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Measurement of Z bosons produced in association with jets via vector boson fusion at 13 TeV with the ATLAS detector

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On behalf of the ATLAS Collaboration

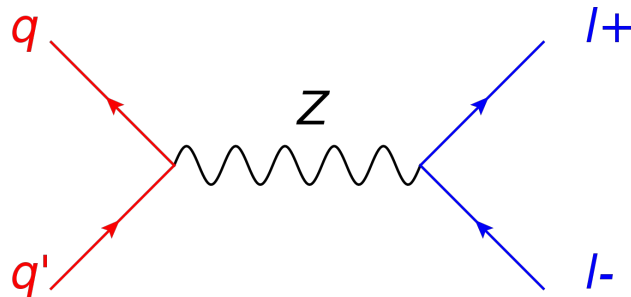


Canadian Association
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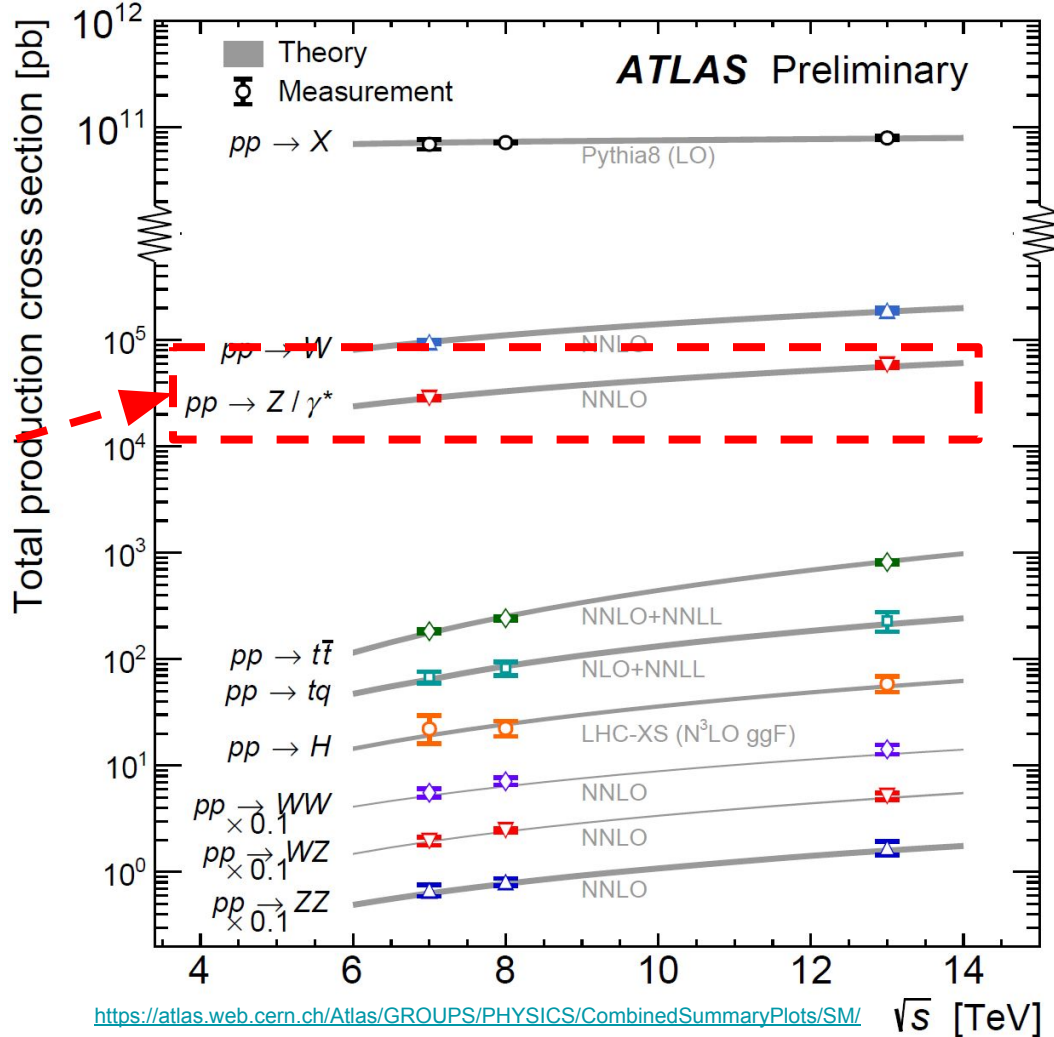
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Motivation

