

# $B \rightarrow P, V$ Form Factors from Light-Cone Sum Rules



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# An overview of the talk

- Introduction: Form Factors (FF) definition
- Motivation: The importance of FFs in physics and B anomalies
- Method: B light-cone Sum Rules (B-LCSR)
- Theoretical Results: New higher twist corrections to the B-LCSR
- Numerical Results:  
How big are these higher twist contributions?

## Introduction: what are the Form Factors?

FFs are functions of  $q^2$  that parametrize exclusive local hadronic elements. They are defined as follows:

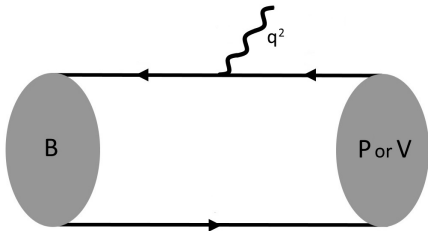
- for  $B \rightarrow P$

$$\langle P(k) | \bar{q}_1 \gamma_\mu b | \bar{B}(k+q) \rangle = 2k_\mu f_{BP}^+(q^2) + q_\mu [f_{BP}^+(q^2) + f_{BP}^-(q^2)]$$

- for  $B \rightarrow V$

$$\begin{aligned} \langle V(k) | \bar{q}_1 \gamma_\mu (1 - \gamma_5) b | \bar{B}(k+q) \rangle = & -i\varepsilon_\mu^* (M_B + M_V) A_1^{BV}(q^2) + i(2k + q)_\mu \frac{A_2^{BV}(q^2)}{M_B + M_V} \\ & + iq_\mu (\varepsilon^* \cdot q) \frac{2M_V [A_3^{BV}(q^2) - A_0^{BV}(q^2)]}{q^2} + \varepsilon_{\mu\nu\rho\sigma} \varepsilon^{*\nu} q^\rho k^\sigma \frac{2V^{BV}(q^2)}{M_B + M_V} \end{aligned}$$

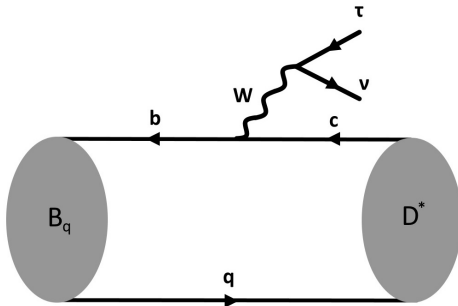
[ $q^2$  is the dilepton mass squared]



# Motivations: why do we need $B \rightarrow P, V$ Form Factors?

In general  $B \rightarrow P, V$  FFs are needed to

- predict decay amplitudes, such as  $B \rightarrow \{P, V\}\bar{l}l$  or  $B \rightarrow \{P, V\}l\nu_l$   
[In this work we consider the final states:  $P = \pi, K, D$  and  $V = \rho, K^*, D^*$ ]
- extract  $|V_{CKM}|$  element from branching ratios
- test the Standard Model and constrain new physics contributions



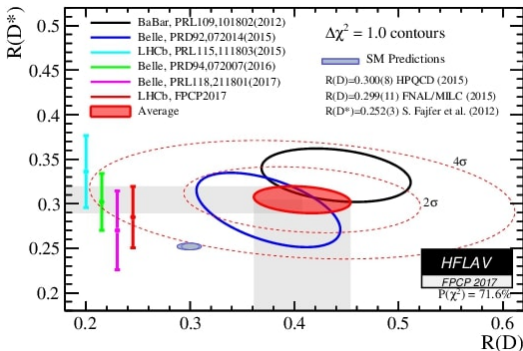
## B anomalies: $R_{K^{(*)}}$ and $R_{D^{(*)}}$

LHCb, Belle and BaBar measurements exhibit a deviation with respect the SM prediction for the observables

$$R_{K^{(*)}} = \frac{BR(B \rightarrow K^{(*)} \mu^+ \mu^-)}{BR(B \rightarrow K^{(*)} e^+ e^-)}$$

$$R_{D^{(*)}} = \frac{BR(B \rightarrow D^{(*)} \tau^+ \nu)}{BR(B \rightarrow K^{(*)} \mu^+ \nu)}$$

In both cases the combination of the results leads to a  $4\sigma$  **deviation** from the SM prediction.



## Method: The techniques to compute the FFs

QCD perturbation theory breaks down at low energy, hence **non perturbative techniques** are needed to compute local hadronic matrix elements, that is the form factors.

The most successful methods to compute  $B \rightarrow P, V$  are FFs are Lattice QCD and light-cone sum rules.

