

EFFECTIVE LAGRANGIAN APPROACH TO TOP DECAY VIA FLAVOR CHANGING NEUTRAL CURRENT

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OBJECTIVES

Possible non-standard tqZ ($q = u, c$) couplings, which induce top decays via Flavor Changing Neutral Current (FCNC), are studied based on following strategies;

- Model-independent analysis is performed using effective Lagrangian.
- All the couplings are handled not only as complex numbers but also as independent parameters.
- Constraints on the couplings are given by comparing current experimental data with the theoretical values derived by varying them at the same time.
- Possible correlation among the parameters is studied.

EFFECTIVE LAGRANGIAN

Assuming that there exists some new physics characterized by an energy scale Λ (e.g., the mass of a typical new particle) and all the non-standard particles are not lighter than the Λ , the standard-model Lagrangian of tqZ interactions describing phenomena around the electroweak scale is extended as

$$\mathcal{L}_{tqZ} = -\frac{g}{2 \cos \theta_W} \left[\bar{\psi}_q(x) \gamma^\mu (f_1^L P_L + f_1^R P_R) \psi_t(x) Z_\mu(x) + \bar{\psi}_q(x) \frac{\sigma^{\mu\nu}}{M_Z} (f_2^L P_L + f_2^R P_R) \psi_t(x) \partial_\mu Z_\nu(x) \right],$$

where g and θ_W is the $SU(2)$ coupling constant and the weak mixing angle, $P_{L/R} \equiv (1 \mp \gamma_5)/2$. $f_{1,2}^{L,R}$ are non-standard couplings including Λ and vacuum expectation value, and treated as complex numbers independent of each other from the viewpoint of model-independent analysis.

EXPERIMENTAL DATA

The following experimental information at 95 % confidence level is used as our input data [ATLAS'17];

- the total decay width of the top quark, Γ^t [GeV].

$$4.8 \times 10^{-2} \leq \Gamma^t \leq 3.5$$

- The upper limits of the branching fractions for $t \rightarrow qZ$ decays

Current	Future expectation (@HL-LHC)
$\text{Br}(t \rightarrow uZ) < 1.7 \times 10^{-4}$	$\text{Br}(t \rightarrow uZ) < 8.5 \times 10^{-5}$
$\text{Br}(t \rightarrow cZ) < 2.3 \times 10^{-4}$	$\text{Br}(t \rightarrow cZ) < 1.2 \times 10^{-4}$

Then, multiplying the minimum (maximum) value of Γ^t by $\text{Br}(t \rightarrow uZ/cZ)$, the partial decay widths for each process, Γ_{tqZ} [GeV], which are input data in our analysis are obtained as

$$0 \leq \Gamma_{tuZ} < 8.1 \times 10^{-6} \quad (5.9 \times 10^{-4}),$$

$$0 \leq \Gamma_{tcZ} < 1.1 \times 10^{-5} \quad (8.0 \times 10^{-4}).$$

DISCUSSION

The maximum and minimum values of flavor-changing neutral tqZ couplings allowed by the present experimental data of the total decay widths and Branching rate were derived by varying all the couplings, $\text{Re}/\text{Im}(f_{1/2}^{L/R})$, independently at the same time.

- The allowed region derived by treating all the coupling constants as complex-number parameters (8 parameter analysis) could be about 1.8 times larger than that region derived by only one coupling being treated as a parameter (1 parameter analysis) because cancellations could happen among the contributions originated from those couplings.

e.g. tuZ couplings case;

