





Probes of hadronization (resonances, (hyper)nuclei, charm baryons, coalescence)

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Evolution of a heavy ion collision



Thermal model fit to LHC data



- Hadron abundances in HI (yields and ratios) can be successfully interpreted in terms of production at chemical equilibrium
- Statistical (thermal) models predict the yields of any particle at chemical freeze-out once that T_{ch} and μ_B are known

At the LHC ($\mu_B = 0$) particle yields of light flavor hadrons (including nuclei) are described within the thermal model with a common chemical freeze-out temperature ($T_{chem} = 156 \pm 2 \text{ MeV}$)

K* not included in the fit

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ALICE Collaboration, arXiv:1710.07531, NPA 971, 1 (2018)

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Resonances

Thermal model fit to ALICE data



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• At the LHC ($\mu_B = 0$) particle yields of light flavor hadrons (including nuclei) are described within the thermal model with a common chemical freeze-out temperature ($T_{chem} = 156 \pm 2 \text{ MeV}$)

K*(892) production not described (yield suppressed)

ALICE Collaboration, arXiv:1710.07531, NPA 971, 1 (2018)

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Hadronic phase



- The yield of resonances produced in heavy ion collisions is expected to be in agreement with T_{ch} chemical freeze-out equilibrium
- Although, can be altered by hadronic interactions between chemical and kinetic freezeouts:
 - rescattering: daughter particles undergo elastic scattering or pseudo-elastic scattering through a different resonance
 → parent particle is not reconstructed → loss of signal
 - regeneration: pseudo-elastic scattering of decay products (πK → K^{*0}, KK → φ, etc.) → increased yields
- Effect of hadronic processes depends on:
 - Ifetime and density of <u>hadronic phase</u>
 - resonance lifetime and scattering cross sections

Light flavoured hadronic resonances



- Resonances have lifetimes comparable to that of the fireball produced in heavy-ion collisions:
 - can be used to study properties and lifetime of the late hadronic phase
- Resonances differ by mass and quark content:
 - insights on the multiplicity-dependent enhancement of strangeness production
 - anomalous baryon-to-meson ratios at intermediate transverse momentum parton energy loss

ALICE Collaboration, arXiv:1805.04365

	ρ(770) ⁰	K*(892) ⁰	۸(1520)	ф(1020)
ст(fm/ <i>c</i>)	1.3	4.2	12.6	46.2

ρ(770)⁰/π[±]

• The ratio of p_{τ} integrated yields divided by particle with similar quark content ρ/π shows clear suppression going from pp and peripheral Pb-Pb collisions to central Pb-Pb



	ρ (770) ⁰	K*(892) ⁰	۸(1520)	ф(1020)
cτ(fm/ <i>c</i>)	1.3	4.2	12.6	46.2

ALICE Collaboration, Phys. Rev. C 91 (2015) 024609

- K*(892)⁰/K[±]
 - K*(892)/K ratio decreases as function of multiplicity

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Suggests that re-scattering is dominant over regeneration



	ρ (770) ⁰	K*(892) ⁰	٨(1520)	ф(1020)
cτ(fm/ <i>c</i>)	1.3	4.2	12.6	46.2

ALICE Collaboration, Phys.Rev. C99 (2019) 024905

STAR Collaboration, Phys.Rev.C78 (2008) 044906 STAR Collaboration, Phys. Rev. Lett. 97 (2006) 132301

- Λ(1520)/Λ
 - A(1520)/A ratio clearly decreases as function of multiplicity in Pb-Pb
 - Clear suppression compared to the thermal model predictions

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ALICE Collaboration, Phys. Rev. C 91 (2015) 024609

- φ(1020)/K[±]
 - φ(1020)/K ratio independent of collision system (energy)
 - The ratio $\phi(1020)/K$ increases at large $\langle dN_{ch}/d\eta \rangle$ because of strangeness enhancement



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ALICE Preliminary

 \Diamond pp s = 7 TeV

□ Pb-Pb \ *s*_{NN} = 5.02 TeV

凸 Xe-Xe V S_{NN} = 5.44 TeV

ALICE • pp \s = 2.76 TeV

• pp s = 7 TeV

■ Pb-Pb \ s_{NN} = 2.76 TeV

STAR

- ★ pp \s = 200 GeV
- ☆ Au-Au \ s_{NN} = 200 GeV

Hadronic Resonances:

- Suppression of ρ^0 , K^{*0}, $\Lambda(1520)$, while ϕ not suppressed
- Qualitative description is obtained with EPOS+UrQMD
 - Hadronic phase (UrQMD) important in FPOS to describe data
- Consistent results for Xe-Xe and Pb-Pb at similar multiplicity
- Indication of the existence of a hadronic phase

ALT-PREL-161421



Light (hyper-)nuclei

Deuteron p_{τ} spectra in Pb-Pb collisions

ALICE-PUBLIC-2017-006



Spectra are extracted in several centrality bins and fitted with blast-wave function for the extraction of yields

E. Schnedermann et al., Phys. Rev. C 48, 2462 (1993)

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⁴He production in Pb-Pb collisions



- Heaviest anti-nucleus observed : 16 candidates in Pb-Pb at 5.02 TeV
- Pre-selection using d*E*/d*x* measured in TPC
- Selection: $\pm 3\sigma$ from the expected value for ⁴He
- Signal extraction from mass squared distribution obtained using TOF

Nuclei production in Pb – Pb collisions



• Thermal model prediction: exponential dependence of the yield

$$\frac{dN}{dy} \propto \exp\left(-\frac{m}{T_{chem}}\right)$$

• The density ratio of a particle with the next heavier one:

$$\frac{n_i}{n_{i+1}} \approx \exp\left(-\frac{\Delta m}{T_{chem}}\right)$$

Thermal model expectation

$$p/d \sim exp(-B) \sim 330^{+70}_{-61}$$

Experimental result

ALICE Collaboration, arXiv:1710.07531, NPA 971, 1 (2018)

Hypernuclei

- A hypernucleus is a nucleus which contains at least one hyperon (a baryon containing one or more strange quarks) in addition to nucleons
- Main goals of hypernuclear physics:
 - Extension of nuclear chart
 - Understand the baryon-baryon interaction in strangeness sector
 - Study the structure of multi-strange systems

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- $^{3}_{\Lambda}$ H is the lightest known hypernucleus and is formed by (p,n, Λ).
 - Mass = $2.991 \, \text{GeV}/c^2$
 - $B_{\Lambda} = 0.13 \pm 0.05 \text{ MeV} (B_{d} = 2.2 \text{ MeV}, B_{t} = 8.5 \text{ MeV}, B_{3He} = 7.7 \text{ MeV})$

 $\binom{3}{\Lambda}\overline{H} \xrightarrow{3}{\Lambda}H$ is unstable under weak decay and branching ratios are not well known \rightarrow Only few theoretical calculations available [1]

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Thermal model comparison



- ³_{\lambda} H/³He ratio compared with different thermal models:
 - Extracted yield is in good agreement with equilibrium thermal model prediction for $T_{chem} = 156$ MeV, such as GSI-Heidelberg model [1] even if B_{Λ} is << T_{ch}

ALICE Collaboration Phys. Lett. B 754 (2016) 360-372

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[1]A. Andronic et al., Phys. Lett. B 697, 203 (2011)

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Thermal model fit



- Opposite to resonances, yields of (hyper)nuclei are described by the common chemical freeze-out temperature (T_{chem} = 156 ± 2 MeV)
- The binding energy of (hyper)nuclei is very small and it is surprising that they do not immediately dissociate in the hadronic phase

•
$$T_{ch}$$
 (160 MeV) < $T_{hadronic}$ < T_{kin} (100 MeV)



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FIG. 1. Schematic for the production of a deuteron in the final state of a relativistic collision between two heavy nuclei.

- Model originally developed to describe <u>light-nuclei production</u> (deuteron, triton...)
 - If baryons at freeze-out are close enough in phase space and match spin state a (anti-)nucleus can be formed
 - Usually, since the nucleus is larger w.r.t. the source, the phase space is reduced to the momentum space
 - Assuming that p an n have the same mass and have the same p_{T} spectra, the yield of any nucleus can be determined as $d^{3}N_{A} = \left(\frac{d^{3}N_{P}}{d^{3}N_{P}} \right)^{A} = \left(\frac{d\pi}{d^{2}N_{P}} \right)^{A}$

$$E_A \frac{\mathrm{d}^3 N_A}{\mathrm{d} p_A^3} = B_A \left(E_\mathrm{p} \frac{\mathrm{d}^3 N_\mathrm{p}}{\mathrm{d} p_\mathrm{p}^3} \right)^A \qquad B_A = \left(\frac{4\pi}{3} p_0^3 \right)^{(A-1)} \frac{1}{A!} \frac{M}{m^A}$$

