

Decay of False Vacua via Impurities in String Theories

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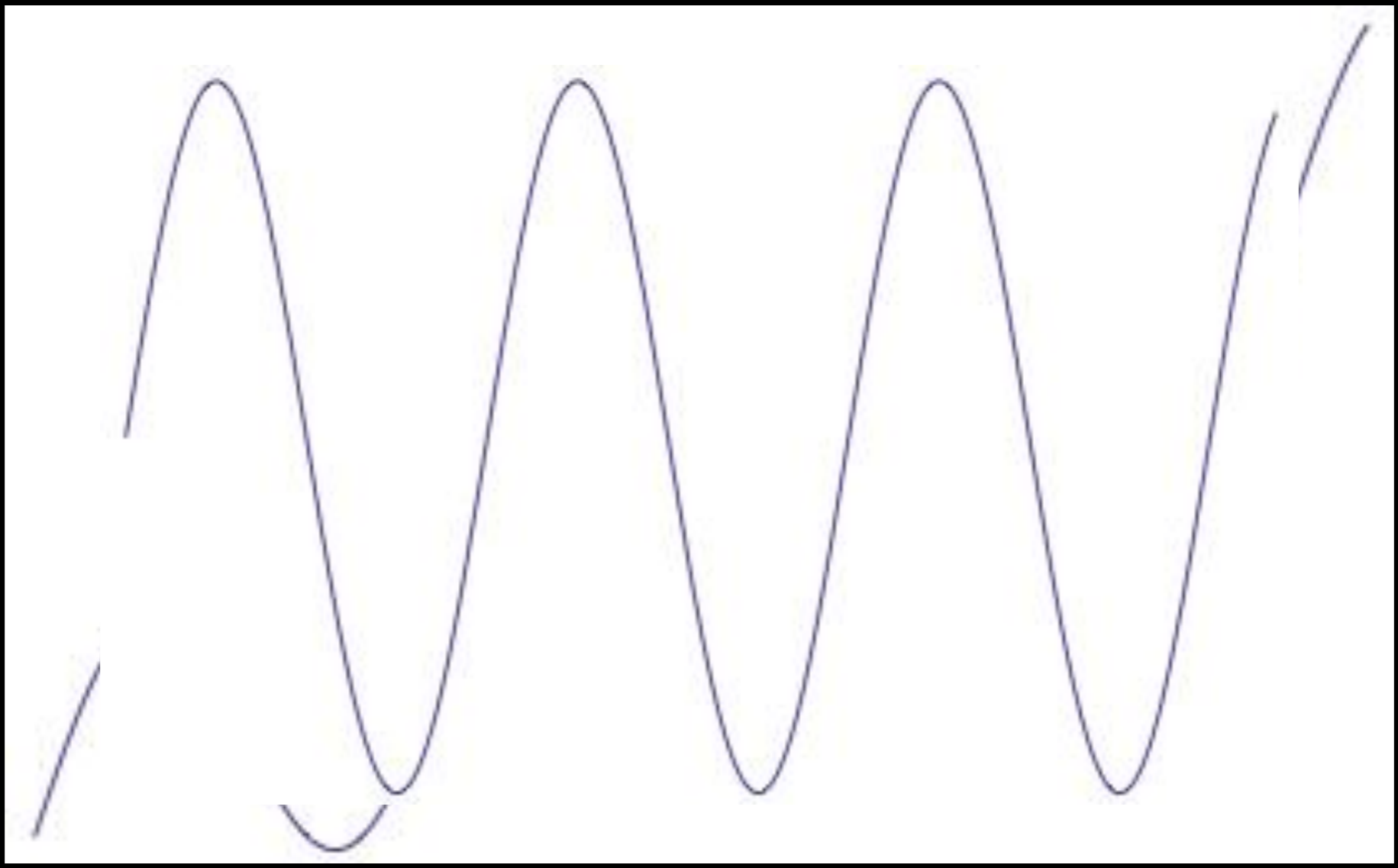
Y. Nakai, YO, N. Tanahashi '18

T. Hiramatsu, M. Eto, K. Kamada, T. Kobayashi, YO, '14

M. Eto, Y. Hamada, K. Kamada, T. Kobayashi, K. Ohashi, YO, '13

SUSY GUT

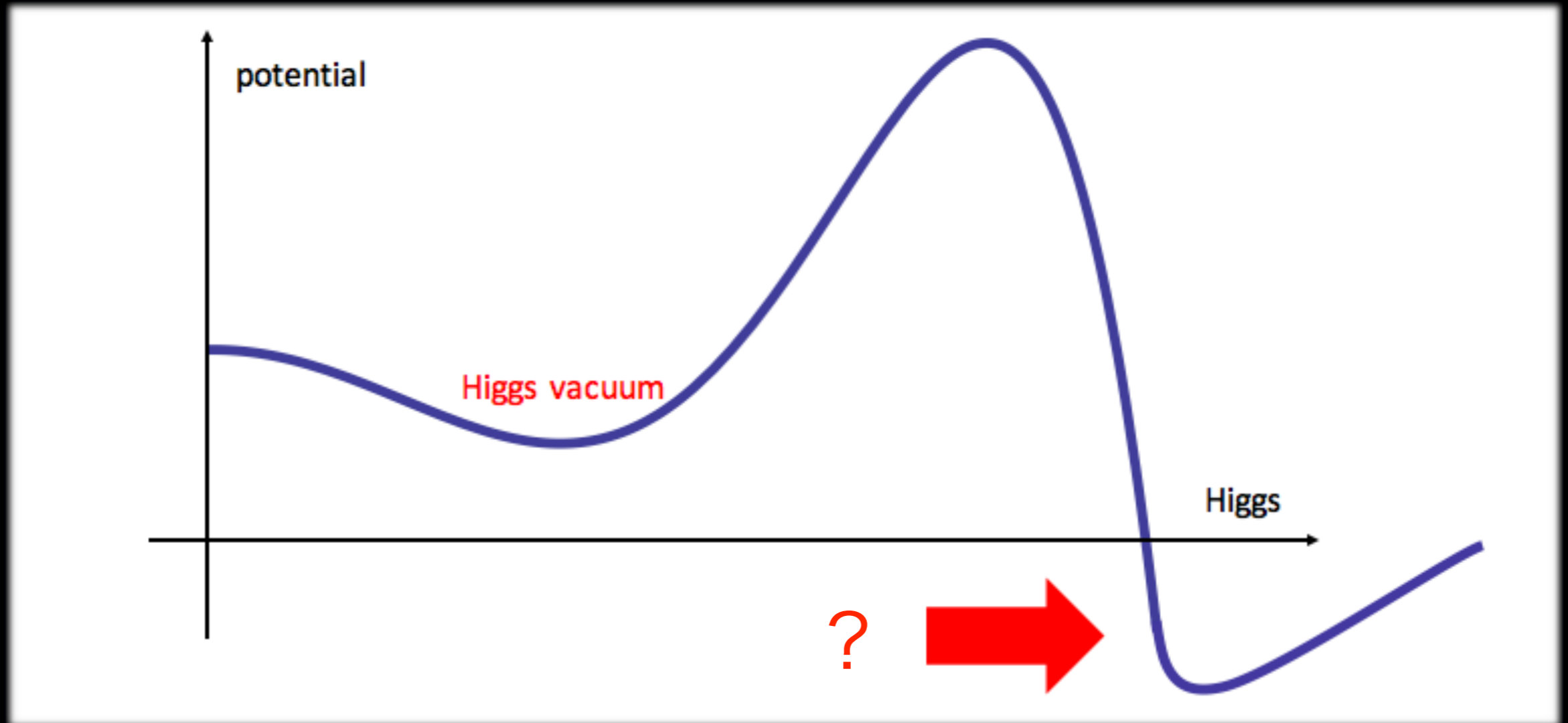
V



$SU(3) \times SU(2) \times U(1)$

$SU(4) \times U(1)$

Higgs Vacuum



Field theory analysis

Original Idea:

P. J. Steinhardt, '81 and Y. Hosotani '83

Simple Toy model

Eto, Hamada, Kamada, Kobayashi, Ohashi, YO, '13

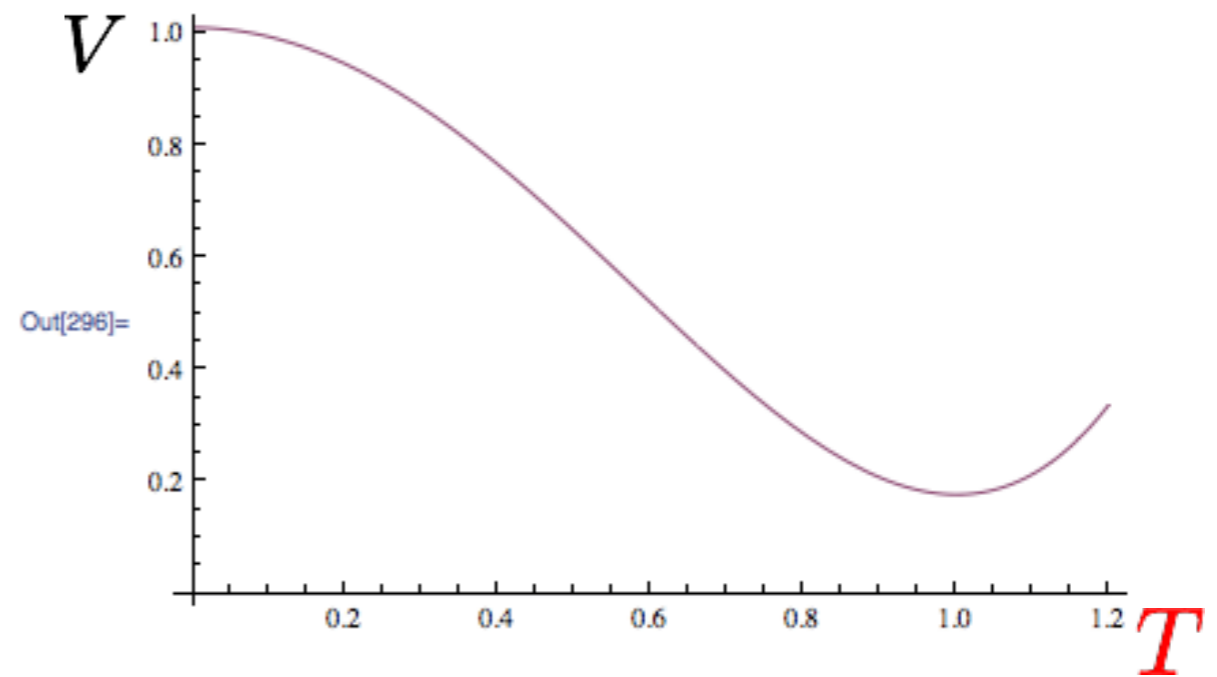
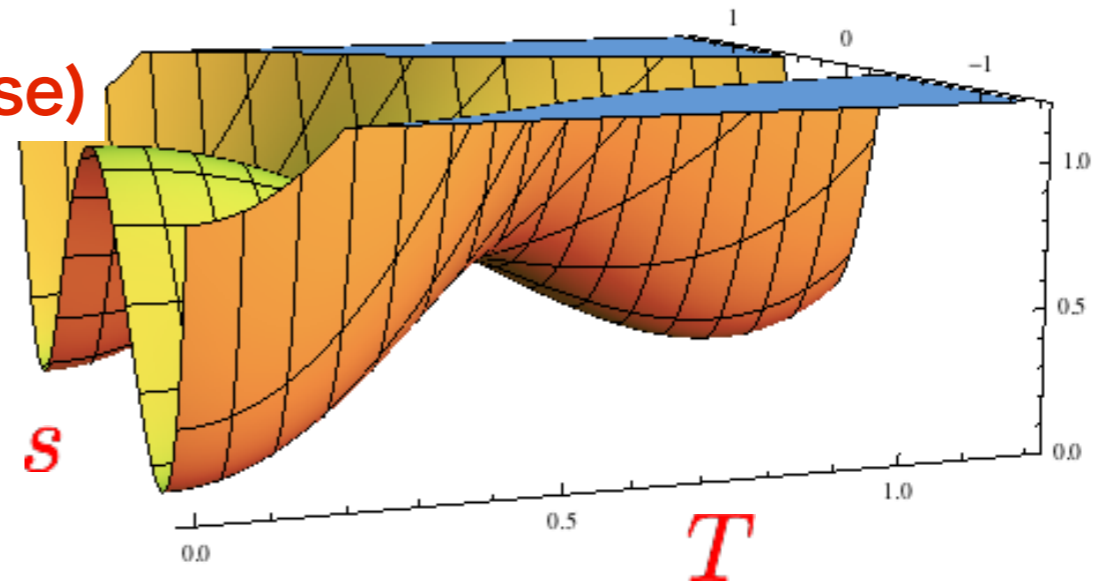
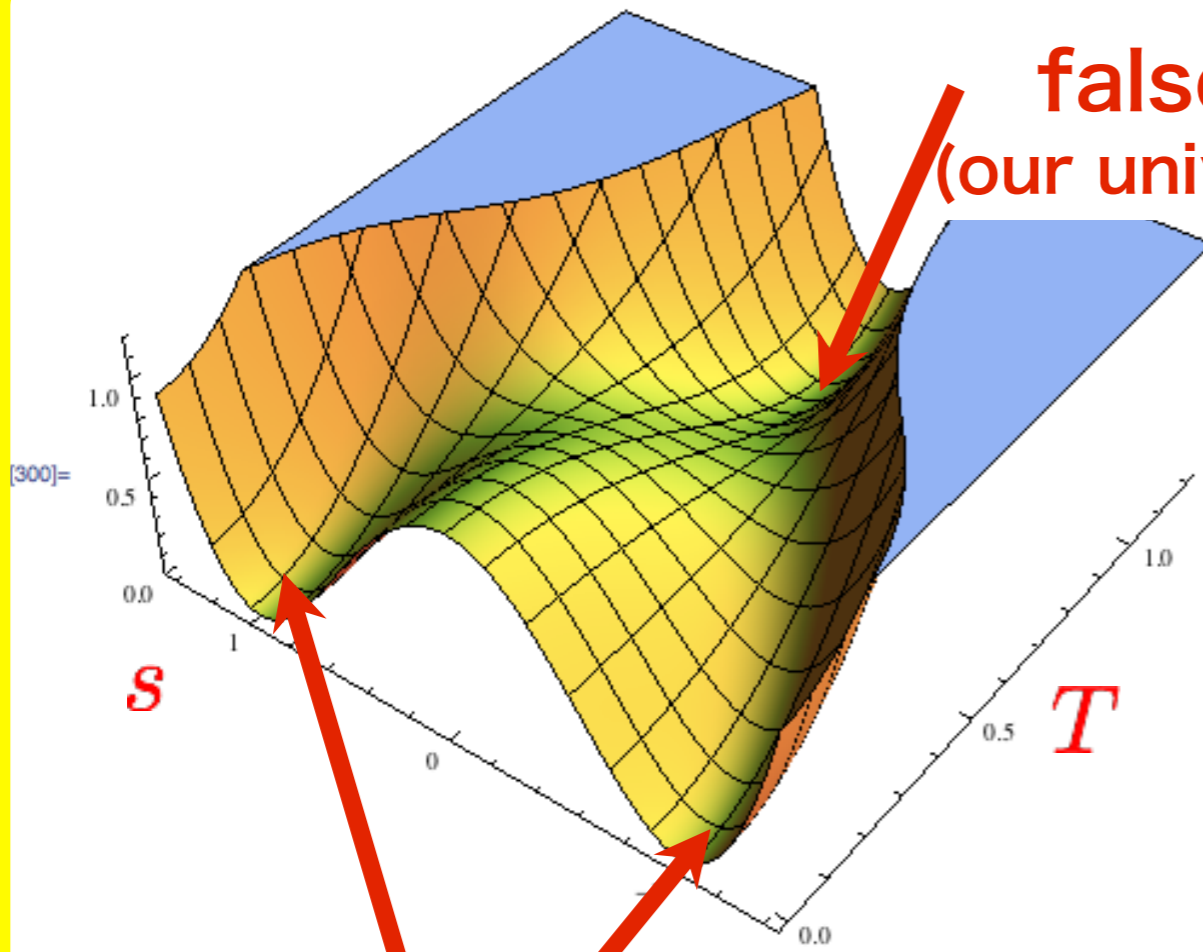
$$\mathcal{L} = \mu^4 \left[\frac{1}{\mathcal{V}(T)} |\tilde{\partial}_\mu T|^2 + \epsilon^2 |\tilde{\partial}_\mu s|^2 - \mathcal{V}(T) |s^2 - 1|^2 - \frac{4}{\epsilon^2} |T|^2 |s|^2 \right].$$