Drell-Yan Z production

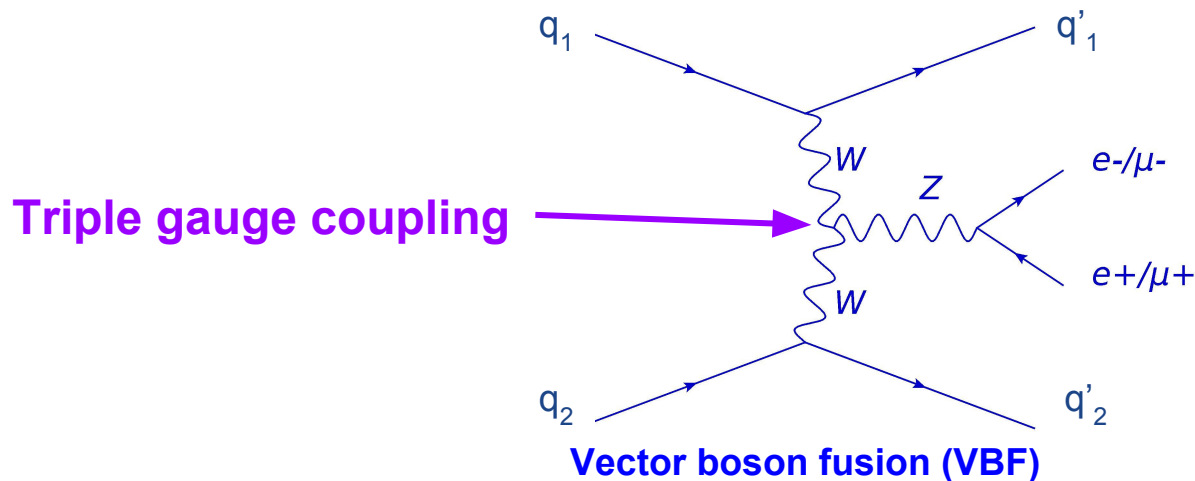


- Well measured leptonic final states with large statistics
- Background for many precision measurements (Higgs, top, diboson) and new physics searches



Signal: Electroweak Z + dijets

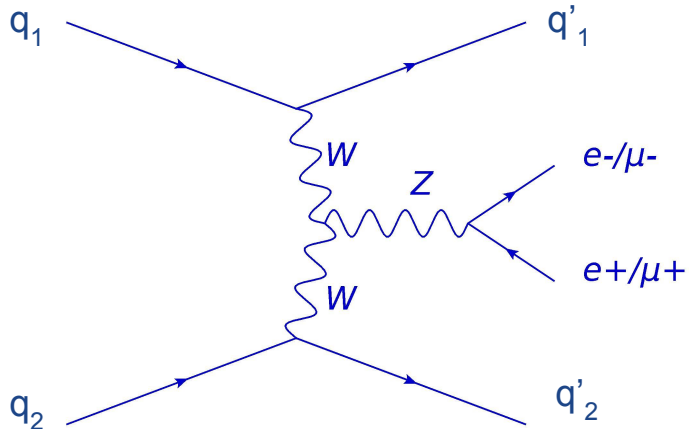
EW Z+dijets includes all processes where there is a **t-channel** exchange of a **W/Z boson** and a **l^+l^-jj** final state



- **Drell-Yan Z+dijets** is produced frequently in pp collisions compared to **EW Z+dijets** (Large Background!)
- VBF Z is a probe for new physics via higher order corrections to the WWZ vertex (**the triple gauge coupling**)

Z+dijets production

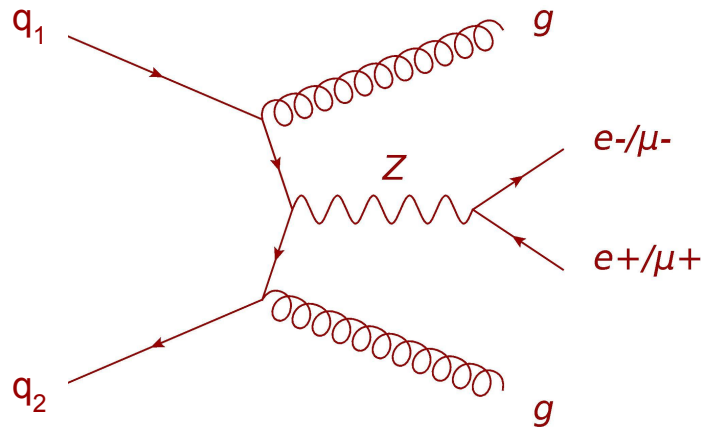
EW Zjj has a **much smaller cross section** compared to the **Drell-Yan** process



SM Prediction: Powheg+Pythia

$$\sigma_{\text{fid}} = 125.2 \pm 3.4 \text{ fb}$$

Phys. Lett. B 775 (2017) 206



SM Prediction: CT14nnlo PDF

$$\sigma_{\text{fid}} = 0.74 \pm 0.03(\text{PDF}) \pm 0.01(\text{scale}) \pm 0.01(\text{other}) \text{ nb}$$

Phys. Lett. B 759 (2016) 601

*A factor of **6000** larger*

Extracting the **EW** signal from the dominant **Drell-Yan** background is challenging
Modeling of the background is crucial

