

Prompt photons from gluon fusion/splitting induced by magnetic fields in heavy-ion collisions

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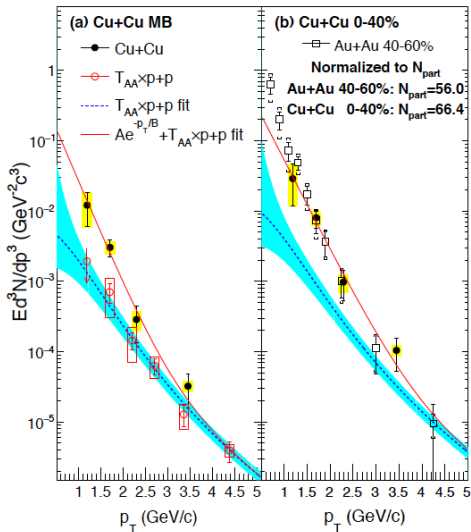
Instituto de Ciencias Nucleares, UNAM

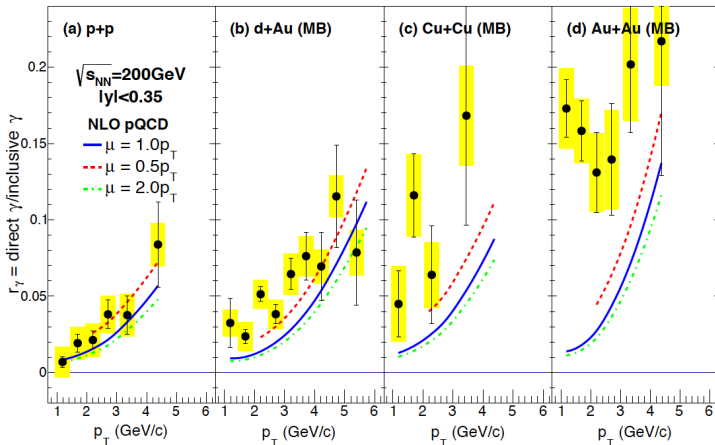
AA, J.D. Castaño-Yepes, I. Domínguez, J. Salinas, M.E. Tejeda-Yeomans, *arXiv:1904.02938*

May, 2019

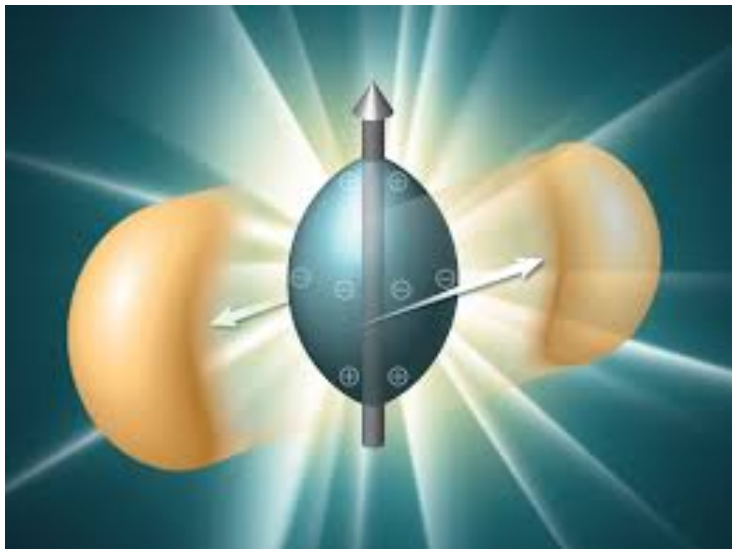


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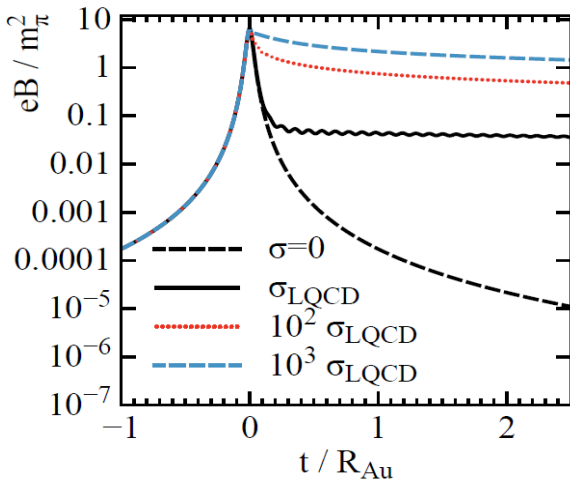




