future collider experiments.

The ILC is upgradable to new physics scale suggested by its own physics outcomes.

Main References

- arXiv:1710.07621v3 [hep-ex] 23 Jan 2018
 Physics Case for the 250 GeV Stage of the International Linear Collider Keisuke Fujii et al.
- 2). Model-independent determination of the triple Higgs coupling at e⁺e⁻ colliders Tim Barklow, Keisuke Fujii, Sunghoon Jung, Michael E. Peskin, and Junping Tian PHYSICAL REVIEW D 97, 053004 (2018)
- 3). Improved formalism for precision Higgs coupling fits Tim Barklow et al, PHYSICAL REVIEW D 97, 053003 (2018)
- 4). Japan-MEXT Commission report on scientific significance of ILC250 etc. Shoji Asai et al.
 <u>http://www.mext.go.jp/b_menu/shingi/chousa/shinkou/038/index.htm</u>
- 5). Physics and Status of ILC250 Junping Tian Joint Kavli IPMU - ICEPP Workshop on New Directions for LHC: Run 2 and Beyond <u>https://indico.cern.ch/event/714089/</u>

Back up



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Abstract

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Contribution ID : 755

Type : Parallel

Precision Higgs Measurements at the 250 GeV ILC

Saturday, 7 July 2018 17:00 (15)

The plan for the International Linear Collider is now being prepared as a staged design, with the first stage at 250 GeV and later stages achieving the full project specifications with 4 ab-1 at 500 GeV. This talk will present the capabilities for precision Higgs boson measurements at 250 GeV and their relation to the full ILC program. It will show that the 250 GeV stage of ILC will already provide many compelling results in Higgs physics, with new measurements not available at LHC, model-independent determinations of key parameters, and tests for and possible discrimination of a variety of scenarios for new physics.

Higgs Properties at the ILC250

O Dark matter: "Hidden sector with Higgs portal model"

Higgs full width Γh and BR_invisible

O Origin of asymmetry on matter and antimatter:

"Baryongenesis & Leptogenesis"

Higgs CP-admixture in an extended Higgs sector

Precision on WIMP in a Hidden-sector

O DM theories constructed with a Hidden sector (particles are SM singlets).

interactions through gravity?

 $\begin{array}{l} & HL\psi \\ \rightarrow \text{Several models as a portal} \\ & \text{Neutrino Portal} \underset{F}{H} \underset{\mu\nu}{F} \underset{\mu\nu}{F} \underset{\mu\nu}{F} \underset{\mu\nu}{F} \underset{\nu}{F} \underset{\nu$

scalar, Majonara, vector WIMP.

Higgs invisible decay BR with 3000 fb⁻¹ATLAS: $BR_{inv} < 10\% \sim 14\%$ CMS: $BR_{inv} < 7\% \sim 11\%$

ATLAS & CMS, Eric Feng (Argonne) @ BSM Higgs Workshop Fermilab Nov 3, 2014



Direct exploration of singlet dark matter is possible up to 2/Mh GeV .

Precision on WIMP in a Hidden-sector



Interactions through gravity ?

 $HL\psi$ $\rightarrow Several models as a portal$ Neutrino Portal $H_{\mu\nu}\tilde{F}^{\mu\nu}$: sterile neutrino $\int f_{S} F_{\mu\nu}\tilde{F}^{\mu\nu} f_{F}^{\mu\nu}$ Higgs Portal $|H|^{2}S^{2}$: Higgs invisible decay $\int f_{F} F_{\mu\nu}\tilde{F}^{\mu\nu}S$ S: a dark matter fields $|H|^{2}\phi^{2}, |H|^{2}S^{2}\chi\chi, |H|^{2}V_{\mu}V^{\mu}$

scalar, Majonara, vector WIMP.

 ILC250 2000 fb-1

 Γ_h 2.5

 $BR(h \rightarrow inv)$ 0.32 [%]

 $BR(h \rightarrow other)$ 1.6 [%]

arXiv:1710.07621 Keisuke Fujii et al.

→ ~ two orders of magnitude higher !



Direct exploration of singlet dark matter is possible up to 2/Mh GeV . with high sensitivity !

