

# Signatures of the Type-I 2HDM at the LHC

# Shoaib Munir KIAS, Seoul

Workshop on the Standard Model and beyond, Corfu Sep. 02, 2018





@ 2-Higgs-Doublet Models

*■ Type-I:* 

- > Electroweak production of light scalarpseudoscalar pairs
- > A (fairly) light charged Higgs boson
- > The promise of multi-photon final states
- $\succ$  EW vs. QCD production of multiple Higgs bosons at the LHC

Conclusions

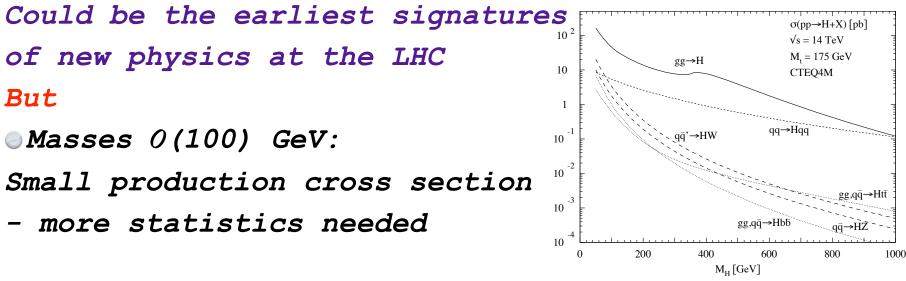


Predicted in a minimalistic new physics contender like the 2-Higgs-Doublet Model and in extended frameworks like Supersymmetry and GUTs

Could provide earliest

signatures of new physics

Predicted in a minimalistic new physics contender like the 2-Higgs-Doublet Model and in extended frameworks like Supersymmetry and GUTs



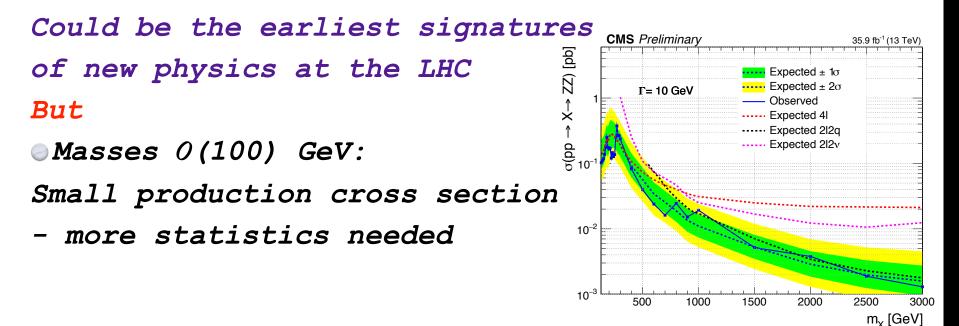
[A. Djouadi, hep-ph/0503173]

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Predicted in a minimalistic new physics contender like the 2-Higgs-Doublet Model and in extended frameworks like Supersymmetry and GUTs



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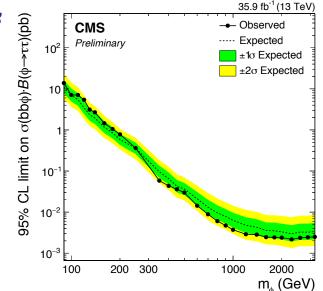
Predicted in a minimalistic new physics contender like the 2-Higgs-Doublet Model and in extended frameworks like Supersymmetry and GUTs

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Could be the earliest signatures
of new physics at the LHC
But
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*Masses 0(100) GeV:* 

Small production cross section

- more statistics needed

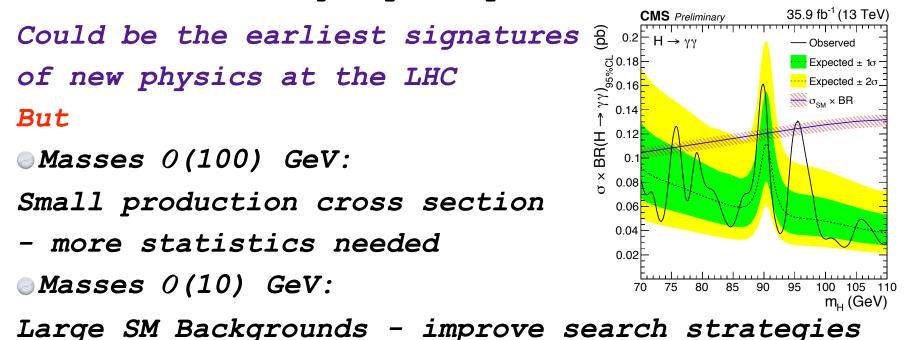


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Predicted in a minimalistic new physics contender like the 2-Higgs-Doublet Model and in extended frameworks like Supersymmetry and GUTs



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Predicted in a minimalistic new physics contender like the 2-Higgs-Doublet Model and in extended frameworks like Supersymmetry and GUTs

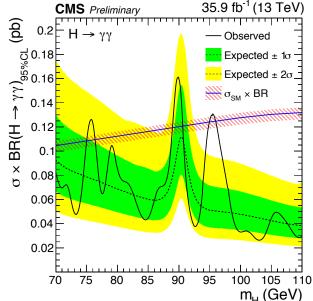
Could be the earliest signatures  $\widehat{e}_{H \rightarrow \gamma\gamma}$ of new physics at the LHC But

*Masses 0(100) GeV:* 

Small production cross section

- more statistics needed

*Masses 0(10) GeV:* 



Large SM Backgrounds - improve search strategies Also (in either case)

decay rates to SM particles may be suppressed

Exploit Higgs-Higgs and Higgs-gauge production

# **2HDM - SCALAR POTENTIAL**

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The Yukawa Lagrangian for the neutral scalars reads  $-\mathcal{L}_{Y} = \overline{Q}_{L}\widetilde{\Phi}_{1}\eta_{1}^{U}U_{R} + \overline{Q}_{L}\Phi_{1}\eta_{1}^{D}D_{R} + \overline{Q}_{L}\Phi_{1}\eta_{1}^{L}L_{R} + \overline{Q}_{L}\widetilde{\Phi}_{2}\eta_{2}^{U}U_{R} + \overline{Q}_{L}\Phi_{2}\eta_{2}^{D}D_{R} + \overline{Q}_{L}\Phi_{2}\eta_{2}^{L}L_{R}$ 

$$M^F = \frac{v}{\sqrt{2}} \left( \eta_1^F \cos\beta + \eta_2^F \sin\beta \right)$$