mass dim

$$\mathcal{V}(T) \equiv 1 - \frac{|T|^2}{2\lambda} + \frac{|T|^4}{4\lambda}.$$

- 2 dim-less fields T, s
- 2 dim-less parameters λ, ϵ

Landscape of vacua



true

$$\epsilon = 1, \lambda = 3/10$$

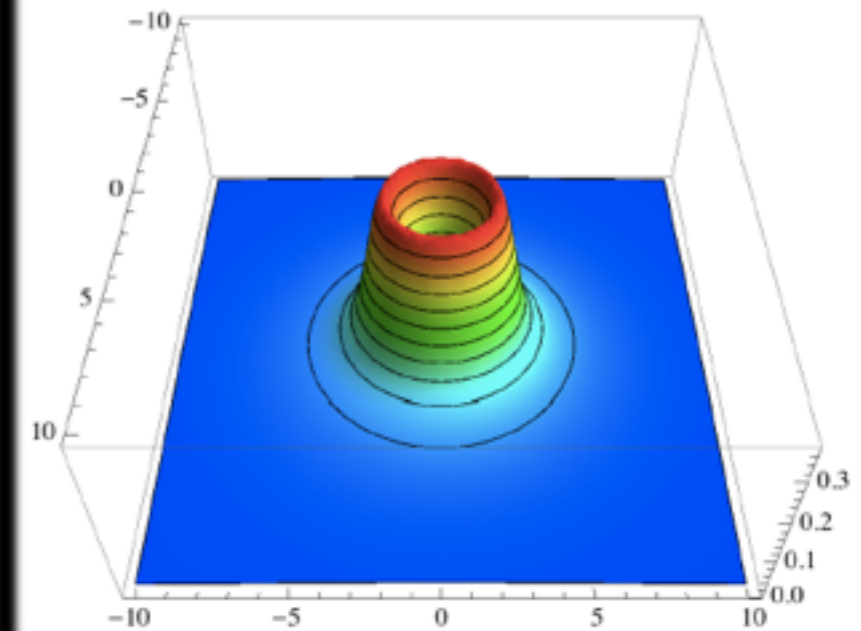
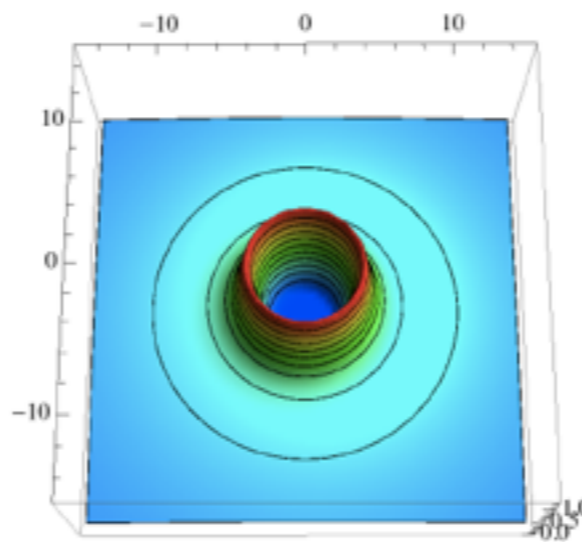
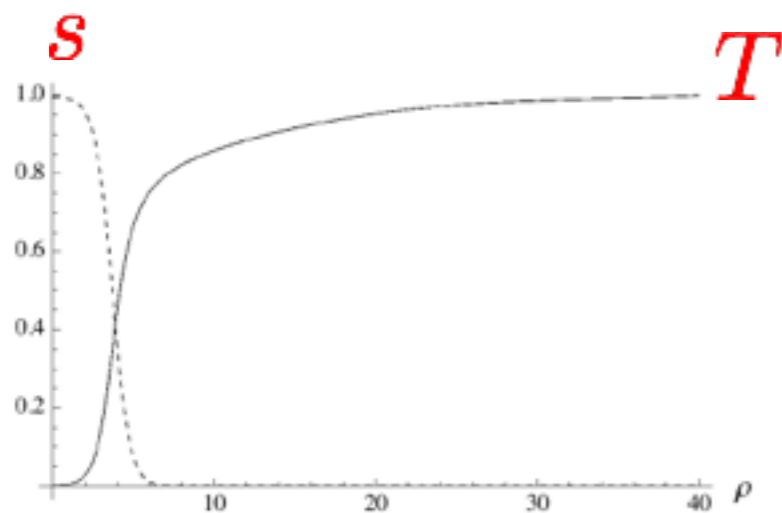
Cosmic R-string (R-tube) solution

$$T = f(\rho)e^{in\theta}, \quad s = \gamma h(\rho),$$

winding number

two vacua

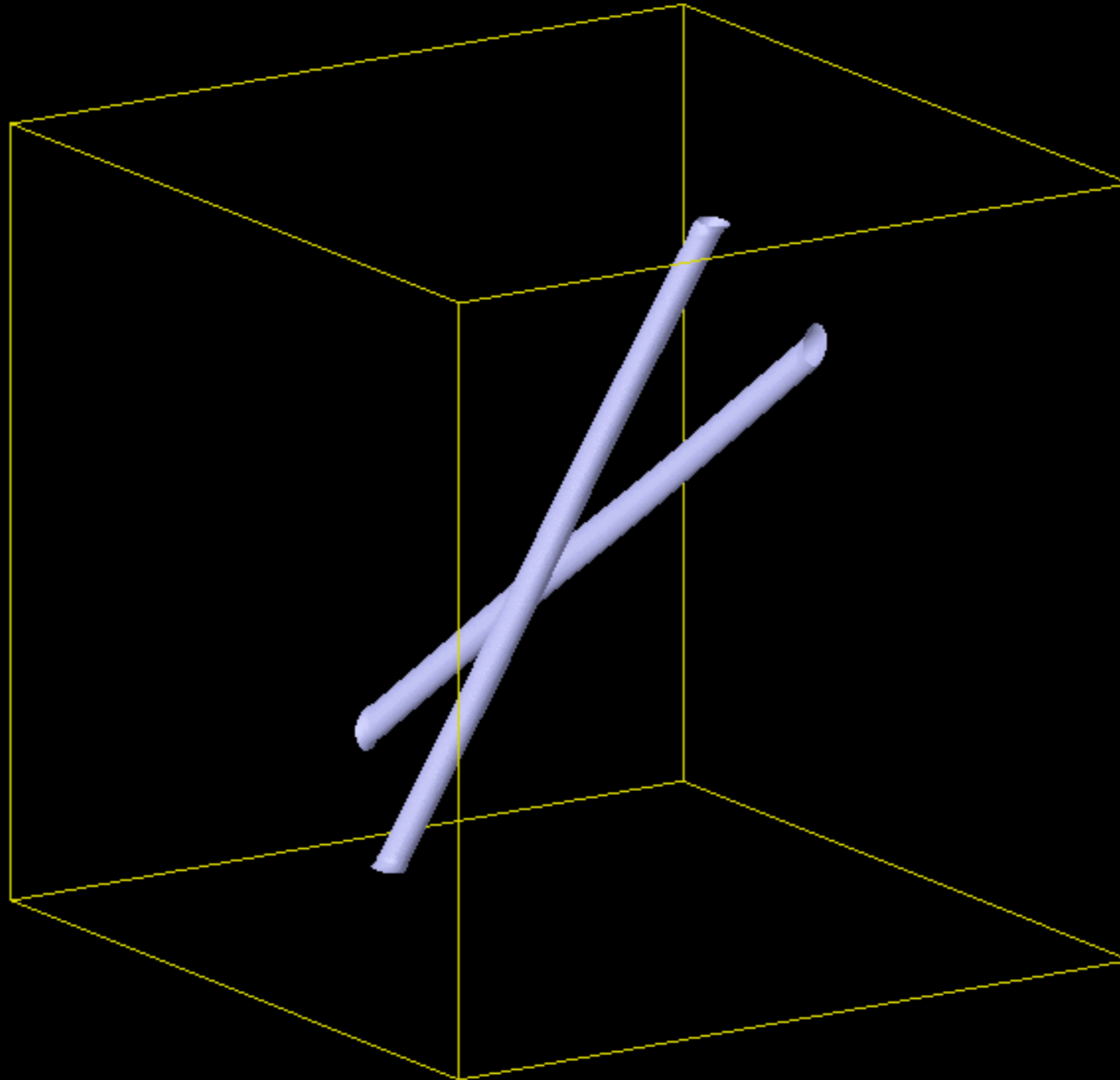
$$\epsilon = 1, \quad \lambda = 27/100$$