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 - Assuming that p an n have the same mass and have the same p_{T} spectra, the yield of any nucleus can be determined as $E_{A} \frac{d^{3}N_{A}}{dn^{3}} \neq B_{A} \left(E_{p} \frac{d^{3}N_{p}}{dn^{3}}\right)^{A} \qquad B_{A} = \left(\frac{4\pi}{3}p_{0}^{3}\right)^{(A-1)} \frac{1}{A!} \frac{M}{m^{A}}$

Measured nucleus p_{T} -spectra

Measured proton p_{T} -spectra

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 - Assuming that p an n have the same mass and have the same p_{T} spectra, the yield of any nucleus can be determined as $d^{3}N_{4} = \left(d^{3}N_{7} \right)^{A} = \left(d^{2}N_{7} \right)^{A}$



- Model originally developed to describe light-nuclei production (deuteron, triton...)
 - If baryons at freeze-out are close enough in phase space and match spin state a (anti-)nucleus can be formed
 - Usually, since the nucleus is larger w.r.t. the source, the phase space is reduced to the momentum space
 - Assuming that p an n have the same mass and have the same $p_{_{\rm T}}$ spectra, the yield of any



Coalescence parameter B_{γ}



Simple coalescence model

• Flat B_2 vs p_T and no dependence on

multiplicity/centrality

 Approximately observed in "small systems": pp, p-Pb and peripheral Pb-Pb

Coalescence parameter B_{2}



F.Bellini and A. P.Kalweit, arXiv:1807.05894 [hep-ph]. R. Scheibl, U. Heinz, PRC 59 (1999) 1585-1602 K. Blum et al., PRD 96 (2017) 103021 Simple coalescence model

- Flat B_2 vs p_T and no dependence on multiplicity/centrality
 - Approximately observed in "small systems": pp, p-Pb and peripheral Pb-Pb

More elaborate coalescence model takes into account the volume of the source:

$$B_2 = \frac{3\pi^{3/2} \langle C_d \rangle}{2m_T R^3(m_T)}$$

• B_2 scales like HBT radii (R)

- decrease with centrality in Pb-Pb is explained as an increase in the source volume
- ▷ increase with p_{T} in central Pb-Pb reflects the k_{T} -dependence of the homogeneity volume (i.e. volume with similar flow properties) in HBT

✓ Qualitative agreement in central Pb-Pb collisions



Light nuclei production: Deuteron to proton ratio



- d/p increases with multiplicity going from pp to peripheral Pb-Pb : consistent with simple coalescence (d \propto p²)
- No significant centrality dependence in Pb-Pb : consistent with thermal model (yield fixed by T_{chem})
- How the two models are connected is not yet fully understood
 - Is there a single particle production mechanism?

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Charm Baryon

Heavy quarks in heavy ion collisions

- Charm and beauty quarks are produced in parton hard scatterings in the initial phase of the heavy-ion collision (production time of cc(bb) pair at rest
- Flavour is conserved in strong interactions \rightarrow Transported through the full system evolution
- What can be tested?
 - In-medium energy loss: colour-charge and quark-mass dependence
 - Heavy quark participation in the collective expansion, thermalisation in the medium
 - Modification of the hadronization mechanisms in the medium





- ALICE measured Λ_{c} in Pb-Pb collisions at 5.02 TeV for 0-80% centrality class
 - 2.5 σ hint of Λ_c / D^o ratio enhanced in Pb-Pb collisions w.r.t. pp and p-Pb collisions
 - Coalescence production mechanism at play

ALICE Collaboration arXiv:1809.10922 DOI: 10.1016/j.physletb.2019.04.046





- ALICE measured Λ_{c} in Pb-Pb collisions at 5.02 TeV for 0-80% centrality class
 - ~2.5 σ hint of Λ_c / D^o ratio enhanced in Pb-Pb collisions w.r.t. pp and p-Pb collisions
 - ~2.0 σ hint of larger R_{AA} of Λ_c than D mesons in 0-10% centrality class
 - Charm quark hadronization via coalescence:
 - hierarchy of the R_{AA} $\Lambda_c > D_s > Non-strange D-meson > pions$



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- CMS measured Λ_c^+ production in 2 centrality intervals (0-30%) and (30-50%) in the p_{τ} interval 10-20 GeV/*c*:
 - The R_{AA} for Λ_c^+ show a hint of suppressed production of Λ_c^+ for $p_T > 10$ GeV/c (0-100%), but no conclusion can be drawn due to the large uncertainty in the pp differential cross section.