# **MINIMAL FLAVOUR VIOLATION**

 $\odot$  To prevent flavour-changing neutral currents, a  $Z_2$ symmetry can be imposed (removes CP-violating  $\lambda_{6,7}$ )

 $= Z_2 - charge$  assignment = four Types

Model	$u_R^i$	$d_R^i$	$e_R^i$
Type I	$\Phi_2$	$\Phi_2$	$\Phi_2$
Type II	$\Phi_2$	$\Phi_1$	$\Phi_1$
Lepton-specific	$\Phi_2$	$\Phi_2$	$\Phi_1$
Flipped	$\Phi_2$	$\Phi_1$	$\Phi_2$



### MINIMAL FLAVOUR VIOLATION

To prevent flavour-changing neutral currents, a  $Z_2$  symmetry can be imposed (removes CP-violating  $\lambda_{6,7}$ )

 $\odot Z_2$ -charge assignment  $\blacksquare$  four Types

Model	$u_R^i$	$d_R^i$	$e_R^i$	
Type I	$\Phi_2$	$\Phi_2$	$\Phi_2$	$\xi_f^h = \cos\alpha / \sin\beta$
Type II	$\Phi_2$	$\Phi_1$	$\Phi_1$	$\xi_f^H = \sin \alpha / \sin \beta$
Lepton-specific	$\Phi_2$	$\Phi_2$	$\Phi_1$	$\zeta_f = \sin \alpha / \sin \beta$
Flipped	$\Phi_2$	$\Phi_1$	$\Phi_2$	

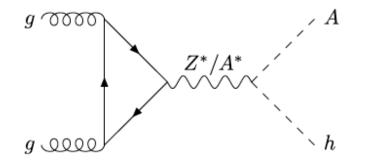
$$\cos \alpha = \sin \beta \sin(\beta - \alpha) + \cos \beta \cos(\beta - \alpha)$$

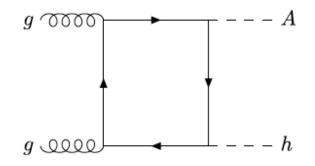
$$-\mathcal{L}_{\text{Yukawa}}^{\text{2HDM}} = \sum_{f=u,d,\ell} \frac{m_f}{v} \left( \xi_f^h \overline{f} fh + \xi_f^H \overline{f} fH - i\xi_f^A \overline{f} \gamma_5 fA + \left\{ \frac{\sqrt{2}V_{ud}}{v} \overline{u} \left( m_u \xi_u^A \mathbf{P}_L + m_d \xi_d^A \mathbf{P}_R \right) dH^+ + \frac{\sqrt{2}m_\ell \xi_\ell^A}{v} \overline{\nu_L} \ell_R H^+ + \text{h.c} \right\},$$



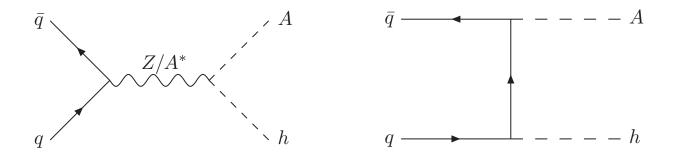
# A LIGHT SCALAR-PSEUDOSCALAR PAIR KIAS S IN STITUTE FOR

Landau-Yang theorem forbids the contribution of a resonant Z boson to the QCD production of a hA pair





but not to EW production: enhanced cross sections?



Numerically scanning of the parameter space (trading  $\lambda_{1-5}$  for the physical Higgs boson masses as input parameters), with the following constraints

	$m_{_H} = 125 \text{ GeV}$
$m_h \; ({\rm GeV})$	10 - 80
$m_A \; ({\rm GeV})$	$10 - (M_Z - m_h)$
$m_{H^{\pm}} (\text{GeV})$	90 - 500
$\sin(\beta - \alpha)$	-1 - 1
$m_{12}^2 \; ({\rm GeV}^2)$	$0 - m_A^2 \sin\beta\cos\beta$
aneta	2,25

*imposed:* Code: 2HDMC [D. Eriksson, J. Rathsman, O. Stal, 0902.0851]

- Unitarity, perturbativity and vacuum stability
- Oblique parameters S, T and U
- Flavour physics



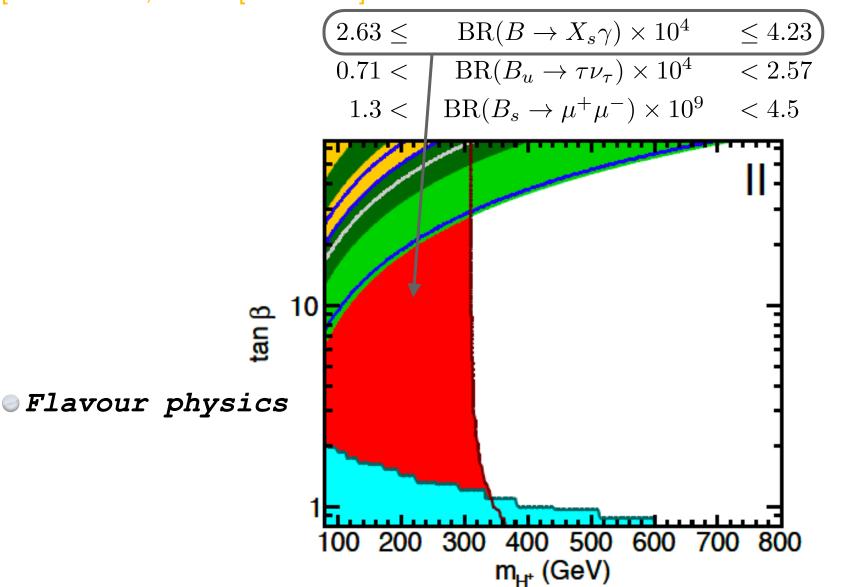
#### Superlso Manual [F. Mahmoudi, 0808.3144]