Higgs CP-admixture in hττ

Eibun Senaha *et al,* Phys. Lett. B 762, 315 (2016) <u>modified parameters</u>



Phys. Rev. D 13 June (2018)

Higgs CP-admixture in hVV

arXiv:1710.07621 Keisuke Fujii et al.

O CP-odd Higgs induced by **Diemnntion-6 operators**

$$\Delta \mathcal{L}_{hZZ} = \frac{1}{2} \frac{\tilde{b}}{v} h Z_{\mu\nu} \tilde{Z}^{\mu\nu}$$

$$ILC250\ 2000\ fb^{-1}$$

$$hZZ : bt_z \sim 0.5\%$$
(hWW : bt_w ~ 11\%)

O LHC ATLAS : EFT analysis

$$\mathcal{L}_{0}^{V} = \left\{ \kappa_{\mathrm{SM}} \left[\frac{1}{2} g_{HZZ} Z_{\mu} Z^{\mu} + g_{HWW} W_{\mu}^{+} W^{-\mu} \right] - \frac{1}{4} \left[\kappa_{Hgg} g_{Hgg} G_{\mu\nu}^{a} G^{a,\mu\nu} + \tan \alpha \kappa_{Agg} g_{Agg} G_{\mu\nu}^{a} \tilde{G}^{a,\mu\nu} \right] - \frac{1}{4} \frac{1}{\Lambda} \left[\kappa_{HZZ} Z_{\mu\nu} Z^{\mu\nu} + \tan \alpha \kappa_{AZZ} Z_{\mu\nu} \tilde{Z}^{\mu\nu} \right] - \frac{1}{2} \frac{1}{\Lambda} \left[\kappa_{HWW} W_{\mu\nu}^{+} W^{-\mu\nu} + \tan \alpha \kappa_{AWW} W_{\mu\nu}^{+} \tilde{W}^{-\mu\nu} \right] \right\} X_{0}$$

Expected and observed confidence intervals at 95% CL with 36.1 fb⁻¹ of data at $\sqrt{s} = 13$ TeV.

BSM coupli	ng Fit	Expected	Observed
K _{BSM}	configuration	conf. inter.	conf. inter.
KAgg	$(\kappa_{Hgg} = 1, \kappa_{SM} = 1)$	[-0.47, 0.47]	[-0.68, 0.68]
K _{HVV}	$(\kappa_{Hgg} = 1, \kappa_{SM} = 1)$	[-2.9, 3.2]	[0.8, 4.5]
κ_{HVV}	$(\kappa_{Hgg} = 1, \kappa_{SM} \text{ free})$	[-3.1, 4.0]	[-0.6, 4.2]
KAVV	$(\kappa_{Hgg} = 1, \kappa_{SM} = 1)$	[-3.5, 3.5]	[-5.2, 5.2]
K _{AVV}	$(\kappa_{Hgg} = 1, \kappa_{SM} \text{ free})$	[-4.0, 4.0]	[-4.4, 4.4]



overall HVV precision (no individually) :

translate to the ILC condition 36fb-1 bt_v [-200%, 200%] HL-LHC bt_v [-20%, 20%]

Synergy with the HL-LHC + other experiments



- "SuperKEK-B" will give some conclusion on flavor physics,
- "LISA" project will give observation in later 2030s .

Higgs sector will be elucidated, and a direction of new physics be uncovered

M CRSS



ILC operation plan



Figure 2: Run plan for the staged ILC starting with a 250-GeV machine under two different assumptions on the achievable instantaneous luminosity at 250 GeV. Both cases reach the same final integrated luminosities as in Fig. 1.

arXiv:1506.05992v2 [hep-ex] 26 Jun 2015 Physics Case for the International Linear Collide

arXiv:1710.07621v2 [hep-ex] 3 Nov 2017

Physics Case for the 250 GeV Stage of



7 observables

TABLE I. Values and uncertainties for precision electroweak observables used in this paper. The values are taken from Ref. [42], except for the averaged value of A_{ℓ} , which corresponds to the averaged value of $\sin^2 \theta_{\text{eff}}$ in Ref. [43]. The best -fit values are those of the fit in Ref. [42]. For the purpose of fitting Higgs boson couplings as described in Sec. VII, we use improvements in some of the errors expected from the LHC [44] and ILC [45]. The improved estimate of the W width is obtained from $\Gamma_W = \Gamma(W \to \ell \nu)/\text{BR}(W \to \ell \nu)$.

