

(Mis)guided advice towards facing the unknown

Gilad Perez

Weizmann

(1st) Meeting on the future of *Experimental HEP* post LHC (in particular in Israel)

Dec. 2017. Weizmann



This is an experimental discussion; I don't have answers, just try to give perspective

- ◆ Pragmatic, th.-free approach & its limitation - the “what have learnt” reduction problem.
- ◆ Observationally-based: a special moment in HEP, the logarithmic crisis.
- ◆ Why things are naturally-speaking worse ? From LHC-data to relaxion-th..
(tiny encouragement: relaxion and the edge of the log scale ...)

(◆ Standard candles)



◆ Summary.



pragmatic approach #1: th.-free search strategy

- ◆ Right now @ the LHC we are searching in model indep' manner...

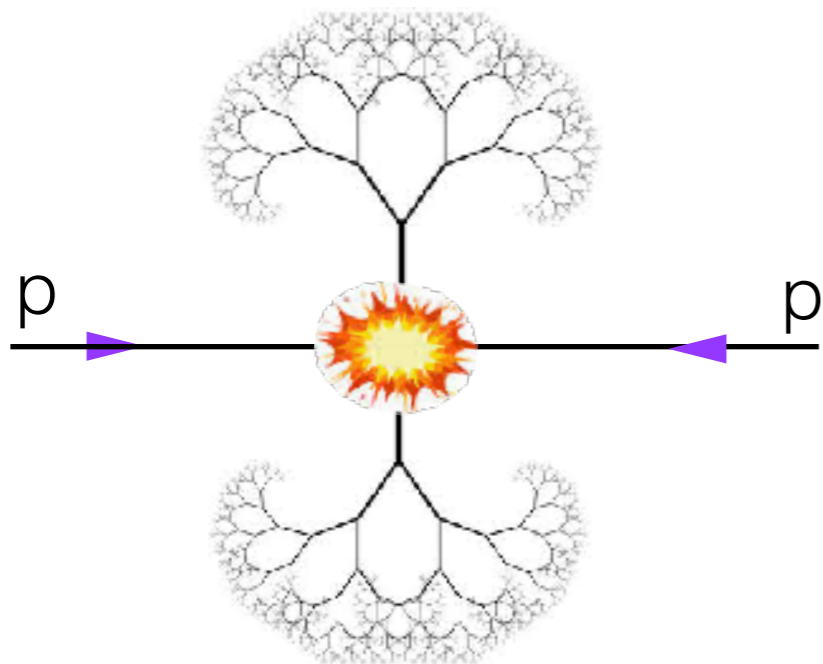
pragmatic approach #1: th.-free search strategy

- ◆ Right now @ the LHC we are searching in model indep' manner...

Not true: we filter $\sim 10^{10}$ Hz of data down to 10^3 Hz (because it's boring ...)

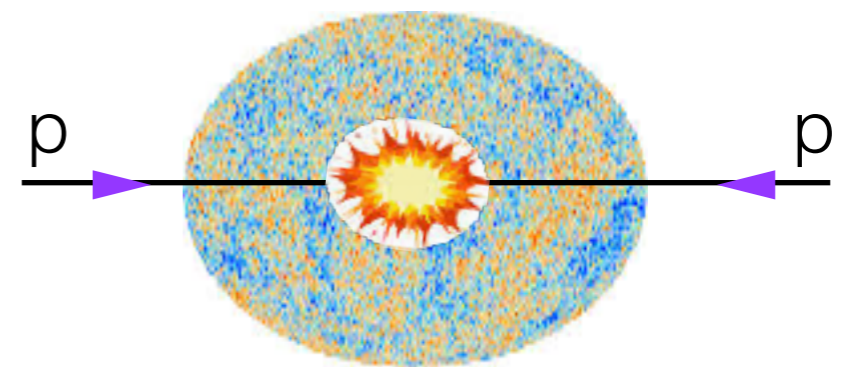
We mostly throw away everything that is soft; what if soft is everything?

Feynman - showering as self similar process



Wolf, Dremin & Kittel (95); Strassler & Zurek (06)

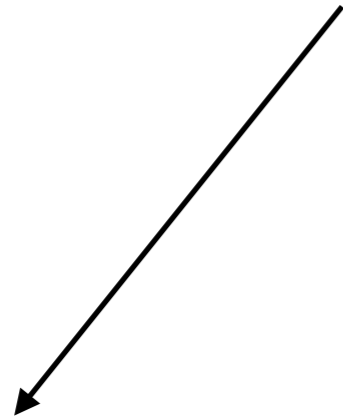
conformal hidden sector



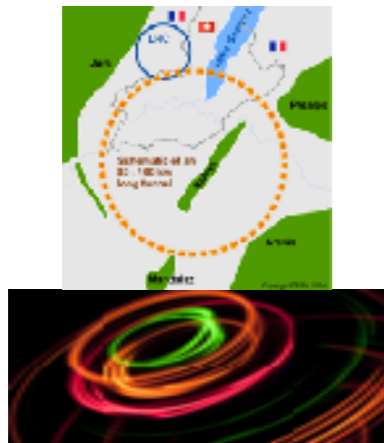
Hofman & Maldacena (08); Georgi (07);
Cai, Cheng, Medina & Terning (09);

Pragmatic #2: th.-free approach to future exp.

How to compare different possibilities?



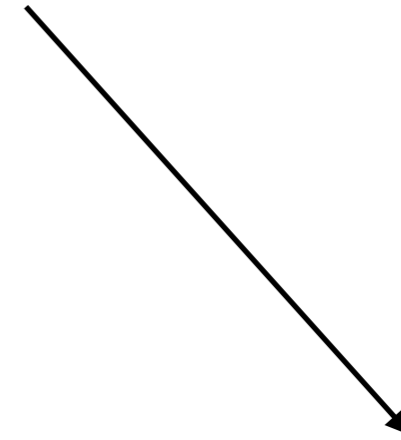
Energy ?



FCC



Luminosity ?



Precision ?



Pragmatic #3: just compare amount of data

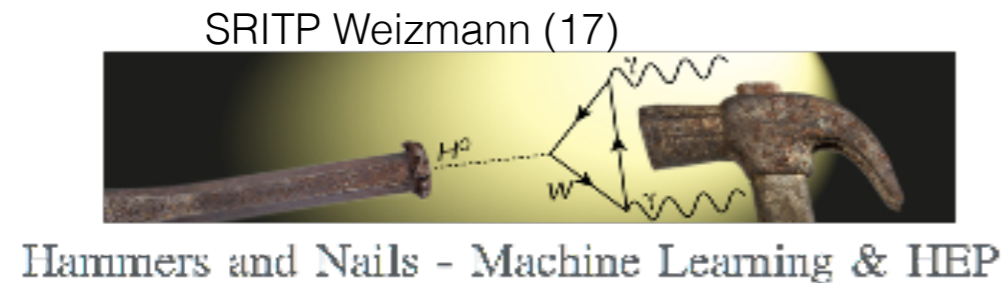
◆ Number of different type of bits per exp’:

# of bits	LHC	NA62	SHiP	FCChh	ILC	SARAF
Ps/es/ POT	“ 10^{16} ”	10^{18}	10^{20}	“ 10^{18} ”		10^{22}
Kaon Produced		10^{14}	10^{16}		smaller	0
B’s Produced	10^{12-14}	10^8	0	10^{16} ?	smaller	0
Higgsses S/B	10^5 10^{-3}	0	0	10^9 10^{-4} ?	10^5 $10^{0,-1}$	0

However, comparing # of protons to # of (clean) Higgsses is insane ...

pragmatic approach #4: machine learning, automatization will improve our large-data-searches

- ◆ Machine learning (ML) can potentially significantly boost the field.



- ◆ In what sense (from experts):

K. Cranmer: Jets, Higgs-EFT, <https://indico.physics.lbl.gov/indico/event/546/>, Berkeley.

Brehmer, Cranmer, Kling & Plehn (16)

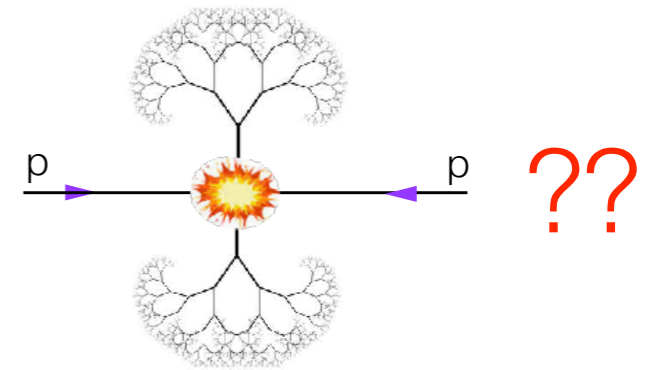
E. Gross: Jets, improving detector simulation ...

S. Bressler: (not necessarily related to ML, but linked to above theme)
“generating signal hypotheses from the data - a data focused paradigm”

Two thoughts on “pragmatism”

◆ At present I don't see a phase transition.

But, it might be too early.



◆ Generically, there is no attempt to deal with the following problem:

What do we learn if there is a null result ?!

Understanding the potential impact of large operation is a must and requires some level of reductionism.

Two thoughts on “pragmatism”

Decay Mode	Projected/Current 2σ Limit on Br(\mathcal{F}_i) 7+8 [14] TeV	Production Mode	Limit on $\frac{\text{Br}(\mathcal{F}_i)}{\text{Br}(\text{non-SM})}$ $\frac{\sigma}{\sigma_{\text{SM}}} \cdot \text{Br}(\text{non-SM})$ 7+8 [14] TeV
$jjjj$	> 1 [0.1*]	W	0.25
$llll$	[?]	G	0.09
$j\mu\mu$	$0.002 - 0.008$ [5 - 20] [14]	W	0.1
$b\bar{b}\mu\mu$	$(2 - 7) \cdot 10^{-4}$ [(0.6 - 2) · 10 ⁻⁴]	G	0.0

◆ At present I don't see a phase transition.

But, it might be too early.

Decay Mode	Projected/Current 2σ Limit on Br(\mathcal{F}_i) 7+8 [14] TeV	Production Mode	Limit on $\frac{\text{Br}(\mathcal{F}_i)}{\text{Br}(\text{non-SM})}$ $\frac{\sigma}{\sigma_{\text{SM}}} \cdot \text{Br}(\text{non-SM})$ 7+8 [14] TeV	quarks allowed	Limit on 7+8 [14] TeV	Decay Mode	Projected/Current 2σ Limit quarks suppressed on Br(\mathcal{F}_i) Limit on 7+8 [14] TeV	Production Mode	Limit on $\frac{\text{Br}(\mathcal{F}_i)}{\text{Br}(\text{non-SM})}$ $\frac{\sigma}{\sigma_{\text{SM}}} \cdot \text{Br}(\text{non-SM})$ 7+8 [14] TeV	Limit on $\frac{\text{Br}(\mathcal{F}_i)}{\text{Br}(\text{non-SM})}$ $\frac{\sigma}{\sigma_{\text{SM}}} \cdot \text{Br}(\text{non-SM})$ 7+8 [14] TeV
$jjjj$	> 1 [0.1*]	W	0.25			$(b\bar{b}) \cancel{E}_T$	> 1 [0.2*]	Z	> 1 [?]	$h \rightarrow \psi\psi \rightarrow \bar{f}_1 f_1 + \bar{f}_2 f_2 + \cancel{E}_T$
$llll$	[?]	G	0.09			$(\tau\tau) \cancel{E}_T$	> 1 [> 1*]	Z	> 1 [> 1*]	$h \rightarrow \psi\psi \rightarrow \bar{f}_1 f_1 + \bar{f}_2 f_2 + \cancel{E}_T$
$j\mu\mu$	$0.002 - 0.008$ [5 - 20] [14]	W	0.1	0.8	0.9 [0.2]	0 $(\tau\tau) \cancel{E}_T$	> 1 [> 1*]	Z	0.15	> 1 [> 1*]
$b\bar{b}\mu\mu$	$(2 - 7) \cdot 10^{-4}$ [(0.6 - 2) · 10 ⁻⁴]	G	0.0	0.1	> 1 [1]	0 $ll \cancel{E}_T$	0.07 [?]	G	0.30	0.2 [?]
$b\bar{b}\mu\mu$	$(2 - 7) \cdot 10^{-4}$ [(0.6 - 2) · 10 ⁻⁴]	G	0.0	$0.5 - 1$	$0.5 - 1$ [0.5 - 1]	0 $ll \cancel{E}_T$	> 1 [?]	G, V	-	-

◆ Generically, there is no attempt to deal with:

What do we learn if there is a negative result ?!

Understanding the potential impact of large operation is a must and requires some level of reductionism.

Decay Mode	Projected/Current 2σ Limit on Br(\mathcal{F}_i) 7+8 [14] TeV	Production Mode	Limit on $\frac{\text{Br}(\mathcal{F}_i)}{\text{Br}(\text{non-SM})}$ $\frac{\sigma}{\sigma_{\text{SM}}} \cdot \text{Br}(\text{non-SM})$ 7+8 [14] TeV	Decay Mode	Projected/Current 2σ Limit on Br(\mathcal{F}_i) 7+8 [14] TeV	Production Mode	Limit on $\frac{\text{Br}(\mathcal{F}_i)}{\text{Br}(\text{non-SM})}$ $\frac{\sigma}{\sigma_{\text{SM}}} \cdot \text{Br}(\text{non-SM})$ 7+8 [14] TeV
$\{\mu\mu\}\{\mu\mu\}$	$1 \cdot 10^{-5}$ (5 · 10 ⁻³) [?]	G	$3 \cdot 10^{-5}$	$\tau\tau\mu\mu$	[?]	G	[?]
$\{ee\}\{ee\}$	limit unclear [?]	W, G		$\mu\mu\mu\mu$	$1 \cdot 10^{-4}$	G	$1 \cdot 10^{-4}$
$\{\mu\mu\} \cancel{E}_T$	0.03 [?]	W		$\mu\mu\mu\mu$	$1 \cdot 10^{-4}$	G	$1 \cdot 10^{-4}$
$\{\mu\mu\}\{\mu\mu\} \cancel{E}_T$	$1 \cdot 10^{-5}$ (5 · 10 ⁻³) [?]	G		$\mu\mu\mu\mu$	$1 \cdot 10^{-4}$	G	$1 \cdot 10^{-4}$
$\{ee\}\{ee\} \cancel{E}_T$	limit unclear [?]	W, G		$jjjj$	> 1 [0.1*]	W	0.99 [0.1*]
$\{\tau\tau\}\{\mu\mu\}$	$(3 - 7) \cdot 10^{-4}$ [?]	G		$jjjj$	> 1 [0.1*]	W	0.99 [0.1*]
$\{\gamma\gamma\}\{\gamma\gamma\}$	0.01 [?]	G		$\gamma\gamma jj$	0.04 [0.01*]	W	0.008 [1*]
$\{\gamma\gamma\} \cancel{E}_T$?[?]			$\gamma\gamma jj$	0.04 [0.01*]	W	0.008 [1*]
$\{gg\}\{gg\}$	> 1 [0.7]	W		$\gamma\gamma\gamma\gamma$	$2 \cdot 10^{-4}$ [3 · 10 ⁻⁵ *]	G	$1 \cdot 10^{-5}$ [1*]
$\{b\bar{b}\}\{b\bar{b}\}$	0.7 [0.2]	W		$\gamma\gamma\gamma\gamma$	$2 \cdot 10^{-4}$ [3 · 10 ⁻⁵ *]	G	$1 \cdot 10^{-5}$ [1*]

Some guidance: Implications of
conflict with observations
(moving to log scale)

The Standard Model (SM) is incomplete

New particle and forces must exist:

Masses of right handed neutrinos $10^{-9} - 10^{15}$ GeV

Mass of Dark Matter particle $10^{-31} - 10^{20}$ GeV

Mass of new particles required for baryogenesis $10^{-2} - 10^{15}$ GeV

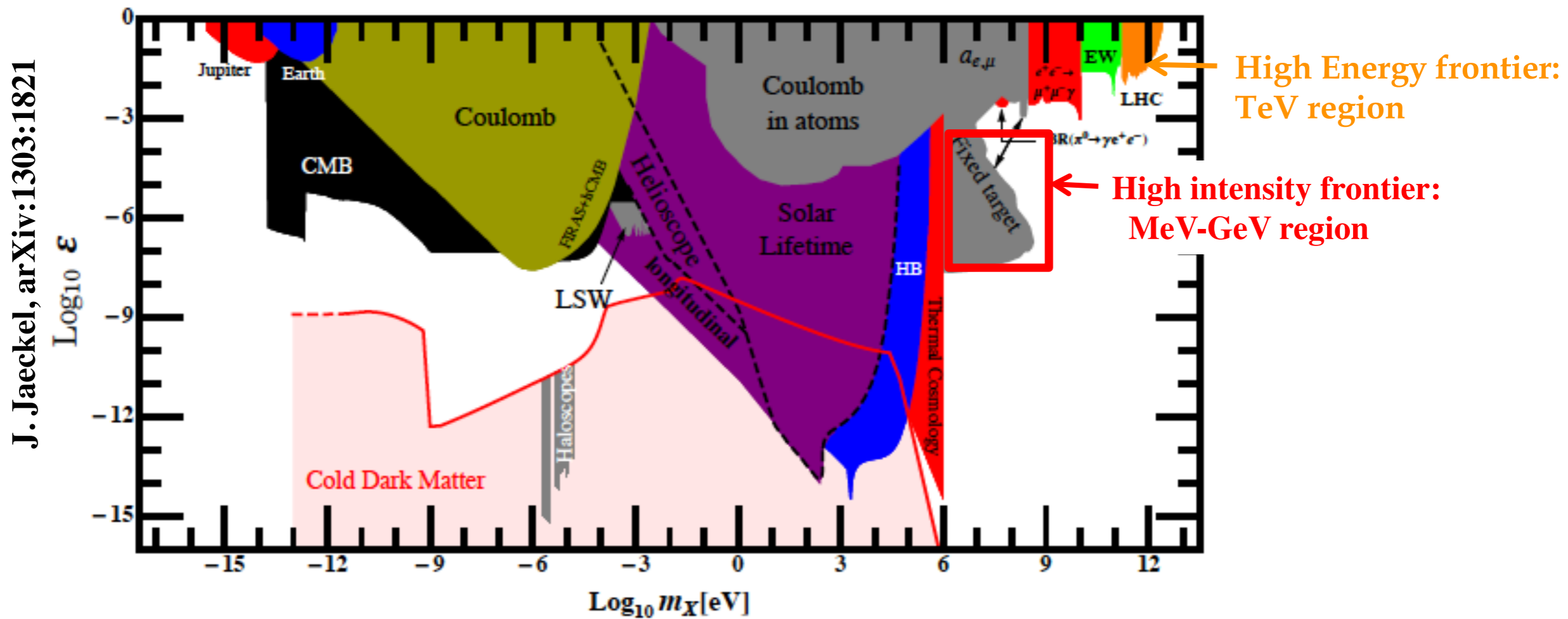
Scale of inflation probably (could be) very high, but at largely unknown ...