String collision Numerical simulation

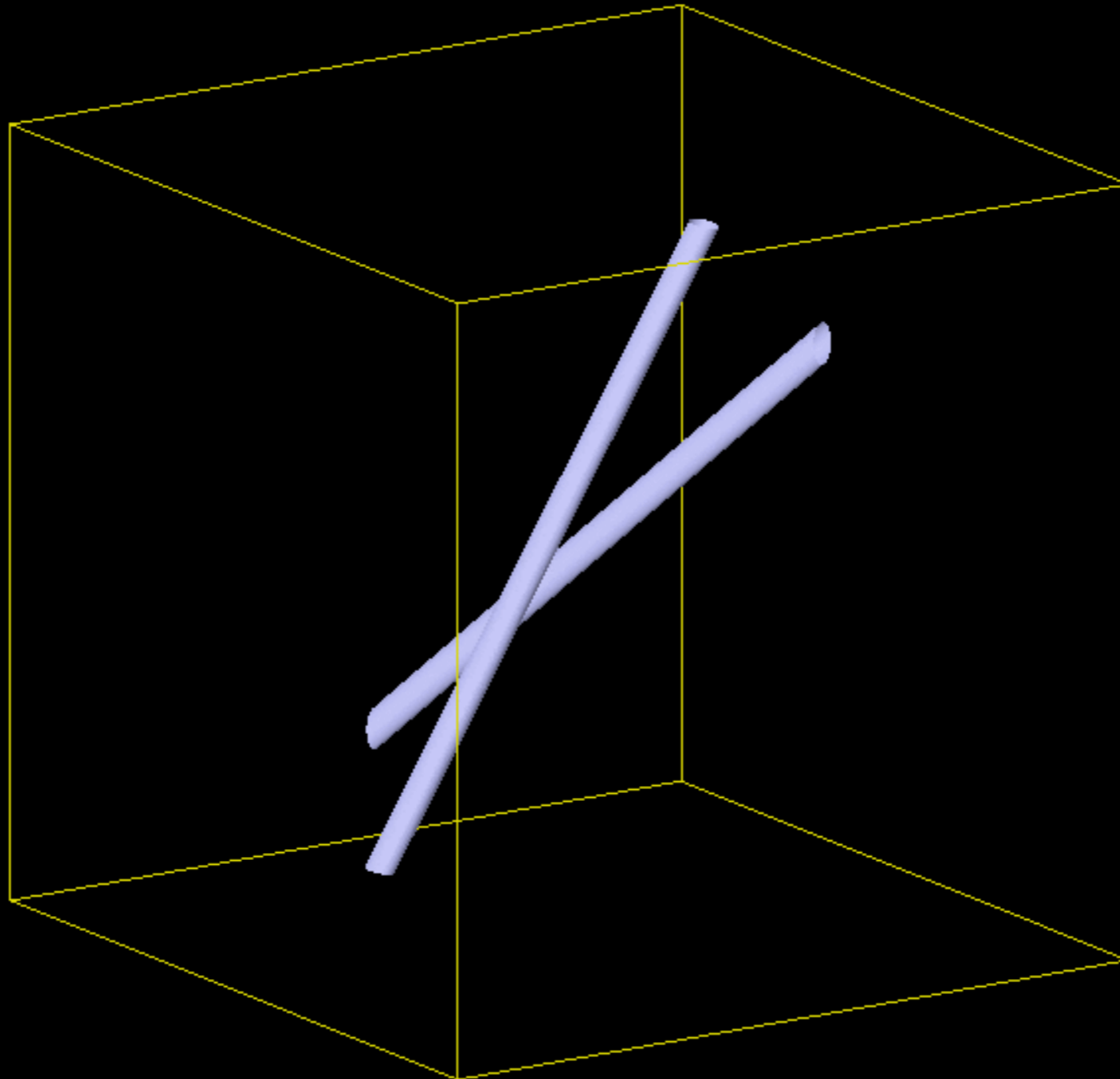
String($n=1$)/String($n=1$) collision 60% speed of light

Hiramatsu, Eto, Kamada, Kobayashi, YO, '14

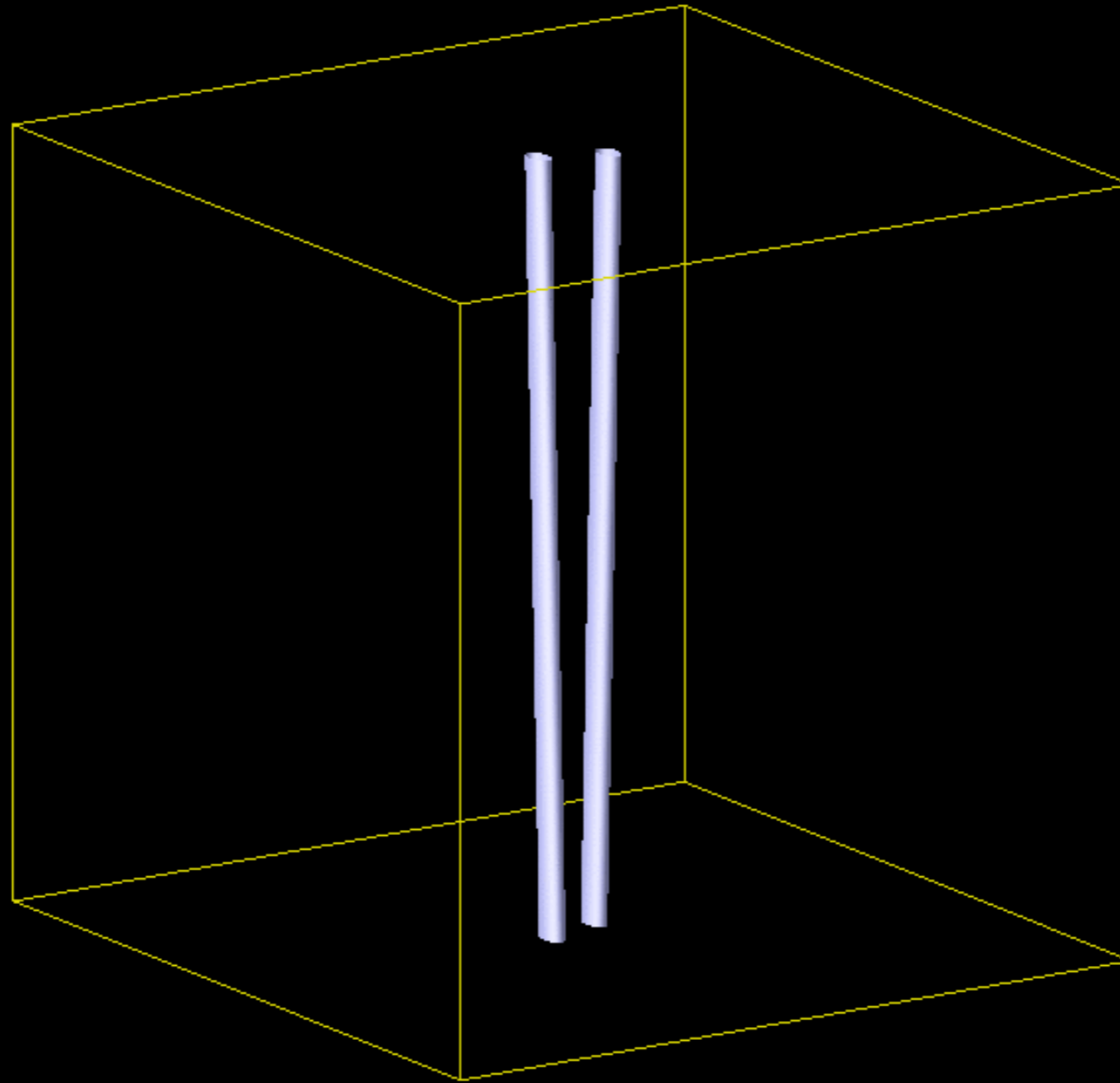


String($n=1$)/String($n=1$) collision
80% speed of light, **large angle**

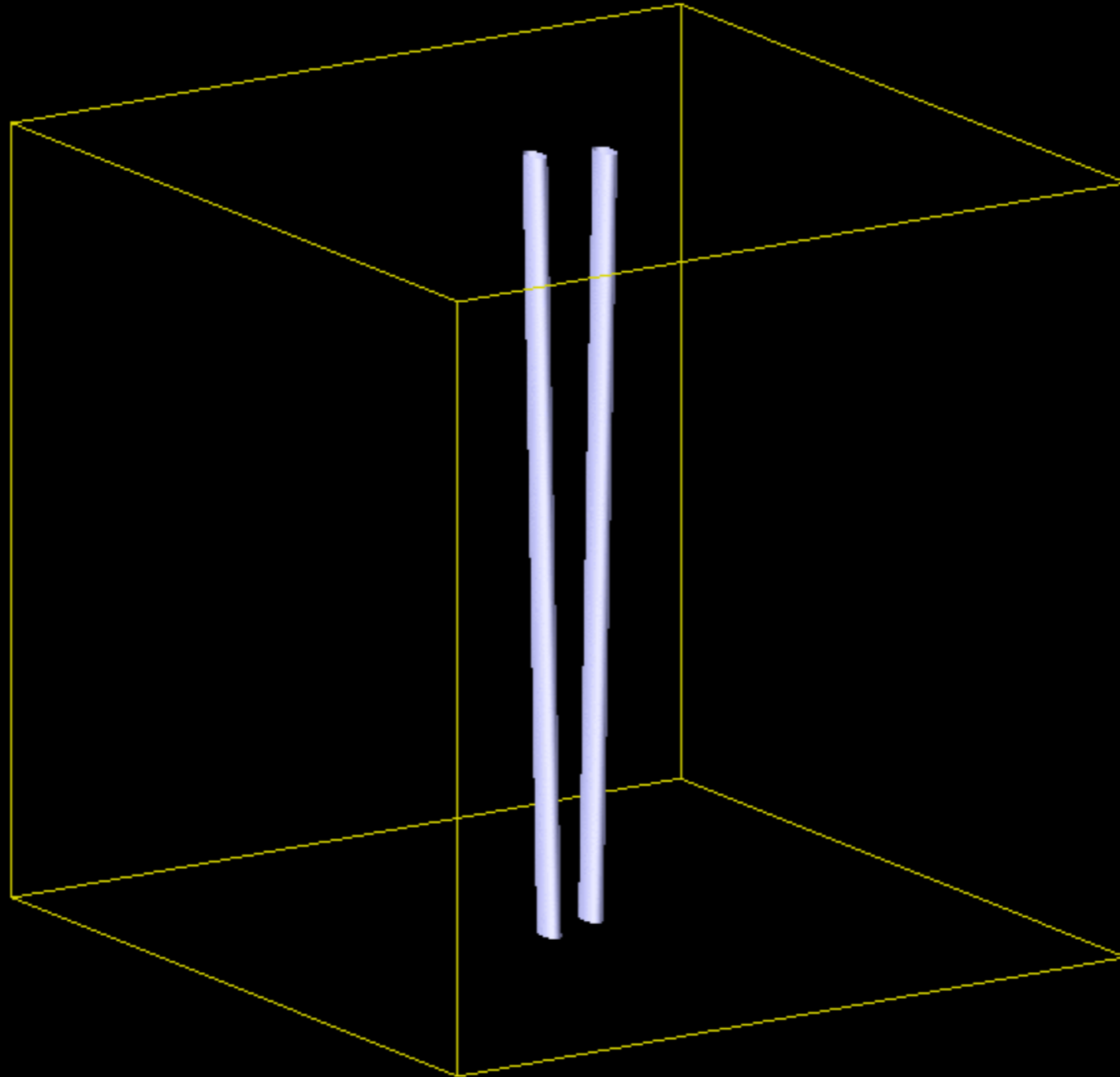
Hiramatsu, Eto, Kamada, Kobayashi, YO, '14