- The Λ_c / D^o ratio in Pb-Pb collisions, is consistent with the result from pp collisions:
 - Result in contrast w.r.t ALICE observation of a large enhancement in the Λ_c / D^o ratio in the p_{τ} range of 6-12 GeV/c
 - > may suggest that there is no significant contribution from the coalescence process for $p_{T} > 10 \text{ GeV}/c$ in Pb-Pb collisions

Conclusions

- Thermal models are able to describe the production yields of light flavor hadrons and (hyper)nuclei within a common chemical freeze-out temperature
- Resonance results support the existence of a hadronic phase in central heavy-ion collisions that lasts long enough to cause a significant reduction of the reconstructed yields of short lived resonances: surprisingly nuclei seem to survive thru this phase
- Light nuclei measurement reveals system size dependence of hadronization:
 - Evolution of B_A and d/p ratio vs multiplicity: is there a single particle production mechanism?
- Charmed-baryon Λ_c is less suppressed than D mesons: favored production mechanism is quark **coalescence** for $p_T < 10 \text{ GeV}/c$, for higher momenta this has not been observed
- Beautiful picture, but there is still a lot to do and to understand!
 - New and more precise data can be expected from the LHC on the presented topics in the next years

Extra Material

Production of light (anti-)(hyper)nuclei

After the Long Shutdown 2 data will be collected with better performance at higher luminosity

- Expected integrated luminosity: ~10 nb^{-1} (~ 8x10⁹ collisions in the 0-10% centrality class)
 - Precision test of coalescence / thermal production models
 - Sensitive to size ratio of the object and the source
 - Search for rarely produced anti- and hypermatter: Insights on the strength of the hyperon-nucleon interaction, relevant for nuclear physics and neutron stars.





Hadronization of Heavy Quark



- Charm quark number is conserved in strong interaction
- Hadronization chemistry: crucial for the interpretation of the heavy flavor hadron spectra
- Run 3+4 data will allow the first comprehensive survey of this effect
- D_s , B_s and Λ_c spectra from low to high p_T in Pb-Pb collisions
- Provide the necessary statistical accuracy to see the emergence of the effect at low $p_{_{
 m T}}$

Identification of nuclei

Low momenta: specific energy loss in the TPC

- Nuclei identification via d*E*/d*x* measurement in the TPC:
 - Excellent separation of (anti-)nuclei from other particles over a wide range of momenta



ALICE Collaboration, Phys. Rev. C 93, 024917 (2016)

Higher momenta: time-of-flight measurement in the TOF

- Velocity measurement with the Time Of Flight detector is used to evaluate the m² distribution
- Excellent TOF performance: $\sigma_{TOF} \approx 85$ ps in Pb-Pb collisions



ALICE Collaboration, Int. J. Mod. Phys. A 29 (2014) 1430044

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"More elaborate" coalescence model

- For "large" systems, the size of the emitting volume (V_{eff}) has to be taken into account:
 - the larger the distance between the protons and neutrons which are created in the collision, the less likely it is that they coalesce
- The source can be parameterized as rapidly expanding under radial flow (hydro)
- The coalescence process is governed by the same correlation volume ("length of homogeneity") which can be extracted from HBT interferometry
- The source radius enters in the $B_{\rm A}$ and in the quantummechanical correction $\langle C_{\rm A} \rangle$ factor that accounts for the size of the object being produced (d, ³He, ...)

$$B_A = \frac{2J_A + 1}{2^A} A \left\langle \mathcal{C}_A \right\rangle \frac{V_{\text{eff}}(A, M_t)}{V_{\text{eff}}(1, m_t)} \left(\frac{(2\pi)^3}{m_t V_{\text{eff}}(1, m_t)} \right)^{A-1}$$

R. Scheibl, U. Heinz, PRC 59 (1999) 1585-1602 K. Blum et al., PRD 96 (2017) 103021

F.Bellini and A. P.Kalweit, arXiv:1807.05894 [hep-ph].