The 21st century frustration:

we know that new physics exists but we don't know where ...

Search for motivated new physics with no preferred energy scale - the *logarithmic crisis*

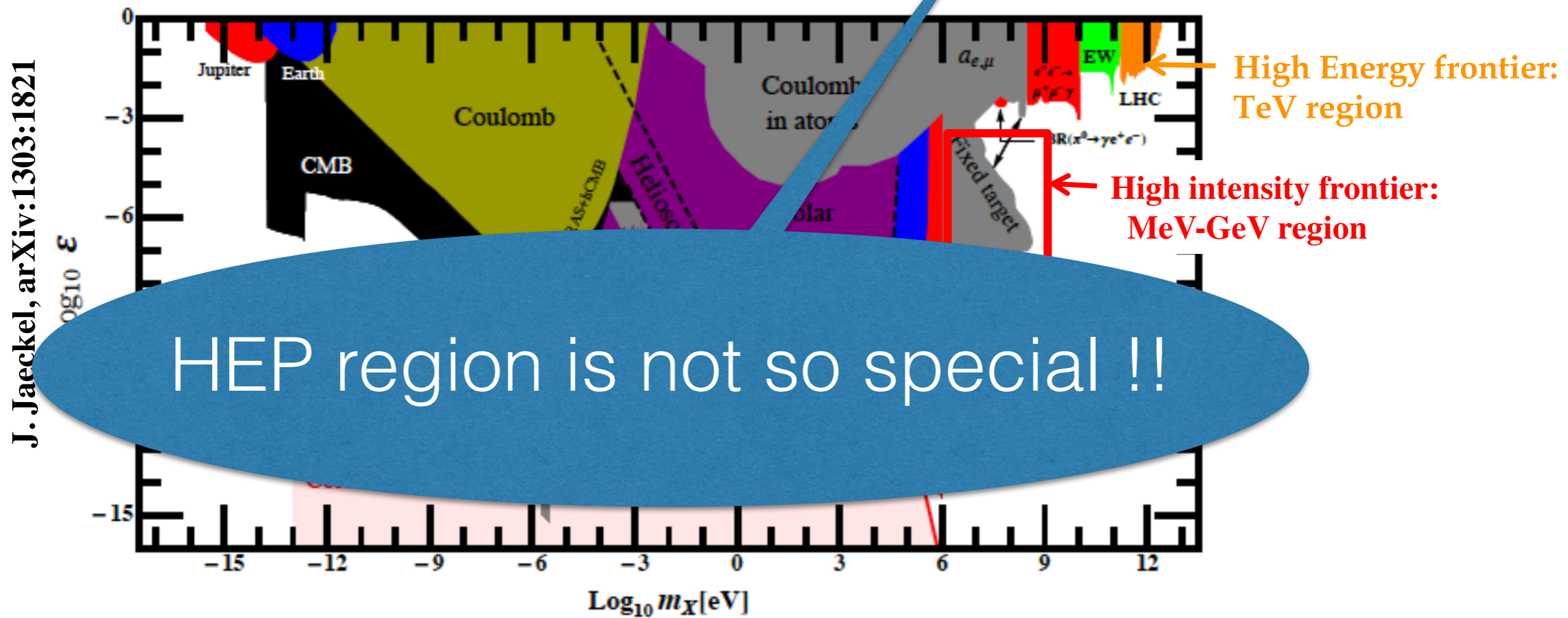
kinetic mixing of a dark photon-ordinary photon versus mass of dark photon



.... *lost among the orders of magnitude*

Search for motivated new physics with no preferred energy scale - the *logarithmic crisis*

kinetic mixing of a dark photon-ordinary photon versus mass of dark photon



.... *lost among the orders of magnitude*

Theoryfull strategy, status of naturalness

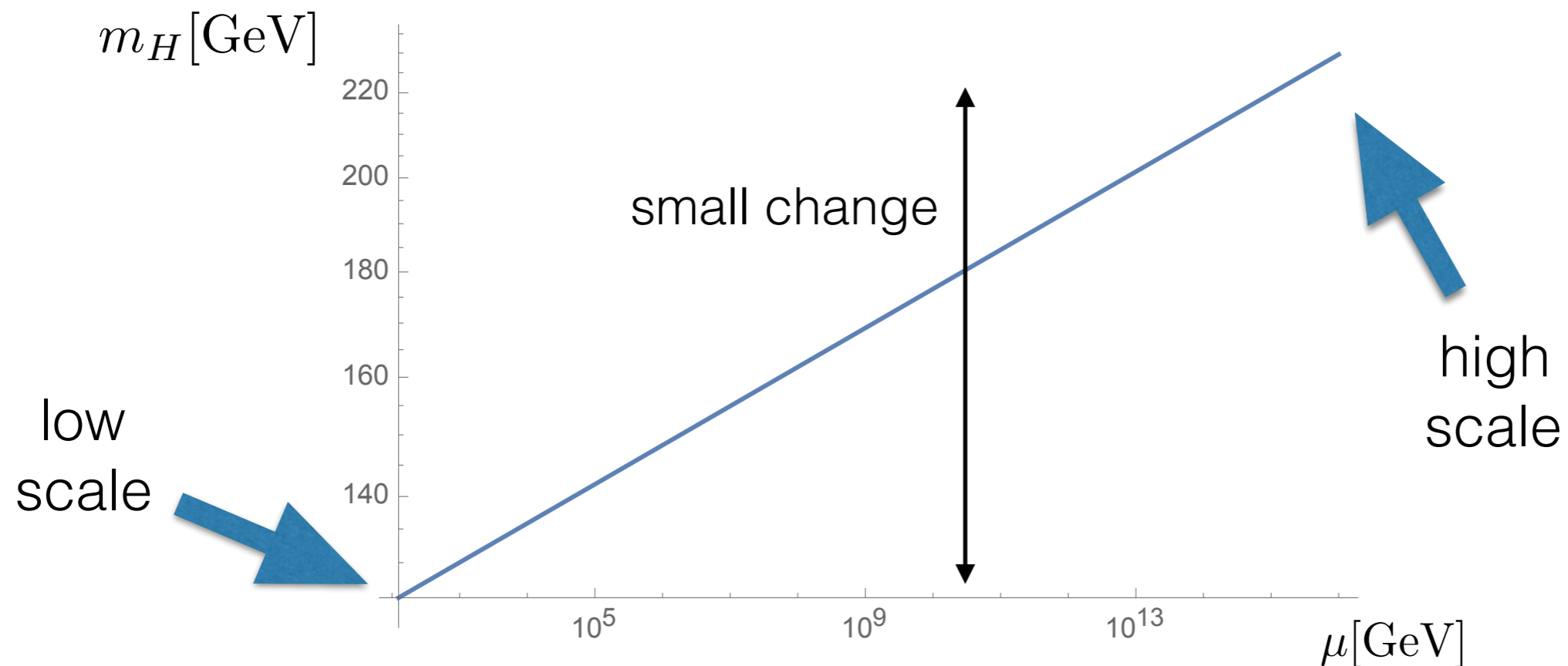
- ◆ Few slides to demystify naturalness.
- ◆ Despite results why is it embarrassingly still the best motivation for new physics in HEP.
- ◆ Why things are possibly actually worse conceptually ?
The relaxion & the naturalness logarithmic crisis.

The Higgs hierarchy/naturalness/fine-tuning problem

Giudice (13)

- ◆ Without new-dynamical scale the fine tuning problem is ill defined;
=> Higgs naturalness is a UV problem.

Higgs RGE, bottom up is natural: $\frac{dm_H^2}{d \ln \bar{\mu}^2} = \beta_m^{\text{SM}} m_H^2$, $\beta_m^{\text{SM}} = \frac{3}{4} \frac{4y_t^2 + 8\lambda - 3g^2 - g_Y^2}{(4\pi)^2}$.



The Higgs hierarchy/naturalness/fine-tuning problem

- ◆ Without new-dynamical scale the fine tuning problem is ill define;
=> Higgs naturalness is a UV problem.

Add particle w coupling λ_N & mass $M \Rightarrow$ Farinaa, Duccio Pappadopulo & Strumia (14)

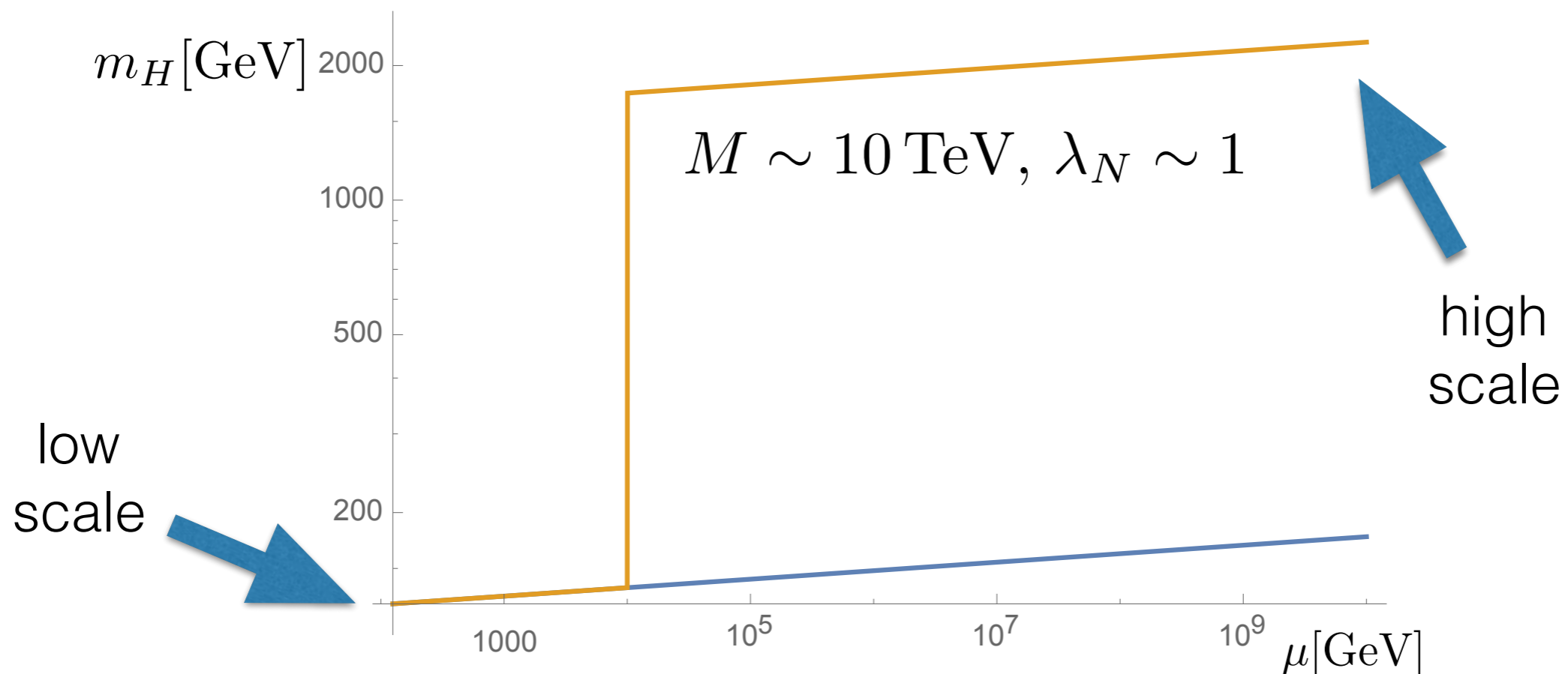
Higgs finite additive correction:
$$\frac{dm_H^2}{d \ln \bar{\mu}^2} = \frac{4\lambda_N^2}{(4\pi)^2} M^2 + \beta_m^{\text{SM}} m_H^2$$

The Higgs hierarchy/naturalness/fine-tuning problem

- ◆ Without new-dynamical scale the fine tuning problem is ill define;
=> Higgs naturalness is a UV problem.

Add particle w coupling λ_N & mass $M \Rightarrow$ Farinaa, Duccio Pappadopulo & Strumia (14)

Higgs finite additive correction:
$$\frac{dm_H^2}{d \ln \bar{\mu}^2} = \frac{4\lambda_N^2}{(4\pi)^2} M^2 + \beta_m^{\text{SM}} m_H^2$$

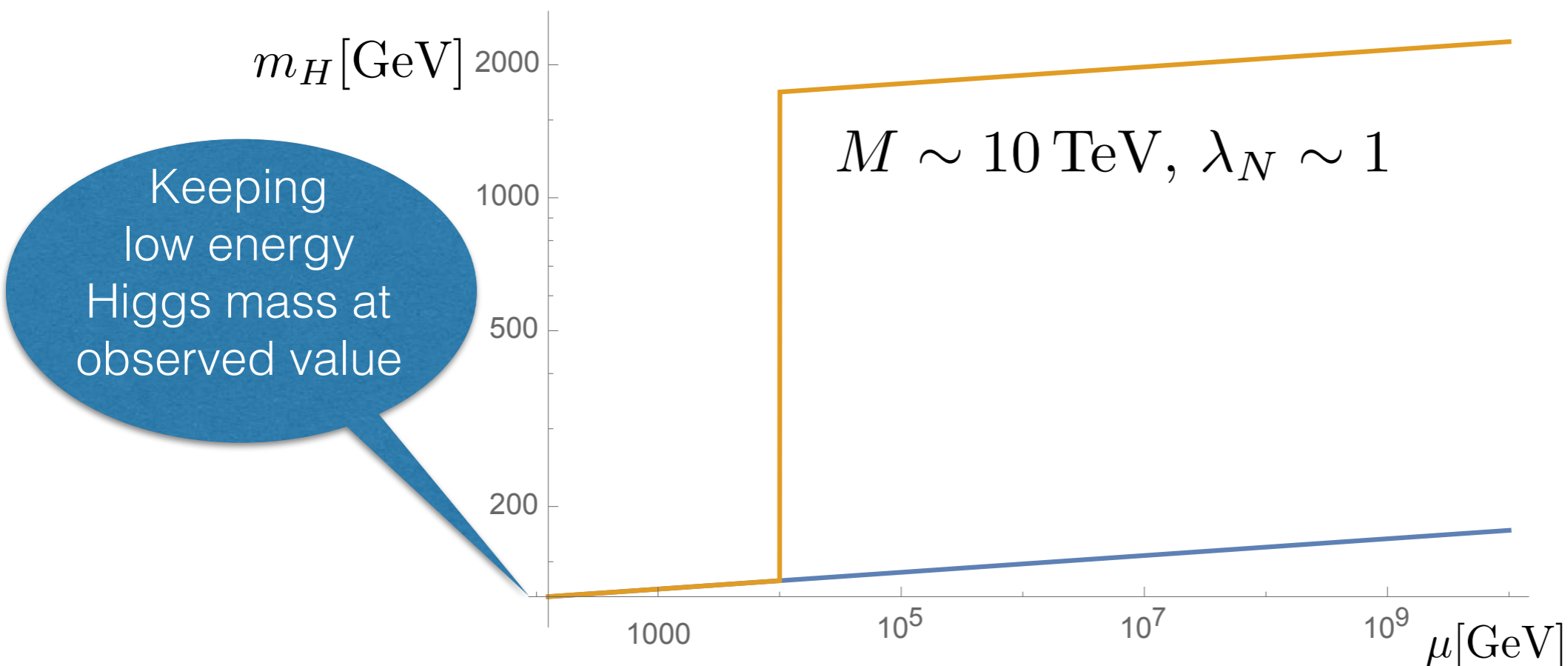


The Higgs hierarchy/naturalness/fine-tuning problem

- ◆ Without new-dynamical scale the fine tuning problem is ill define;
=> Higgs naturalness is a UV problem.