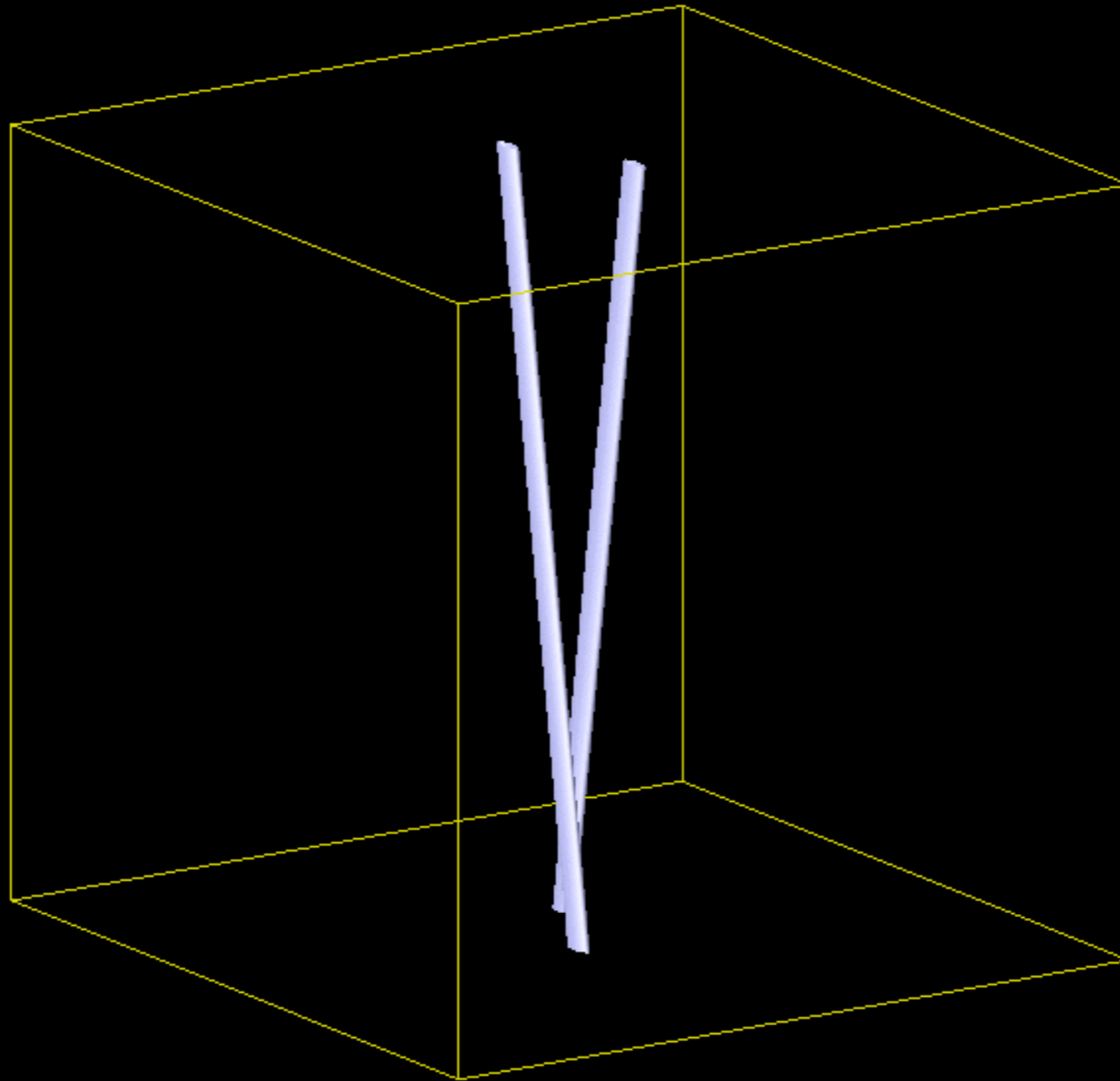
String($n=1$) / String($n=1$) collision
80% speed of light, **small angle**