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$2.63 \leq$	${\rm BR}(B \to X_s \gamma) \times 10^4$	$\leq 4.23$
0.71 <	${\rm BR}(B_u \to \tau \nu_{\tau}) \times 10^4$	< 2.57
1.3 <	$BR(B_s \to \mu^+ \mu^-) \times 10^9$	< 4.5
$-1.7 \times 10^{-2} <$	$\Delta_0(B\to K^*\gamma)$	$< 8.9 \times 10^{-2}$
0.56 <	$R_{ au u_{ au}}$	< 2.70
$2.9 \times 10^{-3} <$	${ m BR}(B  o D^0  au  u_{ au})$	$< 14.2 \times 10^{-3}$
0.151 <	$\xi_{D\ell u}$	< 0.681
	$BR(B_d \to \mu^+ \mu^-)$	$< 1.1 \times 10^{-9}$
0.6257 <	$\frac{\mathrm{BR}(K \to \mu\nu)}{\mathrm{BR}(\pi \to \mu\nu)}$	< 0.6459
Flavour physics $0.985 < $	$R_{\ell 23}$	< 1.013
$4.7 \times 10^{-2} <$	$BR(D_s \to \tau \nu_{\tau})$	$< 6.1 \times 10^{-2}$
$4.9 \times 10^{-3} <$	$BR(D_s \to \mu \nu_\mu)$	$< 6.7 \times 10^{-3}$
$3.0 \times 10^{-4} <$	${ m BR}(D  o \mu  u_{\mu})$	$< 4.6 \times 10^{-4}$
$-2.4 \times 10^{-10} <$	$\delta a_{\mu}$	$< 5.0 \times 10^{-9}$

[F. Mahmoudi, O. Stal [0907.1791]



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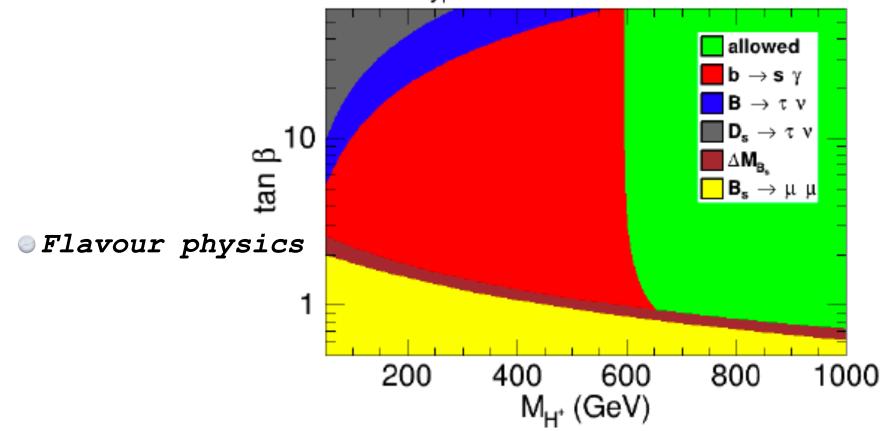


#### [A. Arbey, F. Mahmoudi, O. Stal, T. Stefaniak, [1706.07414]

HFLAV Coll., 1612.07233]  $\begin{array}{rl} 3.32 - 0.15 \leq & \mathrm{BR}(B \to X_s \gamma) \times 10^4 & \leq 3.32 + 0.15 \\ 1.06 \pm 0.19 \leq & \mathrm{BR}(B_u \to \tau^{\pm} \nu_{\tau}) \times 10^4 & \leq 1.06 + 0.19 \end{array}$ 

LHCb Coll., 1703.05747]  $3.0 - 0.85 \le BR(B_s \to \mu^+ \mu^-) \times 10^9 \le 3.0 + 0.85$ 

THDM Type II - Flavour constraints



Numerically scanning of the parameter space (trading  $\lambda_{1-5}$  for the physical Higgs boson masses as input parameters), with the following constraints

	$m_{_H} = 125 \text{ GeV}$
$m_h \; ({\rm GeV})$	10 - 80
$m_A \; ({\rm GeV})$	$10 - (M_Z - m_h)$
$m_{H^{\pm}} (\text{GeV})$	90 - 500
$\sin(eta - lpha)$	-1 - 1
$m_{12}^2 \; ({\rm GeV^2})$	$0 - m_A^2 \sin\beta\cos\beta$
aneta	2, 25

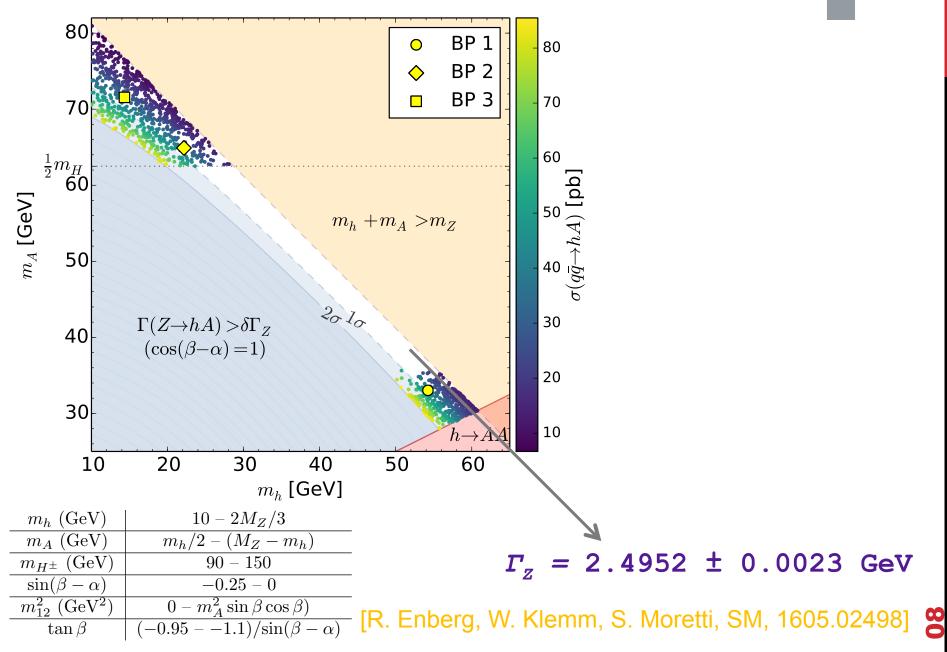
*imposed:* (Code: 2HDMC [D. Eriksson, J. Rathsman, O. Stal, 0902.0851])

• Unitarity, perturbativity and vacuum stability

- Oblique parameters S, T and U
- Flavour physics
- LEP, TeVatron and LHC results for
  - Additional Higgs bosons (HiggsBounds)
  - Measured Higgs signal strengths (HiggsSignals)