Add particle w coupling λ_N & mass $M \Rightarrow$ Farinaa, Duccio Pappadopulo & Strumia (14)

Higgs finite additive correction:
$$\frac{dm_H^2}{d \ln \bar{\mu}^2} = \frac{4\lambda_N^2}{(4\pi)^2} M^2 + \beta_m^{\text{SM}} m_H^2$$

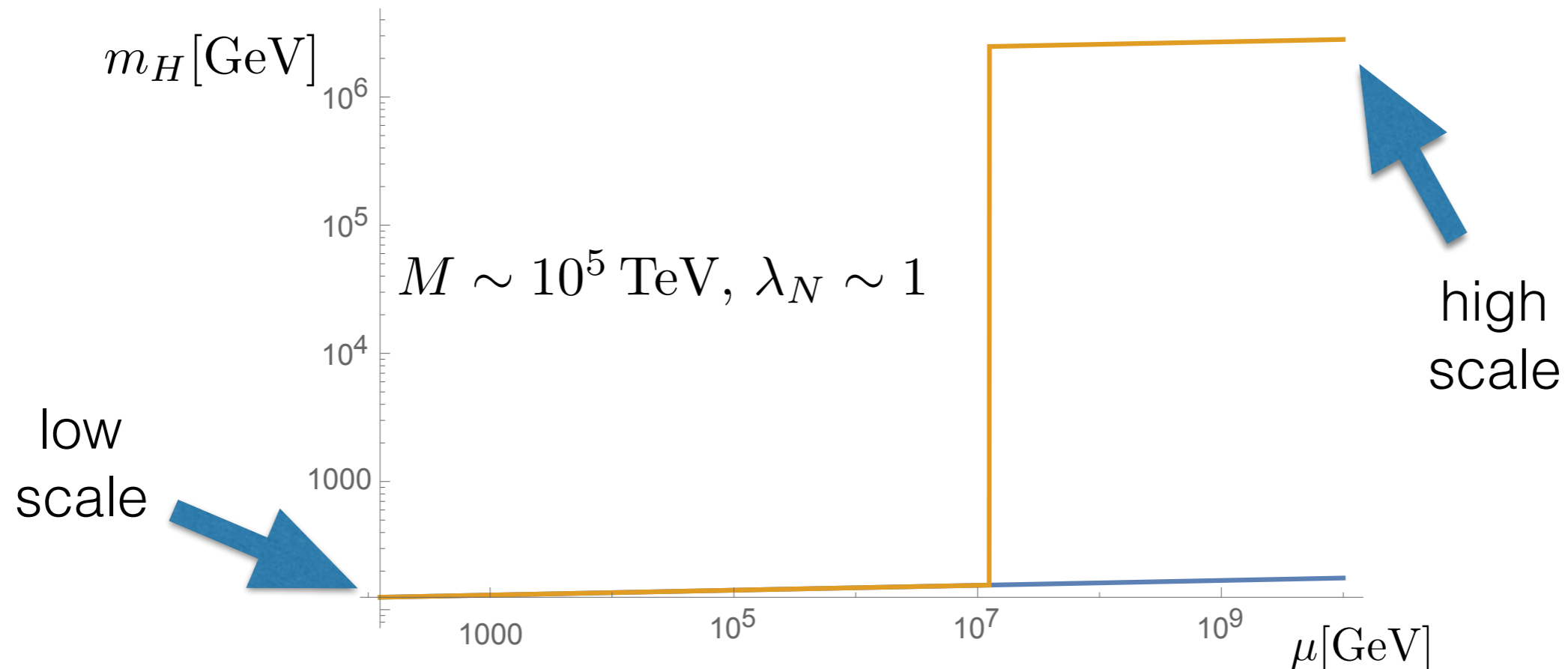


The Higgs hierarchy/naturalness/fine-tuning problem

- ◆ Without new-dynamical scale the fine tuning problem is ill define;
=> Higgs naturalness is a UV problem.

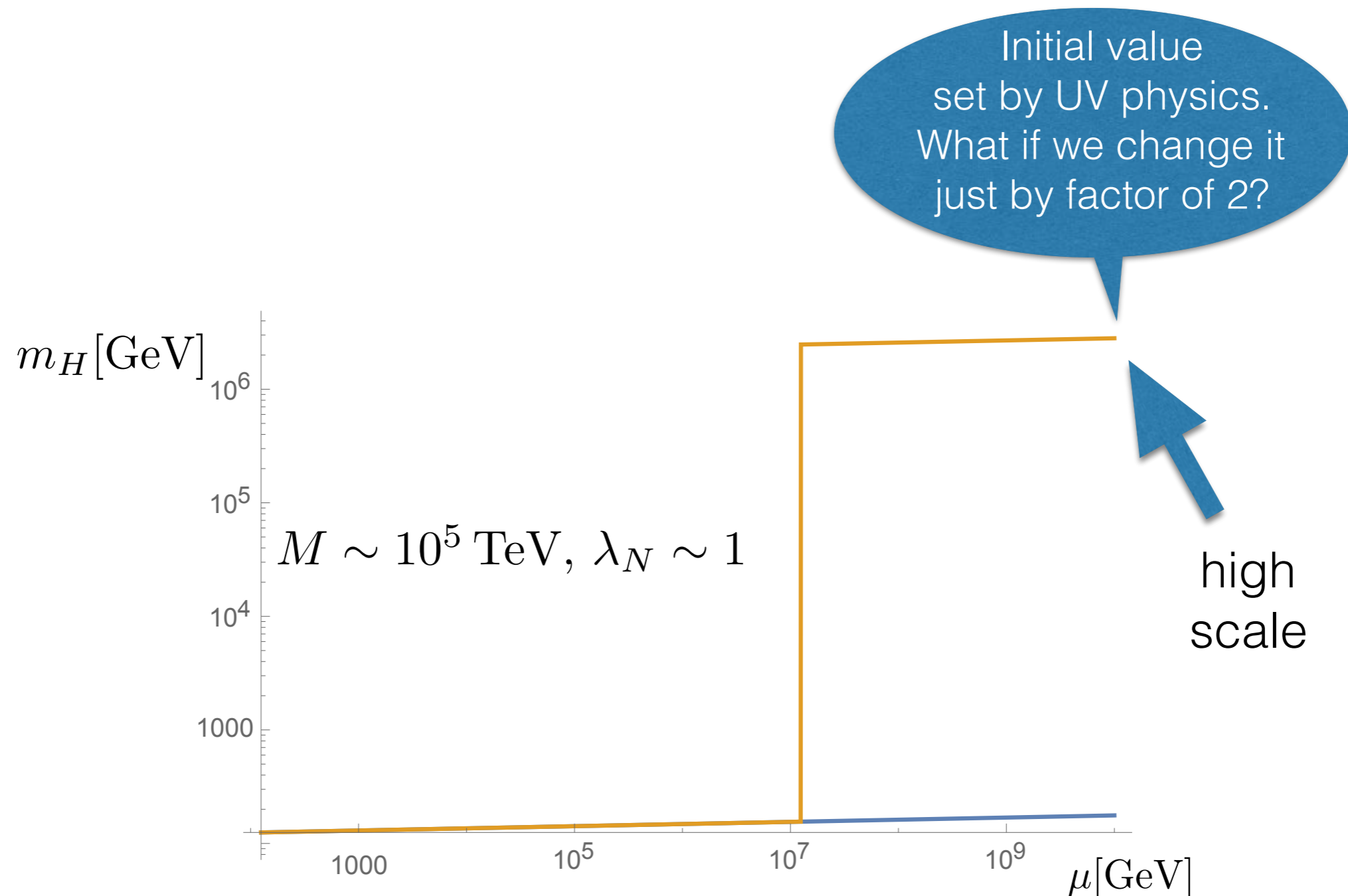
Add particle w coupling λ_N & mass $M \Rightarrow$

Higgs finite additive correction: $\frac{dm_H^2}{d \ln \bar{\mu}^2} = \frac{4\lambda_N^2}{(4\pi)^2} M^2 + \beta_m^{\text{SM}} m_H^2$



The Higgs hierarchy/naturalness/fine-tuning problem

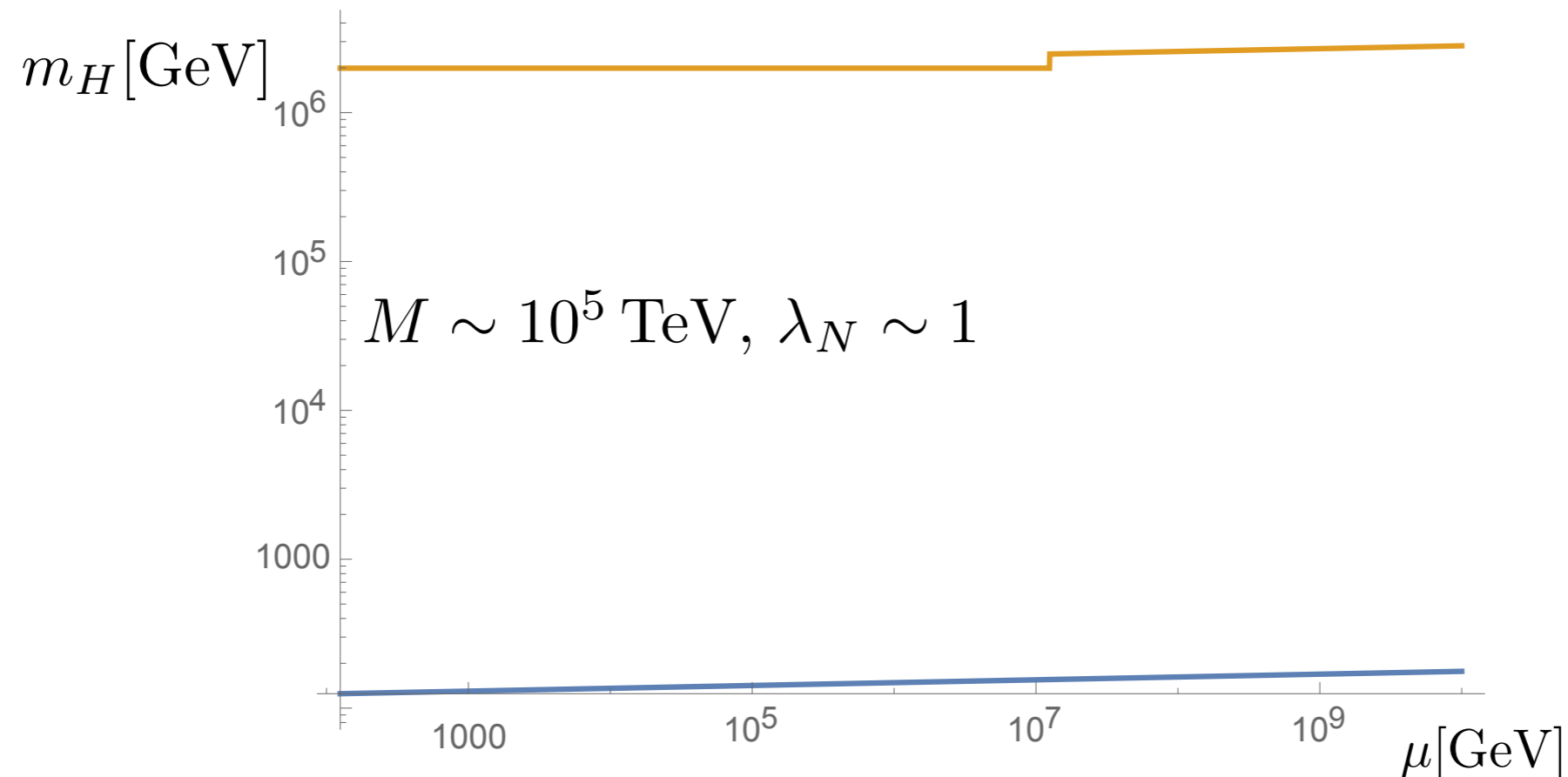
- ◆ Without new-dynamical scale the fine tuning problem is ill define;
=> Higgs naturalness is a UV problem.



The Higgs hierarchy/naturalness/fine-tuning problem

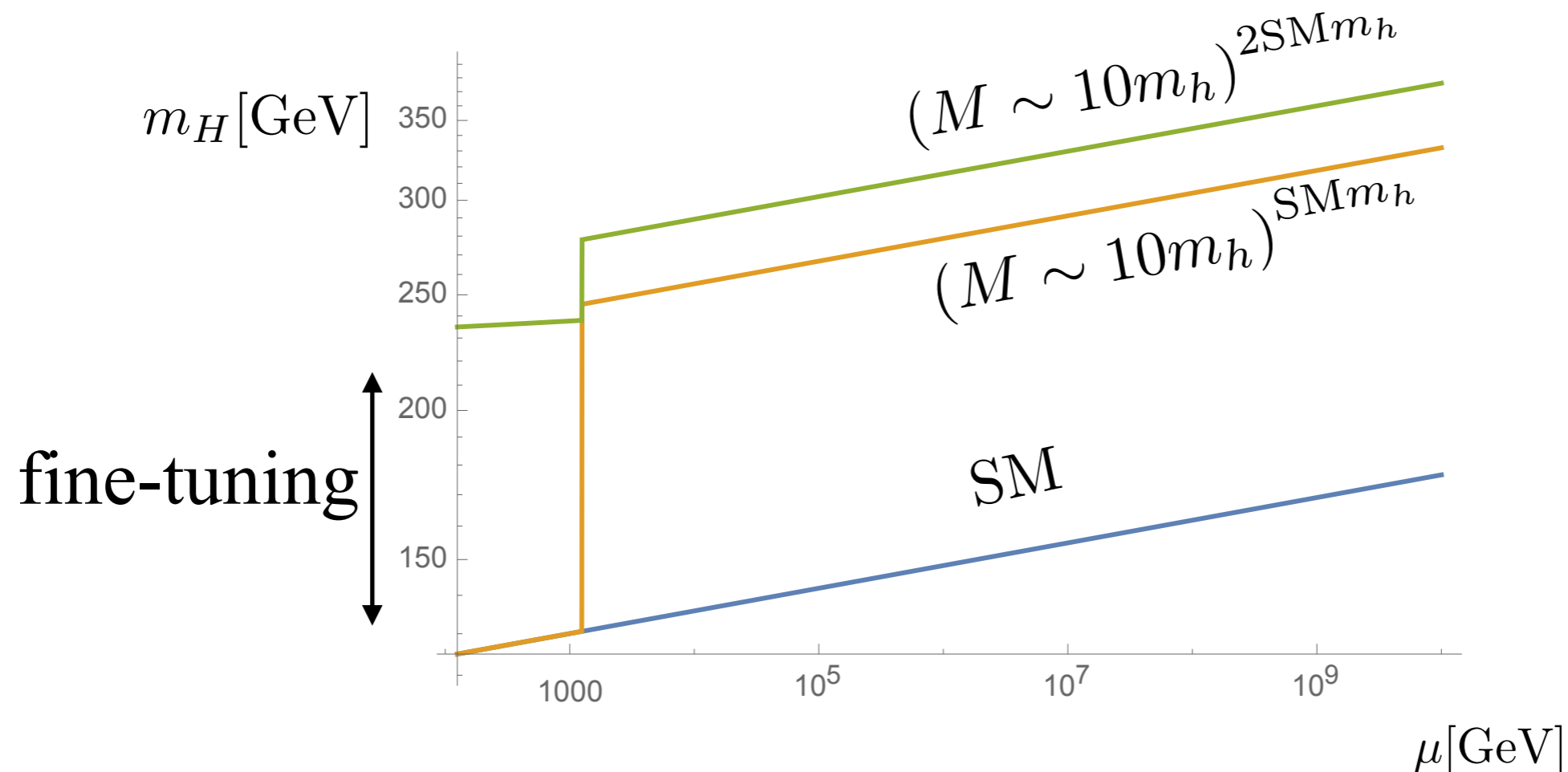
- ◆ Without new-dynamical scale the fine tuning problem is ill define;
=> Higgs naturalness is a UV problem.