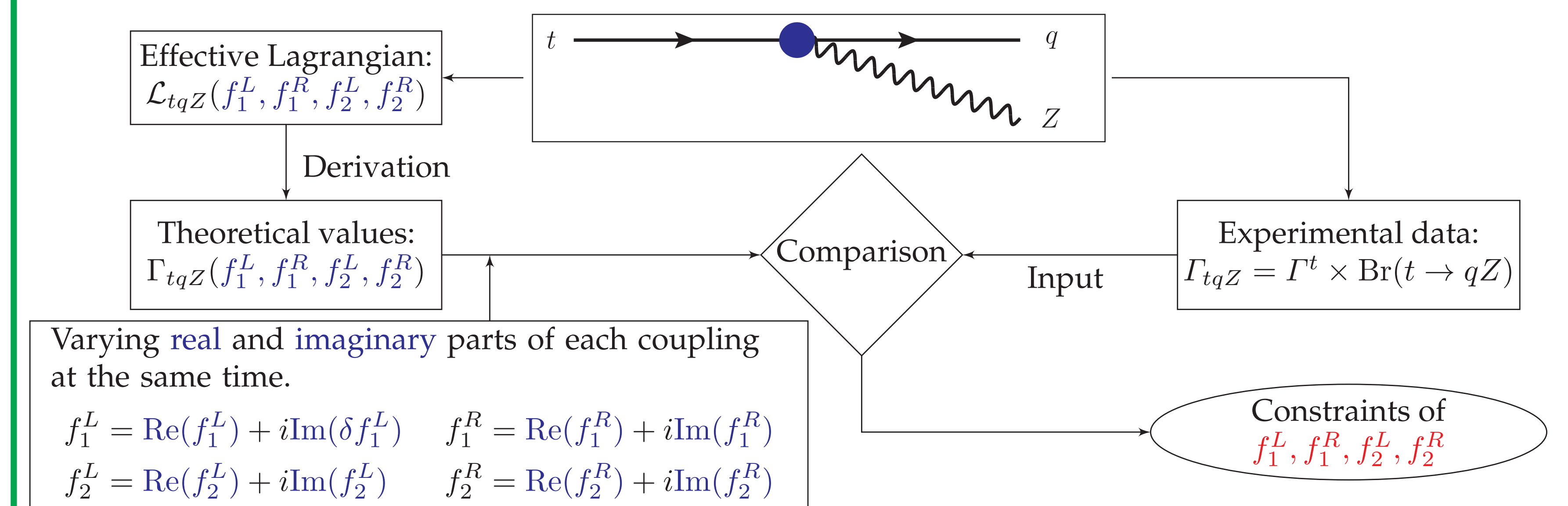
1 parameter analysis	8 parameter analysis
$-3.0 \times 10^{-3} < \text{Re}(f_1^L) < 3.0 \times 10^{-3}$	$-5.5 \times 10^{-3} < \text{Re}(f_1^L) < 5.5 \times 10^{-3}$
$-2.5 \times 10^{-3} < \text{Re}(f_2^L) < 2.5 \times 10^{-3}$	$-4.6 \times 10^{-3} < \text{Re}(f_2^L) < 4.6 \times 10^{-3}$

- There are strong correlations between $\pm \text{Re}/\text{Im}(f_1^{L/R})$ and $\mp \text{Re}/\text{Im}(f_2^{R/L})$;

e.g. Following table is allowed minimum and maximum values of the tuZ couplings for $\Gamma_{tuZ} = 8.1 \times 10^{-6}$ in the case that $\text{Re}(f_1^L)$ is fixed to 5.5×10^{-3} which is allowed maximum value. It is found that a strong correlation exists between $\text{Re}(f_1^L)$ and $\text{Re}(f_2^R)$.

	f_1^L		f_1^R		f_2^L		f_2^R	
	$\text{Re}(f_1^L)$	$\text{Im}(f_1^L)$	$\text{Re}(f_1^R)$	$\text{Im}(f_1^R)$	$\text{Re}(f_2^L)$	$\text{Im}(f_2^L)$	$\text{Re}(f_2^R)$	$\text{Im}(f_2^R)$
Min.	5.5×10^{-3}	-1.0×10^{-3}	-1.0×10^{-3}	-1.0×10^{-3}	-8.0×10^{-4}	-8.0×10^{-4}	-4.2×10^{-3}	-8.0×10^{-3}
Max.	5.5×10^{-3} (Fixed)	1.0×10^{-3}	1.0×10^{-3}	1.0×10^{-3}	8.0×10^{-4}	8.0×10^{-4}	-3.4×10^{-3}	8.0×10^{-3}

SCHEMATIC VIEW OF OUR APPROACH



Varying real and imaginary parts of each coupling at the same time.

$$f_1^L = \text{Re}(f_1^L) + i\text{Im}(\delta f_1^L) \quad f_1^R = \text{Re}(f_1^R) + i\text{Im}(f_1^R)$$