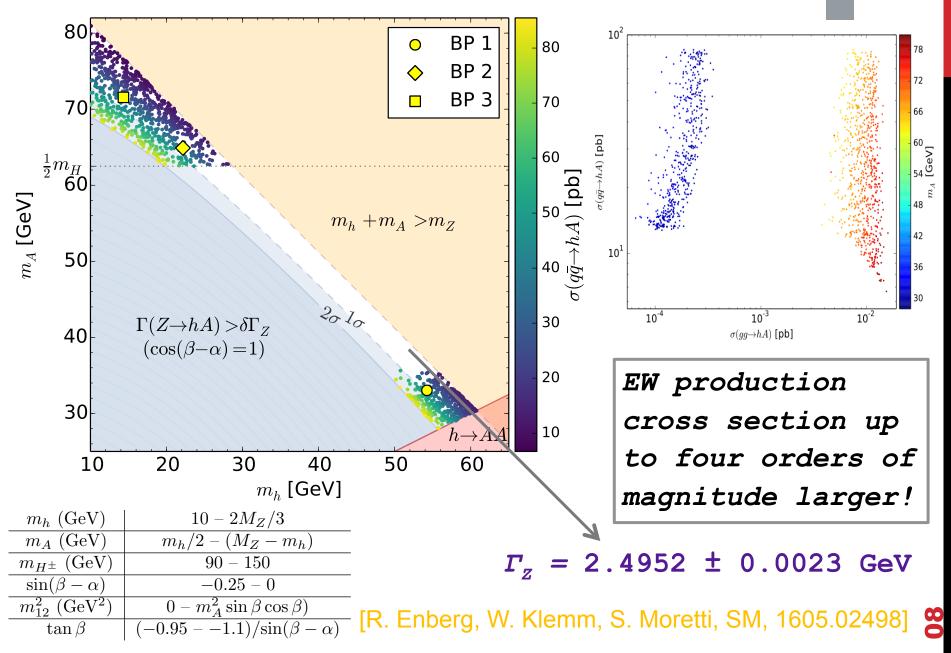
# m<sub>h</sub>+m<sub>A</sub>< m<sub>Z</sub> IN TYPE-I 2HDM



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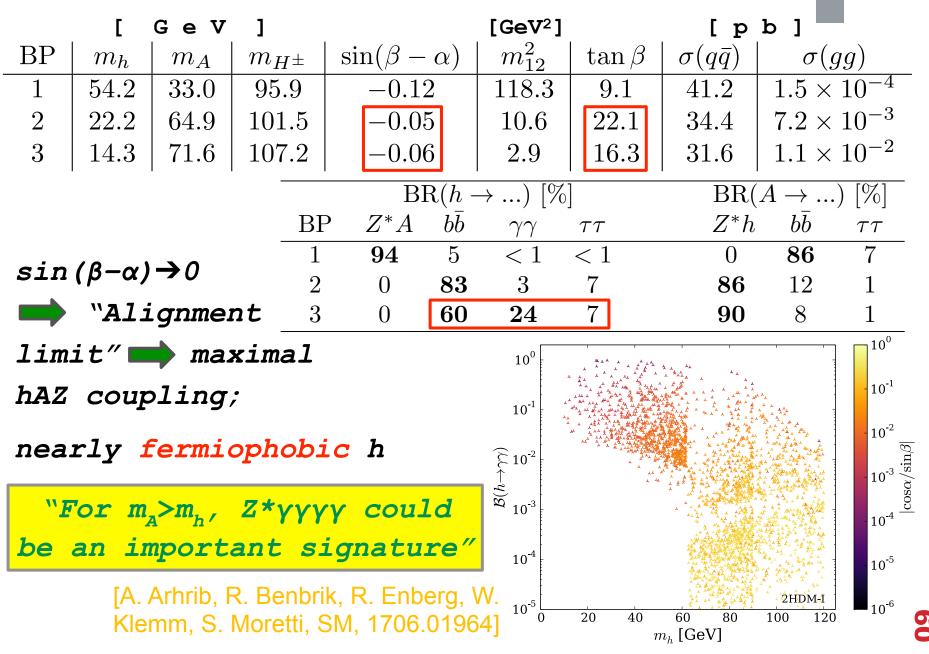
# m<sub>h</sub>+m<sub>A</sub>< m<sub>z</sub> IN TYPE-I 2HDM



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# **DOMINANT SEARCH CHANNELS**



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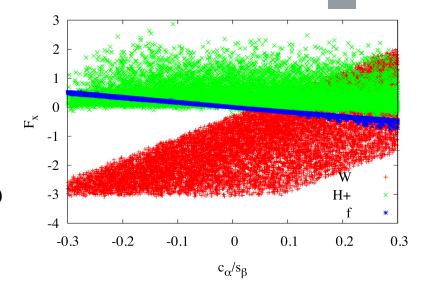
# LIGHT FERMIOPHOBIC HIGGS BOSON

$$F_{f} = \sum_{i} \frac{-2}{\tau_{f}^{2}} N_{f} Q_{f}^{2} \xi_{f}^{h} (\tau_{f} + (\tau_{f} - 1)I(\tau_{f})),$$

 $sin(\beta-\alpha) \rightarrow 0$ 

"Alignment

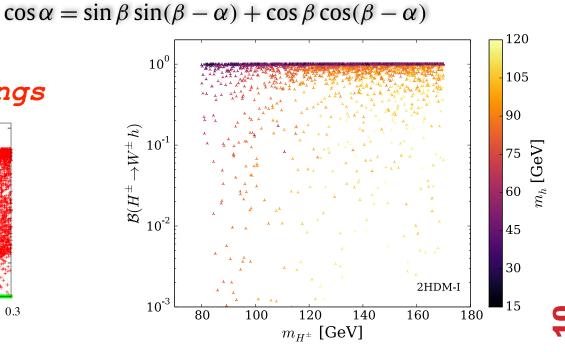
$$F_{H^{\pm}} = \frac{g_{hH^{\pm}H^{\mp}}}{\tau_{H^{\pm}}^{2}} \frac{m_{W}^{2}}{m_{H^{\pm}}^{2}} (\tau_{H^{\pm}} - I(\tau_{H^{\pm}})),$$
  
$$F_{W} = \frac{\sin(\beta - \alpha)}{\tau_{W}^{2}} (2\tau_{W}^{2} + 3\tau_{W} + 3(2\tau_{W} - 1)I(\tau_{W}))$$