The Higgs mass pushed to M => inconsistent \w w data.

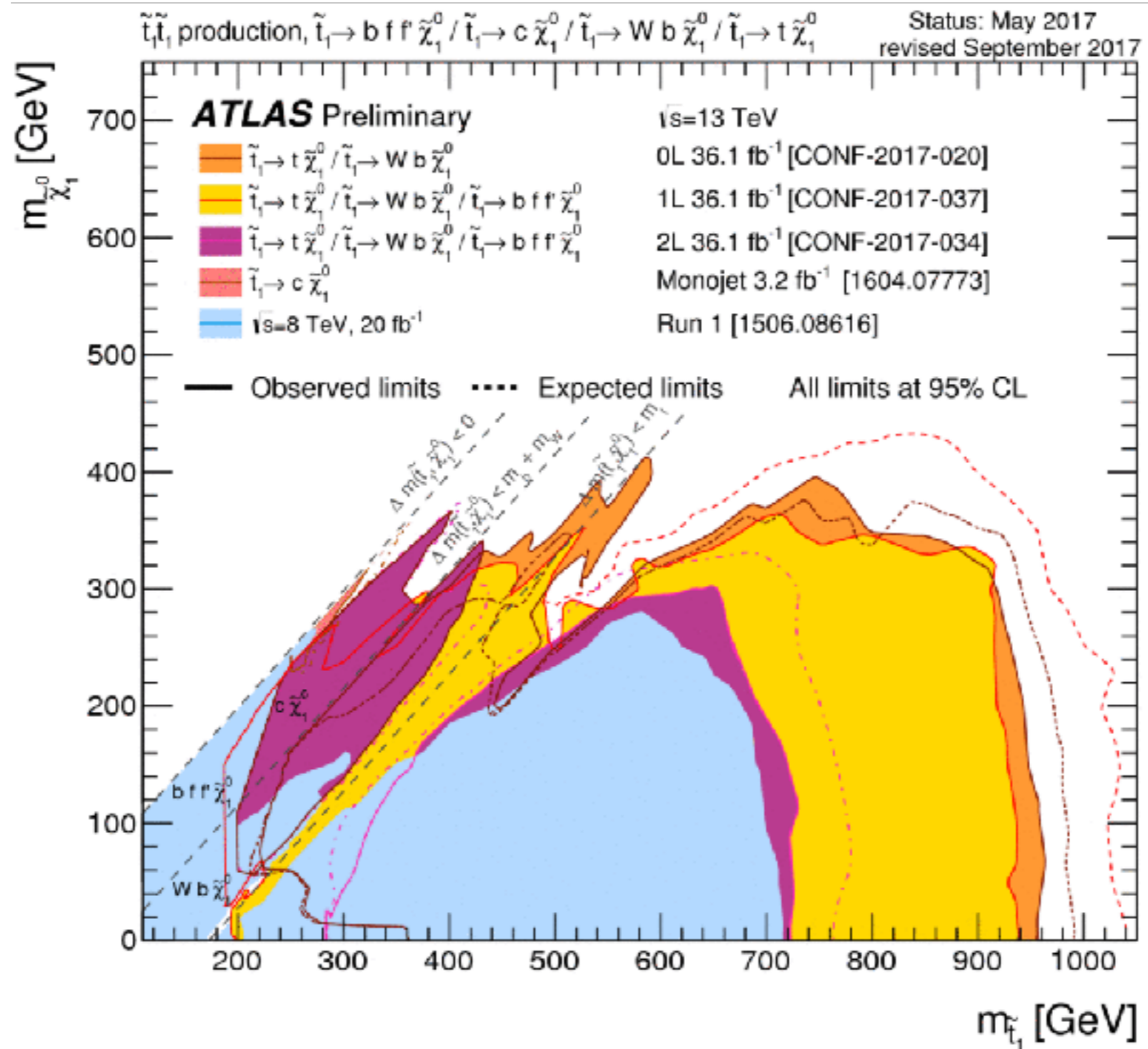


Naturalness gives us motivation to look for \lesssim TeV new particles

Higgs mass evolution in natural theories



The LHC depression - “naturalness is dead”



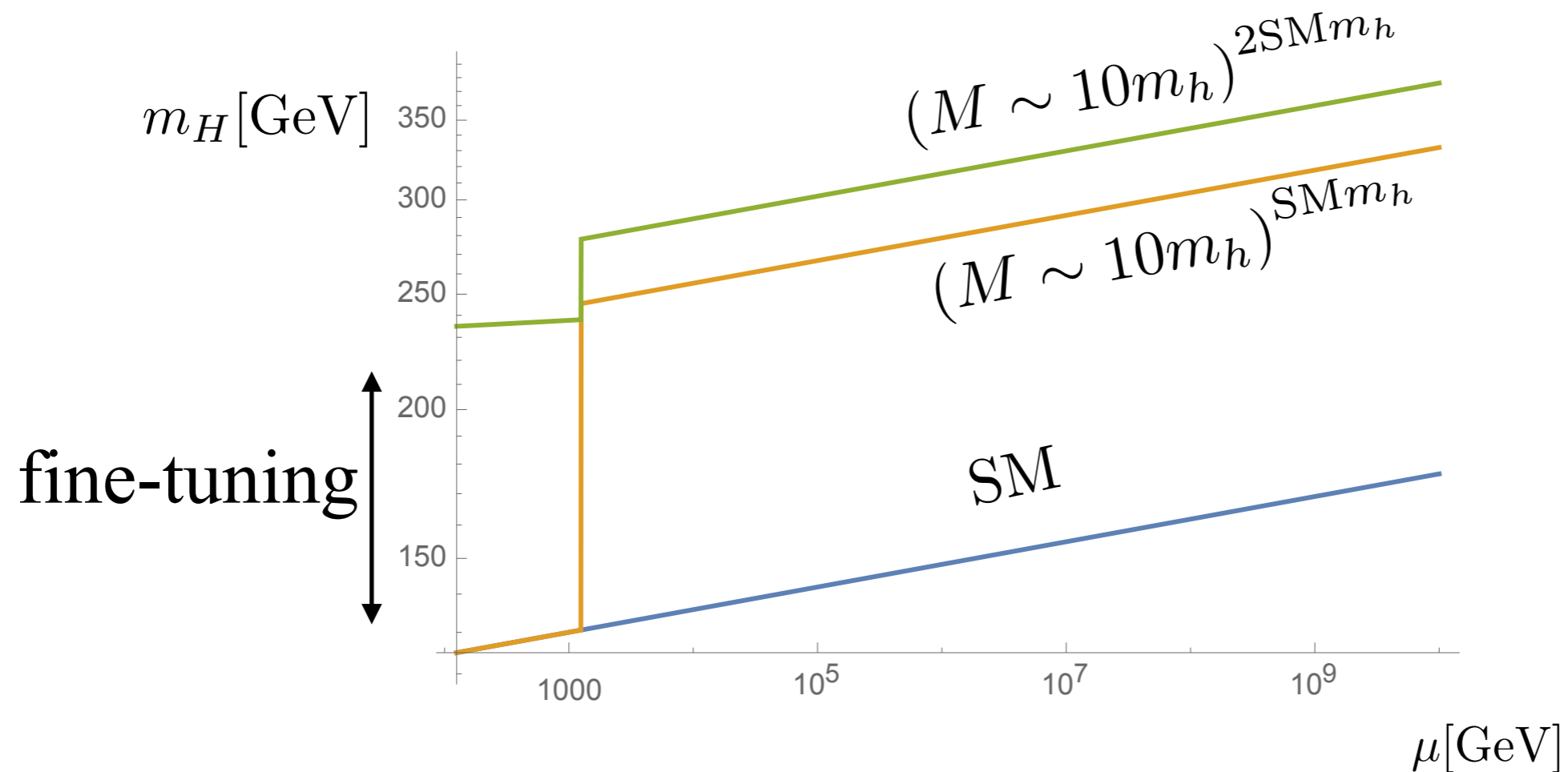
However:

Two reasons for why this logic is not bullet proof:

(i) for the high-energy colliders

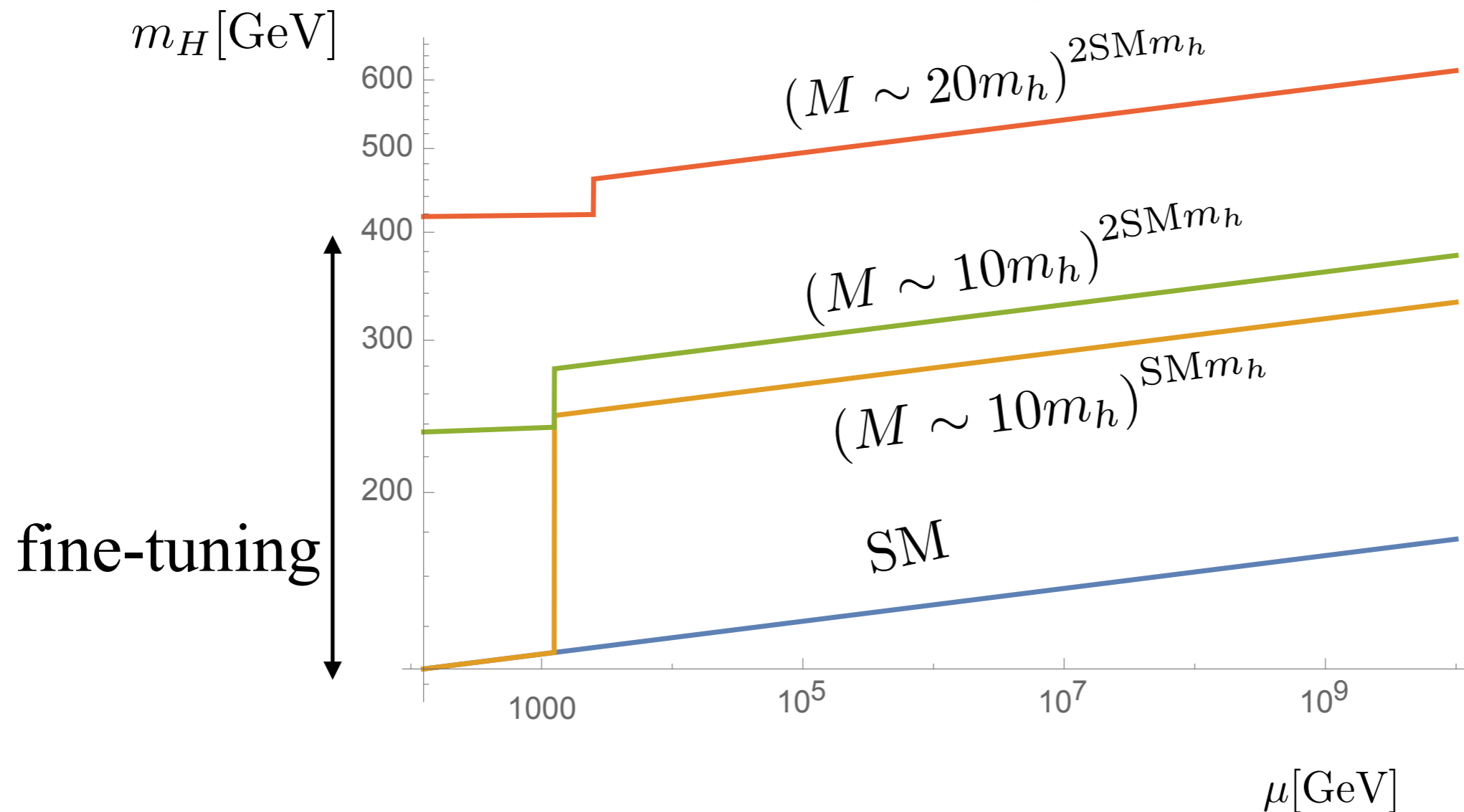
(ii) new & welcome lepton-colliders

First: naturalness is *not* black & white measure



First: naturalness is not black & white measure

Beyond the
HL-LHC reach



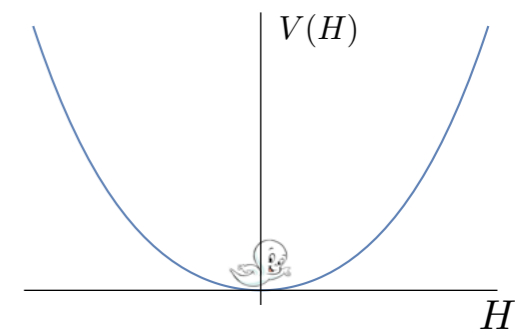
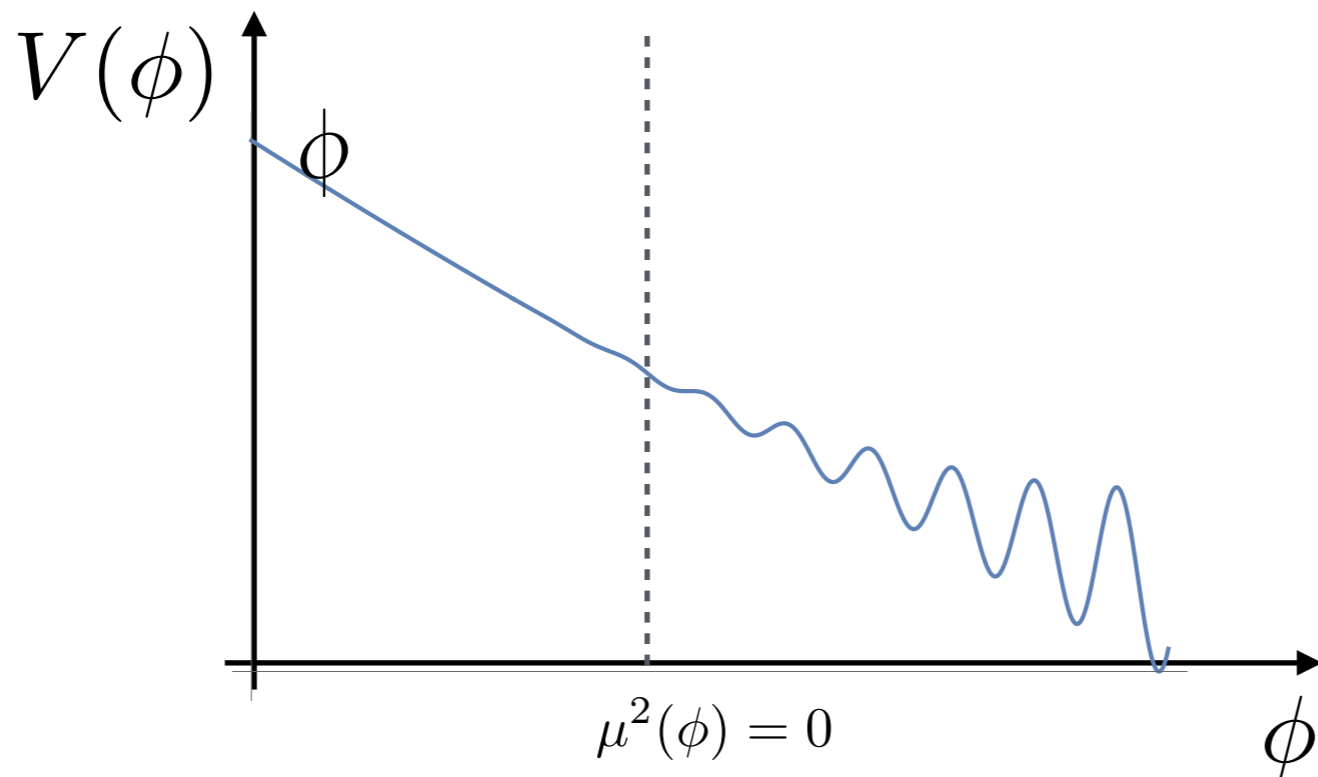
Second: relaxion - naturalness $\not\Rightarrow$ TeV new physics

Graham, Kaplan & Rajendran (15)

◆ A dynamical solution/amelioration of the Higgs fine-tuning problem:

(i) Add a scalar (relaxion) Higgs dependent mass: $(\Lambda^2 - \overbrace{\mu^2(\phi)}^{\mu^2(\phi)}) H^\dagger H$.

(ii) ϕ rolls till μ^2 changes sign $\Rightarrow \langle H \rangle \neq 0 \Rightarrow$ stops rolling.