Measurement: Cross section

SR = Signal Region

Drell-Yan Zjj prediction
~99% of total bkg

ttbar, dibosons ...
~1% of total bkg

EW Z+dijets cross
section of bin i

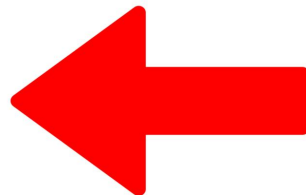
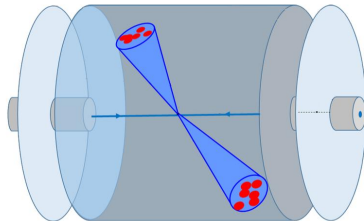
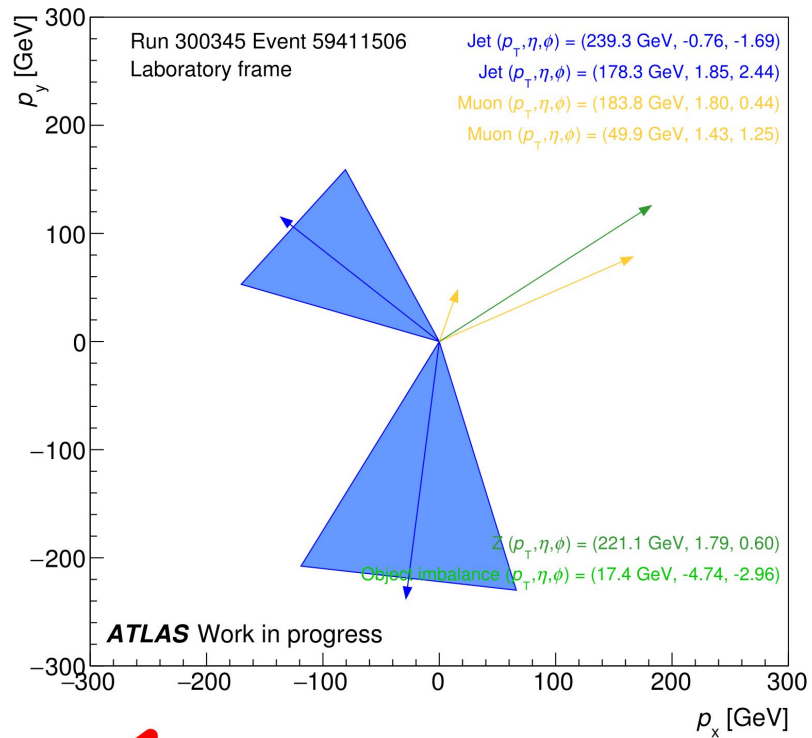
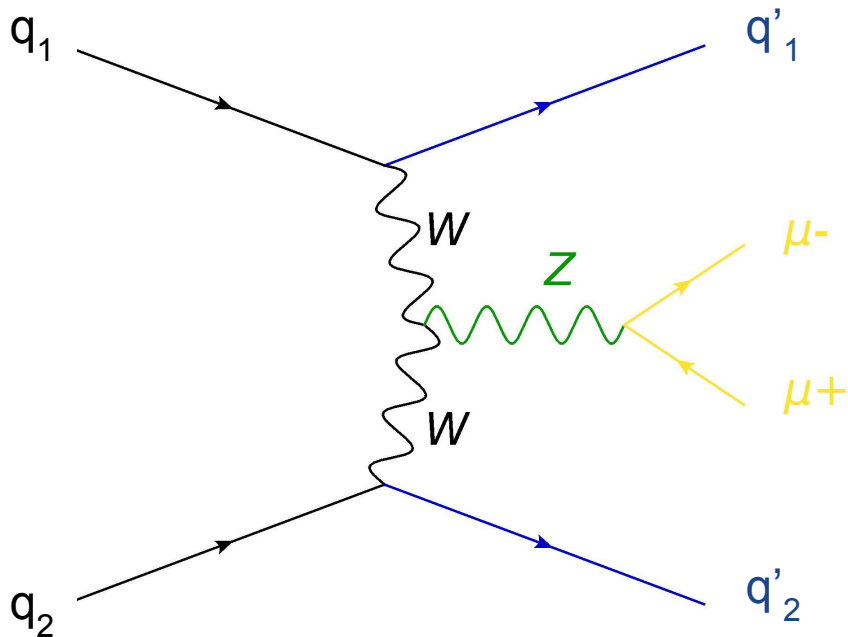
$$\sigma_{\text{fid},i} = \frac{N_{\text{SR},i}^{\text{data}} - N_{\text{SR},i}^{\text{strong}} - N_{\text{SR},i}^{\text{non-Z}}}{C_i \mathcal{L}}$$