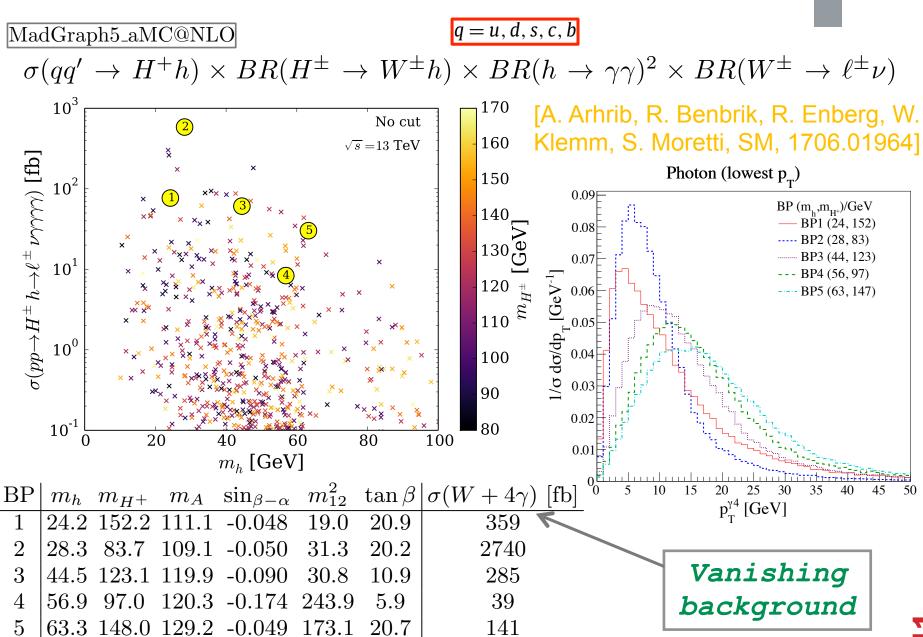
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limit" 페 maximal hAZ and  $hH^+W^-$  couplings bb 0.8  $Br(h \rightarrow xx)$ 0.6 0.4 0.2 0 -0.3 0.2 -0.2 -0.1 0 0.3 0.1 $c_{\alpha}/s_{\beta}$ 



# **DISCOVERY POTENTIAL**



l/σ dσ/dp<sub>m</sub> [GeV<sup>-1</sup>

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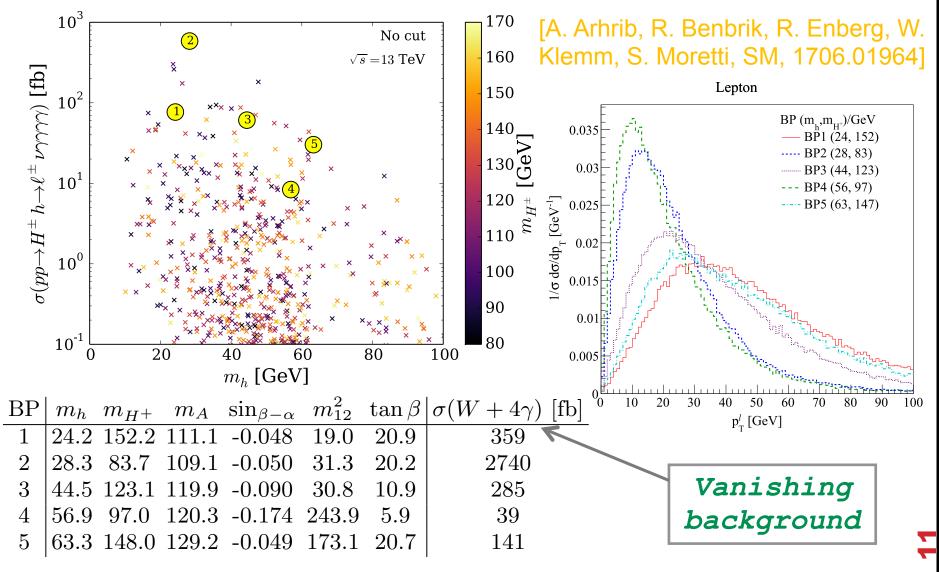
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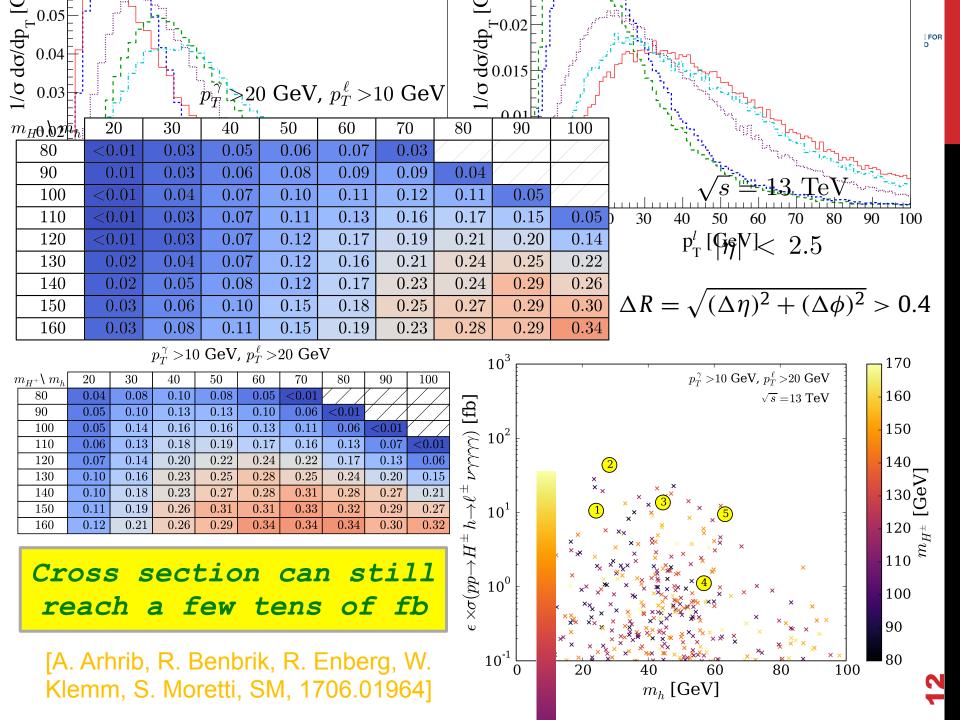
# **DISCOVERY POTENTIAL**