String($n=1$)/String($n=1$) collision
86% speed of light

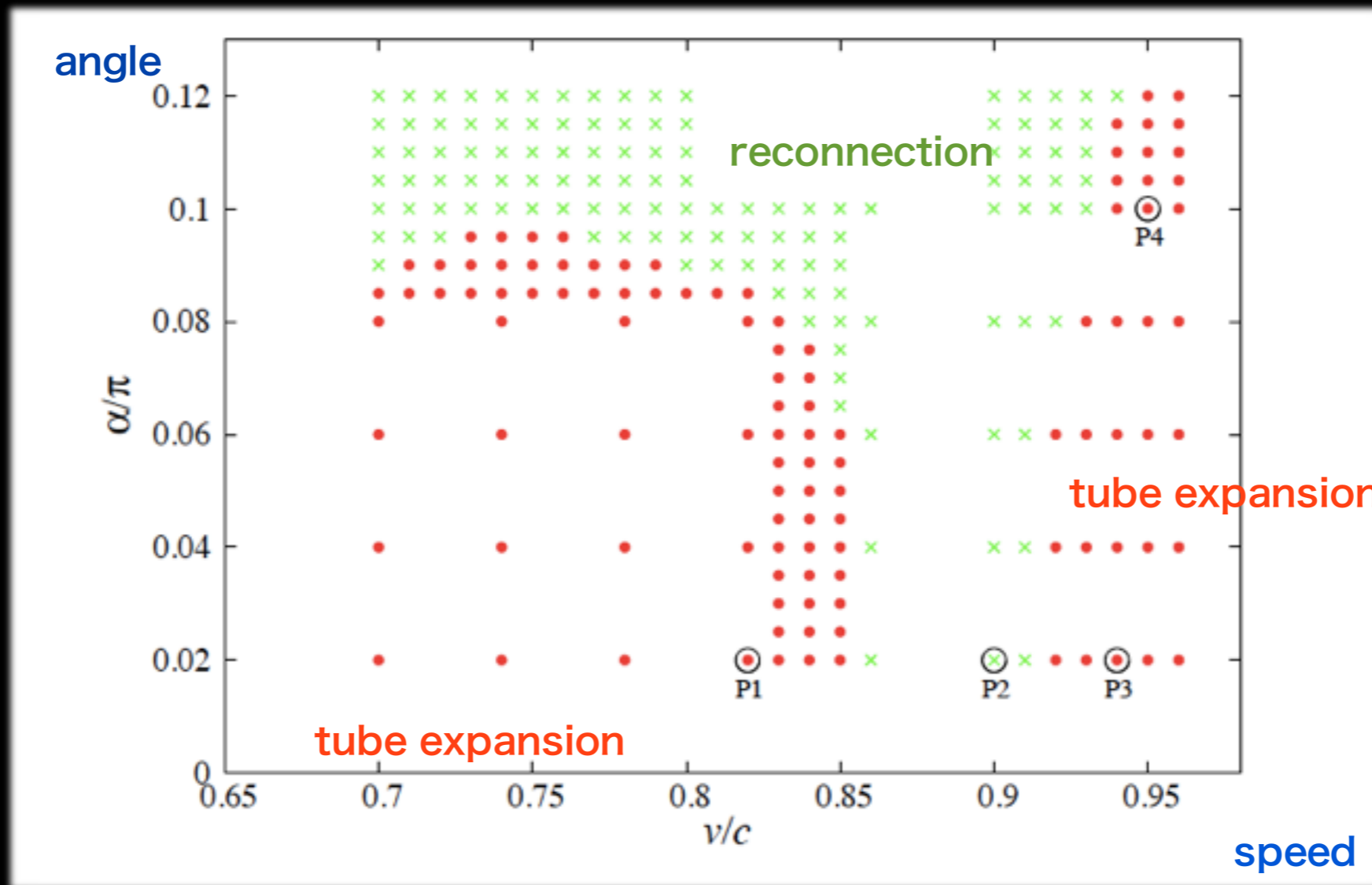


String($n=1$) / String($n=1$) collision
94% speed of light

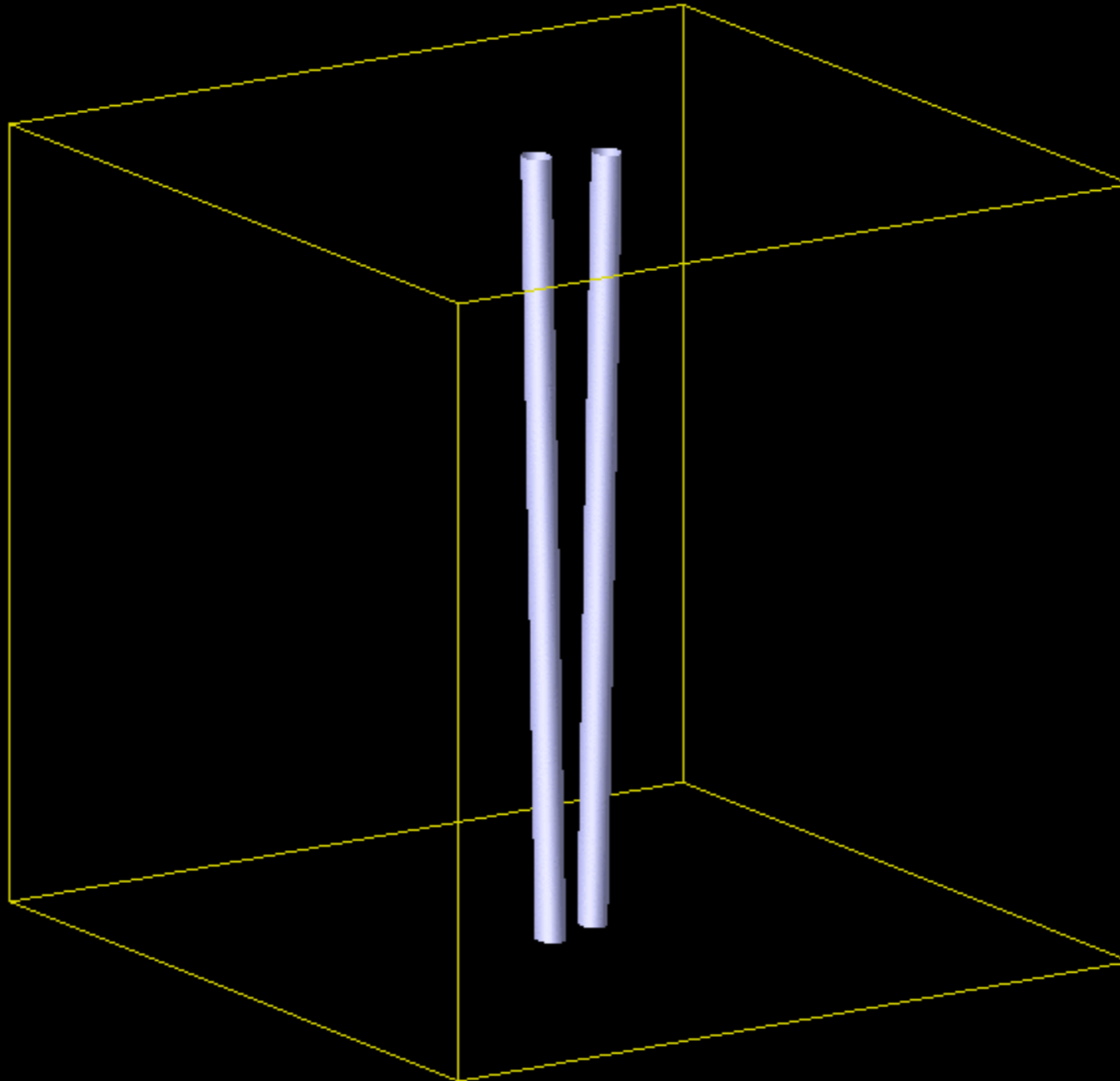


String/String collision

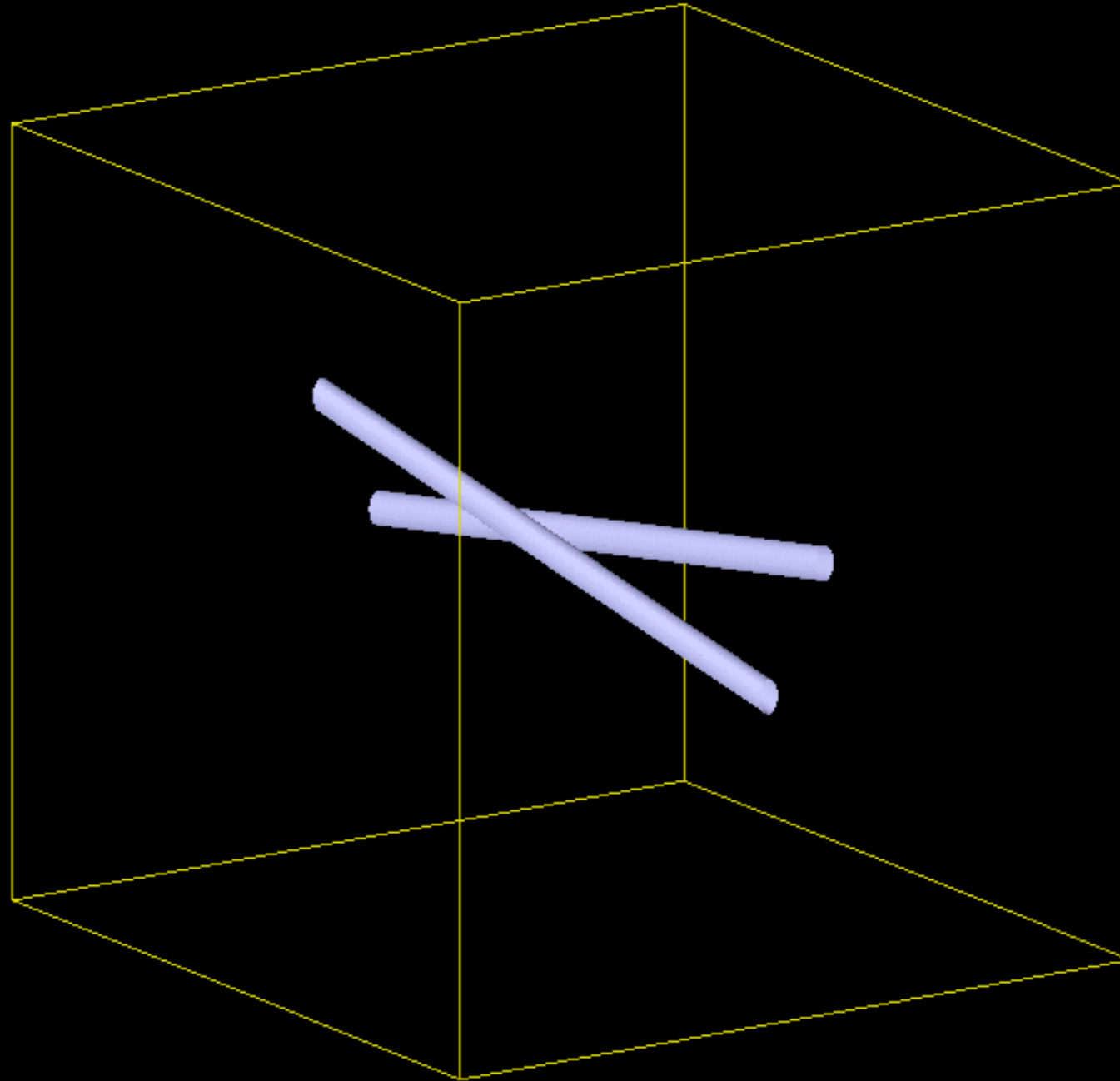
we observed three patterns



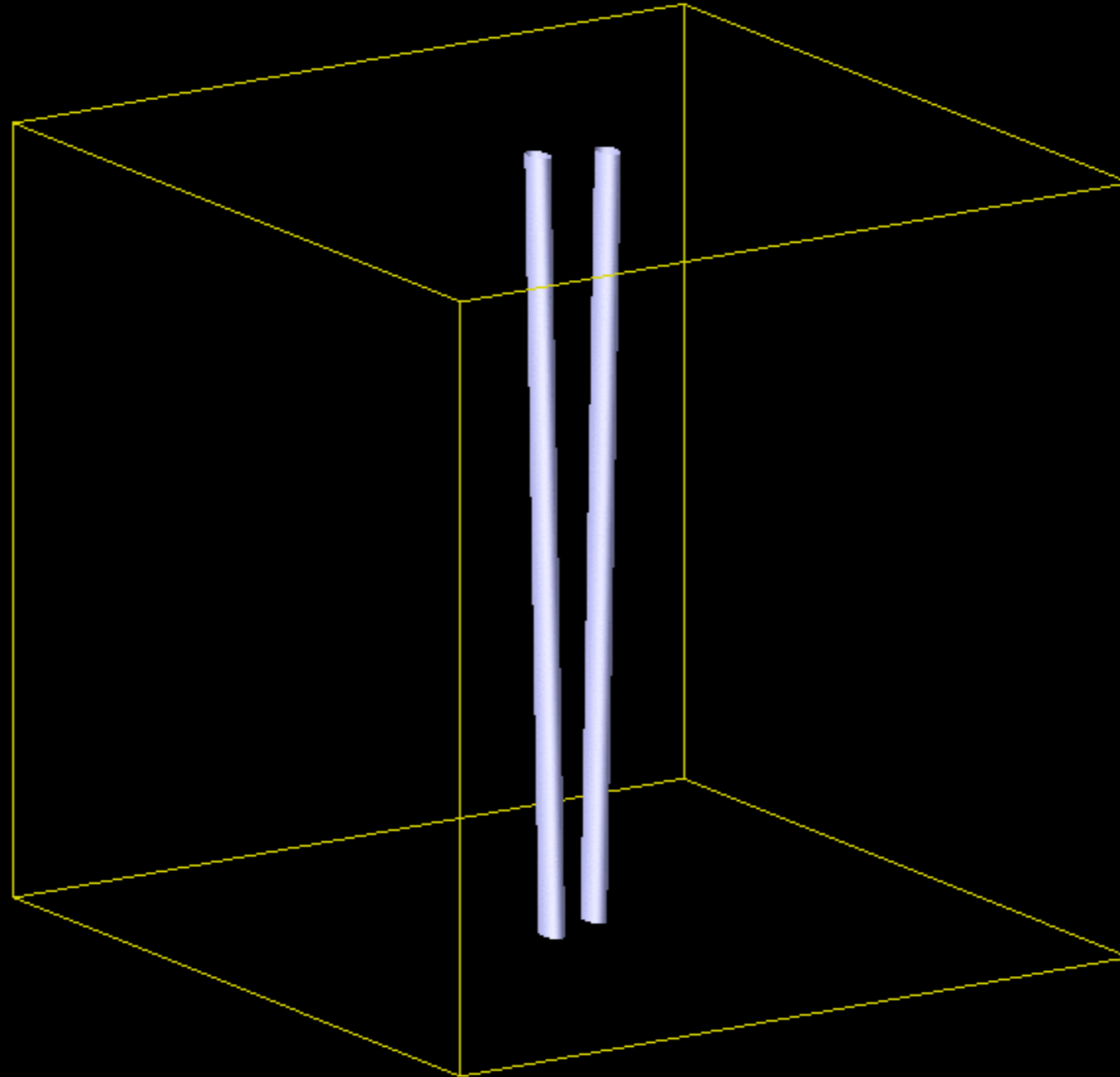
String($n=1$)/Anti-String($n=-1$) collision
60% speed of light



String($n=1$)/Anti-String($n=-1$) collision
80% speed of light, **large angle**

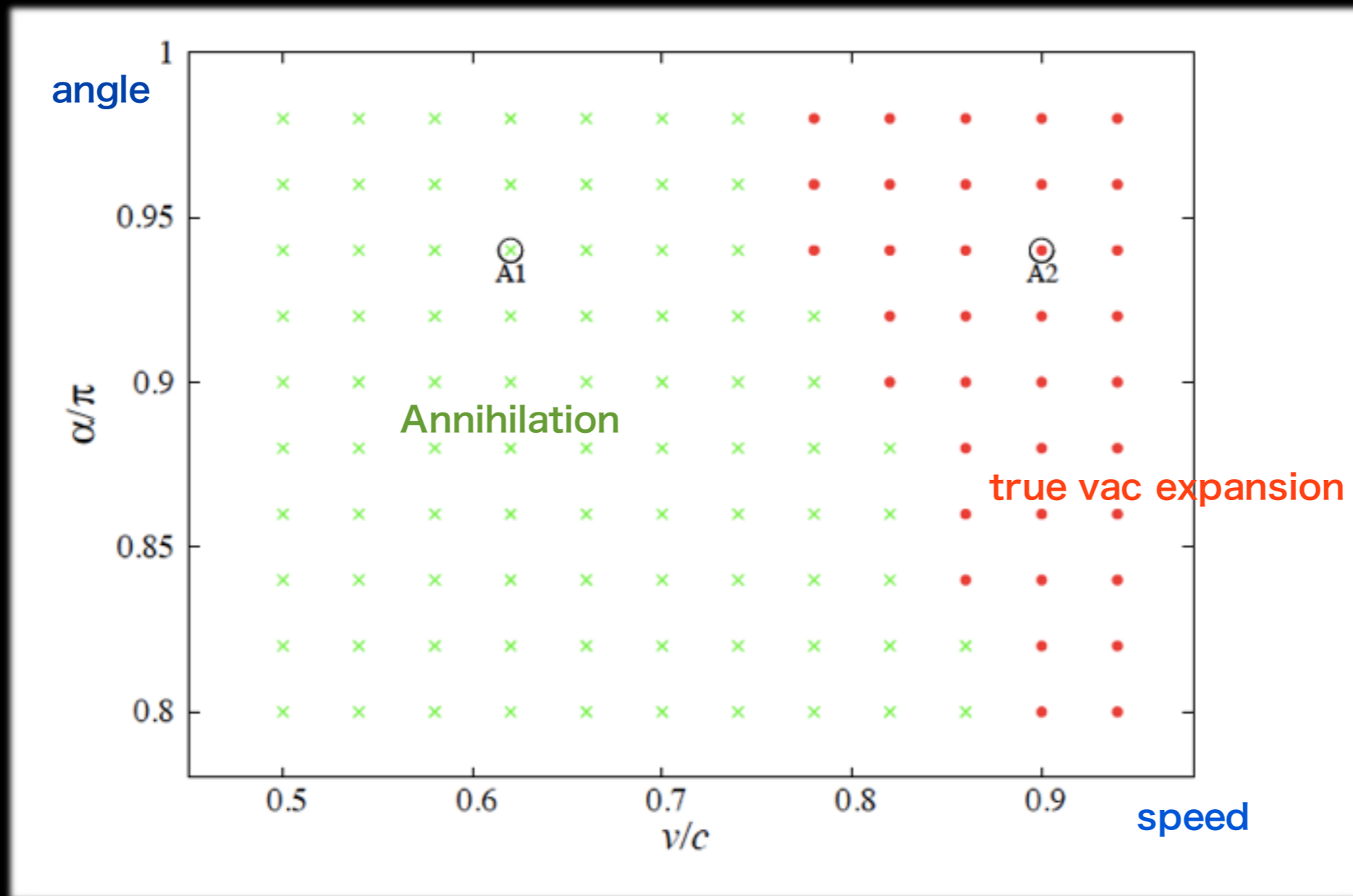


String($n=1$) / Anti-String($n=-1$) collision
80% speed of light, **small angle**