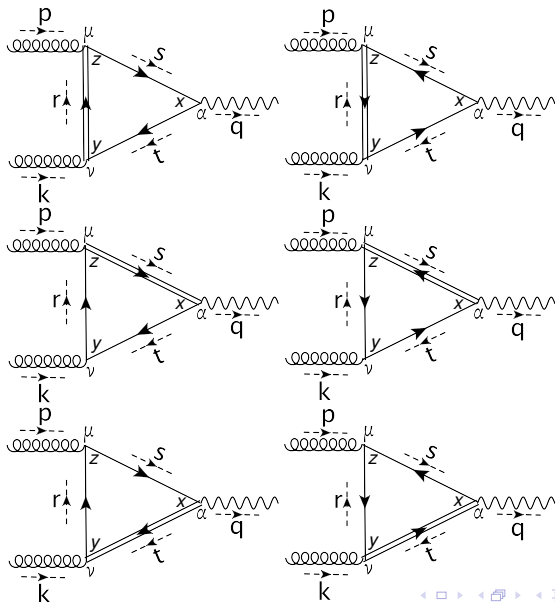
Magnetic field generated in peripheral HICs



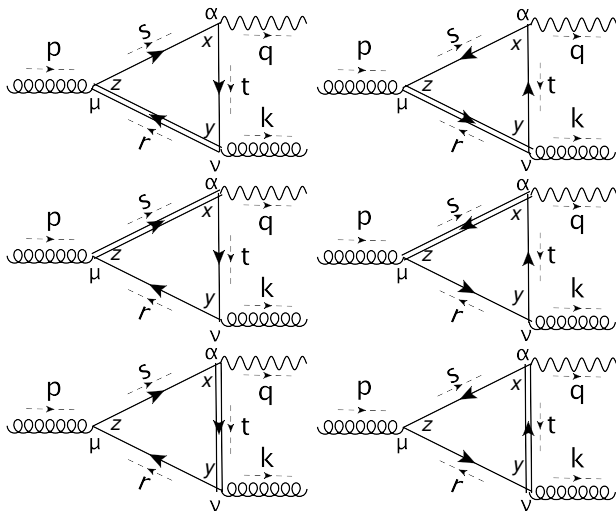
Magnetic field generated in peripheral HICs



Gluon fusion



Gluon splitting



Gluon fusion Matrix Element

$$\begin{aligned}\widetilde{\mathcal{M}}_{gg \rightarrow \gamma} &= - \int d^4x d^4y d^4z \int \frac{d^4r}{(2\pi)^4} \frac{d^4s}{(2\pi)^4} \frac{d^4t}{(2\pi)^4} \\ &\times e^{-it \cdot (y-x)} e^{-is \cdot (x-z)} e^{-ir \cdot (z-y)} e^{-ip \cdot z} e^{-ik \cdot y} e^{iq \cdot x} \\ &\times \left\{ \text{Tr} \left[ieq_f \gamma_\alpha iS(s) ig\gamma_\mu t^c iS(r) ig\gamma_\nu t^d iS(t) \right] \right. \\ &+ \left. \text{Tr} \left[ieq_f \gamma_\alpha iS(t) ig\gamma_\nu t^d iS(r) ig\gamma_\mu t^c iS(s) \right] \right\} \\ &\times \Phi(x, y) \Phi(y, z) \Phi(z, x) \epsilon^\mu(\lambda_p) \epsilon^\nu(\lambda_k) \epsilon^\alpha(\lambda_q) + \text{C.C.}\end{aligned}$$