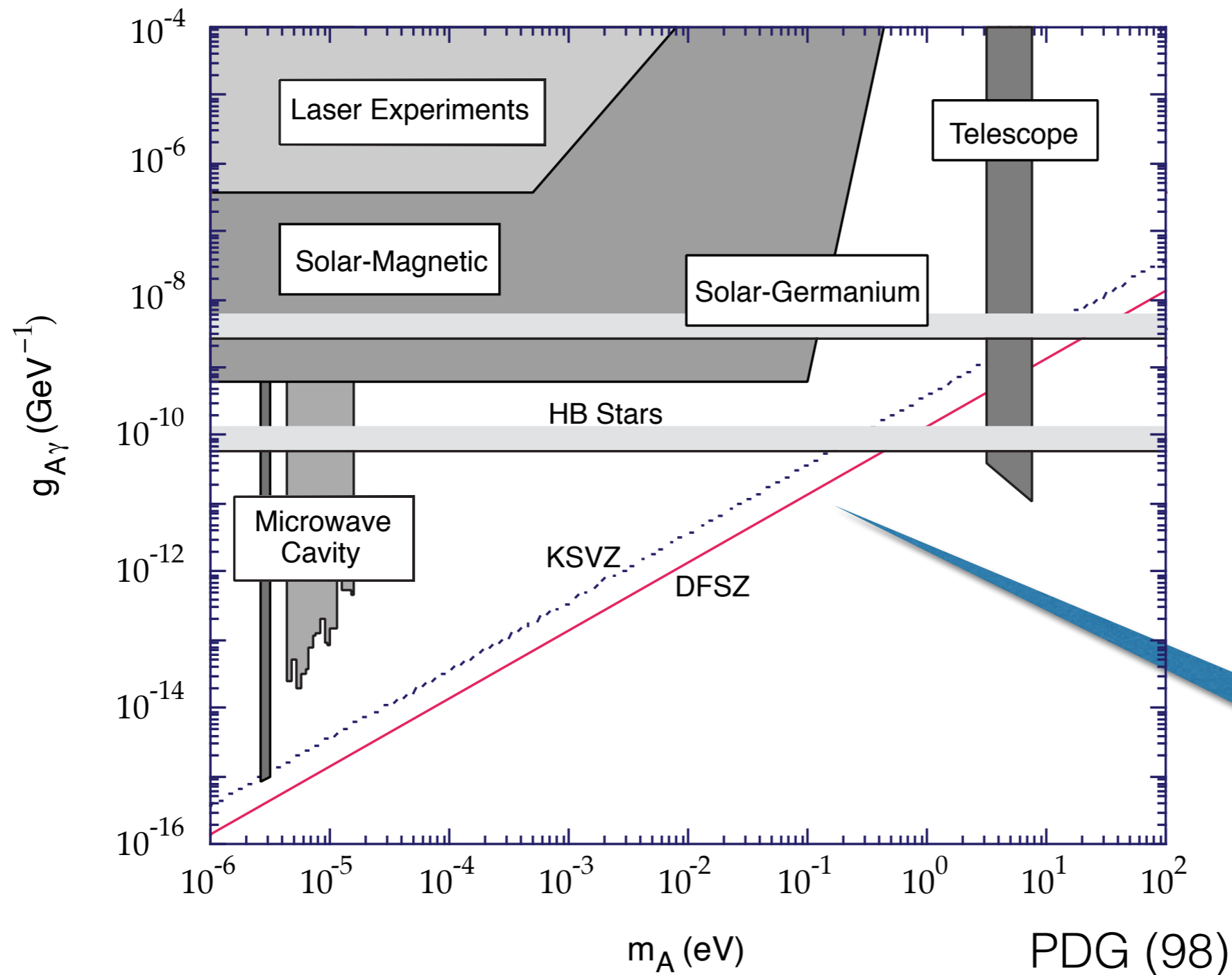
Focus shifts from Higgs to relaxion dynamics

- ◆ Can we even search for relaxion? Yes! How?
 - ◆ Different pheno', no partners. (stops/t', gauginos/KK Z's ...)
 - ◆ In most (but not all) cases, the relaxion is a **pseudo-Nambu-Goldstone-Boson** that (due to CP violation) **mixes w the Higgs**.
- Flacke, Frugiuele, Fuchs, Gupta & GP; Choi & Im (16)
- ◆ It implies that we can simplify its coupling to two parameters, mass & Higgs-mixing angle, with:

$$\sin \theta \lesssim \frac{m}{v} \lesssim 10\%, \quad m \lesssim \frac{v^2}{f} \lesssim 30 \text{ GeV} \times \frac{f}{\text{TeV}}. \quad (v = 174 \text{ GeV})$$

The relaxion parameter space, first analogy w axion

(back to *log* scale }-;)



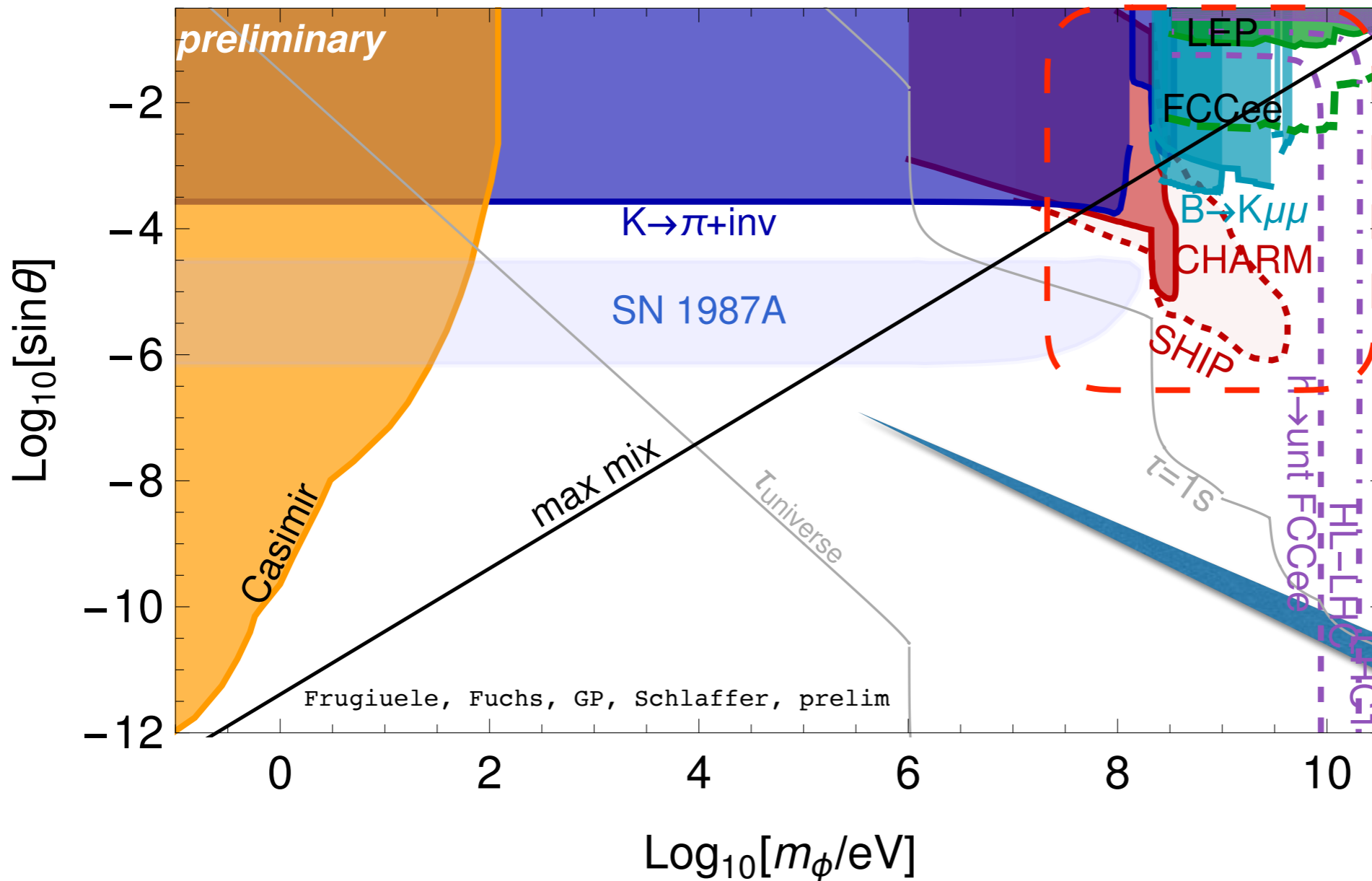
“Physical”
region:
on or below
diagonal

The relaxion parameter space, overview plot

Frugiuele, Fuchs, GP, Schlaffer, in preparation.

Bad news: *log* scale naturalness searches;

Good news: it seems that hep/colliders can probe physical region!



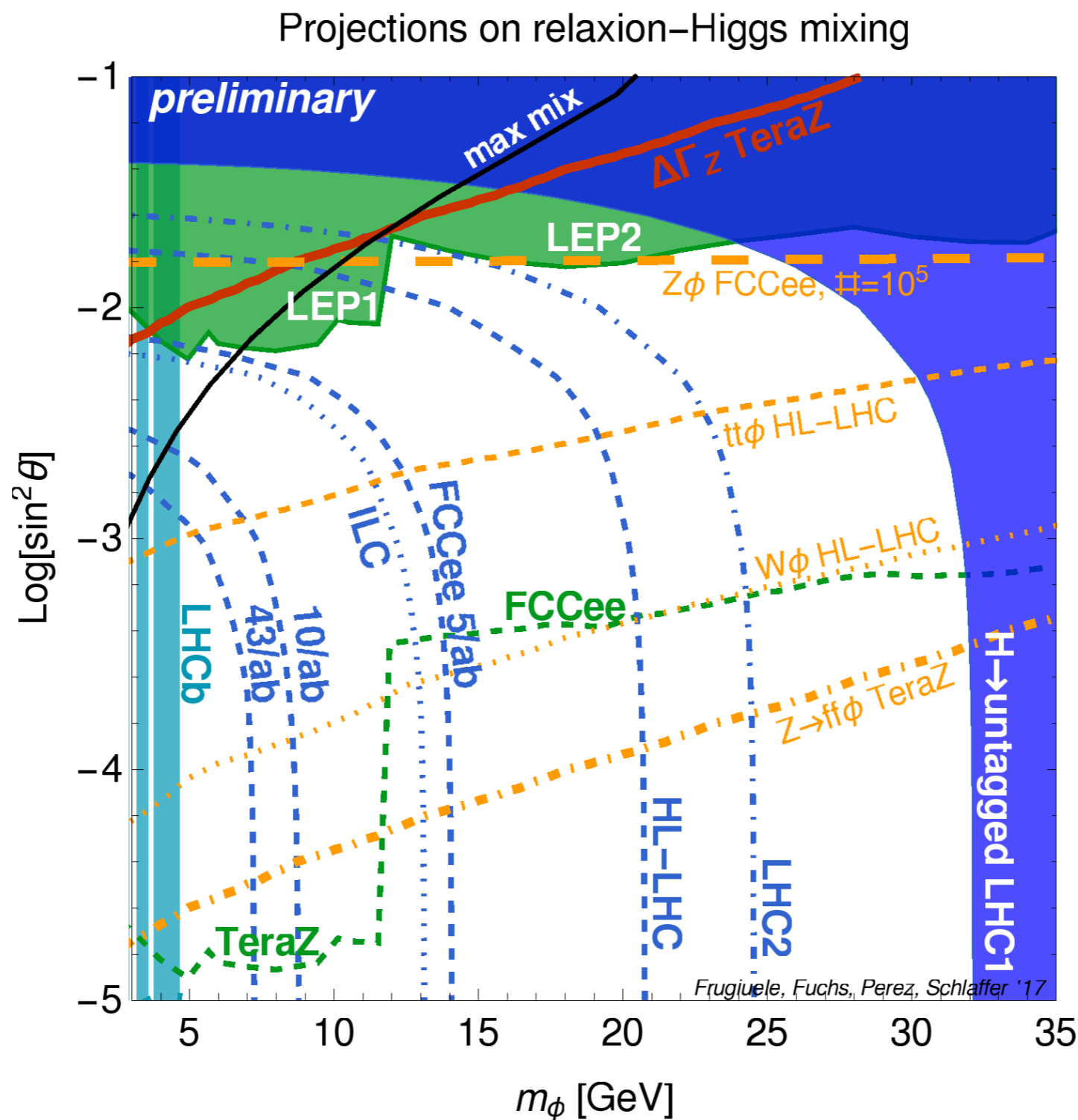
$$\left(\sin\theta \lesssim \frac{m}{v} \lesssim 10\%, \quad m \lesssim \frac{v^2}{f} \lesssim 30 \text{ GeV} \times \frac{f}{\text{TeV}} \right)$$

“Physical”
region:
on or below
diagonal

Zooming in on the region accessible to colliders

Frugiuele, Fuchs, GP, Schlaffer, in preparation.

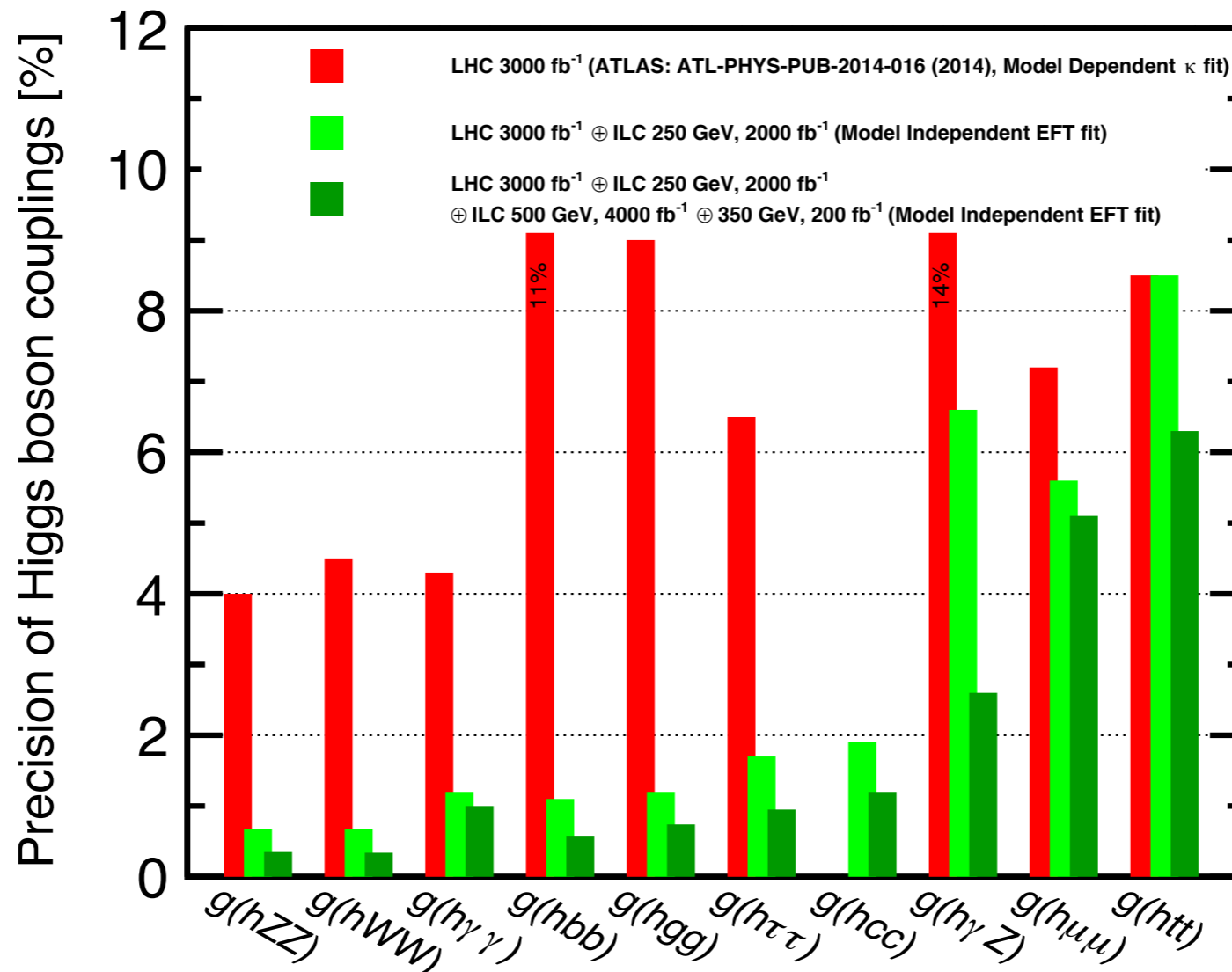
Lepton machines with large # of Z/H probe sizable region!
Same holds for HL-LHC & SHiP.



One slide on standard candles

Several SM coupling have not yet observed directly

◆ A conservative approach would be to attempt and measure these, in particular within the Higgs sector.



Fujii et al (2017)

◆ Lighter generation couplings required to establish the SM flavor picture are not shown.