Bin-by-bin correction
factor for reconstruction
inefficiency

Integrated Luminosity

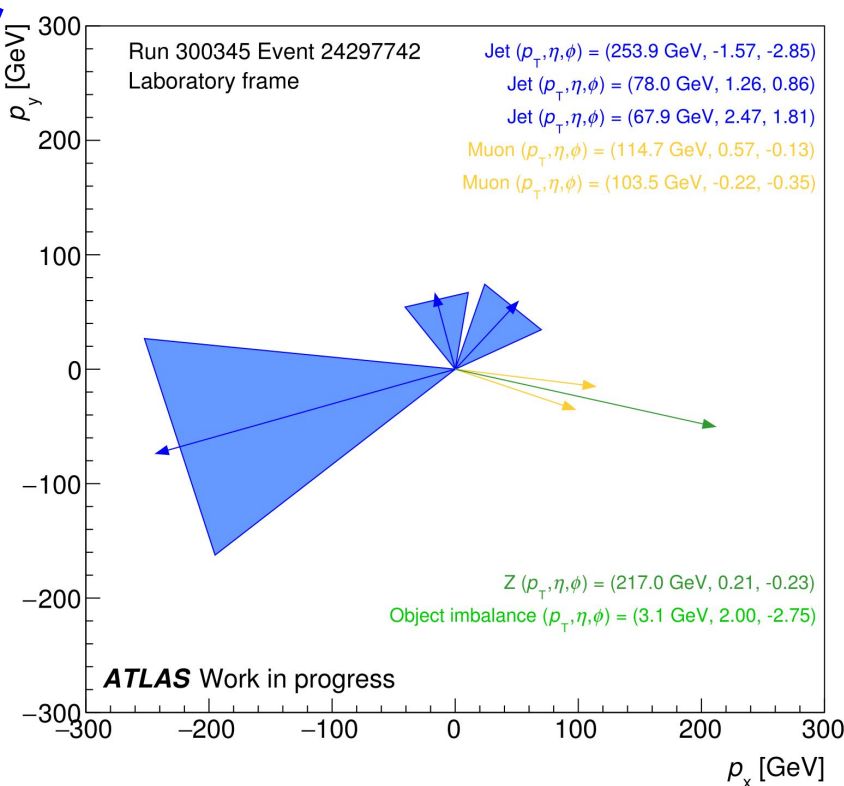
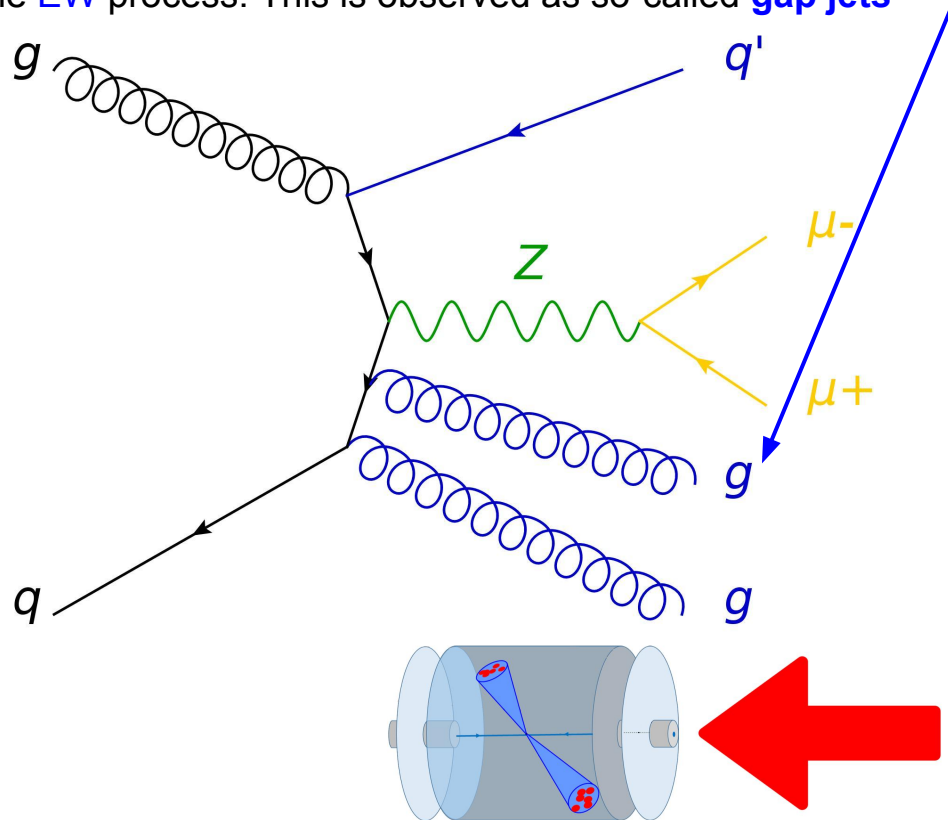
- The **Drell-Yan Z+dijets** accounts for the vast majority of events
- Crucial to understand this process to extract the **EW Z+dijets signal**

EW Z+dijets: What we see with the ATLAS detector



DY Z+dijets: What we see with the ATLAS detector

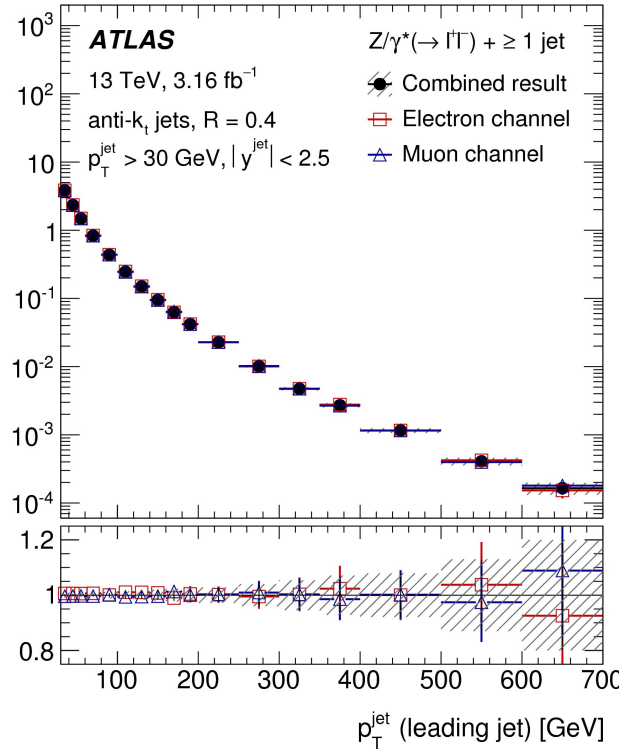
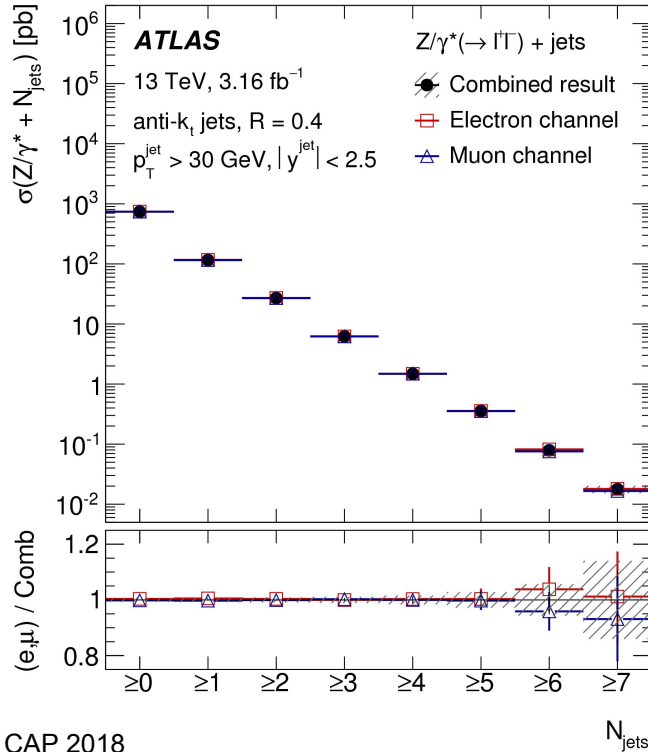
The **Drell-Yan** process is more likely to have additional hadronic activity between the two leading jets than the **EW** process. This is observed as so-called **gap jets**