Good description of the data

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Centrality of the collisions



Centrality = degree of overlap of the 2 colliding nuclei

Central collisions:

- small impact parameter b
- high number of participant nucleons \rightarrow high multiplicity

Peripheral collisions:

- large impact parameter b
- low number of participant nucleons \rightarrow low multiplicity

Centrality connected to observables via Glauber model



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Centrality of the collisions: p-Pb and pp

Multiplicity estimator: slices in VZERO-A (VOA) amplitude



ALI-PERF-51387

Central collision



Peripheral collision



Correlation between impact parameter and multiplicity is not as straight-forward as in Pb-Pb

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Statistical thermal model

- Thermodynamic approach to particle production in heavy-ion collision: all the particles are produced at chemical freeze-out
- Starting point: Grand Canonical partition function (Z) for a relativistic ideal quantum gas of hadrons of particle type i (i = pion, proton,... \rightarrow full PDG)
- Thermal model can predict also the yields of any particle at chemical freeze-out
- Exponential dependence of the particle yield:

$${dN\over dy} \propto e^{\left(-{m\over T_{_{chem}}}
ight)}$$

- The thermal model predicts an exponential decrease of particle yields with increasing mass at a given temperature
- The density ratio of a particle with the next heavier one:

$$\frac{n_i}{n_{i+1}} \approx \exp\left(-\frac{\Delta m}{T}\right)$$

Thermal model



Nature 561 (2018) no.7723, 321-330 arXiv:1710.09425 [nucl-th]

Statistical hadronization model: thermal emission from equilibrated source

Particle abundances fixed at chemical freeze-out

$$N_i = \frac{g_i V}{2\pi^2} \int_0^{+\infty} \frac{p^2 dp}{\exp\left[-\left(\frac{E-\mu_B}{T_{\rm chem}}\right)\right] \pm 1}$$

• Primordial yields modified by hadron decays:

- Contribution obtained from calculations based on known hadron spectrum
- Excellent agreement with data with only 2 free parameters: $\rm T_{\rm chem}$, V

Coalescence parameter B_3



 B_3 of $(\bar{t})t$ and $({}^{3}He){}^{3}He$ measured in pp and Pb-Pb collisions First ever measurements of the B_3 of \bar{t} and ${}^{3}He$ in pp collisions Increasing trend with p_{τ} and centrality observed in Pb-Pb collision

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Hypertriton $(^{3}_{\Lambda}H)$

p n ${}^{3}_{\Lambda}$ H is the lightest known hypernucleus and is formed by (p,n, Λ).

- Mass = $2.991 \, \text{GeV}/c^2$
- **B**_A = 0.13 ± 0.05 MeV (B_d = 2.2 MeV, B_t = 8.5 MeV, B_{3He} = 7.7 MeV)

 $({}^{3}_{\Lambda}\overline{H}){}^{3}_{\Lambda}H$ is unstable under weak decay. Possible decay modes:

$${}^{3}_{\Lambda}H \rightarrow {}^{3}He + \pi^{-} \quad (\sim 25\%)$$

$${}^{3}_{\Lambda}H \rightarrow {}^{3}H + \pi^{0} \quad (\sim 13\%)$$

$${}^{3}_{\Lambda}H \rightarrow d + p + \pi^{-} \quad (\sim 41\%)$$

$${}^{3}_{\Lambda}H \rightarrow d + n + \pi^{0} \quad (\sim 21\%)$$

- Branching ratios are not well known
 - Only few theoretical calculations[1] available

[1]Kamada et al., Phys. Rev. C57(1998)4

Δ



$${}^{3}_{\overline{\Lambda}}\overline{H} \rightarrow {}^{3}\overline{He} + \pi^{+} \quad (^{2}5\%)$$

$${}^{3}_{\overline{\Lambda}}\overline{H} \rightarrow {}^{3}\overline{H} + \pi^{0} \quad (^{1}3\%)$$

$${}^{3}_{\overline{\Lambda}}\overline{H} \rightarrow \overline{d} + \overline{p} + \pi^{+} \quad (^{2}41\%)$$

$${}^{3}_{\overline{\Lambda}}\overline{H} \rightarrow \overline{d} + \overline{n} + \pi^{0} \quad (^{2}21\%)$$