### Lattice QCD

spacetime discretization

smaller error

High  $q^2$  (low recoil)

### Light-cone sum rules

quark-hadron duality

bigger error

Low  $q^2$  (large recoil)

# Light-Cone Sum Rules in a nutshell

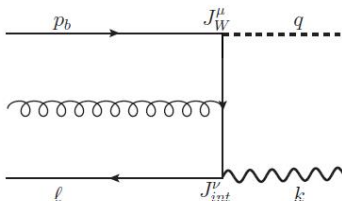
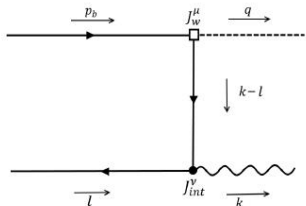
- LCSR are used to determine products of exclusive hadronic matrix elements from an artificial, less-exclusive, non-local hadronic matrix element  $\Pi(k^2, q^2)$
- $\Pi(k^2, q^2)$  is then **expanded near the light-cone**

$$\Pi(k^2, q^2) = f_B m_B \int ds \sum_{n,t} \frac{J_{n,t}(s, q^2)}{[k^2 - s]^n} \phi_t(s)$$

- $J_{n,t}$  can be computed from a hard scattering kernel
- B-meson Light-Cone Distribution Amplitudes (LCDAs)  $\phi_t$  are necessary non-perturbative input
  - general  $B \rightarrow V$ ,  $B \rightarrow P$  results available [Khodjamirian et al. '06 + '08]
  - new insights on LCDAs triggered our revisiting of these sum rule results [Braun/Ji/Manashov '17]

# Theoretical results

- LCSRs together with lattice results and Heavy Quark expansions have been used in present analyses
- B-LCSRs have  $1/m_b$  corrections (related to twist expansion)
- We present **new twist 4 corrections** to the  $B \rightarrow P, V$  LCSRs, higher twists are expected to give corrections only of the order  $O(1/m_b^2)$
- $O(\alpha_s)$  corrections are not considered





# Numerical Results and Comparison

$B \rightarrow D^* FF$	FKKM2008	GKvD2018	NEW Contrib.
	2pt tw2+3 + 3pt	2pt tw2+3 + 3pt*	2pt tw4
$A_0(q^2 = 0)$	0.78	0.79	-10%
$A_1(q^2 = 0)$	0.73	0.73	-12%
$A_2(q^2 = 0)$	0.66	0.65	-17%
$A_0(0)/A_1(0)$	1.07	1.09	+3%

[using the same input parameters, with  $q^2$  the dilepton mass square]

$\phi_+, \phi_-$  2-particle L+NL twist contributions [Faller/Khodjamirian/Klein/Mannel '06]

$g_+$  new 2-particle NNL twist contributions [Gubernari/Kokulu/van Dyk w.i.p.]

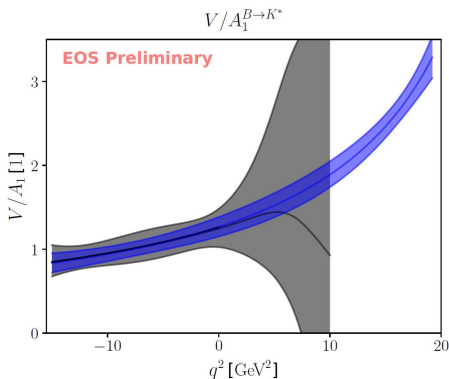
$\phi_3, \phi_4$  new and self-consistent 3-particle NL+NNL twist contr.

# Combining Lattice with LCSRs

- Uncertainties for the LCSRs are of similar size as in previous works
- To improve the knowledge of the FFs over the whole  $q^2$  range, we fit LCSR and Lattice results to the BSZ2015 parametrization.

[Bharucha, Straub, Zwicky 2015]

- The fit to the BSZ2015 parametrization makes uncertainties smaller



Thank you for your attention!

## Power corrections

- correlator is calculated with on-shell  $B$  meson, using its Light-Cone Distribution Amplitudes (LCDAs)
- $B$ -meson LCDAs are defined for bi-local currents involving an HQET field  $h_v$
- power corrections to this involve power of the covariant derivative  $iD^\mu$
- strings of the type  $iD^{\mu_1} iD^{\mu_2} \dots iD^{\mu_n}$  are incorporated in LCDAs of increasing (collinear) twist

## Benefits of the Braun et al. basis

- $\phi_3, \phi_4, \dots$  are LCDAs of definite collinear twist 3, 4,  $\dots$
- LCDAs of twists  $\geq 5$  are expected to contribute *beyond* the next-to-leading  $1/m_b$  corrections! [Braun/Ji/Manashov '17]
- inserting a gluon field adds at least one unit of twist
  - 2-particle LCDAs start at twist 2, and are included in our results (up to and including twist 4)
  - 3-particle LCDAs start at twist 3, and are included in our results (up to and including twist 4)
  - 4-particle LCDAs start at twist 4, and are *not included in our results*
  - 4-particle LCDAs have autonomous RG behaviour, *do not mix with 3-particle LCDAs*

[Braun/Ji/Manashov '17]