#### q=u,d,s,c,b

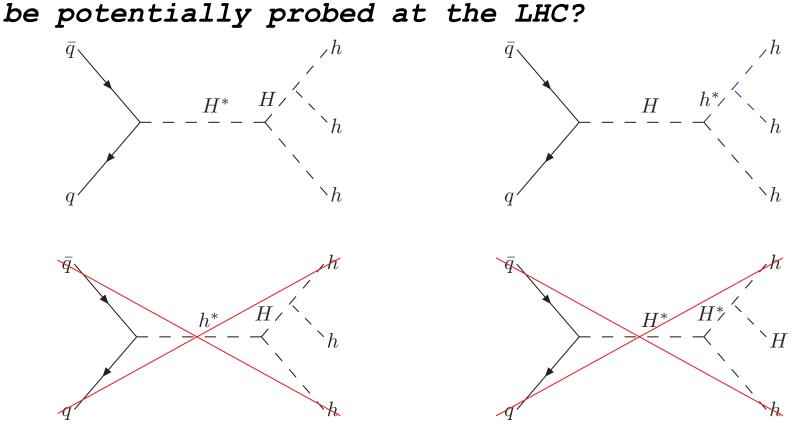
 $\sigma(qq' \to H^+h) \times BR(H^{\pm} \to W^{\pm}h) \times BR(h \to \gamma\gamma)^2 \times BR(W^{\pm} \to \ell^{\pm}\nu)$ 



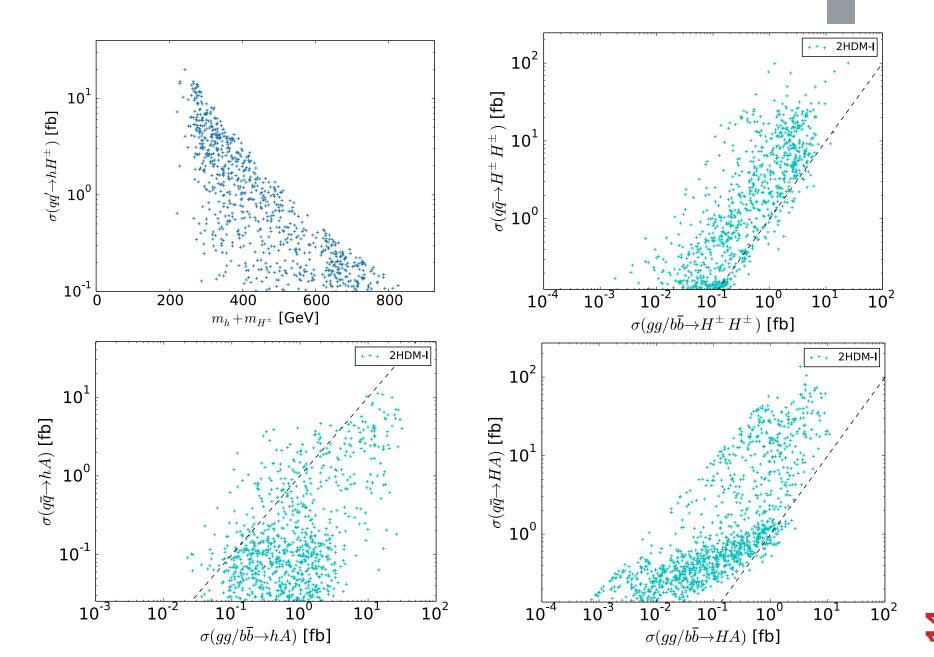


# **MULTI-HIGGS (EW) PRODUCTION**

Electroweak production of all possible 2-body and 3-body Higgs-Higgs/gauge states in the Type-I 2HDM © Can it dominate over QCD production? © Which Higgs-Higgs and Higgs-gauge couplings can



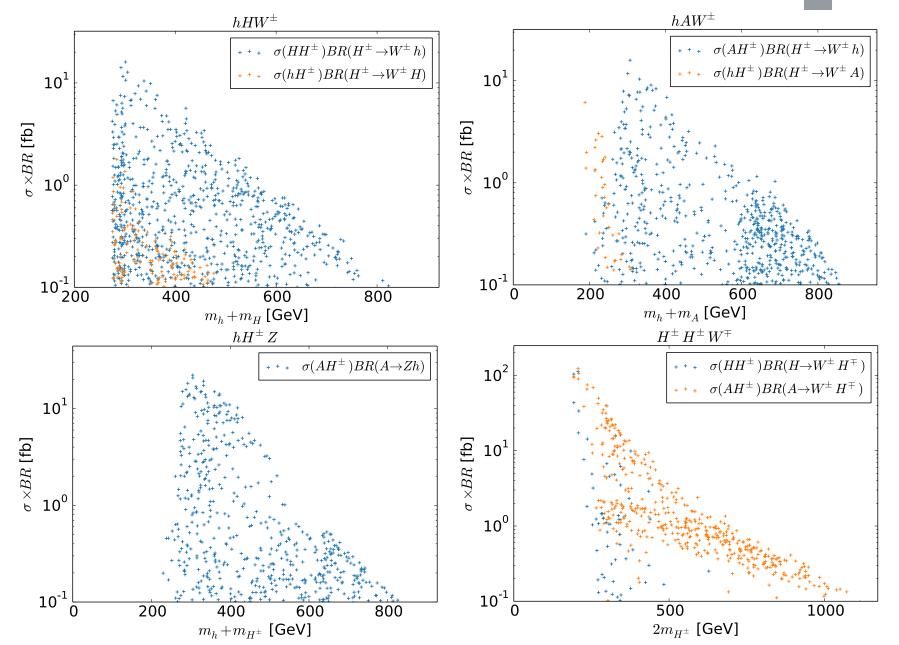
#### **2-BODY FINAL STATES**



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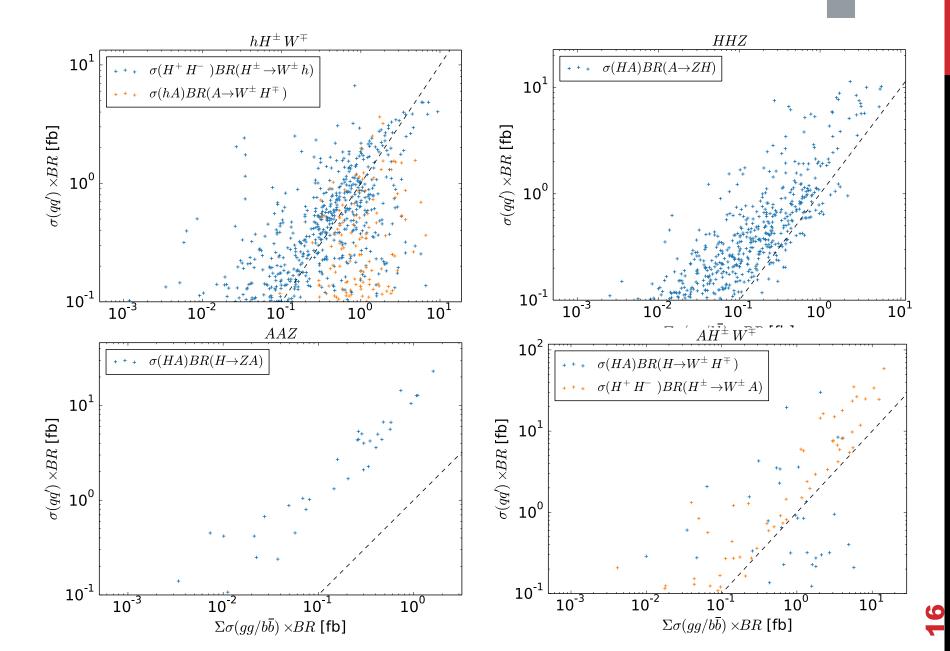
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# **3-BODY FINAL STATES**