Observable	Current value	Current σ	Future σ	SM best-fit value	
$\overline{\alpha^{-1}(m_z^2)}$	128.9220	0.0178		(same)	
$G_F (10^{-10} \text{ GeV}^{-2})$	1166378.7	0.6		(same)	
m_W (MeV)	80 385	15	5	80 361	
m_Z (MeV)	91 187.6	2.1		91 188.0	
m_h (MeV)	125 090	240	15	125 110	
A_{ℓ}	0.14696	0.0013		0.147937	
Γ_{ℓ} (MeV)	83.984	0.086		83.995	
Γ_Z (MeV)	2495.2	2.3		2494.3	
Γ_W (MeV)	2085	42	2	2088.8	

EWPOs and Renormalization EFT

PHYS. REV. D 97, 053004 (2018)

 $\alpha(m_Z), G_F, m_W, m_Z, m_h, A_{LR}(\ell), \Gamma(Z \to \ell^+ \ell^-)$

7 observables

 $\alpha(m_Z), G_F, m_W, m_Z, m_h, A_{LR}(\ell), \Gamma(Z \to \ell^+ \ell^-)$

$$\begin{split} \delta e &= \delta (4\pi\alpha(m_Z^2))^{1/2} = s_w^2 \delta g + c_w^2 \delta g' + \frac{1}{2} \delta Z_A \\ \delta G_F &= -2\delta v + 2c'_{HL} \\ \delta m_W &= \delta g + \delta v + \frac{1}{2} \delta Z_W \\ \delta m_Z &= c_w^2 \delta g + s_w^2 \delta g' + \delta v - \frac{1}{2} c_T + \frac{1}{2} \delta Z_Z \\ \delta m_h &= \frac{1}{2} \delta \overline{\lambda} + \delta v + \frac{1}{2} \delta Z_h \\ & \longrightarrow \\ \delta g, \ \delta g', \ \delta v, \ \delta \lambda, \ C_T \end{split} \qquad \begin{aligned} \delta \Gamma_\ell &= \delta m_Z + 2 \frac{g_L^2 \delta g_L + g_R^2 \delta g_R}{g_L^2 + g_R^2} \\ \delta A_\ell &= \frac{4g_L^2 g_R^2 (\delta g_L - \delta g_R)}{g_L^4 - g_R^4} \\ \delta A_\ell &= \frac{4g_L^2 g_R^2 (\delta g_L - \delta g_R)}{g_L^4 - g_R^4} \\ g_L &= \frac{g_L}{c_w} \Big[(-\frac{1}{2} + s_w^2)(1 + \frac{1}{2} \delta Z_Z) - \frac{1}{2} (c_{HL} + c'_{HL}) - s_w c_w \delta Z_{AZ} \Big] \\ g_R &= \frac{g}{c_w} \Big[(+s_w^2)(1 + \frac{1}{2} \delta Z_Z) - \frac{1}{2} c_{HE} - s_w c_w \delta Z_{AZ} \Big] \\ CHL + C'HL, \ CHE \end{aligned}$$

TGC : 3 parameters

$$\Delta \mathcal{L}_{TGC} = ig_V \Big\{ V^{\mu}(\hat{W}_{\mu\nu}^{-}W^{+\nu} - \hat{W}_{\mu\nu}^{+}W^{-\nu}) + \kappa_V W^{+}_{\mu} W^{-}_{\nu} \hat{V}^{\mu\nu} + \frac{\lambda_V}{m_W^2} \hat{W}_{\mu\nu}^{-\rho} \hat{W}_{\rho\nu}^{+} \hat{V}^{\mu\nu} \Big\}$$

$$\overset{\text{wf}}{\stackrel{\text{s}}{}_{2}} = \overset{\text{wf}}{\stackrel{\text{s}}{}_{2}} + \overset{\text{wf}}{\stackrel{\text{s}}{}_{2}} \overset{\text{wf}}{\stackrel{\text{s}}{}_{2}} \Big\}_{2}$$