Quark propagator in a magnetic field

$$\begin{aligned}S(x, x') &= \Phi(x, x') \int \frac{d^4 p}{(2\pi)^4} e^{-p \cdot (x-x')} S(p), \\iS(p) &= \int_0^\infty \frac{d\tau}{\cos(|eq_f B|\tau)} e^{i\tau [p_\parallel^2 - p_\perp^2 \frac{\tan(\tau)}{\tau} - m_f^2 + i\epsilon]} \\&\times \left\{ [\cos(\tau) + \gamma_1 \gamma_2 \sin(\tau)] (m_f + \not{p}_\parallel) - \frac{\not{p}}{\cos(\tau)} \right\}, \\ \Phi(x, x') &= e^{ieq_f \int_{x'}^x d\xi^\mu [A_\mu + \frac{1}{2} F_{\mu\nu} (\xi - x')^\nu]}.\end{aligned}$$

LLL and 1LL massless quarks

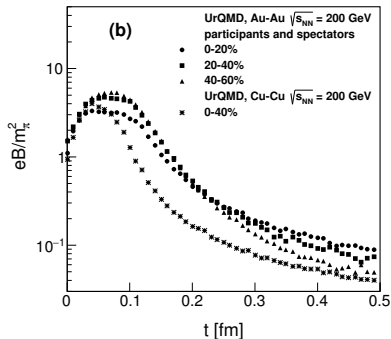
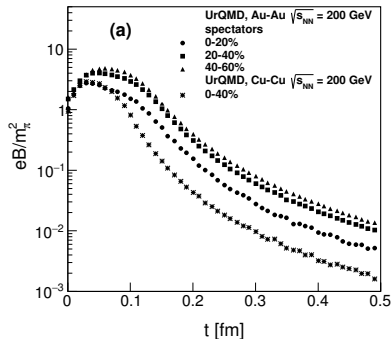
$$iS^{(0)}(p) = i \frac{e^{-p_{\perp}^2/|eq_f B|}}{p_{\parallel}^2} \not{p}_{\parallel} \mathcal{O}^-$$

$$iS^{(1)}(p) = -2i \frac{e^{-p_{\perp}^2/|eq_f B|}}{p_{\parallel}^2 - 2|eq_f B|}$$

$$\times \left[2\not{p}_{\parallel} \mathcal{O}^- \left(1 - \frac{2p_{\perp}^2}{|eq_f B|} \right) - 2\not{p}_{\parallel} \mathcal{O}^+ + 4\not{p}_{\perp} \right]$$

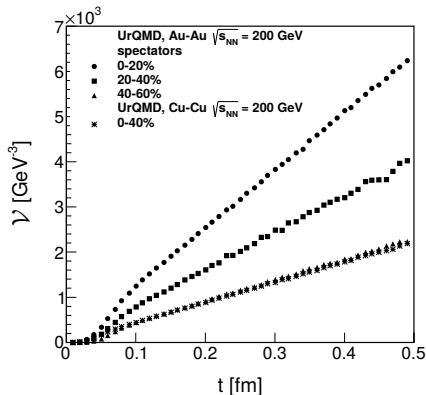
$$\mathcal{O}^{\pm} = [1 \pm i\gamma^1 \gamma^2 \text{sign}(eq_f B)] / 2$$

Field strength as a function of time: Lienard-Wiechert potential UrQMD



$$e\mathbf{B}(\mathbf{x}, t) = \alpha_{em} \sum_j \frac{(1 - v_j^2) \mathbf{v}_j \times \mathbf{R}_j}{R_j^3 \left[1 - \frac{(\mathbf{v}_j \times \mathbf{R}_j)^2}{R_j^2} \right]^{3/2}}$$