# **3-BODY FINAL STATES - COMPARISON**



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# **HIGGS TRIPLE-COUPLINGS**

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2BFS1. hh $\checkmark$ (hhh) $\checkmark$ 2. HH(hhH) $\checkmark$ (3. AA(hhH) $\checkmark$ (4. H <sup>+</sup> H <sup>-</sup> 5. hH(hhH) $\checkmark$ (hhh)5. hH(hhH) $\checkmark$ (hhh) $\checkmark$ (6. hA(hhA)(7. hH <sup>±</sup> (hhH <sup>±</sup> )(8. HA(hhA)(	hHH (hHH) (hhH) hH <sup>+</sup> H <sup>-</sup> )	d. $hAA$ ( $hAA$ ) ( $hAA$ ) ( $hAA$ ) ( $hH^+H^-$ ) ( $HAA$ ) ( $HAA$ ) ( $AAA$ ) ( $AAH^\pm$ ) ( $hHA$ ) ( $hHA$ )	e. $hH^+H^-$ $(hH^+H^-)$ $\checkmark (hH^+H^-)$ $(HH^+H^-)$ $(AH^+H^-)$ $(hhH^{\pm})$ $(H^+H^-H^{\pm})$	f. HHH ✓ (HHH) (hHH) (HHA)	g. $HAA$ ( $HAA$ ) $\checkmark$ ( $HAA$ ) ( $hAA$ ) ( $hHA$ )	h. $HH^+H^-$ ( $HH^+H^-$ ) $\checkmark$ ( $HH^+H^-$ ) ( $hHH^\pm$ )
2. HH       (hhH) $\checkmark$ (         3. AA	$(hhH)$ $hH^+H^-)$	$\checkmark (hAA)$ $(hH^+H^-)$ $(HAA)$ $(HAA)$ $(AAA)$ $(AAH^{\pm})$	$(hH^+H^-)$ $(HH^+H^-)$ $(AH^+H^-)$ $(hhH^{\pm})$	(hHH)	$\checkmark (HAA)$ $(hAA)$ $(hHA)$	<ul> <li>✓ (HH<sup>+</sup>H<sup>−</sup>)</li> <li>(hHH<sup>±</sup>)</li> </ul>
3. $AA$	$(hhH)$ $hH^+H^-)$	$(hH^+H^-)$ $(HAA)$ $(HAA)$ $(AAA)$ $(AAH^{\pm})$	$(HH^+H^-)$ $(AH^+H^-)$ $(hhH^{\pm})$	(hHH)	$\checkmark (HAA)$ $(hAA)$ $(hHA)$	<ul> <li>✓ (HH<sup>+</sup>H<sup>−</sup>)</li> <li>(hHH<sup>±</sup>)</li> </ul>
$4. H^+H^ (hhH)$ $\checkmark$ $(hhh)$ $\checkmark$ $5. hH$ $(hhH)$ $\checkmark$ $(hhh)$ $\checkmark$ $6. hA$ $(hhA)$ $\checkmark$ $(hhA)$ $7. hH^{\pm}$ $(hhH^{\pm})$ $\bullet$ $8. HA$ $(hhA)$ $(hhA)$ $9. HH^{\pm}$ $(hhH^{\pm})$ $(hhH^{\pm})$ $10. AH^{\pm}$ $\bullet$	$hH^+H^-)$	$(hH^+H^-)$ $(HAA)$ $(HAA)$ $(AAA)$ $(AAH^{\pm})$	$(HH^+H^-)$ $(AH^+H^-)$ $(hhH^{\pm})$		(hAA) (hHA)	$(hHH^{\pm})$
$5.hH$ $(hhH)$ $\checkmark$ $(hhh)$ $\checkmark$ $6.hA$ $(hhA)$ (hhA) $7.hH^{\pm}$ $(hhH^{\pm})$ $8.HA$ $(hhA)$ ( $9.HH^{\pm}$ $(hhA)$ ( $10.AH^{\pm}$ (	$hH^+H^-)$	$(HAA)$ $(HAA)$ $(AAA)$ $(AAH^{\pm})$	$(HH^+H^-)$ $(AH^+H^-)$ $(hhH^{\pm})$		( <i>hHA</i> )	$(hHH^{\pm})$
$5.hH$ $(hhH)$ $\checkmark$ $(hhh)$ $\checkmark$ $6.hA$ $(hhA)$ $(hhA)$ $(hhA)$ $7.hH^{\pm}$ $(hhH^{\pm})$ $(hhH)$ $(hhA)$ $8.HA$ $(hhA)$ $(hhA)$ $(hhA)$ $9.HH^{\pm}$ $(hhH^{\pm})$ $(hhH^{\pm})$ $10.AH^{\pm}$ $(hhH)$ $(hhH)$	$hH^+H^-)$	$(hhA)$ $(AAA)$ $(AAH^{\pm})$	$(AH^+H^-)$ $(hhH^{\pm})$		( <i>hHA</i> )	
$7.hH^{\pm}$ $(hhH^{\pm})$ $8.HA$ $(hhA)$ $9.HH^{\pm}$ $(hhH^{\pm})$ $10.AH^{\pm}$	(hHA)	<ul> <li>✓ (AAA)</li> <li>(AAH<sup>±</sup>)</li> </ul>	$(hhH^{\pm})$		( <i>HHA</i> )	
8. $HA$ ( $hhA$ )(9. $HH^{\pm}$ ( $hhH^{\pm}$ )( $hhH^{\pm}$ )10. $AH^{\pm}$ ( $hhH^{\pm}$ )( $hhH^{\pm}$ )	(hHA)			(HHA)		
9. $HH^{\pm}$ ( $hhH^{\pm}$ ) ( $h$	(hHA)	(hHA)		(HHA)	(HHA)	
10. <i>AH</i> <sup>±</sup>				(11111)	$\checkmark (AAA)$	$(AH^+H^-)$
	$(hHH^{\pm})$		$(hHH^{\pm})$	$(HHH^{\pm})$	$(AAH^{\pm})$	$\checkmark \begin{array}{c} (HHH^{\pm}) \\ (H^{+}H^{-}H^{\pm}) \end{array}$
11.hZ $(hhZ)$			$(hAH^{\pm})$		$(HAH^{\pm})$	$(HAH^{\pm})$
		(AAZ)	$(H^+H^-Z)$			
$12. hW^{\pm} \qquad (hhW^{\pm})$		$(AAW^{\pm})$	$(H^+H^-W^\pm)$			
13. HZ  (hhZ)  (	(hHZ)			(HHZ)	(AAZ)	$(H^+H^-Z)$
$14. HW^{\pm} \qquad \qquad (hhW^{\pm}) \qquad (h$	$(hHW^{\pm})$			$(HHW^{\pm})$	$(AAW^{\pm})$	$(H^+H^-W^\pm)$
15. AZ		(hAZ)			(HAZ)	
16. $AW^{\pm}$		$(hAW^{\pm})$			$(HAW^{\pm})$	
$17. H^{\pm}Z$						
$18. H^+W^-$			$(hH^+W^-)$			