$$g_Z = gc_w (1 + \frac{1}{2} \delta Z_Z + \frac{s_w}{c_w} \delta Z_{AZ})$$

$$\kappa_A = 1 + (8c_{WB})$$

$$\lambda_A = -6g^2 c_{3W}$$

$$\overset{\text{wf}}{\stackrel{\text{s}}{}_{2}} + \overset{\text{wf}}{\stackrel{\text{s}}{}_{2}} \Big\}_{2} + \overset{\text{wf}}{\stackrel{\text{s}}{}_{2}} \Big\}_{2}$$

$$\begin{split} \delta g_{Z,eff} &= \delta g_Z + \frac{1}{c_w^2} ((c_w^2 - s_w^2) \delta g_L + s_w^2 \delta g_R - 2 \delta g_W) \\ \delta \kappa_{A,eff} &= (c_w^2 - s_w^2) (\delta g_L - \delta g_R) + 2 (\delta e - \delta g_W) + (8 c_{WB}) \\ \delta \lambda_{A,eff} &= -6 g^2 c_{3W} \end{split}$$

$$g_W = g \left(1 + c'_{HL} + \frac{1}{2}\delta Z_W\right)$$
 29

Synergy with HL-LHC

EFT input from

HL-LHC:
$$\frac{\text{BR}(h \to \gamma \gamma)}{\text{BR}(h \to ZZ^*)} \qquad \frac{\text{BR}(h \to \gamma Z)}{\text{BR}(h \to ZZ^*)}$$

(synergy with HL-LHC)

 $\delta\Gamma(h \to \gamma\gamma) = 528\,\delta Z_A - c_H + 4\delta e + 4.2\,\delta m_h - 1.3\,\delta m_W - 2\delta v$

$$\delta\Gamma(h \to Z\gamma) = 290 \,\delta Z_{AZ} - c_H - 2(1 - 3s_W^2)\delta g + 6c_w^2 \delta g' + \delta Z_A + \delta Z_Z + 9.6 \,\delta m_h - 6.5 \,\delta m_Z - 2\delta v$$

 $\delta\Gamma(h \to ZZ^*) = 2\eta_Z - 2\delta v - 13.8\delta m_Z + 15.6\delta m_h - 0.50\delta Z_Z - 1.02C_Z + 1.18\delta\Gamma_Z$

$$\delta Z_A = s_w^2 \left((8c_{WW}) - 2(8c_{WB}) + (8c_{BB}) \right) \qquad \delta Z_{AZ} = s_w c_w \left((8c_{WW}) - (1 - \frac{s_w^2}{c_w^2})(8c_{WB}) - \frac{s_w^2}{c_w^2}(8c_{BB}) \right)$$

60

Discrimination of BSM models

	Model	$b\overline{b}$	$c\overline{c}$	gg	WW	au au	ZZ	$\gamma\gamma$	$\mu\mu$
1	MSSM	+4.8	-0.8	- 0.8	-0.2	+0.4	-0.5	+0.1	+0.3
2	Type II 2HD	+10.1	-0.2	-0.2	0.0	+9.8	0.0	+0.1	+9.8
3	Type X 2HD	-0.2	-0.2	-0.2	0.0	+7.8	0.0	0.0	+7.8
4	Type Y 2HD	+10.1	-0.2	-0.2	0.0	-0.2	0.0	0.1	-0.2
5	Composite Higgs	-6.4	-6.4	-6.4	-2.1	-6.4	-2.1	-2.1	-6.4
6	Little Higgs w. T-parity	0.0	0.0	-6.1	-2.5	0.0	-2.5	-1.5	0.0
7	Little Higgs w. T-parity	-7.8	-4.6	-3.5	-1.5	-7.8	-1.5	-1.0	-7.8
8	Higgs-Radion	-1.5	- 1.5	+10.	-1.5	-1.5	-1.5	-1.0	-1.5
9	Higgs Singlet	-3.5	-3.5	-3.5	-3.5	-3.5	-3.5	-3.5	-3.5

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Discrimination of BSM models