Z+jets @ 13 TeV

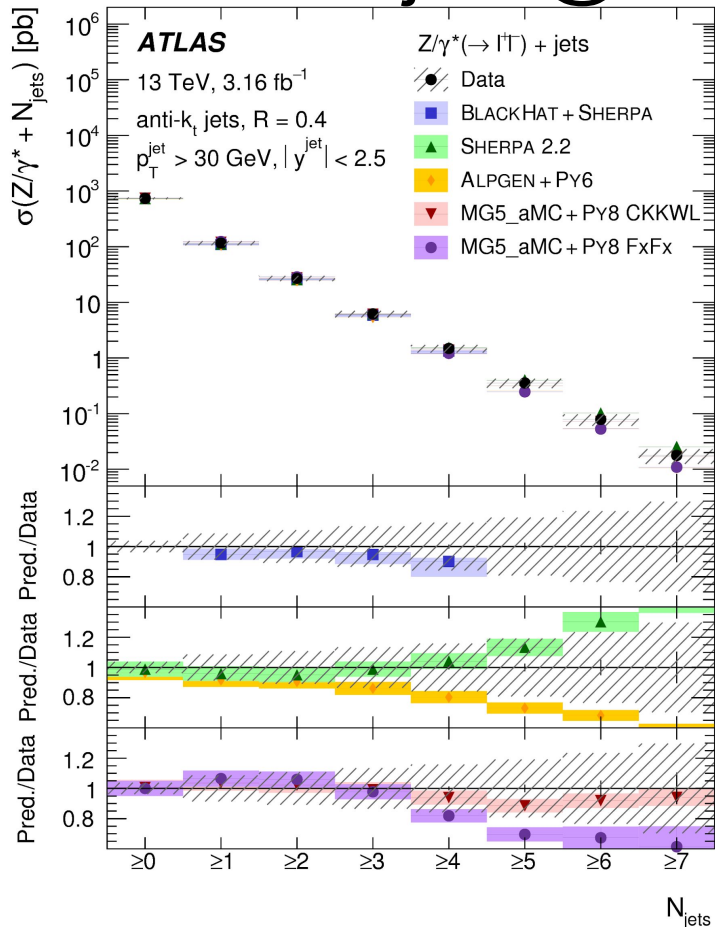
Z+jets @ 13 TeV Eur. Phys. J. C 77 (2017) 361

- Using the 2015 dataset, 3.16 fb^{-1}
- Cross section measured differentially as a function of **characteristic variables**
 - N_{jets} , jet p_{T} (shown) + others



- e^+e^- and $\mu^+\mu^-$ measurements are compatible
- Combination improves stat. uncertainty by $\sim 30\%$ for $N_{\text{jets}} \geq 4$ or greater

Z+jets @ 13 TeV Eur. Phys. J. C 77 (2017) 361



This differential measurement can be used to test various **theory predictions** provided by different **MC generators**

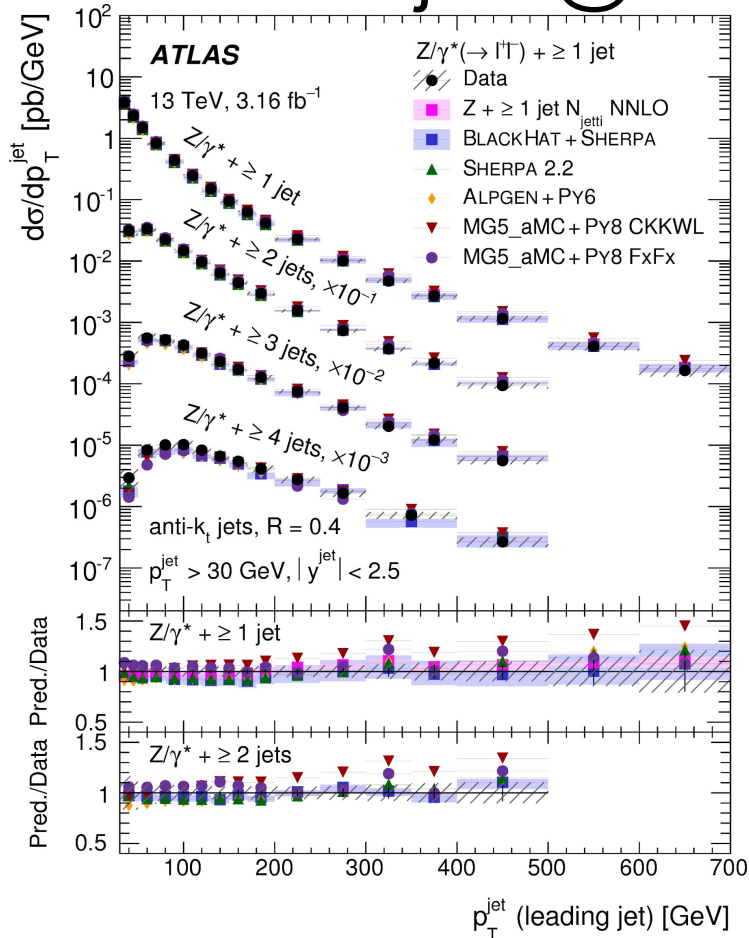
Cross section at **high jet multiplicity** provides sensitivity to differences in **MC generators**