String/Anti-string collision

we observed two patterns



Application to String Theory

Type IIB String Theory compactified on Calabi-Yau manifold

Set up of geometry

M. Aganagic, C. Beem, J. Seo and C. Vafa, Nucl. Phys. B 789, 382 (2008)

Kasai, YO 16

Use the Calabi-Yau manifold

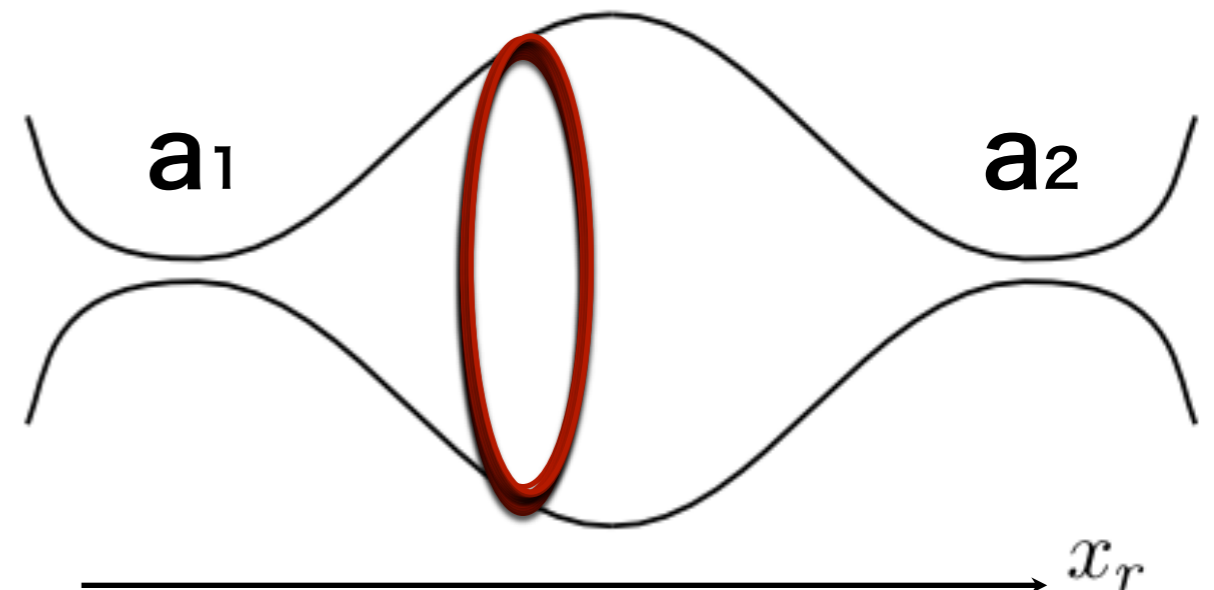
$$W'(x)^2 + y^2 + z^2 + w^2 = 0$$

$$W'(x) = g(x - a_1)(x - a_2).$$

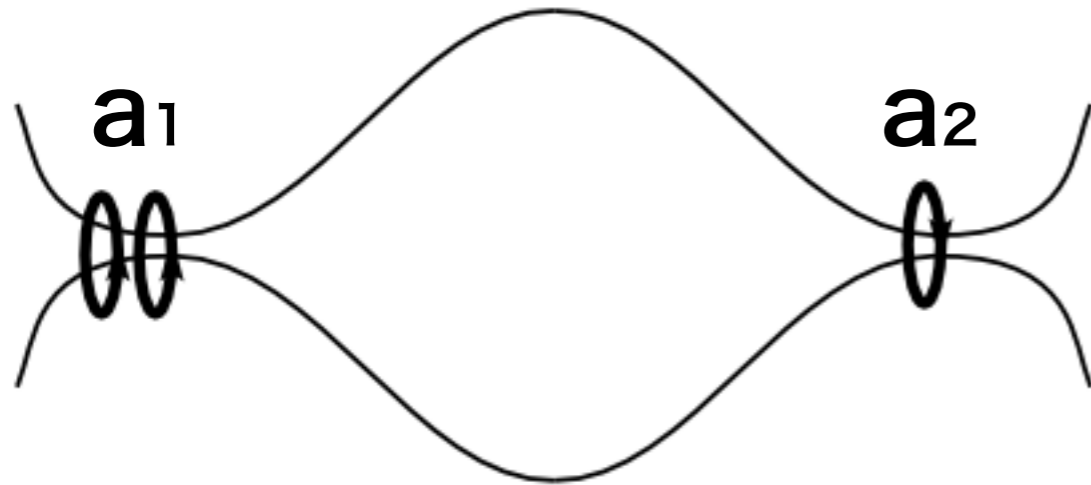
Important 3-dim sub-manifold

$$y_i^2 + z_i^2 + w_i^2 = g^2(x_r - a_1)^2(x_r - a_2)^2$$

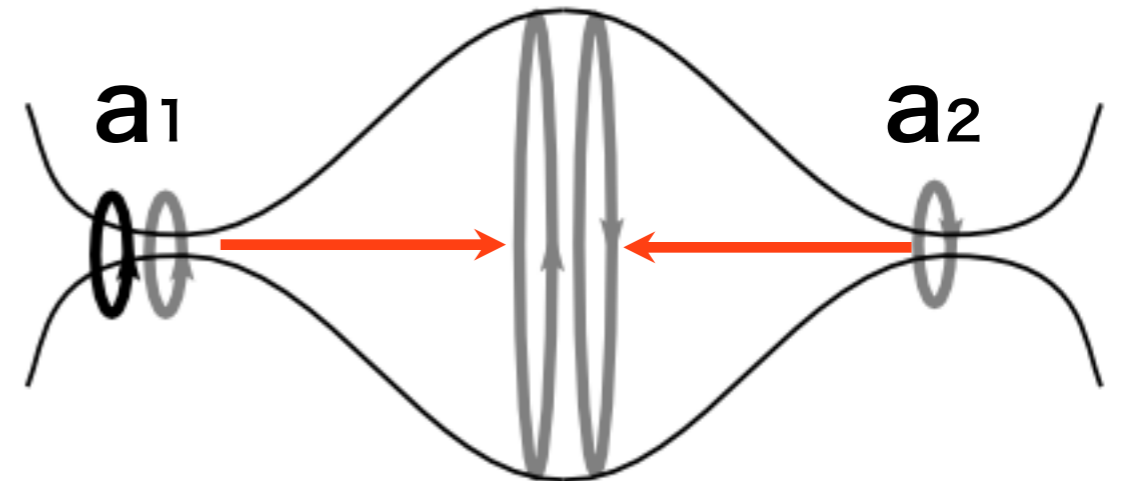
3-cycle=set of S^2



True and false vacua



false vacuum



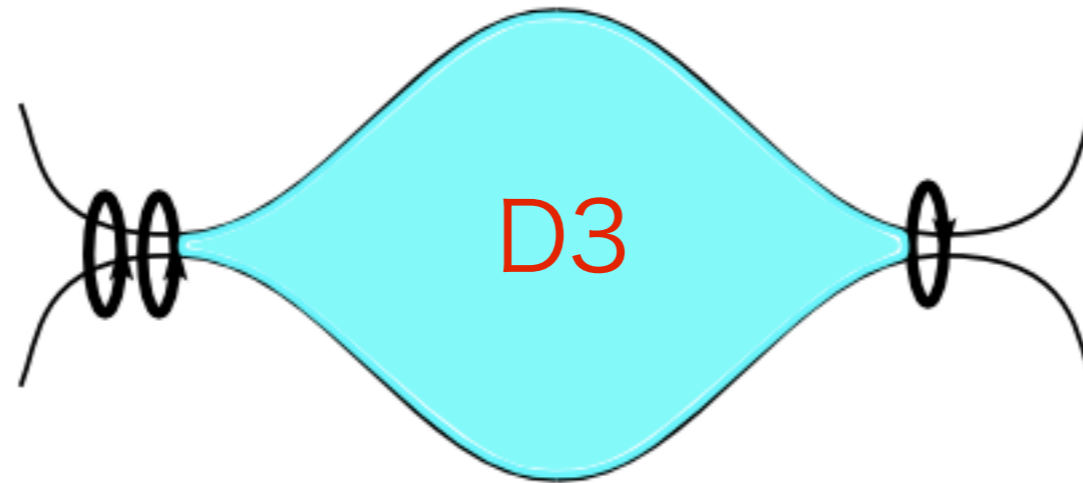
true vacuum

Wrapping n D5 branes at a_1

Wrapping n anti-D5 branes at a_2

→ D5/anti-D5
annihilation

Add D3 brane as an impurity



D3 plays a role of monopole

	0	1	2	3	4	5	6	7	8	9
D3	○	×	×	×	○	○	○	×	×	×
DW	○	△	△	△	○	○	○	×	×	×

Bound state of D3/DWD5

bound state energy

$$E_{D5/D3} = \sqrt{(T_{D5}4\pi R^2 V_3)^2 + (n_{D3}T_{D3}V_3)^2}.$$



total energy of system

$$E_{\text{total}} = \sqrt{(T_{D5}4\pi R^2 V_3)^2 + (n_{D3}T_{D3}V_3)^2} - \frac{4\pi}{3}\Delta V R^3 + \text{const},$$

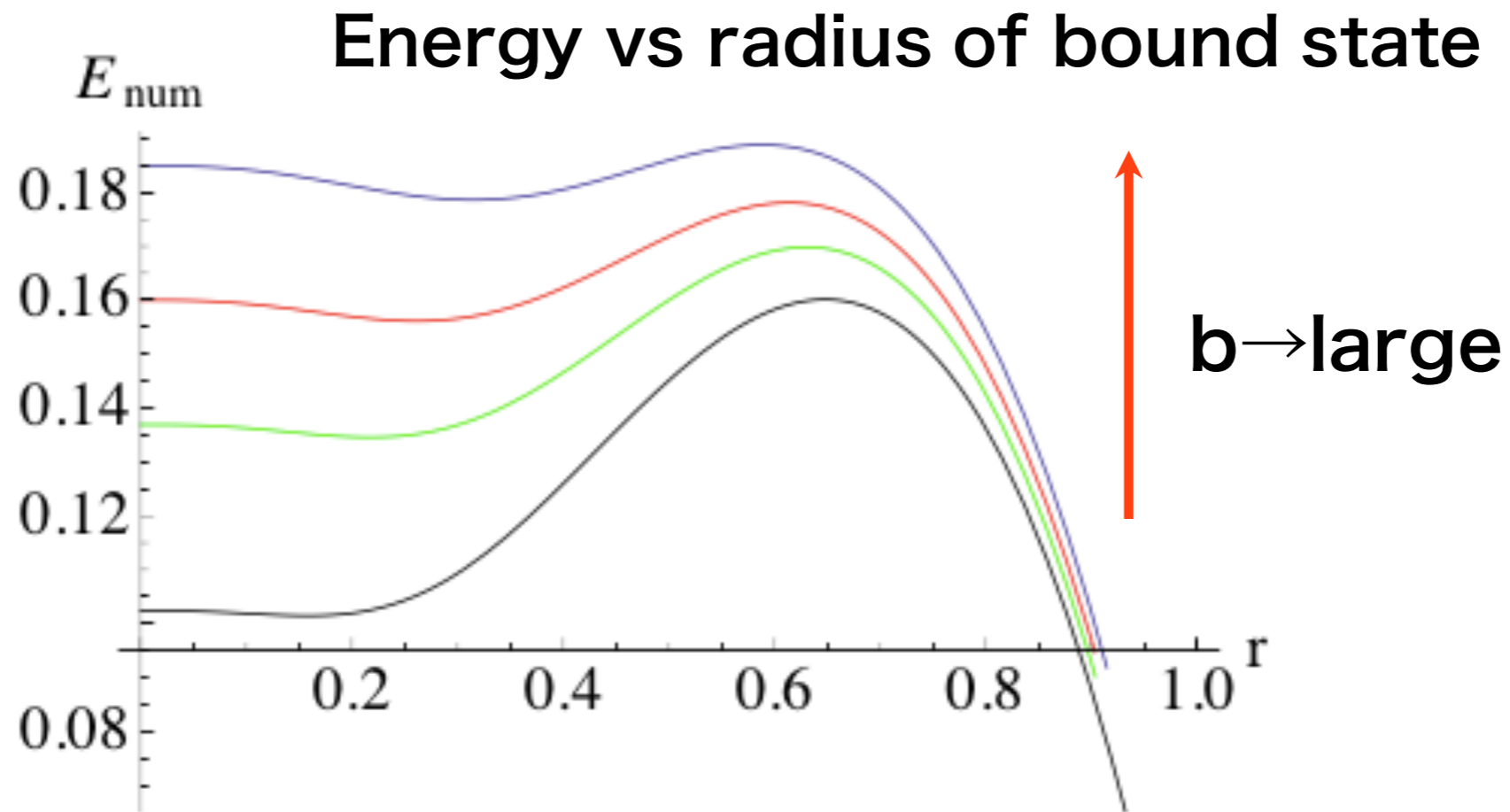
dimensionless parameters

$$r = \frac{\Delta V}{3T_{DW}}R,$$

winding number of monopole

$$b^4 = \left(\frac{\Delta V}{3T_{DW}}\right)^4 \left(\frac{T_{D3}n_{D3}}{4\pi T_{D5}}\right)^2.$$