- a PMSSM model with b squarks at 3.4 TeV, gluino at 4 TeV
- a Type II 2 Higgs doublet model with $m_A = 600 \text{ GeV}, \tan \beta = 7$
 - a Type X 2 Higgs doublet model with $m_A = 450 \text{ GeV}, \tan \beta = 6$
 - a Type Y 2 Higgs doublet model with $m_A = 600 \text{ GeV}, \tan \beta = 7$
 - a composite Higgs model MCHM5 with $f = 1.2 \text{ TeV}, m_T = 1.7 \text{ TeV}$
 - a Little Higgs model with T-parity with $f = 785 \text{ GeV}, m_T = 2 \text{ TeV}$
 - \bullet A Little Higgs model with couplings to 1st and 2nd generation with $f=1.2~{\rm TeV}, m_T=1.7~{\rm TeV}$
 - A Higgs-radion mixing model with $m_r = 500 \text{ GeV}$
- a model with a Higgs singlet at 2.8 TeVcreating a Higgs portal to dark matter and large λ for electroweak baryogenesis

Discrimination of BSM models

○ chi2 discrimination test

$$(\chi^2)_{AB} = (g_A^T - g_B^T) [VCV^T]^{-1} (g_A - g_B)$$

Unit n [
$$\sigma$$
] $n = \sqrt{\chi^2}$

- gA, gB: vector of couplings in Model A, B
- C: covariance matrix of EFT coefficients
- Vij: linear dependence of coupling g_i on EFT coefficient c_j



Distinguishable with more than 3σ for Most of the BSM models !

Anomalous ZZH

$$\mathcal{L}_{ZZH} = M_Z^2 \left(\frac{1}{v} + \frac{a_Z}{\Lambda}\right) Z_\mu Z^\mu H + \frac{b_Z}{2\Lambda} \hat{Z}_{\mu\nu} \hat{Z}^{\mu\nu} H + \frac{\tilde{b}_Z}{2\Lambda} \hat{Z}_{\mu\nu} \hat{Z}^{\mu\nu} H + \frac{\tilde{b}_Z}{2\Lambda} \hat{Z}_{\mu\nu} \hat{Z}^{\mu\nu} H$$
(A=1TeV)



Higgs couplings precision

arXiv:1506.05992v2 [hep-ex] 26 Jun 2015 Physics Case for the International Linear Collider

Topic	Parameter	Initial Phase	Full Data Set	units	ref.
Higgs	m_h	25	15	MeV	[15]
	g(hZZ)	0.58	0.31	%	[2]
	g(hWW)	0.81	0.42	%	[2]
	$g(hb\overline{b})$	1.5	0.7	%	[2]
	g(hgg)	2.3	1.0	%	[2]
	$g(h\gamma\gamma)$	7.8	3.4	%	[2]
		1.2	1.0	%, w. LHC results	[17]
	$g(h\tau\tau)$	1.9	0.9	%	[2]
	$g(hc\overline{c})$	2.7	1.2	%	[2]
	$g(ht\bar{t})$	18	6.3	%, direct	[2]
		20	20	$\%, t\bar{t}$ threshold	[34]
	$g(h\mu\mu)$	20	9.2	%	[2]
	g(hhh)	77	27	%	[2]
	Γ_{tot}	3.8	1.8	%	[2]
	Γ_{invis}	0.54	0.29	%,95% conf. limit	[2]
Тор	m_t	50	50	MeV $(m_t(1S))$	[33]
	Γ_t	60	60	MeV	[34]
	g_L^γ	0.8	0.6	%	[42]
	g_R^γ	0.8	0.6	%	[42]
	g_L^Z	1.0	0.6	%	[42]
	g_R^Z	2.5	1.0	%	[42]
	F_2^{γ}	0.001	0.001	absolute	[42]
	F_2^Z	0.002	0.002	absolute	[42]
W	m_W	2.8	2.4	MeV	[62]
	g_1^Z	$8.5 imes 10^{-4}$	6×10^{-4}	absolute	[63]
	κ_γ	$9.2 imes 10^{-4}$	7×10^{-4}	absolute	[63]
	λ_γ	7×10^{-4}	2.5×10^{-4}	absolute	[63]
Dark Matter	EFT Λ : D5	2.3	3.0	TeV, 90% conf. limit	[61]
	EFT Λ : D8	2.2	2.8	TeV, 90% conf. limit	[61]