UrQMD interaction volume as a function of time



$$V(t) = 2t\pi r_A^2 \left(\frac{N_{\text{part}}}{2N} \right)^{2/3}$$

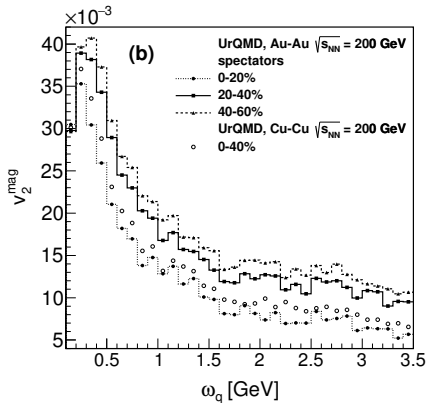
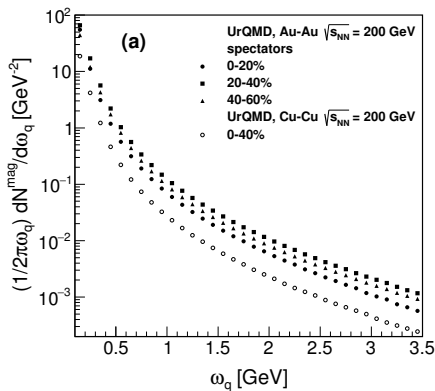
Matrix element squared

$$\overline{\sum_{c,p,f} |\widetilde{\mathcal{M}}|^2} = \nu \Delta \tau (2\pi)^4 \overline{\sum_{c,p,f} \left[\delta^{(4)}(q - k - p) |\mathcal{M}_{gg \rightarrow \gamma}|^2 + \delta^{(4)}(q + k - p) |\mathcal{M}_{g \rightarrow g\gamma}|^2 \right]},$$

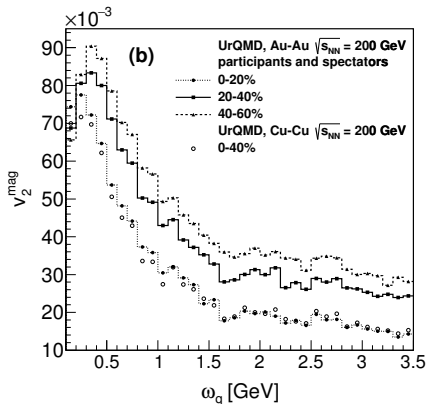
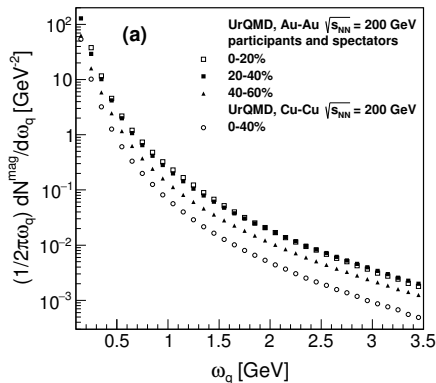
where

$$\begin{aligned} \overline{\sum_{c,p,f} |\mathcal{M}_{gg \rightarrow \gamma}|^2} &= \overline{\sum_{c,p,f} |\mathcal{M}_{g \rightarrow g\gamma}|^2} \\ &= \frac{2\alpha_{\text{em}} \alpha_s^2 q_{\perp}^2}{\pi \omega_q^2} \sum_f q_f^2 (2\omega_p^2 + \omega_k^2 + \omega_p \omega_k) \\ &\times \exp \left\{ -\frac{q_{\perp}^2}{\omega_q^2} (\omega_p^2 + \omega_k^2 + \omega_p \omega_k) \right\} \end{aligned}$$

Photon yield and v_2 from magnetic fields (spectators)