Bound state of D3/DWD5

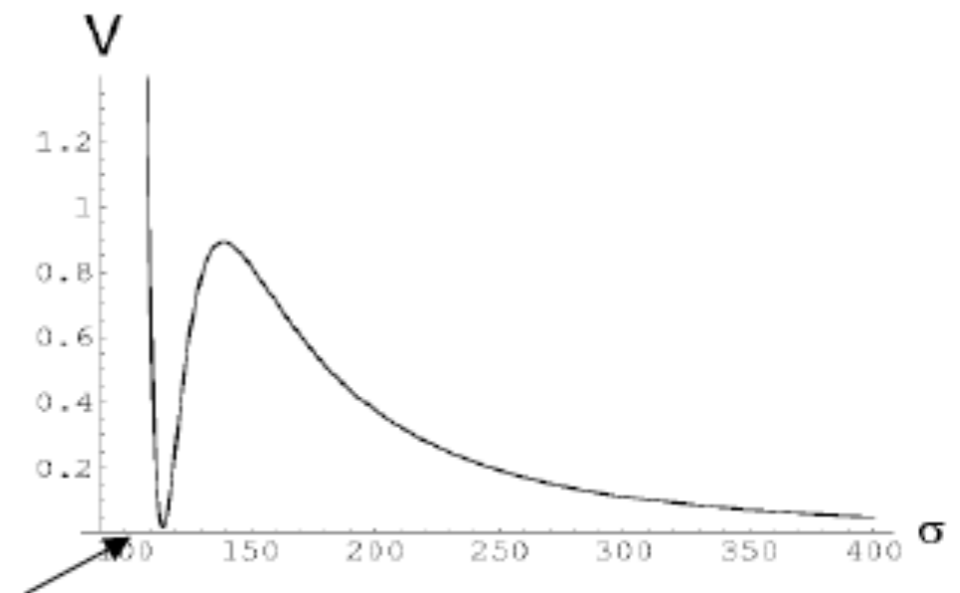
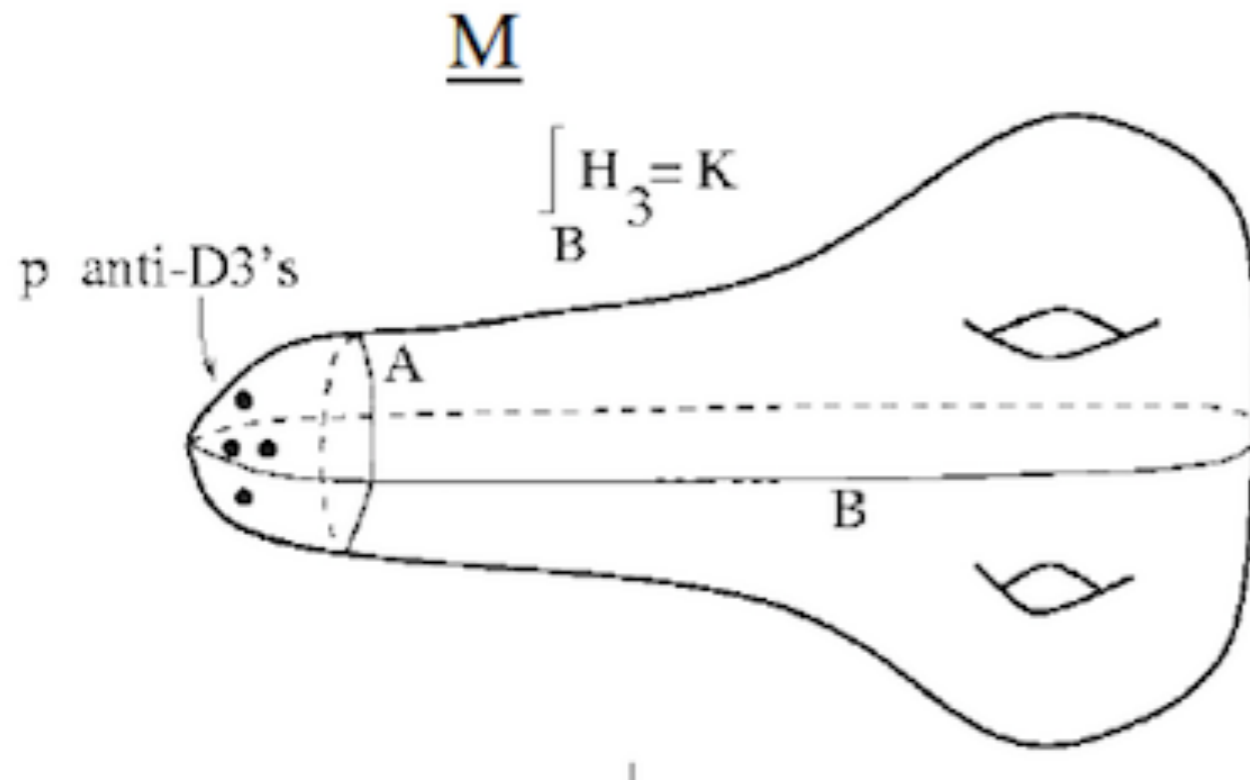
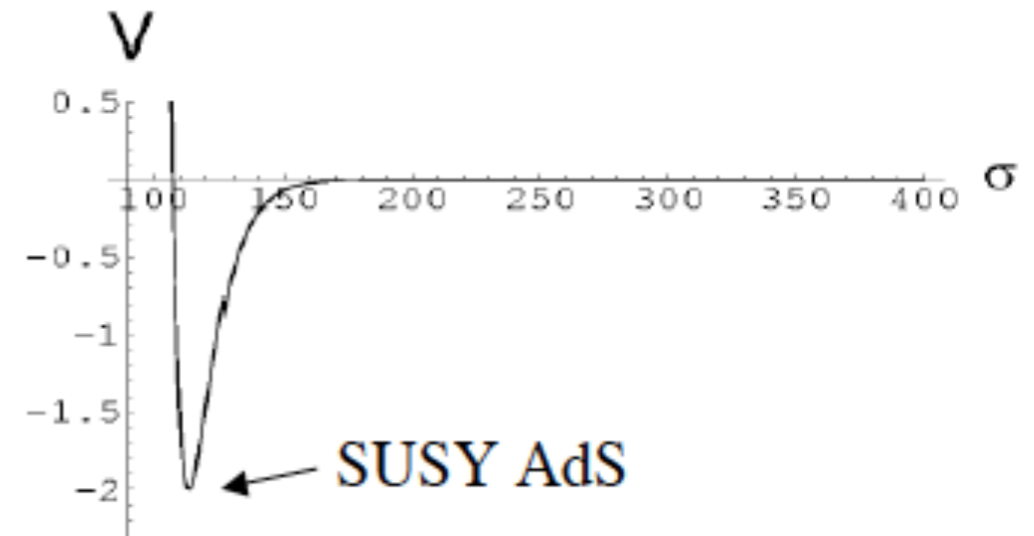


When b is large, the bubble is unstable

$$b > \frac{\sqrt{2}}{3} \quad (\text{unstable}). \quad \rightarrow \text{bubble expand without bound}$$

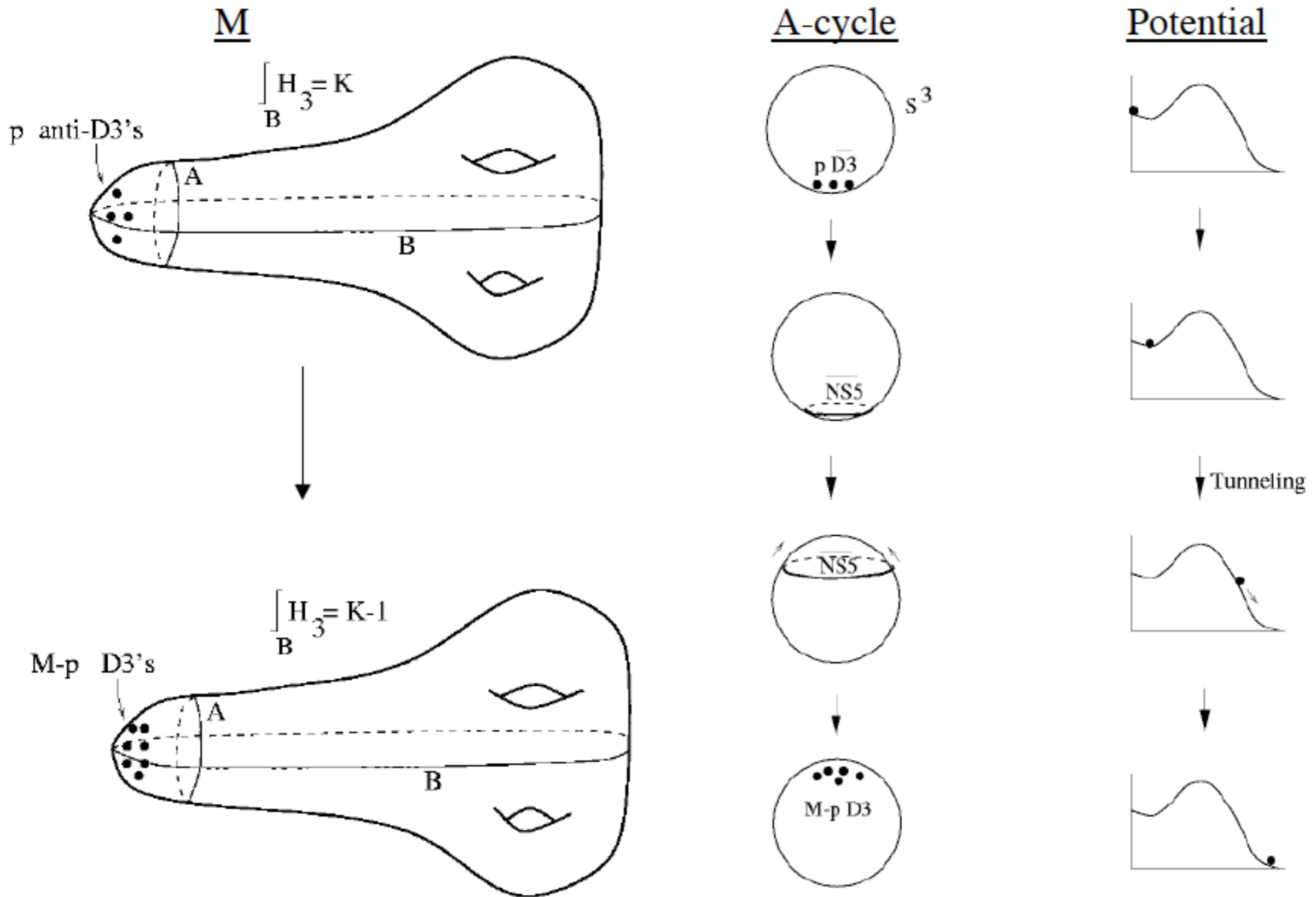
KKLT model

- ✓ compact Calabi-Yau
- ✓ stabilized moduli
- ✓ up-lifted by anti-D3 branes



small positive cosmological constant

KKLT model



KKLT model

We take non-compact limit for estimation of the decay rate

$$\frac{1}{G_{10}} \int d^{10}x(\dots) \simeq \frac{V_6}{G_{10}} \int d^4x$$

$$V_6 \rightarrow \infty$$

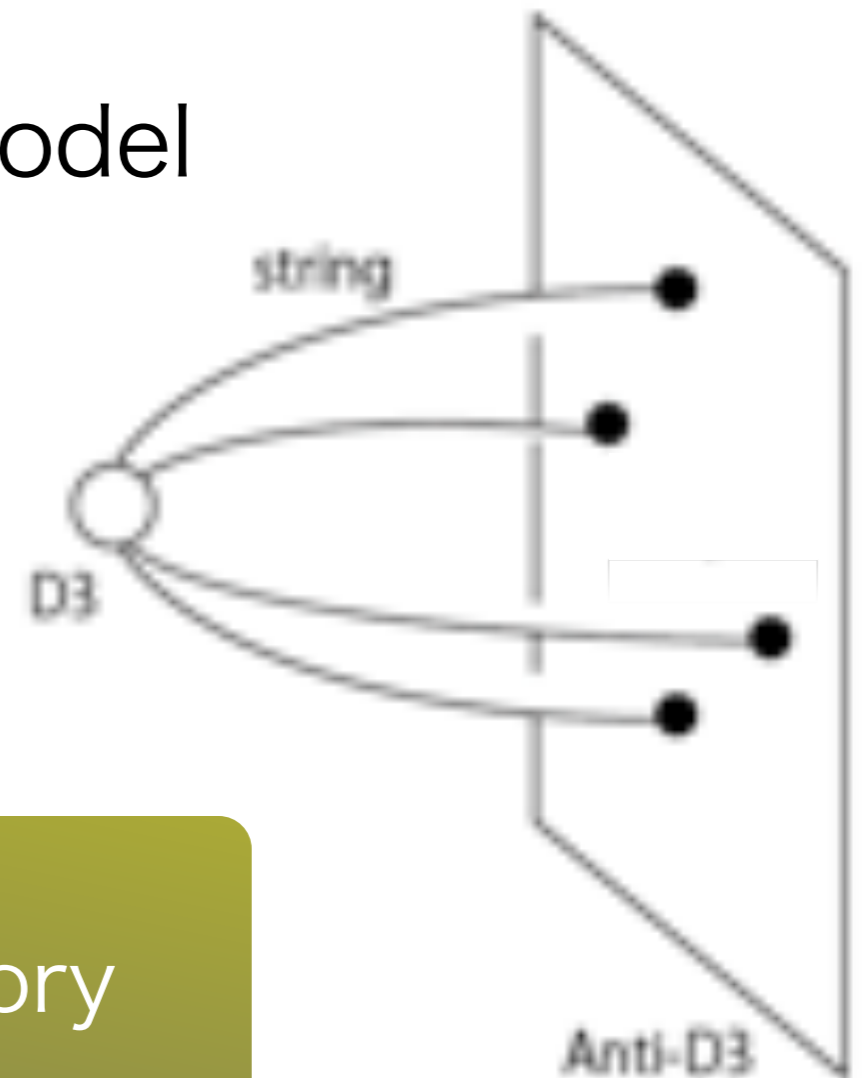
gravity effect is negligible

Dyon as an impurity

Let us apply our idea to the KKLT model

In this model, a **dyon** can be a seed for catalytic decay

Nakai, YO, Tanahashi : [arXiv 1808.10235](https://arxiv.org/abs/1808.10235)