Table 1: Projected accuracies of measurements of Standard Model parameters at the two stages of the ILC program proposed in the report of the ILC Parameters Joint Working Group [7]. This program has an initial phase with 500 fb⁻¹ at 500 GeV, 200 fb⁻¹ at 350 GeV, and 500 fb⁻¹ at 250 GeV, and a luminosity-upgraded phase with an additional 3500 fb⁻¹ at 500 GeV and 1500 fb⁻¹ at 250 GeV. Initial state polarizations are taken according to the prescriptions of [7]. Uncertainties are listed as 1σ errors (except where indicated), computed cumulatively at each stage of the program. These estimated errors include both statistical uncertainties and theoretical and experimental systematic uncertainties. Except where indicated, errors in percent (%) are fractional uncertainties relative to the Standard Model values. More specific information for the sets of measurements is given in the text. For each measurement, a reference describing the technique is given.

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Physics Case for the 250 GeV Stage of the International Linear Collider

	ILC250		+ILC500	
	κ fit	EFT fit	κ fit	EFT fit
g(hbb)	1.8	1.1	0.60	0.58
g(hcc)	2.4	1.9	1.2	1.2
g(hgg)	2.2	1.7	0.97	0.95
g(hWW)	1.8	0.67	0.40	0.34
g(h au au)	1.9	1.2	0.80	0.74
g(hZZ)	0.38	0.68	0.30	0.35
$g(h\gamma\gamma)$	1.1	1.2	1.0	1.0
$g(h\mu\mu)$	5.6	5.6	5.1	5.1
$g(h\gamma Z)$	16	6.6	16	2.6
g(hbb)/g(hWW)	0.88	0.86	0.47	0.46
$g(h\tau\tau)/g(hWW)$	1.0	1.0	0.65	0.65
g(hWW)/g(hZZ)	1.7	0.07	0.26	0.05
Γ_h	3.9	2.5	1.7	1.6
$BR(h \to inv)$	0.32	0.32	0.29	0.29
$BR(h \rightarrow other)$	1.6	1.6	1.3	1.2

Table 1: Projected relative errors for Higgs boson couplings and other Higgs observables, in %, for fits in the κ and EFT formalisms. The ILC250 columns assume a total integrated luminosity of 2 ab⁻¹ at $\sqrt{s} = 250$ GeV, shared by (-+, +-, --, ++) = (45%, 45%, 5%, 5%)as described in Section 2. The ILC500 columns assume, in addition, a total integrated luminosity of 200 fb⁻¹ at $\sqrt{s} = 350$ GeV, shared as (45%, 45%, 5%, 5%), and a total integrated luminosity of 4 ab⁻¹ at $\sqrt{s} = 500$ GeV, shared as (40%, 40%, 10%, 10%). Three observables at the HL-LHC, $BR_{\gamma\gamma}/BR_{ZZ}$, $BR_{\gamma Z}/BR_{\gamma\gamma}$ and $BR_{\mu\mu}/BR_{\gamma\gamma}$, are included in all of the fits. The effective couplings g(hWW) and g(hZZ) are defined as proportional to the square root of the corresponding partial widths. The last two lines give 95% confidence upper limits on the exotic branching ratios. The detailed formulae used in the EFT fit, and the resulting covariance matrix, can be found in [15].

Higgs direct couplings to cc and gg at ILC250



e+e- -> ZH -> ff(jj): b-likeness .vs. c-likeness

Oclean environment at e+e-; excellent b- and c-tagging performance
 Obb/cc/gg modes can be separated simultaneously by template fitting

Ono, et. al, Euro. Phys. J. C73, 2343; F.Mueller, PhD thesis (DESY)

Higgs —> invisible at ILC250

$$e^+ + e^- \rightarrow ZH \rightarrow l^+l^-/q\bar{q} + \text{Missing}$$



recoil technique: Higgs mass fully reconstructed even it decays invisibly
 right-handed beam polarization helps: much lower background
 OBR(H—>inv.) < 0.3% (CL95%)

ILC site in North Japan