Differential cross section as a function of jet multiplicity

- **BlackHat+Sherpa** is a fixed order NLO calculation of Z with ≤4 partons
- **Sherpa 2.2** is NLO for ≤2 partons + LO for ≤4 partons+PS
- **Alpgen+Py6** is LO for ≤5 partons+PS
- **MG5_aMC+Py8CKKWL** is LO for ≤4 partons+PS
- **MG5_aMC+Py8FxFx** is NLO for ≤2 partons+PS

Observe that predictions vary substantially once the **number of jets** exceeds the number of **partons** included in the matrix element calculation

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Differential cross section as a function of jet p_T

- Additional generator $Z+\geq 1$ jet N_{jetti} NNLO is a fixed order **NNLO** prediction
 - Precise agreement with data for $Z+\geq 1$ jet
- Jet p_T is a fundamental observable and a probe for perturbative QCD over a wide range of scales
- The kinematics of jets in events with Z bosons is essential to model the backgrounds of other SM processes and BSM searches
- Observe that the LO generator **MG5_aMC+Py8CKKWL** models a too hard p_T spectrum
- The other NLO generators and fixed order predictions model the spectrum well over the full range of p_T

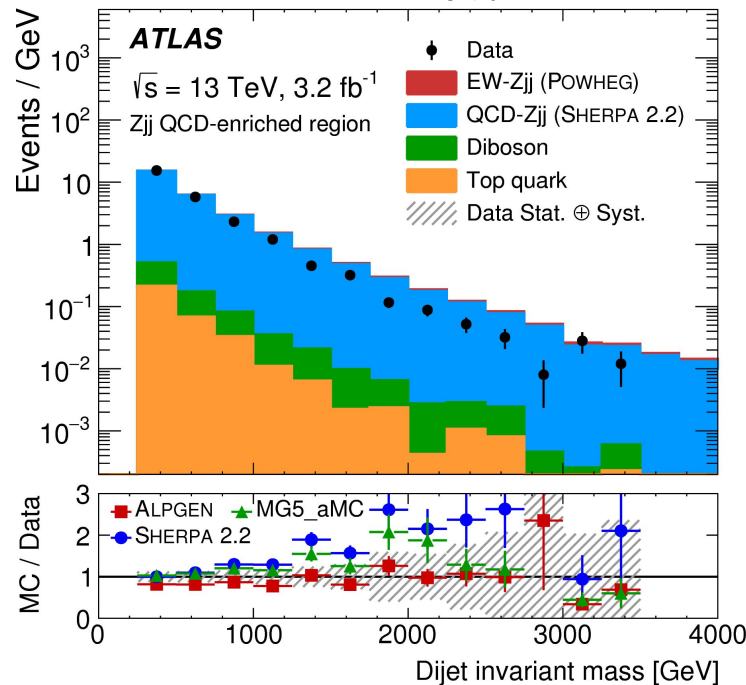
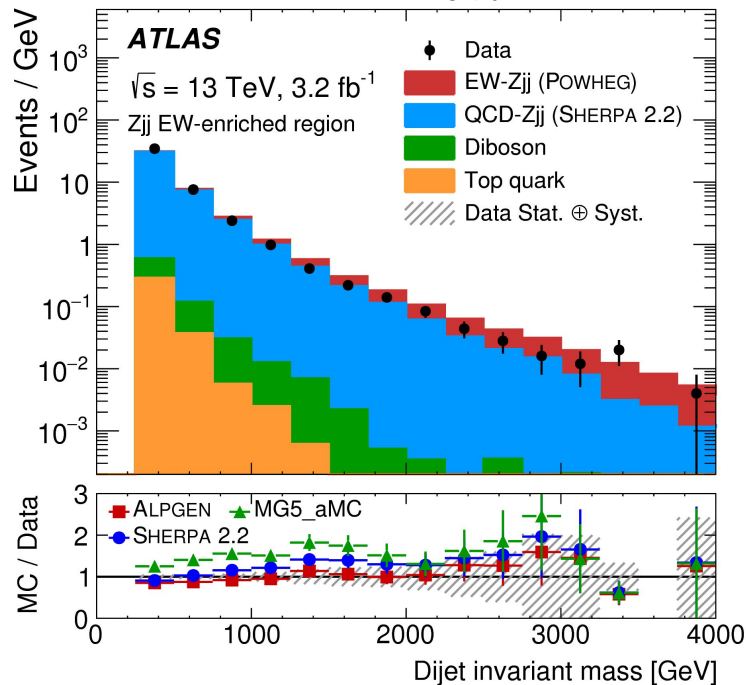
Electroweak Z+dijets @ 13 TeV

Electroweak Z+dijets @ 13 TeV Phys. Lett. B 775 (2017) 206

Define **two regions** to model background and extract signal:

EW enhanced: $N_{\text{gap jets}} = 0$

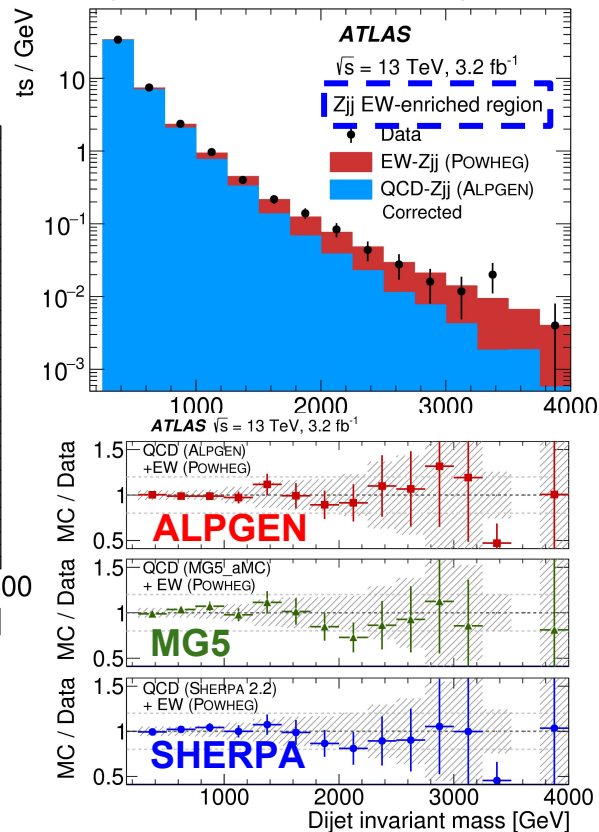
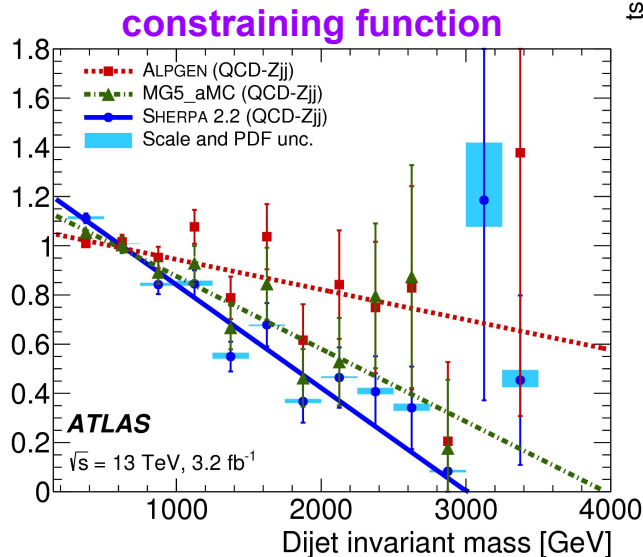
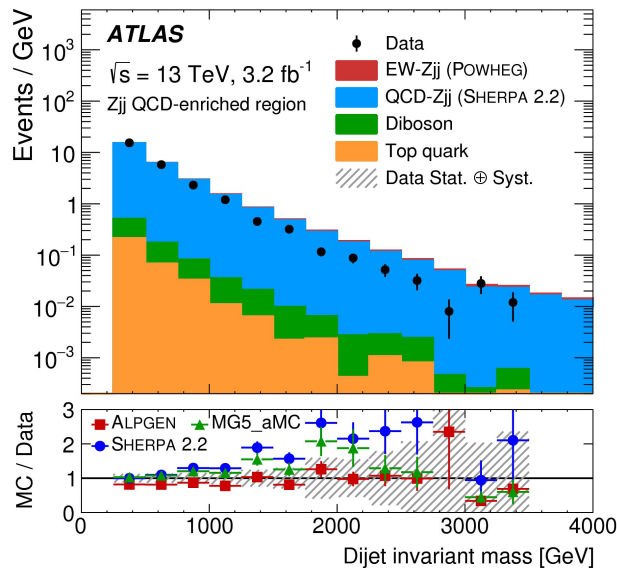
DY enhanced: $N_{\text{gap jets}} \geq 1$



- All 3 **Drell-Yan** MC generators tested **mismodel** Zjj cross section at very high m_{jj}
- Differential m_{jj} prediction constrained using data-driven approach

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- Derive **constraining function** for each MC generator by constraining **Drell-Yan** Zjj to the data in the **QCD-enhanced** region
- Apply the **constraining function** in the **EW-enhanced** region to correctly model the **DY background**
- Subtract the **DY background** to extract the **EW signal**

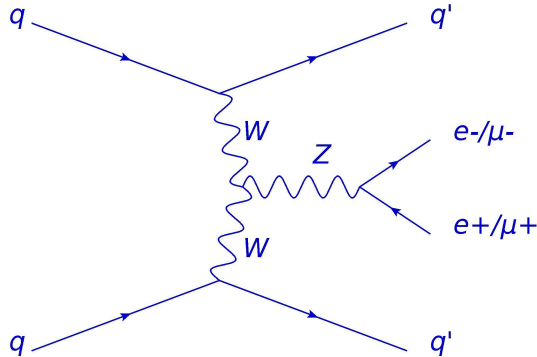


Cross section extracted in **EW-enhanced** region:

$$\sigma_{\text{fid}} = 119 \pm 16(\text{stat}) \pm 20(\text{syst}) \pm 2(\text{lumi}) \text{ fb}$$

Summary

- Differential cross section of Z bosons produced in association with jets
 - Comparing different MC generators at **high jet multiplicity** to observe differences in theory predictions
- Electroweak Z+jets measurement at 13 TeV
 - Extract EW signal from QCD background by applying **data-driven constraint**
 - With more statistics it will be possible to extract a differential EW cross section and search for new physics in the **VVZ** vertex (**anomalous triple gauge coupling**)



Future Outlook:

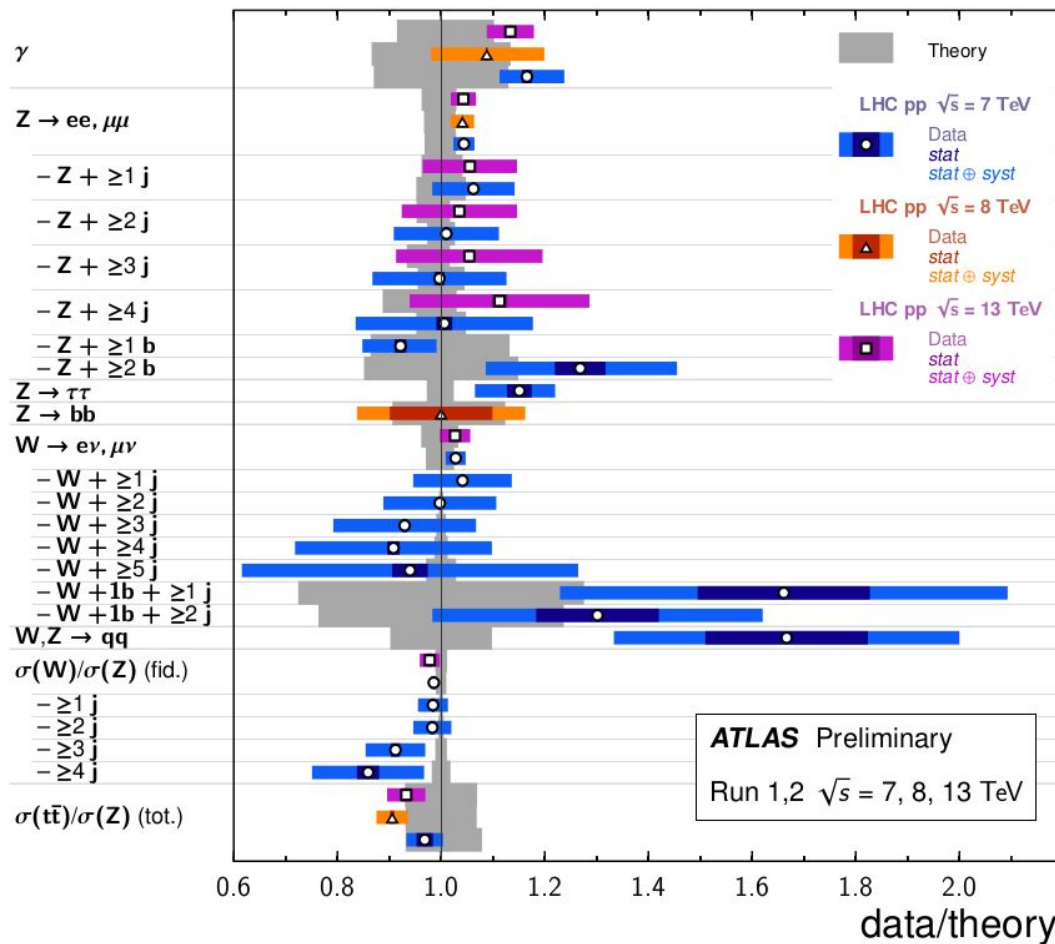
- Work is in progress to measure the **EW Z+dijets** cross section differentially
- Will study the full Run 2 dataset (2015-2018)

$$\mathcal{L} \sim 150 \text{ fb}^{-1}$$

BACKUP

Vector Boson + X fid. Cross Section Measurements

Status: March 2018



The effective Lagrangian, \mathcal{L} , for aTGCs can be written as

$$\frac{\mathcal{L}}{g_{WWZ}} = i \left[g_{1,Z} \left(W_{\mu\nu}^\dagger W^\mu Z^\nu - W_{\mu\nu} W^{\dagger\mu} Z^\nu \right) + \kappa_Z W_\mu^\dagger W_\nu Z^{\mu\nu} + \frac{\lambda_Z}{m_W^2} W_{\rho\mu}^\dagger W_\nu^\mu Z^{\nu\rho} \right] \quad (8.2)$$

if only those terms that conserve charge conjugation and parity are retained from the general expression [68]. Here, $g_{WWZ} = -e \cot \theta_W$, e is the electric charge, θ_W is the weak mixing angle, W^μ and Z^μ are the W -boson and Z -boson fields, $X_{\mu\nu} = \partial_\mu X_\nu - \partial_\nu X_\mu$ for $X = W$ or Z , and $g_{1,Z}$, κ_Z and λ_Z are dimensionless couplings. The SM values of these dimensionless couplings are $g_{1,Z}^{SM} = 1$, $\kappa_Z^{SM} = 1$ and $\lambda_Z^{SM} = 0$.

The tree-level S -matrix for this effective Lagrangian violates unitarity at large energy scales. Unitarity is restored in the full theory by propagator (form factor) effects. A typical approach is to modify the couplings by a dipole form factor

$$a(\hat{s}) = \frac{a_0}{(1 + \hat{s}/\Lambda^2)^2} \quad (8.3)$$

Table 7. The 95% confidence intervals obtained on the aTGC parameters from counting the number of events with $m_{jj} > 1$ TeV in the *search* region. Observed and expected intervals, labelled ‘obs’ and ‘exp’ respectively, are presented for unitarisation scales of $\Lambda = 6$ TeV and $\Lambda = \infty$. The parameter $\Delta g_{1,Z}$ refers to the deviation of $g_{1,Z}$ from the SM value.

aTGC	$\Lambda = 6$ TeV (obs)	$\Lambda = 6$ TeV (exp)	$\Lambda = \infty$ (obs)	$\Lambda = \infty$ (exp)
$\Delta g_{1,Z}$	$[-0.65, 0.33]$	$[-0.58, 0.27]$	$[-0.50, 0.26]$	$[-0.45, 0.22]$
λ_Z	$[-0.22, 0.19]$	$[-0.19, 0.16]$	$[-0.15, 0.13]$	$[-0.14, 0.11]$