Universal Lagrangian of Z_p gauge theory

$$p \int_{4D} B_2 \wedge F_2 \quad F_2 = dA$$

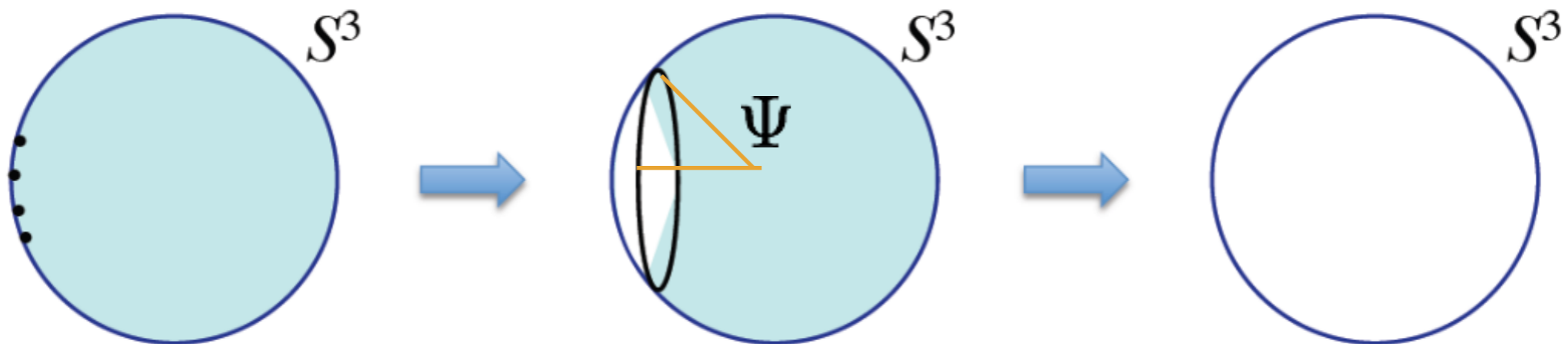
[Banks-Seiberg '10]

Action of the NS-five brane

$$S = -\frac{T_{NS}}{g_s^2} \int d^6 \xi \sqrt{-\det \left(g_{ab} - 2\pi g_s \alpha' \tilde{\mathcal{F}} \right)} - T_{NS} \int B_6 .$$

$$\tilde{\mathcal{L}}_D = -\sqrt{V_2(\Psi)^2(\tilde{r}^4 + \mathcal{B}^2) + \mathcal{D}^2 \sqrt{1 - \dot{\Psi}^2 + \Psi'^2}} + \frac{\tilde{r}^2}{\pi} \left(\Psi - \frac{1}{2} \sin 2\Psi \right)$$

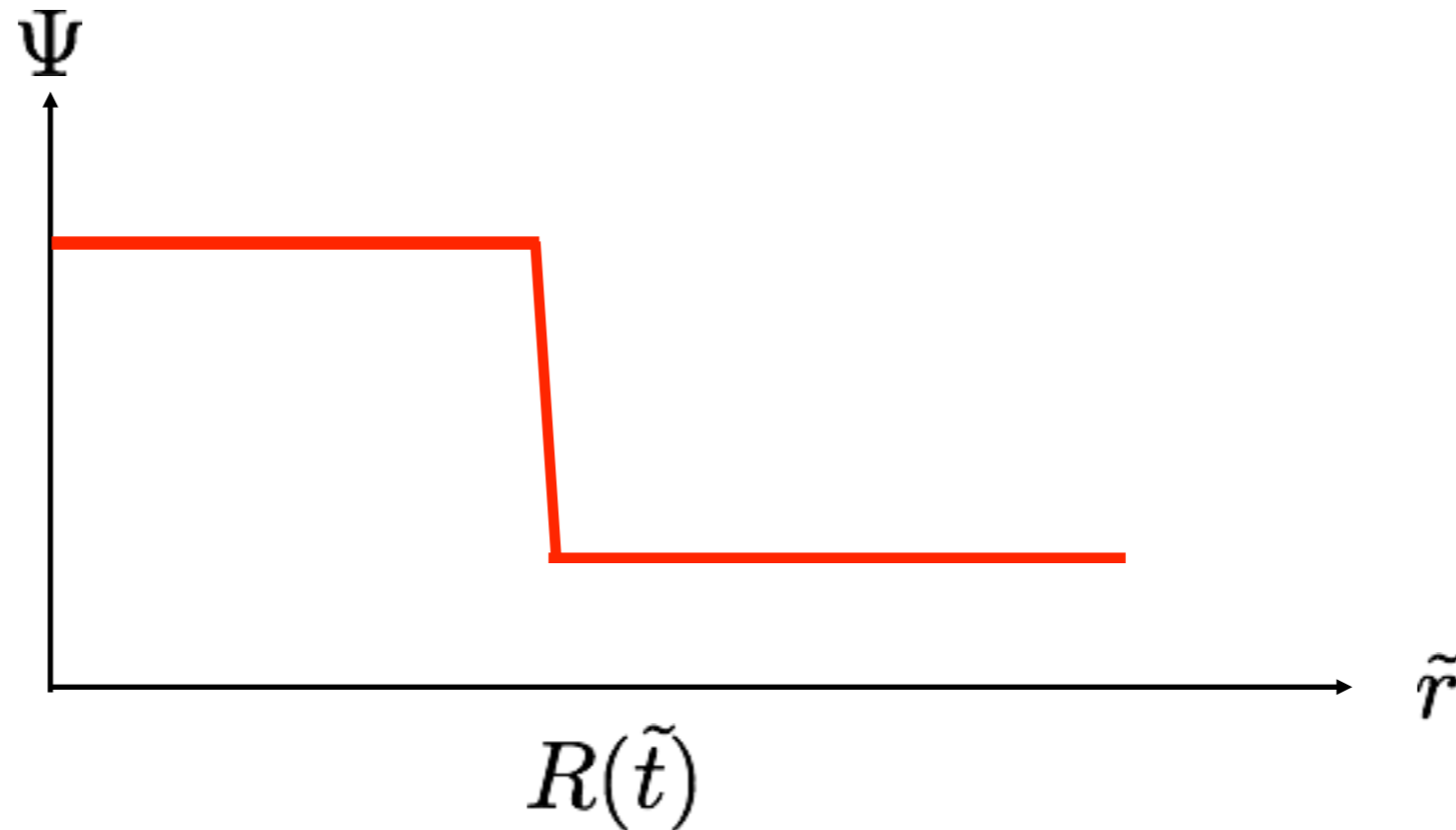
$$V_2(\Psi) = \frac{1}{\pi} \sqrt{b_0^4 \sin^4 \Psi + \left(\frac{\pi p}{M} - \left(\Psi - \frac{1}{2} \sin 2\Psi \right) \right)^2} .$$



Life-time of KKLT vacuum

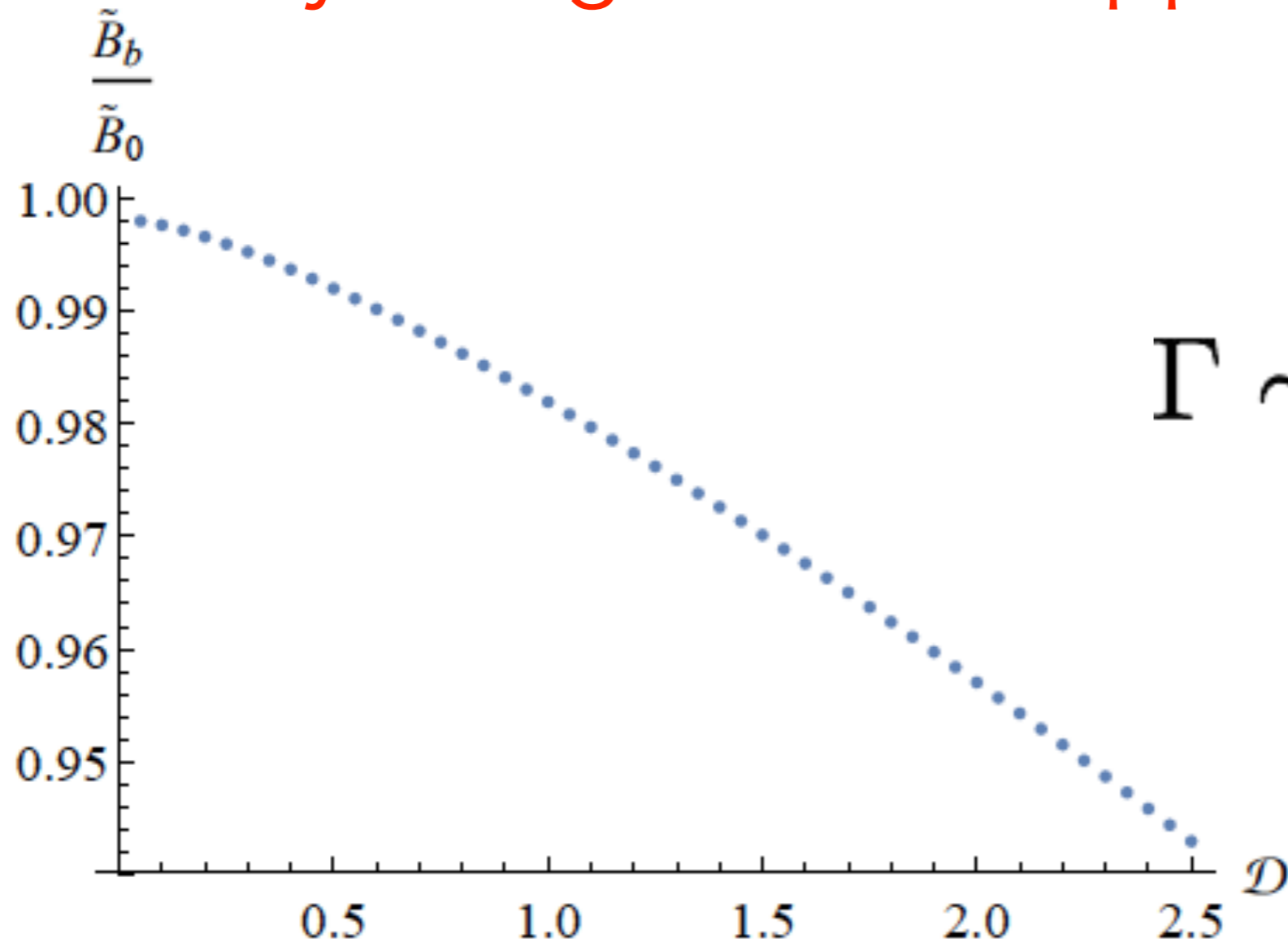
by using thin-wall approximation

$$\Psi = (\Psi_{\max} - \Psi_{\min}) [1 - \theta(\tilde{r} - R(\tilde{t}))] + \Psi_{\min} .$$



Life-time of KKLT vacuum

by using thin-wall approximation



$$\Gamma \sim \exp\left(-\tilde{B}_b\right)$$

$$\tilde{B}_0 > \tilde{B}_b$$

Life-time is drastically shorter!

Summary

Claim 1) Vacuum instability by **catalysis** plays crucial role in string landscape

Claim 2) In estimating the life-time of vacuum, the shape of the potential is not enough

Claim 3) Bubble can be a giant monopole

Claim 4) Topology of a bubble is not necessarily sphere

Claim 5) KKLT de Sitter vacuum may be unstable