Summary

- ◆ HEP **was** special \Rightarrow energy (E) frontier guaranteed discoveries.
- ◆ Reaction \Rightarrow discard theory - pragmatically move to signature based strategy - limited when planning to the future (what learnt?).
- ◆ Observationally-driven approach \Rightarrow log-crisis.
- ◆ Naturalness \Rightarrow possibly best motivation for E-frontier.
- ◆ Relaxion \Rightarrow undermine the energy frontier however motivates for Higgs precision frontier via relaxion-Higgs mixing.

Backups

Relaxion's basic structure

◆ QFT consistent constructions are of the form:

Choi, Kim & Yun (2014) Choi & Im; Kaplan & Rattazzi; Gupta, Komargodski, GP & Ubaldi (15)

(GKR: $g \sim M/f$)

$$V(\phi, H) = H^\dagger H [\Lambda^2 - \overbrace{M^2} \cos(\phi/f)] + r\Lambda^2 M^2 \cos(\phi/f) + H^\dagger H \tilde{M}^2 \cos(n\phi/f).$$

$$V'(\phi_*, v) = 0 \Rightarrow r\Lambda^4 \sin(\phi_*/f) \simeq v^2 n \tilde{M}^2 \sin(n\phi_*/f) \Rightarrow \phi_* \text{ is generic. } (M \sim \Lambda)$$

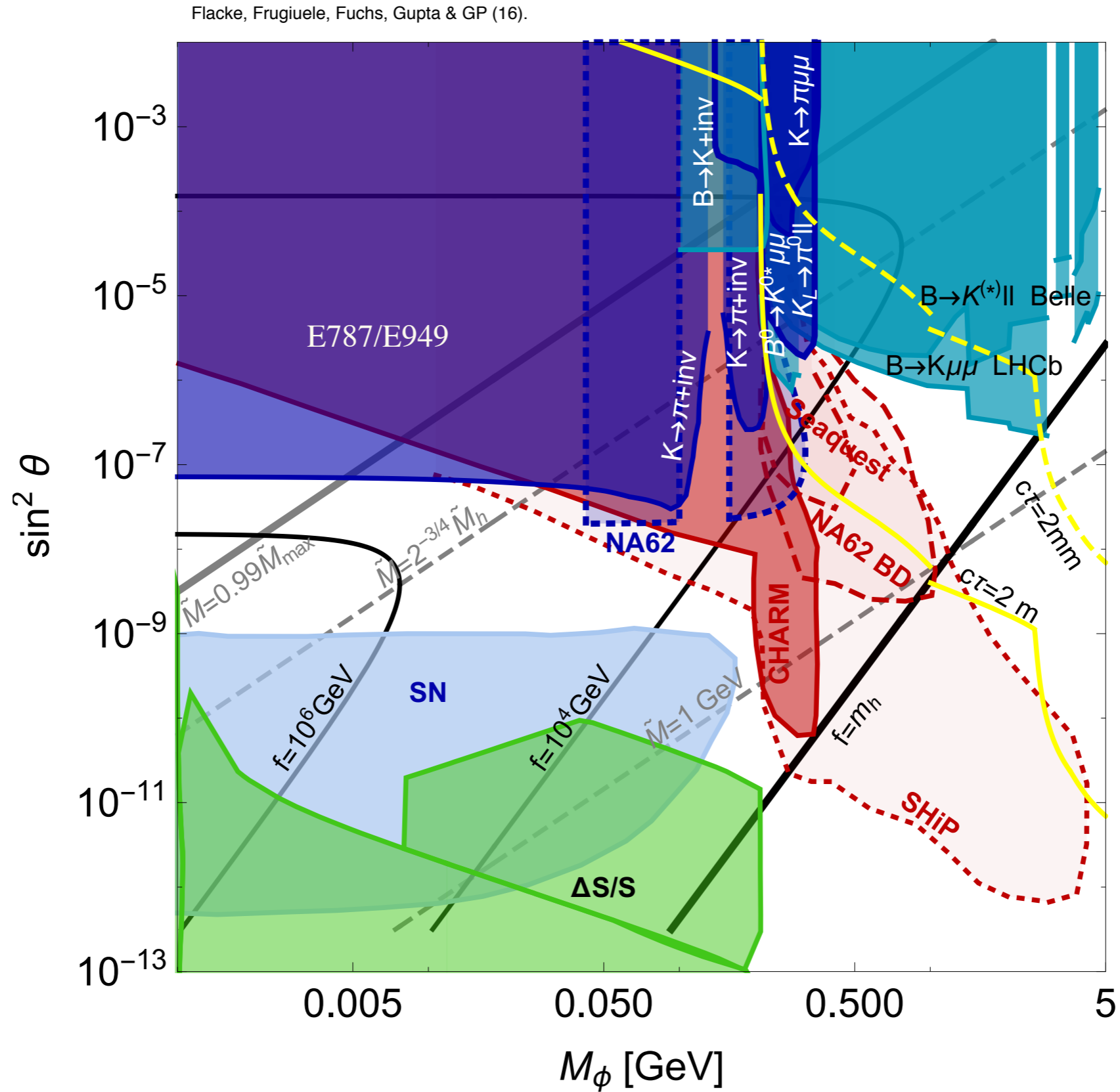
◆ It implies that generically:

(i) CP violation is spontaneously induced (problematic for axion-relaxion models);

(ii) Higgs-relaxion mixing is induced:

$$V_{\text{mix}} \sim \frac{nv\tilde{M}^2}{f} \sin(n\phi_*/f) \times H\phi_{\text{phys}} \simeq \frac{r\Lambda^4}{vf} \sin(\phi_*/f) \times H\phi_{\text{phys}}.$$

Relaxion beams, relaxion flavor



Back to Original Relaxion Proposal

Graham, Kaplan & Rajendran (15)

$$(\Lambda^2 - g^2 \phi^2) H^\dagger H \Rightarrow \phi_{\text{relaxed}} \sim \frac{\Lambda}{g}; \quad (\text{assume : } \Lambda \gg v)$$

$$V(\phi) = r^2 g^2 \Lambda^2 \phi^2 - v^n M_X^{4-n} \cos(\phi/f) \quad (\text{expect : } M_X < 4\pi v; 4 \geq n > 0)$$

$$V'(\phi) = 0 \Rightarrow \frac{\phi_{\text{relaxed}}}{f} \gtrsim \left(\frac{\Lambda}{4\pi v}\right)^4 \times r^2 \quad \left[g \lesssim \left(\frac{4\pi v}{\Lambda}\right)^4 \times \frac{\Lambda}{f r^2} \right]$$

Gupta, Komargodski, GP & Ubaldi (15)

$\Lambda \gg \text{TeV} \Rightarrow \langle \phi \rangle \gg f$ required to be physical.

The (compact) Relaxion Proposal

Gupta, Komargodski, GP & Ubaldi (15)

$\Lambda \gg \text{TeV} \Rightarrow \langle \phi \rangle \gg f$ required to be physical.

However, finite dim' EFT: pNGB \Rightarrow compact manifold.

Again: $\phi \rightarrow \phi + 2n\pi f$ ($n \in \mathbb{Z}$) lead to same physics.

This is a redundant description of the theory \Leftrightarrow discrete gauge sym' (no example w/ local operator that breaks it)

As long as relaxion potential is controlled by global internal sym' EFT locality seems to imply compactness of pNGB manifold:

$$\langle \phi \rangle \lesssim f.$$

Brief: Comments on the Relaxion Proposal

Gupta, Komargodski, GP & Ubaldi (15)

Hence upon the identification:

axion $\leftrightarrow \phi$, $U(1) \leftrightarrow PQ$, and

$y_u H f_\pi^3$ or $y_u^2 H^\dagger H f_\pi^2 \leftrightarrow m_L m_N y y_c H^\dagger H$,

we expect a similar bound to hold:

$$\Lambda \lesssim 10 \text{ TeV} \times \frac{[(y v)^{1,2} f_\pi^{3,2}]^{\frac{1}{4}}}{4\pi v} \times \left(\frac{1}{4\pi r}\right)^{\frac{1}{2}} \times \left(\frac{n}{10}\right)^{\frac{1}{4}}$$

Note: axion realisations also suffer from inflated $n \Rightarrow$ irrelevant operators \Rightarrow tiny backreaction/fine tuning/monstrous beta function.

The hierarchy: relaxion-familon-Nelson-Barr model

Davidi, Gupta, GP, Redigolo & Shalit (17)

Solving the CP problem - relaxion-Nelson-Barr

- ◆ The problem: relaxion spontaneously lead to order one CP Breaking.
- ◆ Nelson-Barr models solves the strong CP problem based on:
 - (i) CP being spontaneously broken.
 - (ii) Structure such the $O(1)$ CP phase induce the $O(1)$ CKM phase
(unlike lore a potential advantage compare to axion models ...)

If relaxion can be integrated to Nelson-Barr's structure => perfect marriage:

- (i) Nelson-Barr reminder; (ii) Relaxion-Nelson-Barr.

real
complex, relaxation the only source of CP breaking.

$$\mathcal{L}_{\text{NB}} = (\psi, Q) \begin{pmatrix} (\mu)_{1 \times 1} & (B_N(\phi))_{1 \times 3} \\ (0)_{3 \times 1} & (HY^d)_{3 \times 3} \end{pmatrix} \begin{pmatrix} \psi^c \\ d_i^c \end{pmatrix} + Y_u Q u^c H + h.c.$$

real

$$\text{Arg}(\det M_q^{7 \times 7}) = \text{Arg}(\det M_u) \text{Arg}(\mu \cdot \det(vY^d)) = 0.$$

No strong CP phase.

CKM: Integrating out heavy fermions, assuming $\mu^2 + B_n B_n^* \gg v^2 Y_{ik}^d Y_{jk}^d$,

$$\left[M_d^{eff} M_d^{eff \dagger} \right]_{ij} \sim v^2 Y_{ik}^d Y_{jk}^d - \frac{v^2 Y_{ik}^d B_k B_\ell^* Y_{j\ell}^d}{\mu^2 + B_n B_n^*}$$

O(1) CKM phase.

Relaxion-Nelson-Barr (NB)

- ◆ Challenge - how to transmit relaxion's complex VEV to quark matrix:

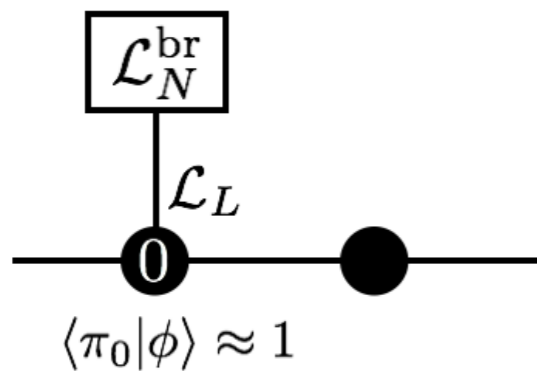
Original model - $B_i \leftrightarrow (g\phi + \tilde{g}\phi^*) \psi d^c$.

Relaxion-NB - $B_i \leftrightarrow (g\phi_N + \tilde{g}\phi_N^*) \psi d^c$.

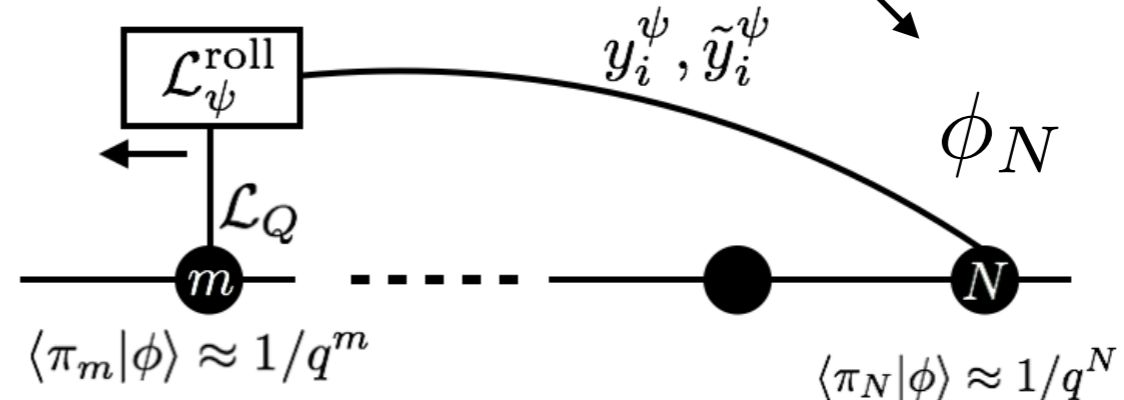
$$\phi_N \sim \rho_N \exp(i\phi/3^N f)$$

Inducing a phase + rolling potential!

Leptons &
Majorana neutrinos



Quark &
Nelson-Barr



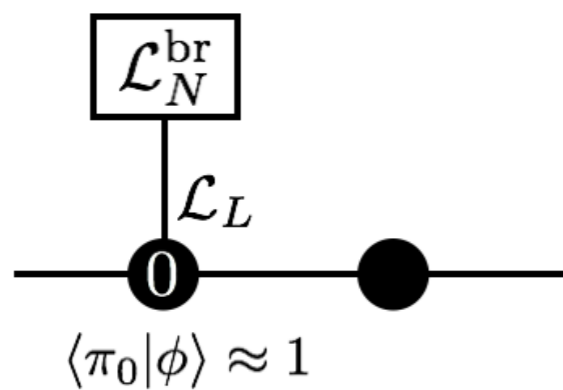
Relaxion-familon

- ◆ U(1) preserving int' on site 1st & m 'th \Rightarrow quark+lepton hierarchies.
- ◆ Traceless quark charges to avoid generating theta term.
- ◆ Explicit breaking on lepton sector a la Gupta et al yield backreaction.

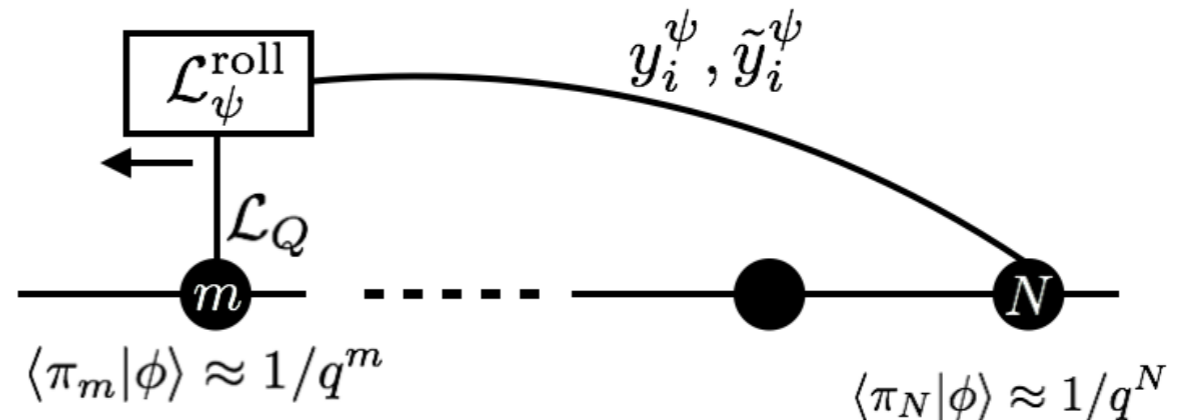


Hierarchion = Relaxion-familon-Nelson-Barr model of hierarchies

Leptons &
Majorana neutrinos



Quark &
Nelson-Barr



The main idea

Graham, Kaplan & Rajendran (15)