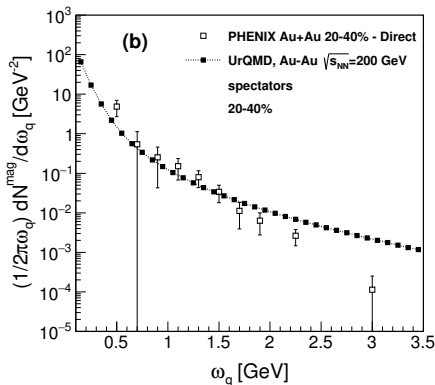
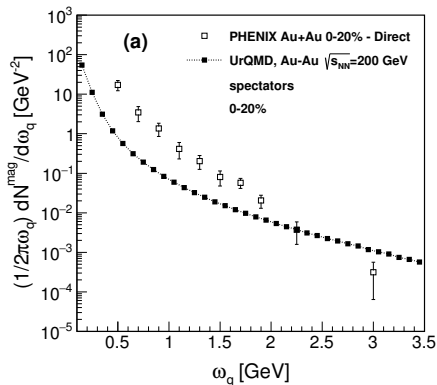


Photon yield and v_2 from magnetic fields (spectators & participants)



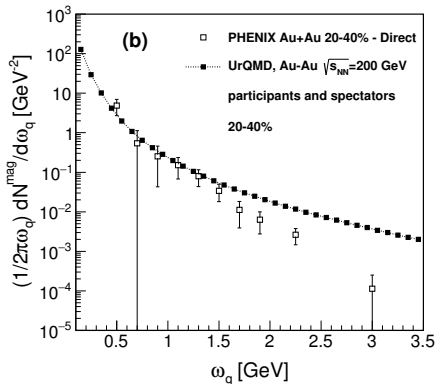
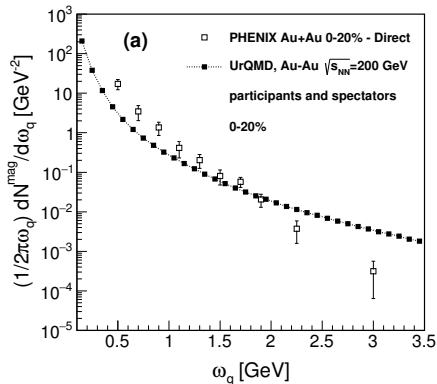
Excess photon yield compared to PHENIX

- Direct photons from Paquet *et. al*, Phys. Rev. C **93**, 044906 (2016).

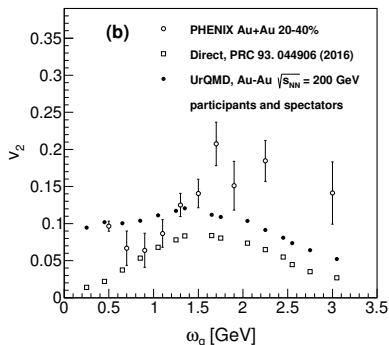
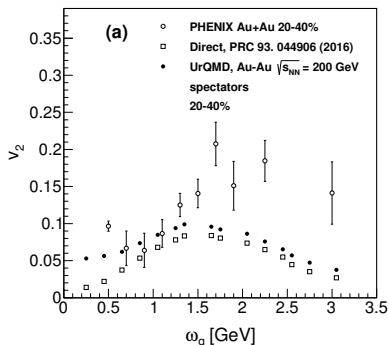


Excess photon yield compared to PHENIX

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Weighed v_2 average (prompt and magnetic)



$$v_2(\omega_q) = \frac{\sum_{i=1}^m \left[\frac{dN}{d\omega_q} \right]_i [v_2^{\text{mag}}(\omega_q)]_i + \frac{dN^{\text{direct}}}{d\omega_q}(\omega_q) v_2^{\text{direct}}(\omega_q)}{\sum_{i=1}^m \left[\frac{dN}{d\omega_q} \right]_i + \frac{dN^{\text{direct}}}{d\omega_q}(\omega_q)},$$

- Intense magnetic field effects during the very early stages of a Relativistic Heavy-Ion Collision.
- High gluon occupation number during early stages together with the magnetic field opens up gluon fusion/splitting channel for photon production.
- We used UrQMD to compute the time evolution of the field intensity and volume reaction to include magnetic field effects in the strong field limit.
- Photon yield and elliptic flow excess better described with magnetic field effects at low energy for semi-central Au + Au and Cu + Cu collisions.