[R. Enberg, W. Klemm, S. Moretti, SM, 1809.XXXXX]

#### [R. Enberg, W. Klemm, S. Moretti, SM, 1809.XXXXX]

Coupling 2BFS	$\mathrm{m.}hAZ$	n. $HAZ$	o. $H^+H^-Z$	p. $hH^+W^-$	q. $HH^+W^-$	$r. AH^+W^-$	s. $hZZ$	t. <i>HZZ</i>	u. $hW^+W^-$	v. $HW^+W^-$
1. <i>hh</i>	(hAZ)			$(hH^+W^-)$			(hZZ)		$(hW^+W^-)$	
2. HH		(HAZ)			$(HH^+W^-)$			(HZZ)		$(HW^+W^-)$
3.AA	(hAZ)	(HAZ)				$(AH^+W^-)$				
$4. H^+ H^-$			$\checkmark(H^+H^-Z)$	$(hH^+W^-)$	$(HH^+W^-)$	$(AH^+W^-)$				
5. hH	(HAZ)	(hAZ)		$(HH^+W^-)$	$(hH^+W^-)$		(HZZ)	(hZZ)	$(HW^+W^-)$	$(hW^+W^-)$
6. hA	$\checkmark  \begin{array}{c} (hhZ) \\ (AAZ) \end{array}$	(hHZ)		$(hH^+W^-)$ $(AH^+W^-)$			(AZZ)		$(AW^+W^-)$	
$7.hH^{\pm}$	$(AH^{\pm}Z)$			$\checkmark \frac{(hhW^{\pm})}{(H^+H^-W^{\pm})}$	$(hHW^{\pm})$	$(hAW^{\pm})$	$(H^{\pm}ZZ)$		$(H^{\pm}W^{+}W^{-})$	
8. HA	(hHZ)	$\checkmark \begin{pmatrix} HHZ \\ (AAZ \end{pmatrix}$			$(HH^+W^-)$ $(AH^+W^-)$			(AZZ)		$(AW^+W^-)$
$9.HH^{\pm}$		$(AH^{\pm}Z)$		$(hHW^{\pm})$	$\checkmark \frac{(HHW^{\pm})}{(H^+H^-W^{\pm})}$	$(HAW^{\pm})$		$(H^{\pm}ZZ)$		$(H^{\pm}W^{+}W^{-})$
$10. AH^{\pm}$	$(hH^{\pm}Z)$	$(HH^{\pm}Z)$		$(hAW^{\pm})$	$(HAW^{\pm})$	$\checkmark \frac{(AAW^{\pm})}{(H^+H^-W^{\pm})}$				
11. hZ	$\checkmark \begin{array}{c} (hhA) \\ (AZZ) \end{array}$	(hHA)		$(H^+ZW^-)$			$\checkmark$ ( <i>hhZ</i> )	(hHZ)		
$12. hW^{\pm}$				$(hhH^{\pm})$ $\checkmark$ $(hHH^{\pm})$ $(H^{\pm}W^{+}W^{-})$		$(hAH^{\pm})$			✓ $(hhW^{\pm})$	$(hHW^{\pm})$
13. <i>HZ</i>	(hHA)	$\checkmark \frac{(HHA)}{(AZZ)}$			$(H^+ZW^-)$			$\checkmark$ ( <i>HHZ</i> )		
14. <i>HW</i> <sup>±</sup>				$(hHH^{\pm})$	$\checkmark \frac{(HHH^{\pm})}{(H^{\pm}W^{+}W^{-})}$	$(HAH^{\pm})$			$(hHW^{\pm})$	$\checkmark$ (HHW <sup>±</sup> )
15. <i>AZ</i>	$\checkmark \begin{array}{c} (hAA) \\ (hZZ) \end{array}$	$\checkmark \begin{array}{c} (HAA) \\ (HZZ) \end{array}$				$(H^+ZW^-)$	(hAZ)	(HAZ)		
16. $AW^{\pm}$				$(hAH^{\pm})$	$(HAH^{\pm})$	$\checkmark (H^{\pm}W^{+}W^{-})$			$(hAW^{\pm})$	$(HAW^{\pm})$
$17. H^{\pm}Z$	$(hAH^{\pm})$	$(HAH^{\pm})$	√				$(hH^{\pm}Z)$	$(HH^{\pm}Z$		
$18. H^+W^-$				$\checkmark \begin{array}{c} (hH^+H^-) \\ (hW^+W^-) \end{array}$	$\checkmark \begin{array}{c} (HH^+H^-) \\ (HW^+W^-) \end{array}$	$\checkmark \begin{array}{c} (AH^+H^-) \\ (AW^+W^-) \end{array}$			$(hH^+W^-)$	$(HH^+W^-)$
			Г				Morott		1000 V	

# **HIGGS-GAUGE COUPLINGS**



# **CONCLUSIONS**

- KIAS S ROPEA INSTITUTE FOR ADVANCED STUDY
- Additional Higgs bosons are predicted in most new physics frameworks - can be lighter or heavier than 125 GeV
- Even when light, they are difficult to detect at the LHC in the conventional channels, owing to generally reduced couplings to the SM
- Their pair-production can provide crucial probes
- In the Type-I 2HDM, a light scalar-pseudoscalar pair as well as a light H<sup>±</sup> could be accessible in multi-photon final states
- EW pair-production essential when a charged Higgs boson is involved - can dominate over QCD even for certain neutral Higgs boson combinations



# THANK YOU! 감사합니다!