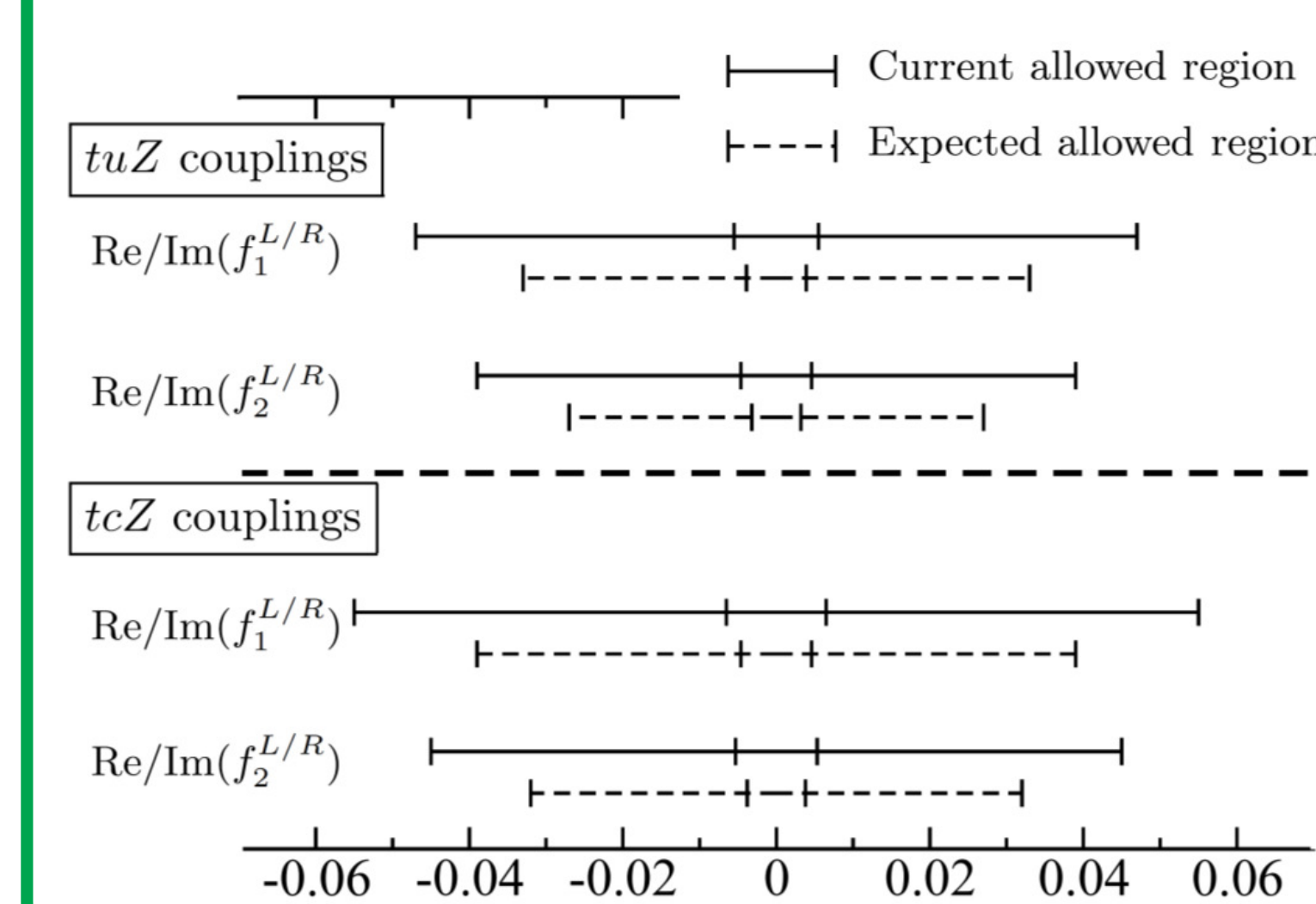
$$f_2^L = \text{Re}(f_2^L) + i\text{Im}(f_2^L) \quad f_2^R = \text{Re}(f_2^R) + i\text{Im}(f_2^R)$$

ALLOWED REGIONS OF tqZ

- Current constraints on the tuZ -coupling parameters: Those over (under) the dashed lines in the rows denoted as Min. and Max. are the minimum and maximum of the allowed ranges coming from $\Gamma_{tuZ} = 8.1 \times 10^{-6}$ (5.9×10^{-4}).

	f_1^L		f_1^R		f_2^L		f_2^R	
	$\text{Re}(f_1^L)$	$\text{Im}(f_1^L)$	$\text{Re}(f_1^R)$	$\text{Im}(f_1^R)$	$\text{Re}(f_2^L)$	$\text{Im}(f_2^L)$	$\text{Re}(f_2^R)$	$\text{Im}(f_2^R)$
Min.	-5.5×10^{-3}	-5.5×10^{-3}	-5.5×10^{-3}	-5.5×10^{-3}	-4.6×10^{-3}	-4.6×10^{-3}	-4.6×10^{-3}	-4.6×10^{-3}
	-4.7×10^{-2}	-4.7×10^{-2}	-4.7×10^{-2}	-4.7×10^{-2}	-3.9×10^{-2}	-3.9×10^{-2}	-3.9×10^{-2}	-3.9×10^{-2}
Max.	5.5×10^{-3}	5.5×10^{-3}	5.5×10^{-3}	5.5×10^{-3}	4.6×10^{-3}	4.6×10^{-3}	4.6×10^{-3}	4.6×10^{-3}
	4.7×10^{-2}	4.7×10^{-2}	4.7×10^{-2}	4.7×10^{-2}	3.9×10^{-2}	3.9×10^{-2}	3.9×10^{-2}	3.9×10^{-2}

- Summary of Current and expected constraints on the tuZ and tcZ couplings



- The allowed regions of tqZ couplings are within $|f_{1/2}^{L/R}| < \mathcal{O}(10^{-3}) \sim \mathcal{O}(10^{-2})$.
- The tuZ couplings are more strongly restricted than the tcZ couplings.
- Both the real and imaginary parts of $f_1^{L/R}$ and $f_2^{L/R}$ in each of the tuZ and tcZ couplings have the same minimum and maximum limits, respectively.
- The allowed regions are expected to be narrowed by about 30 % if the assumed branch fractions are realized at the HL-LHC with 3000fb^{-1} luminosity.

REFERENCES

This presentation is based on following references;

- Z. Hioki, K. Ohkuma and A. Uejima, "Studying flavor-changing neutral tqZ couplings: Current constraints and future prospects," arXiv:1809.01389 [hep-ph]. To appear in Modern Physics Letters A
- Z. Hioki, K. Ohkuma and A. Uejima, Work in progress

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