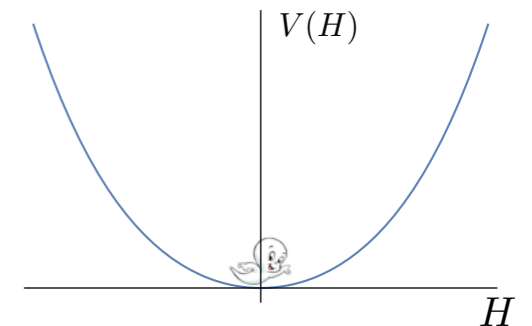
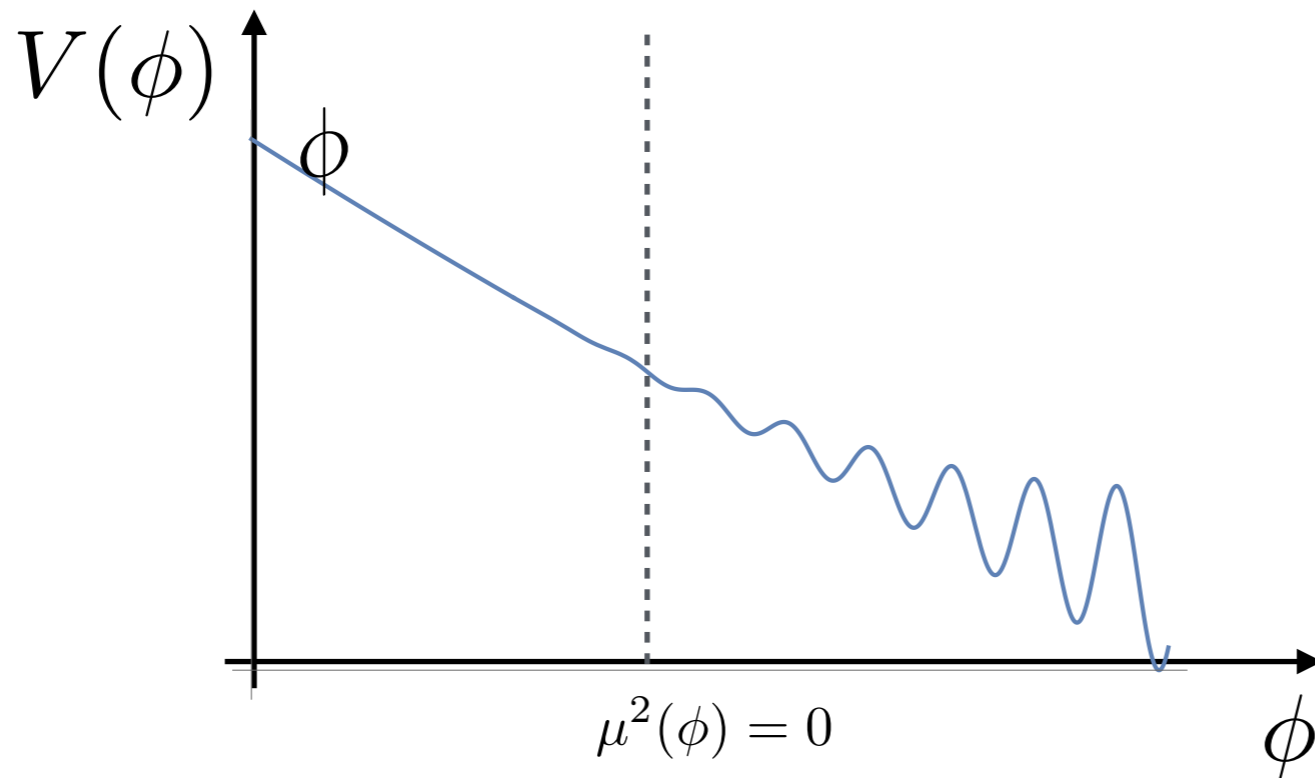
Relaxion's physics

Graham, Kaplan & Rajendran (15)

◆ A dynamical solution/amelioration of the Higgs fine-tuning problem:

(i) Add a scalar (relaxion) Higgs dependent mass: $(\Lambda^2 - \overbrace{g^2 \phi^2}^{\mu^2(\phi)}) H^\dagger H$.

(ii) ϕ rolls till μ^2 changes sign $\Rightarrow \langle H \rangle \neq 0 \Rightarrow$ stops rolling.



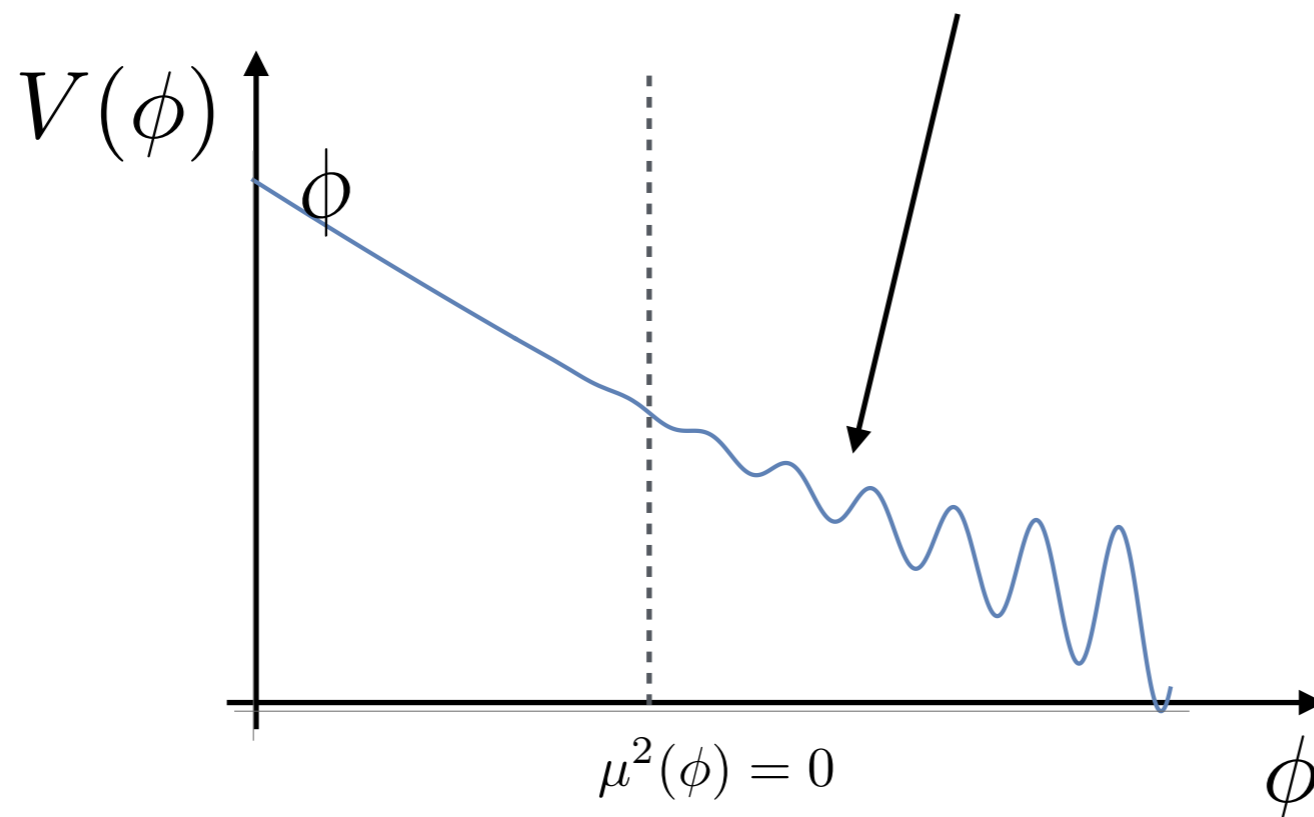
Relaxion mechanism

Graham, Kaplan & Rajendran (15)

◆ A dynamical solution/amelioration of the Higgs fine-tuning problem:

(i) Add a scalar (relaxion) Higgs dependent mass: $(\Lambda^2 - \overbrace{g^2 \phi^2}^{\mu^2(\phi)}) H^\dagger H$.

(ii) ϕ rolls till μ^2 flips sign $\Rightarrow \langle H \rangle \neq 0 \Rightarrow$ backreaction stops ϕ .



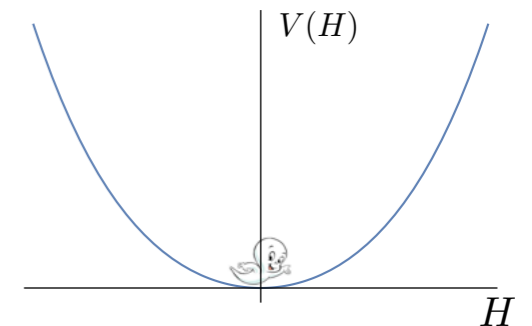
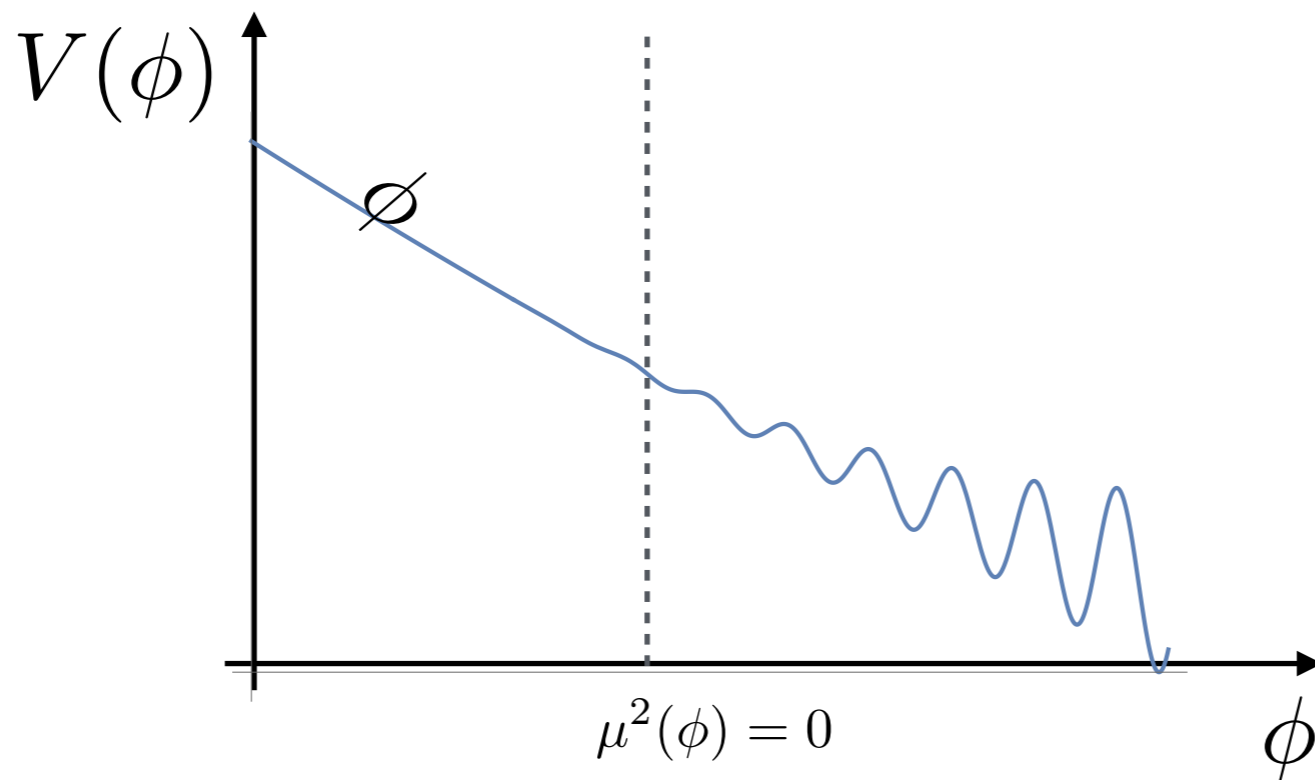
Relaxion's physics

Graham, Kaplan & Rajendran (15)

◆ A dynamical solution/amelioration of the Higgs fine-tuning problem:

(i) Add a scalar (relaxion) Higgs dependent mass: $(\Lambda^2 - \overbrace{g^2 \phi^2}^{\mu^2(\phi)}) H^\dagger H$.

(ii) ϕ rolls till μ^2 changes sign $\Rightarrow \langle H \rangle \neq 0 \Rightarrow$ stops rolling.



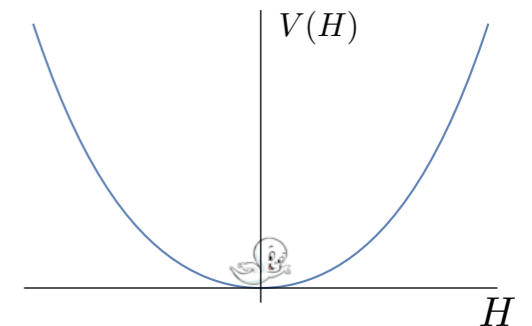
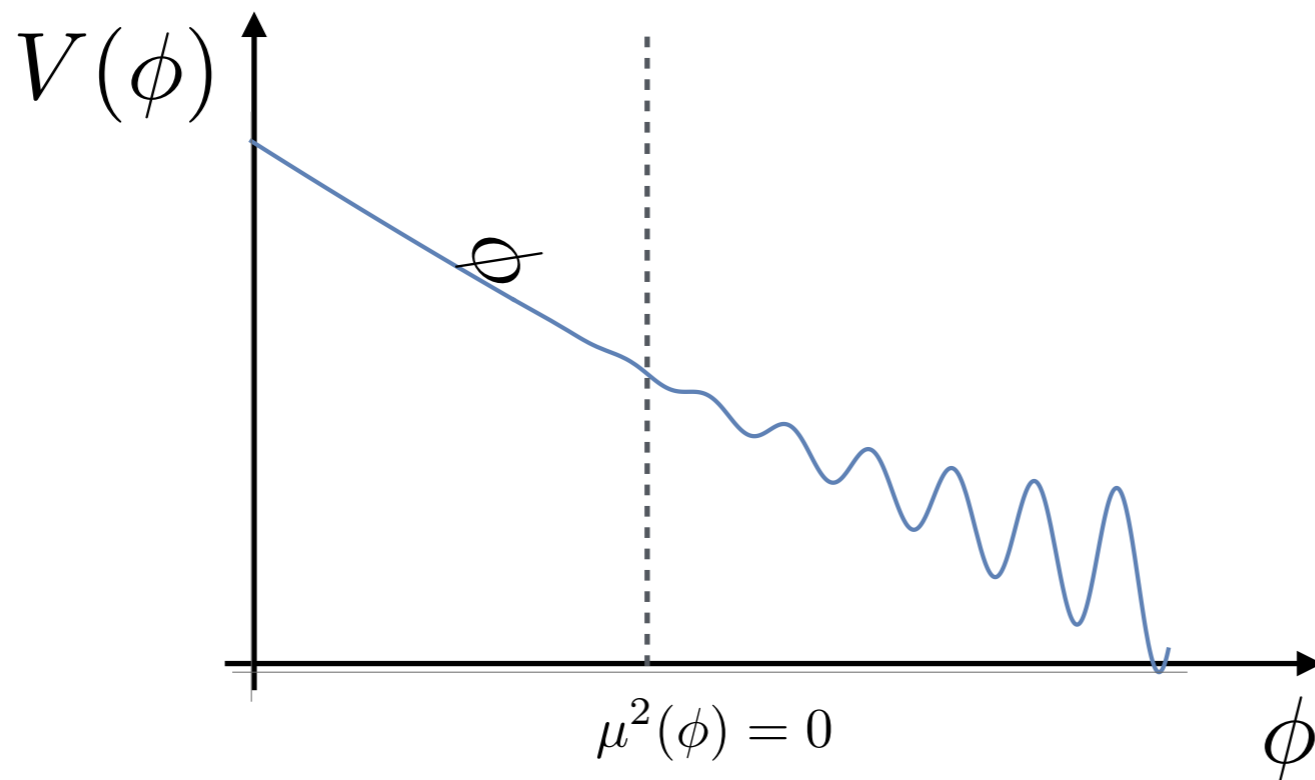
Relaxion's physics

Graham, Kaplan & Rajendran (15)

◆ A dynamical solution/amelioration of the Higgs fine-tuning problem:

(i) Add a scalar (relaxion) Higgs dependent mass: $(\Lambda^2 - \overbrace{g^2 \phi^2}^{\mu^2(\phi)}) H^\dagger H$.

(ii) ϕ rolls till μ^2 changes sign $\Rightarrow \langle H \rangle \neq 0 \Rightarrow$ stops rolling.



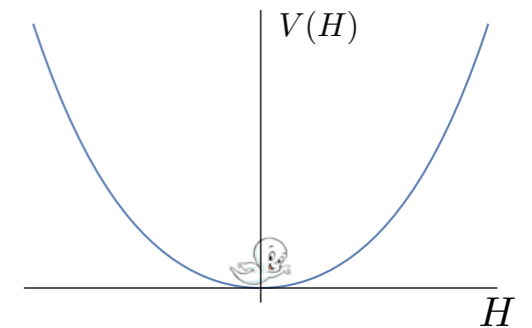
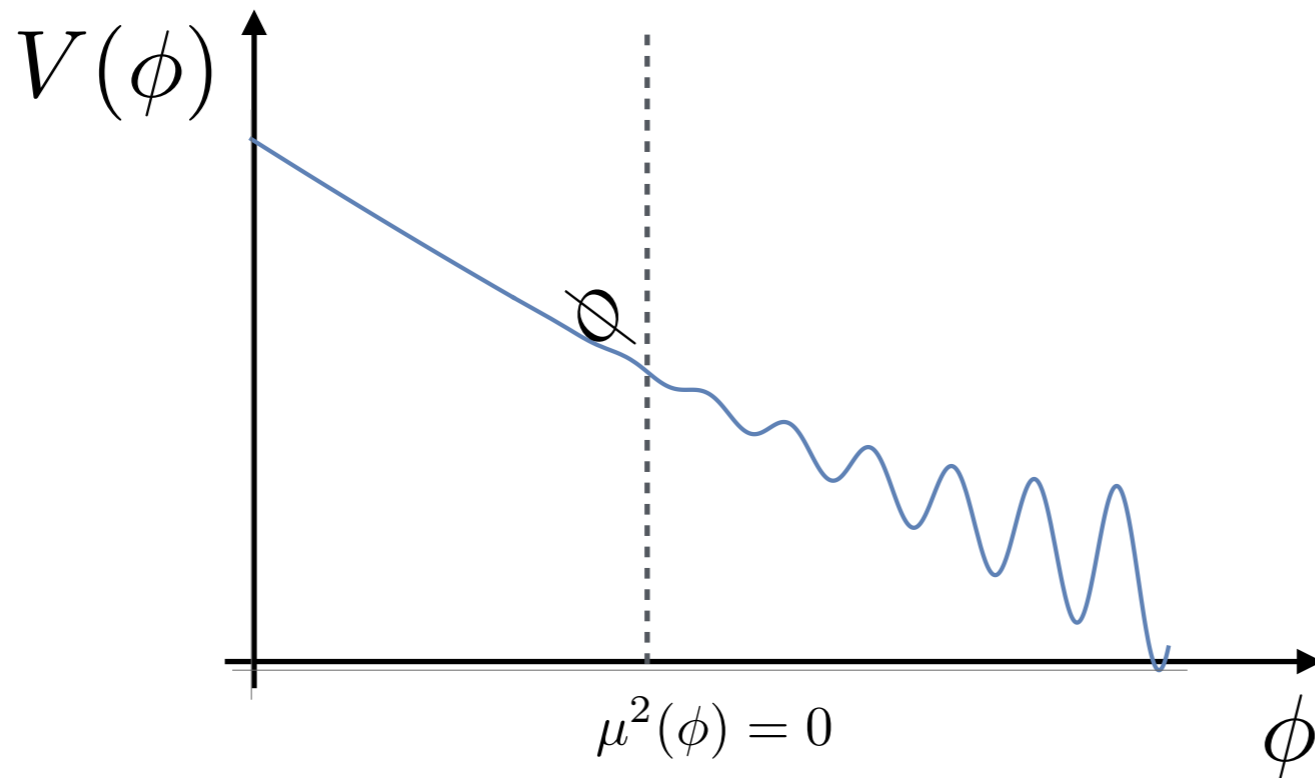
Relaxion's physics

Graham, Kaplan & Rajendran (15)

◆ A dynamical solution/amelioration of the Higgs fine-tuning problem:

(i) Add a scalar (relaxion) Higgs dependent mass: $(\Lambda^2 - \overbrace{g^2 \phi^2}^{\mu^2(\phi)}) H^\dagger H$.

(ii) ϕ rolls till μ^2 changes sign $\Rightarrow \langle H \rangle \neq 0 \Rightarrow$ stops rolling.



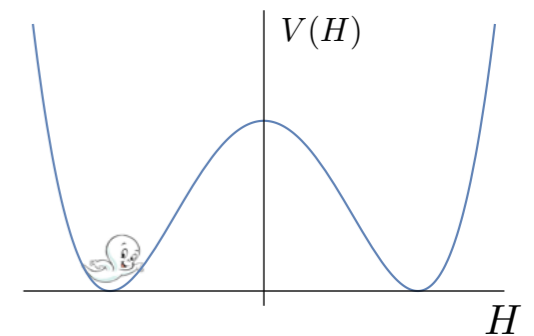
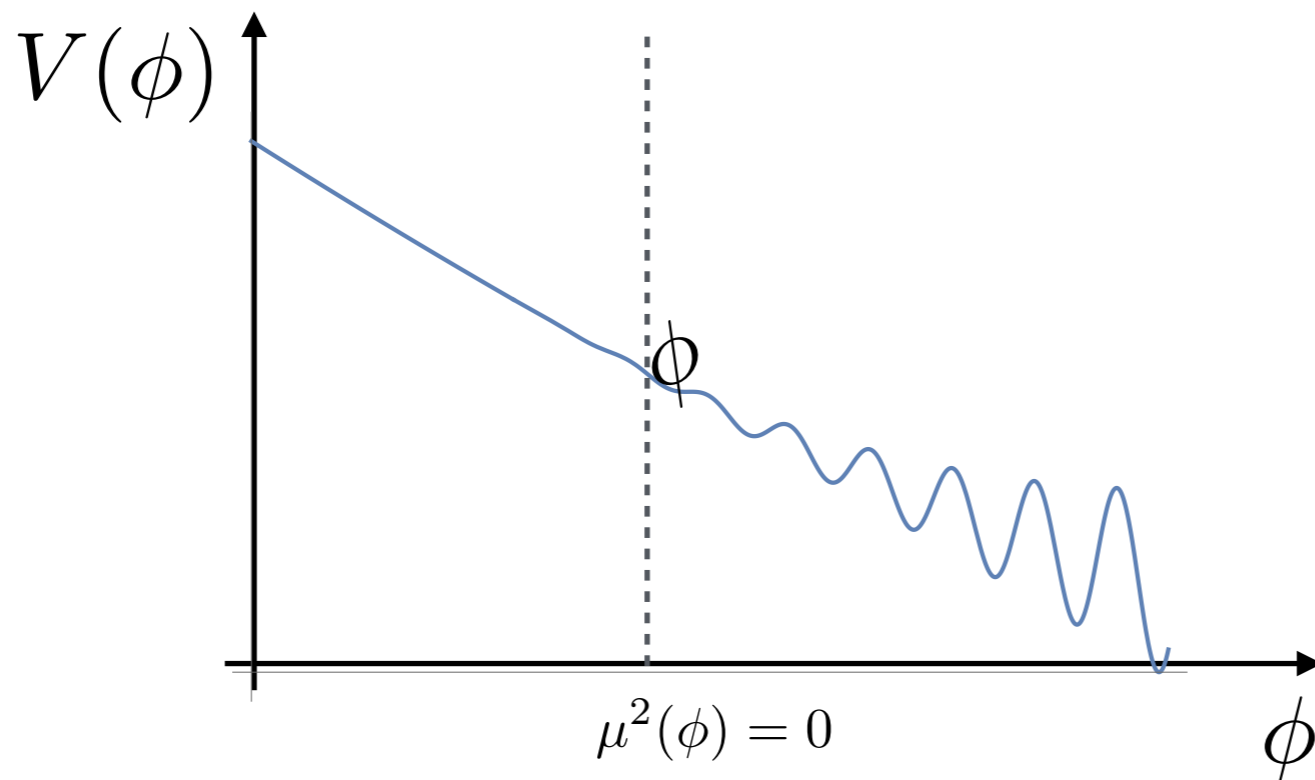
Relaxion's physics

Graham, Kaplan & Rajendran (15)

◆ A dynamical solution/amelioration of the Higgs fine-tuning problem:

(i) Add a scalar (relaxion) Higgs dependent mass: $(\Lambda^2 - \overbrace{g^2 \phi^2}^{\mu^2(\phi)}) H^\dagger H$.

(ii) ϕ rolls till μ^2 changes sign $\Rightarrow \langle H \rangle \neq 0 \Rightarrow$ stops rolling.



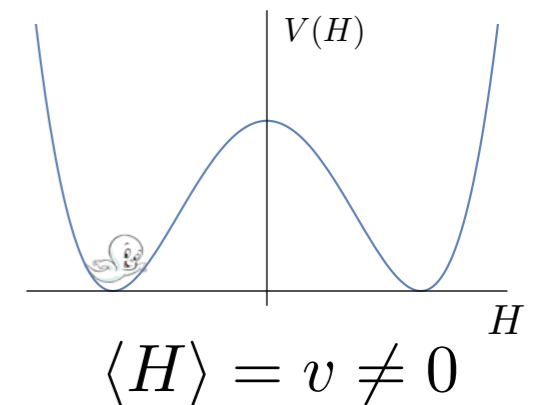
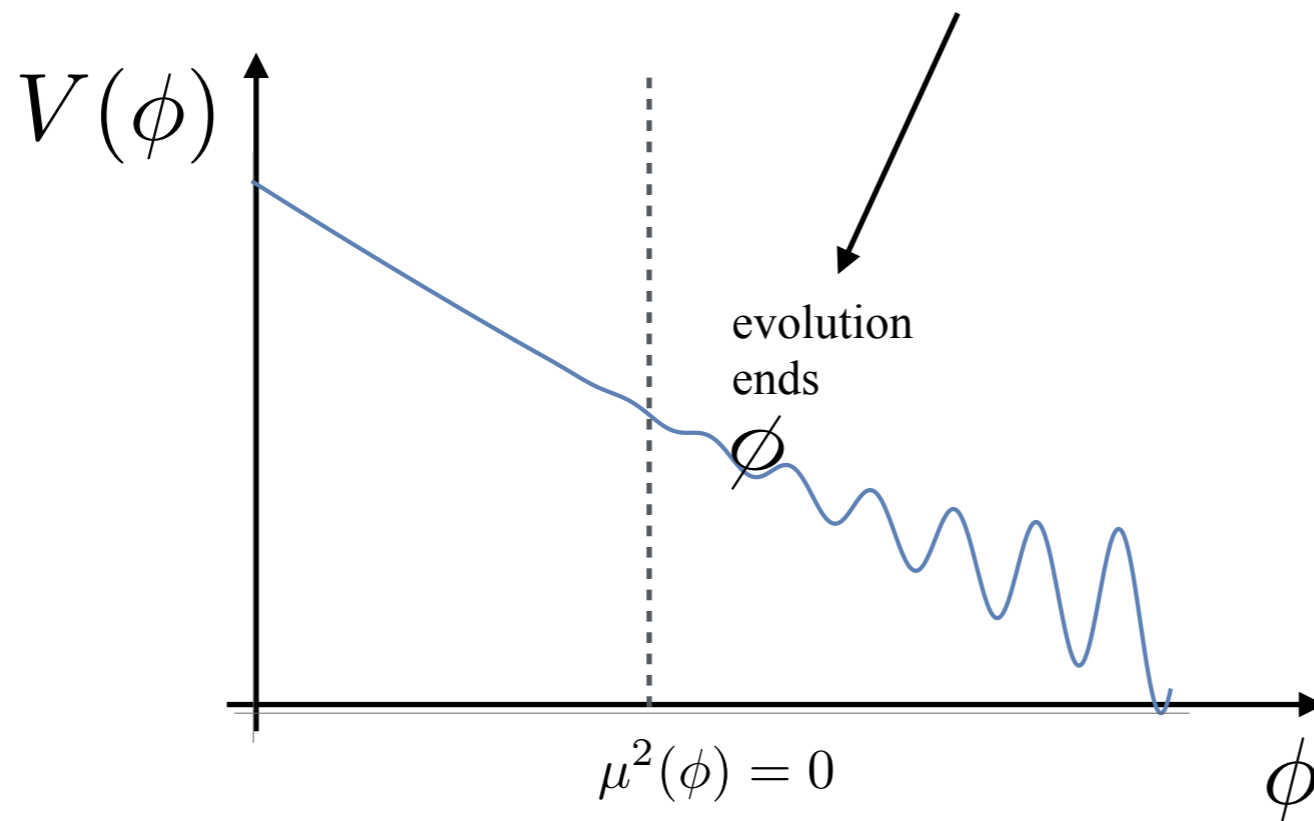
Relaxion's physics

Graham, Kaplan & Rajendran (15)

◆ A dynamical solution/amelioration of the Higgs fine-tuning problem:

(i) Add a scalar (relaxion) Higgs dependent mass: $(\Lambda^2 - \overbrace{g^2 \phi^2}^{\mu^2(\phi)}) H^\dagger H$.

(ii) ϕ rolls till μ^2 changes sign $\Rightarrow \langle H \rangle \neq 0 \Rightarrow$ stops rolling.



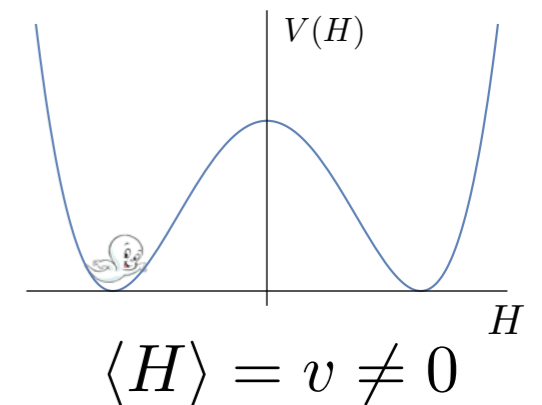
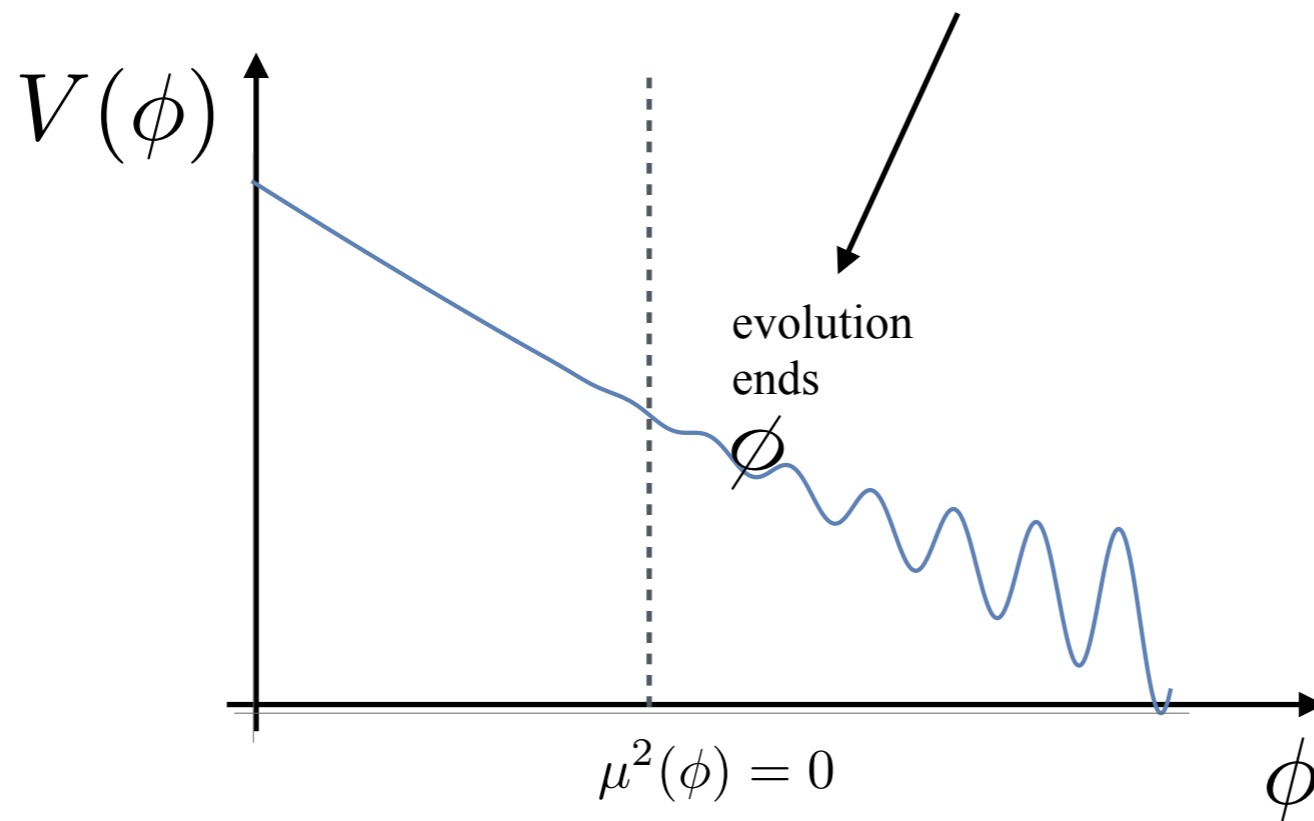
Relaxion's physics

Graham, Kaplan & Rajendran (15)

◆ A dynamical solution/amelioration of the Higgs fine-tuning problem:

(i) Add a scalar (relaxion) Higgs dependent mass: $(\Lambda^2 - \overbrace{g^2 \phi^2}^{\mu^2(\phi)}) H^\dagger H$.

(ii) ϕ rolls till μ^2 changes sign $\Rightarrow \langle H \rangle \neq 0 \Rightarrow$ stops rolling.



Roughly $\Lambda/v \lesssim n^{1/4} \sim (f_{\text{UV}}/f_{\text{IR}})^{1/4} \sim 3^{N_{\text{clock}}/4}$.

For $v \ll \Lambda$ progress achieved.